

# Analysis of $B_s^0 \rightarrow \phi\phi$ decay mode

Giacomo Artoni, Benedetto Di Ruzza, Mirco Dorigo  
Lorenzo Ortolan, Marco Rescigno, Anna Maria Zanetti



**CDF ITALY**

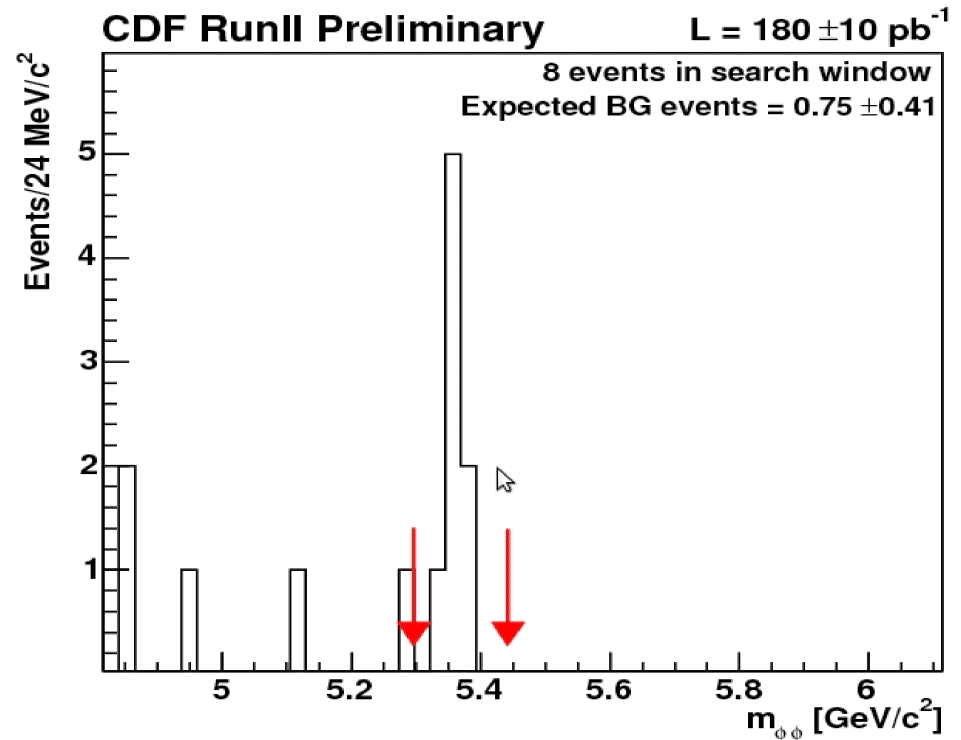
Fermilab 2 settembre 2009



# INTRODUCTION

The  $B_s^0 \rightarrow \phi \phi$  decay was observed for the first time by CDFII in 2005:  
*D.Acosta et al. (CDF Collaboration), Phys. Rev. Lett. 95,031801 (2005).*

Using 180 pb<sup>-1</sup> of data available at that time 8 events were observed.



Now this measurement was done using 2.9 fb<sup>-1</sup> of data available at CDFII (2002-up to April 2008).



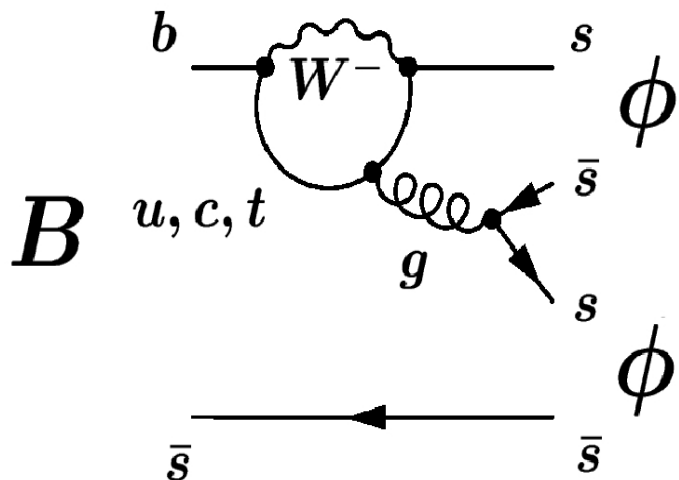
# OUTLINE

- 1) Motivation.
- 2) Theoretical predictions:  
the polarization puzzle.
- 3) Analysis strategy.
- 4) Signal selection.
- 5) Background evaluation.
- 6) Branching ratio measurement.
- 7) Conclusions and perspectives.



# MOTIVATION

The  $B_s^0 \rightarrow \phi \phi$  is a vector vector decay mode



The dominant decay element is the **penguin**  $b \rightarrow s \bar{s} s$

The final state is self - conjugate.

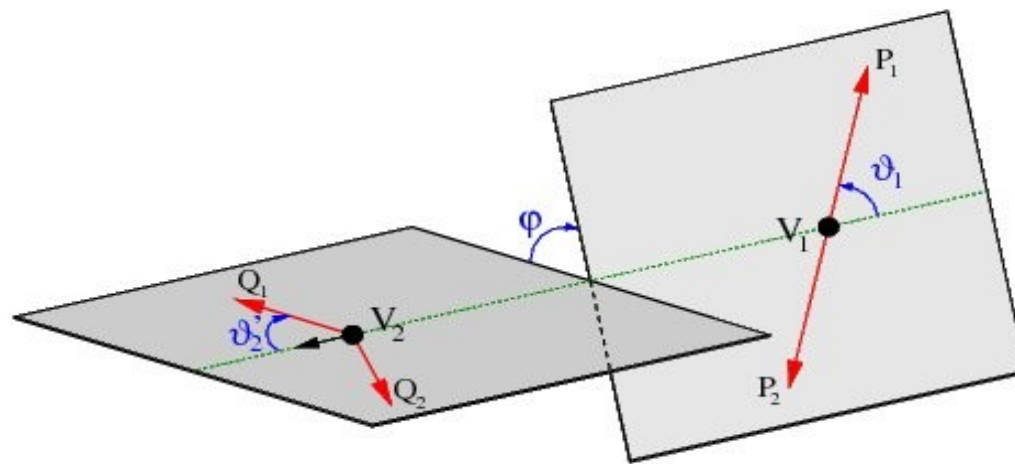
It can be used to measure  $\Delta \Gamma_s$ , CKM studies, and test of decay polarization prediction.



# MOTIVATION

The relation between **differential cross-section** and the observable **angular distribution** in the B meson and V rest frame using a helicity basis is:

$$\begin{aligned} \frac{d\Gamma_{B \rightarrow V_1 V_2 \rightarrow \dots}}{d\cos\vartheta_1 d\cos\vartheta_2 d\varphi} &\propto |\mathcal{A}_0|^2 \cos^2\vartheta_1 \cos^2\vartheta_2 + \frac{1}{4} \sin^2\vartheta_1 \sin^2\vartheta_2 (|\mathcal{A}_+|^2 + |\mathcal{A}_-|^2) \\ &\quad - \cos\vartheta_1 \sin\vartheta_1 \cos\vartheta_2 \sin\vartheta_2 [\operatorname{Re}(e^{-i\varphi} \mathcal{A}_0 \mathcal{A}_+^*) + \operatorname{Re}(e^{+i\varphi} \mathcal{A}_0 \mathcal{A}_-^*)] \\ &\quad + \frac{1}{2} \sin^2\vartheta_1 \sin^2\vartheta_2 \operatorname{Re}(e^{2i\varphi} \mathcal{A}_+ \mathcal{A}_-^*), \end{aligned}$$





