

# Intertwined Quantum Phase Transitions in the Zr Isotopes

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# Quantum Phase Transitions (QPTs)

- **Type I:** Single configuration (shape-phase transitions)

$$\hat{H} = (1 - \xi) \hat{H}_1 + \xi \hat{H}_2$$

neutron number 90 region: Nd-Sm-Gd

- **Type II:** Two configurations A , B (normal-intruder states)

$$\hat{H} = \begin{bmatrix} \hat{H}_A(\xi^A) & \hat{W}(\omega) \\ \hat{W}(\omega) & \hat{H}_B(\xi^B) \end{bmatrix}$$

nuclei near shell-closure: Cd, Pb-Hg

- Intertwined quantum phase transitions (IQPTs)

Type II QPT and Type I QPT coexist

configuration crossing accompanied by pronounced individual shape-evolutions

Microscopic approach: NRMF (Robledo ), RMF (Vretenar), MCSM (Togashi, Otsuka 2016)

Algebraic approach: THIS TALK

Zr isotopes: EXP (Kremer 2016; Ansari 2017; Paul 2017; Witt 2018; Singh 2018)

# Algebraic approach to shape coexistence near shell closure

- Multiparticle-multihole intruder excitations across shell gaps
- Interacting boson model with configuration mixing (IBM-CM) [Duval, Barrett, PLB 81]

0p-0h, 2p-2h, 4p-4h,...  $\rightarrow$   $[N] \oplus [N+2] \oplus [N+4] \dots$     normal  $\oplus$  intruder states

• **Hamiltonian**

$$\hat{H} = \begin{bmatrix} \hat{H}_A(\xi^A) & \hat{W}(\omega) \\ \hat{W}(\omega) & \hat{H}_B(\xi^B) \end{bmatrix} = \hat{H}_A^{(N)} + \hat{H}_B^{(N+2)} + \hat{W}$$

$$\hat{H}_A = \epsilon_d^{(A)} \hat{n}_d + \kappa^{(A)} \hat{Q}_\chi \cdot \hat{Q}_\chi$$

$$\hat{Q}_\chi = d^\dagger s + s^\dagger \tilde{d} + \chi (d^\dagger \times \tilde{d})^{(2)}$$

$$\hat{H}_B = \epsilon_d^{(B)} \hat{n}_d + \kappa^{(B)} \hat{Q}_\chi \cdot \hat{Q}_\chi + \kappa'^{(B)} \hat{L} \cdot \hat{L} + \Delta_p^{(B)}$$

$$\hat{T}(E2) = e^{(A)} \hat{Q}_\chi^{(N)} + e^{(B)} \hat{Q}_\chi^{(N+2)}$$

$$\hat{W} = \omega \left[ (d^\dagger \times d^\dagger)^{(0)} + (s^\dagger)^2 + \text{H.c.} \right]$$

• **Geometry**

$$E(\beta, \gamma) = \begin{pmatrix} E_N(\beta, \gamma) & W(\beta, \gamma) \\ W(\beta, \gamma) & E_{N+2}(\beta, \gamma) \end{pmatrix}$$

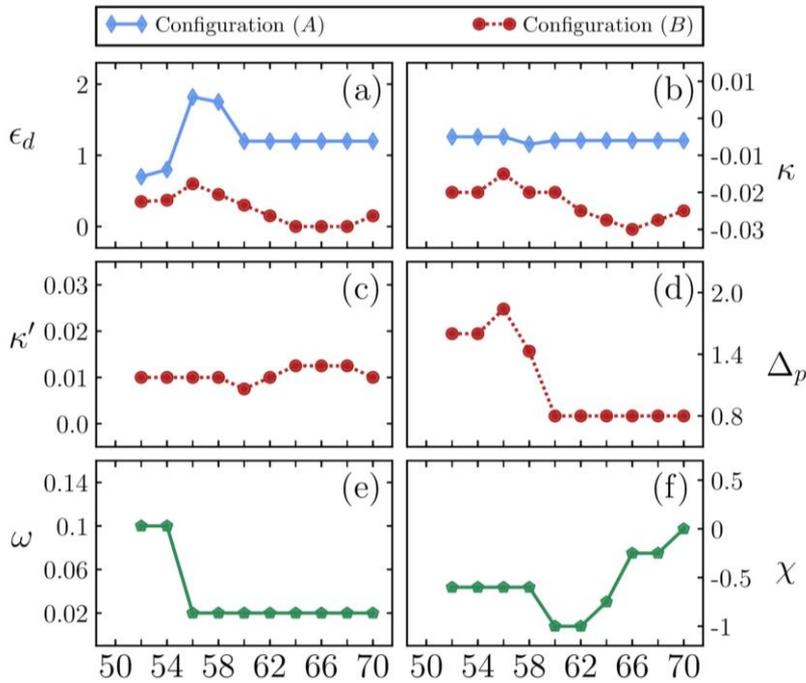
Matrix coherent states

$E_\pm(\beta, \gamma)$  Eigenpotentials

[Frank, Van Isacker, Vargas, PRC 2004;  
Frank, Van isacker, Iachello, PRC 2006]

