

Forbidden transitions in nuclear weak processes relevant to neutrino detection, nucleosynthesis and evolution of stars

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GT transitions in nuclear weak rates at stellar environments

e-capture and β -decay rates in sd- and pf-shell nuclei are updated

Nuclear URCA process in O-Ne-Mg core and cooling of stars

Synthesis of iron-group elements in Type Ia SNe

ν -nucleus reaction cross sections updated for ^{12}C , ^{13}C , ^{40}Ar , ^{56}Fe , ^{56}Ni ; updated

ν -process nucleosynthesis, ν detection

FF (first-forbidden) transitions in β -decay half-lives in N=126 isotones

r-process nucleosynthesis in CCSNe and binary NSM (neutron star mergers)

1. Spin-dipole strength in ^{16}O and ν -induced reactions on ^{16}O

- Spin-dipole strengths in ^{16}O and ν - ^{16}O reaction cross sections with new shell-model Hamiltonians
- Synthesis of light elements in SNe based on the cross sections
- Supernova ν detection by ν - ^{16}O reactions
- MSW ν -oscillation effects

Dependence of ν - ^{16}O cross sections on ν mass hierarchy

Suzuki, Chiba, Yoshida, Takahashi, and Umeda, Phys. Rev. C98, 034613 (2018)

Nakazato, Suzuki, and Sakuda, PTEP 2018, 123E02 (2018)

2. Electron-capture and β -decay rates of sd-shell nuclei

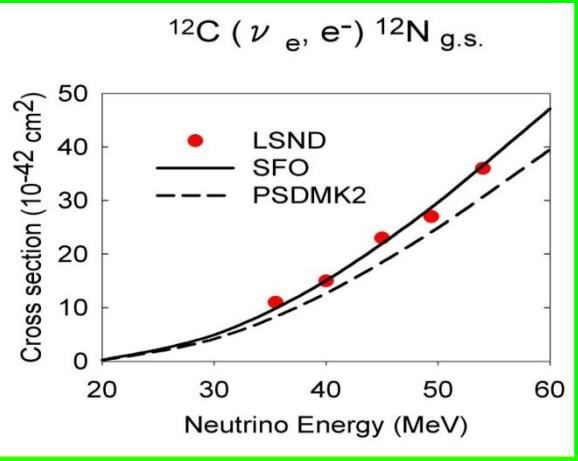
Electron-capture rates for a **2nd-forbidden transition** in ^{20}Ne

▪ ν -induced reactions on ^{16}O

SFO: p-sd shell

$\text{CK}(\text{p}) + \text{MK}(\text{p-sd}) + \text{G}(\text{sd})$

Monopole ($\text{p}_{1/2}-\text{p}_{3/2}$; $T=0$) enhanced



$$g_A^{\text{eff}}/g_A = 0.95$$

$B(\text{GT: } ^{12}\text{C})_{\text{cal}}$
= experiment

Magnetic moments
Of p-shell nuclei are
well reproduced

$^{12}\text{C} (\nu, \nu'), (\nu_e, e^-)$ GT+ SD

SFO reproduces exclusive and inclusive
DAR cross sections

▪ Modification of SFO \rightarrow SFO-tls

Full inclusion of tensor force

p-sd: tensor- $\rightarrow \pi+\rho$

LS $\rightarrow \sigma+\rho+\omega$

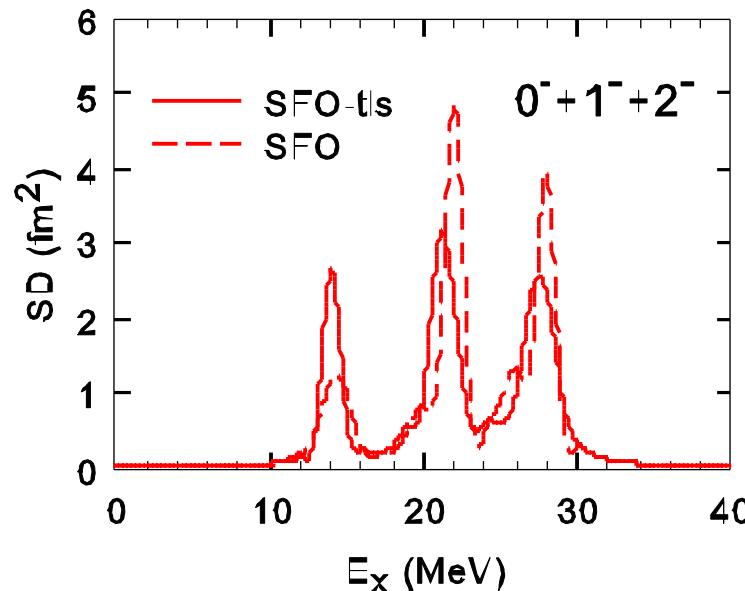
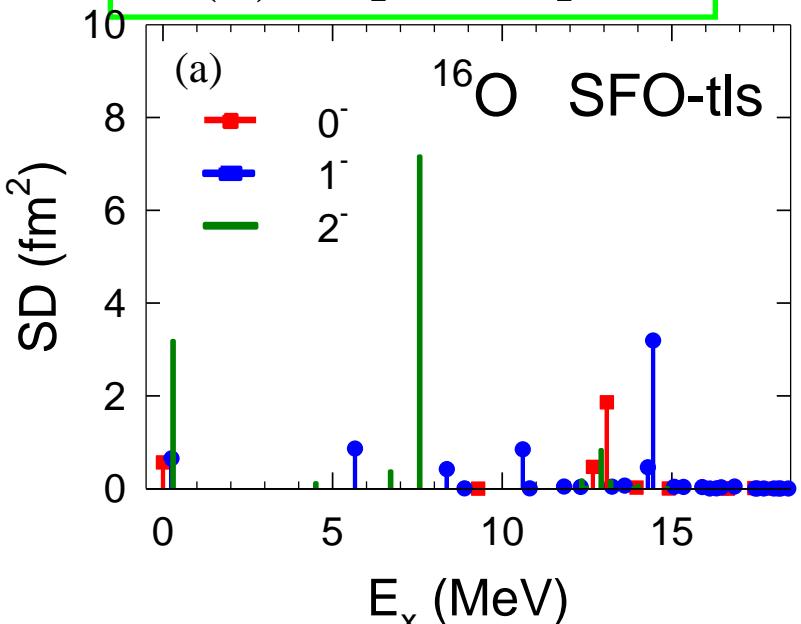
$$V = V_C + V_T + V_{LS}$$

$$V_T = V_\pi + V_\rho$$

$$V_{LS} = V_{\sigma+\omega+\rho}$$

Spin-dipole strength in ^{16}O

$$O(\lambda) = r[Y^1 \times \sigma]^\lambda t_-$$



Spin-dipole sum

$$B(SD\lambda)_{\mp} = \frac{1}{2J_i + 1} \sum_f |\langle f \parallel S_{\mp}^{\lambda} \parallel i \rangle|^2$$

$$S_{\mp,\mu}^{\lambda} = r [Y^1 \times \vec{\sigma}]_{\mu}^{\lambda} t_{\mp}$$

NEWS-rule: $S_-^{\lambda} - S_+^{\lambda} = \langle 0 \mid [\hat{S}_-^{\lambda}, \hat{S}_+^{\lambda}] \mid 0 \rangle = \frac{2\lambda + 1}{4\pi} (N\langle r^2 \rangle_n - Z\langle r^2 \rangle_p)$