## MOTIVATION

In Experimental analyses, observables are preferably defined in terms of the transversity as they have definite  $CP$  transformation properties. A typical set of observable consists of the branching fraction, two out of the three polarization fractions  $f_L$ ,  $f_{\parallel}$ ,  $f_{\perp}$ , and two phases  $\phi_{\parallel}$ ,  $\phi_{\perp}$ , where

$$f_{L,\parallel,\perp}^B = \frac{|\mathcal{A}_{0,\parallel,\perp}|^2}{|\mathcal{A}_0|^2 + |\mathcal{A}_{\parallel}|^2 + |\mathcal{A}_{\perp}|^2}; \quad \phi_{\parallel,\perp}^B = \arg \frac{\mathcal{A}_{\parallel,\perp}}{\mathcal{A}_0}$$

and naively, within the Standard Model, one expects

$$\frac{f_T}{f_L} \propto \frac{m_V}{m_B}$$



## Polarization puzzle

Within the standard model one expect:

$$\frac{f_T}{f_L} \propto \frac{m_V}{m_B}$$

And for the  $B_s^0 \rightarrow \phi \phi$  decay:  $\frac{f_T}{f_L} \propto \frac{1020}{5369}$

However it was observed in some decays, for example

$$B \rightarrow \varphi K^* \quad \text{and} \quad B \rightarrow \rho K^*$$

that **this relation** is not valid

Babar collaboration : hep-hex: 0705.1798 :  
 $f_T = 0.21_{\pm 0.05_{\pm 0.02}}$      $f_L = 0.49_{\pm 0.05_{\pm 0.03}}$



# Polarization puzzle

Explanations of this puzzle inside the SM are in:

Ali & Kramer hep-ph/070316: on recent BR and polarization prediction in the perturbative QCD approach.

Fleisher & Gronau in hep-ph/07094013

Some of them predict the Branching ratio too:

	BR[10 <sup>-6</sup> ]
QCD Factorisation	21.8 <sup>+1.1+30.4</sup> <sub>-1.1-17.0</sub>
	19.5 <sup>+1.0+13.1</sup> <sub>-1.0-8.0</sub>
QCD Factorisation	13.1
Naive Factorisation	9.05

M. Beneke et al., hep-hex/0612290.

X. Li et al. P.R.L. D.68, 114015(2003);  
D71 019902(2005);  
hep-hex/030936.





## MEASUREMENT STRATEGY

The measurement consists in the determination of this ratio:

$$\frac{BR(B_s \mapsto \phi\phi)}{BR(B_s \mapsto J/\psi \phi)} = \frac{N_{\phi\phi}}{N_{J/\psi\phi}} \cdot \frac{\epsilon_{(B_s \mapsto J/\psi\phi)}}{\epsilon_{(B_s \mapsto \phi\phi)}} \cdot \frac{BR(J/\psi \mapsto \mu\mu)}{BR(\phi \mapsto KK)} \cdot \epsilon_{\mu}$$

From MC (red arrow) points to  $\frac{\epsilon_{(B_s \mapsto J/\psi\phi)}}{\epsilon_{(B_s \mapsto \phi\phi)}}$   
From PDG (blue arrow) points to  $\frac{BR(J/\psi \mapsto \mu\mu)}{BR(\phi \mapsto KK)}$   
From data (red arrow) points to  $\epsilon_{\mu}$

$$\phi \rightarrow K^+ K^-$$

Where

$$J/\psi \rightarrow \mu\mu$$

We have performed a normalized branching ratio and not absolute because:

- 1) The two decays are topologically very similar.
- 2) In the ratio we can simplify a lot of sistematics.



# Selection Optimization

In order to choose the variables to be used in the event selection and select the best interval values of these variables, we followed the

procedure of the “maximization of the score function”  $\frac{S}{\sqrt{(S+B)}}$

$B_s \longrightarrow \phi\phi$	
Variable	cut

$L_{xy}$	$> 330\mu m$
$P_T^{K min}$	$> 0.7 \text{ GeV}/c$
$\chi_{xy}^2$	$< 17$
$d0(B)$	$< 65\mu m$
$d0_{max}^{\phi}$	$> 85\mu m$

$B_s \longrightarrow J/\psi\phi$	
Variable	cut

$L_{xy}$	$> 290\mu m$
$P_T^{\phi}$	$> 1.4 \text{ GeV}/c$
$\chi_{xy}^2$	$< 15$
$d0(B)$	$< 80\mu m$
$P_T^{J/\psi}$	$> 2.0 \text{ GeV}/c$



# Backgrounds

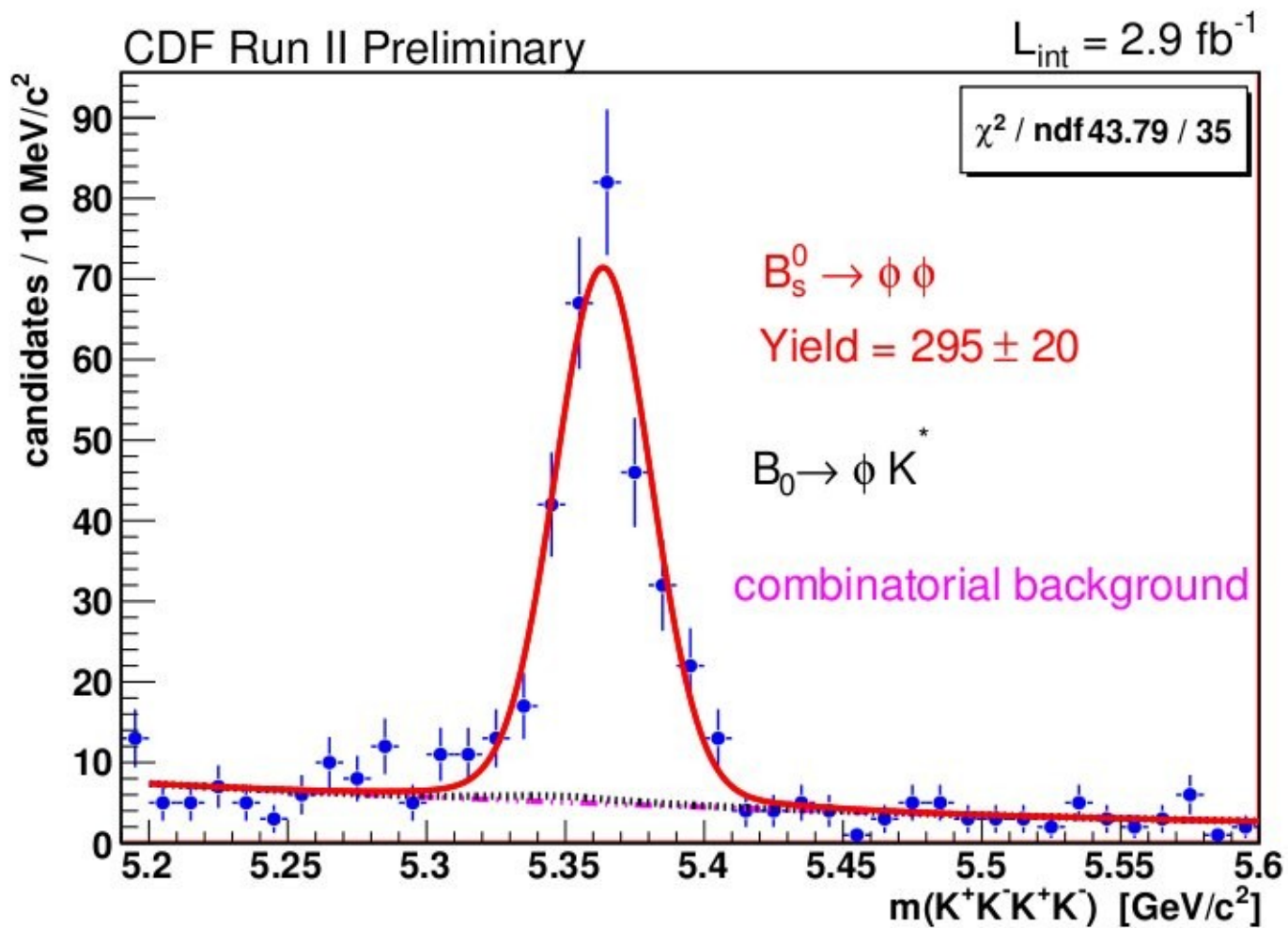
Expected background are:

