

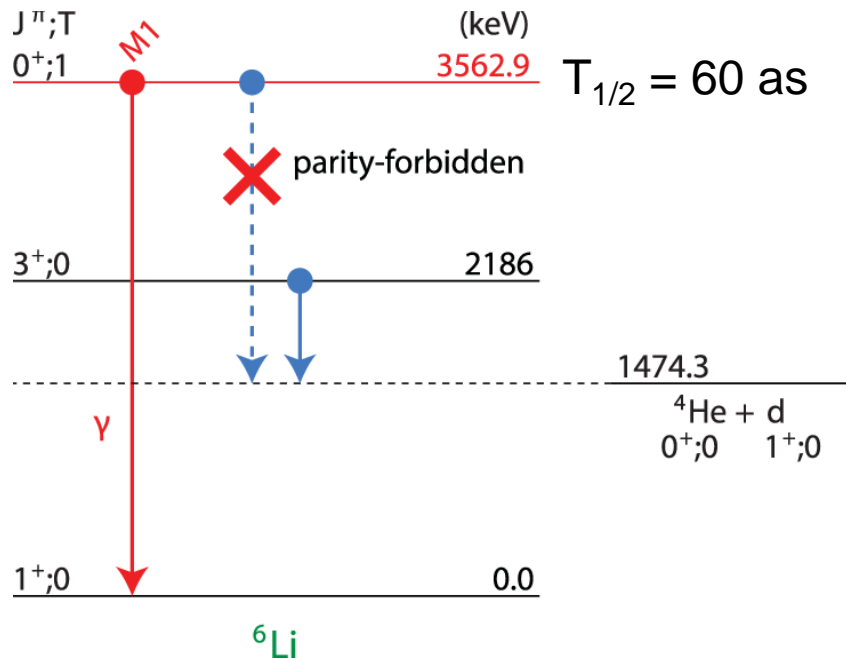
First High-Precision Measurement of the Isovector M1 Strength in ${}^6\text{Li}$ at the Photon Point



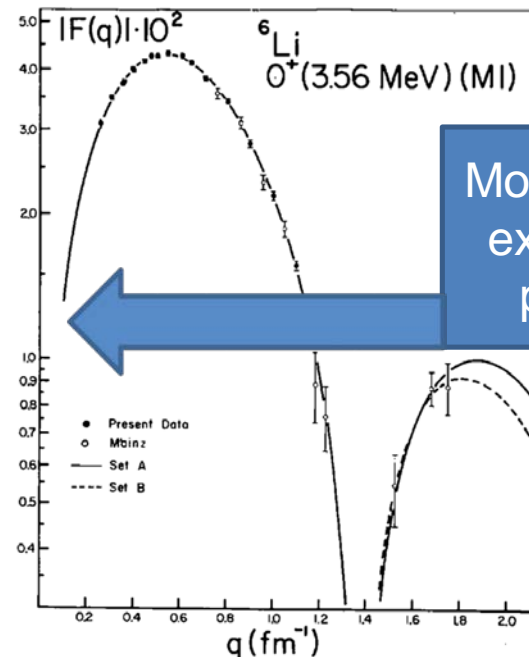
TECHNISCHE
UNIVERSITÄT
DARMSTADT

Norbert Pietralla + group, TU Darmstadt

Udo Gayer, Christoph Romig, Volker Werner et al.



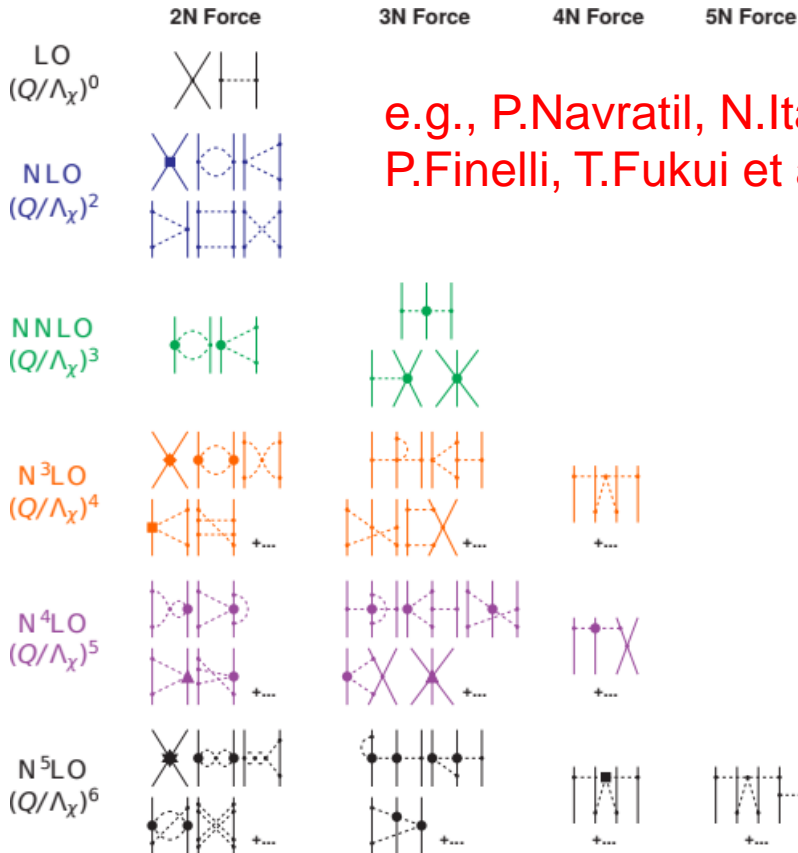
Newest data: Bergstrom (1975)



Model-dependent
extrapolation to
photon point

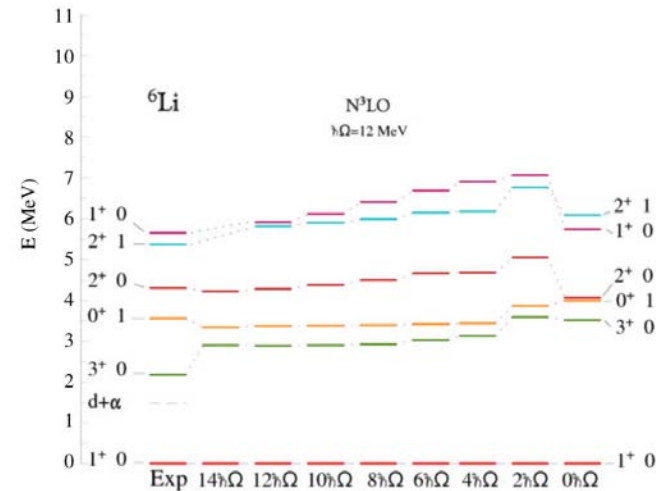
**Claimed
precision
2.3%**

Towards ab-initio Nuclear Theory



e.g., P.Navratil, N.Itaco,
P.Finelli, T.Fukui et al.

