

Rare and radiative B decays at LHCb

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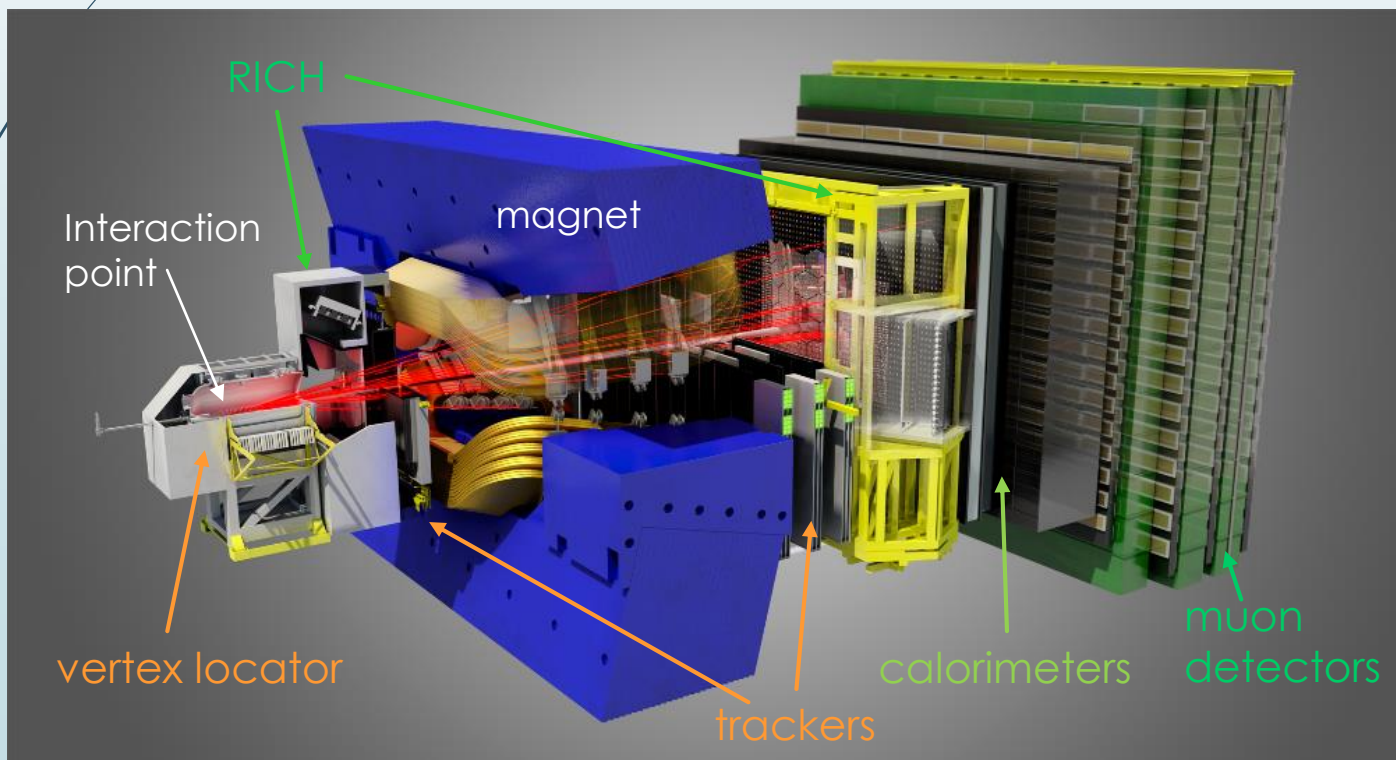
On behalf of the LHCb collaboration

PHOTON 2019

3-7 June, INFN – LFN, Frascati

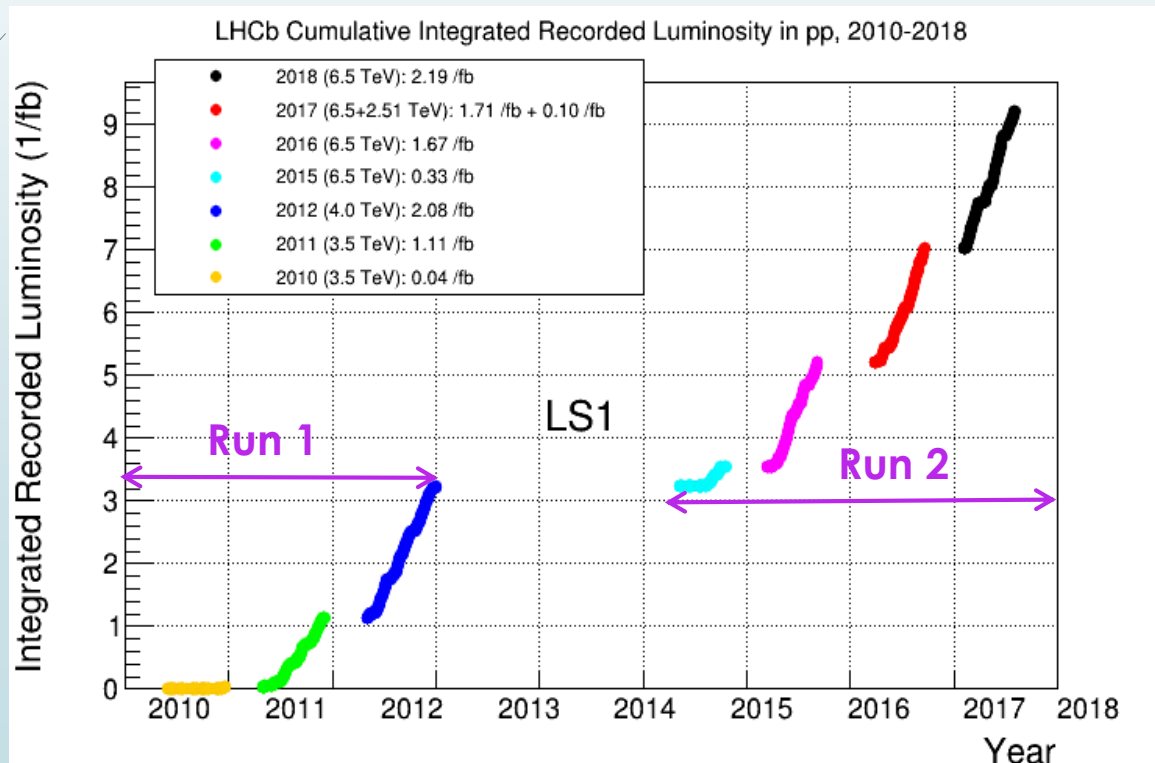
The LHCb detector

- Forward arm spectrometer to study b- and c-hadron decays ($2 < \eta < 5$).
 - Good vertex and impact parameter resolution ($\sigma(IP) = 15 + 29/p_T \text{ } \mu\text{m}$).
 - Excellent momentum resolution ($\sigma(mB) \sim 25 \text{ MeV}/c^2$ for 2-body decays).
 - Excellent particle identification (μ ID 97% for $(\pi \rightarrow \mu)$ misID of 1-3%).
 - Versatile and efficient trigger.



