

Chromatic Effects on Plasma Channels Formed with an Axicon Lens

A J Ross, A Alejo, A von Boetticher, J Cowley, J Holloway, J Jonnerby, A Picksley
R Walczak, S M Hooker

John Adams Institute for Accelerator Science and Department of Physics, University of Oxford, Denys Wilkinson Building, Keble Road, Oxford OX1 3RH, UK

Introduction

- To reach 10s of GeV in a single laser-plasma accelerator stage, the driving laser will need to remain focused over 100s of mm.
- One way to achieve this is with a HOFI plasma channel [1] which can be formed using an axicon, as successfully demonstrated by the Oxford group using the RAL Astra-Gemini TA2 laser facility [2].
- During the experiment, it was found that ~ 10 higher intensity was required on the axicon to form the plasma channel than predicted.
- To investigate this, a numerical code has been written to model the effect of dispersion through an axicon.

Axicons

A refractive axicon is a conical lens which forms an extended line focus.

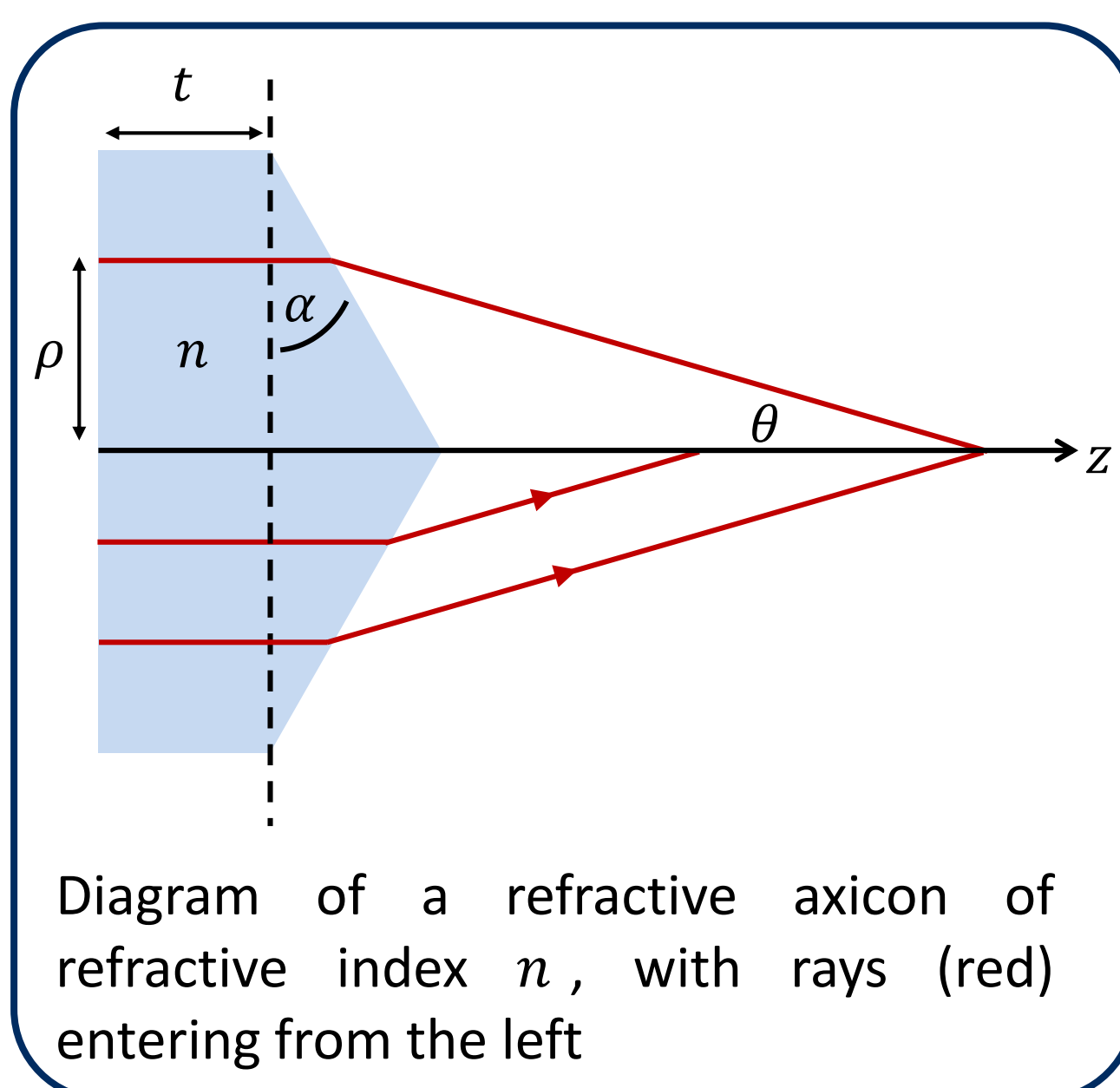
The transmitted beam has a characteristic Bessel shaped profile:

$$U_{\text{out}} \propto J_0[(n-1)k\rho'\alpha]$$

at output radius ρ' for wavelength $\lambda = 2\pi/k$.

Chromatic aberrations may arise in the axicon focus due to;

- Direct k dependence of the Bessel function.
- Dispersion in $n(k)$, causing a variation in approach angle θ and pulse stretching.

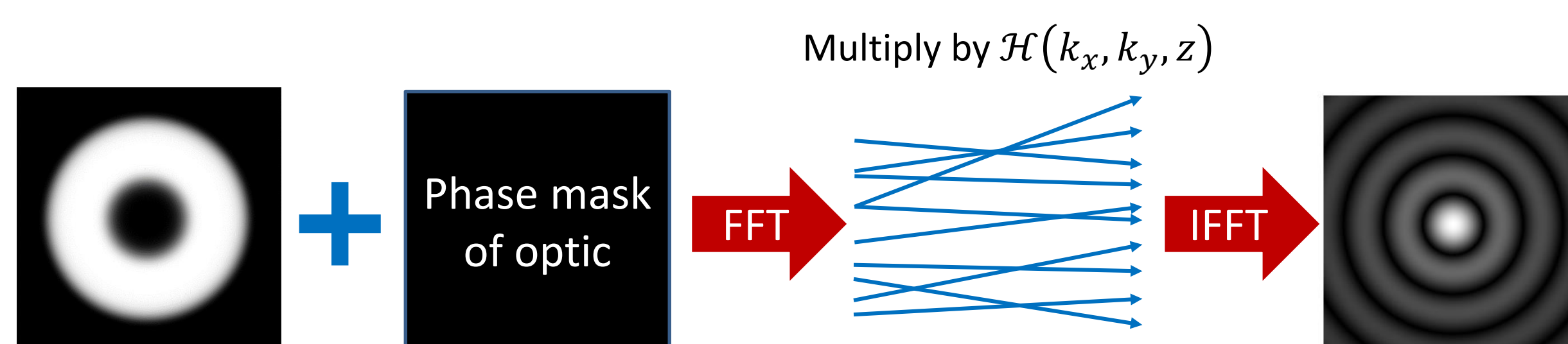


Numerical Code - ASM

The Angular Spectrum Method can be used to propagate an initial transverse field, $U(x, y, 0)$, to output plane z via the following steps;