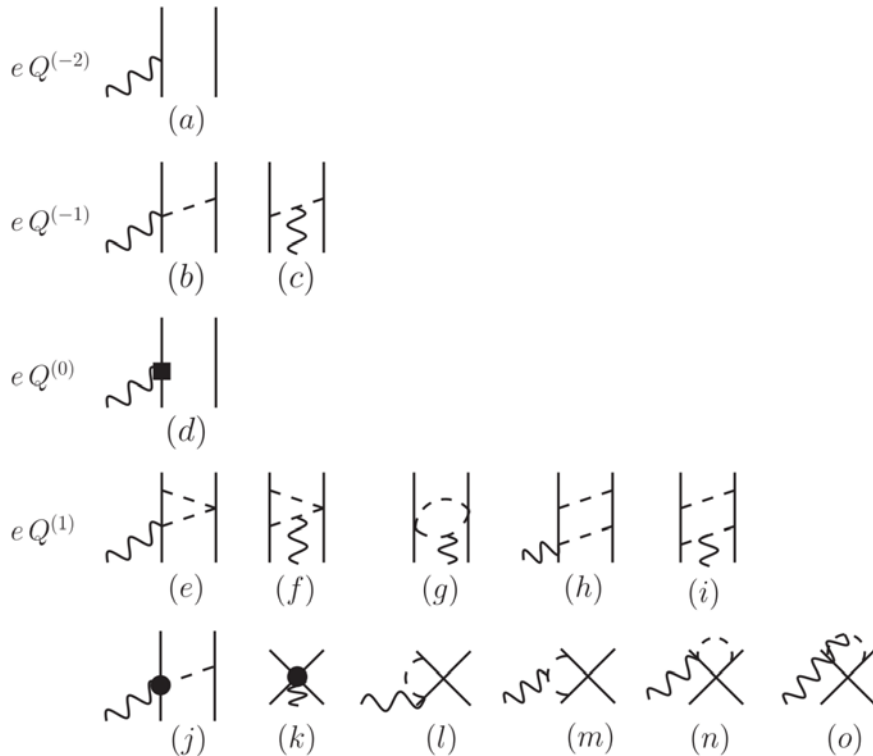
- NN-potential from χ EFT
- → Systematic expansion
- → Theoretical “error bars”
- Increasing complexity
- Ab-initio calculations possible for light nuclei



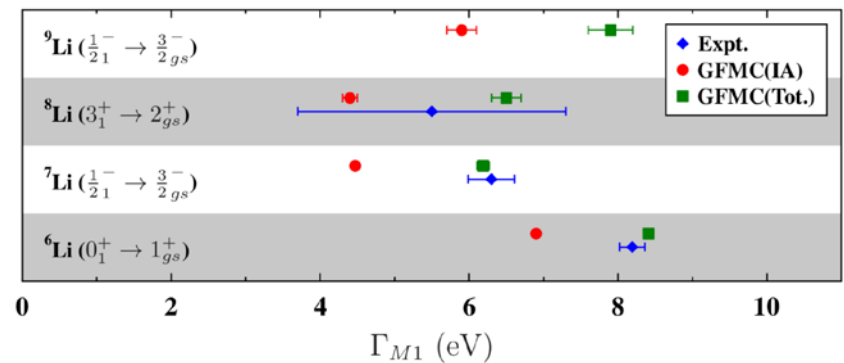
D.R. Entem, R. Machleidt, Y. Nosyk, PRC **96** (2017)

B.R. Barrett, P. Navrátil, J.P. Vary, Prog. Part. Nucl. Phys. **69** (2013)

Requires Consistent Transition Operator



- NN-potential from χ EFT
- \rightarrow Systematic expansion
- \rightarrow Theoretical “error bars”
- Increasing complexity
- Ab-initio calculations possible for light nuclei
- Higher-order “effects” comparable to experimental uncertainties



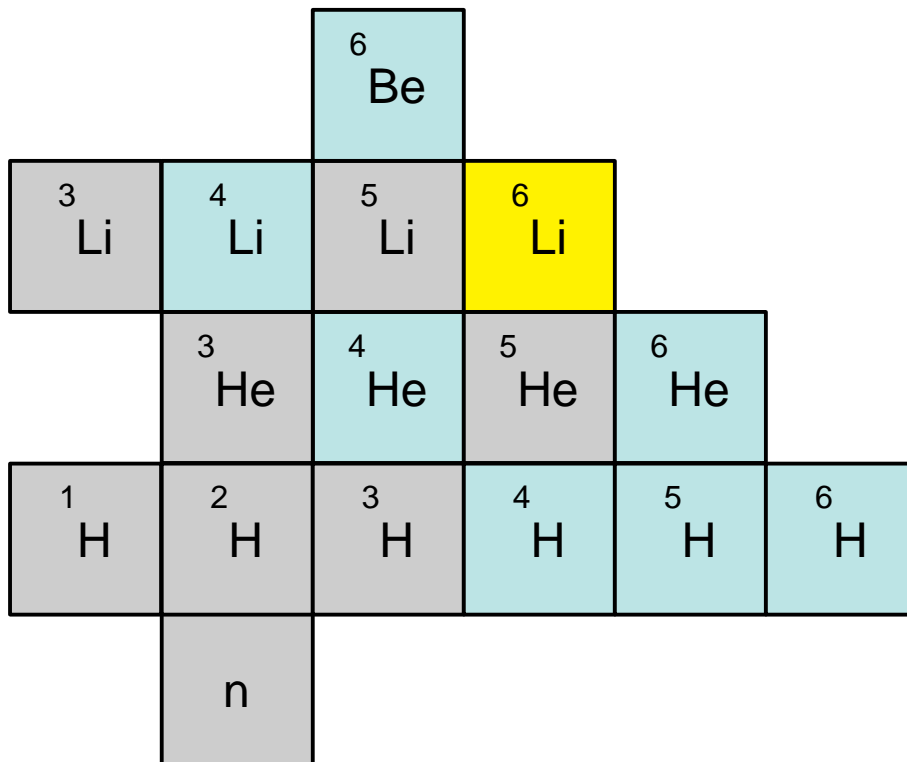
S. Pastore et al., PRC 87 (2013)

Figure: C. Romig, Doctoral Thesis, TU Darmstadt (2015)