For ^{16}O ; N=Z

$$S_{\lambda}(SD) = \sum_{\mu} |\langle \lambda, \mu \mid S_{-, \mu}^{\lambda} \mid 0 \rangle|^2 = \frac{1}{2} \left\{ \begin{array}{l} \frac{3}{4\pi} 4b^2 = 2.99 \text{fm}^2 \quad \lambda^{\pi} = 0^- \\ \frac{3}{4\pi} 12b^2 = 8.98 \text{fm}^2 \quad \lambda^{\pi} = 1^- \\ \frac{3}{4\pi} 20b^2 = 14.96 \text{fm}^2 \quad \lambda^{\pi} = 2^- \end{array} \right. \quad \begin{array}{l} p \rightarrow sd \\ \propto 2\lambda + 1 \end{array}$$

Energy-weighted sum

$$EWS_{\pm}^{\lambda} = \sum_{\mu} |\langle \lambda, \mu \mid S_{\pm, \mu}^{\lambda} \mid 0 \rangle|^2 (E_{\lambda} - E_0),$$

$$\begin{aligned} EWS^{\lambda} &= EWS_-^{\lambda} + EWS_+^{\lambda} \\ &= \frac{1}{2} \langle 0 \mid [S_-^{\lambda\dagger}, [H, S_-^{\lambda}]] + [[S_+^{\lambda\dagger}, H], S_+^{\lambda}] \mid 0 \rangle. \end{aligned}$$

kinetic energy term (K) for $H = \frac{\vec{p}^2}{2m}$

$$EWS_K^\lambda = \frac{3}{4\pi}(2\lambda+1)\frac{\hbar^2}{2m}A[1 + \frac{f_\lambda}{3A} <0| \sum_i \vec{\sigma}_i \cdot \vec{\ell}_i |0>]$$

: $f_\lambda = 2, 1$ and -1 for $\lambda^\pi = 0^-, 1^-$ and 2^- , respectively.

One-body spin-orbit potential term

$$V_{LS} = -\xi \sum_i \vec{\ell}_i \cdot \vec{\sigma}_i,$$

$$EWS_{LS}^\lambda = \frac{3}{4\pi}(2\lambda+1)\frac{f_\lambda}{3}\xi <0| \sum_i (r_i^2 + g_\lambda r_i^2 \vec{\ell}_i \cdot \vec{\sigma}_i) |0>$$

$g_\lambda = 1$ for $\lambda^\pi = 0^-$, 1^- and $g_\lambda = -7/5$ for $\lambda^\pi = 2^-$.

For N=Z, $EWS_-^\lambda = EWS_+^\lambda$, and $EWS_-^2/5 < EWS_-^1/3 < EWS_-^0$

EWS_-^λ	0^-	1^-	2^-	
K+LS	56.4	144.1	155.9	MeV·fm ²
SFO-tls (/K+LS)	73.0 (1.29)	173.2 (1.20)	246.5 (1.58)	
SFO (/K+LS)	76.1 (1.35)	175.0 (1.21)	258.2 (1.66)	

$E_\lambda = EWS_-^\lambda / NEWS_-^\lambda$	0^-	1^-	2^-	
SFO-tls	24.5	25.1	20.1	MeV
SFO	25.8	25.2	21.0	

Tensor interaction:
attractive for 0^- and 2^-
repulsive for 1^-

μ -capture rate on ^{16}O and the quenching factor

The muon capture rate for $^{16}\text{O} (\mu, \nu_\mu) ^{16}\text{N}$ from the 1s Bohr atomic orbit

$$\omega_\mu = \frac{2G^2}{1 + \nu/M_T} |\phi_{1s}|^2 \frac{1}{2J_i + 1} \left(\sum_{J=0}^{\infty} |< J_f \parallel M_J - L_J \parallel J_i >|^2 + |< J_f \parallel T_J^{el} - T_J^{mag} \parallel J_i >|^2 \right),$$

$$|\phi_{1s}|^2 = \frac{R}{\pi} \left(\frac{m_\mu M_T}{m_\mu + M_T} Z \alpha \right)^3 \quad R = 0.79$$

Induced pseudo-scalar current $F_P(q_\mu^2) = \frac{2M_N}{q_\mu^2 + m_\pi^2} F_A(q_\mu^2)$

$f = g_A^{eff}/g_A = 0.95$

SFO $10.21 \times 10^4 \text{ s}^{-1}$ (SFO/exp = 0.995)

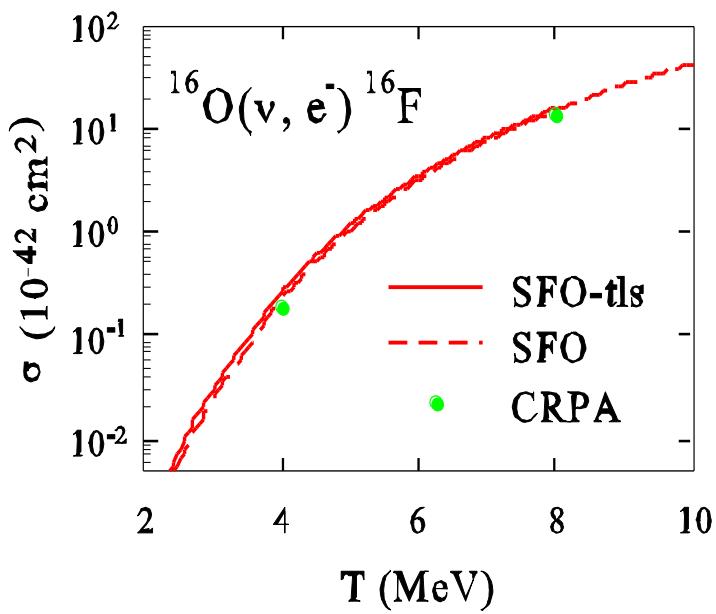
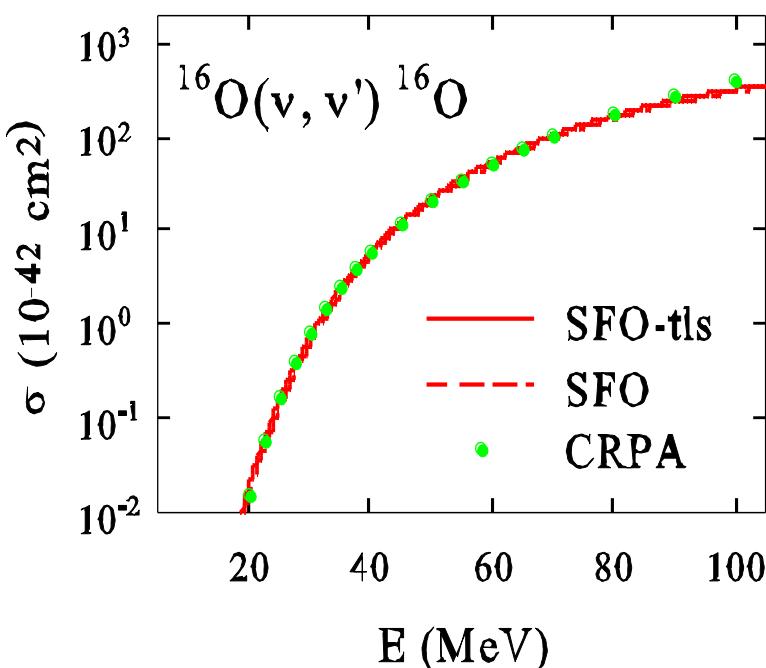
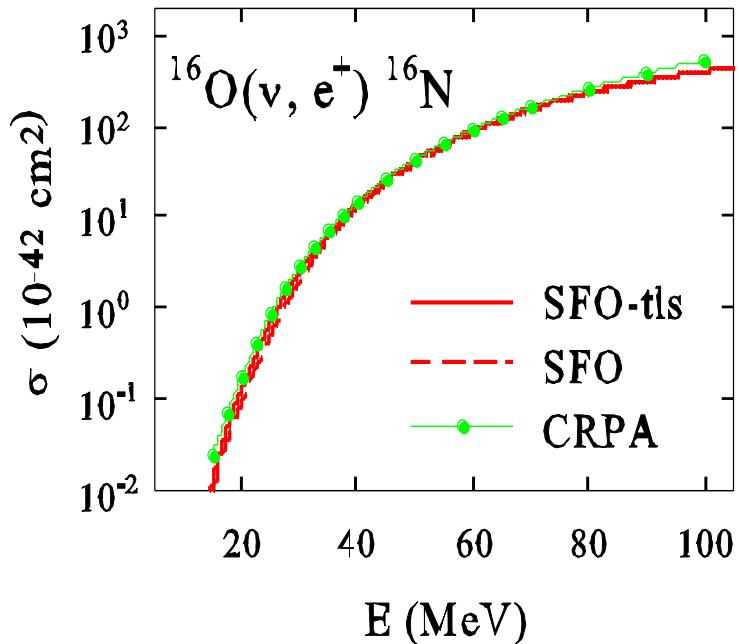
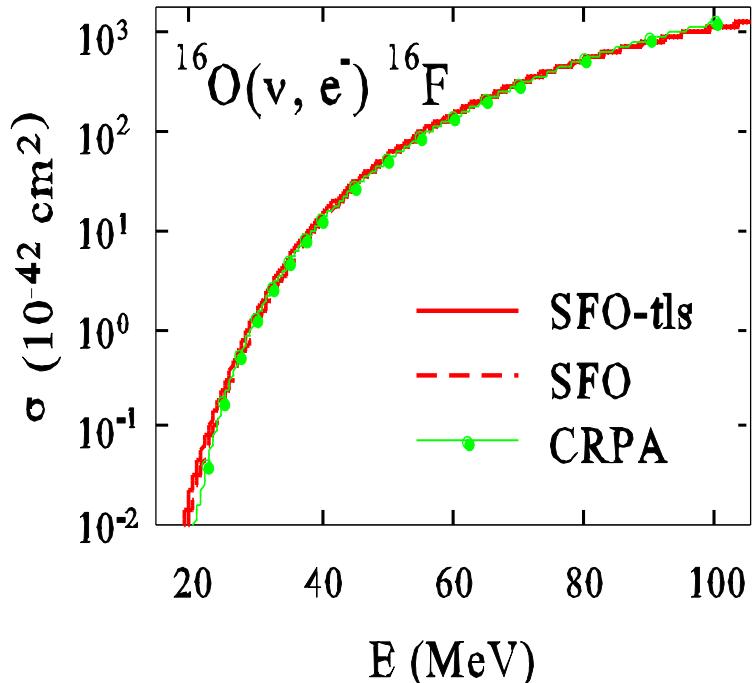
SFO-tls, $11.20 \times 10^4 \text{ s}^{-1}$ (SFO-tls/exp = 1.092)