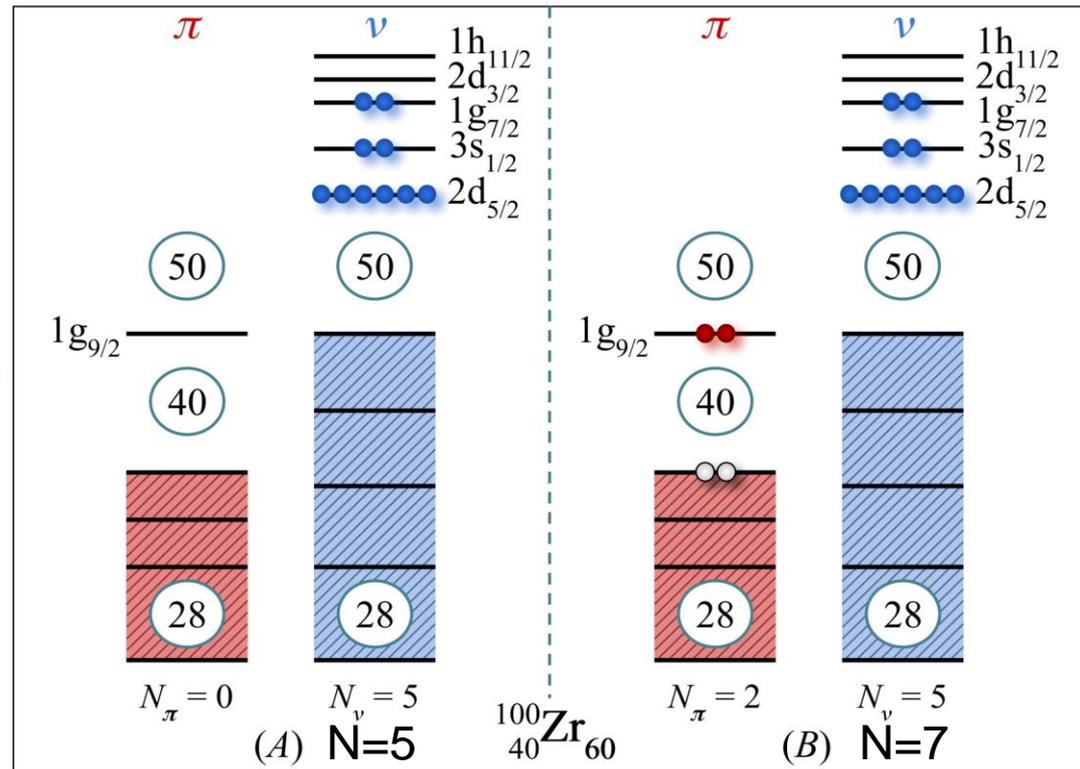
Applications: Mo, Po, Hg, Pb.. [Sambataro, Molnar, Garcia-Ramos, Heyde, Van Isacker, Nomura, ... ]

# IBM-CM in the Zr chain ( $z = 40, n = 52-70$ )



$$\hat{H} = \hat{H}_A^{(N)} + \hat{H}_B^{(N+2)} + \hat{W}$$

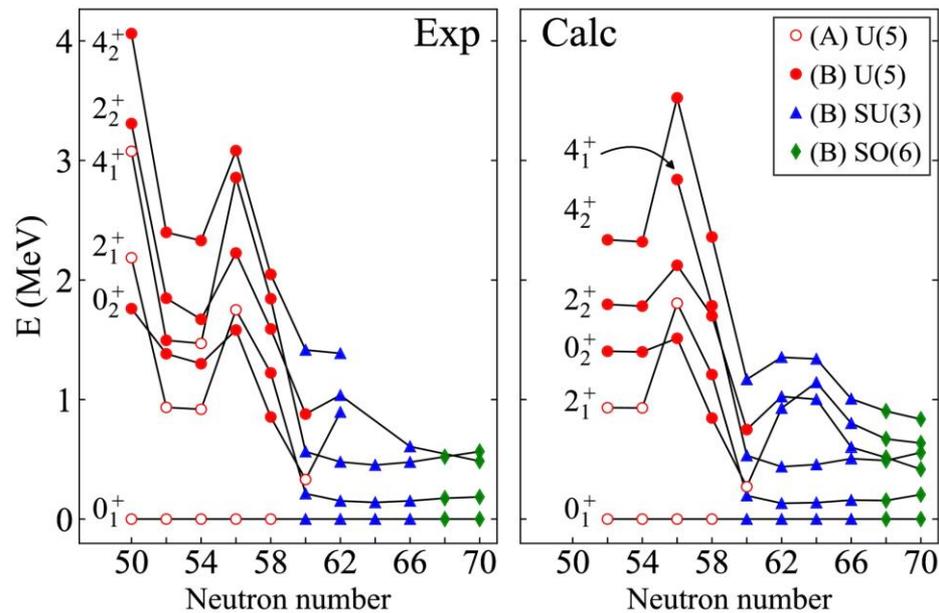
$$|\Psi; L\rangle = a |\Psi_A; [N], L\rangle + b |\Psi_B; [N+2], L\rangle$$



