

The PTOLEMY experiment: an opportunity for nuclear physics

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International Conference on Cosmic Neutrino Background Detection and Dark Matter Searches with PTOLEMY

The “basic model” of nuclear theory

- Goal of nuclear theory: comprehensive description of nuclear systems
 - NN scattering data: thousands of exp. data ($d\sigma/d\Omega \dots$)
 - Spectra and static properties: binding energies, radii, mag. mom. ...
 - Nucleonic matter EoS: neutron stars ...
- Inputs
 - Degrees of freedom (nucleons, pions, ...)
 - Many-body interactions between the constituents

$$H = \sum_{i=1}^A \frac{\mathbf{p}_i^2}{2m_i} + \sum_{i < j=1}^A v_{ij} + \sum_{i < j < k=1}^A V_{ijk} + \dots$$

One-body Two-body Three-body
 NN NNN

- Electroweak current operators

$$J^{EW} = \sum_{i=1}^A j_i + \sum_{i < j=1}^A j_{ij} + \sum_{i < j < k=1}^A j_{ijk} + \dots$$

One-body Two-body Three-body

Few-nucleon systems: the *ab-initio* approach



- ***Ab-initio* method** → obtain X by solving the relevant quantum many-body equations, without any uncontrolled approximation
- controlled approximations are allowed (expansion on a certain basis)
→ converged results = *ab-initio* results
- comparison of *ab-initio* results obtained with different *ab-initio* methods
→ **benchmark calculations**
- comparison of *ab-initio* results with data
→ **test of H & J^{EW}**
→ **predictions for “un-measurable” observables**

How few is “few”? ⇒ Few-body \leftrightarrow *ab-initio* methods

$A = 3$ is few!

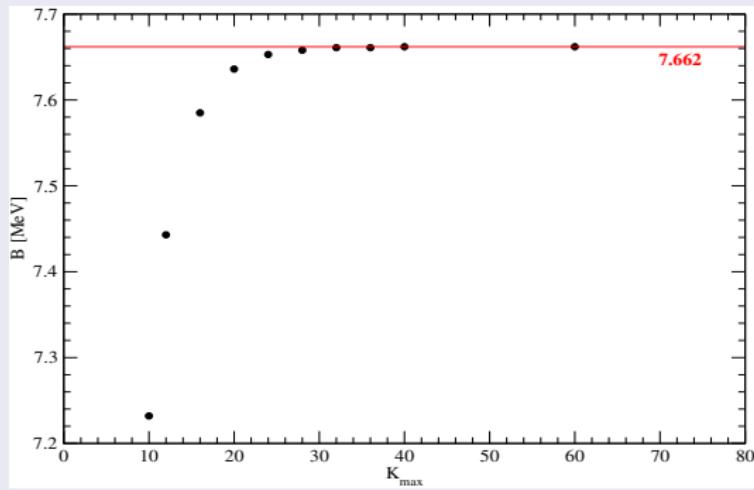
An *ab-initio* approach: the Hyperspherical Harmonics (HH) Method

Bound states

$$\Psi^{JJ_z} = \sum_{\mu} c_{\mu} \Psi_{\mu}$$

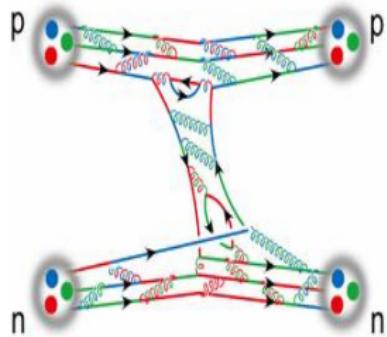
- $\Psi_{\mu} \rightarrow$ known functions (spin-isospin HH functions)
- Rayleigh-Ritz var. principle: $\delta_c \langle \Psi^{JJ_z} | H - E | \Psi^{JJ_z} \rangle = 0$
⇒ Solve for E and c_{μ}

Convergence of the method: $B(^3\text{H})$ with first 3 spin-isospin channel – N3LO



→ bound/scattering states $A \leq 4$, low-energy scattering states (astrophysical interest)

