In search for Matter Creation[©] The landscape of double beta decay

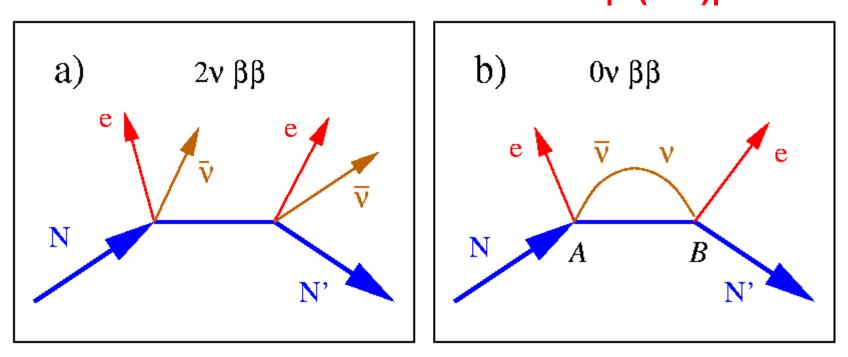
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Chiara Brofferio
Mark Chen
Steve Elliott
Kunio Inoue
Yuri Kolomenski
Suzanne Mertens
John Wilkerson

There are two varieties of $\beta\beta$ decay

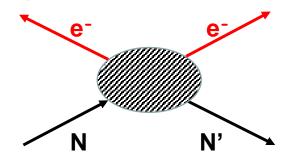
2v mode: a conventional 2nd order process in nuclear physics 0v mode: a hypothetical process can happen only if: $M_v \neq 0$ $\overline{v} = v$ $|\Delta L|=2$ $|\Delta (B-L)|=2$



"Black box" theorem*: "0vββ decay always implies new physics"

There is no scenario in which observing 0vββ decay would not be a great discovery

- ➔ Majorana neutrinos
- ➔ Lepton number violation
- → Probe new mass mechanism up to the GUT scale
- Probe key ingredient in generating cosmic baryon asymmetry



Neutrino masses have to be non-zero for $0v\beta\beta$ to be possible.

Because the distinction between Dirac and Majorana particles is only observable for particles of non-zero mass.

Strictly speaking, this is the ONLY connection with neutrino masses relevant to discover new physics.

Hence it is appropriate to think of the sensitivity to new physics as scaling with $T_{1/2}$, irrespective of the neutrino mass scenarios. A $T_{1/2}$ sensitivity increase from ~10²⁶ to ~10²⁸ yr (~100x), should be compared, e.g., to the \sqrt{s} increase from Tevatron to LHC (~20), although, admittedly, with a smaller array of channels for new physics.

* J. Schechter, and J. W. F. Valle, Phys. Rev. D25, 2951 (1982).

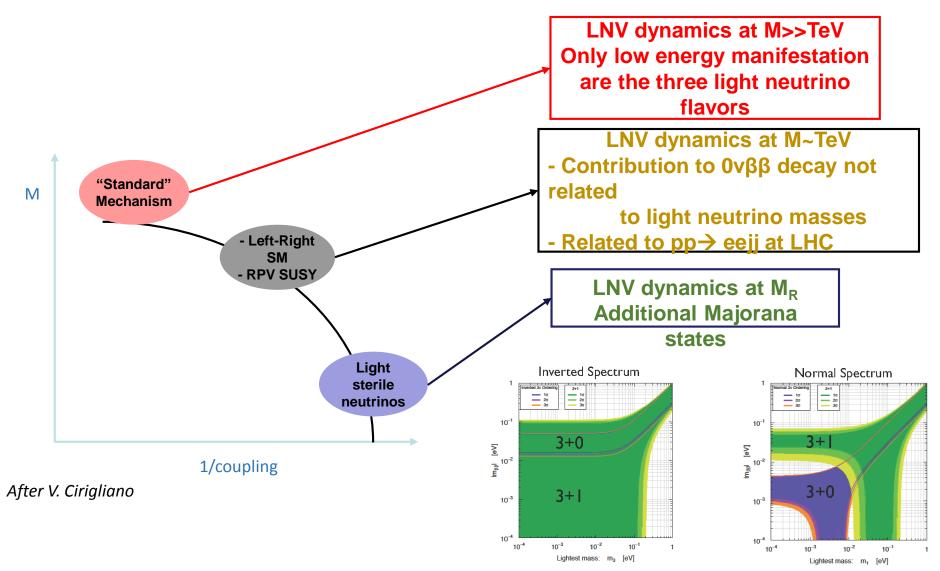
The connection with the v mass also means that the observation of 0vββ decay can provide information on the v mass scale, provided that:

- The mechanism producing the decay is understood
- The nuclear matrix element is calculated with sufficiently small uncertainty
- The appropriate value of g_A to be used is clarified

This is of course an important bonus, but these uncertainties do not affect the discovery potential of tonne-scale experiments.

It is also a convenient, although imperfect, metric to compare isotopes and experiments.

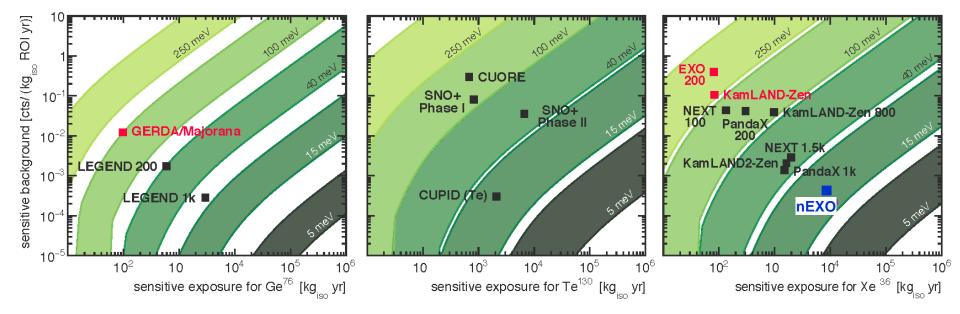
Different physics can contribute to the decay



