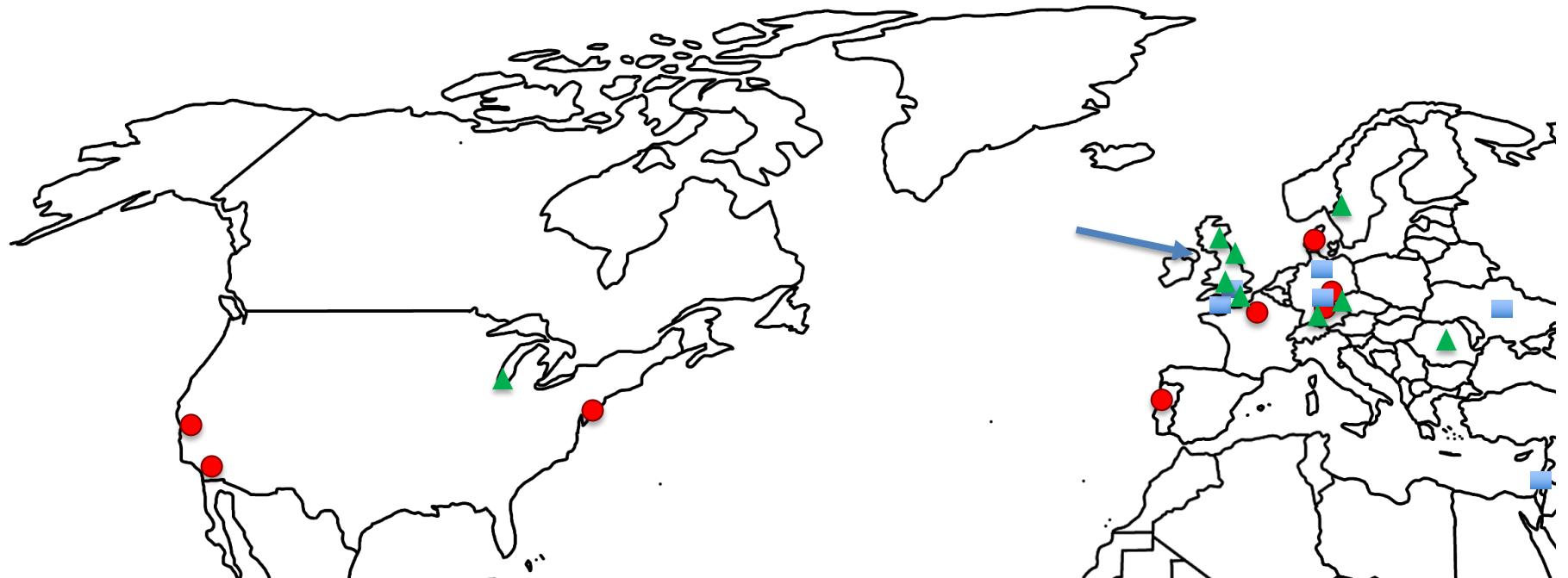


High-intensity quantum electrodynamics in the field of an ultra-intense laser

Gianluca Sarri

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Main Collaborators



● E-320 experiment at FACET-II 

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■ LUXE experiment at Eu-XFEL



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 מִן וַיְצָמֵן לְבָנָע
WEIZMANN INSTITUTE OF SCIENCE

 UCL

 KINR
KIEV INSTITUTE FOR NUCLEAR RESEARCH

 TEL AVIV
UNIVERSITY

▲ Laser-driven experiments



Imperial College
London

of York

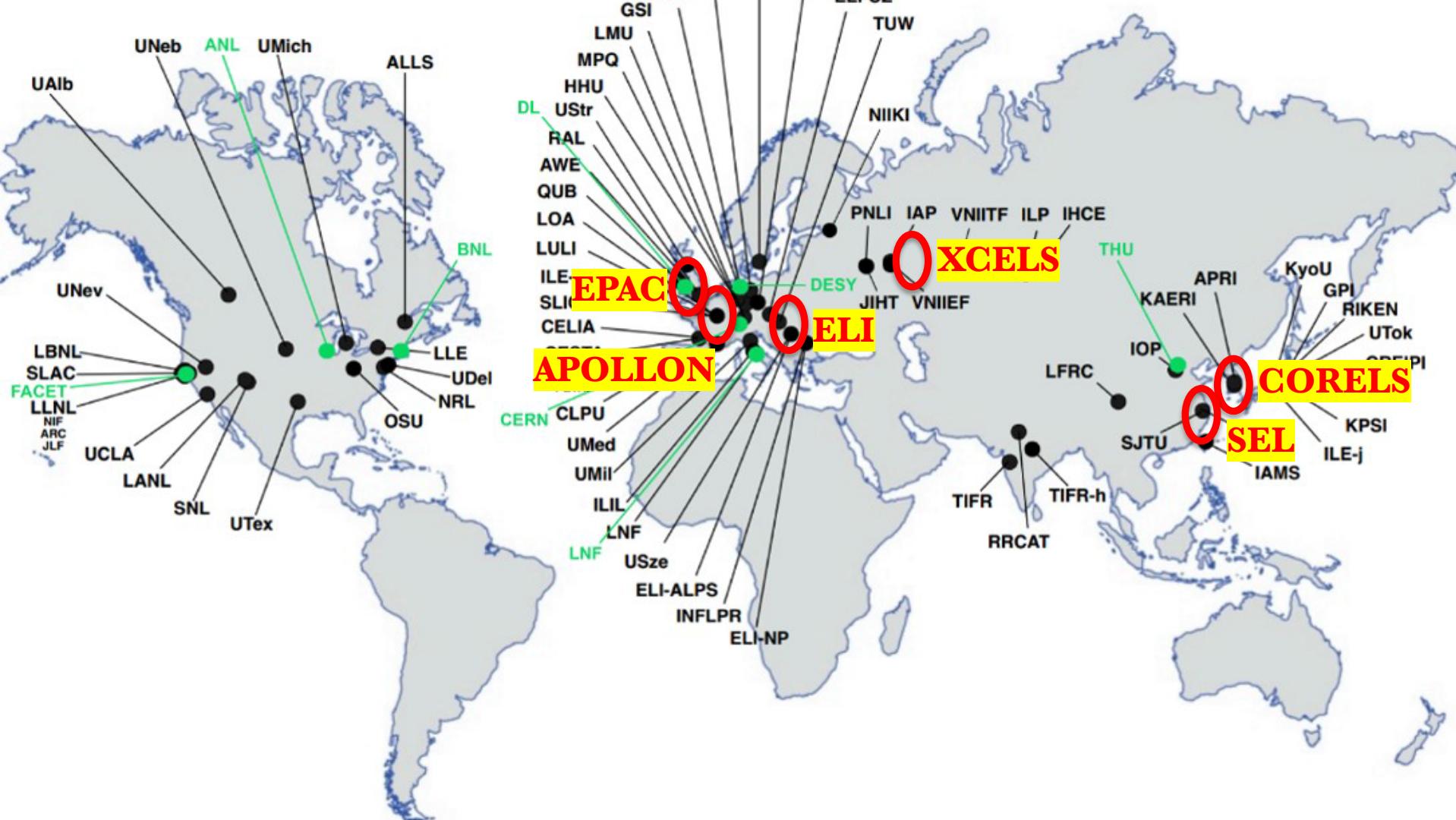


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 CHALMERS
UNIVERSITY OF TECHNOLOGY

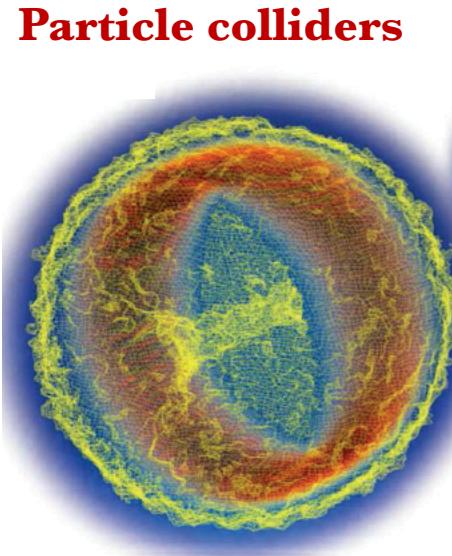




Part 1

Introduction

Particle colliders



250 GeV cm energy
 strong coupling
 $\alpha\chi^{2/3} > 1$

V. Yakimenko et al. PRL 2019

Astrophysics

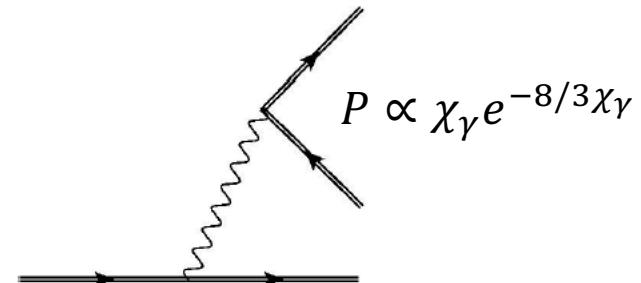
surface magnetic field of magnetars
 $\sim 10^{-4} - 10 B_{\text{cr}}$

B. Cerutti Space Sci. Rev. 2017

Fundamental physics

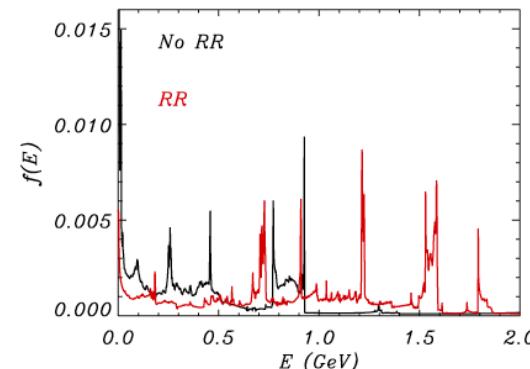
perturbative QED: $\alpha \ll 1$

non-perturbative phenomena:



V. I. Ritus J. Russ. Laser Res. 1985

Plasma physics



M. Tamburini et al. NJP 2010

⇒ **Radiation Reaction** is one of the oldest and most fundamental problems in electromagnetism:
How do we correctly model the electron dynamics if we include radiative losses?

0. Classical Lorentz force

$$m \frac{du^u}{ds} = e F^{uv} u_v$$

X No energy loss

1. LAD Equation

$$m \frac{du^u}{ds} = e F^{uv} u_v + \frac{2}{3} e^2 \left(\frac{d^2 u^u}{ds^2} + \frac{du^v}{ds} \frac{du^v}{ds} u^u \right)$$

Schott's term

- ✓ Damping force (radiation reaction term)
- X** Classical renormalisation (point-like electron)
- X** Runaway solutions! (diverging acceleration even without external field)

2. LL Equation

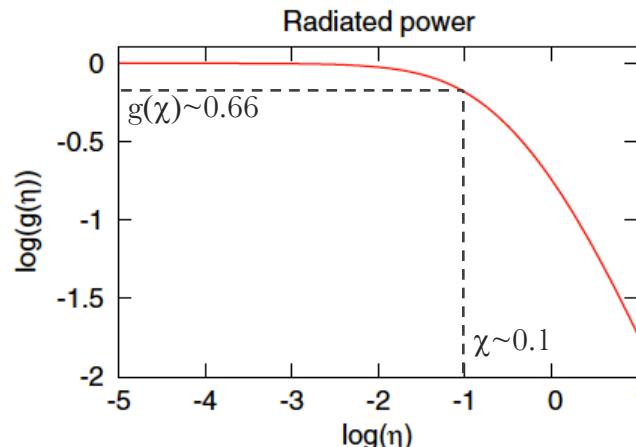
$$m \frac{du^u}{ds} = e F^{uv} u_v + \frac{2}{3} e^2 \left(\frac{e}{m} (\partial_\alpha F^{uv}) u^\alpha u_v - \frac{e^2}{m^2} F^{uv} F_{\alpha v} u^\alpha + \frac{e^2}{m^2} (F^{\alpha v} u_v) (F_{\alpha \lambda} u^\lambda) u^u \right)$$

- ✓ No runaway solutions
- ✓ Valid in special relativity

$\lambda \gg \alpha \lambda_C$ (localised wavefunction)
 $F \ll F_{cr}/\alpha$ (classical critical field)

→ The classical treatment of radiation reaction neglects three main additional phenomena:

1. The energy of a single emitted photon can not exceed that of the electron



Generally speaking, this leads to a classical overestimate of the total energy loss experienced by the electron ($\chi = \gamma F_L / F_S$)

$$g(\chi) \sim (3.7\chi^3 + 31\chi^2 + 12\chi + 1)^{-4/9}$$

J. G. Kirk et al., PPCF 2013

A. G. R. Thomas et al., PRX 2012

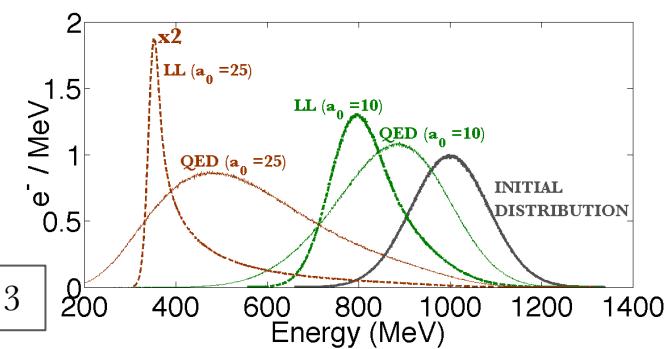
2. Photon emission is probabilistic

2.a $a_0 \gg 1$

2.b constant cross-field approximation
(instantaneous photon emission))