- $N = N_\pi + N_\nu$  ; Beyond mid-shell ( $n=66$ ), bosons counted as holes
- Parameters smooth function of  $n$ , or constant, similar to Mo (Sambataro 82) and Ge (Duval 83, Padilla 06)
- Some fluctuations due to sub-shell closure ( $n=56$ ;  $2d_{5/2}$  Talmi 65)
- Descending cross shell-gap excitations and onset of deformation ( $n>56$ ;  $V_{pn}$  between protons and neutrons occupying  $1g_{9/2}$  and  $1g_{7/2}$  Federman Pittel 79)

Normal (A) configuration  
Intruder (B) configuration

90-100Zr



n=50-56: config. (A) **spherical** (seniority-like)  $R_{4/2}^{(A)} \sim 1.6$   
config. (B) **weakly-deformed**  $R_{4/2}^{(B)} \sim 2.3$

From n=58: pronounced drop in energy for states of config. (B)

n=60: two configurations exchange role  $\Rightarrow$  **Type II QPT**  
config. (B) at critical point of **U(5)-SU(3) Type I QPT**  
first excited  $L=0^+_B$  low ( $E(0^+_3)$  in  $^{100}\text{Zr}$ ), similar to Sm-Gd

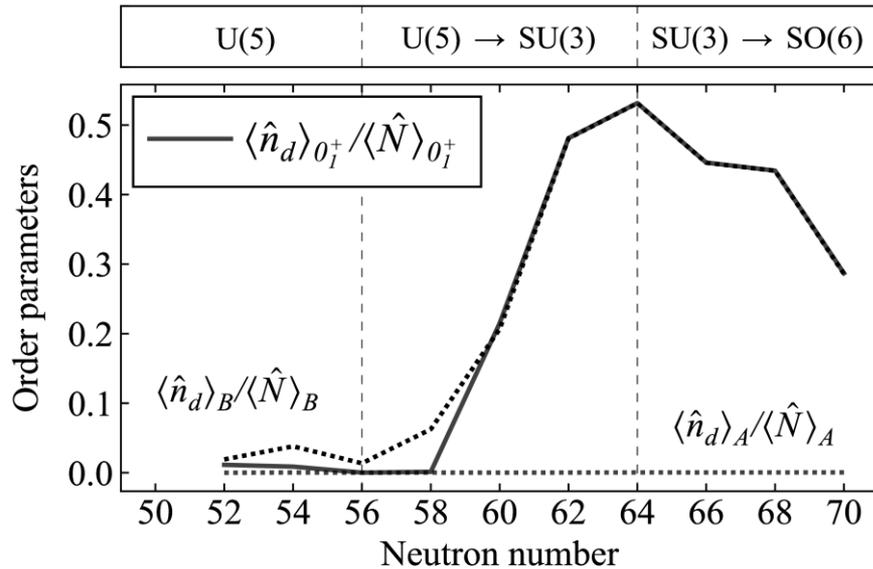
n>60: intruder config. (B) strongly **deformed [SU(3)]**  
 $^{104}\text{Zr}$ :  $E(2_1) = 139.3 \text{ keV}$ ,  $R_{4/2}^{(B)} = 3.24$

n=66: g.s. becomes  **$\gamma$ -unstable** (or triaxial) **SU(3)  $\rightarrow$  SO(6) crossover**

$^{106}\text{Zr}$ :  $E(2_2) = 607$ ,  $E(4_1) = 476.5 \text{ keV}$

$^{110}\text{Zr}$ :  $E(4_1) = 565$ ,  $E(2_2) = 485 \text{ keV}$  (Paul PRL 2017) **SO(6) signature**

# Order parameters



Config. (A): spherical for all n

Config. (B): weakly-deformed for n=52-58

clear jump between n=58 & 60 from (A) to (B)  
 $\Rightarrow$  1<sup>st</sup> order **Type II QPT**

further increase at n=60-64  
 $\Rightarrow$  U(5)-SU(3) **Type I QPT**

decrease at n=66  
 $\Rightarrow$  SU(3) to SO(6) crossover & shift of particle to hole bosons at mid-shell

$$|\Psi; L = 0_1^+\rangle = a |\Psi_A; [N], L = 0_1^+\rangle + b |\Psi_B; [N + 2], L = 0_1^+\rangle$$