S. Bacca, S. Pastore, J. Phys. G **41**, 123002 (2014)

Testing Theory for Consistent EM Transitions

Predominant decay modes of excited states of known $A \leq 6$ nuclei



no bound excited state

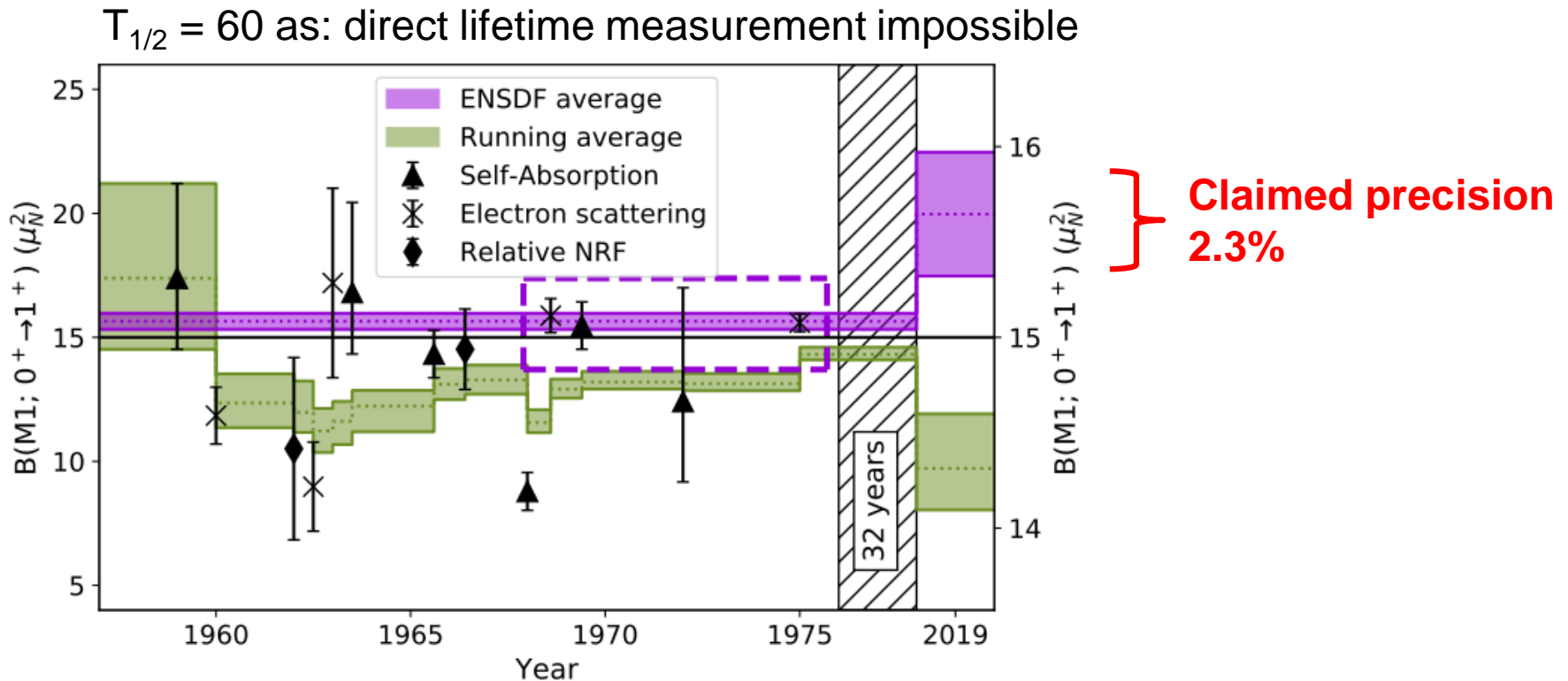
particle emission

particle + γ emission

- ${}^6\text{Li}$ (3.562 MeV) is the lightest system that decays predominantly via EM interaction
- $B(M1; 0^+ \rightarrow 1^+)$ one of the strongest known M1 transitions
- Ideal testing ground for ab-initio nuclear theory

Decay data from: <https://www.nndc.bnl.gov/> (03/25/2019)

History of Experimental Data



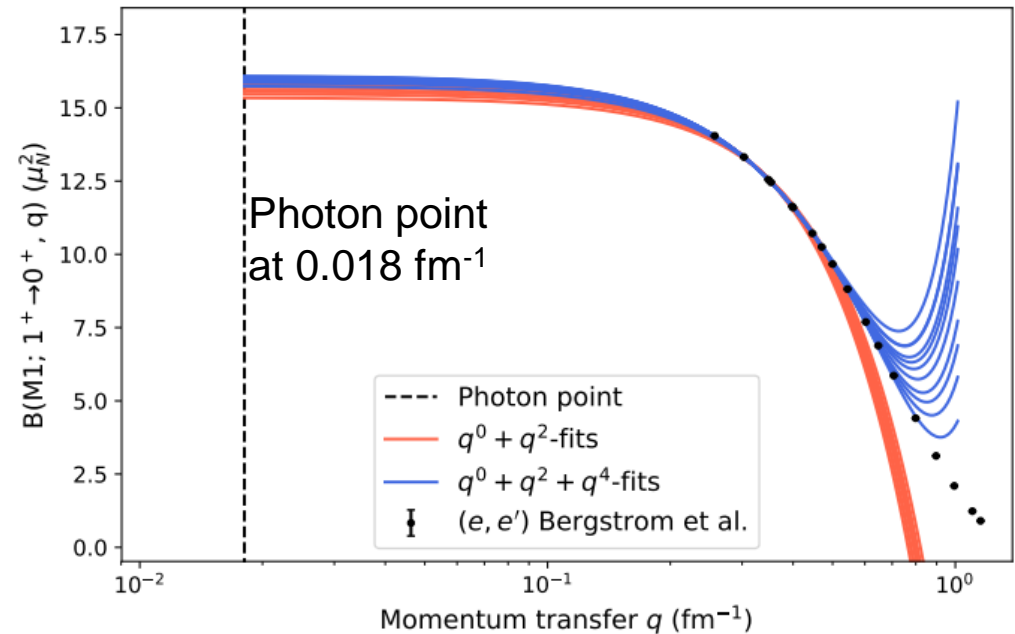
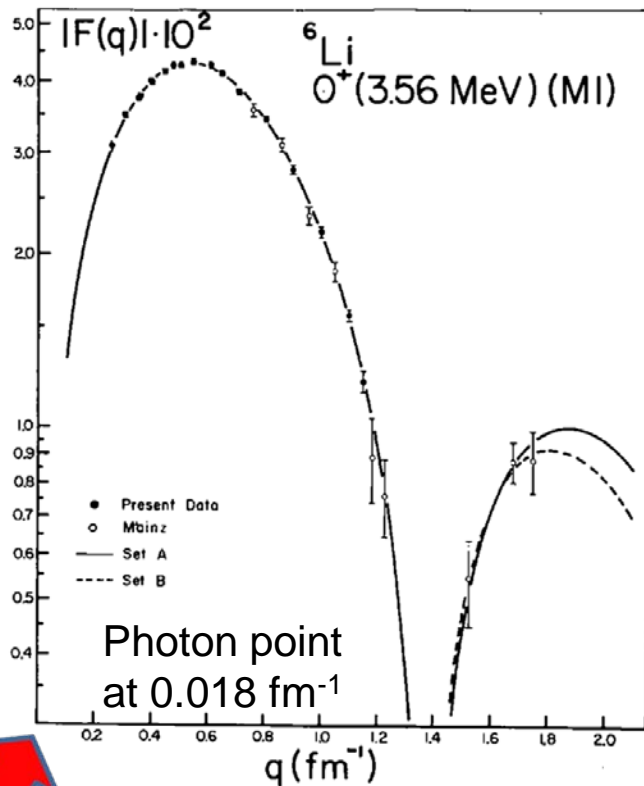
- Current ENSDF literature value based on three selected experiments; in conflict with running world-average

D.R. Tilley *et al.*, Nucl. Phys. **A708** (2002)

ENSDF Value based on these Data

Plane-Wave Born Approximation

$$B(M1, q) = B(M1)(1 - c_1 q^2 + c_2 q^4 + \dots)$$



- Most precise value of $B(M1) = 15.6(4) \mu_N^2$ from extrapolation of (e,e') form factor, **claimed 2.3%**

