



Uniwersytet
Wrocławski



Numerical study of fabrication tolerances for dielectric laser acceleration (DLA) structures

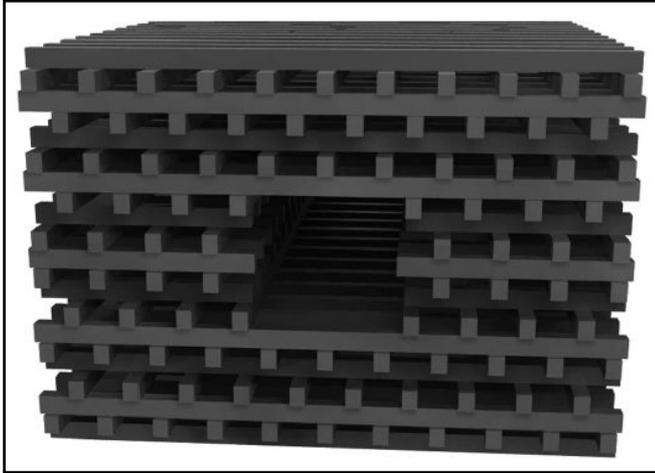
Andrzej Szczepkowicz

*Institute of Experimental Physics
University of Wrocław*

EAAC 2019, Elba
Working group 3



1. Literature on tolerances in DLA
2. Simple numerical models:
 - double column setup
 - Bragg mirror
3. Adjusting of a fabricated structure?
 - electro-mechanical: piezoelectric effect
 - electro-optical: Pockels effect
4. Conclusions



J. England *et al.*, Rev. Mod. Phys. 86, 1337 (2014);
B. Cowan *et al.*, IPAC2010 proceedings;

“Fabrication of 17-layer woodpile accelerator structures was successfully achieved using the layer-by-layer approach [...]. Fabrication tolerances were within 5–8% of the layer thickness, rod width, layer-to-layer alignment, and taper angle. Final alignment of the two half structures reveals a one-third period offset. Improved accuracy and automation in the alignment system is important for future fabrication runs.”



F. Mayet *et al.*, *First order sensitivity analysis in dual grating type dielectric laser acceleration structures*, Proceedings of IPAC2018

- Influence of laser beam parameters and input electron beam parameters on output beam parameters

partial variances is used. This method is based on a large number of Monte Carlo runs of a given model. This way also numerical models can be used. The resulting sensitivity measure is referred to as the *first order sensitivity index* of a given input parameter to the system. It is given by

$$S_i = \frac{V_{a_i}(E_{\sim a_i}(f(\mathbf{a})|a_i))}{V(f(\mathbf{a}))} \in [0, 1], \quad (4)$$

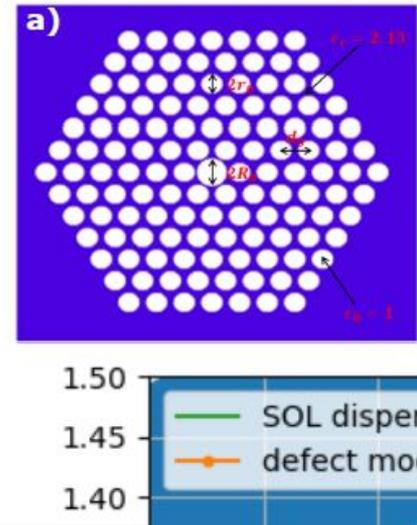
where – sticking to the literature – V is the variance, E the expectation value and $\sim a_i$ means “all but a_i ” (cf. [3]). The numerator can be read as “the variance of the expectation value of $f(\mathbf{a})$ for fixed (known) a_i ”. One additional advantage of this measure is the fact that the whole configuration space is explored instead of focusing on one fixed a_i . In

[A. Saltelli, *Global Sensitivity Analysis: The Primer*, Wiley, Chichester, (2008).]

“Finally it has to be noted that this kind of analysis does not have to be based on simulations. It can also be performed on experimental data. If the data acquisition of all relevant machine parameters is time synchronized, recorded data then corresponds to the Monte Carlo runs of a given model.”

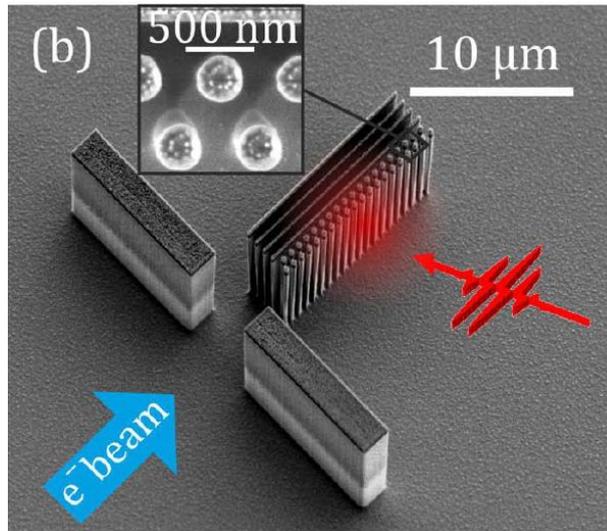
Abstract

Laser-driven hollow core photonic bandgap (PBG) fibers were proposed by Lin in 2001 as high-gradient accelerators. The central defect in the transversely periodic lattice supports an accelerating mode for synchronous acceleration in the ultra-relativistic regime. The optical frequencies in such dielectric laser accelerators motivate a sensitivity and tolerance study to overcome manufacturing imperfections. Finally we discuss the propagation characteristics of Lin-fibers and find that small-bandwidth (\sim ns) pulses would be needed for efficient acceleration over longer distances.