$$\langle \hat{N} \rangle_A = N, \quad \langle \hat{N} \rangle_B = N + 2$$

$$\langle \hat{n}_d \rangle_A, \quad \langle \hat{n}_d \rangle_B$$

shapes of configurations A, B

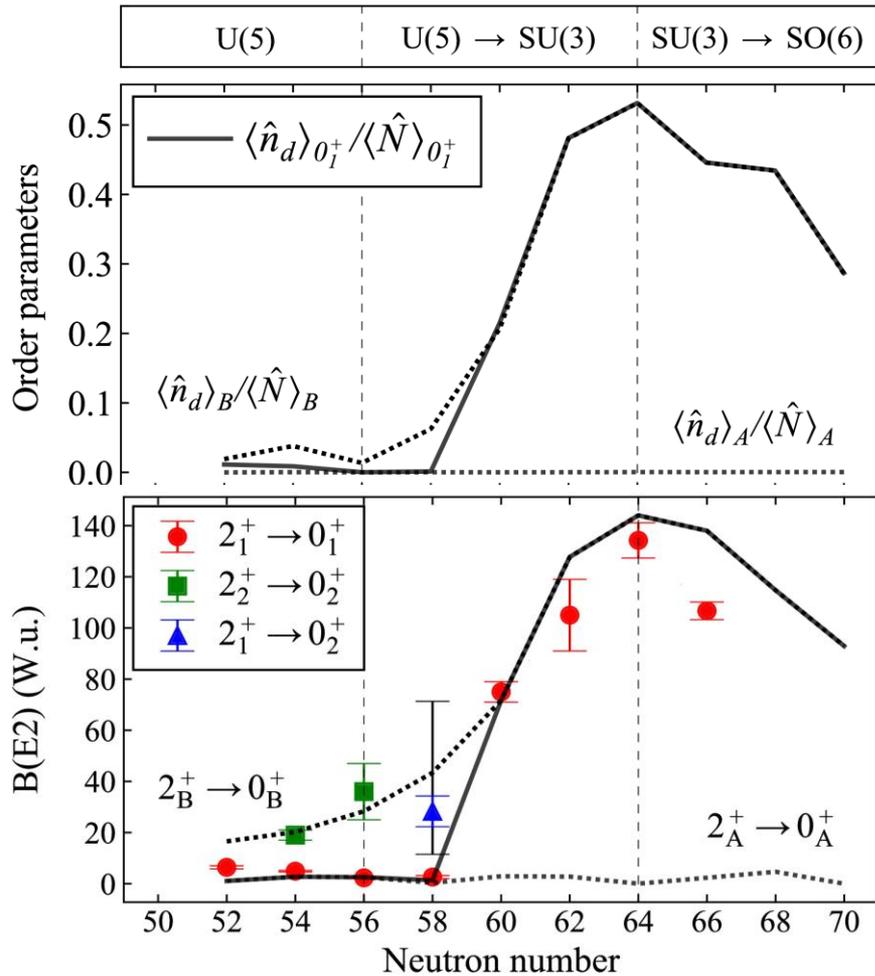
$$\langle \hat{N} \rangle_{0_1^+} = a^2 N + b^2 (N + 2)$$

$$\langle \hat{n}_d \rangle_{0_1^+} = a^2 \langle \hat{n}_d \rangle_A + b^2 \langle \hat{n}_d \rangle_B$$

normal-intruder mixing

Coexisting **Type I QPT** (shape-evolution) and **Type II QPT** (configuration crossing)

# Order parameters and B(E2)'s



Config. (A): spherical for all n

Config. (B): weakly-deformed for n=52-58

clear jump between n=58 & 60 from (A) to (B)  
 $\Rightarrow$  1<sup>st</sup> order **Type II QPT**

further increase at n=60-64  
 $\Rightarrow$  U(5)-SU(3) **Type I QPT**

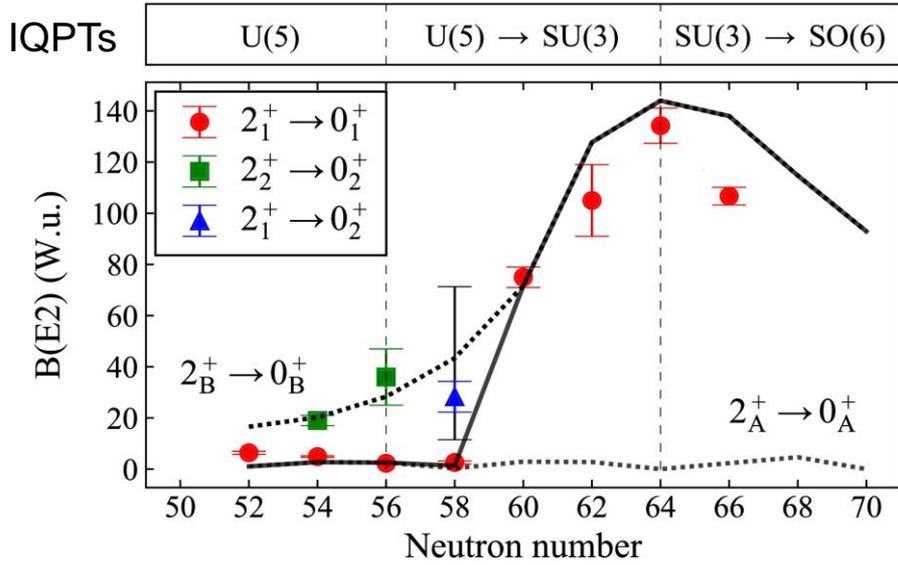
decrease at n=66  
 $\Rightarrow$  SU(3) to SO(6) crossover & shift of particle to hole bosons at mid-shell