Exp. $10.26 \times 10^4 \text{ s}^{-1}$

PCAC

Goldberger-Treiman

$$-2M_N F_A = \sqrt{2} g_\pi F_\pi$$



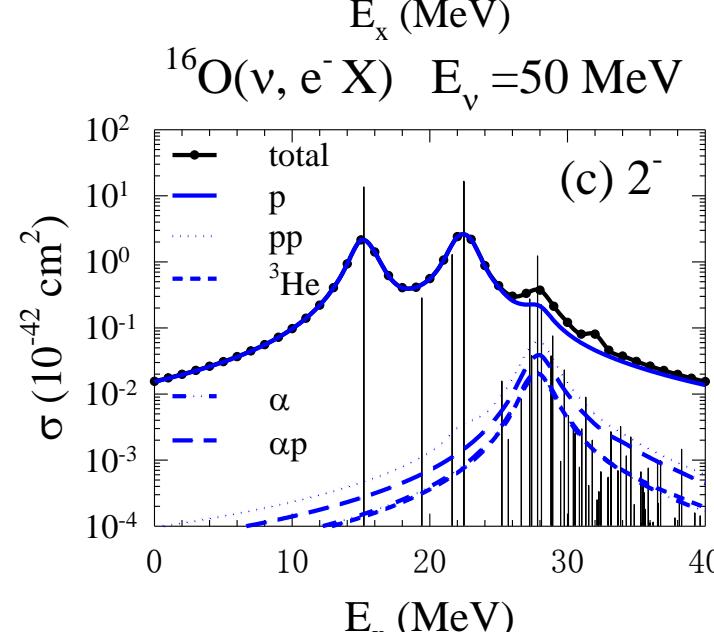
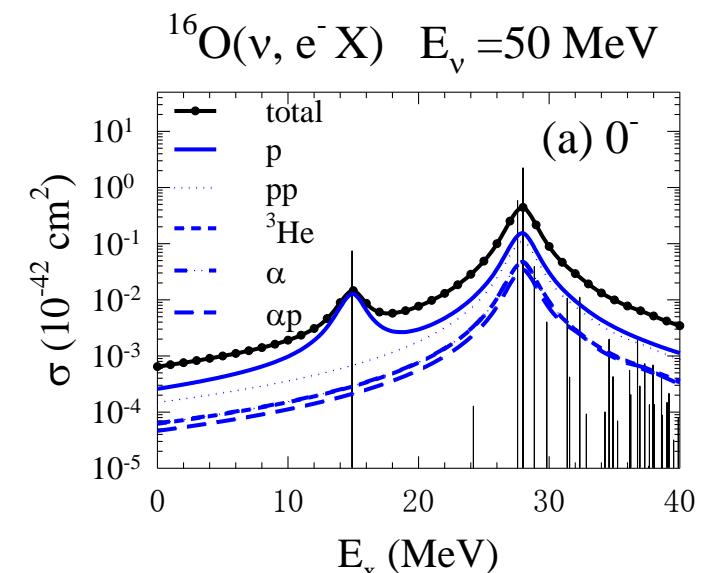
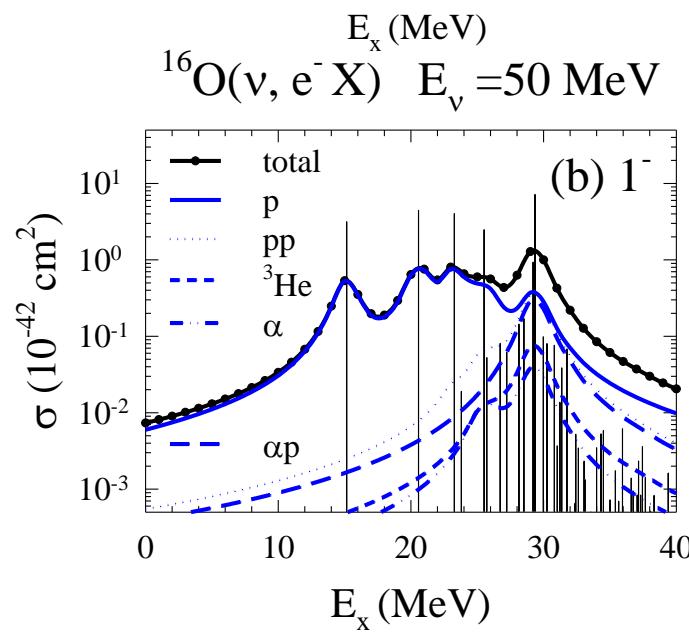
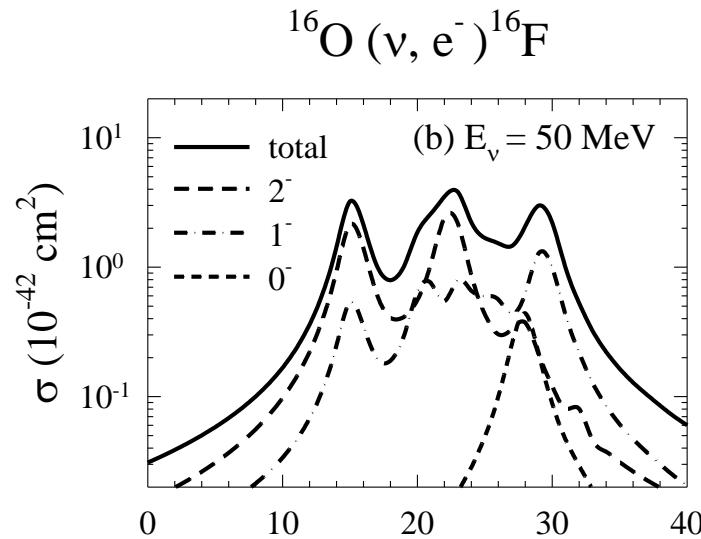
CRPA: Kolbe, Langanke & Vogel, PR D66 (2002)