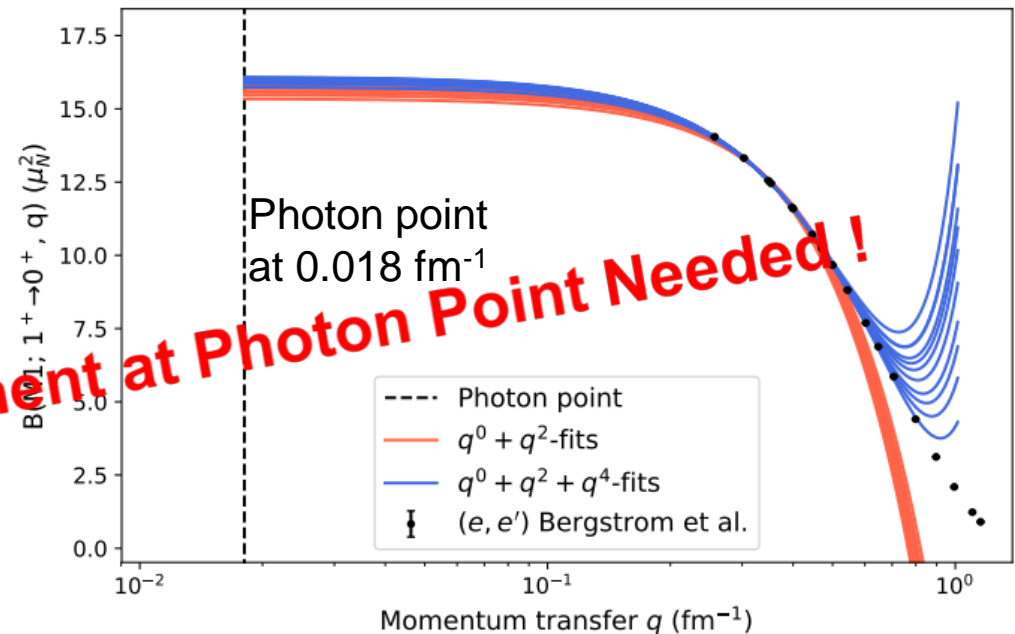
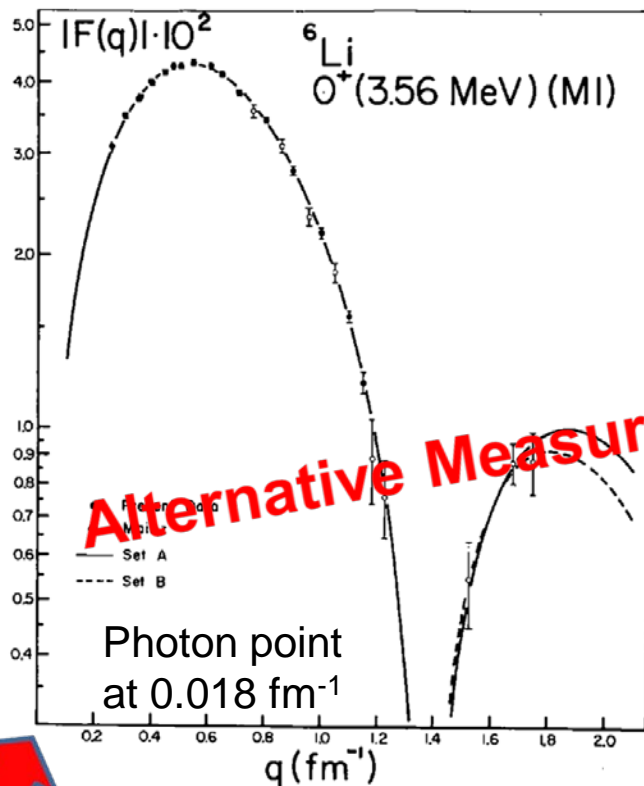
J.C. Bergstrom, I.P. Auer, R.S. Hicks, Nucl. Phys. **A251** (1975)

Extrapolate to 0.018 fm^{-1} !

ENSDF Value based on these Data

Plane-Wave Born Approximation

$$B(M1, q) = B(M1)(1 - c_1 q^2 + c_2 q^4 + \dots)$$

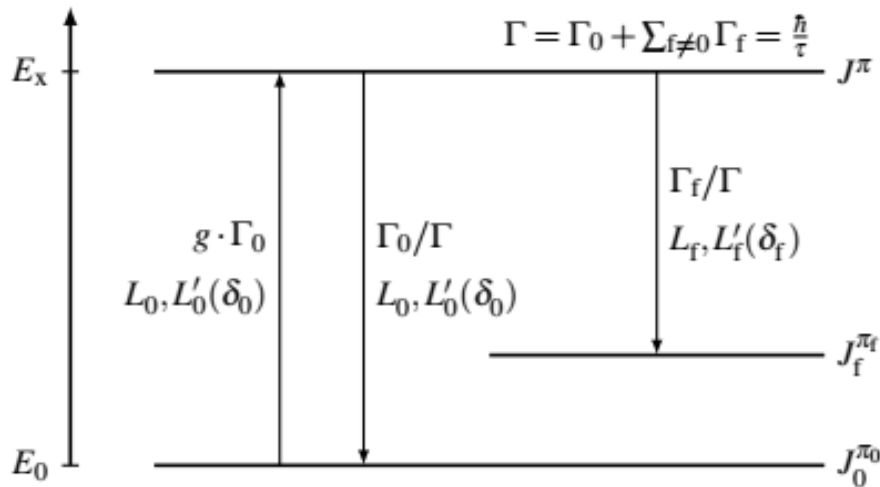


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J.C. Bergstrom, I.P. Auer, R.S. Hicks, Nucl. Phys. **A251** (1975)

Extrapolate to 0.018 fm^{-1} !

Use Photons! Nuclear Resonance Fluorescence



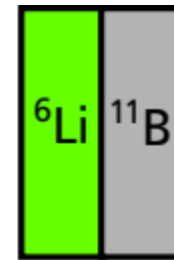
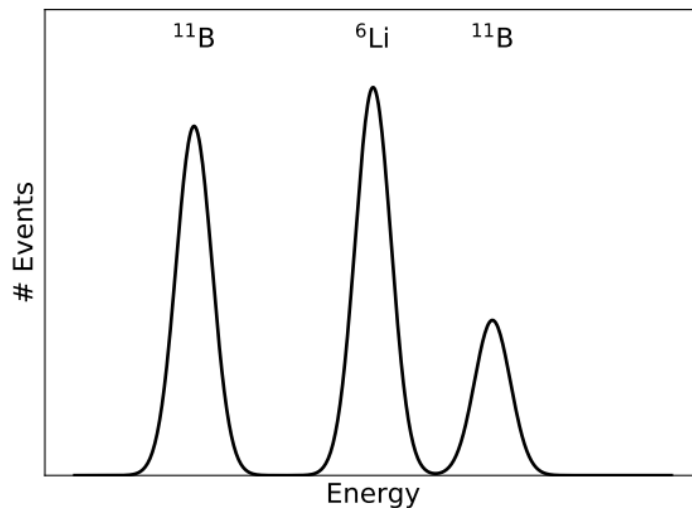
- Nuclear resonance fluorescence (NRF)
- Measurement at the photon point
- → no extrapolation
- Proportionality of cross section to level width / $B(M1)$ (in this case)

$$I \propto \Gamma \propto \frac{\hbar}{2\pi} / \tau$$

$$\propto B(M1)$$

F.R. Metzger, Prog. Nucl. Phys. 7 (1959)