Coexisting **Type I QPT** and **Type II QPT**  $\Rightarrow$  **IQPTs**

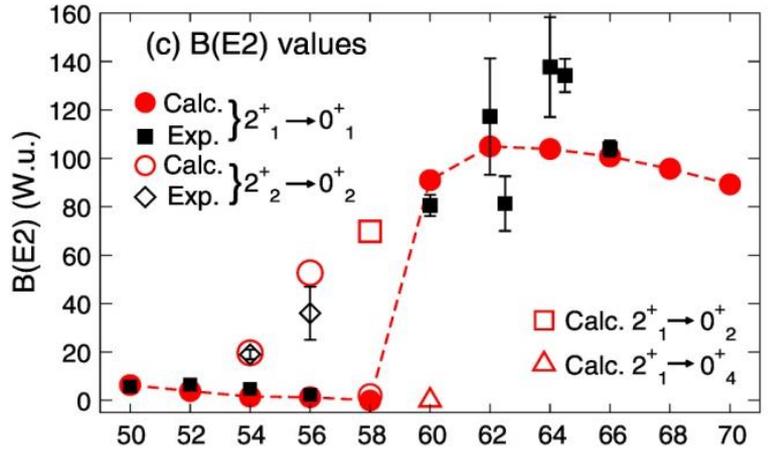
$$\hat{T}(E2) = e^{(A)} \hat{Q}_\chi^{(N)} + e^{(B)} \hat{Q}_\chi^{(N+2)}$$

$$\hat{Q}_\chi = d^\dagger s + s^\dagger \tilde{d} + \chi (d^\dagger \times \tilde{d})^{(2)}$$