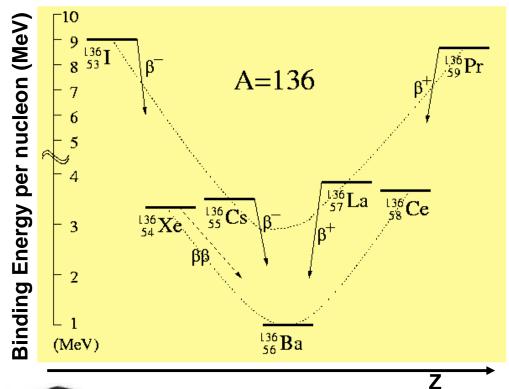
Giunti-Zavanin 1505.00978, JHEP07(2015)171

The discovery potential was recently estimated for various proposals, assuming Type I seesaw, the free value of g_A and using a Bayesian analysis with flatly distributed priors

(Agostini, Benato, Detwiler, PRD 96 (2017) 053001 also A. Caldwell et al., PRD 96 (2017) 073001)



The observable is the half-life of a nuclear state for which the regular β decay is forbidden



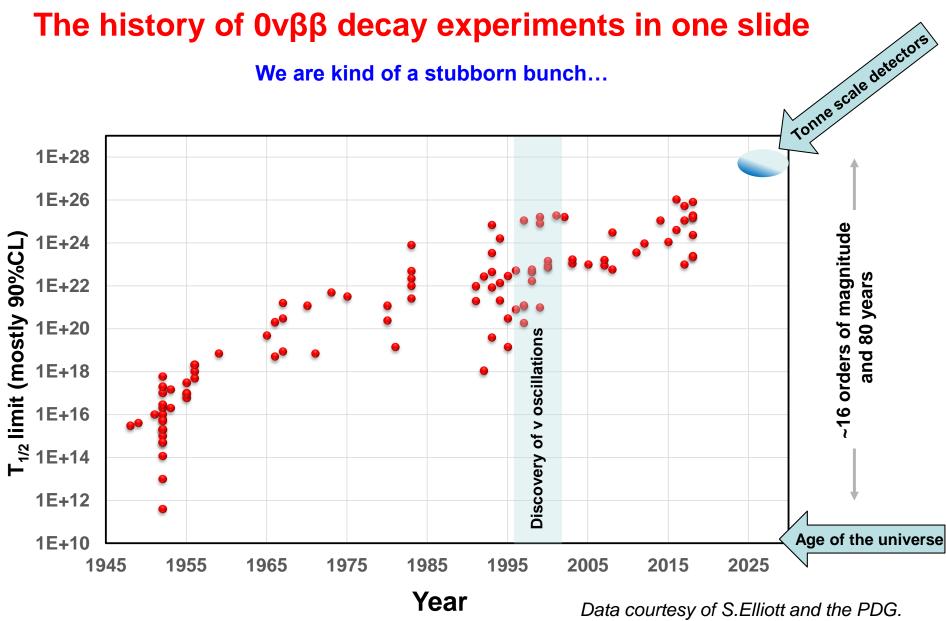
0
900
E
NO
Amedeo Avogadro

1776-1856

...and we owe the remarkable T_{1/2} sensitivity to the magnitude of Avogadro's number!

Examples with Q>2MeV						
Candidate	Q (MeV)	Abundance				
⁴⁸ Ca→ ⁴⁸ Ti	4.271	0.187				
⁷⁶ Ge→ ⁷⁶ Se	2.040	7.8				
⁸² Se→ ⁸² Kr	2.995	9.2				
⁹⁶ Zr→ ⁹⁶ Mo	3.350	2.8				
¹⁰⁰ Mo→ ¹⁰⁰ Ru	3.034	9.6				
¹¹⁰ Pd→ ¹¹⁰ Cd	2.013	11.8				
¹¹⁶ Cd→ ¹¹⁶ Sn	2.802	7.5				
¹²⁴ Sn→ ¹²⁴ Te	2.228	5.64				
¹³⁰ Te→ ¹³⁰ Xe	2.533	34.5				
¹³⁶ Xe→ ¹³⁶ Ba	2.458	8.9				
$^{150}Nd \rightarrow ^{150}Sm$	3.367	5.6				

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Not all results are necessarily shown.

A historical note

Early experiments:

- Geochemical or Radiochemical experiments

- (search for trace amounts of element A in a large amount of B after a long time).
- Can't discriminate between 0v and 2v decays But may see a renaissance in combination with real-time counting
- Counting experiments with gram quantities of candidate isotopes.

"Previous generation" experiments:

- Counting experiments with kg quantities of enriched material

Running and future experiments:

>100kg-class counting experiments (mainly enriched)

Four fundamental requirements for modern experiments:

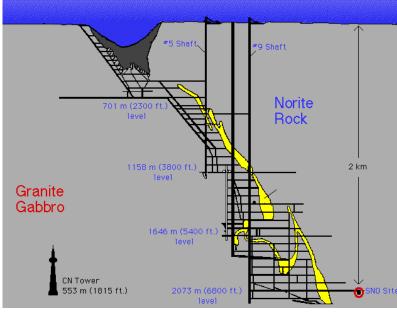
1) Isotopic enrichment of the source material (that is generally also the detector)

100kg – class experiment running or completed. Ton – class experiments under planning.

2) Underground location to shield cosmic-ray induced background

> Several underground labs around the world, Next round of experiments 1-2 km deep.





Four fundamental requirements for modern experiments:

3) Ultra-low radioactive contamination for detector construction components

Materials used $\approx <10^{-15}$ in U, Th (U, Th in the earth crust \sim ppm)

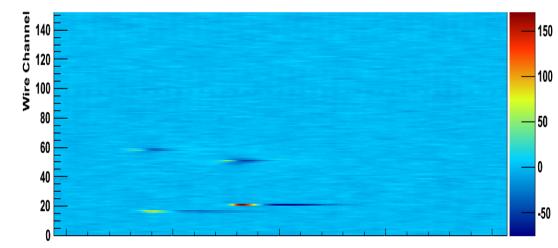


4) New techniques to discriminate signal from background

Non trivial for E~1MeV

Modern detectors have a number of handles.

This gets easier in larger detectors.