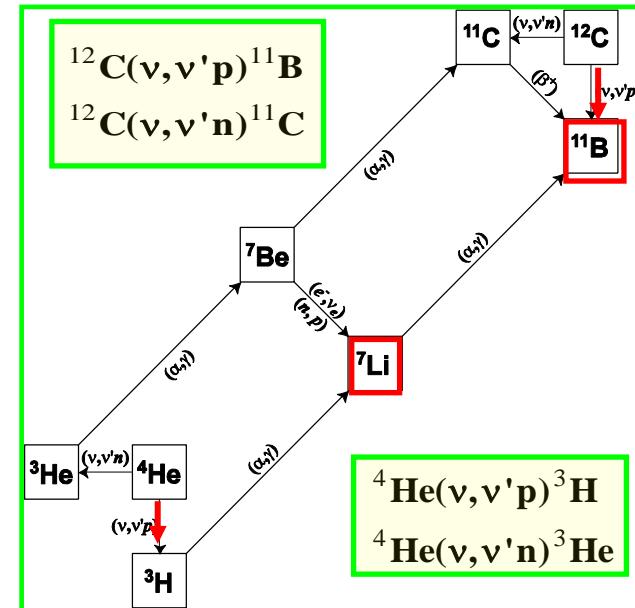
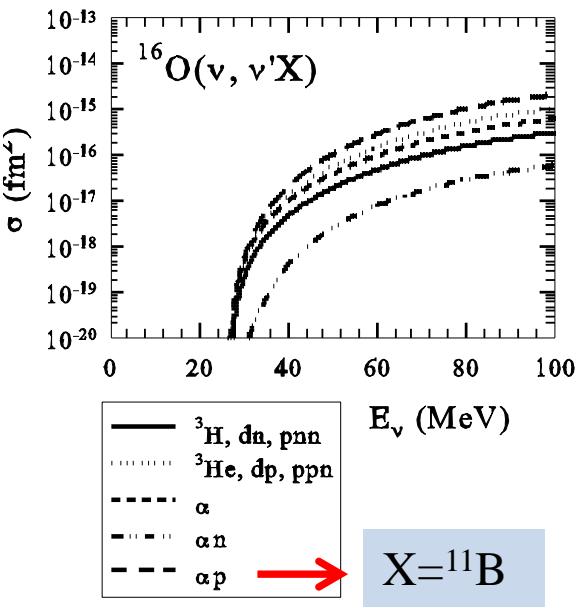
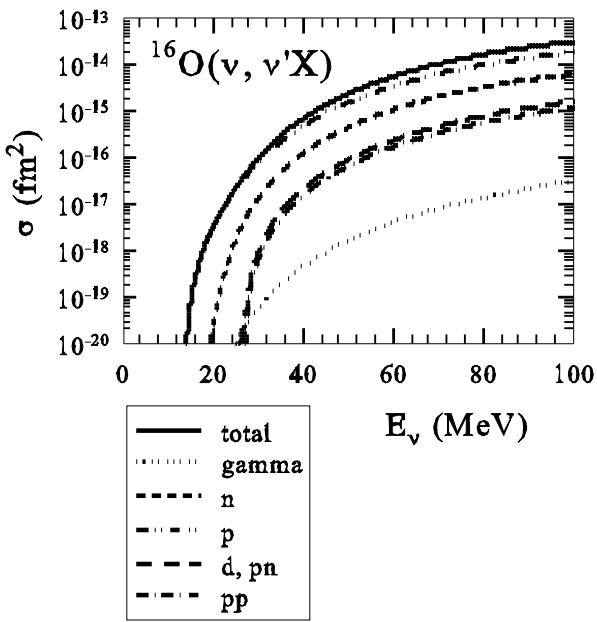
Hauser-Feshbach statistical model

Branching ratios for γ and particle emission channels (with multi-particle emission channels):

γ , n, p, np (d), nn, pp, ${}^3\text{H}$ (nnp), ${}^3\text{He}$ (npp), α , ap, an, ann, anp, app, ...

Isospin conservation is taken into account (S. Chiba)





$$\frac{\sigma(^{16}\text{O}(\nu, \nu'\alpha p)^{11}\text{B})}{\sigma(^{12}\text{C}(\nu, \nu'p)^{11}\text{B})} \approx 10\%$$

Case1: previous branches used in ^{16}O (γ, n, p, α -emissions) and HW92 cross sections

Case2: previous branches, and new cross sections

Case3: multi-particle branches and new cross sections

Production yields of ^{11}B and ^{11}C (10^{-7}M_\odot)

核種生成量	15M $_\odot$ モデル			20M $_\odot$ モデル		
	Case 1	Case 2	Case 3	Case 1	Case 2	Case 3
$M(^{11}\text{B})$	2.94	2.92	3.13	6.77	6.58	7.66
$M(^{11}\text{C})$	2.80	2.71	3.20	9.33	8.91	9.64
$M(^{11}\text{B} + ^{11}\text{C})$	5.74	5.62	6.33	16.10	15.49	17.29

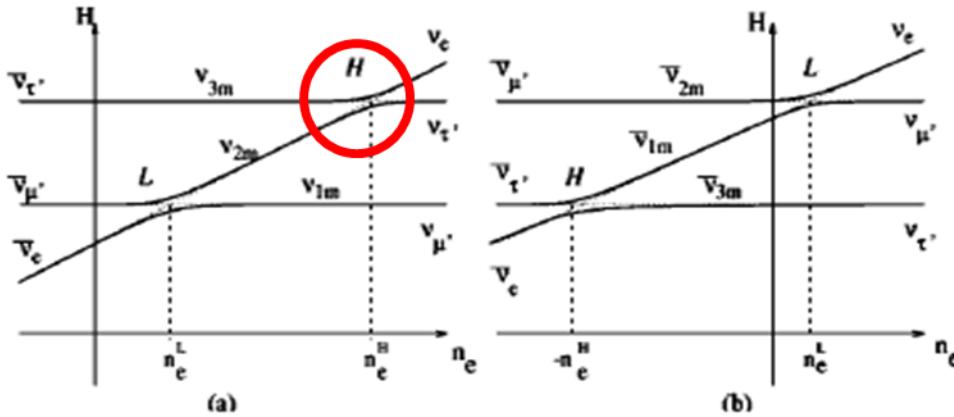
T. Yoshida

ν oscillation effects → ν mass hierarchy

MSW ν oscillations

Normal hierarchy

Inverted hierarchy



Normal – hierarchy : $\nu_\mu, \nu_\tau \rightarrow \nu_e$

Inverted – hierarchy : $\bar{\nu}_\mu, \bar{\nu}_\tau \rightarrow \bar{\nu}_e$

$$N(\nu_e) = P * N^0(\nu_e) + (1-P) * N^0(\nu_x)$$

$$N(\text{anti-}\nu_e) = P' * N^0(\text{anti-}\nu_e) + (1-P') * N^0(\nu_x)$$

Normal hierarchy: $(P, P') = (0, 0.68)$

Inverted hierarchy: $(P, P') = (0.32, 0); \sin^2 \theta_{12} = 0.32$

Dighe and Smirnov, PR D62, 033007 (2000)