The nuclear Hamiltonian: a little bit of history



Nuclear interaction: $V_{NN} + V_{NNN}$

- Until $\simeq 20$ years ago:
phenomenological potentials
 - $V_{NN} + V_{NNN}$ semi-phenomenological
 - V_{NN} with $\simeq 40$ parameters fitted to $A = 2$ data $\rightarrow \chi^2/\text{datum} \simeq 1$
 - V_{NNN} with 2-3 parameters fitted to $B(A = 3, 4)$
- ⇒ **no simple connection to QCD**
- ⇒ **no clear relation for H & J^{EW}**
- Then ... chiral effective field theory (χEFT)

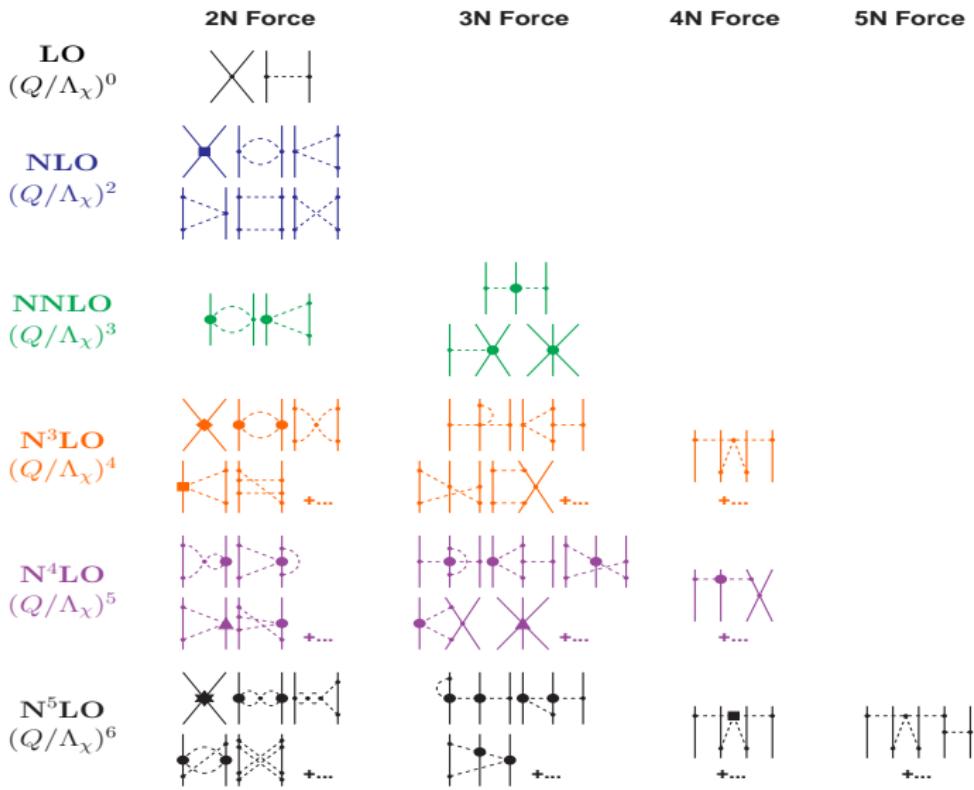
Chiral Effective Field Theory (χ EFT): a short summary

- QCD → quarks and gluons (“high-energy” d.o.f.)
- Nuclear physics → nucleons and pions (“low-energy” d.o.f.)
- EFT → processes with $E \simeq p \simeq m_\pi \ll \Lambda_\chi \sim 1 \text{ GeV}$
 - ★ “h-e” d.o.f. integrated out → contact interactions with “l-e” d.o.f. and **low-energy constants (LECs) obtained from experiment**
 - ★ **perturbative theory**: matrix elements $\propto O(p/\Lambda_\chi)^\nu$
- **χ EFT → EFT with spontaneous breaking of QCD's χ -symmetry**
- Regularization with cutoff function → $\Lambda \simeq 414, 450, 500, 600 \text{ MeV}$