Nuclear Resonance Fluorescence

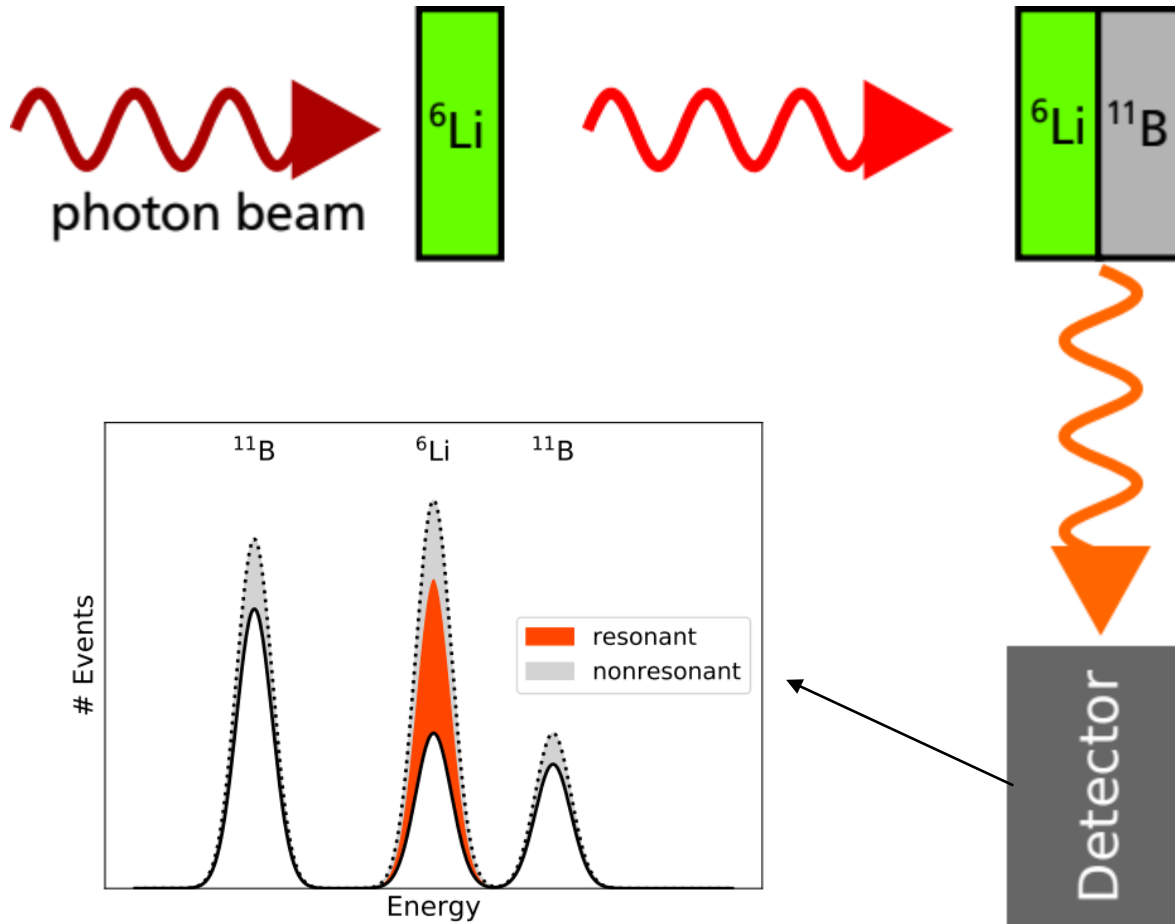


- Nuclear resonance fluorescence (NRF)
- Measurement at the photon point
- → no extrapolation
- Proportionality of cross section to level width / $B(M1)$ (in this case)

$$I \propto \Gamma \propto \frac{h}{2\pi} / \tau \propto B(M1)$$

→ Measurement relative to calibration standard

Relative Self-Absorption



- Relative self-absorption (RSA) based on NRF

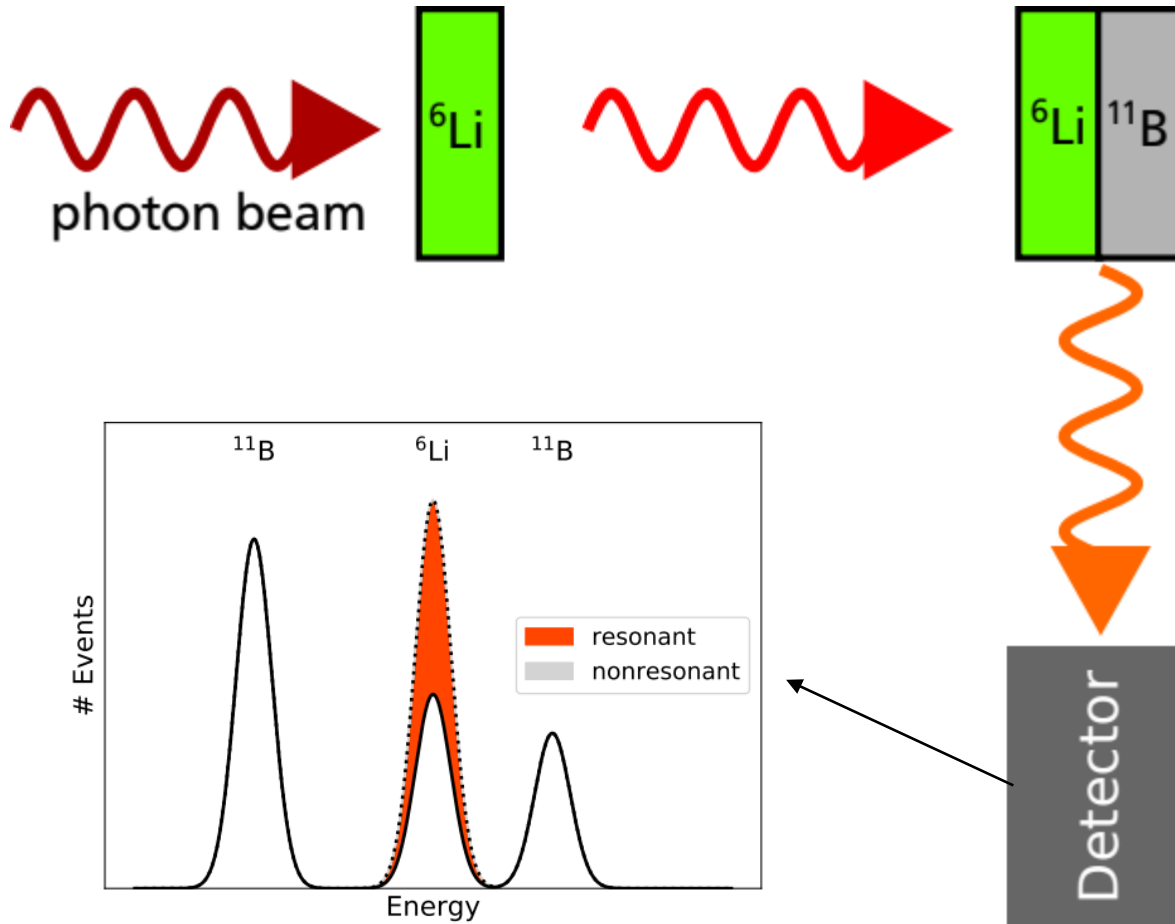
C. Romig *et al.*, Phys. Lett. B **744** (2015)

C. Romig, Doctoral Thesis, TU Darmstadt (2015)

- Reduction of count rate depends on level width
- and atomic absorption

$$R \left(\Gamma_{0^+ \rightarrow 1^+} \right) = \frac{N_{NRF} - N_{RSA}}{N_{NRF}}$$