Resonance condition

$$\rho Y_e = N_e = \frac{\Delta}{2\sqrt{2}EG_F} \cos 2\theta$$

$$= 6.55 \times 10^6 \left(\frac{\Delta m_{ij}^2}{1 \text{ eV}^2} \right) \left(\frac{1 \text{ MeV}}{E_\nu} \right) \cos 2\theta_{ij} \text{ g} \cdot \text{cm}^{-3}$$

H – resonance: θ_{13}

$\rho Y_e = 300 - 3000 \text{ g} \cdot \text{cm}^{-3}$ He/C layer

L – resonance: θ_{12}

$\rho Y_e = 4 - 40 \text{ g} \cdot \text{cm}^{-3}$

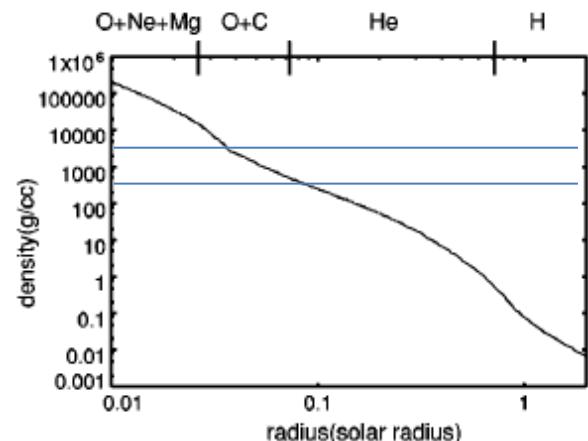
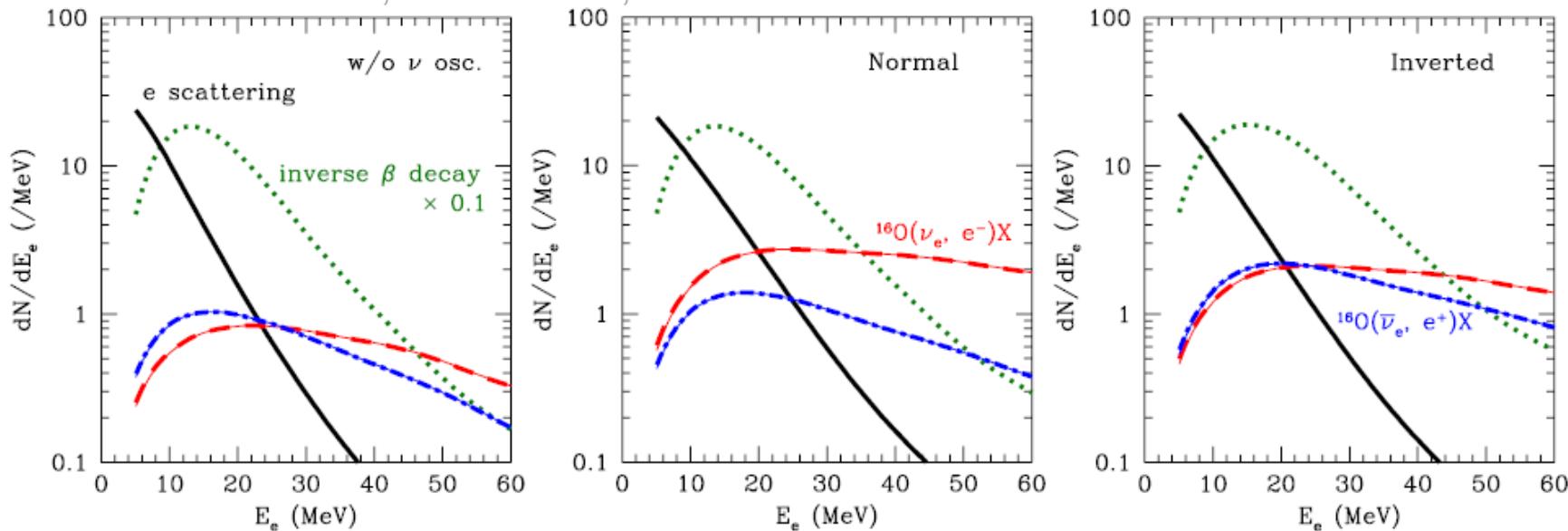


FIG. 3. Density profile of the presupernova star model used in the paper [20]. The progenitor mass is set to be $15 M_\odot$.

Charged current scattering off ^{16}O nucleus as a detection channel of supernova neutrinos

Ken'ichiro Nakazato¹, Toshio Suzuki², and Makoto Sakuda³

PTEP 2018, 123E02 82018)



$(M, Z) = (20M_{\odot}, 0.02)$ Z = metalicity

$\langle E_{\nu_e} \rangle = 9.32 \text{ MeV}$, $\langle E_{\bar{\nu}_e} \rangle = 11.1 \text{ MeV}$, $\langle E_{\nu_X} \rangle = 11.9 \text{ MeV}$

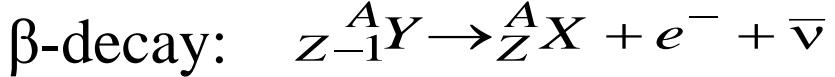
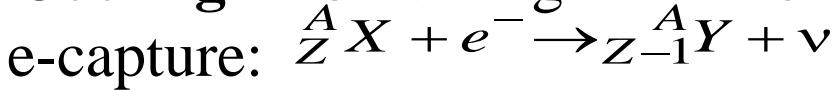
Expected enent numbers

Nakazato et al., ApJ. Suppl. 205, 2 (2013)

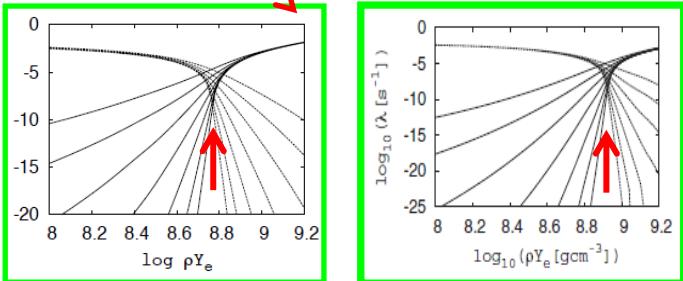
reaction	ordinary supernova			black hole formation		
	no osc.	normal	inverted	no osc.	normal	inverted
$^{16}\text{O}(\nu_e, e^-)\text{X}$	41	178	134	2482	2352	2393
$^{16}\text{O}(\bar{\nu}_e, e^+)\text{X}$	36	58	103	1349	1255	1055
electron scattering	140	157	156	514	320	351
inverse β -decay	3199	3534	4242	17525	14879	9255
total	3416	3927	4635	21870	18806	13054

2. Evolution of O-Ne-Mg cores in stars

- Cooling of O-Ne-Mg core in $8\text{-}10 M_{\odot}$ stars by nuclear URCA process

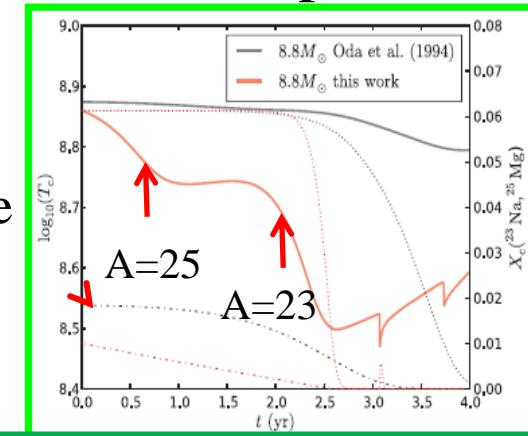


$(^{25}\text{Na}, ^{25}\text{Mg})$ & $(^{23}\text{Ne}, ^{23}\text{Na})$ pairs \rightarrow Cooling of O-Ne-Mg core