- ▶ combinatorics (expected smooth function);
- ▶ decays not correctly reconstructed (reflections) peaking under the signal:
- ▶ for  $B_s \longrightarrow J/\psi\phi$ :
  - ▶  $B_d \longrightarrow J/\psi K^*$  ~ 4.2 % of the signal
  
- ▶ for  $B_s \longrightarrow \phi\phi$ :
  - ▶  $B_d \longrightarrow \phi K^*$  ~ 2.7 % of the signal
  - ▶  $B_s \longrightarrow \bar{K}^* K^*$  ~ 0 % of the signal

The reflections were evaluated using Monte Carlo sample

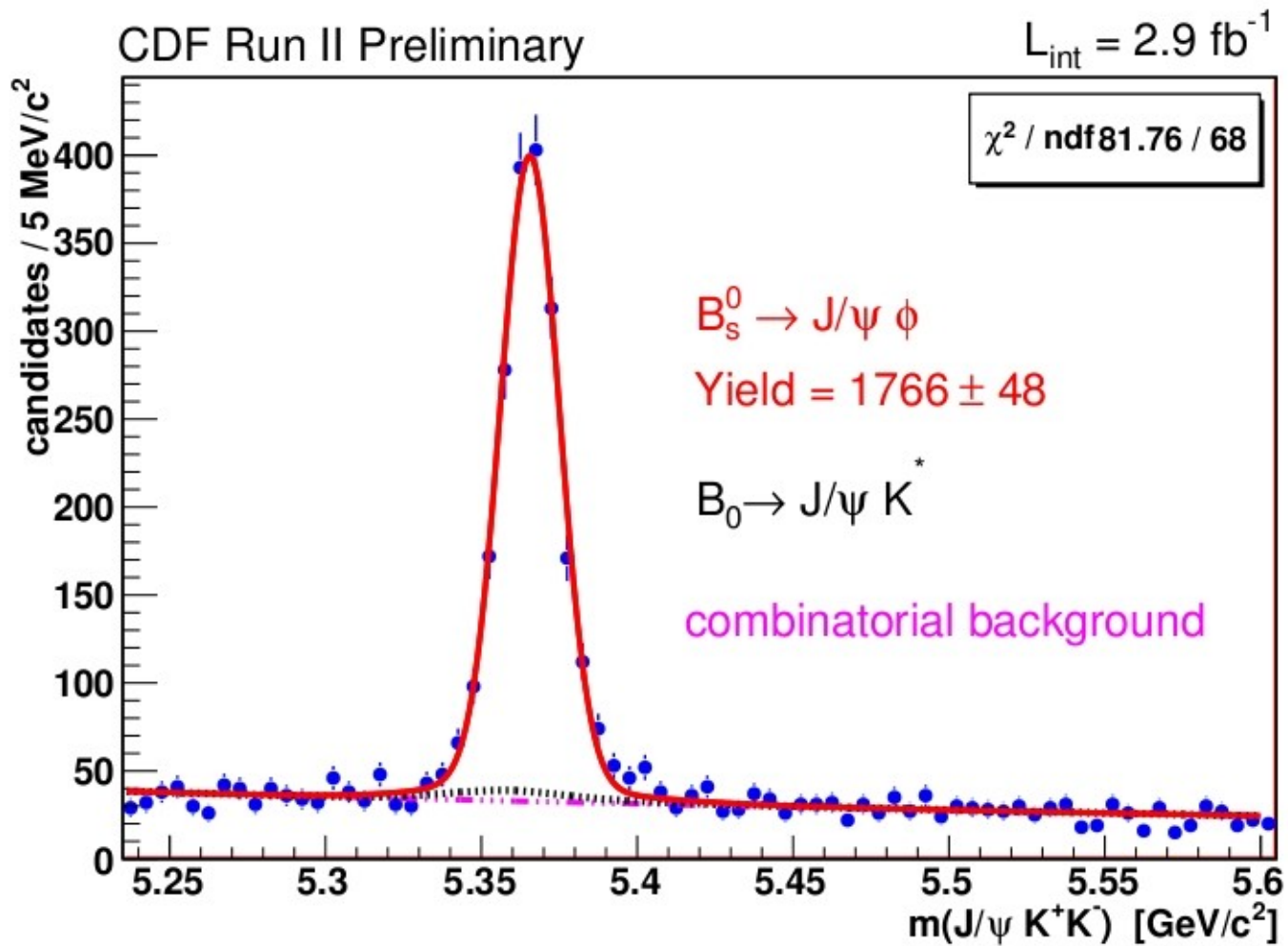


# Selected events





# Selected events





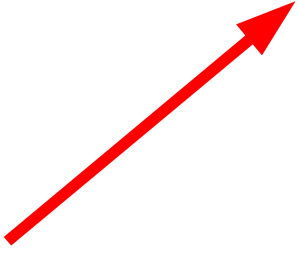
# Efficiencies

**Evaluated using MC**



$$\frac{BR(B_s \mapsto \phi\phi)}{BR(B_s \mapsto J/\psi \phi)} = \frac{N_{\phi\phi}}{N_{J/\psi\phi}} \cdot \frac{\epsilon_{(B_s \mapsto J/\psi\phi)}}{\epsilon_{(B_s \mapsto \phi\phi)}} \cdot \frac{BR(J/\psi \mapsto \mu\mu)}{BR(\phi \mapsto KK)} \cdot \epsilon_{\mu}$$

**Evaluated on data**

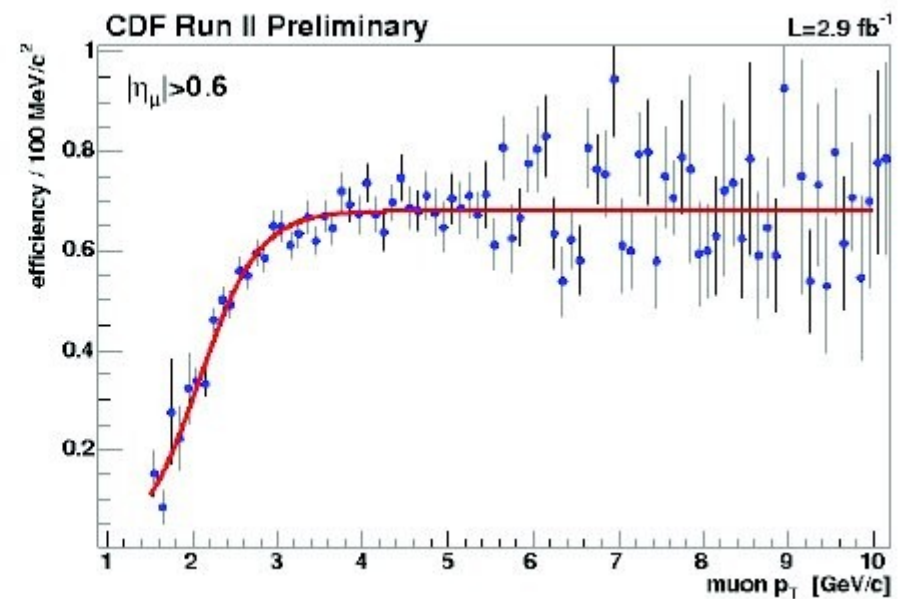
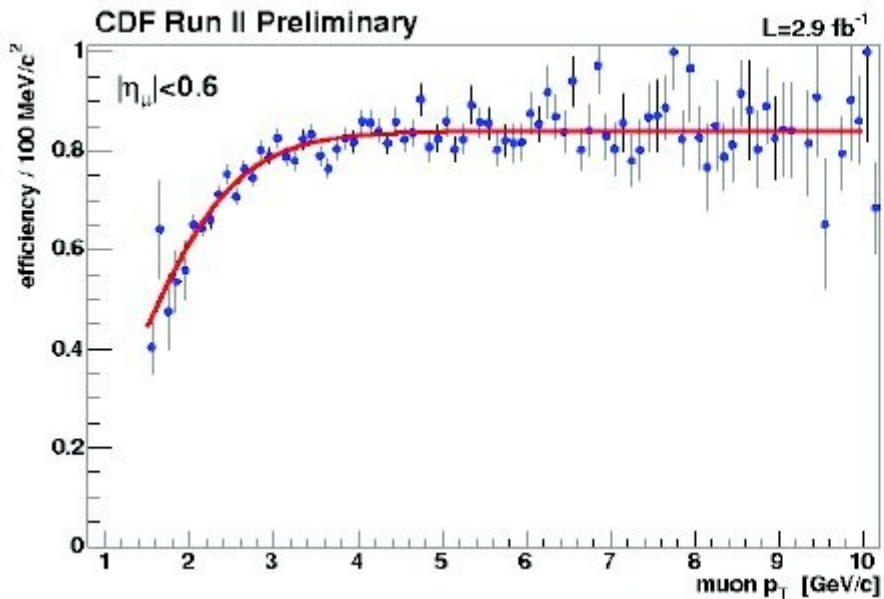