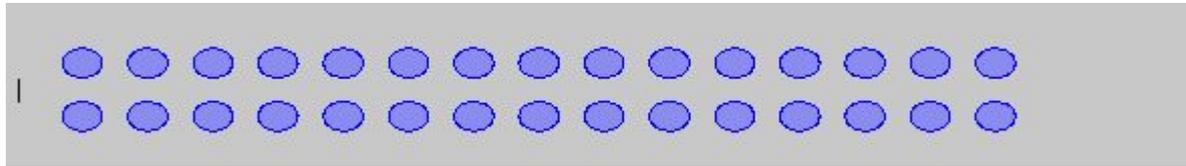
In this paper we have presented a tolerance study assuming realistic manufacturing imperfections and showing that the geometrical PBG fiber properties play a crucial role for getting and confining the optimal accelerating mode into the defect. Moreover, we find a tolerance range of 10% in which the mode properties in the fiber can be recovered by tuning the laser wavelength. Finally, we have pointed out that a

P. Yousefi *et al.*, EAAC2017 proceedings,
P. Yousefi *et al.*, Opt. Lett. 44, 1520 (2019).



“Our final structures have a standard deviation of 10 nm from the designed geometry” [$\sim 0.5\%$ of the laser wavelength] [local defects?]

“We fabricated two sets of structures [...]. The relative difference between geometrical dimensions of the fabricated structures is 0.8%.”



- Geometry and laser parameters from D. Black *et al.*, *Laser-driven electron lensing in silicon microstructures*, PRL 122, 104801 (2019)

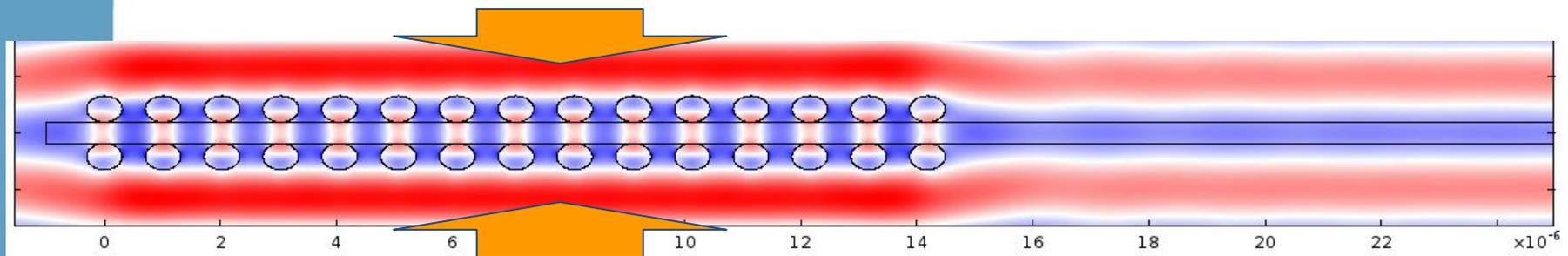
15 pairs of Si columns

electron velocity $\beta = 0.525$

laser wavelength = 1.95 μm , vacuum gap = 0.375 μm

laser field from two sides, amplitude $E_0 = 300 \text{ MV/m}$

- focal length = 18.4 μm



- Full PIC code
- Often used alternative: a code to calculate field, another code to track particles
- This work: all in Comsol (RF Module + Particle Tracing Module)
- Particle tracking approaches based on the transfer properties of a DLA unit cell, utilizing quasi-periodicity of DLA structures.
 - U. Niedermayer et al., IPAC2017 Proceedings;
U. Niedermayer et al., PRAB 20, 111302 (2017).
 - W. Kuroпка, et al, IPAC2017 Proceedings;
F. Mayet et al, IPAC2017 Proceedings
 - A. Szczepkowicz, PRAB 20, 081302 (2017);
A. Szczepkowicz, EAAC2017 Proceedings
 - Beware of boundary effects

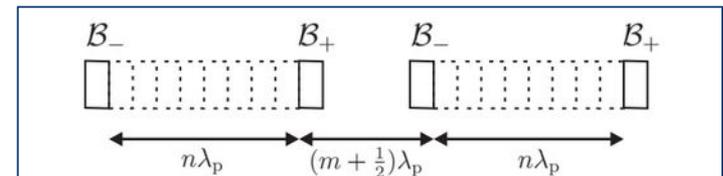
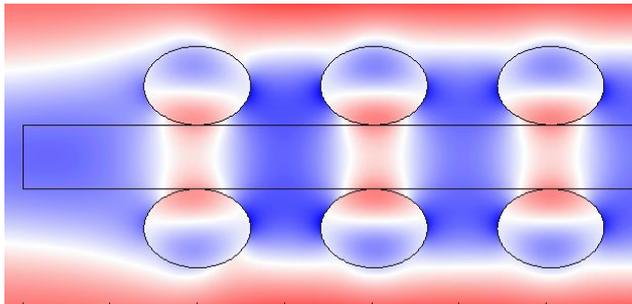
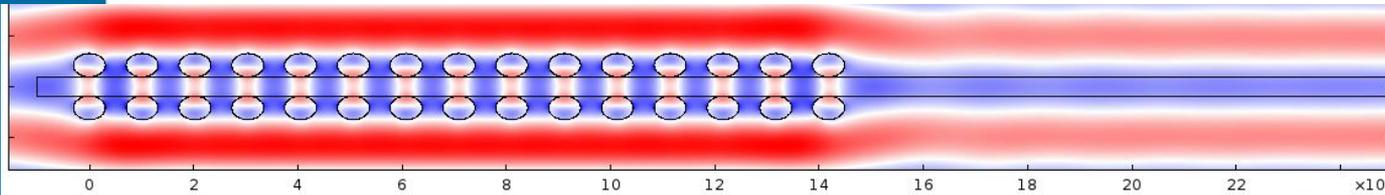


FIG. 8. Boundary field effects can be handled using boundary cells with corresponding \mathcal{B}_{\pm} transfer functions.