Relative Self-Absorption



- Relative self-absorption (RSA) based on NRF

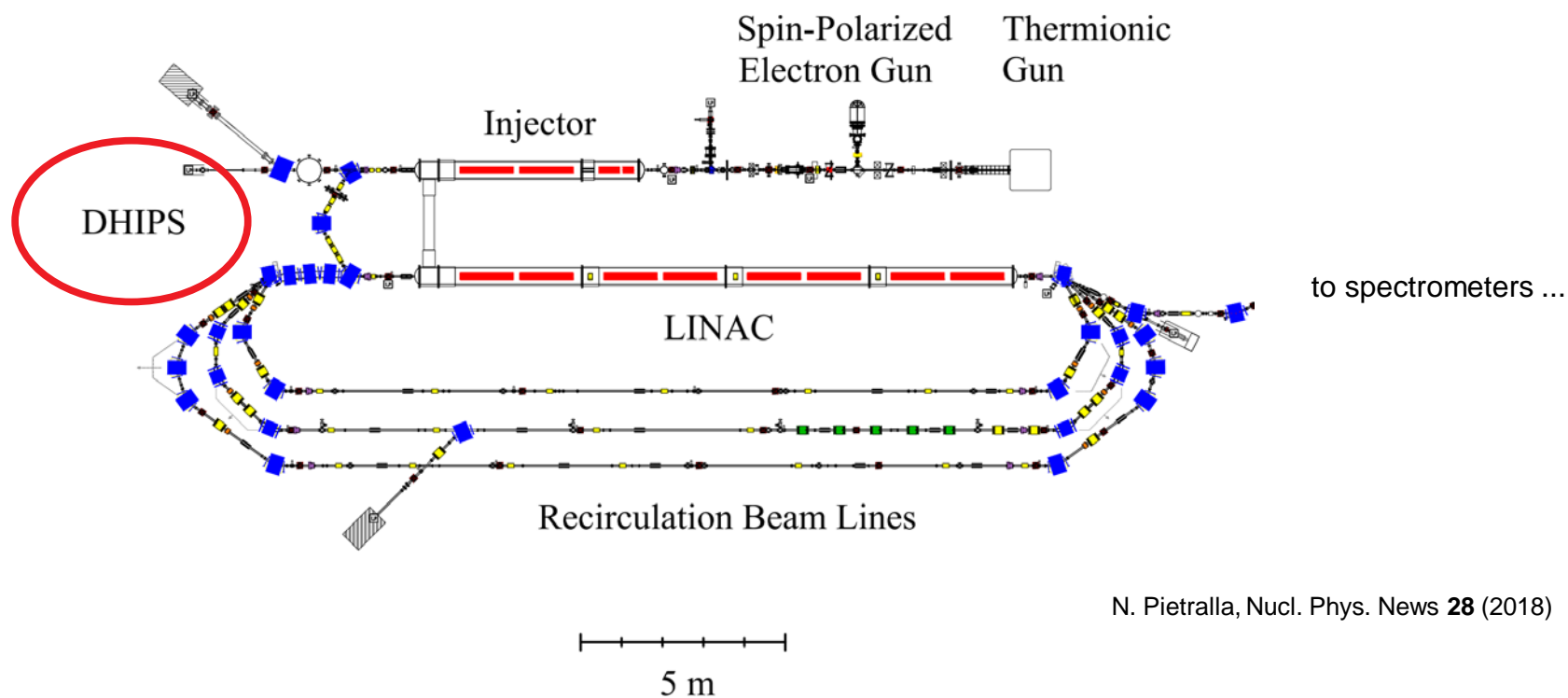
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- Reduction of count rate depends on level width
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$$R \left(\Gamma_{0^+ \rightarrow 1^+} \right) = \frac{N_{NRF} - N_{RSA}}{N_{NRF}}$$

- 'Monitor target' to correct for atomic absorption



N. Pietralla, Nucl. Phys. News **28** (2018)

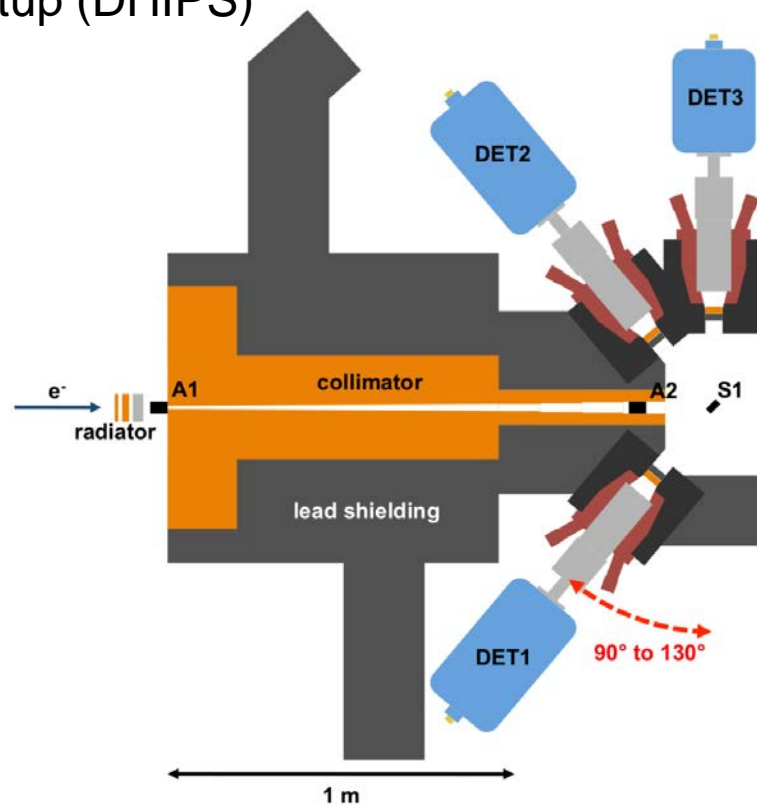
Superconducting-DArmstadt LINEar ACcelerator (S-DALINAC)

- Injector energy ≤ 10 MeV
- Injector current ~ 20 μ A

Darmstadt High-Intensity Photon Setup (DHIPS)

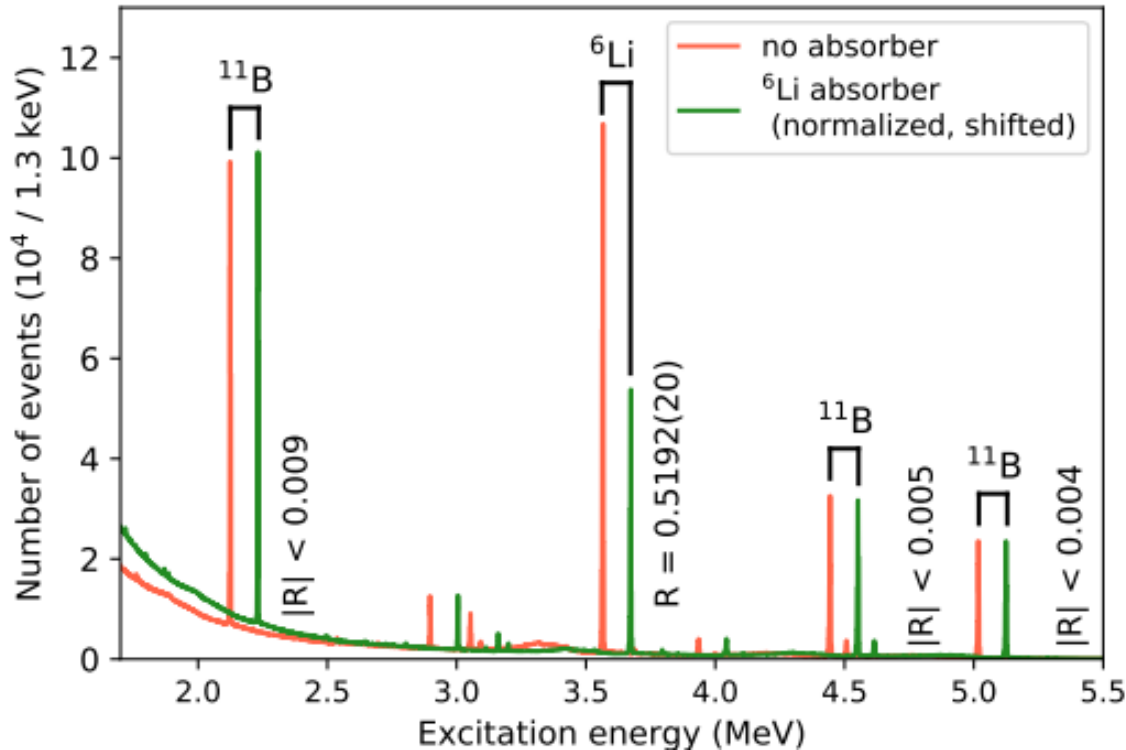
K. Sonnabend et al., NIM **A640** (2011)

- Bremsstrahlungs-photon flux
 $\sim 10^2 \text{ s}^{-1} \text{ eV}^{-1}$
- γ -ray detection by three 120%
High-Purity Germanium (HPGe)
detectors



C. Romig, Doctoral Thesis, TU Darmstadt (2015)

Spectrum



- Normalization of atomic scattering to ${}^{11}\text{B}$ monitor target
- Measuring times:
 - 122h (NRF)
 - 189h (RSA)

$$R(\Gamma_{0^+ \rightarrow 1^+}) = \frac{N_{NRF} - N_{RSA}}{N_{NRF}}$$

$$= 0.5192(20)$$

→ Relative statistical uncertainty of 0.4%!