# Efficiencies

$$\frac{\epsilon_{\psi\phi}}{\epsilon_{\phi\phi}} = 0.939 \pm 0.030 \pm 0.009$$

$$\epsilon_{\psi\phi}^{\mu} = 0.8695 \pm 0.0044(\text{stat})$$





# Systematics Evaluation

We evaluated systematic uncertainty:

- ▶ due to different parametrization:
  - ▶ change in fit range  $\Rightarrow$  take into account possible structures below the mass peak due to unknown partially reconstructed
  - ▶ using one gaussian instead of two for the signal
- ▶ change in background subtraction  $\Rightarrow$  driven by  $B_s \rightarrow J/\psi\phi$
- ▶ due to effect not simulated in MonteCarlo used for example:
  - ▶ the polarization





## BR result

The result is:

$$\frac{BR(B_s \rightarrow \phi\phi)}{BR(B_s \rightarrow J/\psi\phi)} = [1.78 \pm 0.14^{stat} \pm 0.20^{syst}] \cdot 10^{-2}$$

$$BR(B_s^0 \rightarrow \phi\phi) = [24.0 \pm 2.1(stat) \pm 2.7(syst) \pm 8.2(BR)]10^{-6}$$



## BR result

$$\text{BR}(B_s^0 \rightarrow \phi\phi) = [24.0 \pm 2.1(\text{stat}) \pm 2.7(\text{syst}) \pm 8.2(\text{BR})]10^{-6}$$

	BR[10 <sup>-6</sup> ]
Experiment	14 <sup>+6</sup> <sub>-5</sub> (stat.) ± 6(syst.)
QCD Factorisation	21.8 <sup>+1.1+30.4</sup> <sub>-1.1-17.0</sub> 19.5 <sup>+1.0+13.1</sup> <sub>-1.0-8.0</sub>
QCD Factorisation	13.1
Naive Factorisation	9.05



## Conclusion and perspectives

- 1) The final result is in agreement with SM prediction.
- 2) Error is dominated by  $J_{\psi}/\Phi$  Branching ratio error
- 3) The sample selected allow to polarization studies



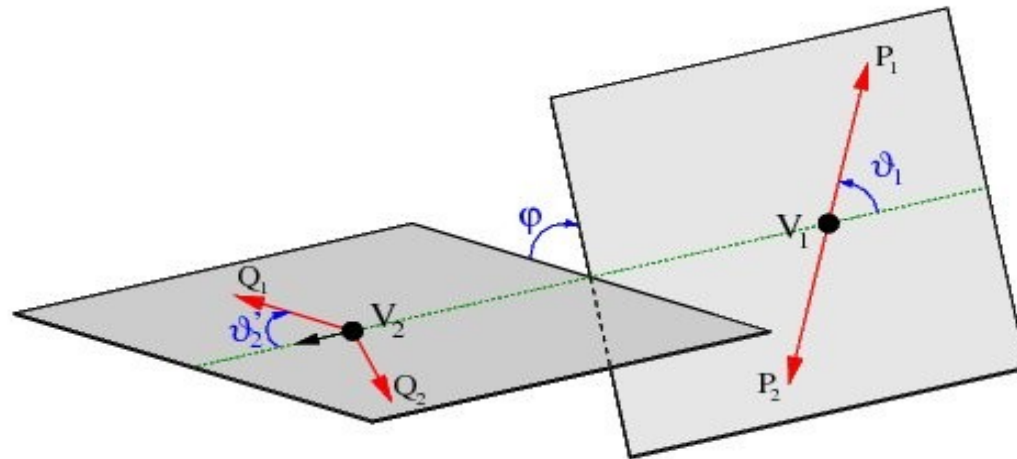
Backup slides



# The $B_s \rightarrow VV$ decays

The relation between differential cross-section and the observable angular distribution in the B meson rest frame is :

$$\begin{aligned} \frac{d\Gamma_{B \rightarrow V_1 V_2 \rightarrow \dots}}{d\cos\vartheta_1 d\cos\vartheta_2 d\varphi} &\propto |\mathcal{A}_0|^2 \cos^2\vartheta_1 \cos^2\vartheta_2 + \frac{1}{4} \sin^2\vartheta_1 \sin^2\vartheta_2 (|\mathcal{A}_+|^2 + |\mathcal{A}_-|^2) \\ &\quad - \cos\vartheta_1 \sin\vartheta_1 \cos\vartheta_2 \sin\vartheta_2 [\operatorname{Re}(e^{-i\varphi} \mathcal{A}_0 \mathcal{A}_+^*) + \operatorname{Re}(e^{+i\varphi} \mathcal{A}_0 \mathcal{A}_-^*)] \\ &\quad + \frac{1}{2} \sin^2\vartheta_1 \sin^2\vartheta_2 \operatorname{Re}(e^{2i\varphi} \mathcal{A}_+ \mathcal{A}_-^*), \end{aligned}$$





# The $B_s \rightarrow VV$ decays

Considering a  $B$  meson with four-momentum  $p_B$  decaying into light vector mesons  $V_1(p_1, \eta^*)$ ,  $V_2(p_2, \epsilon^*)$ , with masses  $m_{1,2}$  of order  $\Lambda_{QCD}$ , the decay amplitude can be decomposed into three scalar amplitude  $S_{1,2,3}$  according to

$$\mathcal{A}_{B \rightarrow V_1 V_2} = i\eta^{*\mu} \epsilon^{*\nu} \left( S_1 g_{\mu\nu} - S_2 \frac{p_{B\mu} p_{B\nu}}{m_B^2} + S_3 i\epsilon_{\mu\nu\rho\sigma} \frac{p_1^\rho p_2^\sigma}{p_1 \cdot p_2} \right),$$

or alternatively an helicity basis can be defined as :

$$\begin{aligned} \mathcal{A}_0 &= \mathcal{A}(B \rightarrow V_1(p_1, \eta_0^*) V_2(p_2, \epsilon_0^*)) = \frac{im_B^2}{2m_1 m_2} \left( S_1 - \frac{S_2}{2} \right), \\ \mathcal{A}_\pm &= \mathcal{A}(B \rightarrow V_1(p_1, \eta_\pm^*) V_2(p_2, \epsilon_\pm^*)) = i(S_1 \mp S_2), \end{aligned}$$

or a transversity amplitude basis can be defined using  $\mathcal{A}_L = \mathcal{A}_0$ , and replacing  $\mathcal{A}_\pm$  with

$$\begin{aligned} \mathcal{A}_\parallel &= \frac{(\mathcal{A}_+ + \mathcal{A}_-)}{\sqrt{2}} \\ \mathcal{A}_\perp &= \frac{(\mathcal{A}_+ - \mathcal{A}_-)}{\sqrt{2}} \end{aligned}$$