A. Szczepkowicz, PRAB 20, 081302 (2017)

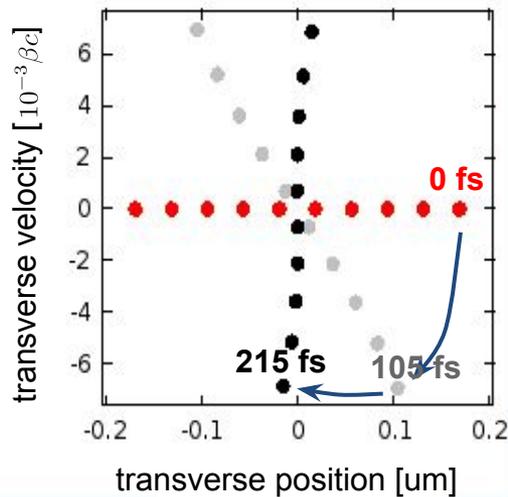
A numerical model: double column DLA segment

1. Sensitivity to column cross-section size



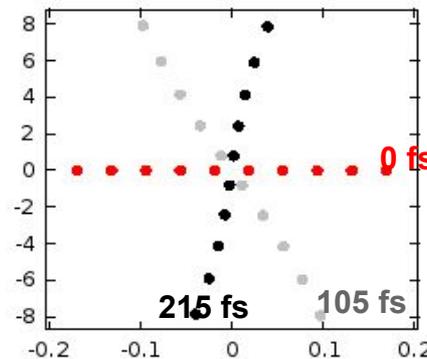
Trace space plots

Design geometry



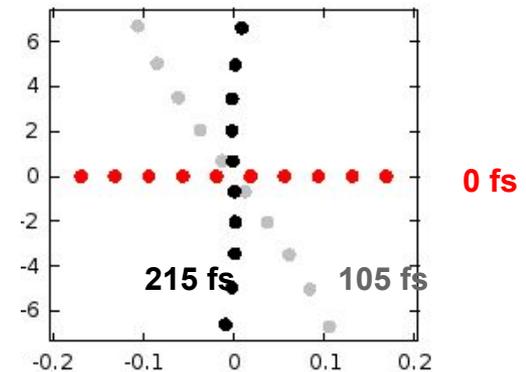
focal length = 18.0 μm

Modified geometry
thinner columns
(-10%)



focal length =
14.8 μm

Modified geometry
thicker columns
(+10%)

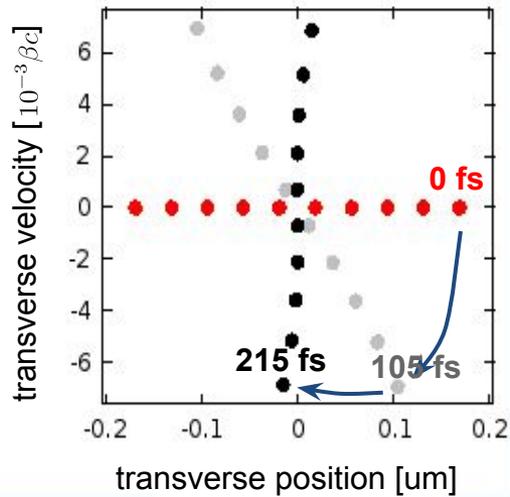


focal length = 18.8 μm

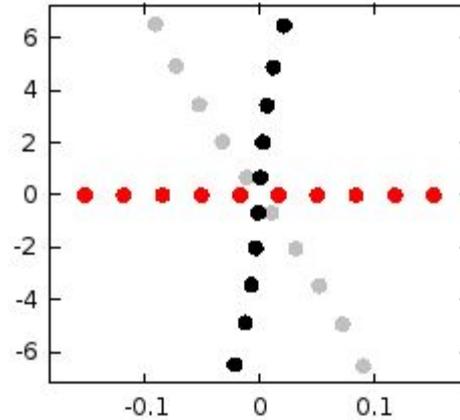
2. Sensitivity to vacuum channel width

3. Sensitivity to laser phase mismatch

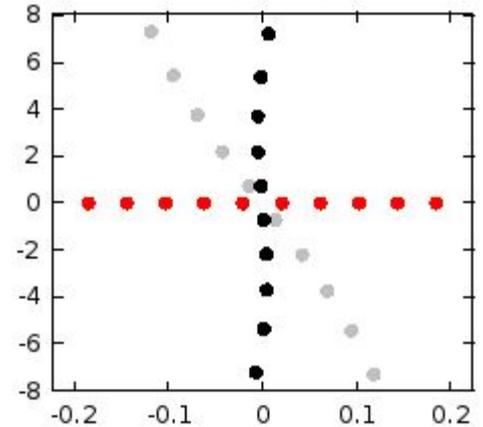
Design geometry,
 $f = 18.0 \text{ } \mu\text{m}$



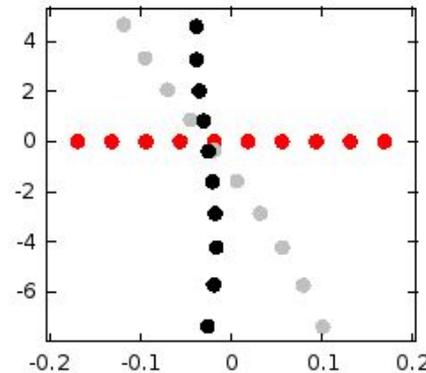
Narrower vacuum
channel (-10%)
 $\Rightarrow f = 16.5 \text{ } \mu\text{m}$