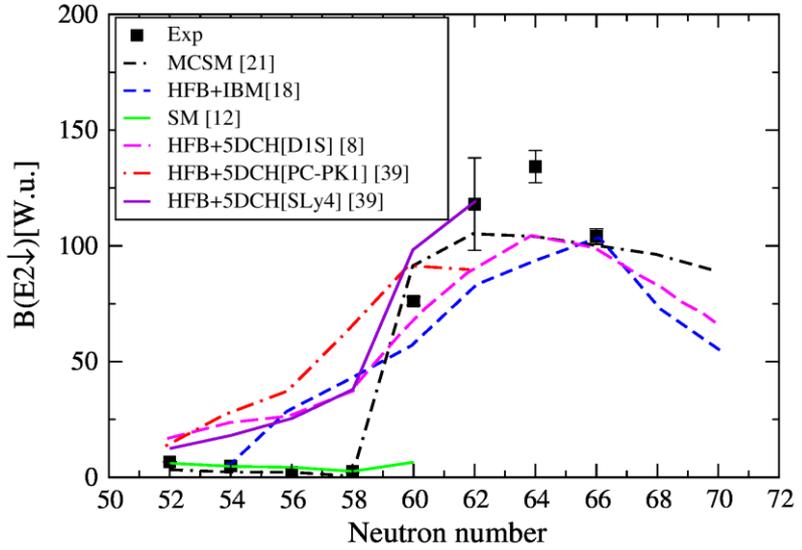
# B(E2)'s in the Zr chain



## MCSM (Togashi et al. PRL 2016)



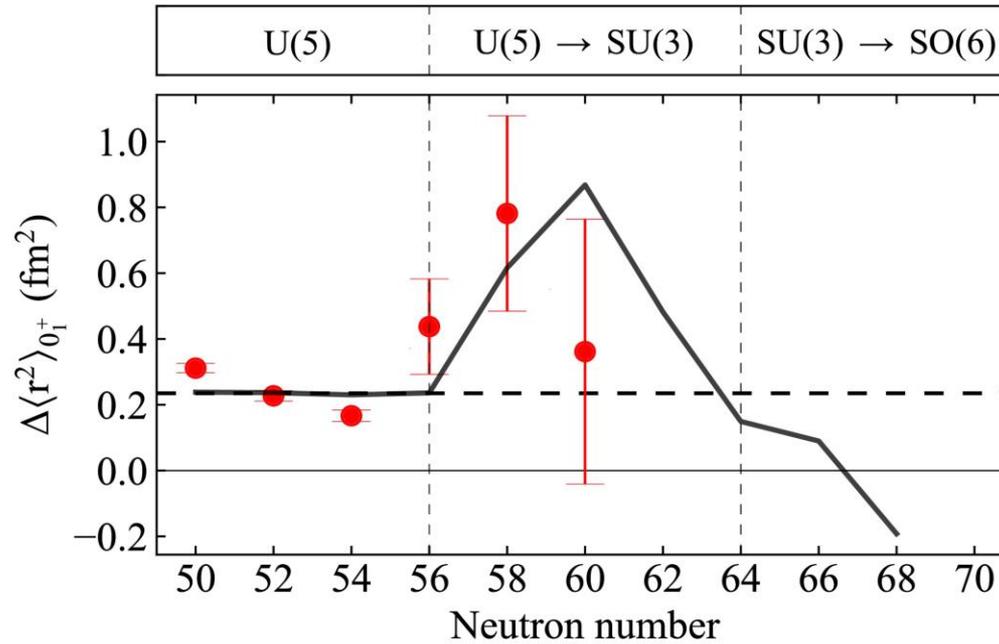
## (Singh et al. PRL 2016)



Both MCSM & IQPTs: large jump in  $B(E2; 2_1 \rightarrow 0_1)$  between  $^{98}\text{Zr}$  and  $^{100}\text{Zr}$   
 $\Rightarrow$  1<sup>st</sup> order QPT

Mean-field based calculations:  
 smooth out the QPT behavior,  
 show no jump at the critical point

## Isotope shift in the Zr chain

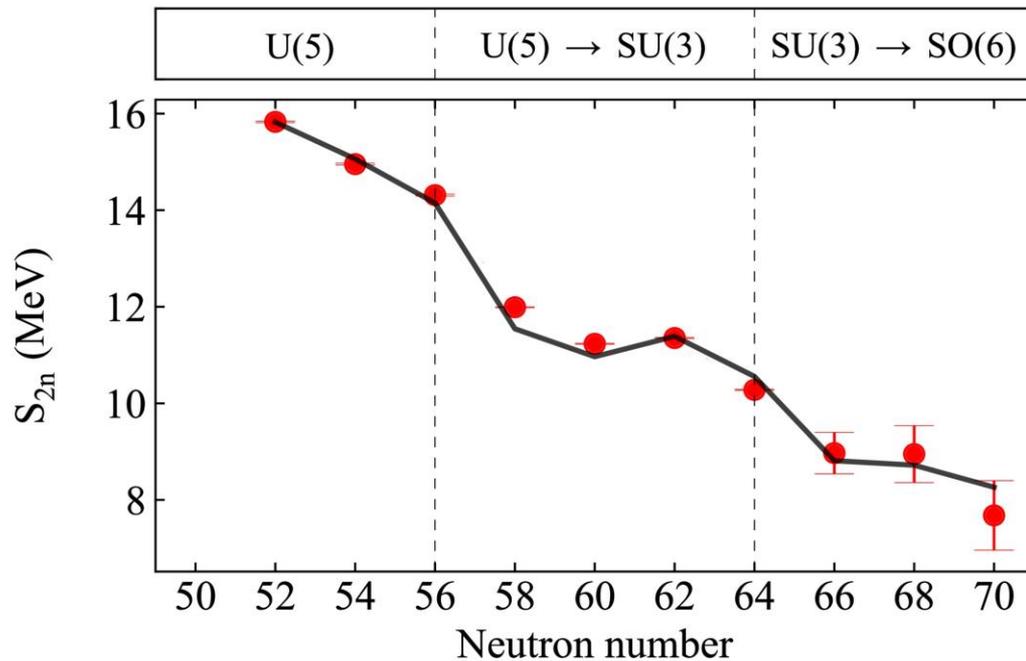


$$\Delta \langle \hat{r}^2 \rangle_{0_1^+} = \langle \hat{r}^2 \rangle_{0_1^+; A+2} - \langle \hat{r}^2 \rangle_{0_1^+; A}$$

$$\langle \hat{r}^2 \rangle = r_c^2 + \alpha N_v + \eta \left[ \langle \hat{n}_d^{(N)} \rangle + \langle \hat{n}_d^{(N+2)} \rangle \right]$$