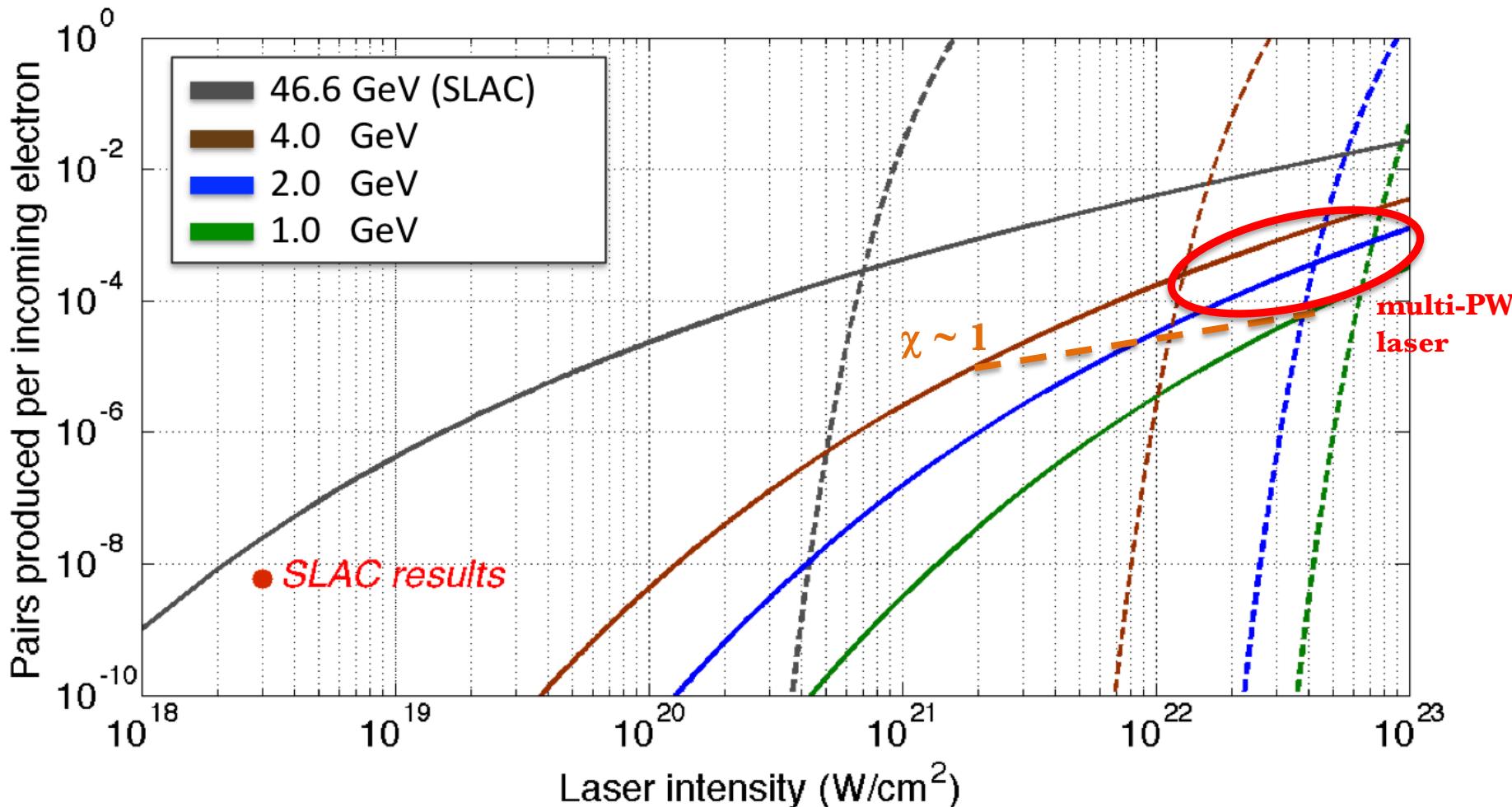
V. I. Ritus, J. Sov. Laser Res. 1985

N. Neitz and A. Di Piazza, PRL 2013



3. Production of electron-positron pairs (important only for $\chi \geq 1$)

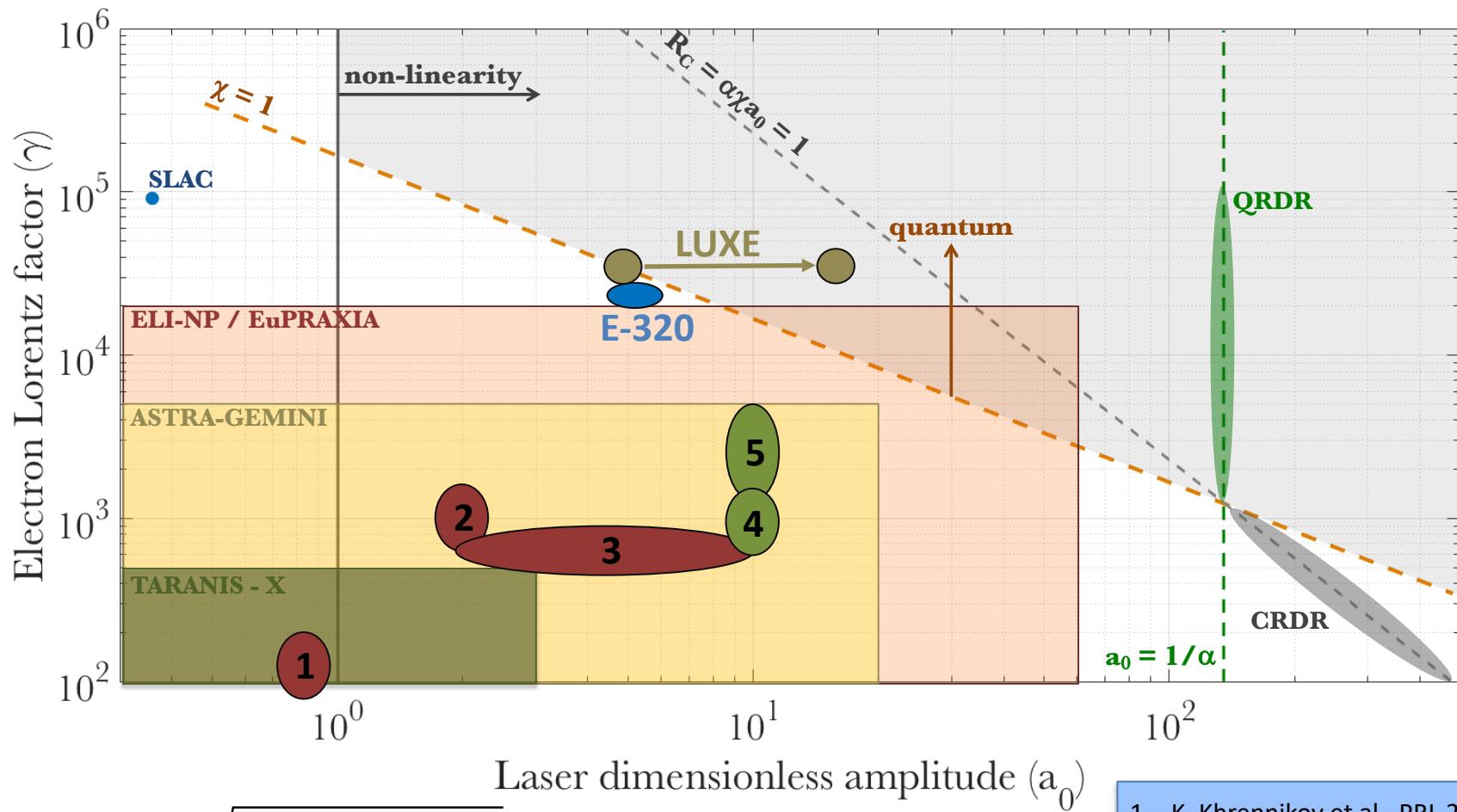
→ Creation of electron-positron pairs becomes significant



Part 2

Recent experiments on

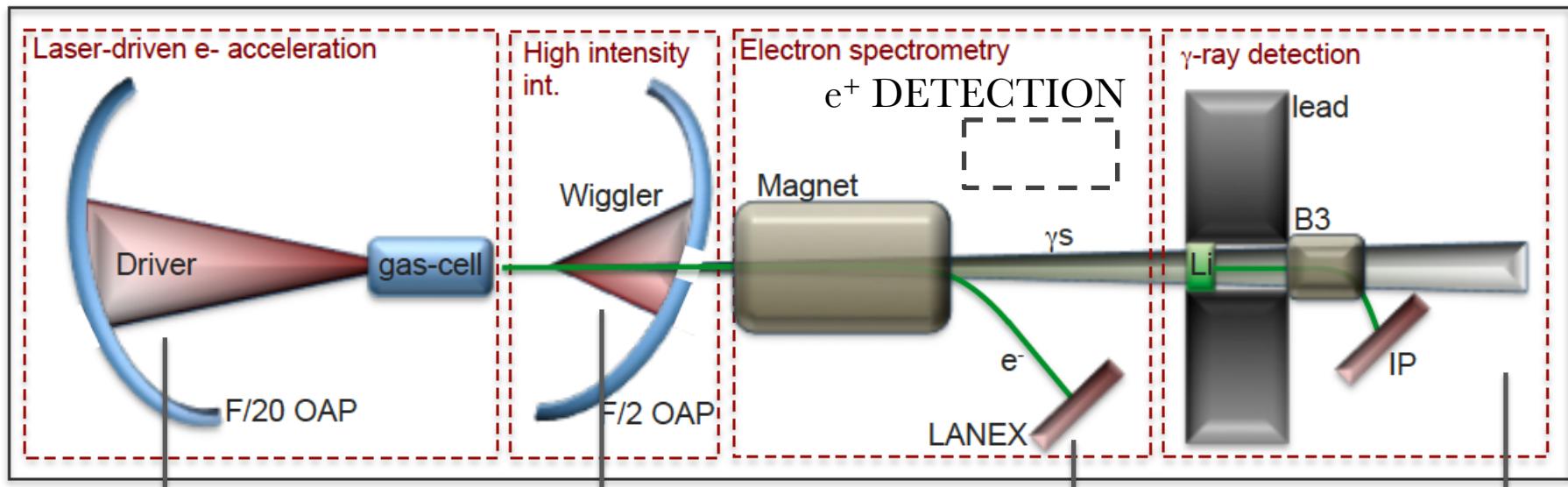
quantum radiation reaction



$$a_0 \sim 6\lambda_L [\mu\text{m}] \sqrt{I_L [10^{20} \text{W/cm}^2]} \text{ Intensity}$$

$$\chi \sim 6.1 \times 10^{-6} \gamma_e a_0 \quad \text{Quantum parameter}$$

1. K. Khrennikov et al., PRL 2015
2. G. Sarri et al., PRL 2014
3. W. Yan et al., Nat. Phot. 2017
4. J. Cole et al., PRX 2018
5. K. Poder et al, PRX 2018



DRIVER LASER:

- $t_L \sim 30\text{-}40 \text{ fs}$
- $E \sim 1 - 10 \text{ J}$
- Long F# ($F/20 - F/40$)
- $a_0 \sim 1.5 - 2$

SCATTERING LASER:

- $t_L \sim 30\text{-}40 \text{ fs}$
- $E \sim 1 - 10 \text{ J}$
- Short F# ($F/2$)
- $a_0 \gg 1 (\sim 10)$

e⁻ SPECTROMETER

γ -ray SPECTROMETER

X Unstable electron beam spectrum

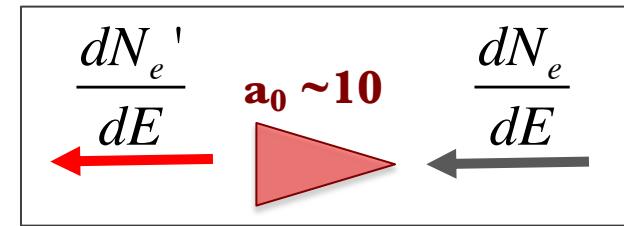
G. Samarin et al., J. Mod. Opt. 65, 1362 (2017)