Time

The last point deserves more discussion, particularly as the size of detectors grows...

Modern detectors use a combination of four parameters/measurements:

- 1. Energy measurement (for small detectors this is ~all there is).
- 2. Event multiplicity (γ's Compton scatter depositing energy in more than one site in large detectors).
- 3. Depth in the detector (or distance from the walls) is (for large monolithic detectors) a powerful parameter for discriminating between signal and (external) backgrounds.
- 4. αs can be distinguished from electrons in many detectors.

Powerful detectors use most (possibly all) these parameters in combination, providing the best possible background rejection and simultaneously fitting for signal and background.

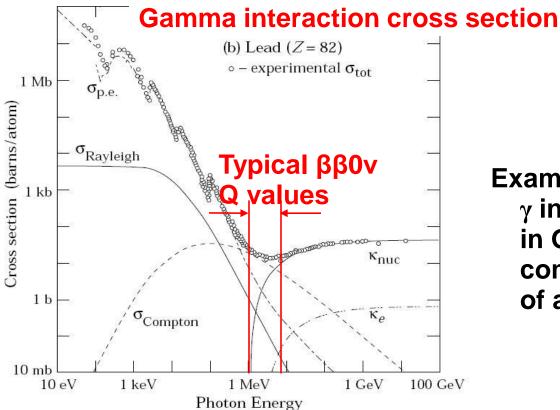
The optimal combination of parameters does depend on the size of the detector and on new techniques constantly being developed.

0vββ searches are quite different from Dark Matter ones

0vββ decay searches	Dark Matter searches
Optimize for energy resolution at O(2MeV)	Optimize for <100keV threshold
Gamma ray shielding essential	Neutron shielding essential
Signal is electron-like	Signal is nuclear recoil-like
Isotopic enrichment	2vββ is a background
Self shielding starts being useful at 100kg scale	Infinite self shielding, external backgrounds generally not important

But there is synergy in the detector development for some of the technologies

Shielding a detector from \sim MeV γ s is difficult!

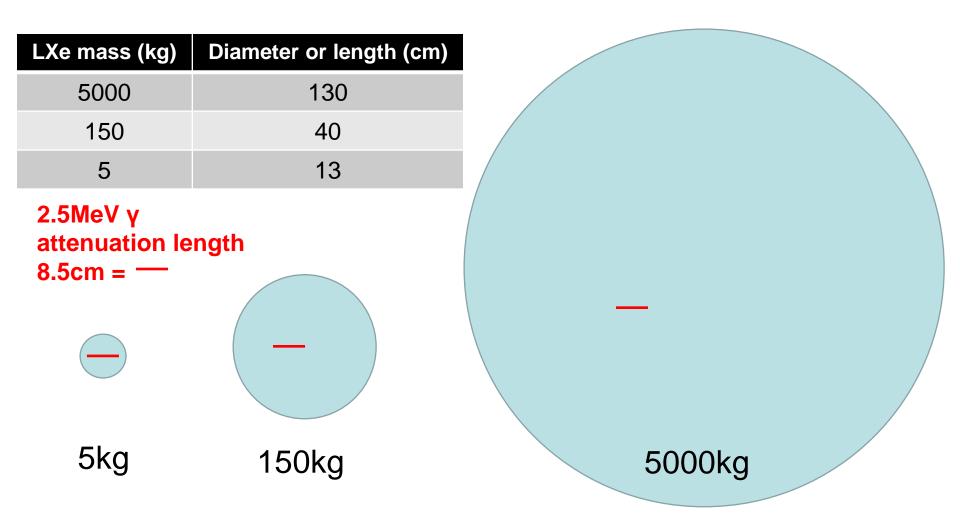


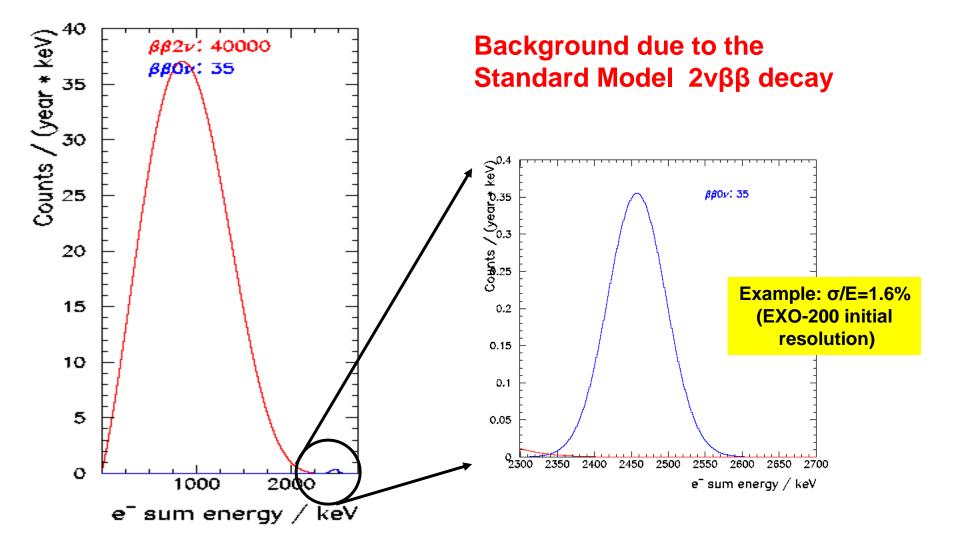
Example:

 γ interaction length in Ge is 4.6 cm, comparable to the size of a germanium detector.

Shielding *ββ* decay detectors is much harder than shielding Dark Matter ones We are entering the "golden era" of $\beta\beta$ decay experiments as detector sizes exceed int lengths

Moving forward, monolithic is very attractive





Subtracting the tail under the $0\nu\beta\beta$ peak is tricky and, irrespective of other background considerations, sufficient energy resolution is required.