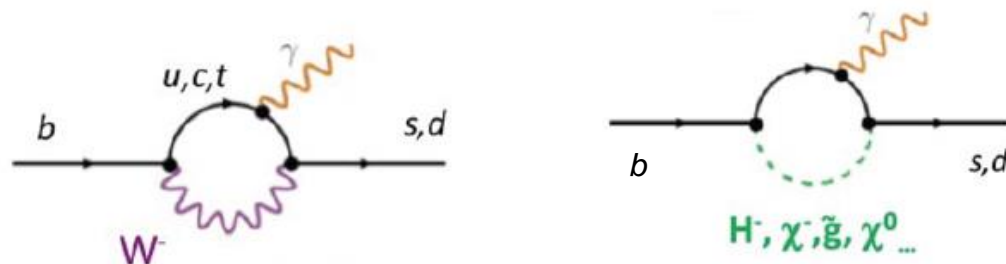
LHCb dataset

- **Run 1** (2010-2012): 7/8 TeV pp collisions; $\int \mathcal{L} = 3 \text{ fb}^{-1}$
- **Run 2** (2015-2018): 13 TeV pp collisions; $\int \mathcal{L} = 6 \text{ fb}^{-1}$
 - Production rates a factor 2 larger \rightarrow x 4 larger yields.



Introduction to radiative B decays

- Radiative b-hadron decays correspond to $b \rightarrow q\gamma$ ($q = d, s$) transitions.
- These Flavour-Changing Neutral Current (FCNC) transitions are forbidden at tree level in the Standard Model (SM).
 - Branching fractions, $\text{BR} < 10^{-4}$.
- Sensitive probes of New Physics (NP) through the study of branching fractions, angular observables, CP asymmetries and measurements of the polarisation of the photon emitted in the decay.
 - Indirect searches can probe NP at larger energy scales.



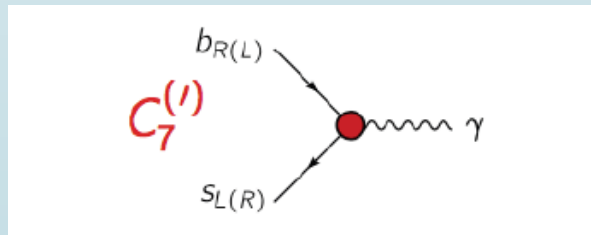
- FCNC are described by an effective Hamiltonian in the form of an Operator Product Expansion, which allows to identify the types of operators (O_i) that enter in each transition, along with their corresponding Wilson coefficients (C_i).

$$\mathcal{H}_{eff} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i \{C_i O_i + \underbrace{C'_i O'_i}_{\text{Right-handed part (supressed in SM)}}\} + \sum_i \frac{C_i^{NP}}{\Lambda^2} O_i^{NP}$$

Right-handed part
(supressed in SM)

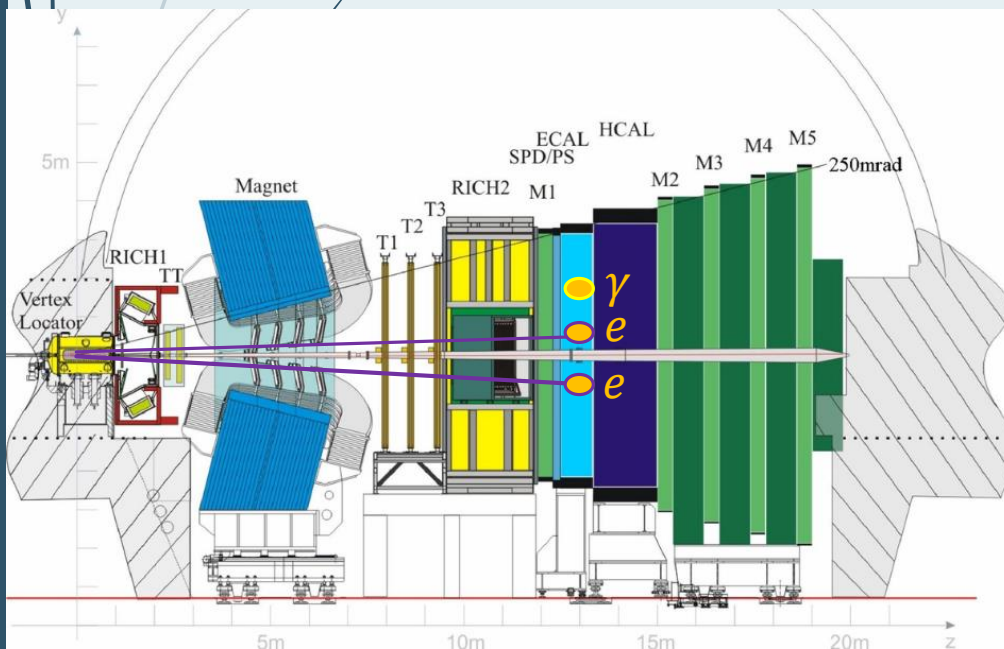
i=1, 2	Tree
i=3-6, 8	Gluon penguin
i=7	Photon penguin
i=9, 10	Electroweak penguin
i=5	Higgs (scalar) penguin
i=P	Pseudoscalar penguin

- Leading operator for radiative decays is the electromagnetic operator O_7 .
- Look for deviations of the values of the Wilson coefficients with respect to the SM predictions as a sign of NP.