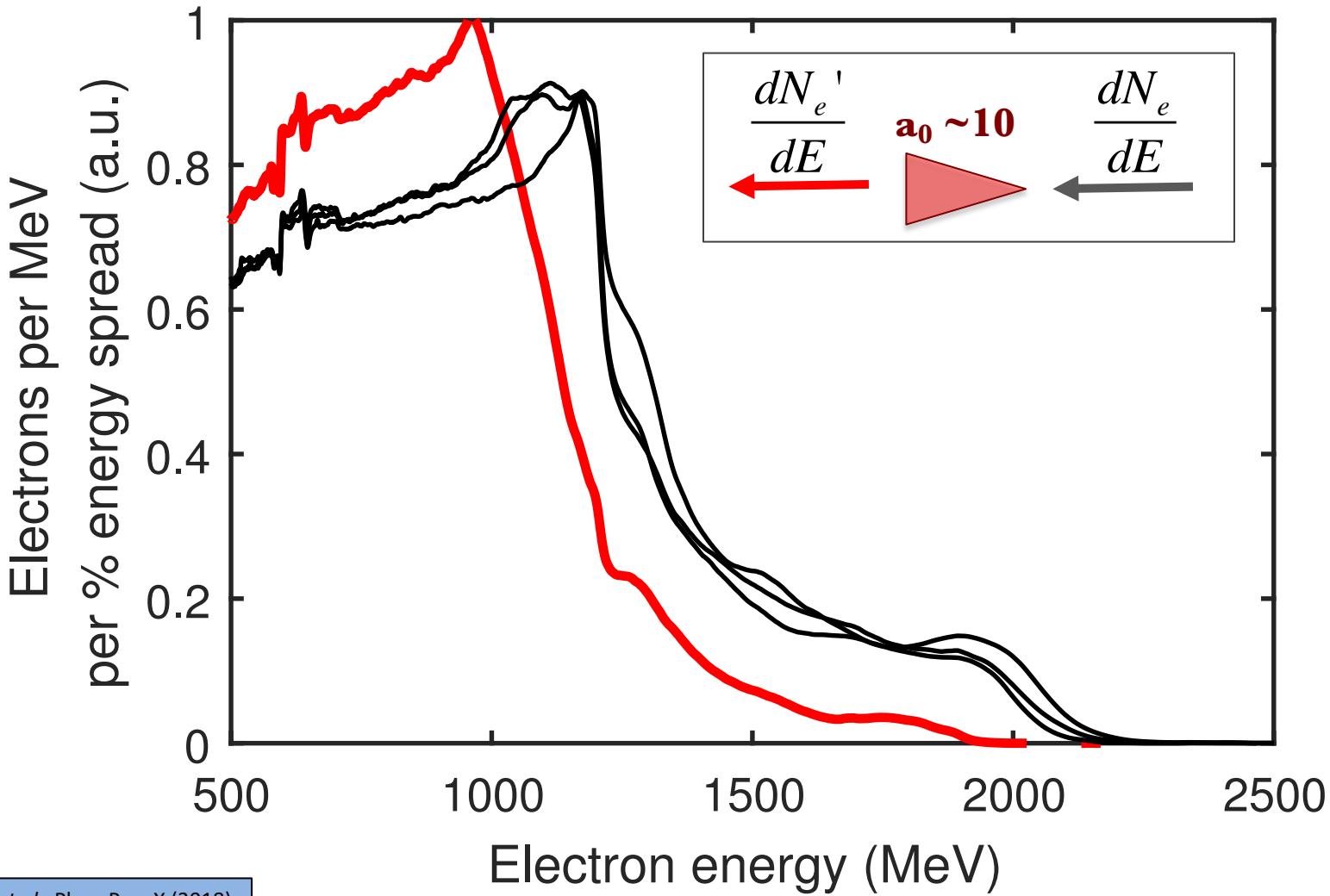
X Pointing fluctuations

X fs-scale synchronisation

D. Corvan et al., Opt. Express 24, 3127 (2016).

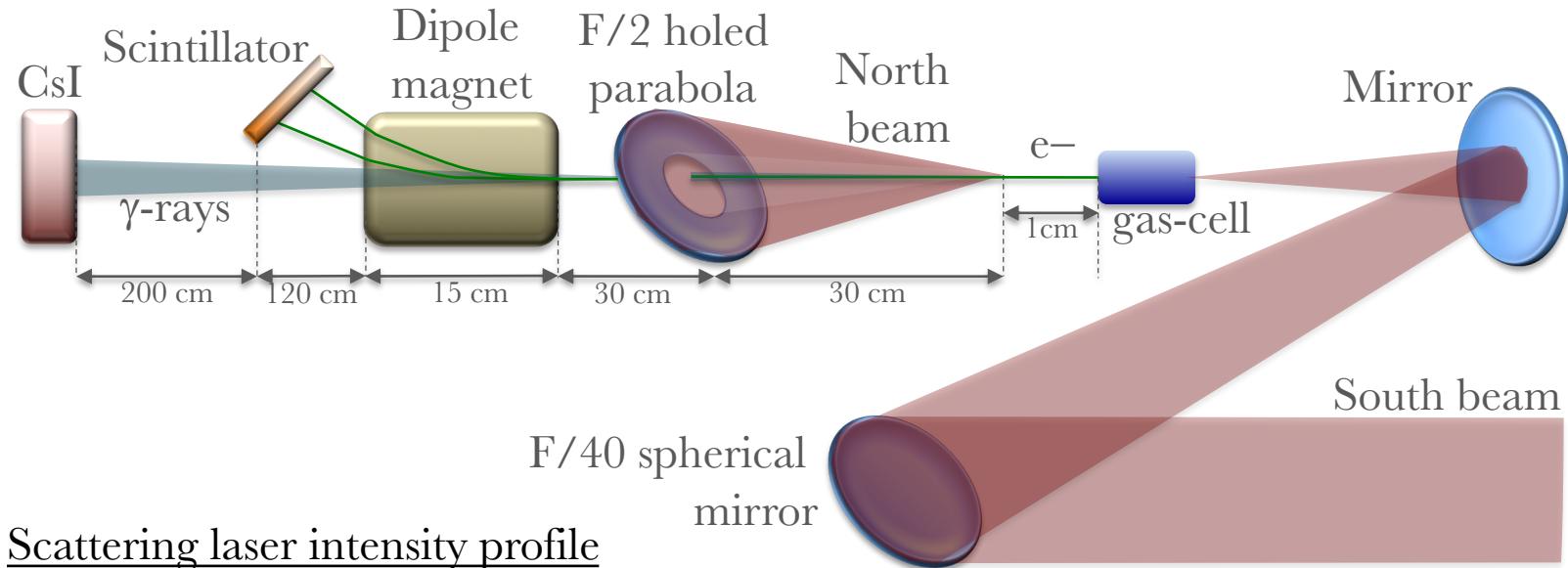
What do we see?



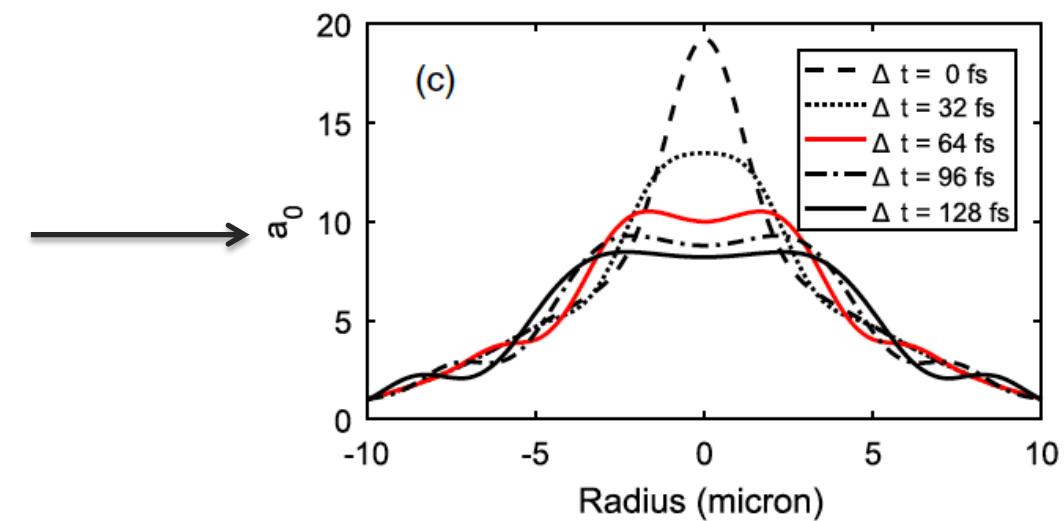
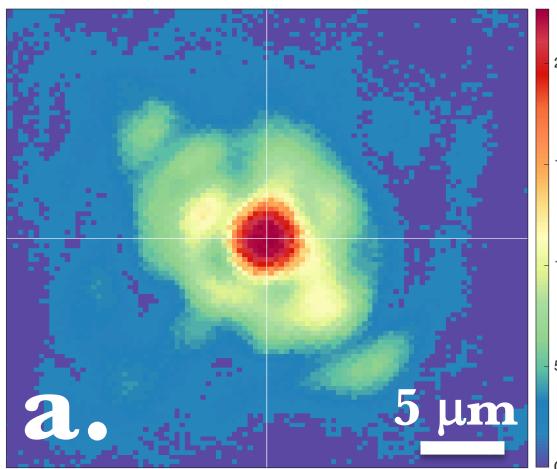


K. Poder *et al.*, Phys. Rev. X (2018)

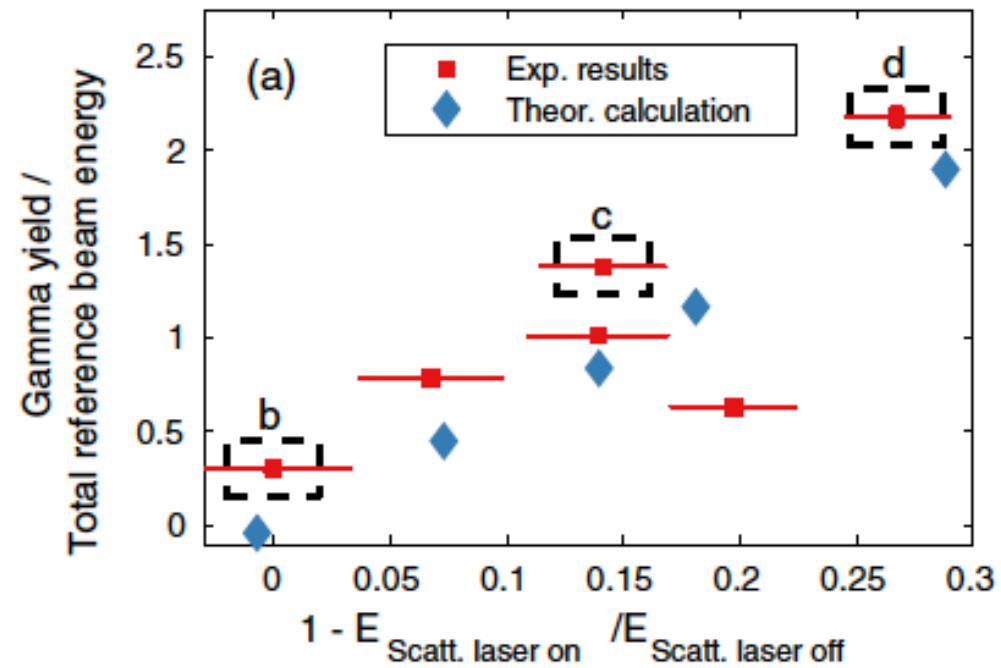
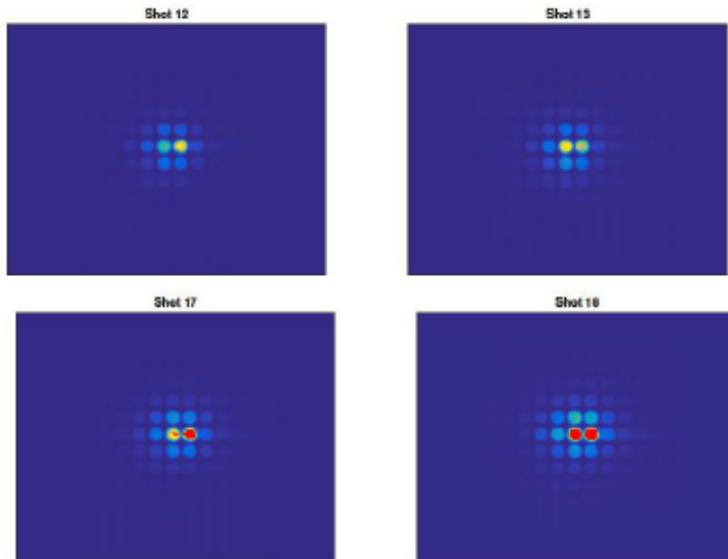
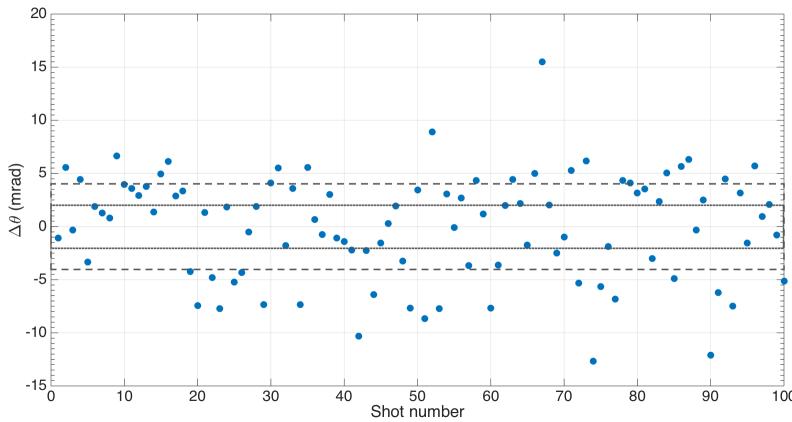
Experimental setup



Scattering laser intensity profile

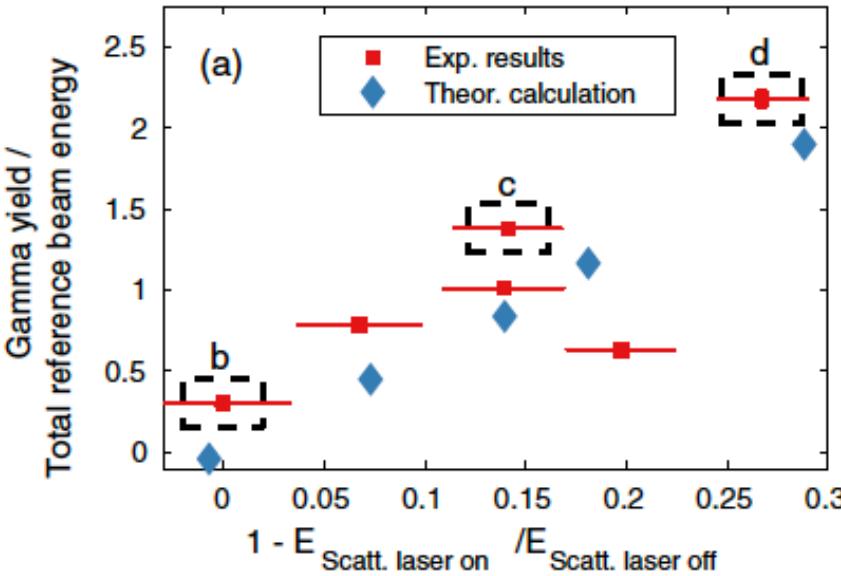


Collision diagnostic

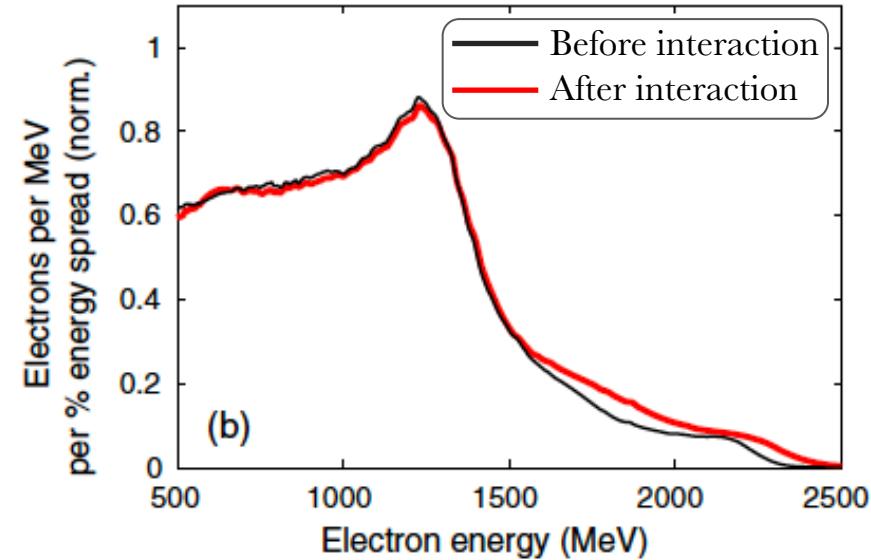
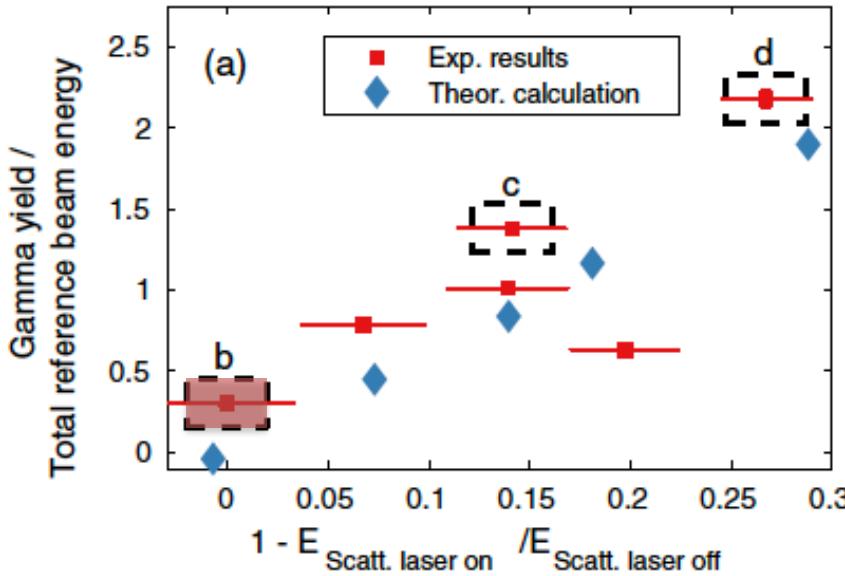