Disadvantage: limited to processes with $E \sim 1 - 2 m_\pi$

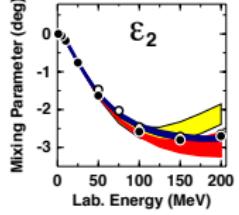
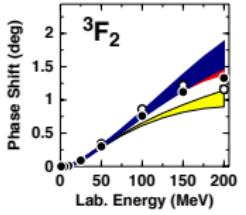
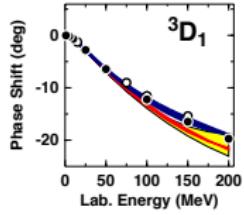
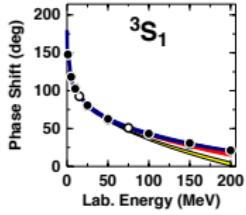
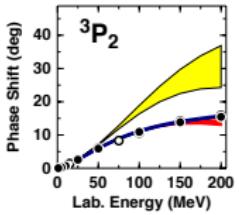
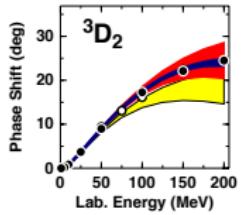
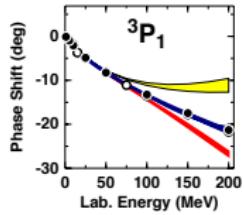
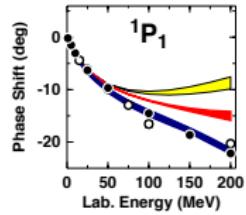
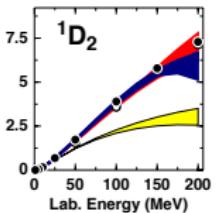
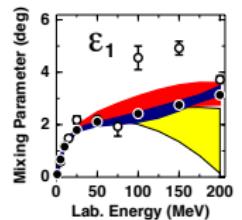
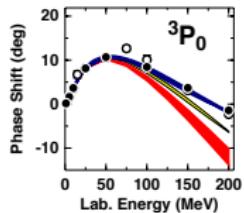
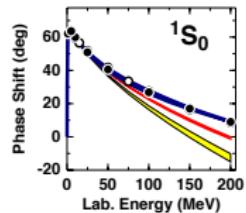
Advantages

- nuclear force “hierarchy” → accurate $V_{NN} + V_{NNN}$
- **consistent framework for H & J^{EW}** (add external EW field among the d.o.f.)



Machleidt and Sammarruca, Physics Scripta **91**, 083007 (2016)

⇒ each • = LECs (20-30) → fit to NN data



Yellow=NLO; Red=N2LO; Blue=N3LO

The J^{EW} operator

EW operators: $\rho^\gamma, \mathbf{j}^\gamma$; $\rho^{V/A}, \mathbf{j}^{V/A}$

$$\text{CVC} \Rightarrow \rho^V / \mathbf{j}^V \rightarrow \rho^\gamma / \mathbf{j}^\gamma$$

Power counting for j^γ

$$\mathcal{O}(Q^{-2}) \quad \left| \begin{array}{c} \text{---} \\ \text{---} \end{array} \right| \quad \mathbf{j}^{(-2)} \propto [e_N(1)(\mathbf{p}'_1 + \mathbf{p}_1) + i\mu_N(1)\sigma_1 \times \mathbf{q}] \times \delta(\mathbf{p}'_2 - \mathbf{p}_2) + 1 \leftrightarrow 2$$

$\mathcal{O}(Q^{-1})$  "standard" one-pion-exchange

$\mathcal{O}(Q^0)$   | ■ = relativistic corrections

PREDICTIONS \Leftarrow EM observables (many) \Leftarrow LECs

Static EM properties for $A = 2, 3$ nuclei

	PhenApp	χ EFT	Exp.
$r_c(d)$ [fm]	2.119	2.126(4)	2.130(10)
$Q(d)$ [fm 2]	0.280	0.2836(16)	0.2859(3)
$r_c(^3\text{He})$ [fm]	1.928	1.962(4)	1.973(14)
$r_m(^3\text{He})$ [fm]	1.909	1.920(7)	1.976(47)
$r_c(^4\text{He})$ [fm]	1.639	1.663(11)	1.681(4)

Marcucci *et al.*, JPG 43, 023002 (2016)