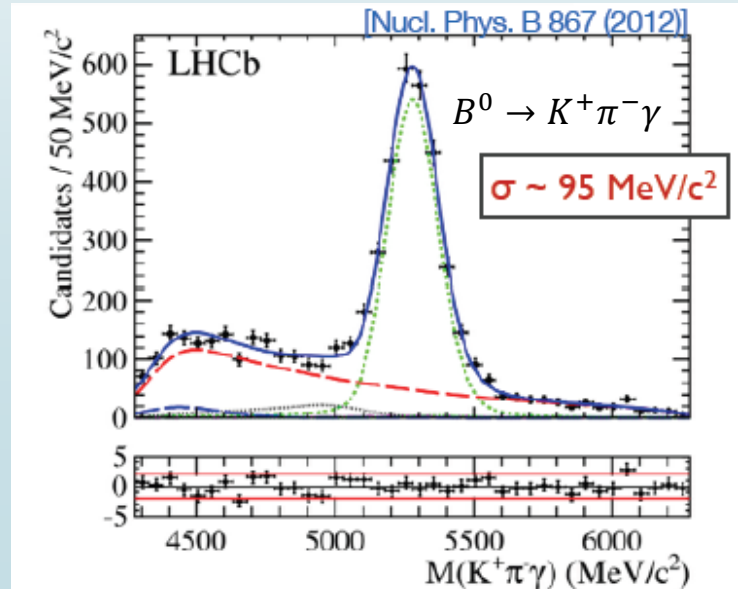
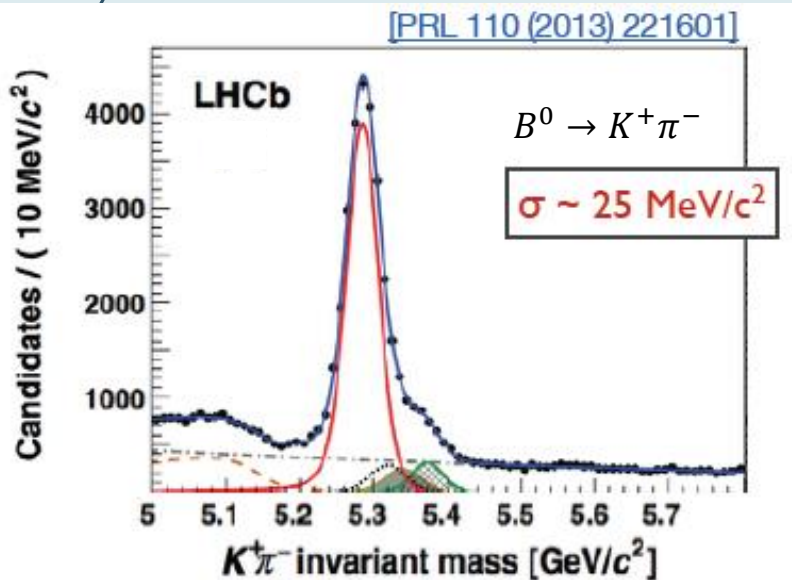
Photon reconstruction

- Photons are mainly reconstructed from energy deposits in the Electromagnetic Calorimeter (of Shashlik type).
 - Or as converted photons $\gamma \rightarrow e^+e^-$ with the tracking system. Better mass resolution (around factor 3), but much lower efficiency than calorimetric photons (around factor 15 in yield).

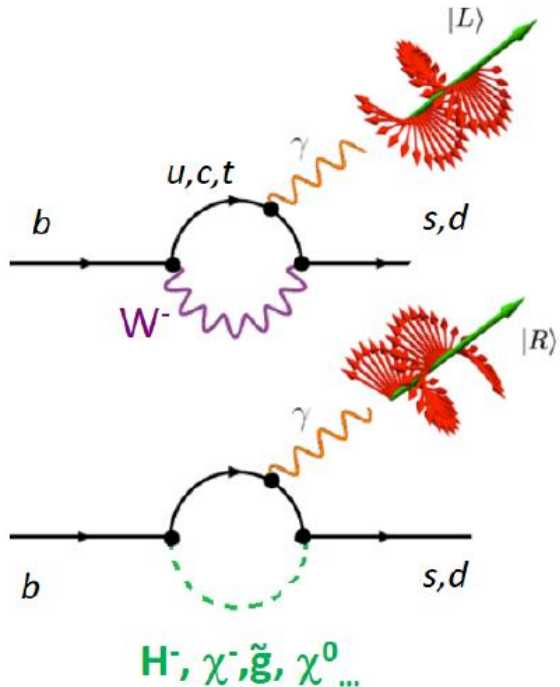


- Due to trigger constraints and large combinatorics the radiative decays mostly rely on high p_T photons.
 - $E_T > 2.5$ or 3 GeV;
 - Efficiency $\sim 30\text{-}40\%$ (vs 80-90% dimuon channels).

- Decays with photons in the final state are more challenging due to:
 - no constraint on vertexing from γ ,
 - large photon multiplicity and
 - limited mass resolution (dominated by the photon reconstruction).
- Radiative reconstruction rates in Run-1:
 - $B^0 \rightarrow K^{*0}\gamma$: ~ 10000 candidates/ fb^{-1}
 - $B_s^0 \rightarrow \phi\gamma$: ~ 1500 candidates/ fb^{-1}



Photon polarisation



- Due to the chiral structure of the W boson, in SM the photon polarisation is predominantly left-handed.

- Amplitude suppressed by m_q/m_W with m_q the mass of the right-handed quark.

$$\frac{C'_7}{C_7} = \frac{A_R(b_L \rightarrow s_R \gamma_R)}{A_L(b_R \rightarrow s_L \gamma_L)} \approx \frac{m_s}{m_b} = 0.02$$

[Descotes-Genon et al., JHEP 06 (2011) 099]

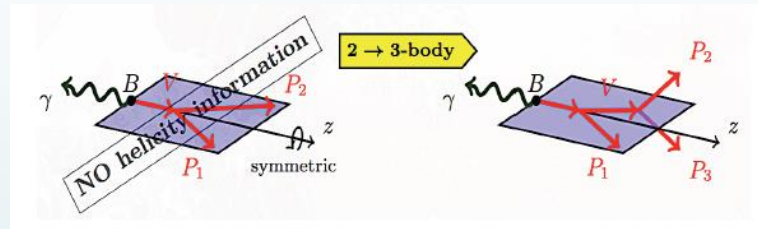
- Right-handed component could be enhanced in NP models.

- For instance, $\left| \frac{A_R}{A_L} \right|$ up to 0.5 in LRSM.