# The $B_s \rightarrow VV$ decays



In Experimental analyses, observables are preferably defined in terms of the transversity as they have definite  $CP$  transformation properties. A typical set of observable consists of the branching fraction, two out of the three polarization fractions  $f_L$ ,  $f_{\parallel}$ ,  $f_{\perp}$ , and two phases  $\phi_{\parallel}$ ,  $\phi_{\perp}$ , where

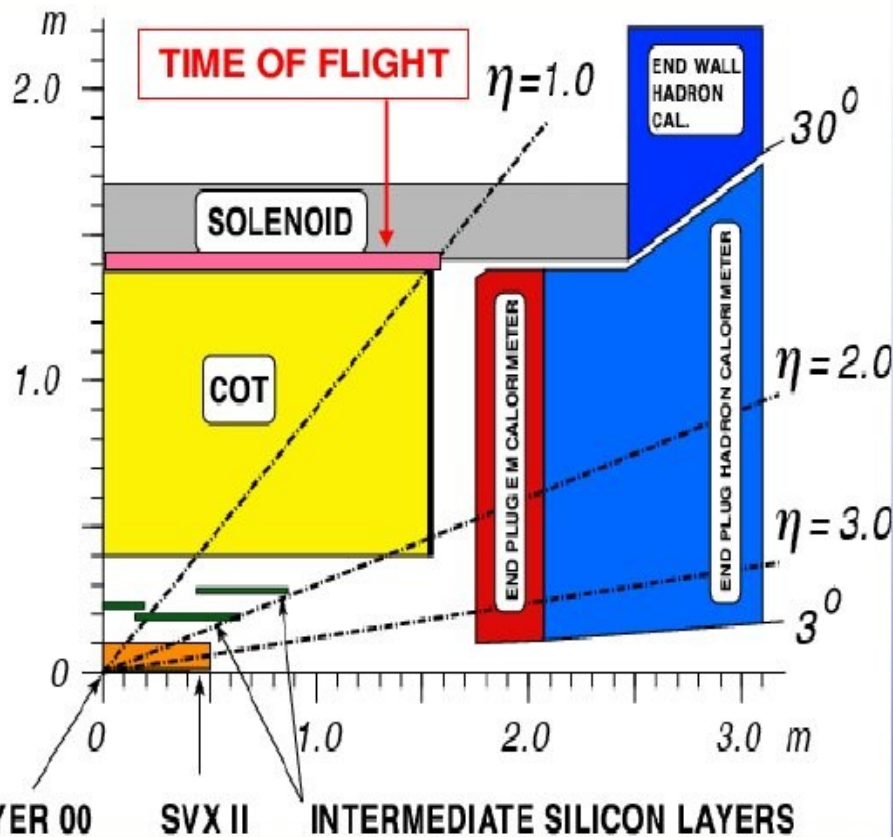
$$f_{L,\parallel,\perp}^B = \frac{|\mathcal{A}_{0,\parallel,\perp}|^2}{|\mathcal{A}_0|^2 + |\mathcal{A}_{\parallel}|^2 + |\mathcal{A}_{\perp}|^2}; \quad \phi_{\parallel,\perp}^B = \arg \frac{\mathcal{A}_{\parallel,\perp}}{\mathcal{A}_0}$$

and naively, within the Standard Model, one expects

$$\frac{f_T}{f_L} \propto \frac{m_V}{m_B}$$



# Tracking system



TOF: 100ps resolution, 2 sigma  $K/\pi$  separation for tracks below 1.6 GeV/c (significant improvement of  $B_s$  flavor tag effectiveness)

COT: large radius (1.4 m) Drift C.

- 96 layers, 200ns drift time
  - Precise  $P_T$  above 400 MeV/c
  - Precise 3D tracking in  $|\eta| < 1$
- $\sigma(1/P_T) \sim 0.1\% \text{GeV}^{-1}$ ;  $\sigma(\text{hit}) \sim 150 \mu\text{m}$
- $dE/dx$  info provides  $>1.3$  sigma  $K/\pi$  separation above 2 GeV

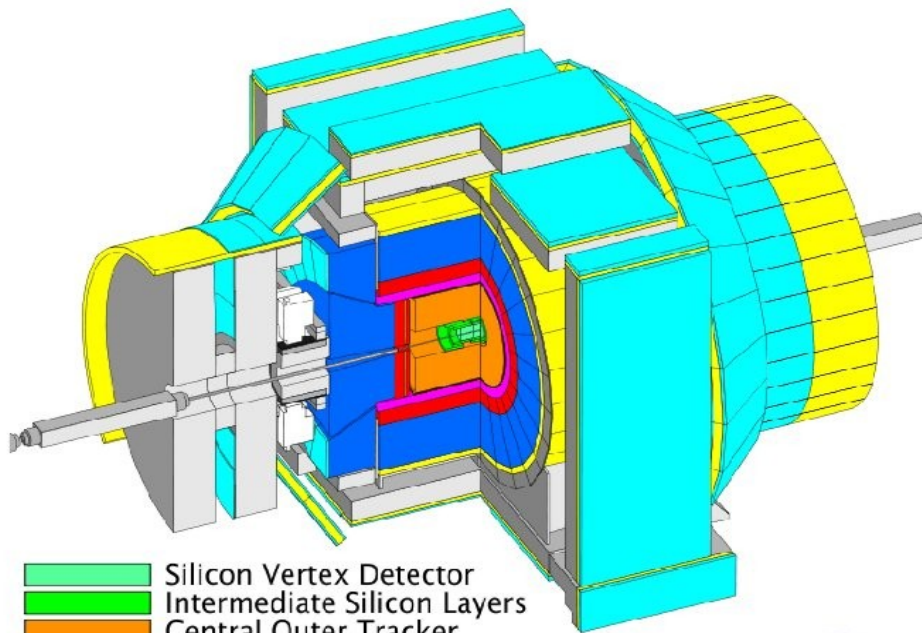
SVX-II + ISL: 6 (7) layers of double-side silicon ( $3\text{cm} < R < 30\text{cm}$ )

- Standalone 3D tracking up to  $|\eta| = 2$
- Very good I.P. resolution:  $\sim 30 \mu\text{m}$  ( $\sim 20 \mu\text{m}$  with Layer 00)

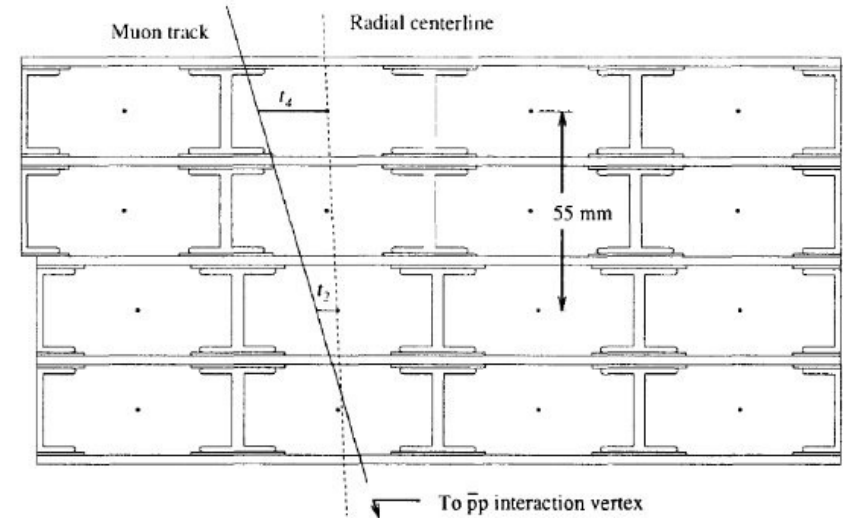




# Muons detectors



- Silicon Vertex Detector
- Intermediate Silicon Layers
- Central Outer Tracker
- Time Of Flight
- 1.4 T Superconducting Solenoid
- EM Calorimeter
- Hadron Calorimeter
- Muon Counters/Chambers