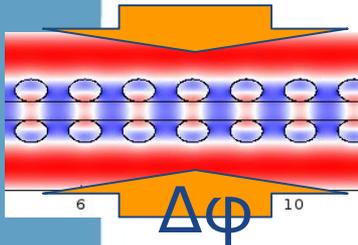
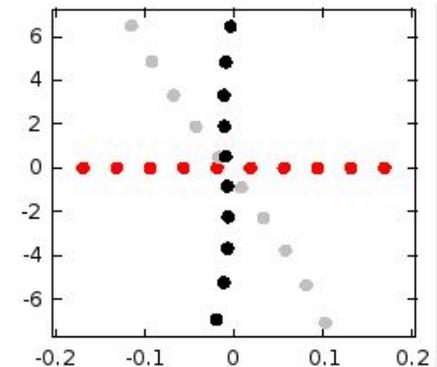
Wider vacuum
channel (+10%)
 $\Rightarrow f = 19.7 \text{ } \mu\text{m}$

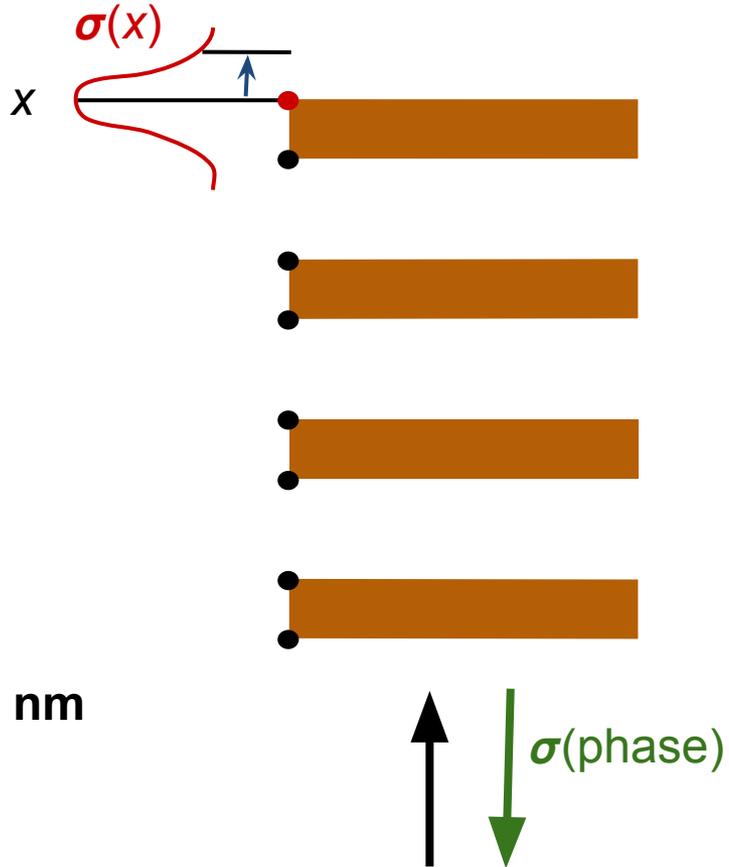


Laser#1 phase change (-30°)
 $\Rightarrow f = 22.2 \text{ } \mu\text{m}$ and
focus transverse shift -26 nm



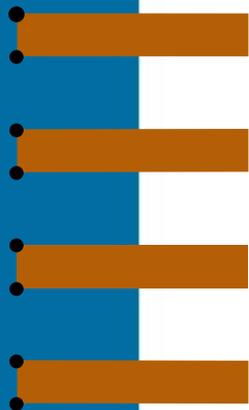
Laser#1 phase change (+30°)
 $\Rightarrow f = 19.6 \text{ } \mu\text{m}$ and
focus transverse shift -8 nm





$\lambda = 2 \text{ } \mu\text{m}$
 $\lambda / 4 = 500 \text{ nm}$

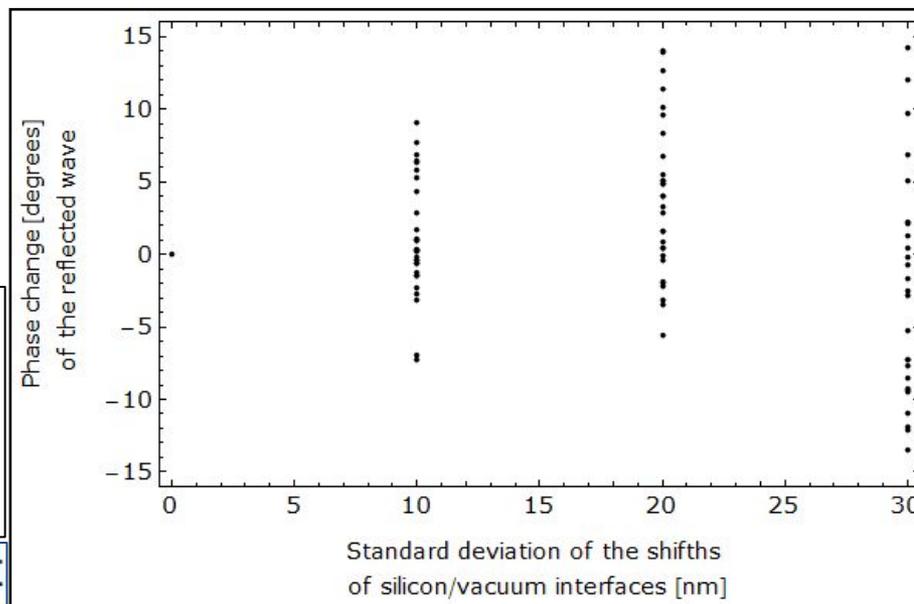
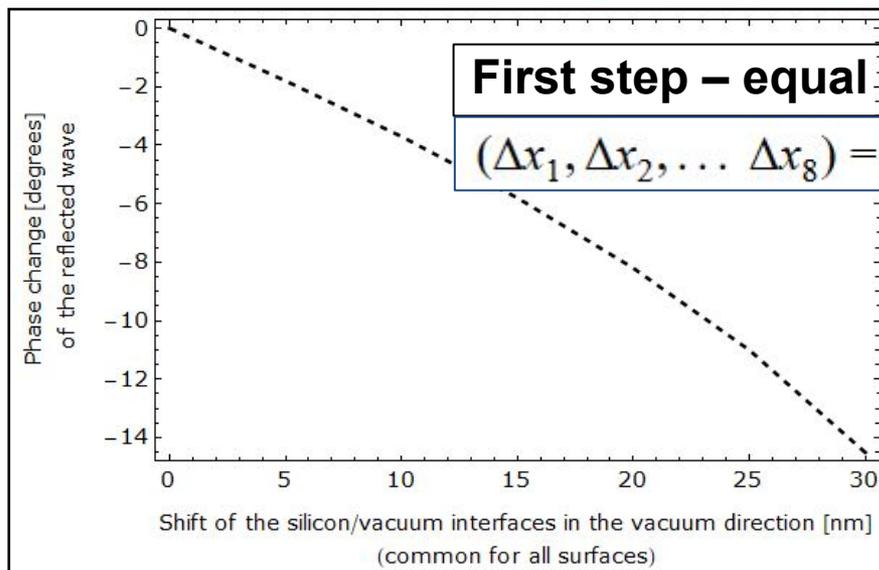
Bragg mirror Systematic vs. random interface shifts