[Atwood et al., PRL 79 (1997) 185-188]

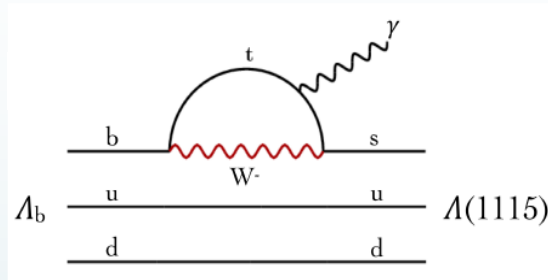
[Yu et al., JHEP 12 (2013) 102]

- Experimentally, the photon polarisation can be extracted from:
 - Angular distribution of radiative decays with 3 charged final state particles, e.g.
 - $B^+ \rightarrow K^+ \pi^+ \pi^- \gamma$ [Phys. Rev. Lett. 112 \(2014\) 161801](#) ◀ First observation that the photon is polarised



- Transverse asymmetry in $B^0 \rightarrow K^{*0} e e$. [JHEP 04 \(2015\) 064](#)
 - Angular analysis at low q^2 of the virtual photon decay (pollution from C_9, C_{10})
 - Time-dependent analysis of $B \rightarrow f_{CP} \gamma$ decays, e.g.
 - $B_s^0 \rightarrow \phi \gamma$ [Phys. Rev. Lett. 118 \(2017\) 021801](#)
 - $B^0 \rightarrow K_S \pi^+ \pi^- \gamma$
 - Angular analysis of radiative b-baryon decays, e.g.
 - $\Lambda_b \rightarrow \Lambda^{(*)} \gamma$
 - $\Xi_b \rightarrow \Xi^{(*)} \gamma$
- New results presented today

First observation of $\Lambda_b^0 \rightarrow \Lambda \gamma$



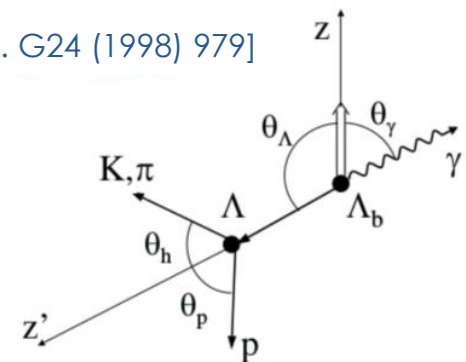
- Search of this decay using 2016 data (1.7 fb^{-1}).
- SM predictions for $\text{BR}(\Lambda_b^0 \rightarrow \Lambda \gamma)$ in the range $(0.06-1) \cdot 10^{-5}$.
 - [Eur. Phys. J. C59 (2009) 861, JHEP 12 (2011) 67, Commun. Theor. Phys. 58 (2012) 872, PRD 96 (2017) 053006]
- Sensitivity to C_7, C_7' through angular distributions.

$$\Gamma(\theta_p) = \frac{1}{4} (1 - \alpha_\gamma \alpha_\Lambda \cos \theta_p)$$

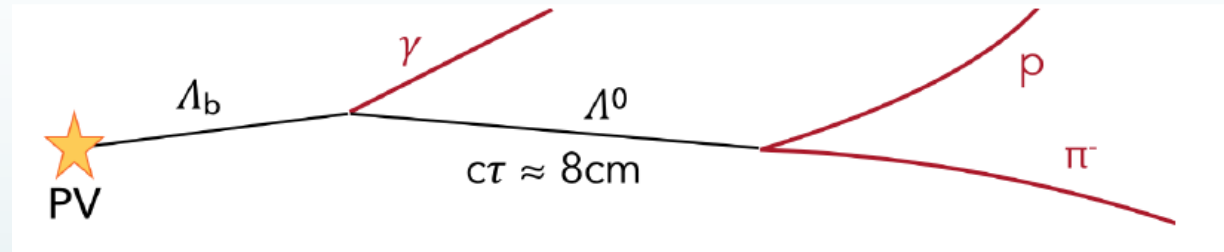
$$\alpha_\Lambda = 0.642 \pm 0.013 \text{ [PDG 2016]}$$

Experimental prospects at LHCb:
[arXiv:1902.04870v2]

[J. Phys. G24 (1998) 979]



- Challenge: no Λ_b vertex because Λ is long lived and no direction from the γ cluster.



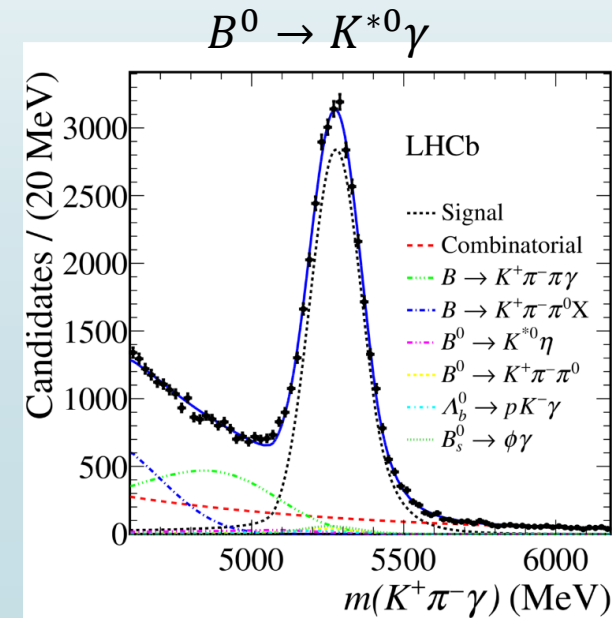
- Signal/background separation using XGBoost.
- Used $B^0 \rightarrow K^{*0}\gamma$ as normalisation mode.
 - 32670 \pm 290 candidates.

$$\frac{N(\Lambda_b^0 \rightarrow \Lambda \gamma)}{N(B^0 \rightarrow K^{*0} \gamma)} = \frac{f_{\Lambda_b^0}}{f_{B^0}} \times \frac{\mathcal{B}(\Lambda_b^0 \rightarrow \Lambda \gamma)}{\mathcal{B}(B^0 \rightarrow K^{*0} \gamma)} \times \frac{\mathcal{B}(\Lambda \rightarrow p \pi^-)}{\mathcal{B}(K^{*0} \rightarrow K^+ \pi^-)} \times \frac{\epsilon(\Lambda_b^0 \rightarrow \Lambda \gamma)}{\epsilon(B^0 \rightarrow K^{*0} \gamma)}$$

From LHCb measurement [arXiv:1902.06794](https://arxiv.org/abs/1902.06794)