good electron-laser overlap → { high photon yield
high energy loss }

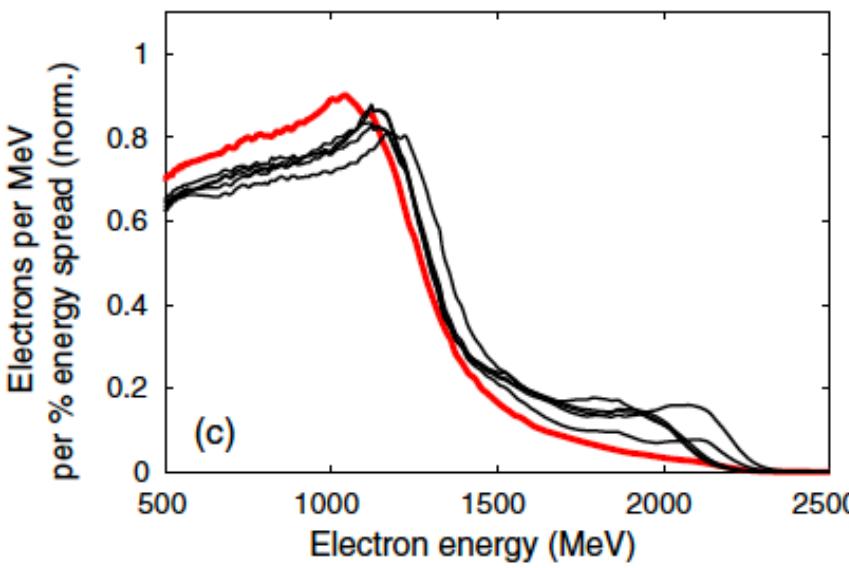
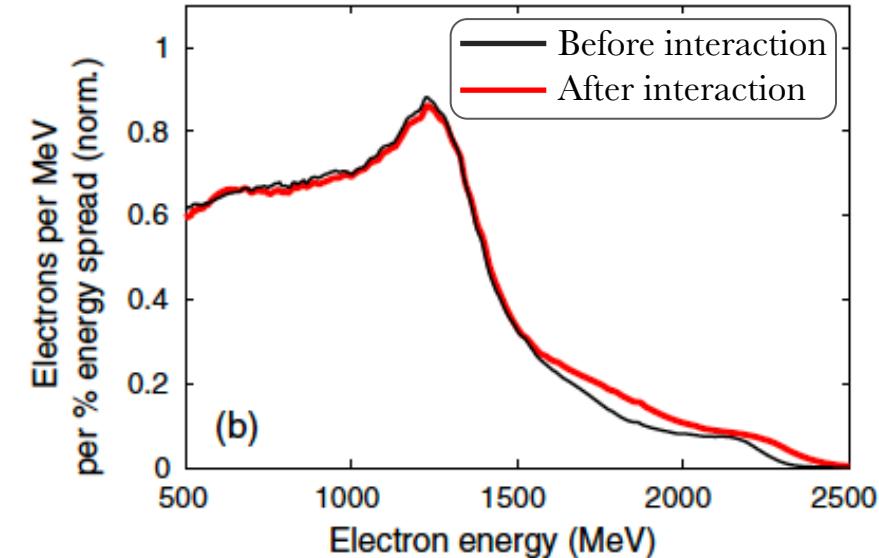
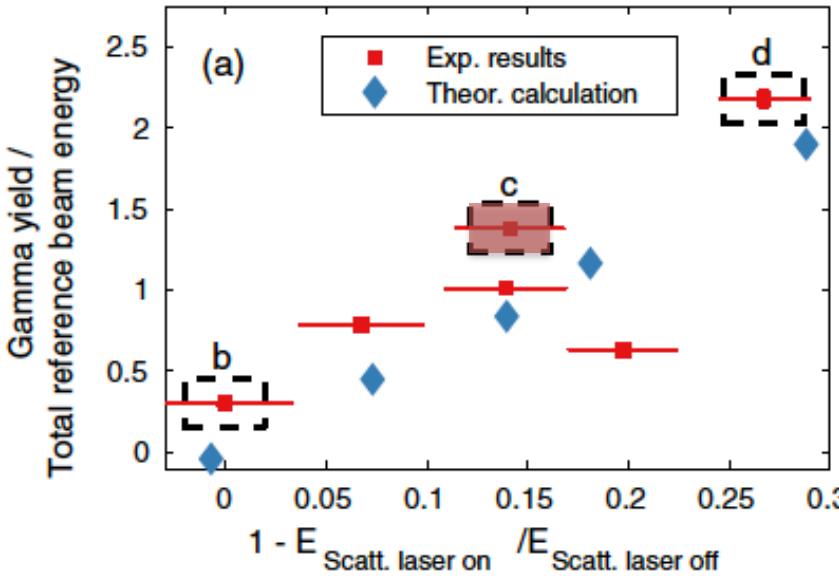
K. Poder *et al.*, Phys. Rev. X (2018)



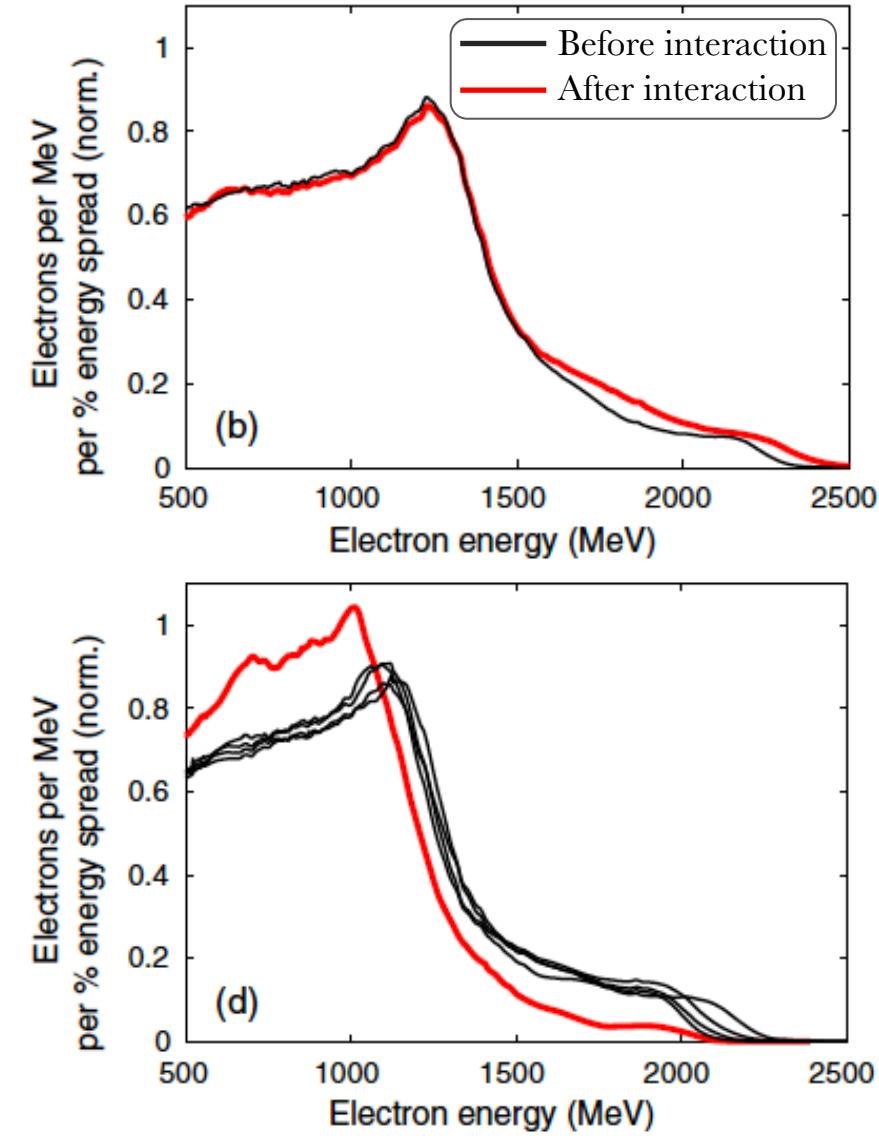
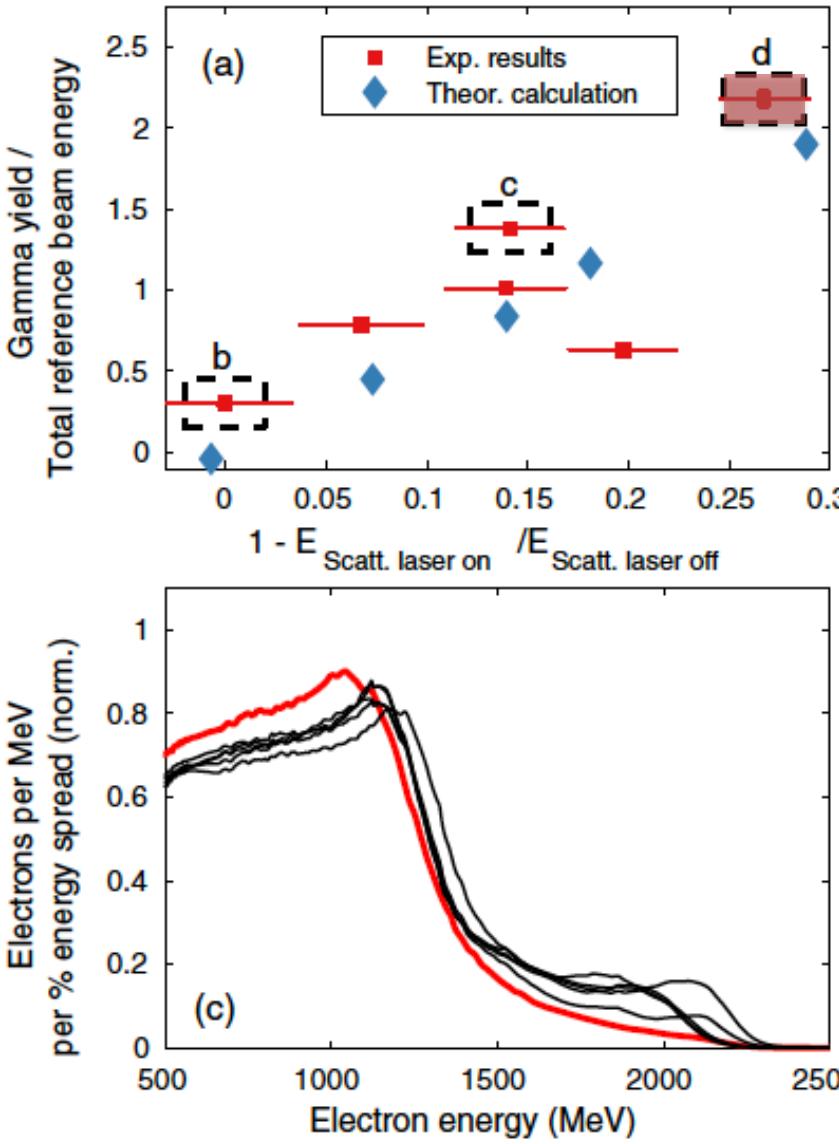
Collision diagnostic