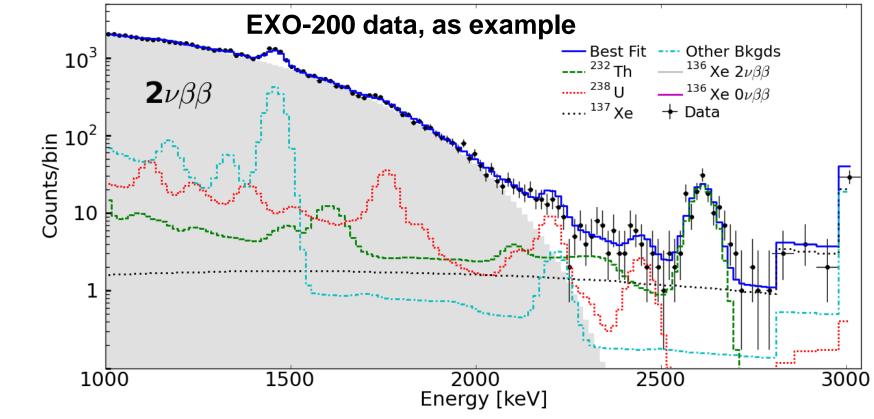
Of course goal of 0vββ experiments is to make a discovery

--setting a limit is the fall back position!

So, making sure that experiments have sensitivity is also very important

One can think of the 2vββ in a more positive way, as Nature's "blind injection"

-these events look like $0v\beta\beta$ events, extending in energy just below the Q-value.



Double-beta decay experiments -- Gratta

A healthy neutrinoless double-beta decay program requires more than one isotope.

This is because:

- There could be unknown γ transitions and a line observed at the "end point" in one isotope does not necessarily imply the 0vββ decay discovery
- Nuclear matrix elements are not very well known and any given isotope could come with unknown liabilities
- Different isotopes correspond to vastly different experimental techniques
- 2 neutrino background is different for various isotopes
- The elucidation of the mechanism producing the decay requires the analysis of more than one isotope

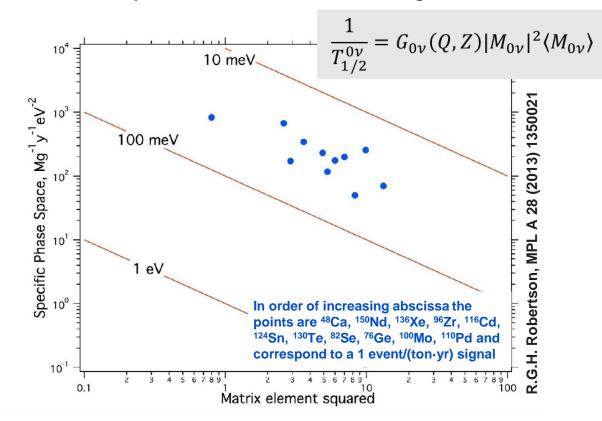
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Can we just concentrate on the best isotope? Many isotopes have comparable sensitivities

(at least in terms of rate per unit neutrino mass using the standard SeeSaw)



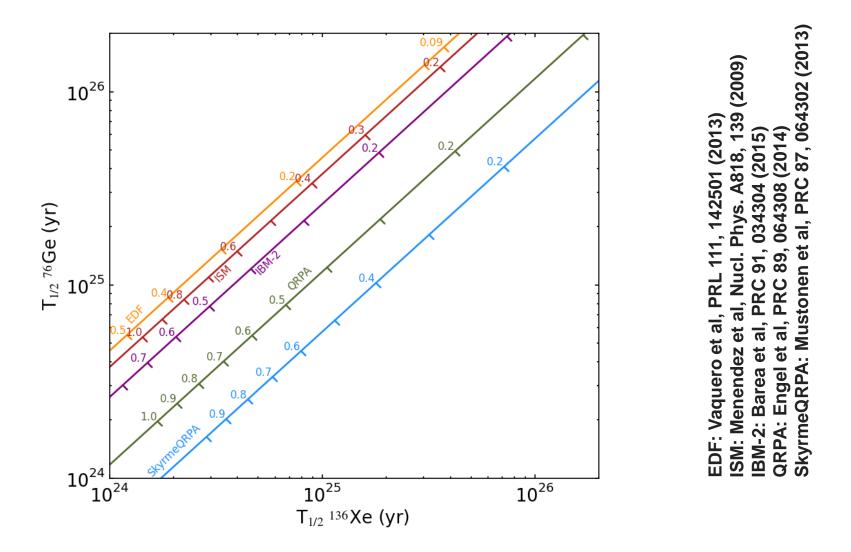
There is an "empirical" anticorrelation between phasespace and NME.

The choice between isotopes is more a choice between different techniques and the performance they have achieved *in actually running* detectors.

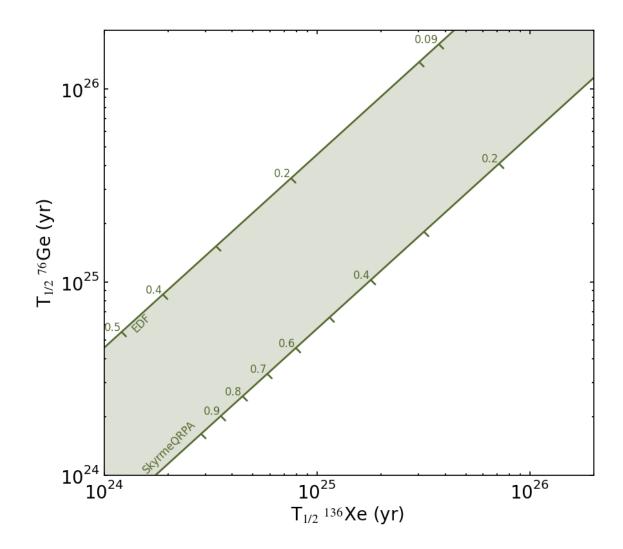
This is essential for the tonne scale, where we are talking about ~100M\$ investments.