From PDG

From simulation and calibration samples

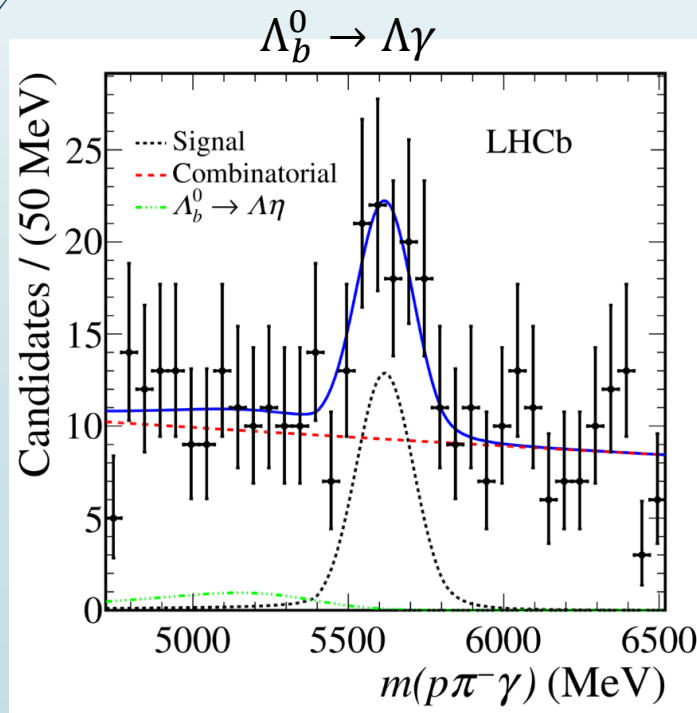


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- Signal excess with 5.6σ significance.
 - 65 ± 13 candidates.
 - **First observation!**
- Branching fraction within the range of SM predictions.

$$\mathcal{B}(\Lambda_b \rightarrow \Lambda \gamma) = (7.1 \pm 1.6 \pm 0.6 \pm 0.7) \cdot 10^{-6}$$

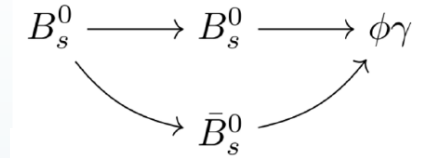
↑ statistical
 ↑ systematic
 ↑ systematic from external measurements



Systematic uncertainties

Source	Uncertainty (%)
Data/simulation agreement	7.7
Λ_b^0 fit model	3.0
$B^0 \rightarrow K^{*0} \gamma$ backgrounds	2.7
Size of simulated samples	1.7
Efficiency ratio	1.4
Sum in quadrature	9.0
$f_{\Lambda_b^0}/f_{B^0}$	8.7
Input branching fractions	3.0
Sum in quadrature	9.2

Time-dependent analysis of $B_s^0 \rightarrow \phi\gamma$



- Access to the photon polarisation through the measurement of mixing-induced and CP-violating observables in the decay rate.

$$\Gamma(t) \propto e^{-\Gamma_s t} \left[\cosh\left(\frac{\Delta\Gamma_s t}{2}\right) - A^\Delta \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) \pm C \cos(\Delta m_s t) \mp S \sin(\Delta m_s t) \right]$$

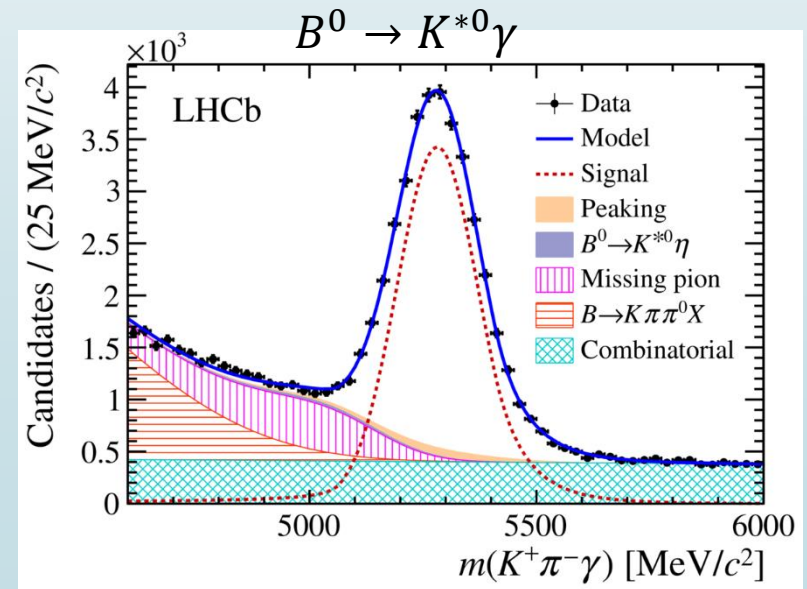
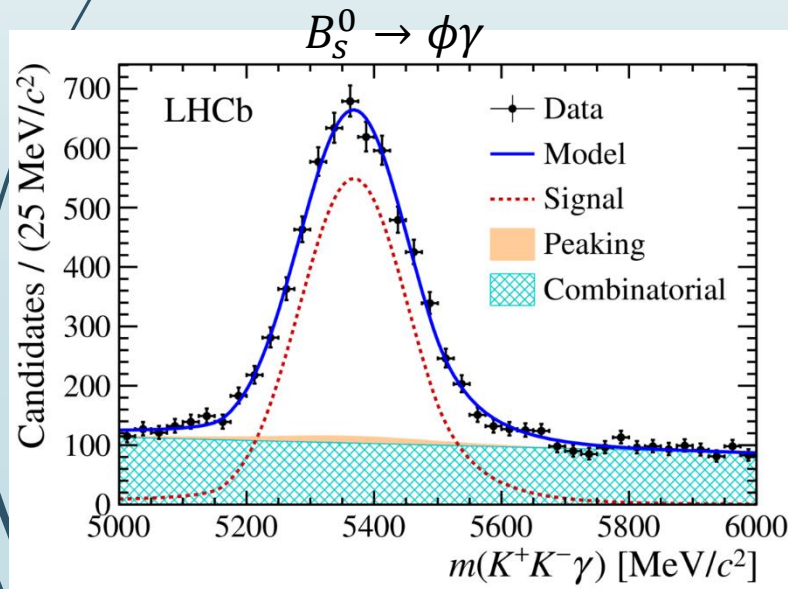
- A^Δ and S are sensitive to the photon polarisation (C_7, C_7'), while C is related to direct CP asymmetry.
 - $\Delta\Gamma_s$ and Δm_s are the decay width and mass differences between B_s^0 CP eigenstates.
- Previous analysis without flavour tagging information (no separation between B_s^0 and \bar{B}_s^0) with Run-1 data:

$$A^\Delta = -0.98_{-0.52}^{+0.46} {}_{-0.20}^{+0.23}$$

[Phys. Rev. Lett. 118 \(2017\) 021801](#)