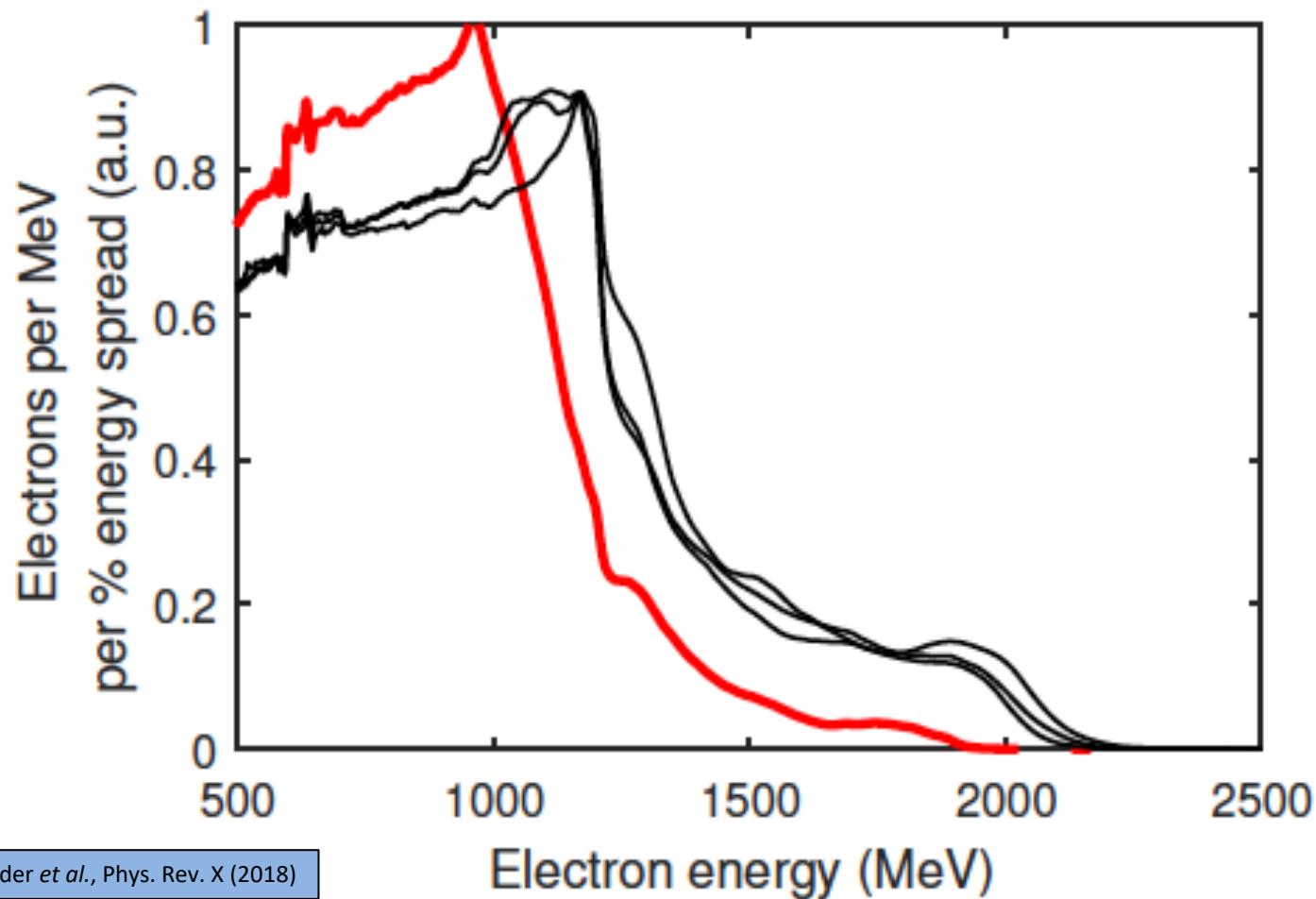


Collision diagnostic

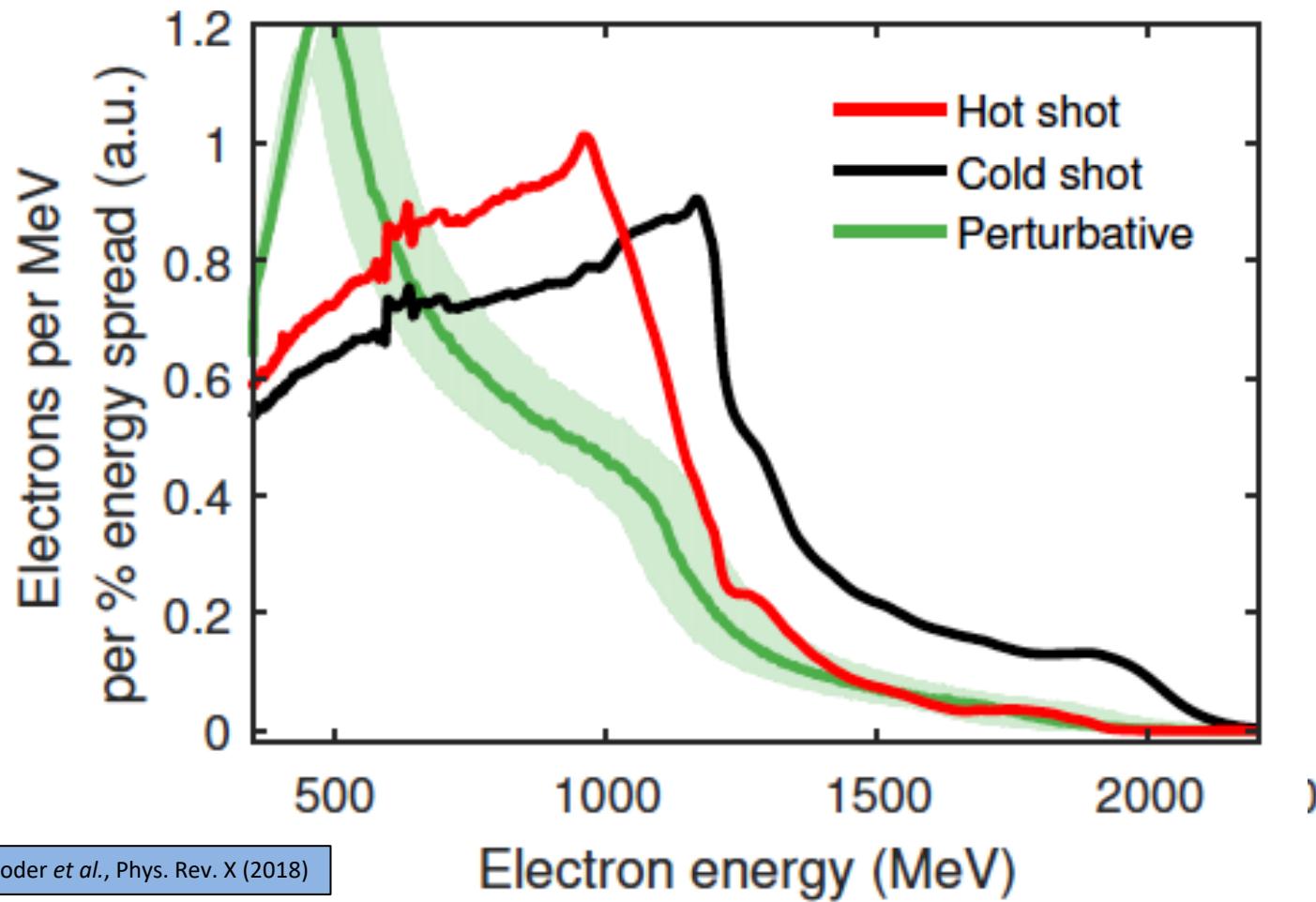


Collision diagnostic

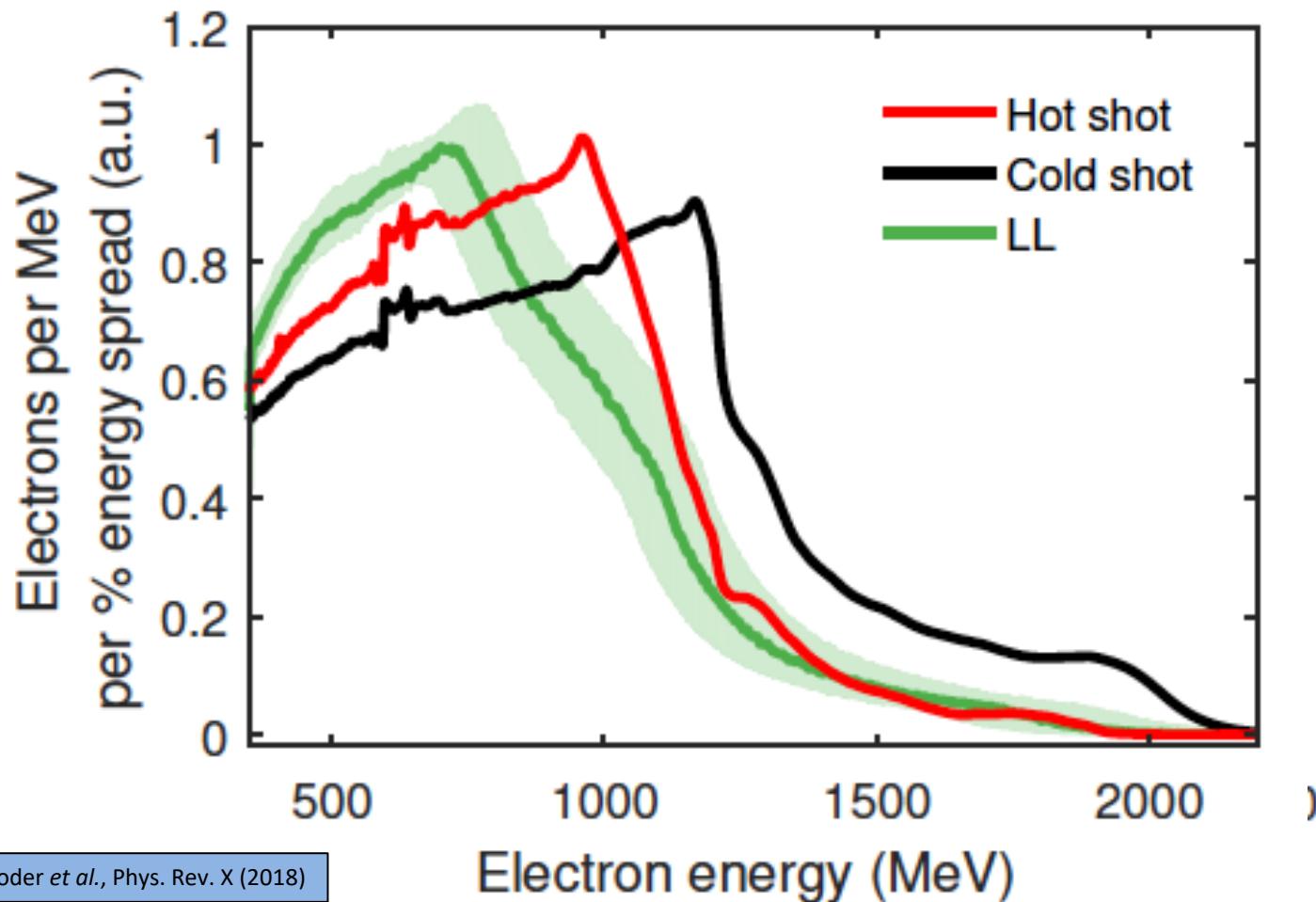




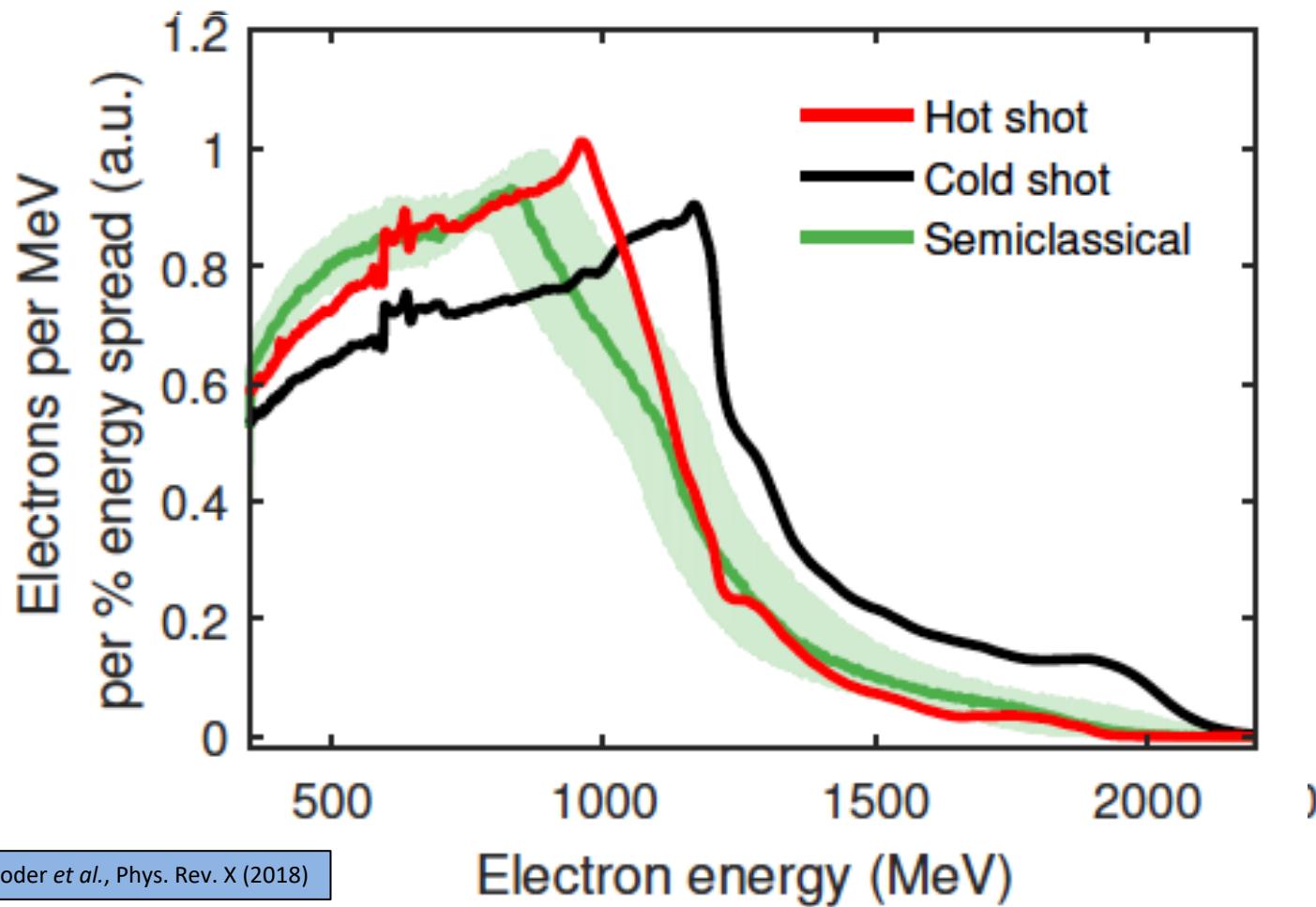
K. Poder *et al.*, Phys. Rev. X (2018)



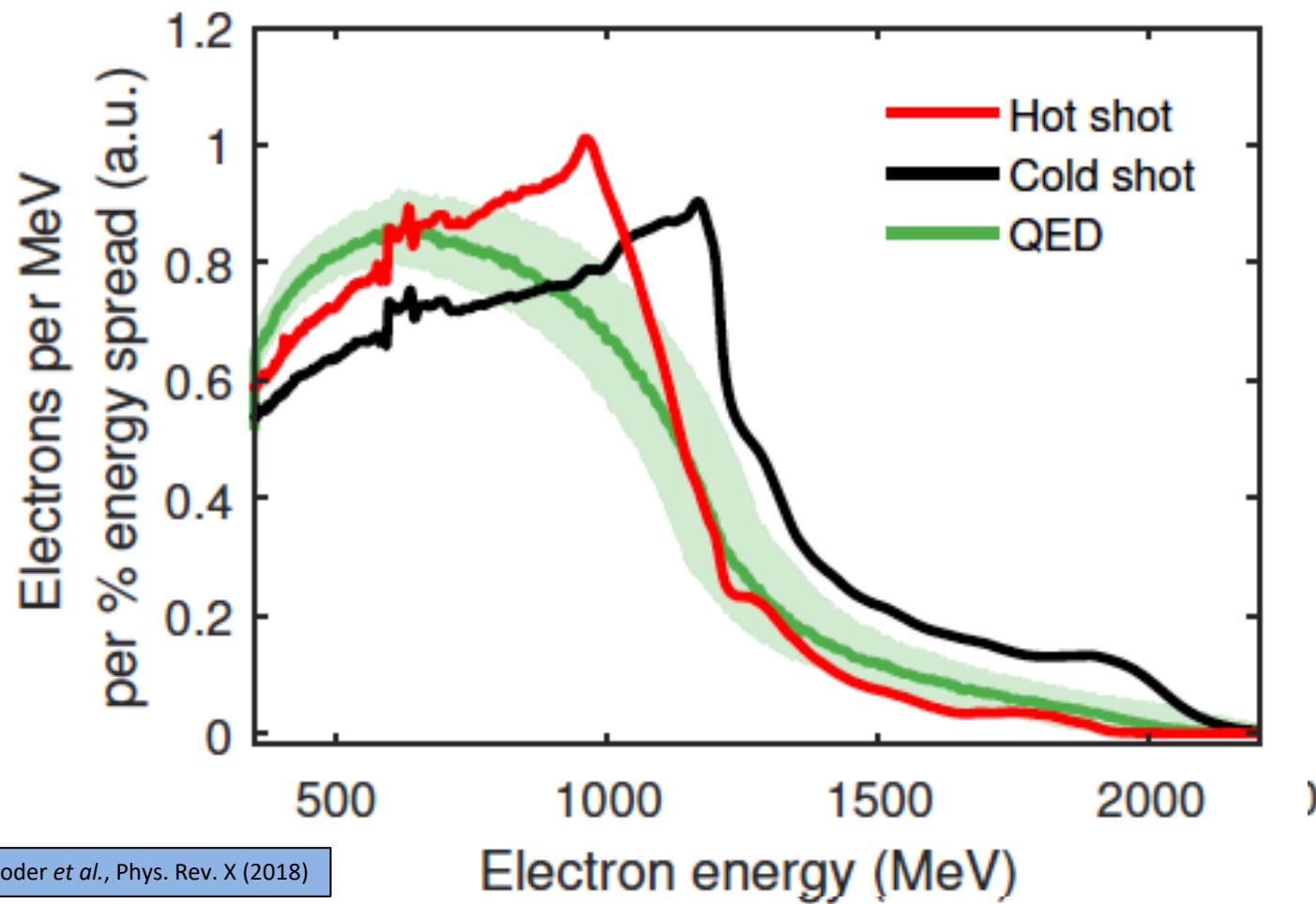
K. Poder *et al.*, Phys. Rev. X (2018)