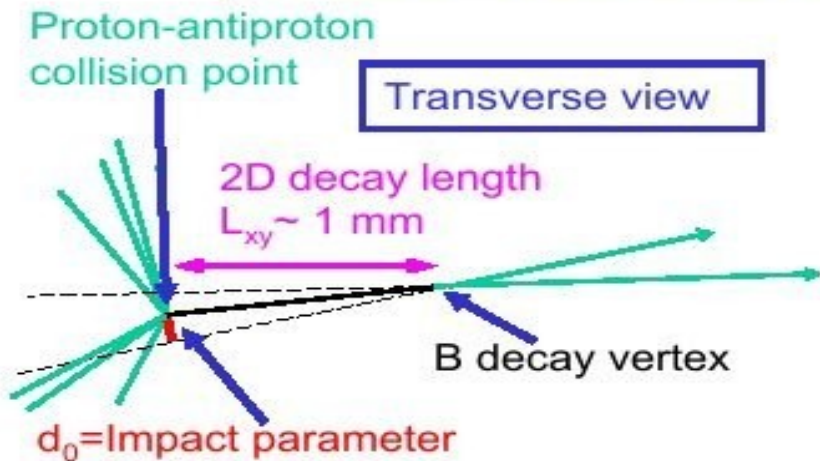




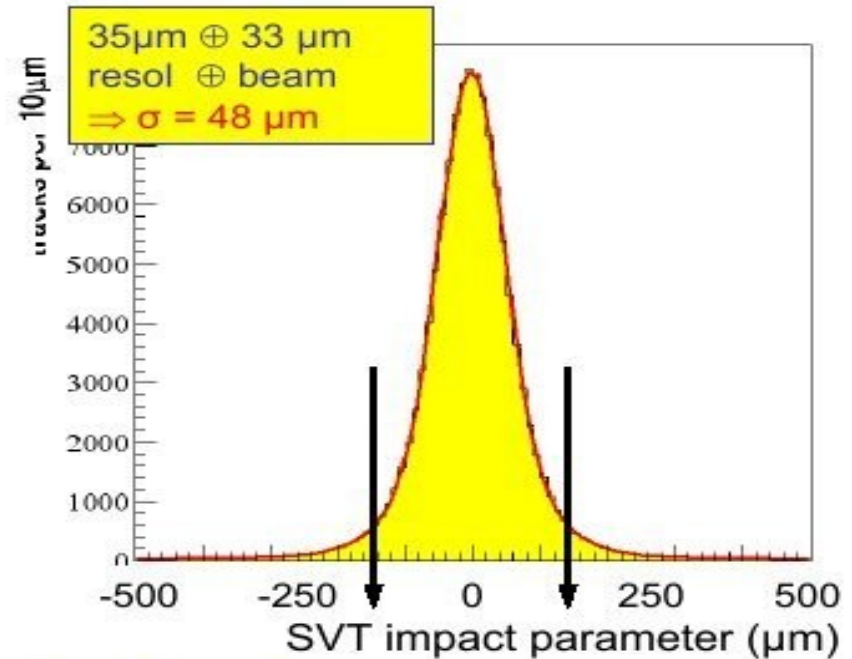
# Relevants detector elements for this analysis:



## CDF Level 2 Silicon Vertex Trigger



**Exploit long b, c lifetimes in Trigger!**  
**L1 track + Si hits = Impact parameter @L2**  
**A first at a hadron collider!**  
**CDF is a charm/ B Factory!**



### Lepton ( $e, \mu$ ) + displaced track trigger

Lepton:  $p_T > 4 \text{ GeV}$

Track:  $p_T > 2 \text{ GeV}$ ,  $d_0 > 120 \mu\text{m}$

Semi-leptonic B decays ( $B \rightarrow \ell \nu X$ )

### Displaced two track trigger

Tracks:  $p_T > 2 \text{ GeV}$ ,  $d_0 > 120 \mu\text{m}$

$\Sigma p_T > 5.5 \text{ GeV}$

Fully hadronic B decays ( $B \rightarrow hh'$ ,  $B_s \rightarrow D_s \pi$ ,  $D \rightarrow K \pi \dots$ )

# systematics

	$B_s^0 \rightarrow \phi\phi$	$B_s^0 \rightarrow J/\psi\phi$
	$\Delta N_{\phi\phi}/N_{\phi\phi}$	$\Delta N_{J/\psi\phi}/N_{J/\psi\phi}$
fit range	3%	-
signal parametrization	3%	2%
background subtraction: error on BRs	1%	1%
	$\Delta \varepsilon_{\phi\phi}/\varepsilon_{\phi\phi}$	$\Delta \varepsilon_{J/\psi\phi}/\varepsilon_{J/\psi\phi}$
polarization in MC	7%	6%
	$\Delta \varepsilon_{\phi\phi}/\varepsilon_{J/\psi\phi}$	
XFT particle dep.	4%	
$p_T$ reweight	0.9%	
	$\Delta \varepsilon_{\mu}/\varepsilon_{\mu}$	
$\eta$ parametrization & correlation	0.9%	

Table 16: *Contributions to the total relative uncertainty from the systematic uncertainty sources considered.*



# MOTIVATION

Considering a  $B$  meson with four-momentum  $p_B$  decaying into light vector mesons  $V_1(p_1, \eta^*)$ ,  $V_2(p_2, \epsilon^*)$ , with masses  $m_{1,2}$  of order  $\Lambda_{QCD}$ , the decay amplitude can be decomposed into three scalar amplitude  $S_{1,2,3}$  according to

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or alternatively an helicity basis can be defined as :

$$\begin{aligned} \mathcal{A}_0 &= \mathcal{A}(B \rightarrow V_1(p_1, \eta_0^*) V_2(p_2, \epsilon_0^*)) = \frac{im_B^2}{2m_1 m_2} \left( S_1 - \frac{S_2}{2} \right), \\ \mathcal{A}_\pm &= \mathcal{A}(B \rightarrow V_1(p_1, \eta_\pm^*) V_2(p_2, \epsilon_\pm^*)) = i(S_1 \mp S_2), \end{aligned}$$

or a transversity amplitude basis can be defined using  $\mathcal{A}_L = \mathcal{A}_0$ , and replacing  $\mathcal{A}_\pm$  with

$$\begin{aligned} \mathcal{A}_\parallel &= \frac{(\mathcal{A}_+ + \mathcal{A}_-)}{\sqrt{2}} \\ \mathcal{A}_\perp &= \frac{(\mathcal{A}_+ - \mathcal{A}_-)}{\sqrt{2}} \end{aligned}$$