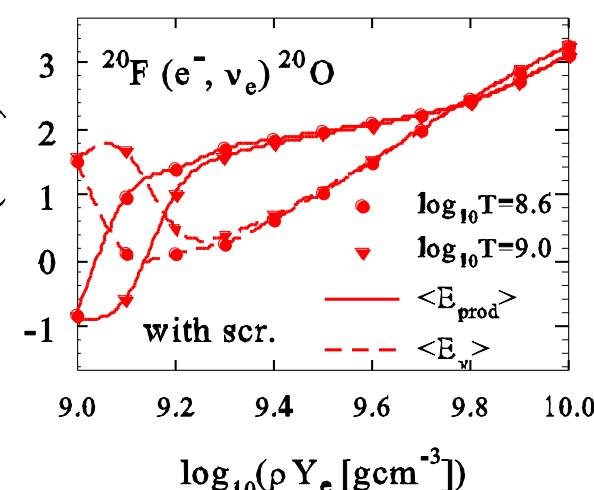
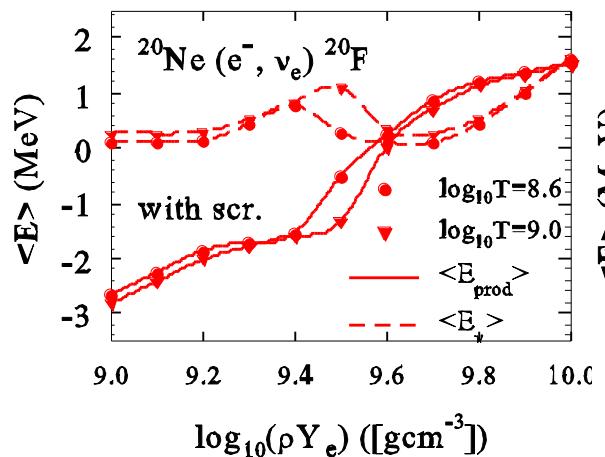
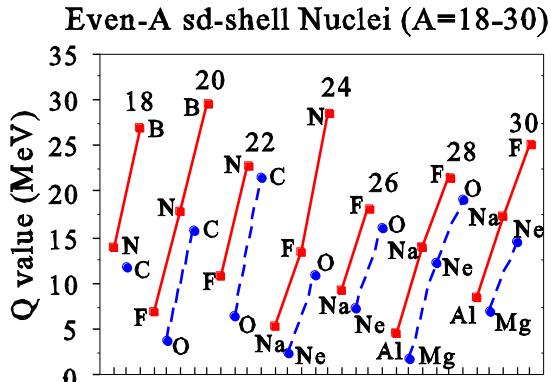
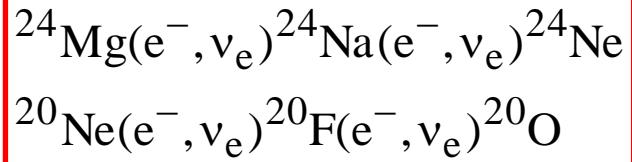
$\log_{10} T = 8$ to 9.2 in steps of 0.2

Toki, Suzuki, Nomoto, Jones and Hirschi, PR C 88, 015806 (2013)
Jones et al., Astrophys. J. 772, 150 (2013)



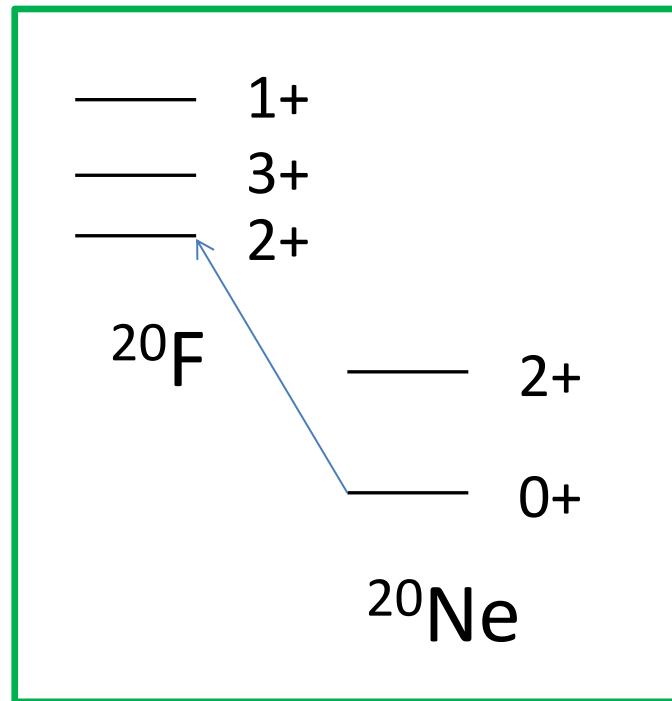
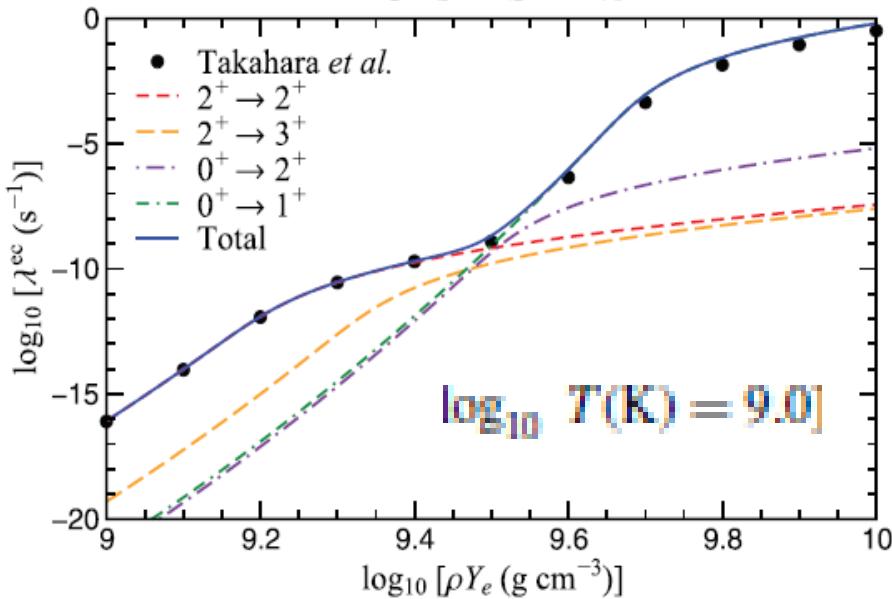
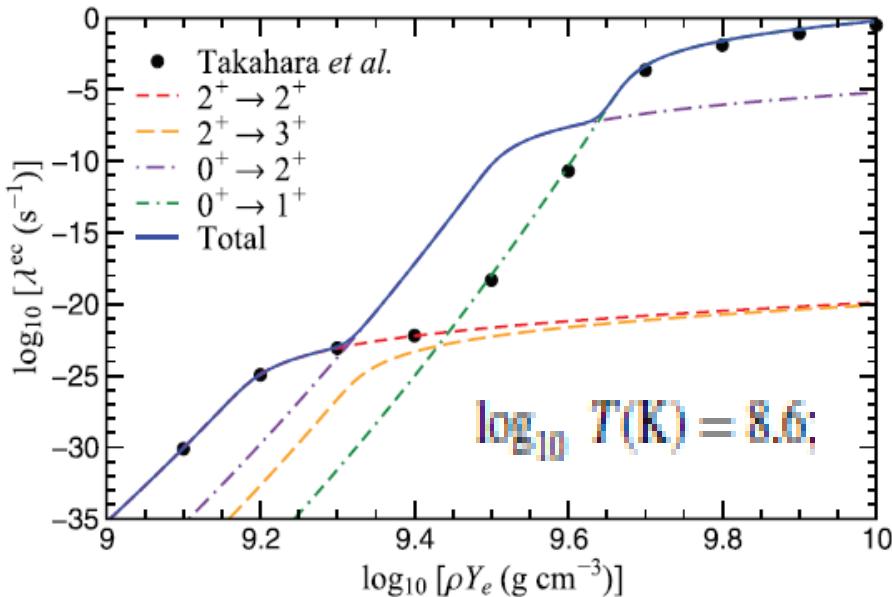
$8.8 M_{\odot}$ star collapses triggered by subsequent e-capture on ^{24}Mg and ^{20}Ne (e-capture SNe)