$\Delta\langle r^2 \rangle$  increases at the transition point and decreases

## Two-neutron separation energies in the Zr chain



$$S_{2n} = -\tilde{A} - \tilde{B}N_v \pm S_{2n}^{\text{def}} - \Delta_n$$

$$S_{2n}^{\text{def}} \quad \langle H \rangle \text{ in g.s.}$$

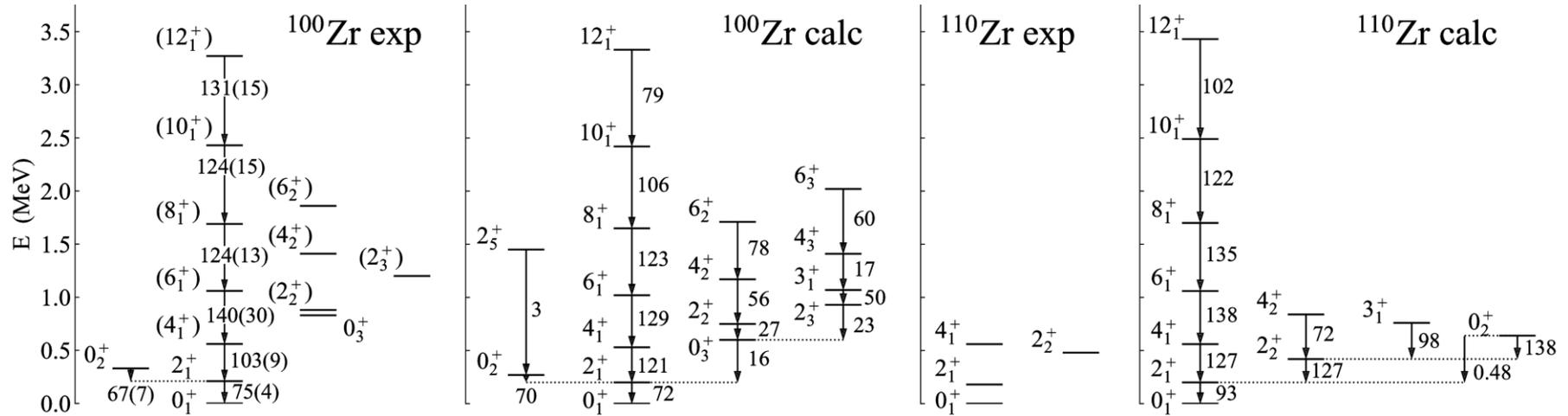
n=52-56:  $S_{2n}$  straight line, g.s. spherical (seniority-like)

After 56:  $S_{2n}$  first goes down (sub-shell closure) then flattens  $\Rightarrow$  1<sup>st</sup> order QPT (as in Sm)

After 62:  $S_{2n}$  goes down (increasing deformation)

n=66-70:  $S_{2n}$  flattens (SU(3) $\rightarrow$  SO(6) crossover)

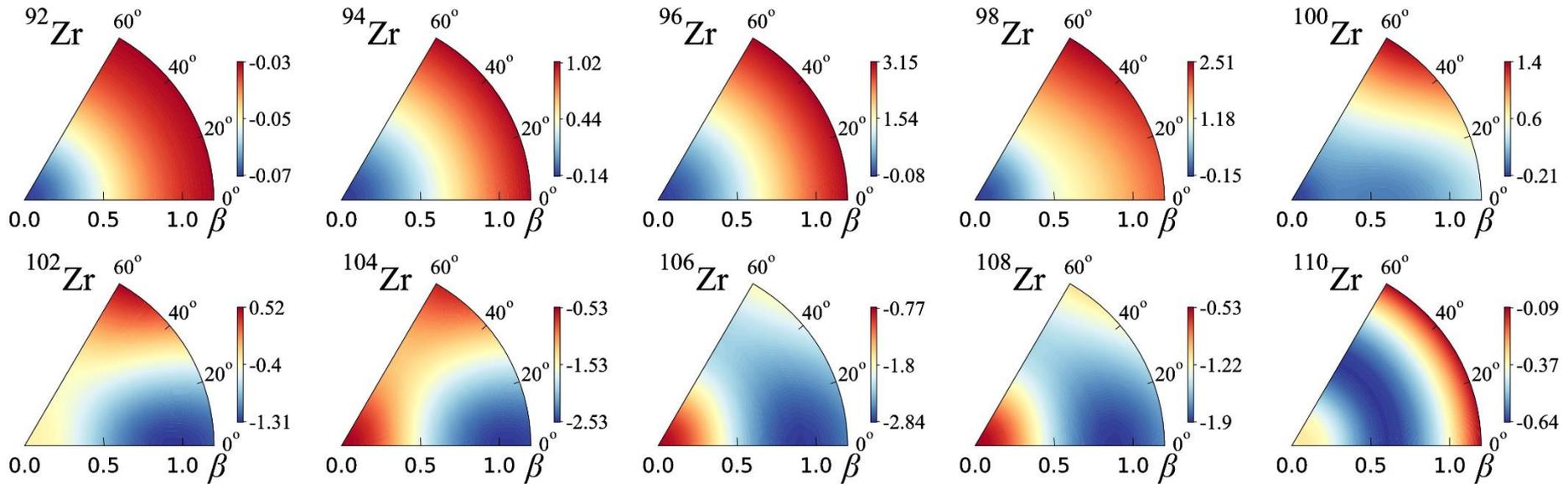
# Quantum analysis



$^{100}\text{Zr}$ : near the **critical point** of both Type I and Type II QPT  
 g.s. band, config. (B): **X(5)** symmetry  
 spherical config. (A): excited band  $0_2$