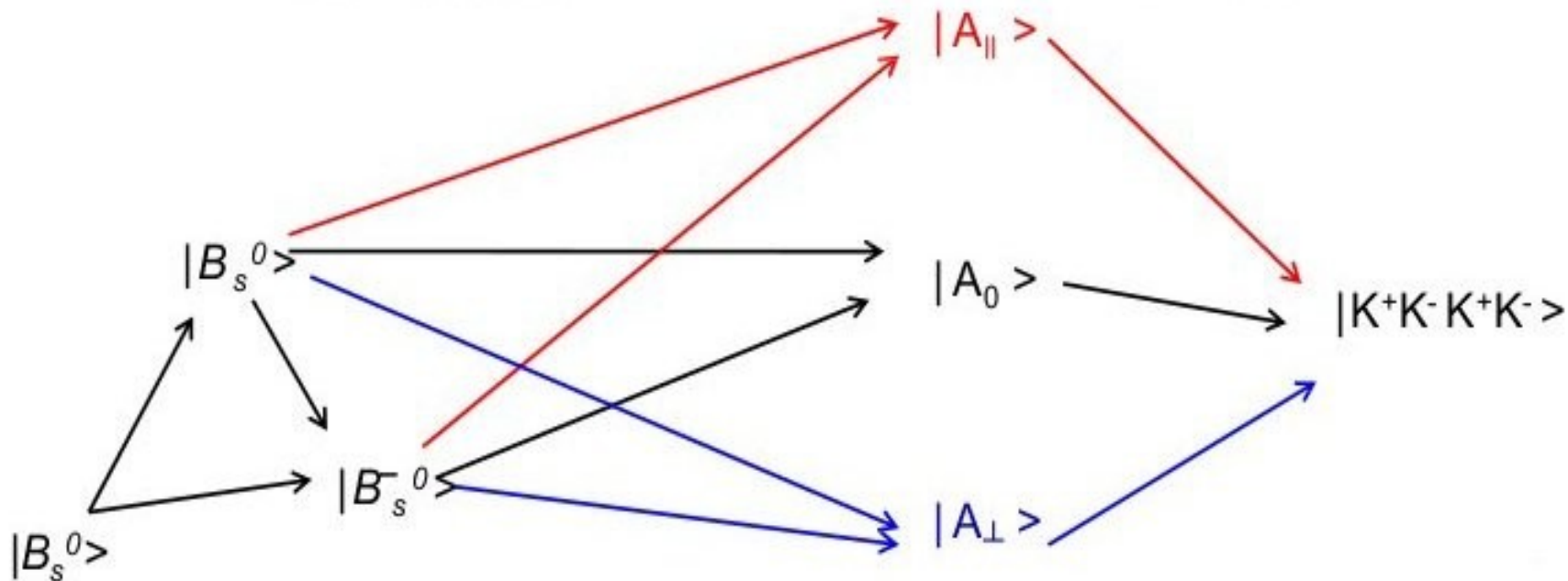


# MOTIVATION

In the transversity basis the amplitudes have defined CP value:

Corresponding decay amplitudes:  $A_0, A_{\parallel}, A_{\perp}$

- transverse ( $\perp$  perpendicular to each other)  $\rightarrow$  CP odd
- transverse ( $\parallel$  parallel to each other)  $\rightarrow$  CP even
- longitudinal (0)  $\rightarrow$  CP even



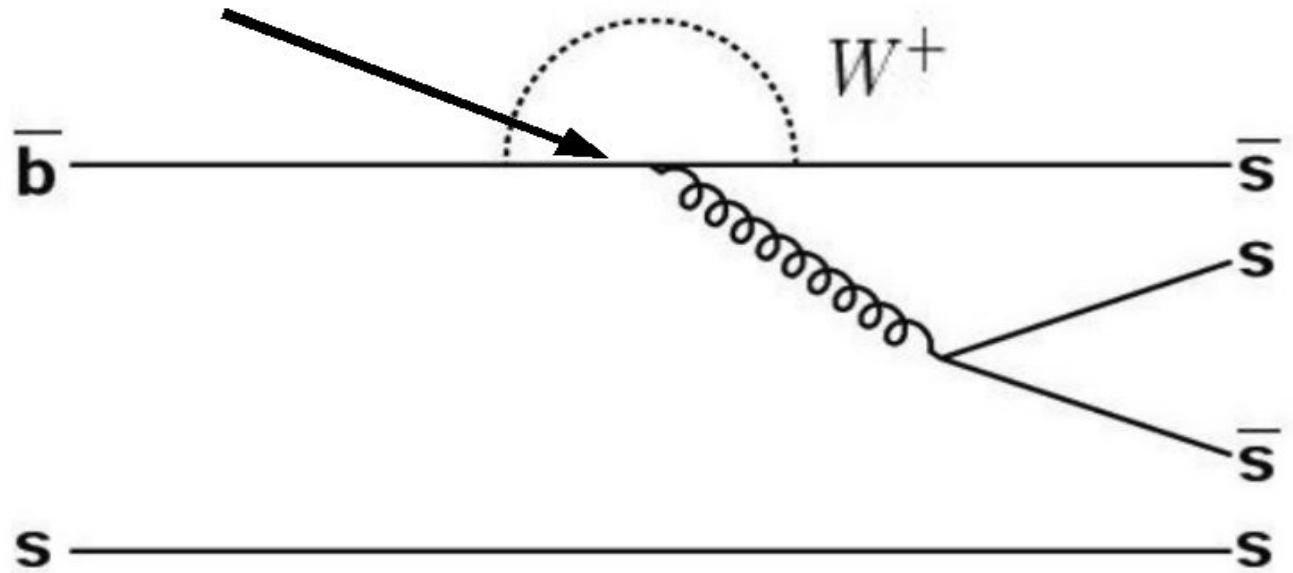


# Polarization puzzle

This result, known also as “Polarization Puzzle”, is more evident in the B meson decays involving penguin diagram decays. Different solutions were proposed for this “puzzle”, beyond and inside the Standard Model.

One solution beyond the SM could be New Physics inside the penguin loop:  
(Grossman [hep-nh/0311022](#), Yang et al [hep-nh/0411211](#))

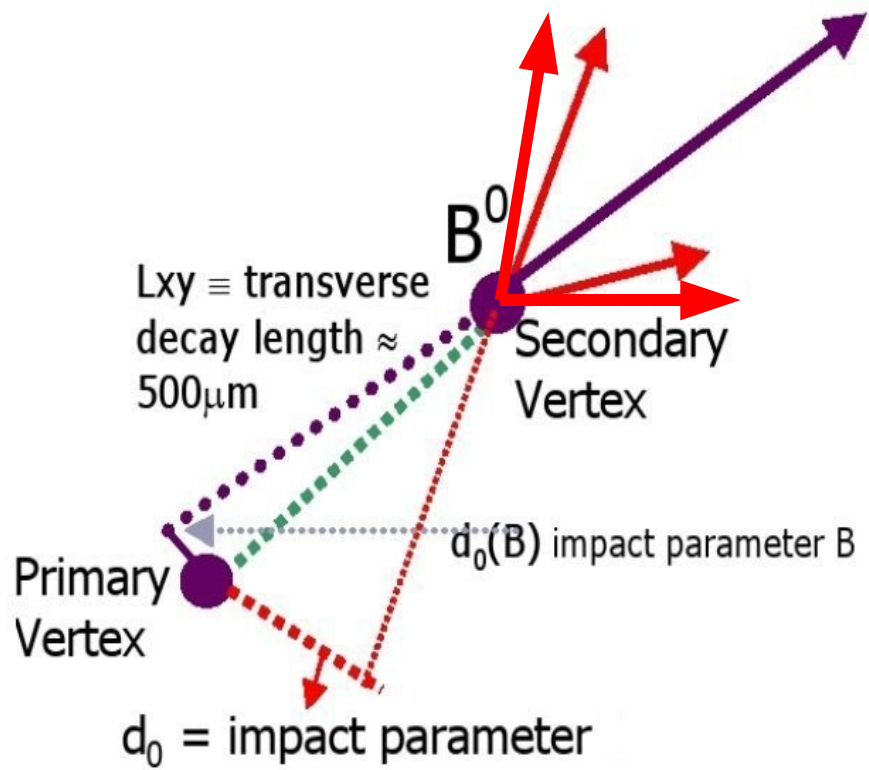
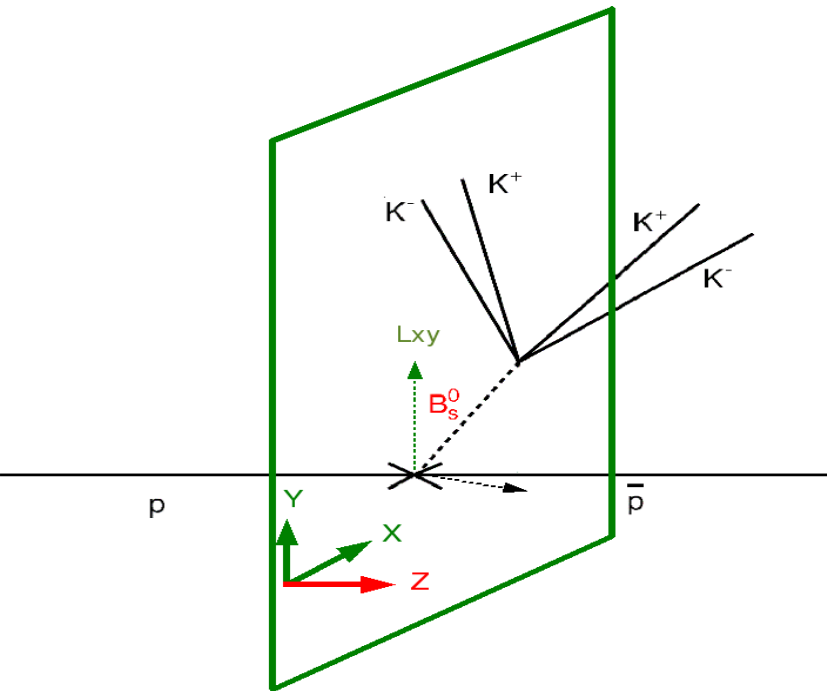
$\bar{u}, \bar{c}, \bar{t}, \dots$  **what else ?**