$\lambda = 2 \text{ } \mu\text{m}$
 $\lambda / 4 = 500 \text{ nm}$

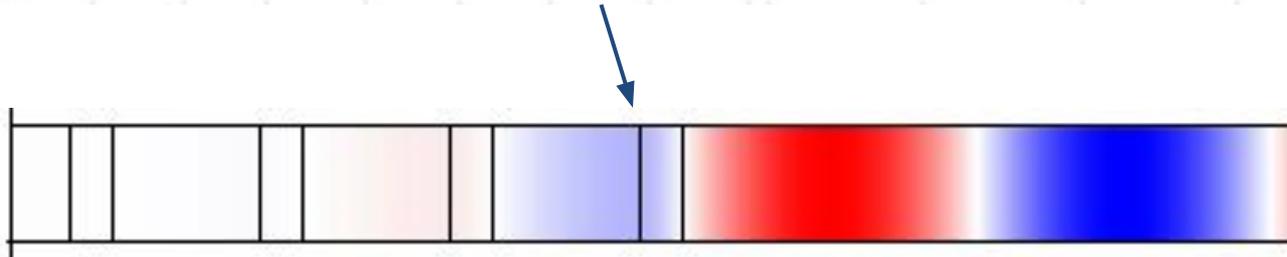
**Random, independent
interface shifts, 30
samples**

$(\Delta x_1, \Delta x_2, \dots, \Delta x_8) = \text{random set}$



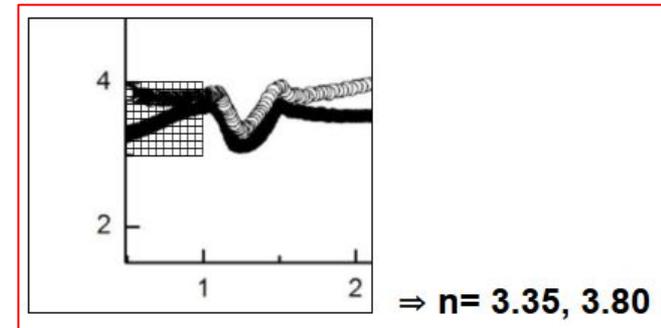
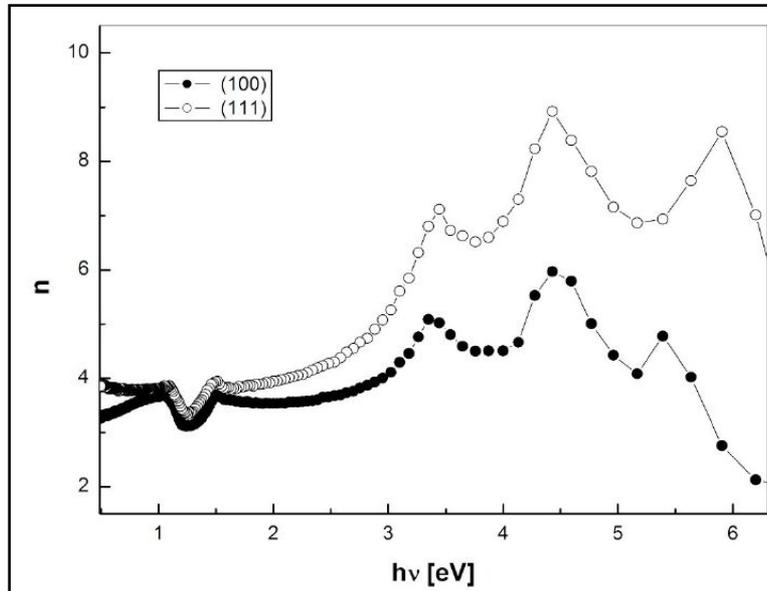
Bragg mirror Shifting only one interface

i8	i7	i6	i5	i4	i3	i2	i1	R (1)	arg(e...	abs(e...
0	0	0	0	0	0	0	30	0.99978	-0.35824	0.99989
0	0	0	0	0	0	30	0	0.99978	-10.15...	0.99989
0	0	0	0	0	30	0	0	0.99978	-0.02929	0.99989
0	0	0	0	30	0	0	0	0.99978	-0.8596	0.99989
0	0	0	30	0	0	0	0	0.99978	-0.00145	0.99989
0	0	30	0	0	0	0	0	0.99978	-0.07166	0.99989
0	30	0	0	0	0	0	0	0.99978	0.00142	0.99989
30	0	0	0	0	0	0	0	0.99978	-0.00505	0.99989



For the phase of the reflected wave, the position of the second interface is crucial ("back of the first layer").