χ EFT \longrightarrow **theoretical error!**

Power counting for j^A

$$\mathcal{O}(Q^{-3}) \quad \left| \begin{array}{c} \diagup \\ \diagdown \end{array} \right|$$

$$\mathcal{O}(Q^{-2}) \quad \left| \begin{array}{c} \diagup \\ \diagup \end{array} \right|$$

$$\mathcal{O}(Q^{-1}) \quad \left| \begin{array}{c} \blacksquare \\ \diagup \\ \diagdown \end{array} \right|$$

$$\mathcal{O}(Q^0) \quad \left| \begin{array}{c} \diagup \\ \diagdown \end{array} \right| \quad \text{---} \quad \text{---}$$

- $\mathcal{O}(Q^1)$ not shown: loop and 2π -exchange contributions (hughly diagrams!)
- Only one LEC - d_R

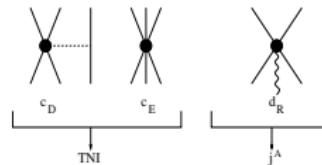
$$d_R = -\frac{M_N}{4\Lambda_\chi g_A} c_D + \frac{1}{3} M_N (c_3 + 2c_4) + \frac{1}{6}$$

Gårdestig and Phillips, PRL **96**, 232301 (2006)

Gazit *et al.*, PRL **103**, 102502 (2009)

Marcucci *et al.*, PRL **121**, 049901(E) (2018)

- fit c_D and c_E (in V_{NNN}) to $B(A=3)$ and Gamow-Teller m.e. of tritium β -decay (GT_{Exp})



⇒ **PREDICTIONS** for other observables

Predictions in the weak sector

^3H β -decay

$$(1 + \delta_R)t_{1/2}f_V = \frac{K/G_V^2}{\langle \mathbf{F} \rangle^2 + f_A/f_V g_A^2 \langle \mathbf{GT} \rangle^2}$$

- $g_A = 1.2723$; $\delta_R = 1.9\%$ outer radiative corrections; $t_{1/2}$ =half-life; $f_{V/A}$ = Fermi functions
 - Exp. values: $K/G_V^2 = (6144.5 \pm 1.9)$ s & $(1 + \delta_R)t_{1/2}f_V = (1134.6 \pm 3.1)$ s
- $\langle \mathbf{F} \rangle = \langle ^3\text{He} | \sum_j \tau_{j,+} | ^3\text{H} \rangle \simeq 0.999$ very stable from theory $\Rightarrow \langle \mathbf{GT} \rangle_{EXP}$

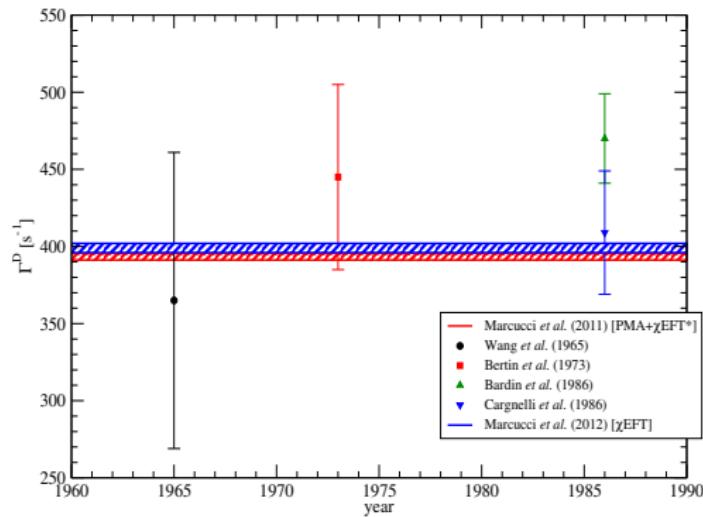
Polarized ^3H β -decay: ${}^3\overrightarrow{\text{H}} \rightarrow {}^3\text{He} + e^- + \bar{\nu}_e$

$$\frac{d\omega}{dE_e \, d\Omega_e \, d\Omega_\nu} \propto \xi [1 + a \vec{\beta} \cdot \hat{\nu} + \hat{P} \cdot (A\vec{\beta} + B\hat{\nu})]$$

- \hat{P} =pol. versor; $\hat{\nu} = \mathbf{p}_\nu/E_\nu$; $\vec{\beta} = \mathbf{p}_e/E_e$;
- $\xi \sim \langle \mathbf{F} \rangle^2 + f_A/f_V g_A^2 \langle \mathbf{GT} \rangle^2$; $a\xi \sim \langle \mathbf{F} \rangle^2 - f_A/f_V g_A^2 \langle \mathbf{GT} \rangle^2/3$
- $A\xi$ and $B\xi$: other combinations of $\langle \mathbf{F} \rangle$ and $\langle \mathbf{GT} \rangle$

Calculation:
work in progress
but NO DATA!