# Event signature

The events are characterized by 4 charged particle in the final state





## Optimization Selection procedure

Event selection is performed using the following variables

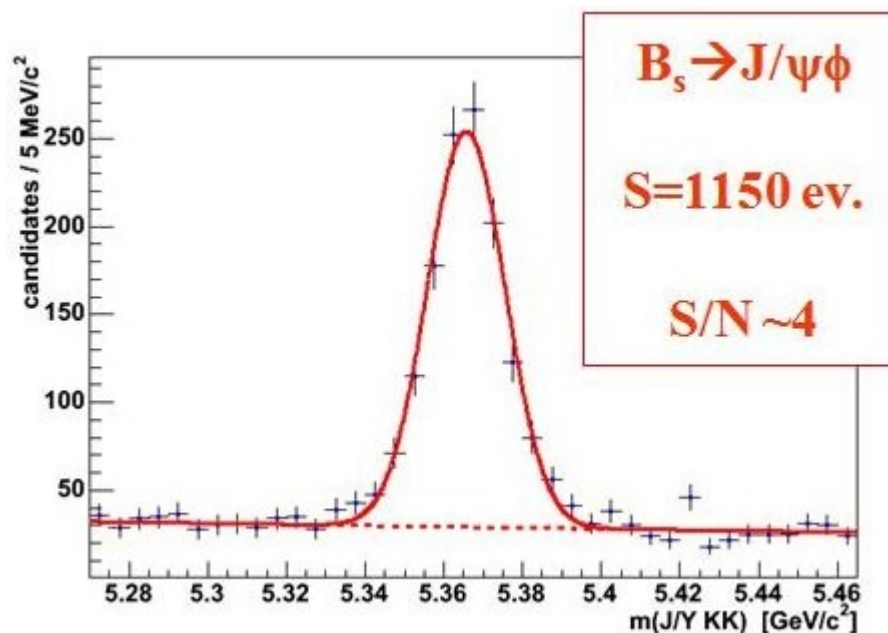
- transverse momentum of the  $B$  meson:  $p_T^B$ ;
- transverse decay length of the  $B$  meson:  $L_{xy}$ ;
- impact parameter of the  $B$  meson:  $d_0^B$ ;
- impact parameter of the more energetic  $\phi$ :  $d_{0max}^\phi$ ;
- impact parameter of the less energetic  $\phi$ :  $d_{0min}^\phi$ ;
- transverse momentum of the  $J/\psi$ :  $p_T^{J/\psi}$ ;
- transverse momentum of the less energetic kaon from  $\phi$  decay:  $p_T^K$ ;
- the bi-dimensional  $\chi^2$  of the primary vertex fit:  $\chi_{xy}^2$ .





## Selected events

Adding more trigger restriction:



These events are only in the trigger sample we used in this analysis (**TwoTrackTrigger**), not in the samples used for other **J/psi phi** measurements (**DI\_MUON Trigger**)



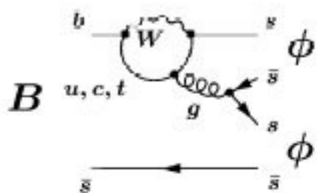
# Systematics Evaluation

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signal parametrization	3%	2%
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	$\Delta \varepsilon_{\phi\phi}/\varepsilon_{\phi\phi}$	$\Delta \varepsilon_{J/\psi\phi}/\varepsilon_{J/\psi\phi}$
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XFT particle dep.	4%	
$p_T$ reweight	0.9%	
	$\Delta \varepsilon_{\mu}/\varepsilon_{\mu}$	
$\eta$ parametrization & correlation	0.9%	

# Measurement of $BR(B_s^0 \rightarrow \phi\phi)$ in $2.9 \text{ fb}^{-1}$

## of Two Track Trigger data

G.Artoni, B. Di Ruzza, M.Dorigo, L.Ortolan, M.Rescigno, A.M. Zanetti

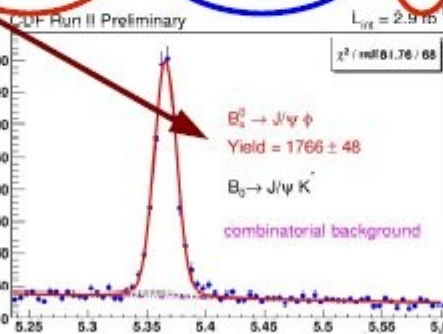
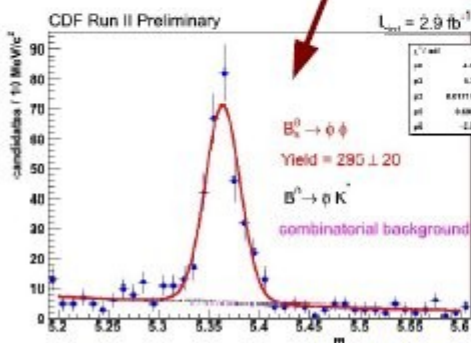


Experimental interest: improve precision w.r.t. to earlier pub.(180 pb-1).  
Expected sensitivity in future  $\sin(2\beta_s)$  measurement in  $B_s \rightarrow \phi\phi$ .

Theoretical interest: test recent QCD factorization calculations in  $B \rightarrow VV$ .

**WE MEASURE:**

$$\frac{BR(B_s \rightarrow \phi\phi)}{BR(B_s \rightarrow J/\psi\phi)} = \frac{N_{\phi\phi}}{N_{J/\psi\phi}} \cdot \frac{\epsilon(B_s \rightarrow J/\psi\phi)}{\epsilon(B_s \rightarrow \phi\phi)} \cdot \frac{BR(J/\psi \rightarrow \mu\mu)}{BR(\phi \rightarrow KK)} \cdot \epsilon_\mu$$



	BR [ $10^{-6}$ ]
Experiment	$14_{-5}^{+6}(\text{stat.}) \pm 6(\text{syst.})$
QCD Factorisation	$21.8_{-1.1-17.0}^{+1.1+30.4}$
QCD Factorisation	$19.5_{-1.0-8.0}^{+1.0+13.1}$
Naive Factorisation	9.05

- a: D. Acosta et al. P.R.L. 95, 031801 (2005)
- b: M. Beneke et al., hep-hex/0612290
- c: X. Li et al. P.R.L. D.68, 114015 (2003); D71 019902(2005); hep-hex/030936

**The result is:**  $\frac{BR(B_s \rightarrow \phi\phi)}{BR(B_s \rightarrow J/\psi\phi)} = [1.78 \pm 0.14^{stat} \pm 0.20^{syst}] \cdot 10^{-2}$

$$BR(B_s^0 \rightarrow \phi\phi) = [24.0 \pm 2.1(stat) \pm 2.7(syst) \pm 8.2(BR)] 10^{-6}$$

**Perspectives:** Measure polarization in this decay through the study of angular distributions