[Yousefi *et al.* 2019] use "phosphorus-doped Si<100>" \Rightarrow crystalline Si



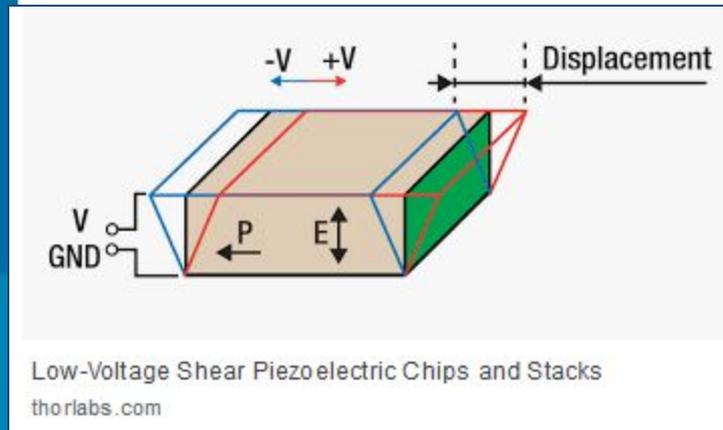
M. El-Nahass and H. Ali, *Estimation of optical parameters of silicon single crystals with different orientations*, Materials Science-Poland, 37, 65 (2019)

$\Rightarrow n = 3.35, 3.80, 13\%$ difference at 2 μm ($h\nu = 0.62$ eV)

Amorphous silicon at 2 μm : $n = 3.45$ (3.51 @ 600 K)

<https://refractiveindex.info/?shelf=main&book=Si&page=Li-293K>

shear piezos



DLA section 1

DLA section 1

How to align DLA segments?

2018 Kozak *et al*, *Ponderomotive...* - Supplemental Material

“One translation stage in the ω_1 interferometer is equipped with a piezo-crystal allowing to change its position with a precision of 10 nm.”

2018 Leedle *et al*. “The relative phase of the two lasers is controlled via a piezo-driven delay stage.”

3 Technical Aspects of Scanning Probe Techniques

Bert Voigtländer

3.8

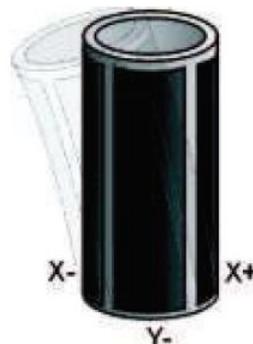
Bert Voigtländer

1.2 Tube Scanner

One central task in scanning probe microscopy is to position the scanning probe with an accuracy of less than one tenth of an Angstrom in all three dimensions. The tube scanner is the most widely used actuator element to move the probe tip or the sample in order to scan a surface (fine motion). This is due to its simple design, its high piezo constants and high resonance frequencies. It consists of a tube, made out of piezoceramics, which is covered inside and outside with metal, which acts as electrodes. The outer electrode is sectioned into four quadrants. The tube is poled in radial direction. One of the great advantages of the tube scanner is that motions in all three dimensions can be controlled by one very simple actuator element. A motion in z-

Technical Aspects of Scanning Probe Techniques

3.9



BaTiO₃ (Barium titanate, BTO)

nature
materials

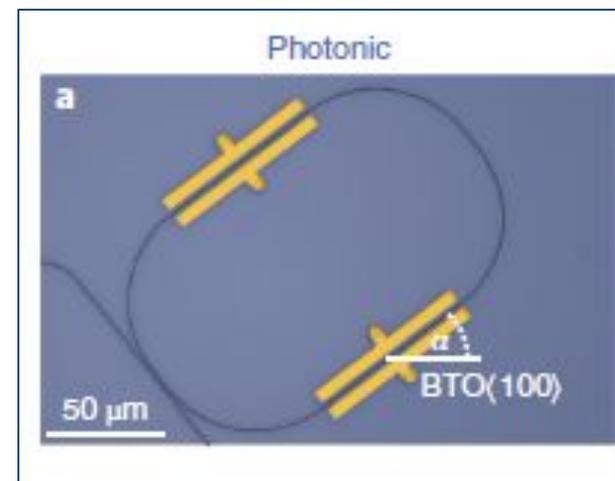
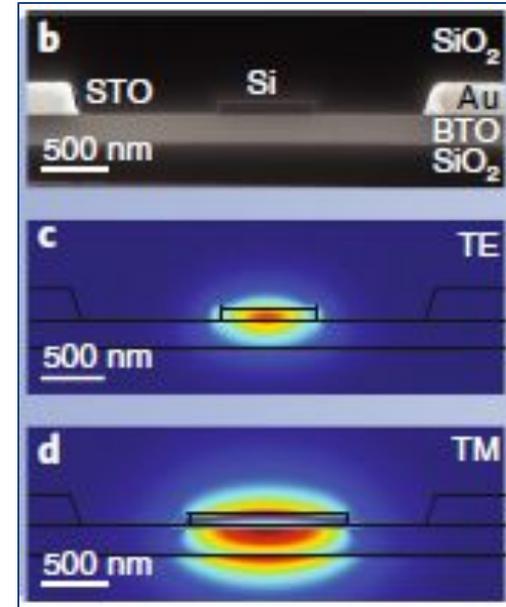
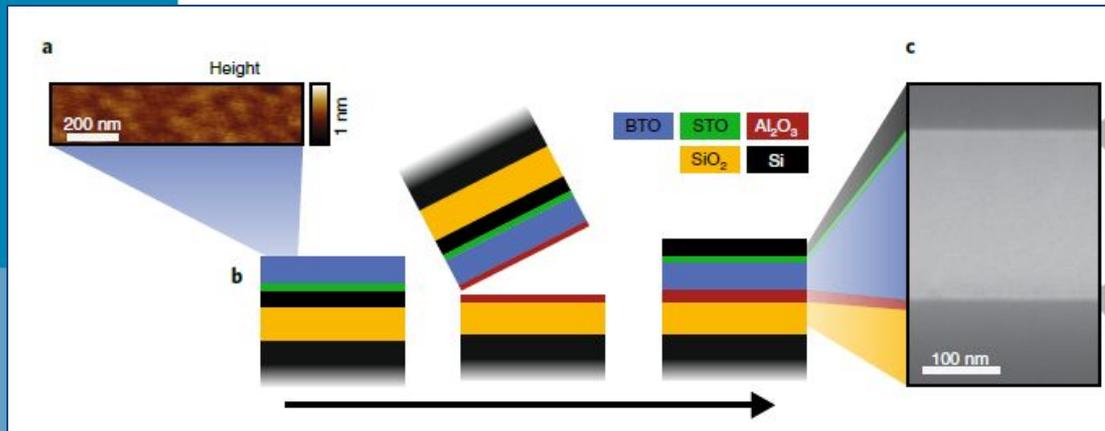
ARTICLES