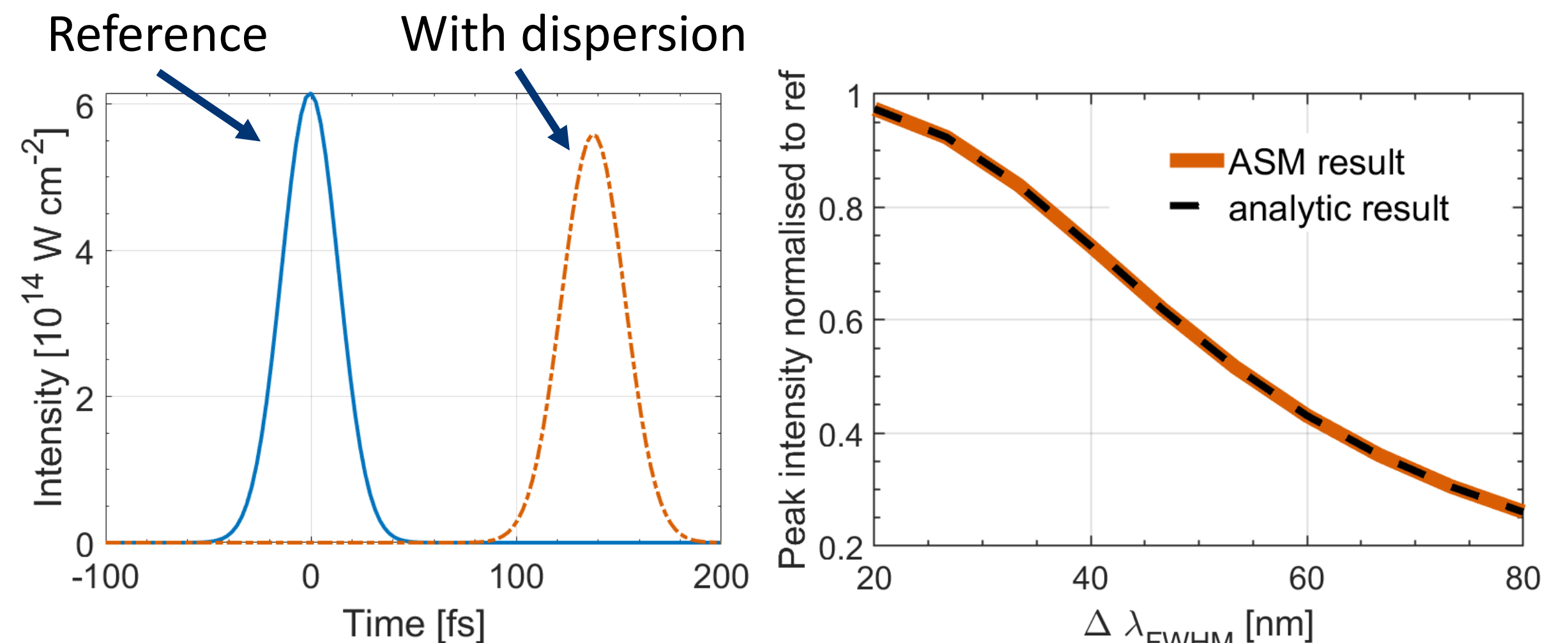
- Transform $U(x, y, 0)$ into wavenumber (k_x, k_y) space.
- Multiply by the propagator function $\mathcal{H}(k_x, k_y, z) = \exp(-i k_z z)$.
- Perform inverse Fourier transform to retrieve new field, $U(x, y, z)$.

This process is for a single wavelength. To model a laser pulse, the code first discretizes the spectrum of wavelengths, propagates each wavelength using ASM and sums the resulting complex fields (at a fixed (x, y) position).



Flow chart illustrating how the ASM numerical code works for each wavelength.