- In agreement with SM prediction $A_{SM}^\Delta = 0.047_{-0.025}^{+0.029}$ [PLB 664 (2008) 174]

- Now including flavour tagging information and re-analysing Run-1 data (3 fb^{-1}) with optimized reconstruction, selection and photon identification algorithms.
 - First measurement of C and S in B_s^0 decays.**
- $5110 \pm 90 \text{ } B_s^0 \rightarrow \phi \gamma$ candidates ($\sim 4.2\text{k}$ in previous analysis).
- Use $B^0 \rightarrow K^{*0} \gamma$ as control channel for the decay time acceptance.
 - 33860 ± 250 candidates.



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- Simultaneous unbinned ML fit to both decays, using per-candidate decay time resolution (σ_t) and flavour-tagging information (q, ω).

$$\Gamma_{B_S^0 \rightarrow \phi \gamma}(t') = \Gamma(t', q | \omega) \otimes \{A(t_i) R(t, t' | \sigma_t)\}$$

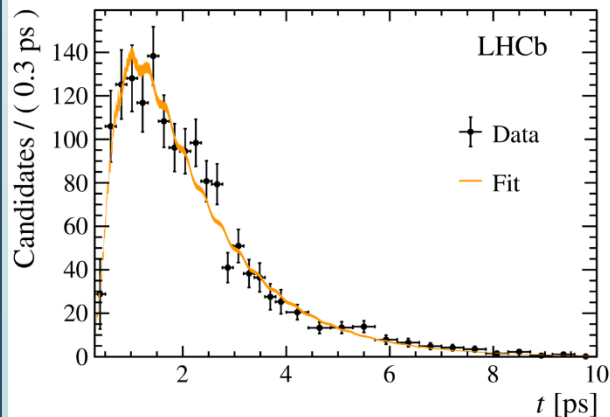
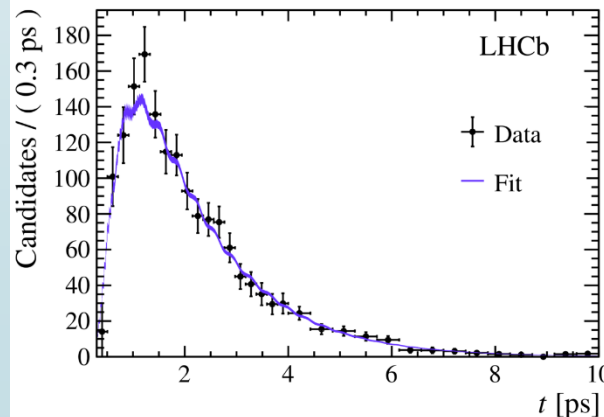
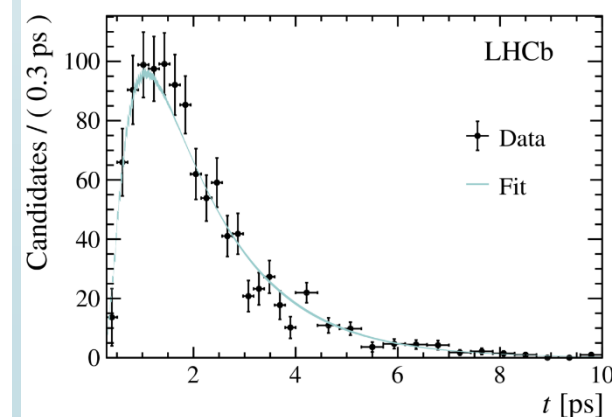
- Results:

$$S = 0.43 \pm 0.30 \pm 0.11$$

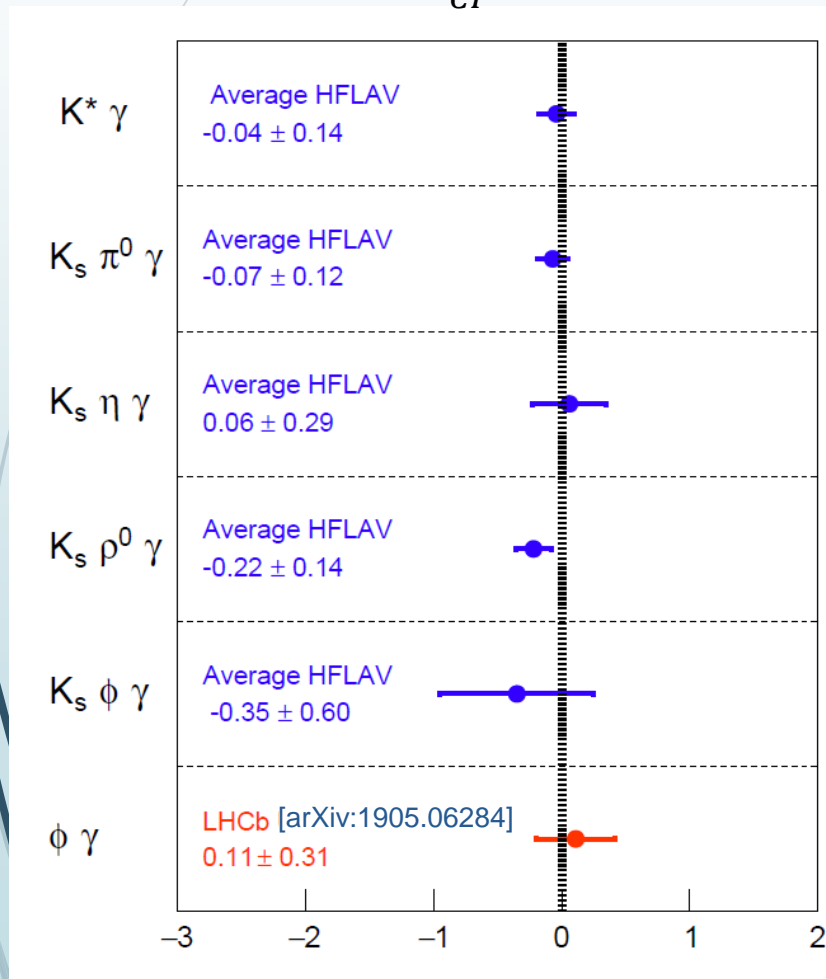
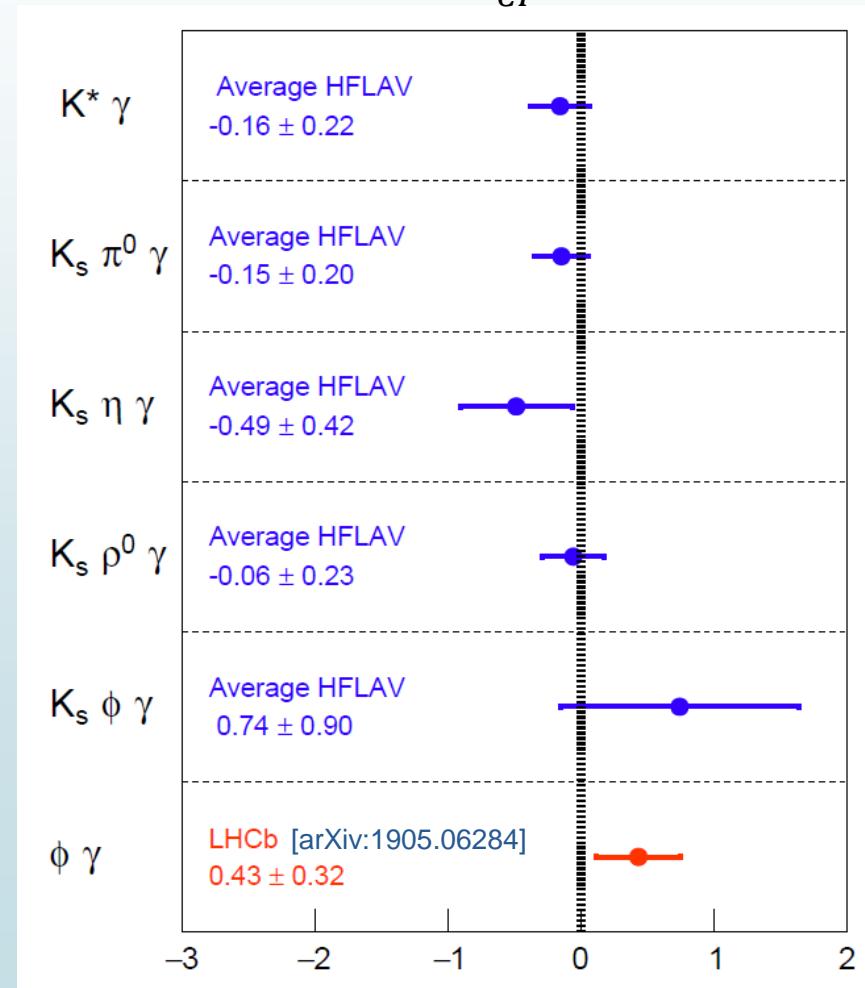
$$C = 0.11 \pm 0.29 \pm 0.11$$

$$A^{\Delta} = -0.67^{+0.37}_{-0.41} \pm 0.17$$

- Compatible with SM [PLB 664 (2008) 174-179] at 1.3, 0.3, 1.7 σ , respectively.
- Statistical limited. Dominant source of systematic uncertainties related to the determination of the decay-time acceptance.

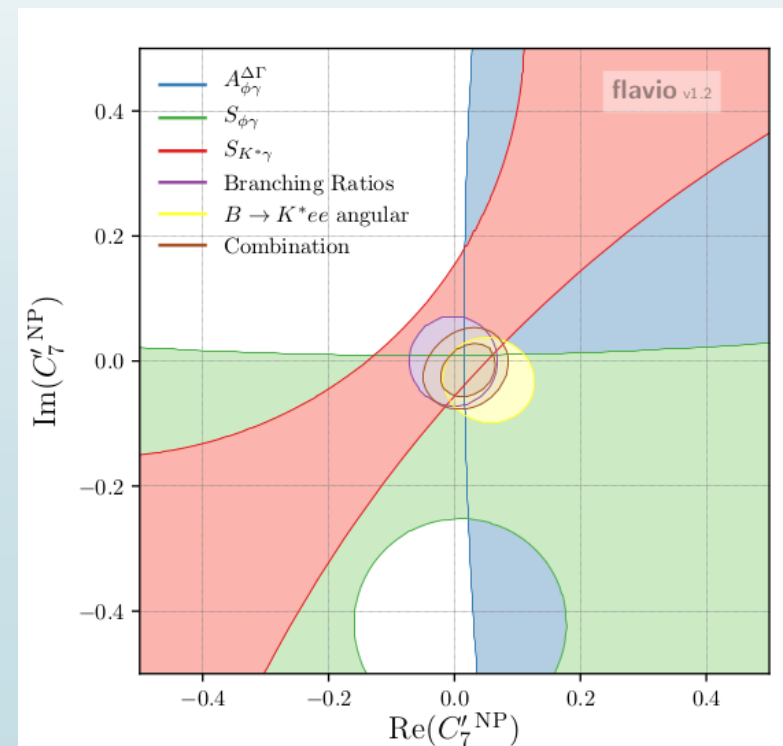
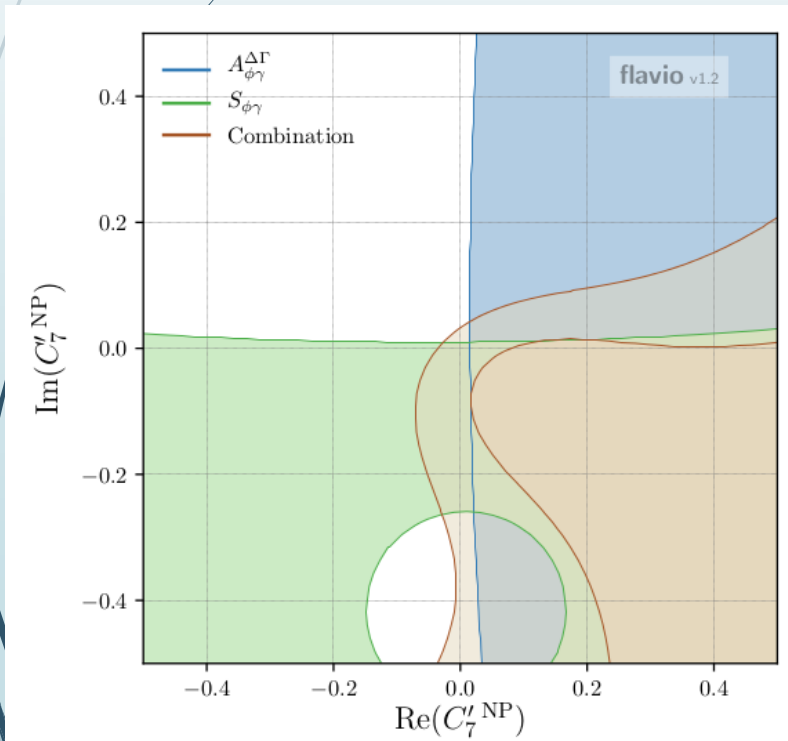
 $B_S^0 \rightarrow \phi \gamma$  $\bar{B}_S^0 \rightarrow \phi \gamma$ untagged $B_S^0 \rightarrow \phi \gamma$ 