<https://doi.org/10.1038/s41563-018-0208-0>

(2019)

Large Pockels effect in micro- and nanostructured barium titanate integrated on silicon

Stefan Abel^{1,7*}, Felix Eltes^{1,7}, J. Elliott Ortmann², Andreas Messner³, Pau Castera⁴, Tino Wagner⁵, Darius Urbonas¹, Alvaro Rosa⁴, Ana M. Gutierrez⁴, Domenico Tulli⁶, Ping Ma^{3*}, Benedikt Baeuerle³, Arne Josten³, Wolfgang Heni³, Daniele Caimi¹, Lukas Czornomaz¹, Alexander A. Demkov², Juerg Leuthold³, Pablo Sanchis^{4*} and Jean Fompeyrine¹



BaTiO₃ (Barium titanate, BTO)

refractive index ~ 2

largest coefficient of the Pockels tensor $r_{42} \sim 1 \text{ nm/V}$

Electric field-induced change of coefficients of the refractive index ellipsoid

$$\Delta \left(\frac{1}{n^2} \right)_i = \sum_{j=1}^3 r_{ij} E_j, \quad i = 1 \dots 6$$

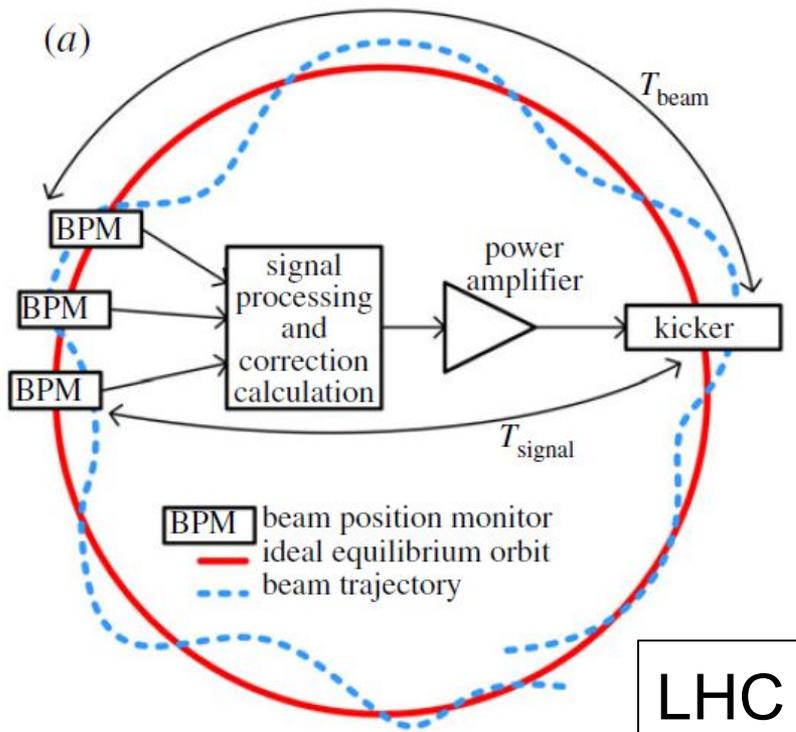
$$\Delta n \sim \frac{n^3}{2} \Delta \left(\frac{1}{n^2} \right) \sim \frac{n^3}{2} r_{42} E_2$$

$$\text{e.g. } \Delta n \sim \frac{2^3}{2} \cdot \frac{1 \text{ nm}}{\text{V}} \cdot \frac{10 \text{ V}}{200 \text{ nm}} = 0.2$$

[www.rp-photonics.com/pockels_effect.html]

10. Transverse beam feedback

The beams in accelerators are inherently unstable. Many mechanisms are used in order to maintain the beam and its quality for many hours of circulation in the vacuum chamber. A crucial feedback system for the LHC is depicted in figure 26a. The transverse oscillations that appear at the onset of instability are measured in one location of the circumference by a suitable beam position monitor (BPM). This signal is processed electronically and a corrective 'kick' signal is generated, amplified and applied to the beam by a transverse electromagnetic field ('kicker'). Figure 26b shows the system as installed in the LHC tunnel.