Muon capture on light nuclei: deuteron



$$\Gamma^D = 399(3) \text{ s}^{-1} \text{ vs. } \Gamma^D(\text{exp}) \dots$$

MuSun \rightarrow 1.5 % accuracy

Muon capture on light nuclei: ${}^3\text{He}$ (I)

- $\mu^- + {}^3\text{He} \rightarrow n + d + \nu_\mu$ (20%) [poor data]
- $\mu^- + {}^3\text{He} \rightarrow n + n + p + \nu_\mu$ (10%) [poor data]
- $\mu^- + {}^3\text{He} \rightarrow {}^3\text{H} + \nu_\mu$ (70%)

$\mu^- + {}^3\text{He} \rightarrow {}^3\text{H} + \nu_\mu$: two hyperfine states, $P(f, f_z) = (1; \pm 1, 0)$ and $(0; 0)$

$$\cos \theta = \hat{\mathbf{z}} \cdot \hat{\mathbf{q}} \leftarrow \text{momentum transfer of the lepton pair}$$

$$\frac{d\Gamma}{d(\cos \theta)} = d = \frac{1}{2} \Gamma_0 [1 + A_v P_v \cos \theta + A_t P_t (\frac{3}{2} \cos^2 \theta - \frac{1}{2}) + A_\Delta P_\Delta]$$

$$P_v = P(1, 1) - P(1, -1)$$

$$P_t = P(1, 1) + P(1, -1) - 2P(1, 0)$$

$$P_\Delta = P(1, 1) + P(1, -1) + P(1, 0) - 3P(0, 0) \equiv 1 - 4P(0, 0)$$

Muon capture on light nuclei: ${}^3\text{He}$ (II)

- Experimental data

- $\Gamma_0^{\text{EXP}} = 1496(4) \text{ s}^{-1}$
- $A_v^{\text{EXP}} = 0.63 \pm 0.09 \text{ (stat.)}^{+0.11}_{-0.14} \text{ (syst.)}$

Ackerbauer *et al.*, PLB **417**, 224 (1998)

Souder *et al.*, NIMA **402**, 311 (1998)

- Theoretical predictions

- $\Gamma_0 = 1492(19) \text{ s}^{-1}$
- $A_v = 0.5435(6)$
- $A_t = -0.355(1); A_\Delta = -0.101(2)$

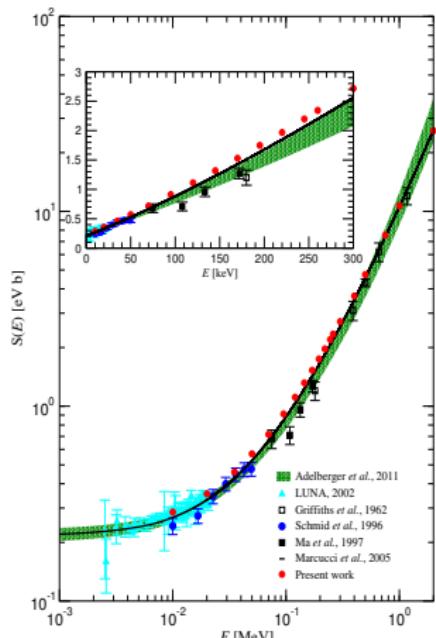
Marcucci *et al.*, PRL **108**, 052502 (2012); Erratum PRL **121**, 049901(E) (2018)

Bottom line

- Theory seems to be ok, but ...
- Need **more** and **more accurate** data in the weak sector to be sure!
- **PTOLEMY can play a role in this game**

An example in the EM sector: the $p + d \rightarrow {}^3\text{He} + \gamma$

- interesting for **BBN**: larger rate \Rightarrow smaller D/H primordial abundance
- on-going measurement by the **LUNA Collab.** at LNGS
- ab-initio* study \Leftarrow initial $p + d$ scattering state (only HH method available)