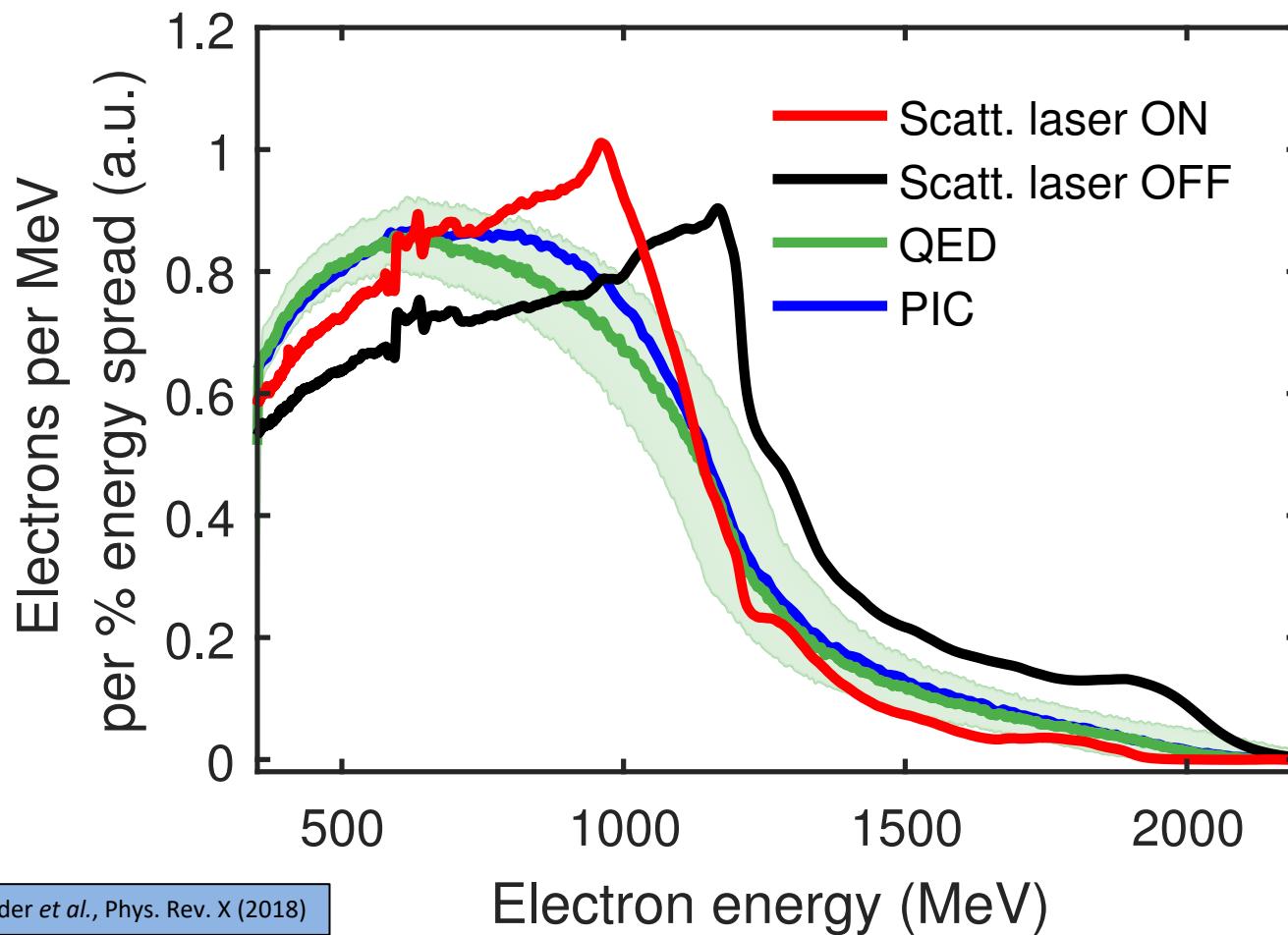
K. Poder *et al.*, Phys. Rev. X (2018)



K. Poder *et al.*, Phys. Rev. X (2018)



K. Poder *et al.*, Phys. Rev. X (2018)



K. Poder *et al.*, Phys. Rev. X (2018)

Why are the semiclassical and QED model not reproducing the data exactly?

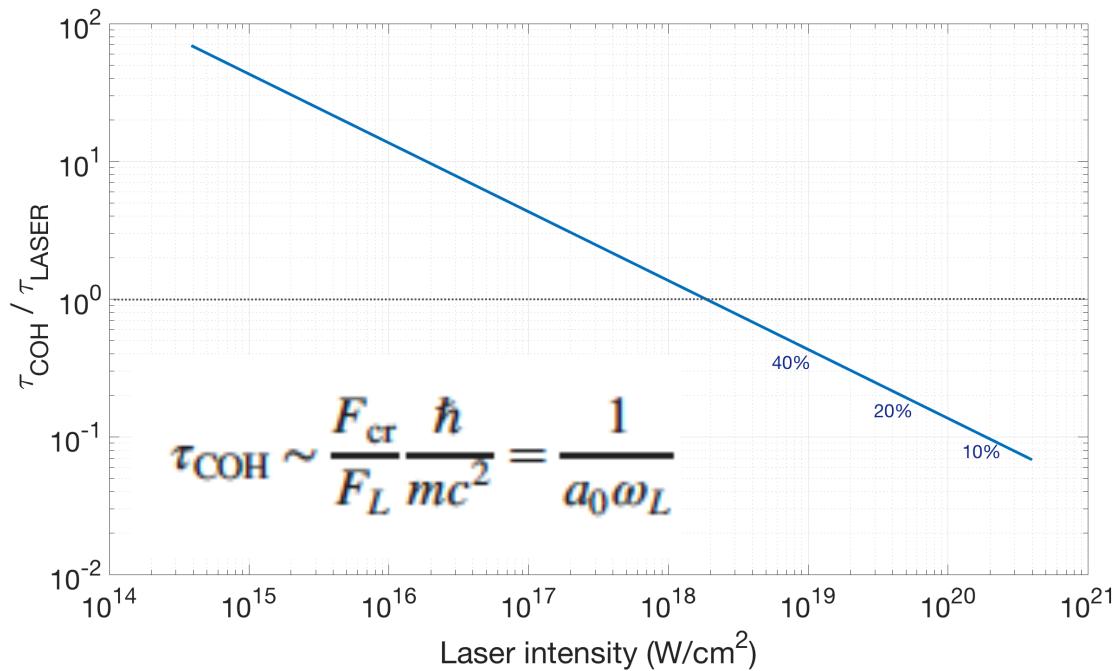
Several possibilities:

- Incomplete knowledge of laser spectral phase
- Incomplete knowledge of longitudinal laser distribution
- ...

Why are the semiclassical and QED model not reproducing the data exactly?

- Several possibilities:
- Incomplete knowledge of laser spectral phase
 - Incomplete knowledge of longitudinal laser distribution
 - ...

OR, we could be in a situation where the **constant cross-field approximation** is not strictly valid



This approximation is used to calculate

$g(\chi)$ in the semiclassical model

Photon emission probability in the QED model

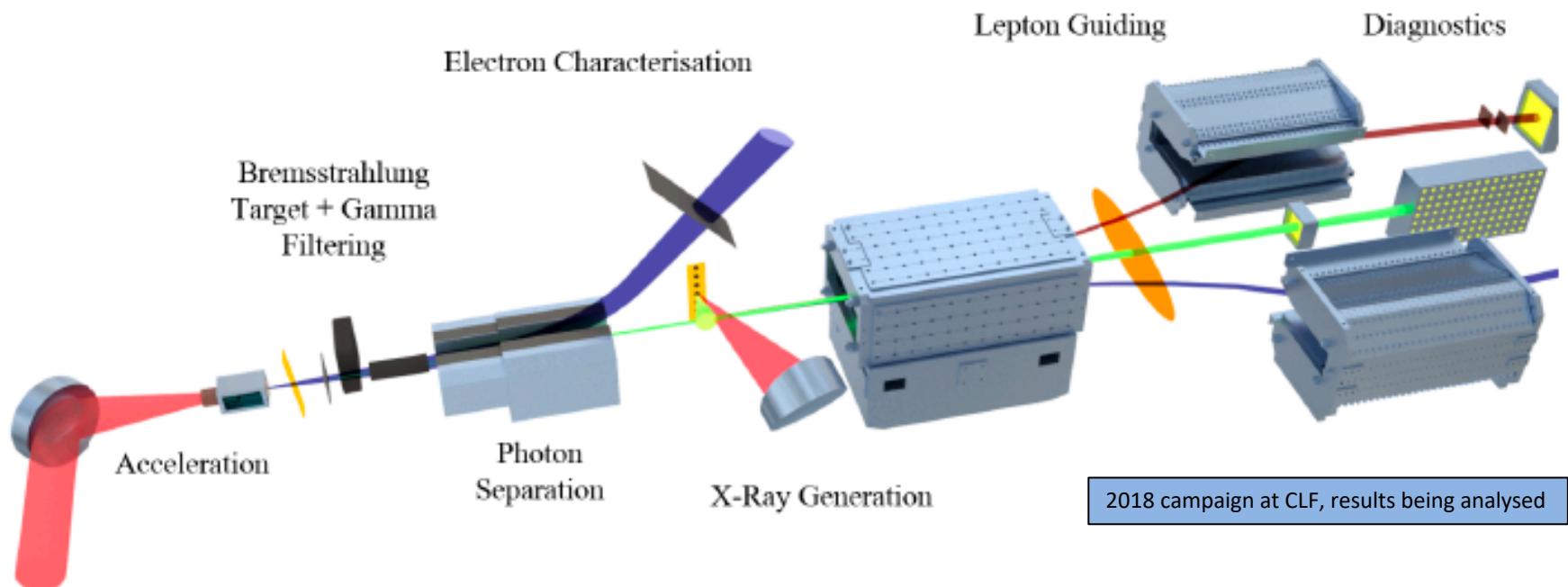
Part 3

Recent experiments on pair production

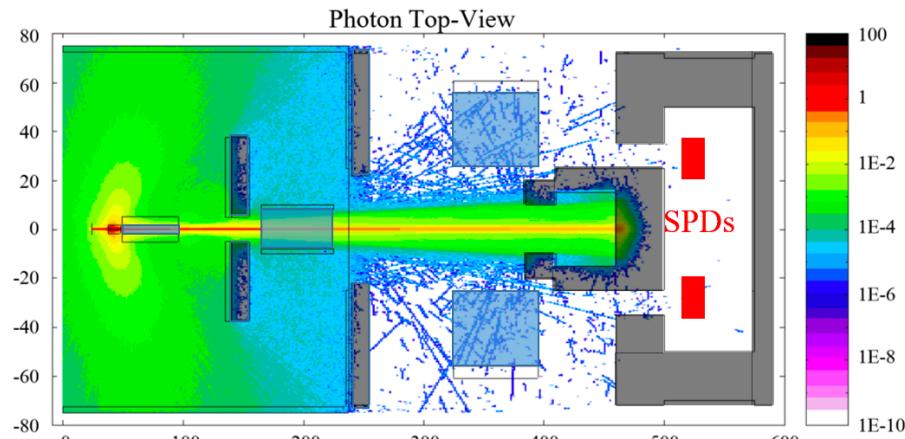
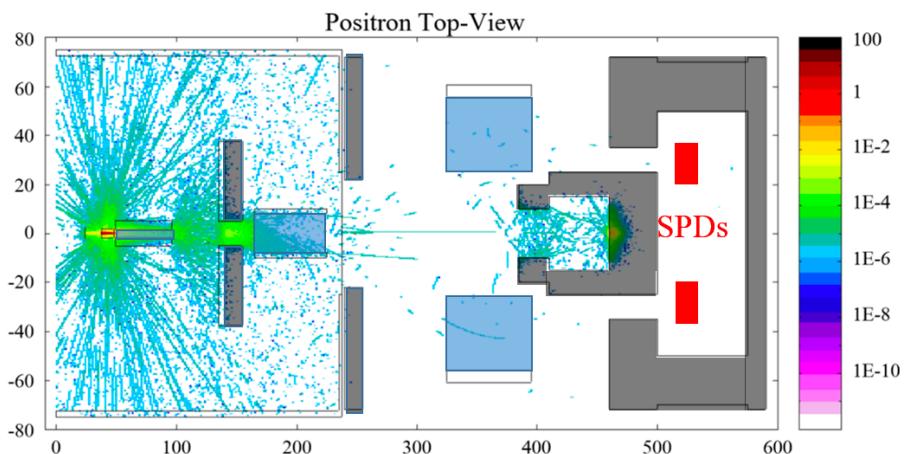
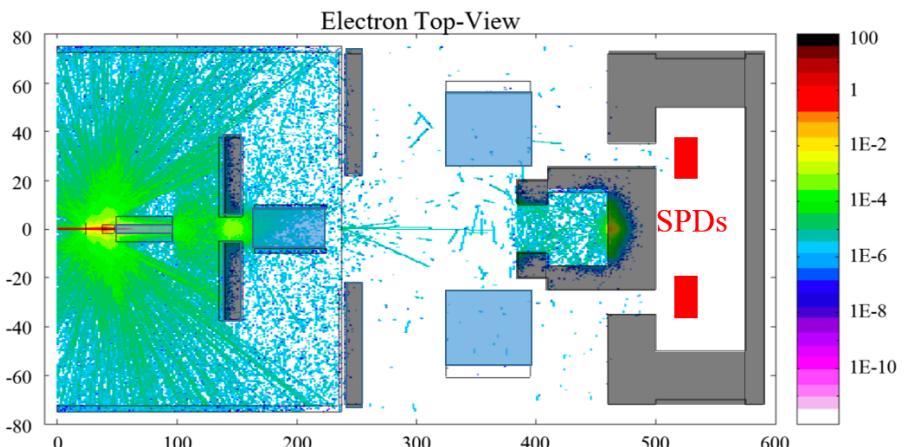
For an electron-positron pair to be produced during the collision of two photons of energy E_1 and E_2 , we need a centre of momentum energy of $\sqrt{E_1 E_2} > mc^2$: **Breit-Wheeler pair production**

If we start with $E_1 = 1.5$ eV (laser photon)
but, with $E_1 = 1.5$ keV (X-ray photon)