On-Chip Laser Power Delivery System for Dielectric Laser Accelerators

Tyler W. Hughes,* Si Tan,† Zhexin Zhao, Neil V. Sapro, Kenneth J. Leedle, Huiyang Deng, Yu Miao, Dylan S. Black, Olav Solgaard, James S. Harris, Jelena Vuckovic, Robert L. Byer, and Shanhui Fan
Stanford University, Stanford, CA 94305

Yun Jo Lee and Minghao Qi
Purdue University, West Lafayette, IN 47907

(ACHIP Collaboration)
(Dated: September 14, 2017)

2017

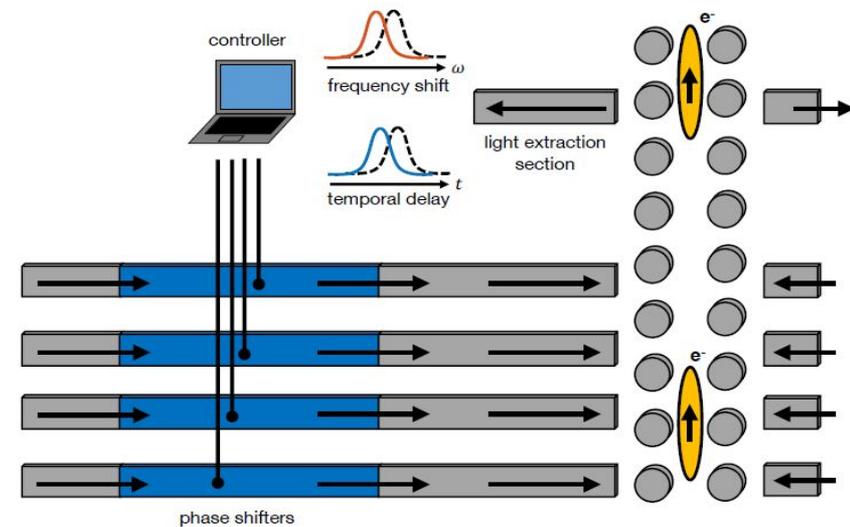
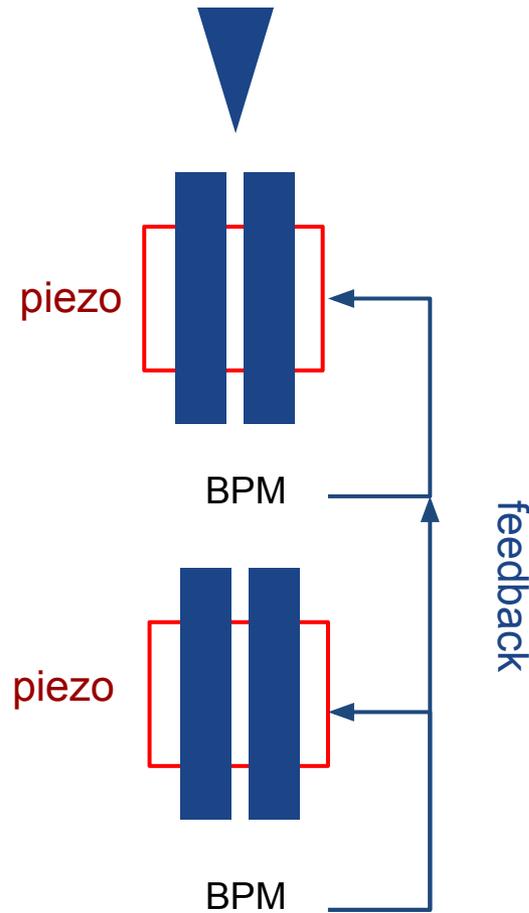


FIG. 7. Feedback system for automatic phase control. A dedicated light extraction section is added to the accelerator. Light is radiated from the electron beam transverse the DLA structures and the frequency content and/or timing of the light is sent to a controller. The phase shifts of each waveguide are optimized with respect to either the frequency or the delay of the signal. After several runs, the system should converge to stable operation.

From: S. Myers, *The engineering needed for particle physics*, Phil. Trans. R. Soc. A370, 3887–3923 (2012)

Structure tuning in DLA?





1. Fabrication tolerances are not an issue for recent proof-of-concept DLA experiments utilizing short structures.
2. Fabrication constraints become tighter with increasing structure length, and post-fabrication adjustment of structures may become necessary:
 - electromechanical: integrated shear piezo crystals? piezotubes?
 - electrooptical: Pockels effect?

Acknowledgments

I am grateful to Joel England, Levi Schächter and Koby Scheuer for discussions on DLA, to Oskar Warmusz for help with particle tracing in Comsol, to Marek Nowicki for his comments on piezo crystals, and to Janusz Przesławski for his comments on the electrooptical effects.

Today's DLA posters →



- 116. Frank Mayet**, Status report on the dielectric laser acceleration experiments at the SINBAD/ARES linac
- 121. Thilo Egenolf**, Transverse Beam Breakup Instability in Dielectric Laser Accelerators
- 159. Andrzej Szczepkowicz**, Numerical calculation of the Purcell-Smith radiation from dielectric laser acceleration (DLA) structures
- 284. Stefanie Kraus**, A compact UHV-compatible high-voltage supply for a small dielectric laser accelerator
- 318. Johannes Illmer**, Alternating phase focusing in dielectric laser acceleration
- 331. Norbert Schönenberger**, Generation and characterization of attosecond micro-bunched electron pulse trains via dielectric laser acceleration



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