PArthENoPE \rightarrow D/H abundance:

$$\text{D/H}|_{TH} = (2.46 \pm 0.06) \times 10^{-5}$$

$\Omega_b h^2 \rightarrow$ Planck 2015 & standard N_{eff}

vs.

$$\text{D/H}|_{Exp} = (2.53 \pm 0.04) \times 10^{-5}$$

Marcucci *et al.*, PRL 116, 102501 (2016)

Collaboration with LUNA: theory used for simulation/rate estimate **but needs to be tested!**

Conclusions

- Goal of nuclear theory: comprehensive description of nuclear systems
- Few-body nuclei ($A = 3$) → *ab-initio* methods
 - **test the theoretical framework (H & J^{EW})**
 - **give solid predictions + theoretical uncertainty (χ EFT)**
- But **few (poor) data in the weak sector**



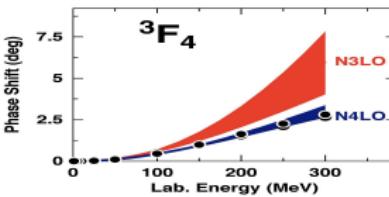
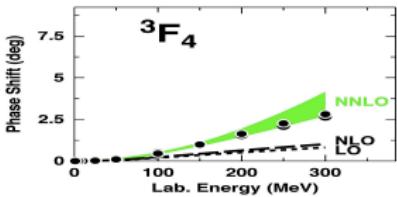
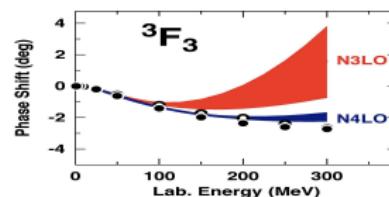
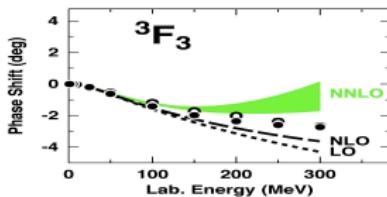
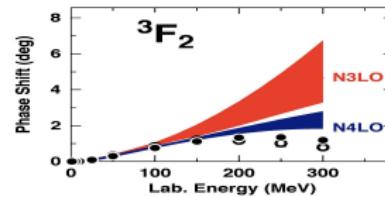
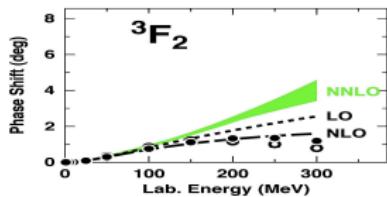
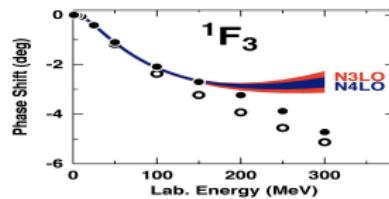
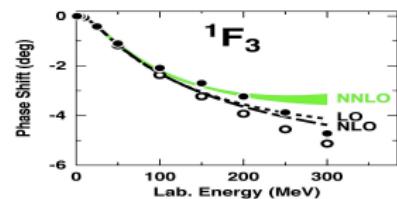
The PTOLEMY experiment: a **GREAT opportunity for nuclear physics**

THANK YOU!



Pon-
Tecorvo
Observatory for
Light,
Early-universe,
Massive-neutrino
Yield

SPARES



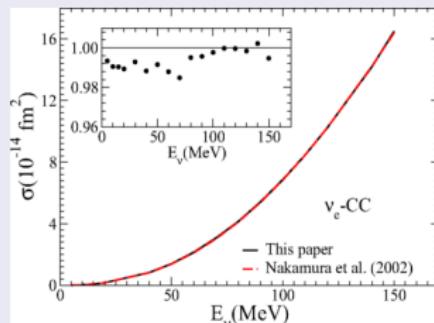
The PTOLEMY reaction: $\nu_e + {}^3\text{H} \rightarrow {}^3\text{He} + e^-$

Preliminary “related” studies

- $\nu_e + d \rightarrow e^- + p + p$

Baroni and Schiavilla,
PRC **96**, 014002 (2017)

in χ EFT



- $\bar{\nu}_e + {}^3\text{He} \rightarrow {}^3\text{H} + e^+$

Golak *et al.*,
PRC **98**, 015501 (2018)
in PhenApp