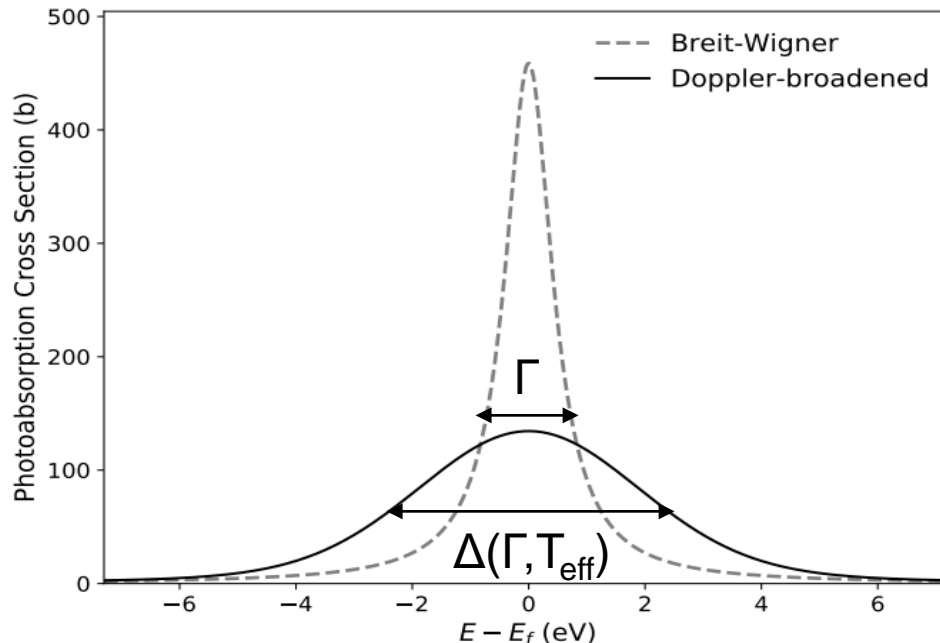
Determination of the level width

- Doppler-broadening due to effective temperature T_{eff} of sample

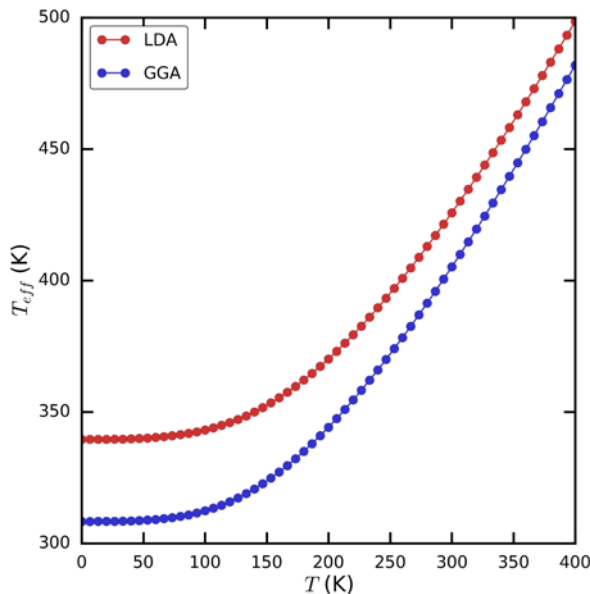
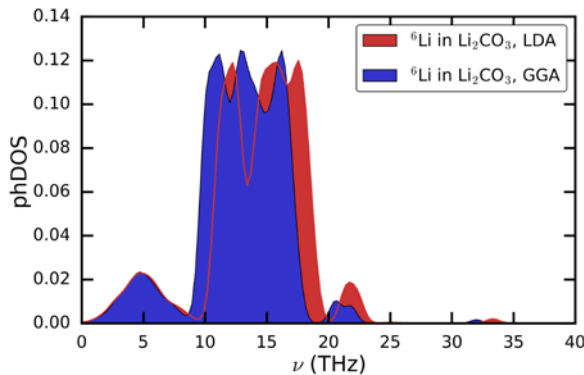
$$\sigma_D(E) = \int_{-\infty}^{\infty} dv \sigma_{BW}(E, \Gamma) \sqrt{\frac{M}{2\pi k_B T_{\text{eff}}}} e^{\frac{-Mv^2}{2k_B T_{\text{eff}}}}$$

$$N = \int_{z_0}^{z_1} dz \int_{-\infty}^{\infty} dE \sigma_D(E, \Gamma, T_{\text{eff}}) e^{-\sigma_D(E, \Gamma, T_{\text{eff}})z}$$

$$R(\Gamma_0, T_{\text{eff}}) = \frac{N_{\text{NRF}} - N_{\text{RSA}}}{N_{\text{NRF}}}$$



Determination of the effective temperature



- T_{eff} for Li_2CO_3 from ab-initio atomic DFT via its relation to the Phonon density of states (phDOS) $g(\nu)$
- Use of different exchange-correlation functionals for uncertainty quantification

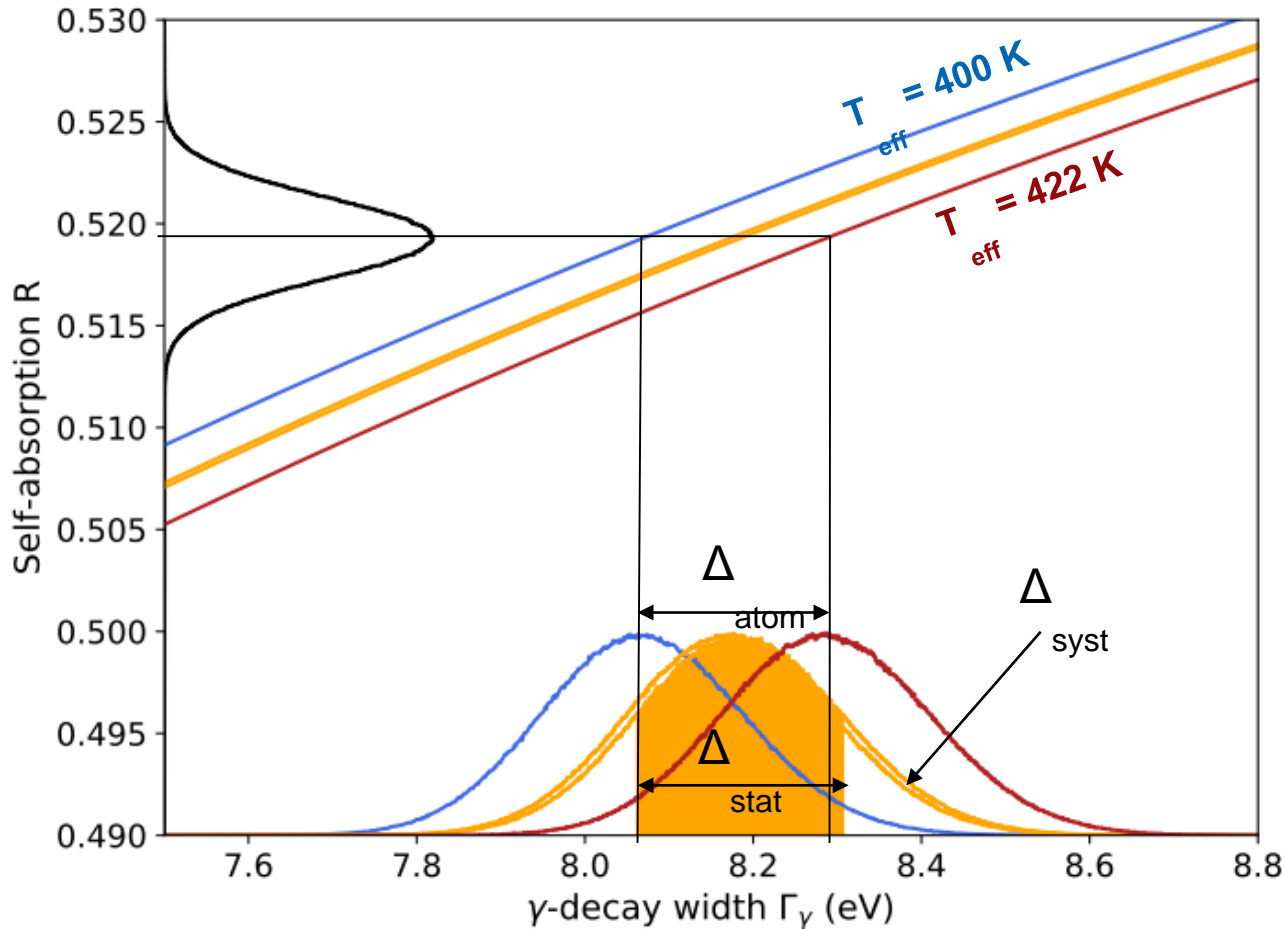
$$k_B T_{\text{eff}} = \int_0^{\infty} h \nu g(\nu) \left(\frac{1}{\exp(h \nu / k_B T)} + \frac{1}{2} \right)$$