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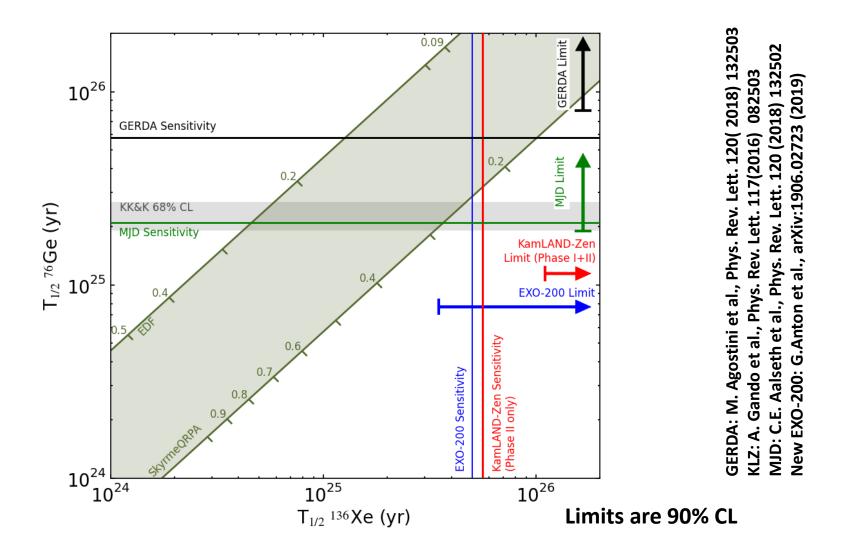
Because of the uncertainties in the 0vββ decay mechanism and the NME, accurate comparisons between different isotopes are non-trivial. Example using ¹³⁶Xe and ⁷⁶Ge (and assuming standard SeeSaw)



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One (my own) possible classification of technologies

Crystals

(keeping in mind that real things are always more complex than classifications!)

- GERDA, Majorana (⁷⁶Ge) - CUORE, CUPID (¹³⁰Te, ¹⁰⁰Mo)

Pros: Superb energy resolution,

2-parameter measurement

Cons: Intrinsically fragmented

Liquid (organic) scintillators

- KamLAND-ZEN (¹³⁶Xe)
- SNO+ (¹³⁰Te)

Pros: "simple", large detectors exist, self-shielding Cons: Not very specific, 2v background

Low density trackers

- NEXT, PandaX (¹³⁶Xe gas TPC)
- SuperNEMO (foils and gas tracking, ⁸²Se)
 Pros: Superb topological information
 Cons: Very large size → expensive

Liquid TPC

- nEXO (¹³⁶Xe)

Pros: Homogeneous with good E resolution and topology Cons: Does not excel in any single parameter

Recent results (>10²⁵ yr half life)

Isotope	Experiment	Isotope exposure (kg yr)	$T_{1/2}^{0\nu\beta\beta}$ average sensitivity (10 ²⁵ yr)	$T_{1/2}^{0\nu\beta\beta}$ limit 90CL (10 ²⁵ yr)	$T_{1/2}^{0\nu\beta\beta}$ limit 90CL (13.8 Gyr)	$\langle m_{ u} angle$ range from NME * (meV)	Reference
70 0	Gerda	82	5.8	>8	> 5.8 •10 ¹⁵	<120-260	Agostini et al., PRL 120 (2018) 132503
⁷⁶ Ge	Majorana	26	4.8	>2.7	>1.9·10 ¹⁵	<200-433	Alvis et al., arXiv:1902.02299 (2019)
¹³⁰ Te	CUORE	24	0.7	>1.3	>9.4·10 ¹⁴	<110-520	Alduino et al., PRL 120 (2018) 132501
136N -	EXO-200 Ne	3 41	5.0	>3.5	> 2.5 •10 ¹⁵	<93-286	Anton et al. arXiv:1906.02723 (2019)
¹³⁶ Xe	KamLAND-ZEN	504	5.6	>10.7	>7.7·10 ¹⁵	<61-165	Gando et al., PRL 117 (2016) 082503

* Note that the range of "viable" NME is chosen by the experiments and uncertainties related to g_A are not included

I will concentrate on these proven techniques in the rest, while just mentioning (or omitting --apologies) many other concepts and ideas.

I will omit discussing bolometers since there is a dedicated talk by F.Bellini (Cuore/Cupid)

The next step:

ton-scale detectors entirely covering the inverted hierarchy

Testing lepton number violation with >100x the current $T_{1/2}$ sensitivity.

Modern 0vββ detectors are truly beautiful machines, with every component carefully optimized to work in harmony with everything else.

A tour of the proposed techniques.

A (possibly controversial) rambling on scalability and unbiased measurements

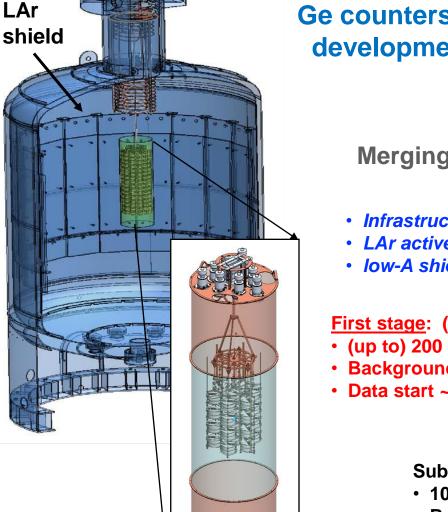
- Much has been said about the appropriateness of performing blind analyses to obtain unbiased experiments. "With moderation" this is, of course, very good.

 Much has been said about the convenience (particularly when seeking funding!) of scalable detectors, whereby the active mass can be increased gradually.

But, these two laudable ideas are not entirely compatible with each other! In some extreme every increase of detector mass is guaranteed not to be able to make a discovery and the discovery occurs adiabatically by means of a never-improving limit and liability to biases.

LEGEND: ⁷⁶Ge detectors





Ge counters played an essential role in the early developments on the field.