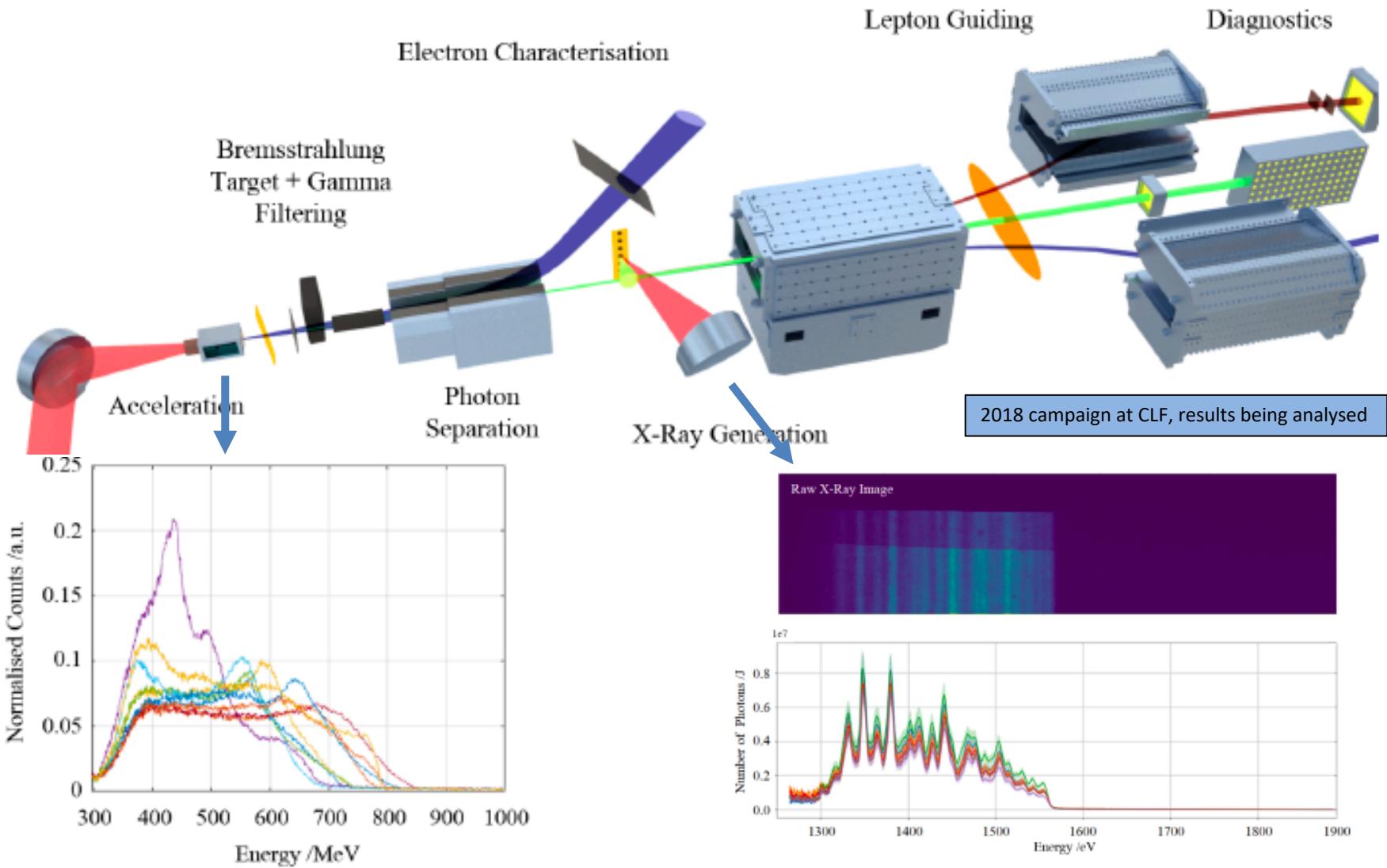
this means $E_2 = 174$ GeV (forget about it...)
this means $E_2 = 174$ MeV!



A major issue in these experiments is that you require incredibly high signal-to-noise.
 You want to measure a single particle (indeed, a fraction of a particle per shot) in an area flooded with secondary particles!



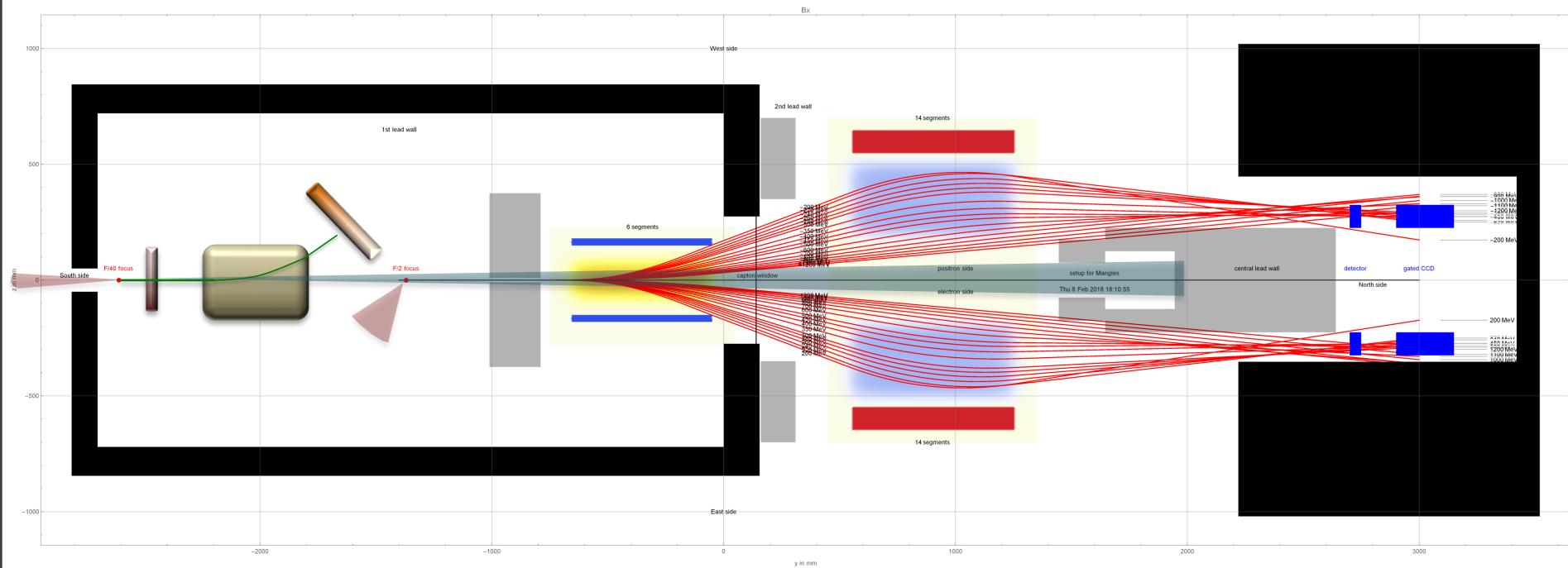
2018 campaign at CLF, results being analysed



However, what if I can get \mathcal{N} photons involved in one event? Then my threshold becomes

$$\sqrt{NE_1E_2} > mc^2 \text{ **Non-linear Breit-Wheeler pair production**}$$

If we start with $E_1 = 1.5$ eV (laser photon) and $E_2 = 1$ GeV (bremsstrahlung), then $N \sim 170$.



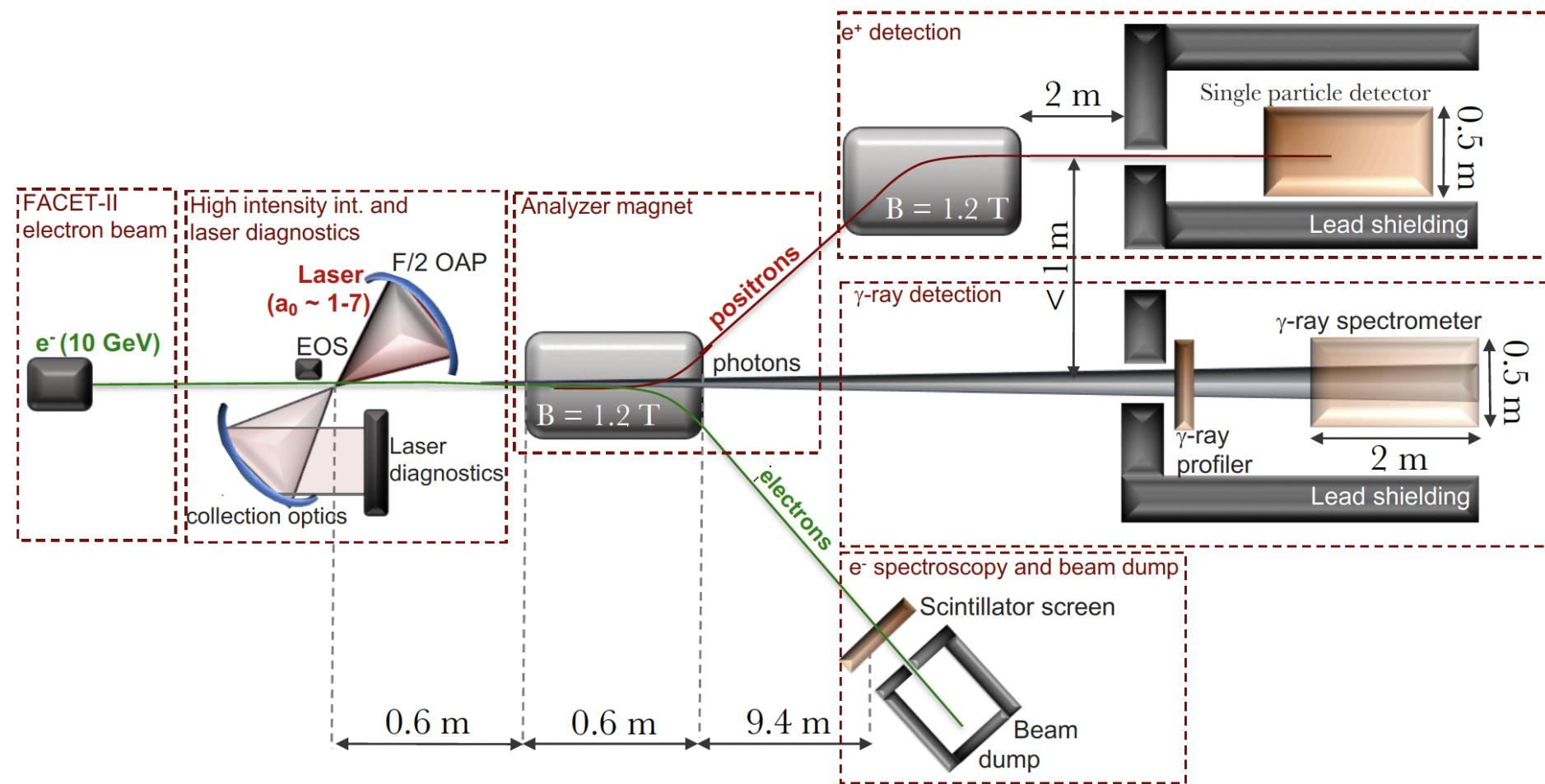
Part 4

Next experiments

E-320 at FACET-II

The E-320 experiment

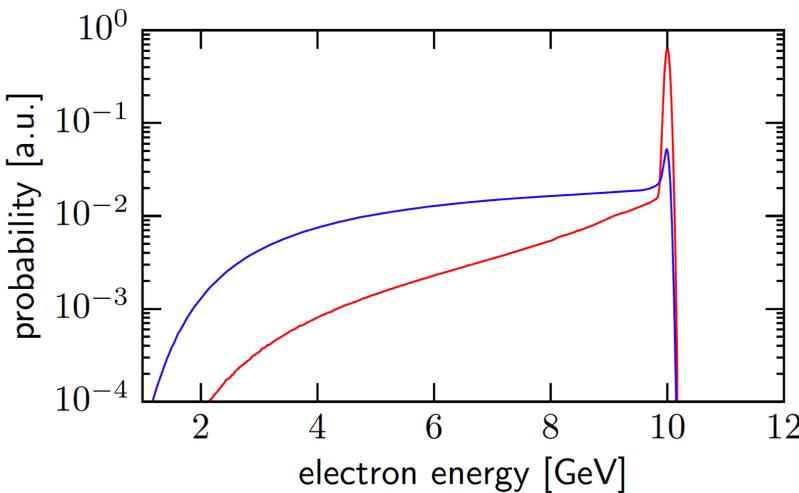
FACET-II provides 13 GeV electron beams, which are already coupled with a low-power laser beam (20 TW)



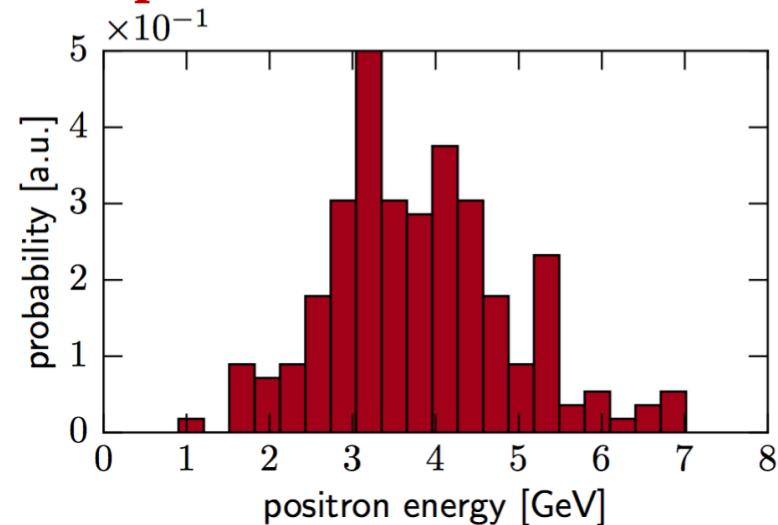
At FACET-II we aim (first shots in 2020) at measuring:

1. pair production in the laser field
2. quantum corrections to radiation reaction
3. non-linear Compton scattering (photon emission above the Compton edge)
4. breakdown of the LCFA

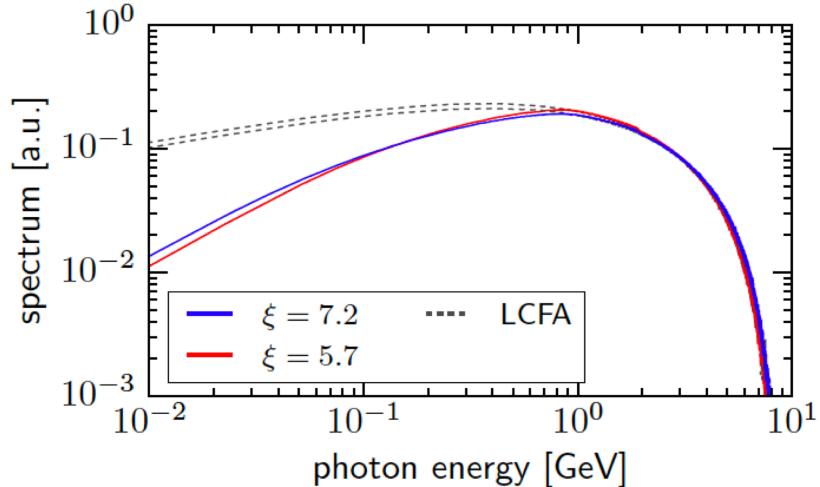
Quantum Radiation Reaction



Pair production



Non-linear Compton scattering



Part 5

Next experiments

LUXE at the EuXFEL

The EuXFEL generates high-quality electron beams with the following characteristics:

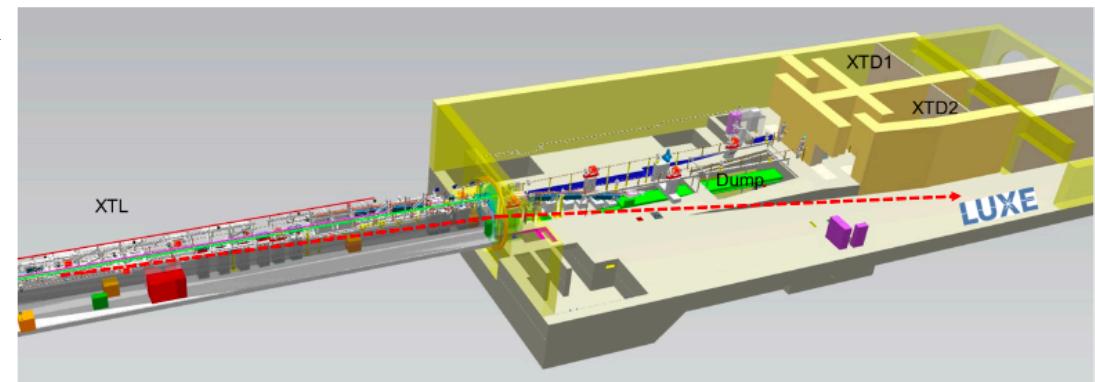
$$E = 17.5 \text{ GeV}$$

$$\Delta E/E = 2 \times 10^{-4}$$

$$N = 10^9$$

$$f = 10 \text{ Hz}$$

What if we couple a high-intensity laser to it?