W.E. Lamb, Phys. Rev. **55** (1938)

$$T_{\text{eff}}(T = 293\text{K}) = 411(11)\text{K}$$

(2.7% rel. unc.)

Uncertainty Budget



B(M1) = 15.61

Δ_{stat} statistics

+0.23 – 0.21

Δ_{atom} atomic theory

+0.19 – 0.21

Δ_{syst} target dimensions

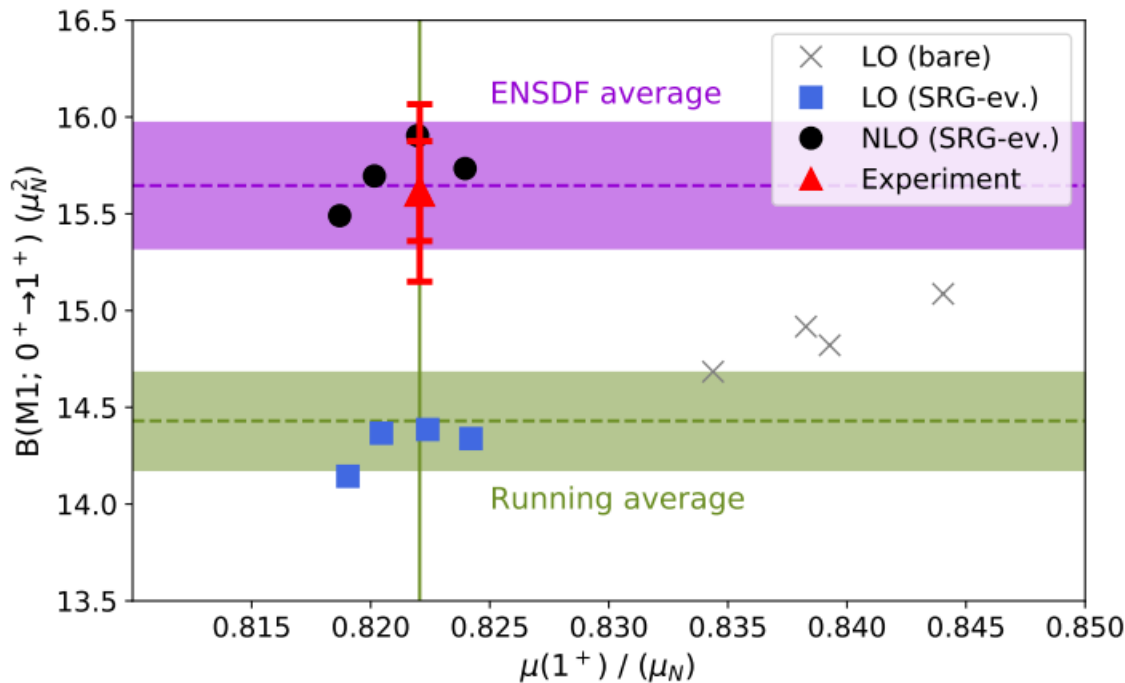
+ -0.04

Δ_{num} numerical evaluation

+ - << 0.01

**= 15.61 +- 0.30
(100% +- 1.9%)**

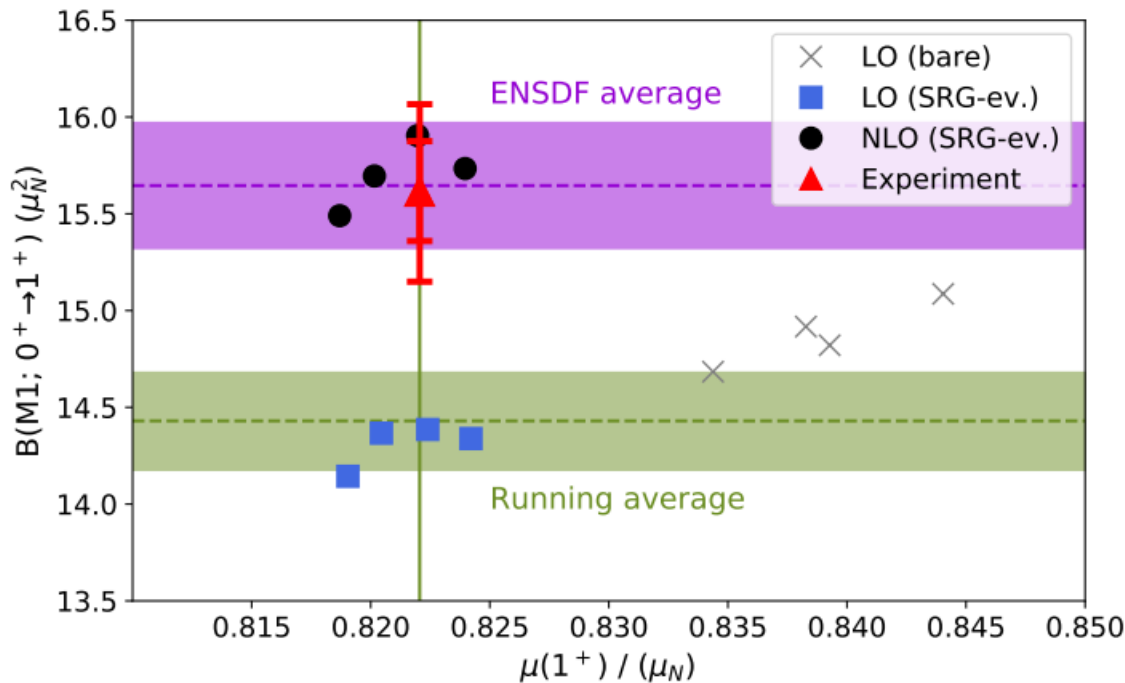
Comparison with theory



- χ EFT calculations of $B(M1; 1^+ \rightarrow 0^+)$ and $\mu(1^+)$ in the no-core shell model
 - SRG-evolved next generation chiral NN+3N interactions up to N4LO+N3LO
D.R. Entem, R. Machleidt, Y. Nosyk, PRC **96** (2017)
 - Unevolved M1 operator, evolved M1 operator at LO, and evolved M1 operator at NLO
- First complete chiral calculation of these observables

High experimental precision crucial to test state-of-the-art theory!

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Thank you!

Thanks to U.Gayer for preparation of slides

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