- Heating of stars: γ emission after double e-captures on ^{24}Mg and ^{20}Ne



Suzuki, Toki and Nomoto, ApJ. 817 (2016)

$^{20}\text{Ne} (\text{e}^-, \nu) ^{20}\text{F}$



$0+ \rightarrow 2+$: 2nd-forbidden transition

Kirsebom et al., arXiv:1805.08149

β -decay exp. $\rightarrow \log ft = 10.47(11)$

cf. NNDC : $\log ft > 10.5$

$$ft = \frac{6147}{B(\text{GT})} \quad B(\text{GT}) = g_A \vec{\sigma} \tau_-$$

e-capture rate

$$\lambda^\infty(T) [\text{s}^{-1}] = \frac{g_V^2 c}{\pi^2 (\hbar c)^3} \int_{E_{\min}}^{\infty} \sigma(E_e, T) E_e p_e c f_e(E_e) dE_e, \quad f_\alpha = \frac{1}{1 + e^{\frac{\alpha - \mu}{k_B T}}},$$

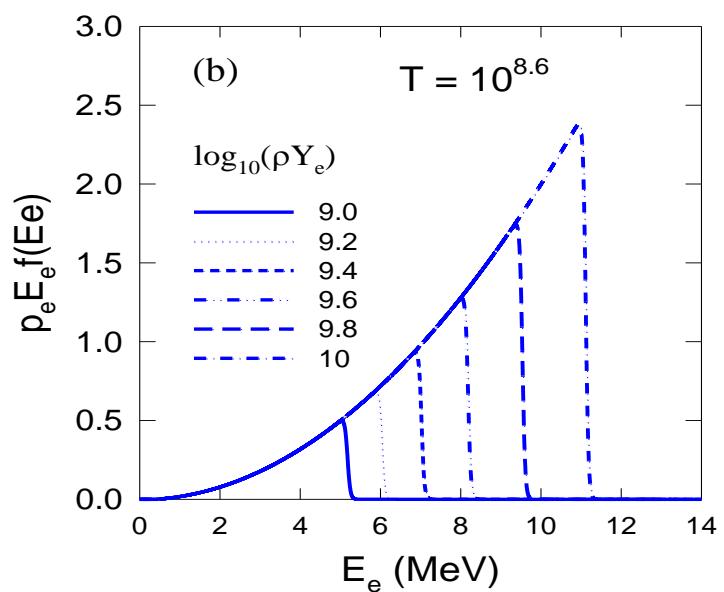
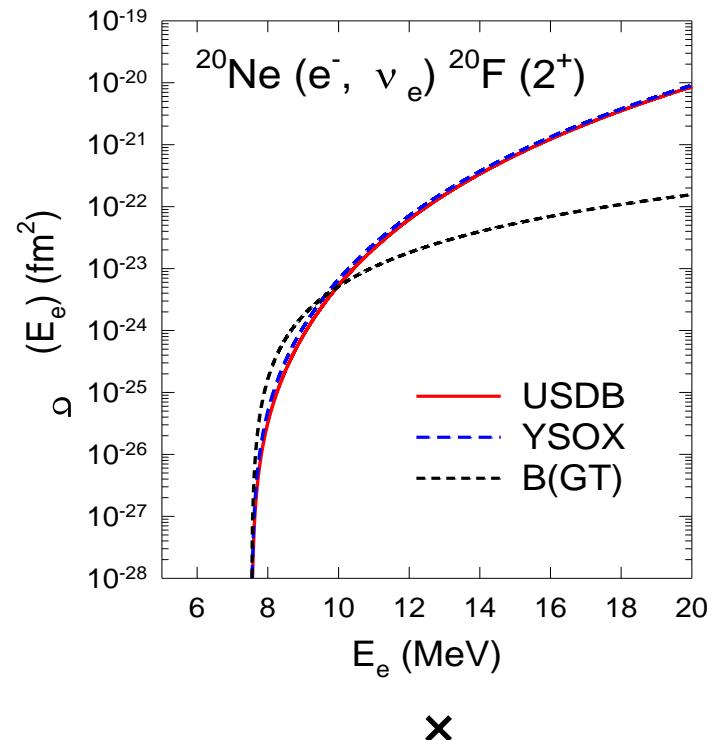
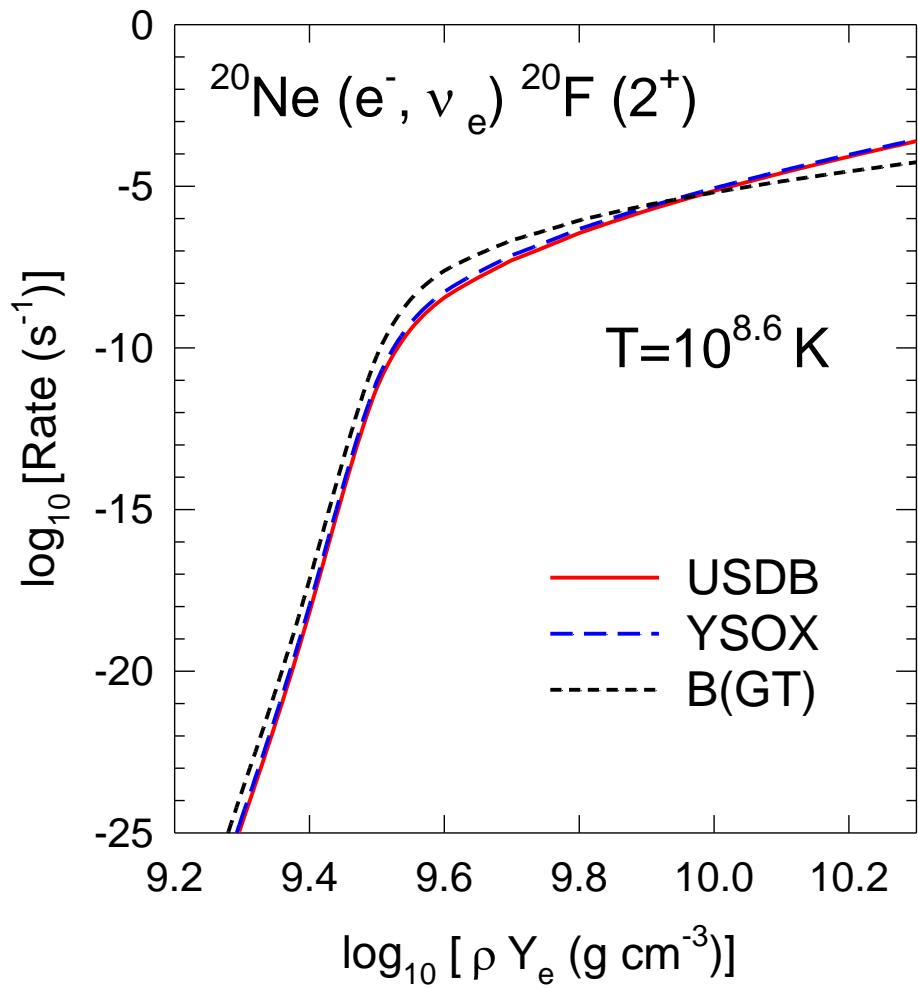
$$\begin{aligned} \frac{d\sigma}{d\Omega} &= \frac{G_F^2 \cos^2 \theta_c}{2\pi} \frac{F(Z, E_e)}{(2J_i + 1)} \\ &\times \left(\sum_{J \geq 1} \mathcal{W}(E_v) \{ (1 - (\hat{\mathbf{v}} \cdot \hat{\mathbf{q}})(\beta \cdot \hat{\mathbf{q}})) \right. \\ &\times [|\langle J_f || \hat{T}_J^{\text{mag}} || J_i \rangle|^2 + |\langle J_f || \hat{T}_J^{\text{el}} || J_i \rangle|^2] \\ &- 2\hat{\mathbf{q}} \cdot (\hat{\mathbf{v}} - \beta) \text{Re} \langle J_f || \hat{T}_J^{\text{mag}} || J_i \rangle \langle J_f || \hat{T}_J^{\text{el}} || J_i \rangle^* \} \\ &+ \sum_{J \geq 0} \mathcal{W}(E_v) \{ (1 - \hat{\mathbf{v}} \cdot \beta) + 2(\hat{\mathbf{v}} \cdot \hat{\mathbf{q}})(\beta \cdot \hat{\mathbf{q}}) \\ &\times \langle J_f || \hat{\mathcal{L}}_J || J_i \rangle |^2 + (1 + \hat{\mathbf{v}} \cdot \beta) \langle J_f || \hat{\mathcal{M}}_J || J_i \rangle |^2 \} \\ &\left. - 2\hat{\mathbf{q}}(\hat{\mathbf{v}} + \beta) \text{Re} \langle J_f || \hat{\mathcal{L}}_J || J_i \rangle \langle J_f || \hat{\mathcal{M}}_J || J_i \rangle^* \} \right), \end{aligned}$$