E.Fiorini et al. Phys Lett 25B (1967) 602

Merging of Gerda and Majorana programs

- Infrastructure
- LAr active veto
- low-A shield

- Electroformed Cu
- Low-E threshold
- Radio-pure low-noise electronics

First stage: (Legend-200)

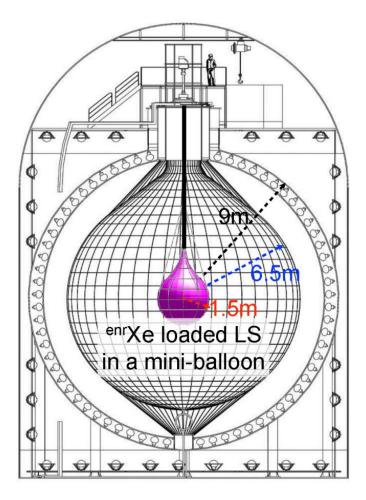
- (up to) 200 kg by upgrading the existing infrastructure at LNGS
- Background goal 0.6 cts/(FWHM t yr)
- Data start ~2021

Subsequent stages:

- 1000 kg, staged via individual payloads
- Background goal <0.1 cts/(FWHM t yr)
- Required depth (Ge-77m) under investigation

Isotope dissolved in a large liquid scintillation detector More advanced detector KamLAND-ZEN, using ¹³⁶Xe





In many ways this is the "extreme opposite" to the Ge and bolometric detectors (poor energy resolution, not very specific, "low tech" but huge and homogeneous)

KamLAND detector:

- 1kton of isoparaffine-pseudocumene liquid scintillator
- ~2000 20"/17" PMTs
- ~2.5m-thick paraffine buffer to shield active volume from PMT activity
- External, active water Cherenkov veto
- "Ballon" separates active (scintillating) volume from non-scintillating buffer
- Taking data since 2002, with glorious contributions to neutrino oscillation (reactors) and geo-neutrino measurements

- Added "miniballoon" to contain the ¹³⁶Xe-doped scintillator



Liquid Xe Time Projection Chambers: EXO-200 and nEXO

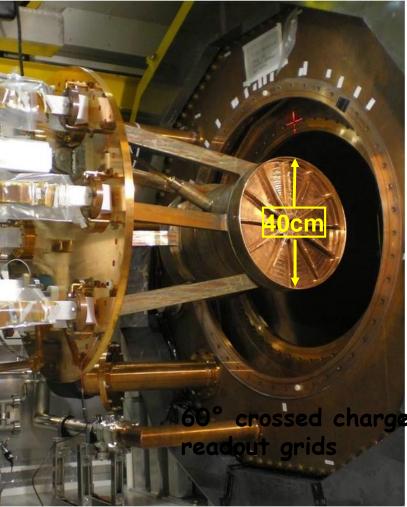


Example: the EXO-200 detector

Cathode (not shown)

