# Positron transport and acceleration in beam-driven plasma accelerators using a plasma column

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#### **Positron acceleration is a challenge**

In the blowout regime

- Plasma electrons return to axis on the same spot, forming a **narrow, high density electron cusp** —
- Positron focusing fields only within the small region of high electron density

Other concepts, using hollow core drive beams,<sup>1</sup> hollow core plasma channels,<sup>2</sup> or positron drive beams<sup>3</sup> are either **inefficient** or **unstable** 

<sup>1</sup>N. Jain *et al.*, Phys. Rev. Lett. 115, 195001 (2015)
<sup>2</sup>S. Gessner *et al.*, Nat. Commun. 7, 11785 (2016)
<sup>3</sup>S. Corde *et al.*, Nature (London) 524, 442 (2015)

Drive beam

2

-2

-2

 $k_p \zeta$ 

0

2

Ω

 $n_p/n_0 \quad n_b/n_0$ 

- 3

-2

2.0

-1.5

 $\cdot 1.0$ 

-0.5

 $E_z/E_0$ 

-0.5

-1.0

-1.5

-2.0

- 3.0

-2.4

-1.8

-1.2

-0.6

0.0

-0.75

-0.50

-0.25

-0.00

-0.25

 $(E_x - B_y)/E_0$ 

Plasma wave

3

2

 $k_p$ 

-1

-2

-3

 $k_p x$ 

 $^{-1}$ 

-2

-3

-6

-4

Focusing field

for positrons

-6

with regard to positron acceleration

#### in pre-ionized plasma columns

In the homogeneous, infinite plasma case:







T. J. Mehrling *et al.*, Plasma Phys. Controlled Fusion 56, 084012 (2014) T. J. Mehrling *et al.*, 2018 IEEE AAC Proceedings

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#### in pre-ionized plasma columns

In the homogeneous, infinite plasma case:





#### in pre-ionized plasma columns

Now, using a finite, pre-ionized plasma column:





6

4 -

#### in pre-ionized plasma columns

Now, using a finite, pre-ionized plasma column:





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modified transverse wakefield

## Elongated plasma electron trajectories induce positron acc. field

#### in pre-ionized plasma columns

- Modified transverse wakefield leads to elongated plasma electron trajectories
- Elongated plasma electron trajectories create long, high-density electron filament
- This leads to a region with accelerating and focusing fields for positrons



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#### Plasma radius can be optimized w.r.t. efficient positron acc.

#### in pre-ionized plasma columns

 The optimal plasma column radius for a bi-Gaussian beam can be calculated by

$$k_p R_p \approx 2\sqrt[3]{2I_b/I_A}$$



Here:  $\sigma_x = 0.3 k_p^{-1}$ ,  $\sigma_z = \sqrt{2} k_p^{-1}$ ,  $I_b/I_A = 1$ 

Simulation performed with cylindricallysymmetric, quasi-static PIC code INF&RNO<sup>1,2</sup>



<sup>1</sup>C. Benedetti *et al.*, AIP Conf. Proc. 1299, 250 (2010). <sup>2</sup>C. Benedetti *et al.*, AIP Conf. Proc. 1812, 050005 (2017).

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# Positron transport and acceleration in plasma columns

Important requirements for collider applications:

- High accelerating gradient
- High witness bunch charge
- Low emittance > emittance needs to be preserved during acceleration

Emittance (normalized to  $m_e c$ ):  $\epsilon_x = \sqrt{\langle x^2 \rangle \langle u_x^2 \rangle} - \langle x u_x \rangle^2$ for a Gaussian beam with no correlation:  $\epsilon_x = \sigma_x \sigma_{u_x}$ 

• Emittance is usually preserved in linear focusing fields

## Emittance preservation achievable with matched beams

• Here: focusing fields are almost like a **step-function**:

$$\frac{E_x - B_y}{E_0} = -\alpha \operatorname{sgn}(x)$$

 For quasi-matched beams emittance growth at saturation can be calculated numerically and even for a simple Gaussian with

$$\sigma_x^3 \simeq 1.72 \ \frac{\epsilon_x^2}{\alpha \gamma}$$

it yields  $\simeq 2\%$ .

 Non-linear matching condition: for a chosen emittance ε<sub>x</sub>, the optimal witness beam rms size σ<sub>x</sub> can be calculated for a given field strength α and beam energy γ



Simulations performed with INF&RNO Page 14

## **Demonstration of emittance-preserving positron acceleration**

**Comparison to PIC simulation** 

- Witness beam parameters:  $\epsilon_x = 0.1 k_p^{-1}$ ,  $\sigma_x = 0.025 k_p^{-1}$ ,  $\gamma = 2000$ With background plasma density  $n_0 = 5 \times 10^{17} cm^{-3}$ : Charge  $Q_w \approx 84 \ pC$ ,  $Q_D \approx 1.5 \ nC$
- Average acceleration gradient:  $30 \ GeV/m$



Emittance growth from simulation: Quasi-matched central slice:  $\simeq 3\%$ total (projected) bunch:  $\simeq 7\%$ 

The drive beam is kept rigid in this simulation

#### Energy spread can be controlled by beam loading



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#### Plasma column can be generated by beam-field-ionization

Space-charge fields of beam ionize a plasma column, wakefield ionizes further behind drive beam

Advantage: self-inherent alignment between drive beam and plasma column



Expanded region of high electron density, possibly allowing for positron acceleration

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 $k_p X_0 \quad n_p/n_0$ -1.75-1.50-1.25 $k_p x$  $k_p x$ -1.00 - 0.75 -2-2- 0.50 -0.25-4-40.00 0 -6-2-6-4 $k_p \zeta$ 

Expanded region of high electron density, possibly allowing for positron acceleration



Not in this case, positron accelerating and focusing field don't match

Solvable by parameter optimization!

 $(E_r - B_u)/E_0$ 

## **Coupled plasma column generation limits accelerating fields**



## Optimized (Gaussian) drive beam parameters yield

#### Problem:

Plasma column generation is **coupled** to drive beam parameters, strongly limiting the usable parameter space!

Additionally: scheme is vulnerable to drive beam variations (e.g. head erosion)

Therefore: Higher accelerating gradient and more stable scheme by using laser-pre-ionized plasma columns, as this **decouples** plasma generation and the drive beam.



- Plasma structure suitable for positron transport and acceleration has been proposed
- The wakefields in these structures were studied and optimized with respect to positron acceleration
- The superior concept of using laser-pre-ionized plasma columns was used to demonstrate qualitypreserving positron acceleration

For more details, see our publication: **Diederichs** *et al.*, **Phys. Rev. Accel. Beams 22, 081301 (2019)** or my master's thesis