$^{110}\text{Zr}$ : **SO(6)** symmetry

# Classical analysis



$^{92-98}\text{Zr}$ : spherical

$^{100}\text{Zr}$ : flat-bottomed potential

$^{102-104}\text{Zr}$ : axially deformed

$^{106-110}\text{Zr}$ :  $\gamma$ -unstable

## Concluding remarks

- Intertwined Quantum Phase Transitions (IQPTs):

QPT involving a crossing of two configurations (Type II) each of which is in itself undergoing a QPT (Type I)

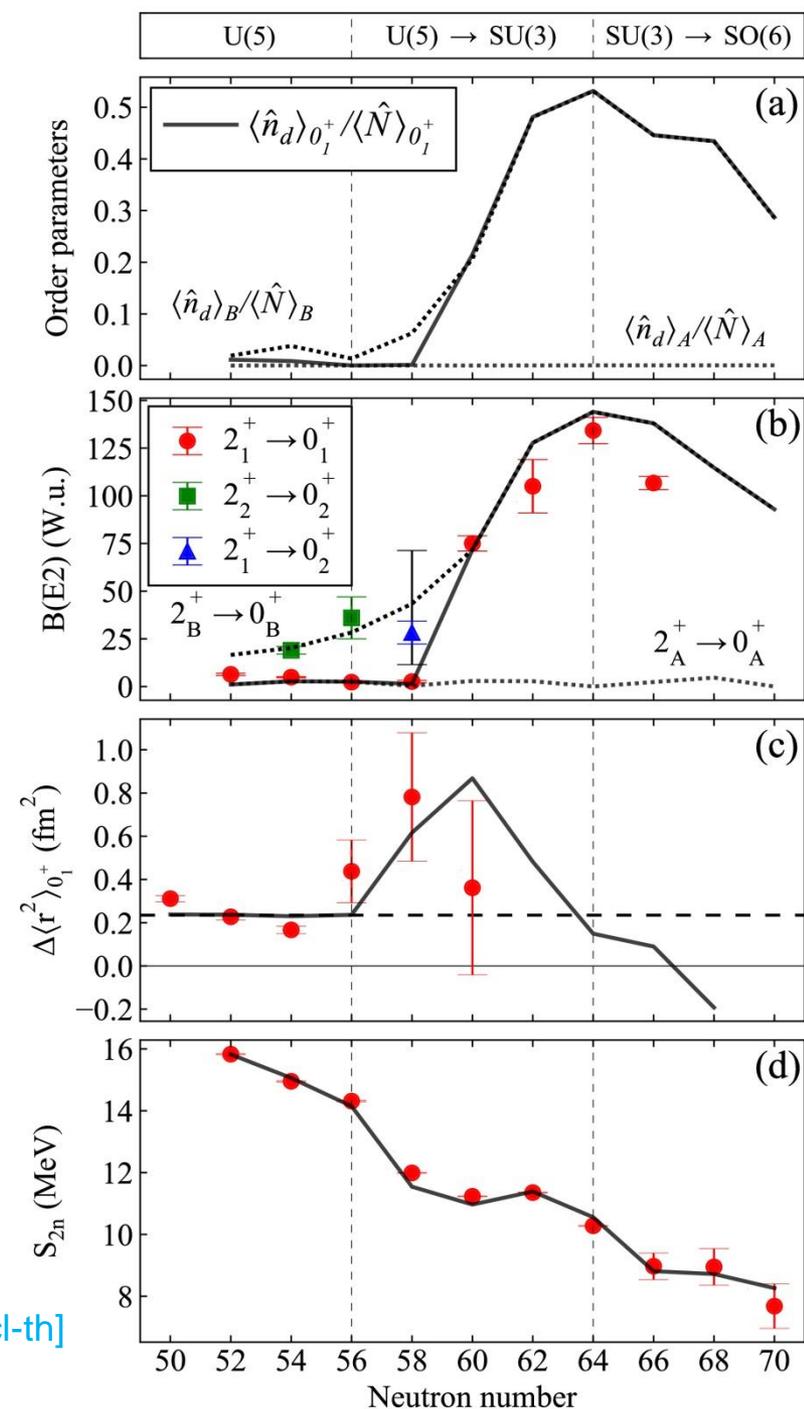
- Comprehensive analysis in the IBM-CM framework of spectra and other observables ( $B(E2)$ ,  $\Delta\langle r^2 \rangle$ ,  $S_{2n}$ ) in the Zr chain ( $n=52-70$ )

- Two configurations, one spherical (A) and the other (B) undergoing first a QPT U(5)-SU(3) and then a crossover SU(3)-SO(6), coexisting with a normal-intruder configuration- crossing

IQPTs appear to be manifested empirically in the Zr chain

- Initial indications for similar features in the Sr chain

N. Gavrielov, A. Leviatan and F. Iachello, arXiv:1904.09919 [nucl-th]



Thank you