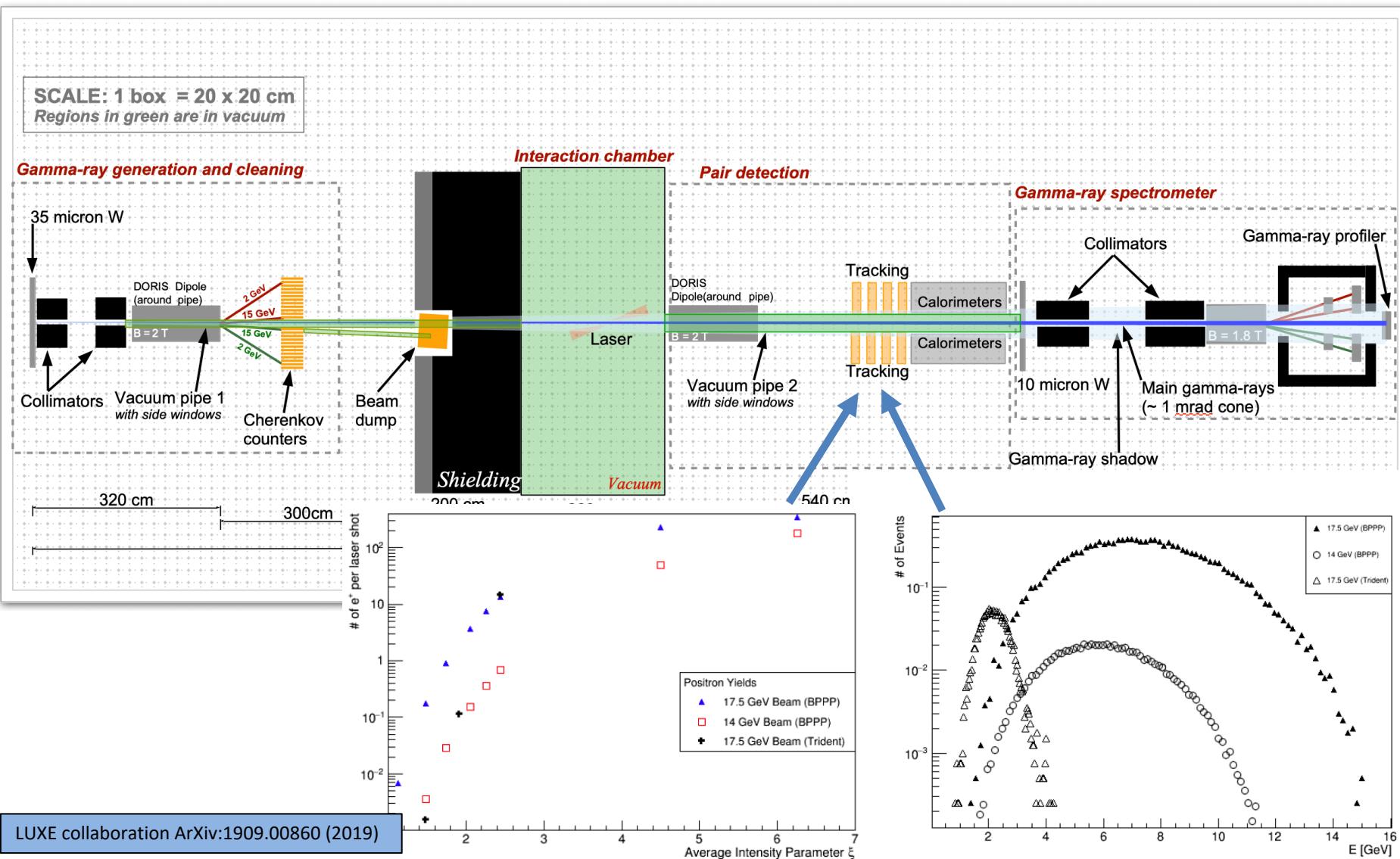
	30 TW, 8 μm	300 TW, 8 μm	300 TW, 3 μm
Laser energy after compression (J)	0.9	9	9
Percentage of laser in focus (%)	40	40	40
Laser energy in focus (J)	0.36	3.6	3.6
Laser pulse duration (fs)	30	30	30
Laser focal spot FWHM (μm)	8	8	3
Peak intensity in focus (Wcm⁻²)	1.6×10^{19}	1.6×10^{20}	1.1×10^{21}
Dimensionless peak intensity, ξ	2	6.2	16
Laser repetition rate (Hz)	1	1	1
Electron-laser crossing angle (rad)	0.35	0.35	0.35

17.5 GeV electrons

Electron Lorentz factor	3.4×10^4	3.4×10^4	3.4×10^4
Quantum parameter χ	0.41	1.26	3.26

LUXE collaboration ArXiv:1909.00860 (2019)

The LUXE experiment



LUXE collaboration ArXiv:1909.00860 (2019)

What is unique about the LUXE experiment?

The only experiments in the area are the E-144 at SLAC and the Gemini experiments

E – 144 and Gemini

- ✗ $\chi \sim 0.2$
- ✗ the E-144 still operated in a quasi-linear regime ($a_0 \sim 0.3$)
perturbative non-linearities in Compton scattering
- ✗ ***lack of parametric studies in intensity***
- ✗ ***no direct photon-photon studies***
- ✗ ***difficult to secure sustained access***

LUXE

- ✓ $\chi \sim 1$ even for a 30 TW laser
- ✓ high a_0 implies ***strong non-linearities*** in Compton scattering and pair production
- ✓ systematic and precision ***parametric studies***
- ✓ ***pure non-linear Breit-Wheeler*** above and below threshold
- ✓ easily upgradable → sustained campaigns
- ✓ exotic physics at the ***intensity frontier!***

Equipment cost for phase 1 ~ 30 M€
input from European and national funding agencies

LUXE collaboration ArXiv:1909.00860 (2019)

Part 6

Next experiments

E6 area at ELI-NP



Assuming 60% of it
in a $3\mu\text{m}$ FWHM focal spot, we get

$$I = 6 \times 10^{22} \text{ Wcm}^{-2} \quad (a_0 > 100)$$

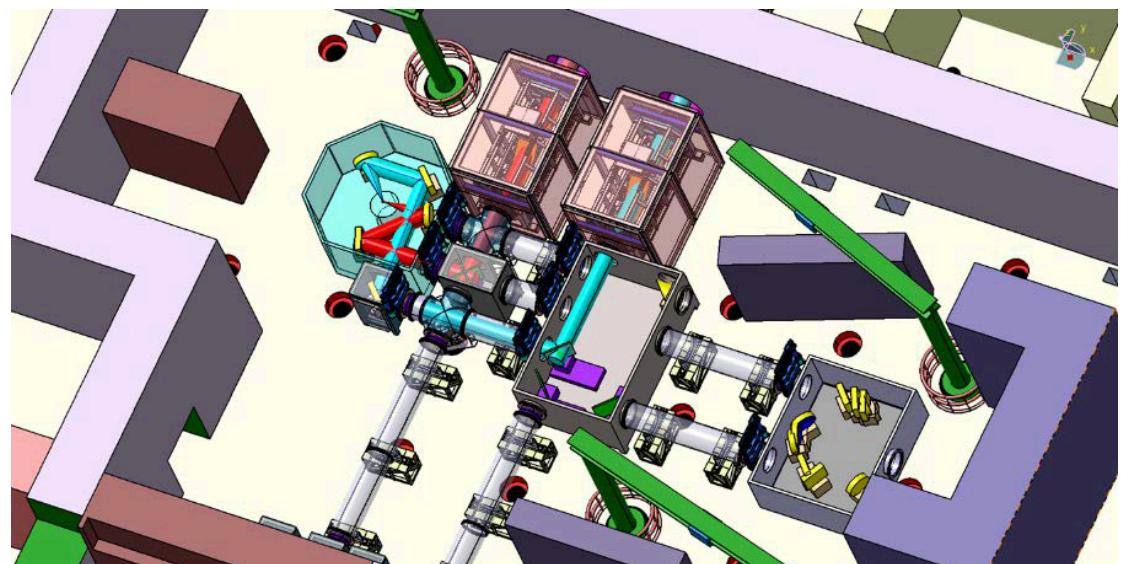
If we couple it with a
5 GeV electron beam, $\chi \sim 7!$

**First commissioning
experiments
end 2020**

LaserFocusWorld®

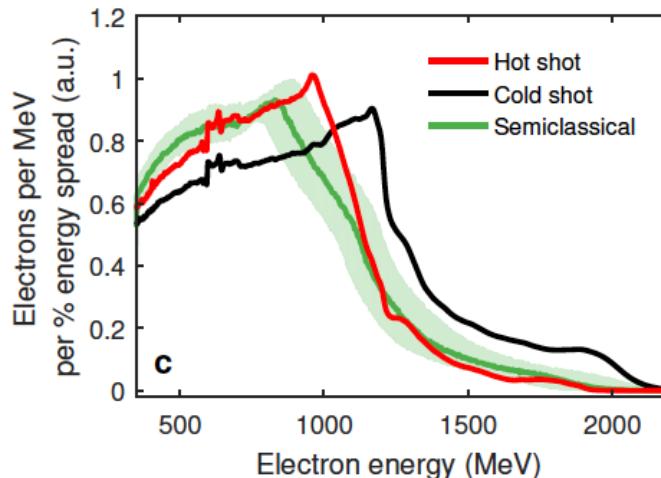
World's most powerful laser, developed by Thales and ELI-NP,
achieves record power level of 10 PW

After delivering pulses of 7 PW for more than 4 h continuously, the Thales system reached 10 PW on 7 March 2019.

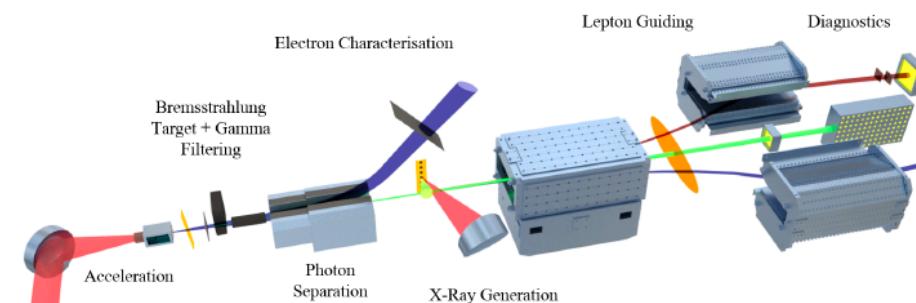


Conclusions

First experiments in CLF showed hints of the quantum nature of radiation reaction ($\chi \sim 0.2$)

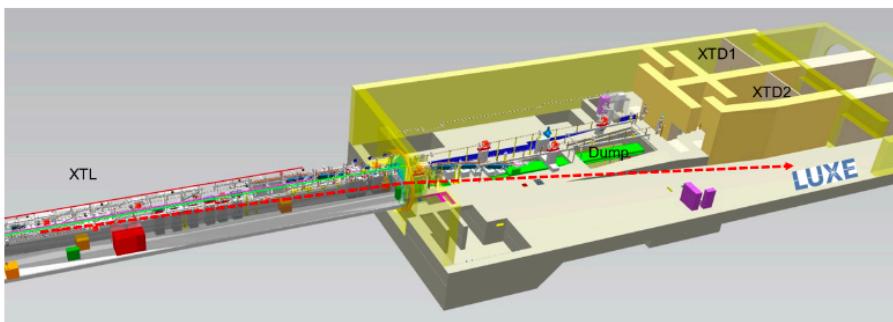


First attempts at studying pair production from photon-photon collisions (**Breit Wheeler**)



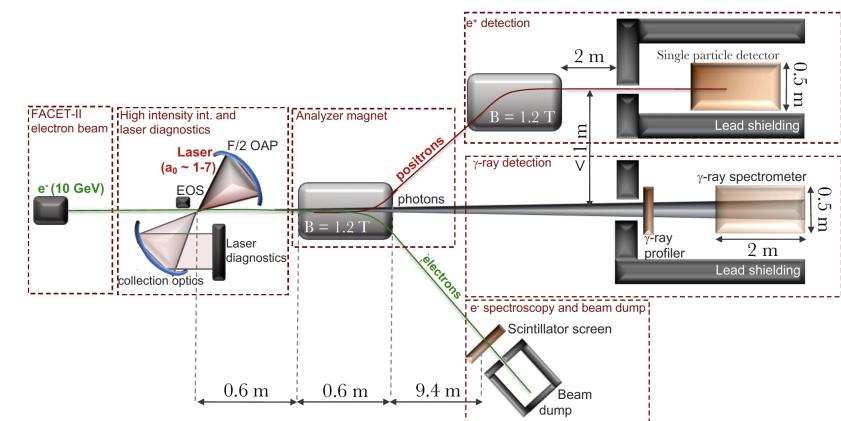
LUXE experiment at EuXFEL

(17.5 GeV electrons and 30 - 300 TW laser)



E-320 experiment at FACET-II

(13 GeV electrons and 20 TW laser)



Thanks for your attention!

Gianluca Sarri
g.sarri@qub.ac.uk

Main publications

- [1] D. J. Corvan et al., Rev. Sci. Instrum. 85, 065119 (2014)
- [2] G. Sarri et al., Phys. Rev. Lett. 113, 224801 (2014)
- [3] D. J. Corvan et al., Optics Express 24, 3127 (2016)
- [4] K. T. Behm et al., Rev. Sci. Instrum 89, 113303 (2018)
- [5] K. Poder et al., Phys. Rev. X 8, 031004 (2018)
- [6] J. Cole et al., Phys. Rev. X 8, 011020 (2018)
- [7] LUXE collaboration ArXiv:1909.00860 (2019)