Parr et al., PRC80, 055801 (2009)

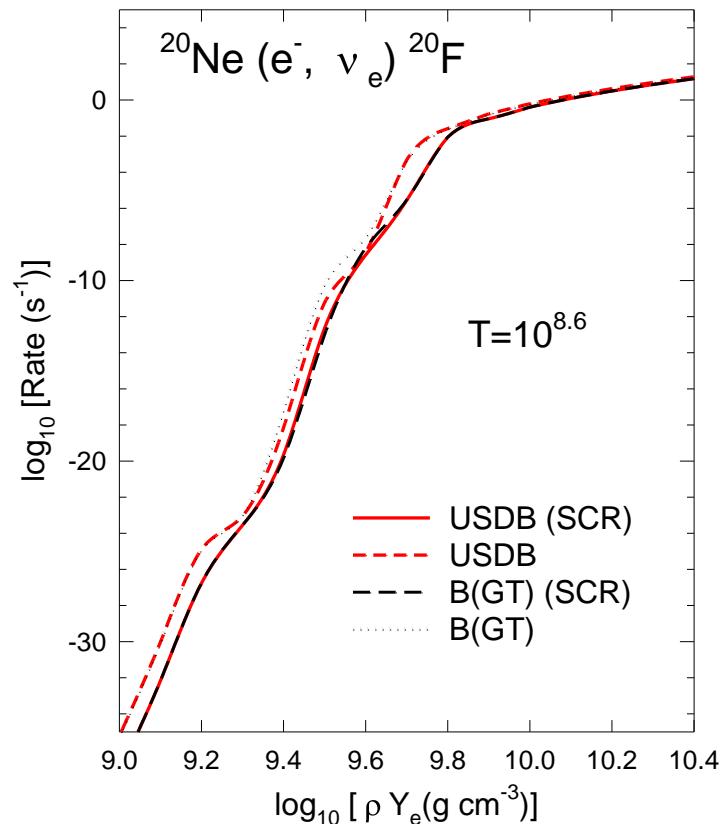
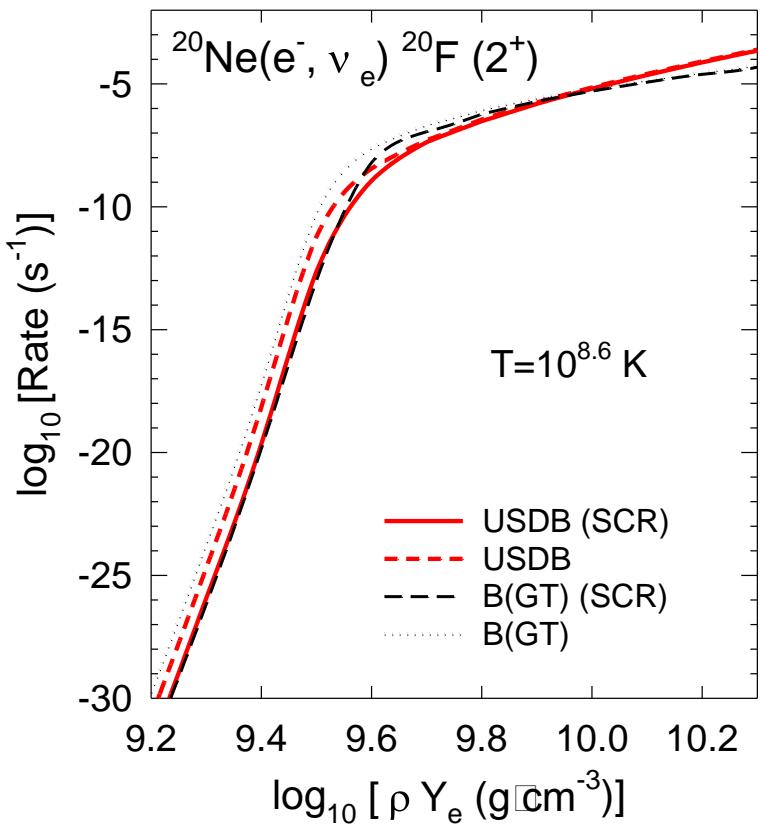
$$J=2^+ \quad C2+L2: \quad F_l^V \frac{q_\mu^2}{q^2} j_2(qr) Y^2$$

$$E2: \quad \frac{q}{M} [F_l^V (\sqrt{\frac{3}{5}} j_1(qr) [Y^1 \times \frac{\nabla}{q}]^2 - \sqrt{\frac{2}{5}} j_3(qr) [Y^3 \times \frac{\nabla}{q}]^2) + \frac{1}{2} \mu_V j_2(qr) [Y^2 \times \sigma]^2]$$

$$M2_5: \quad F_A j_2(qr) [Y^2 \times \sigma]^2$$



Coulomb effects



Summary

1. ν - ^{16}O reactions and ν detection

- Spin-dipole strengths in ^{16}O and ν - ^{16}O reaction cross sections are evaluated with a shell-model Hamiltonian, SFO-tls.

Partial cross sections for γ emission, single- and multi-particle emission channels are obtained by Hauser-Feshbach method with isospin conservation.

- Syntheses of light elements ^{11}B and ^{11}C in SNe are studied with full inclusion of the multi-particle emission processes.
- Charged-current reactions on ^{16}O induced by SNv are studied.

Mass hierarchy dependence of the cross sections is presented; promising for distinguishing mass hierarchies in future SNv detection

2. Evolution of O-Ne-Mg core in stars

- e-capture rates for 2nd-forbidden transition in ^{20}Ne , ^{20}Ne (0^+) \rightarrow ^{20}F (2^+) are evaluated by multipole expansion method with shell-model Hamiltonians; USDB (sd), YSOX (p-sd), and compared with those obtained with a prescription that treats the transition as an allowed GT one. Sizable difference in the rates are noticed between the two methods.

Collaborators

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