Results

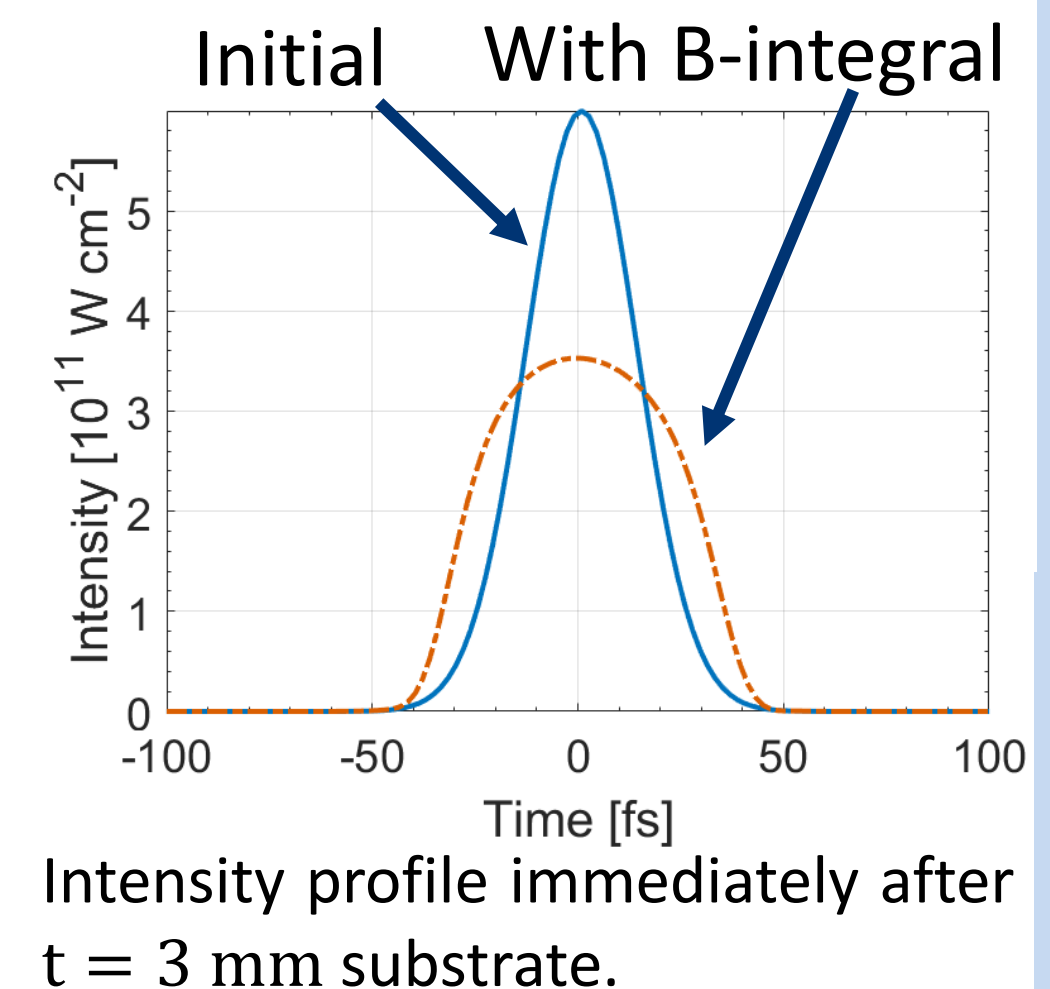


Temporal intensity profiles at focus for initial Gaussian beam with $I_0 = 10^{11} \text{ W cm}^{-2}$, $\lambda = 805 \text{ nm}$, $\Delta\lambda_{\text{FWHM}} = 30 \text{ nm}$ (left). Axicon parameters; $\alpha = 5.52^\circ$, $t = 3 \text{ mm}$, fused silica. Initial beam radius $25.4 \times 0.1 \text{ mm}$. Reference has $n(\lambda) = \text{const}$.

Other Factors

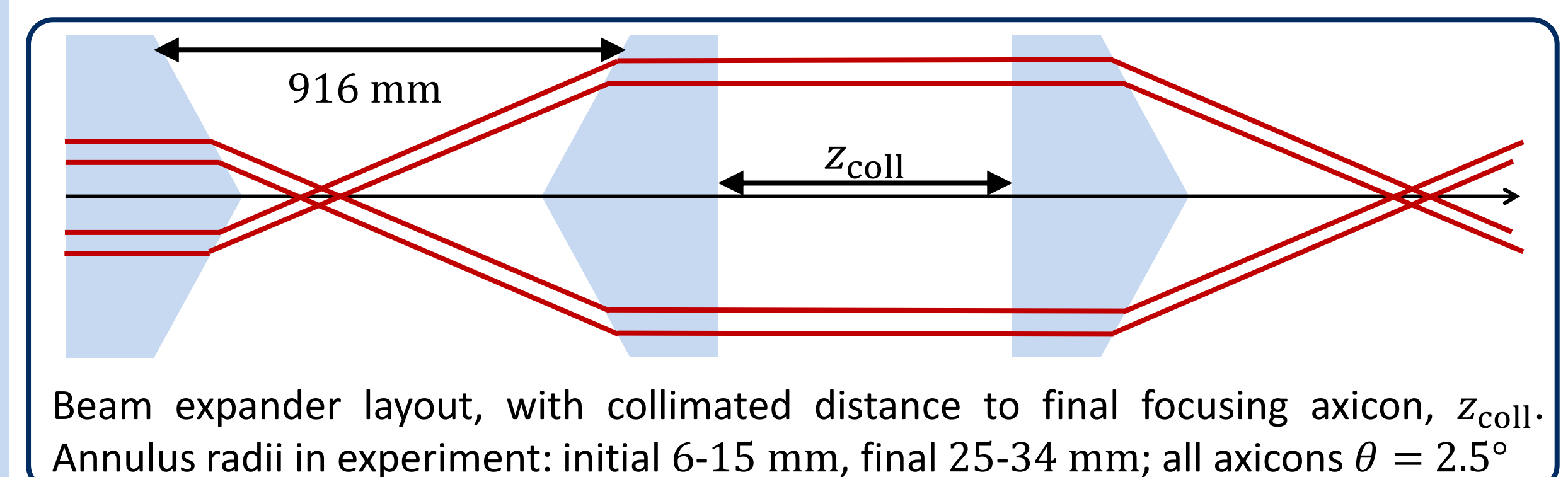
Factors which may have further reduced the on-axis intensity are;

- B-integral through the axicon (Right).
- Imperfections in the experimental near field of the laser.
- The ionisation threshold of the H_2 gas not matching predictions.
- The plasma's own refractive index distorting the Bessel focus.