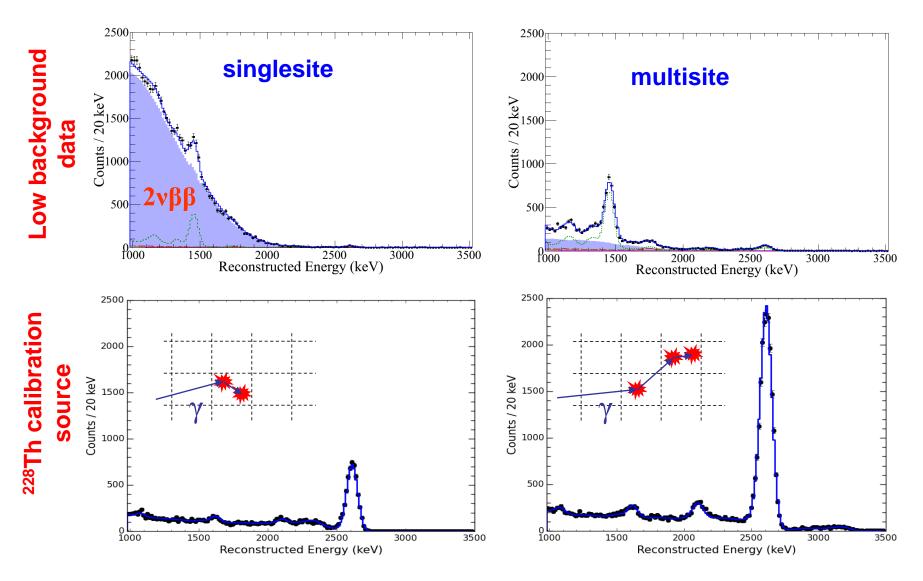
1/2 of 150kg of Liquid enrXe

1/2 of 600 Large Area APDs and charge collection wires

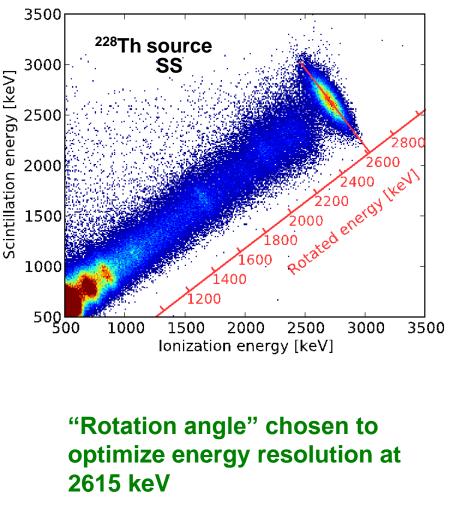


Field-shaping rings

Using event multiplicity to recognize the dominant γ backgrounds

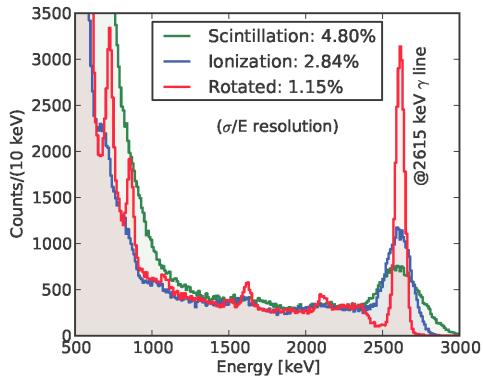


Combining Ionization and Scintillation to obtain the best energy resolution

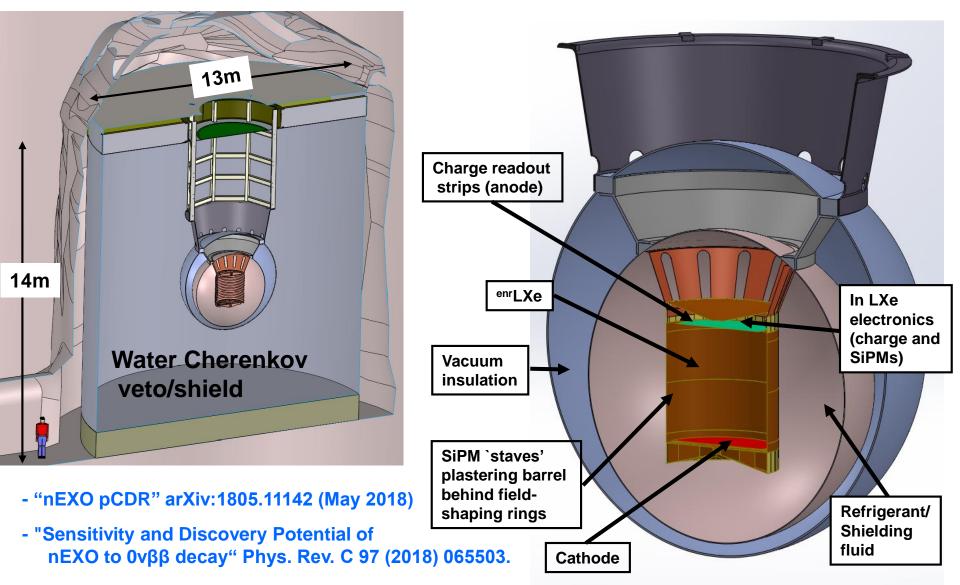


While no one really understands the energy resolution in LXe, scintillation and ionization are anticorrelated and this can be exploited to improve the energy resolution

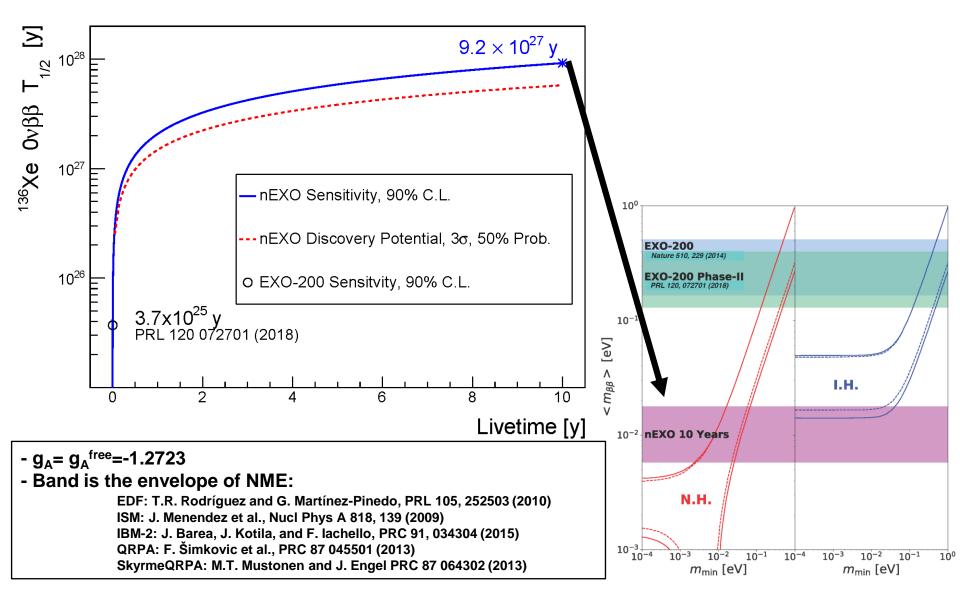
E.Conti et al. Phys Rev B 68 (2003) 054201



Moving towards the tonne scale: the 5000kg ^{enr}Xe nEXO detector



nEXO sensitivity as a function of time for the baseline design



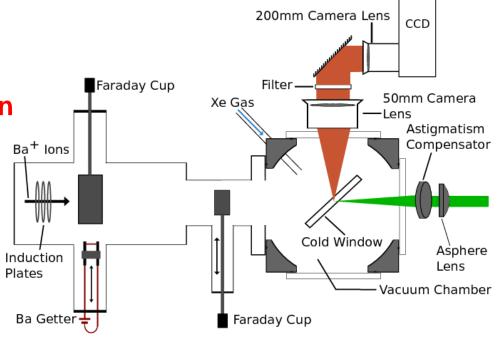
Conclusions on the matter of the creation of matter

- 0vββ searches are discovery science, with connections to many areas of modern physics
- 100 kg yr exposure and few x 10²⁵yr sensitivity are becoming passée
 - No discovery yet
- Looking at more than one isotope is important
- Tonne-scale experiments are being designed and soon built

In fact, ¹³⁶Xe offers the possibility to confirm a ββ decay by retrieving and tagging spectroscopically the Ba atom in the final state.

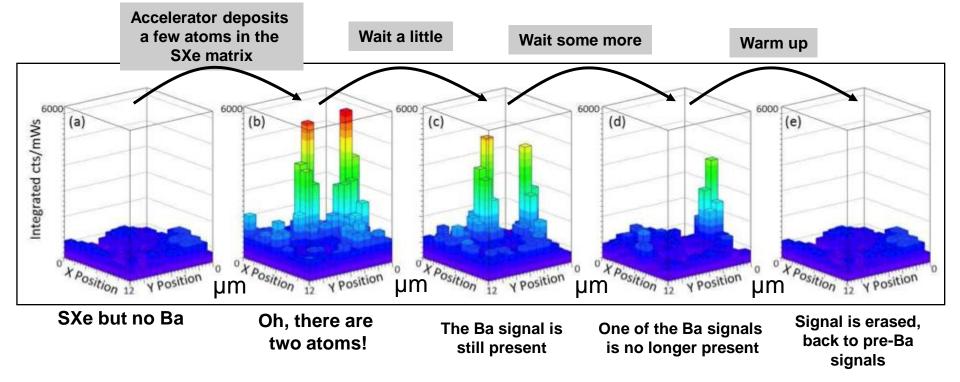
This is not necessary for nEXO to reach its design sensitivity and, indeed, it is not part of the design presented in the pre-CDR.