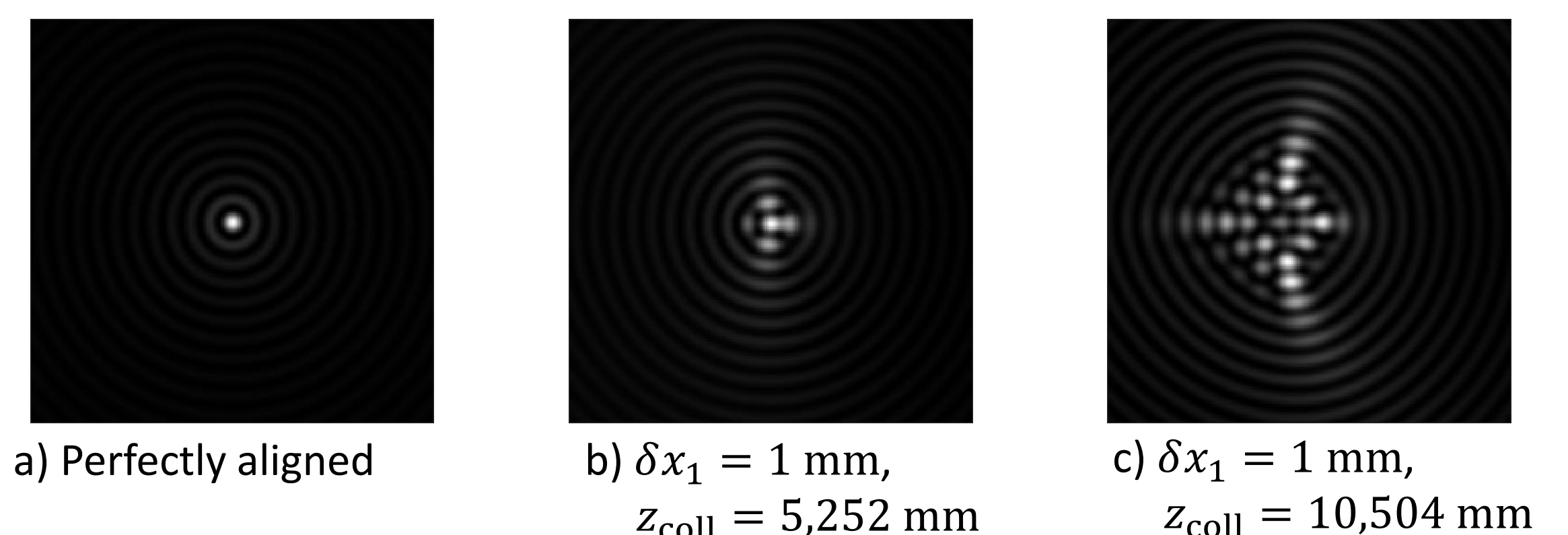


Axicon Beam-Expander

The ASM code can be applied to other optics studies, such as of the aberrations induced by a mis-aligned two axicon beam-expander.



In experiment, focal spot quality was very sensitive to misalignments of the axicons in the expander. The effect of a small offset of the first axicon, δx_1 is modelled below (all distances scaled down by 10).



Conclusions and Future Work

- The study of chromatic effects through refractive axicons has revealed that dispersion reduces the peak intensity by $\sim 10\%$.
- This is **not** enough to fully account for the increased intensity that was required to form the plasma channel on experiment.
- The ASM code can also be applied to other more exotic phenomena, such as curved plasma channels and flying foci.

[1] R.J. Shalloo *et al.*, *Physical Review E*, **97**(5), 1–8 (2018)

[2] R.J. Shalloo *et al.*, *Physical Review Accel. Beams*, **22**(4), 1–6 (2019)

[3] Y. Wang *et al.*, *Journal of the Optical Society of America A* **34**(7), (2017)

[4] J. W. Goodman, *Introduction to Fourier Optics*, (2005)

Work supported by the UK Science and Technology Facilities Council (STFC UK) grant number ST/S505833/1. JAI rolling grant: ST/P002048/1. This material is based upon work supported by the Air Force Office of Scientific Research under award number FA9550-18-1-7005. This work was supported by the European Union's Horizon 2020 research and innovation programme under grant agreement No. 653782. The author acknowledges the support of the EAAC2019 Student Grant.