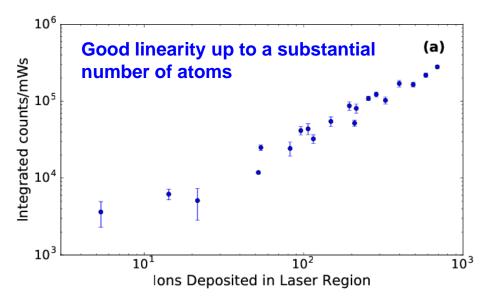
Nevertheless the "physics component" of the technique was recently demonstrated, including the ability to delete "old" Ba atoms (i.e. there is no "memory effect").



This work only addresses the physics feasibility, while the engineering of its implementation has not been explored yet.

Possibly Ba tagging could become a long term upgrade patch, extending the sensitivity of the experiment after a 5 to 10yr run in the baseline configuration.

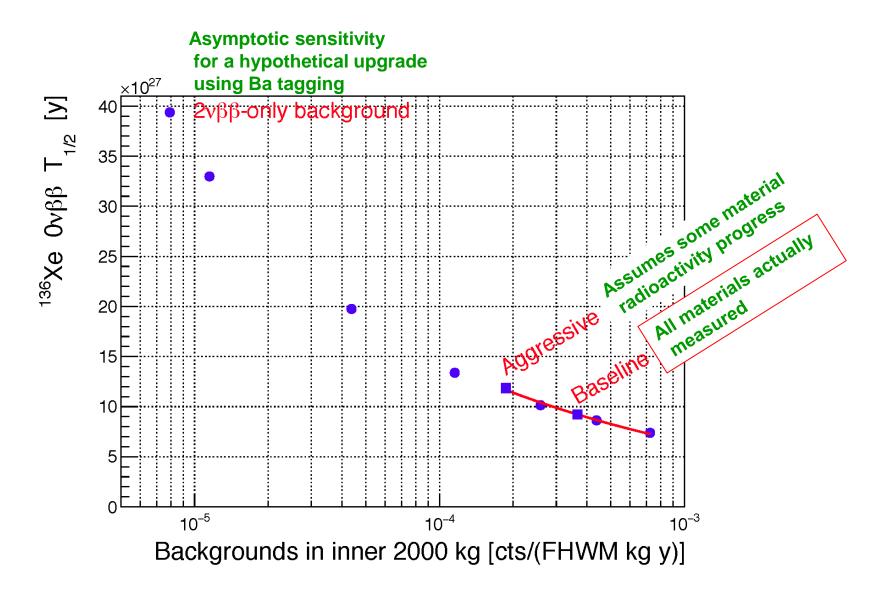




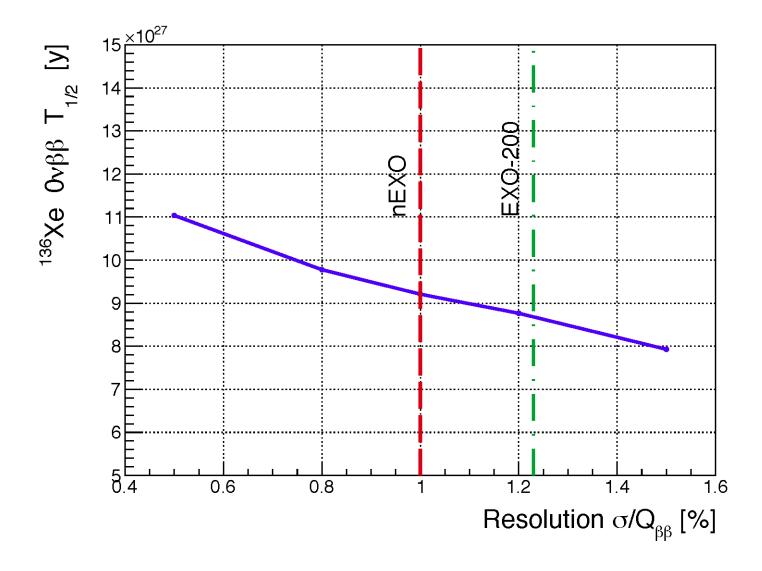
C.Chambers et al. Nature 569 (2019) 203

see also similar result in A.D. McDonald et al., Phys Rev Lett 120 (2018) 132504.

How does the sensitivity scale with background assumptions?



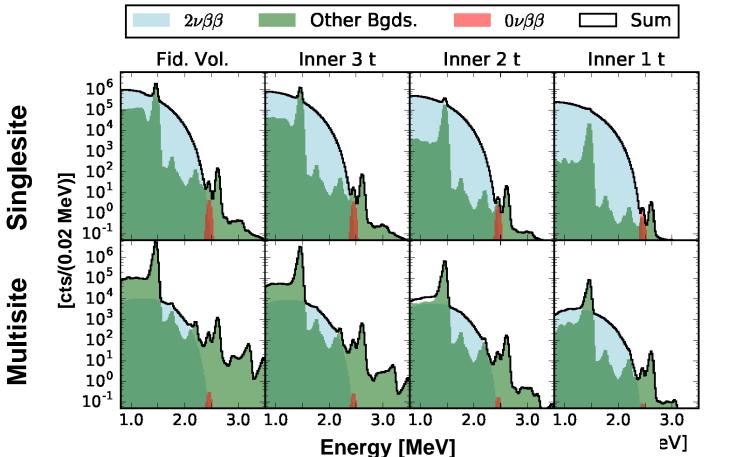
...and with energy resolution



Particularly in the larger nEXO, background identification and rejection fully use a fit considering simultaneously energy, $e-\gamma$ and α - β discrimination and event position.

➔ The power of the homogeneous detector,

this is not just a calorimetric measurement!



Corresponding to 10 yr data, with 0vββ T^{1/2}=5.7x10²⁷ yr

So, a simple "background index" is not the entire story.

- The innermost LXe mostly measures signal
- The outermost LXe mostly measures background
- The overall fit knows all this (and more) very well and uses all the information available to obtain the best sensitivity

Nevertheless, for the aficionados of "background index", here it is, as a function of depth in the TPC. For the inner 3000 kg this is better than 10⁻³ (kg yr FWHM)⁻¹