- Competitive with previous measurements from B-factories.

 C_{CP}

 S_{CP}


NP constraints

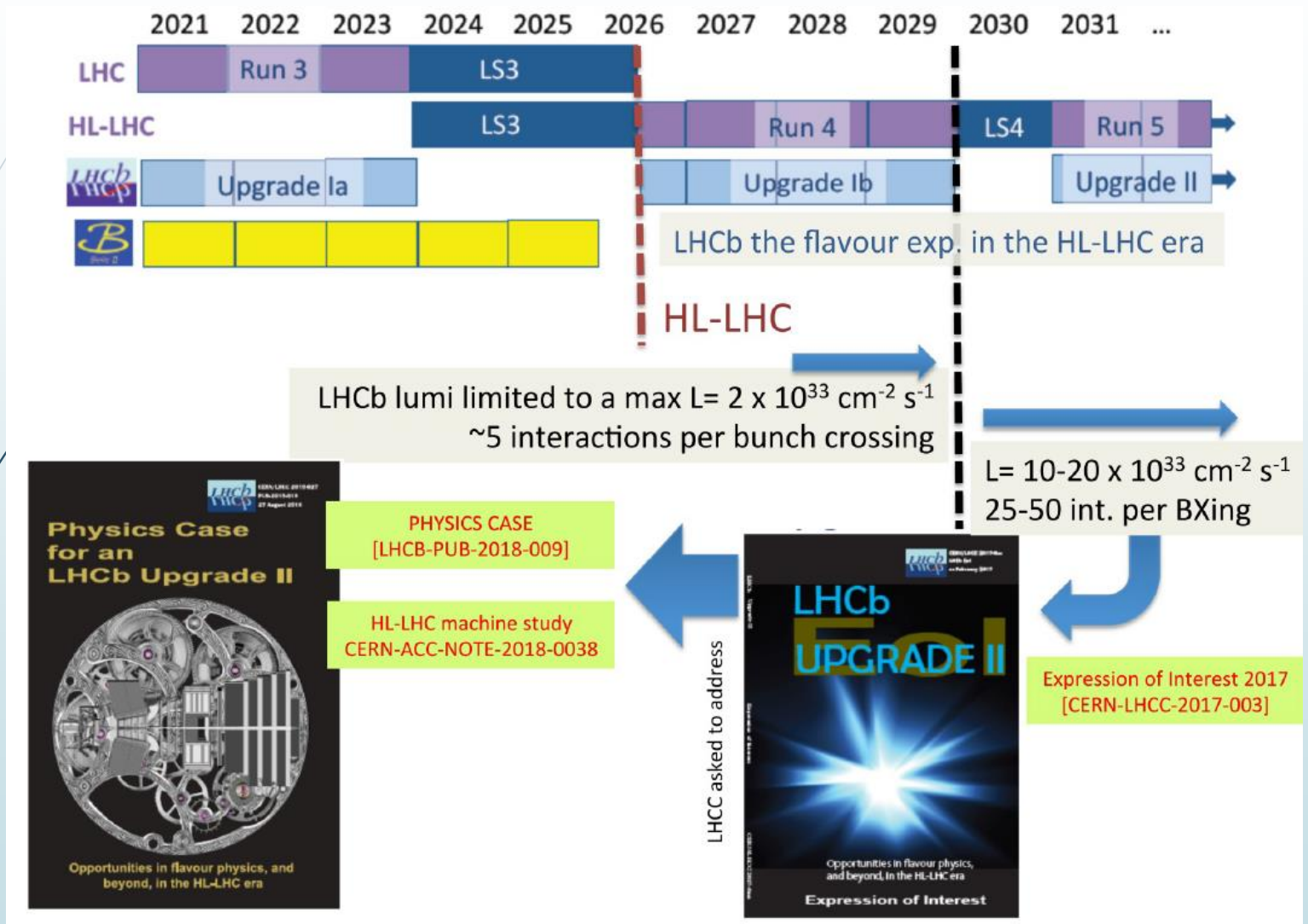
- A^Δ and S provide complementary constraints in the complex plane.
- Constraints to C'_7 using Flavio [arXiv:1810.08132]
 - Inputs: from LHCb [arXiv:1905.06284, JHEP 04 (2015) 064] and HFLAV.



Summary

- Radiative b-decays allow to probe NP at large energy scales through indirect measurements.
- The photon polarisation can be measured in several ways.
 - Interesting observable that puts constraints on C'_7 (right-handed part).
- Latest results from LHCb:
 - First observation of $\Lambda_b^0 \rightarrow \Lambda \gamma$, opening the possibility of measuring the photon polarisation in b-baryon decays.
 - Time-dependent analysis of $B_s^0 \rightarrow \phi \gamma$. Constraints on C'_7 and first measurement of C and S in a B_s^0 decay.
- Many more results with Run-2 data coming soon.

Backup slides



Physics case for and LHCb Upgrade II

Decay mode	Yield (300 fb ⁻¹)	Statistical sensitivity on measurement
$B_s \rightarrow \phi \gamma$	800k	0.02 on A^Δ
$B^0 \rightarrow K_s^0 \pi \pi \gamma$	200k	Competitive on S_{CP}
$B^+ \rightarrow K \pi \pi \gamma$		Photon polarisation <1%
$B^0 \rightarrow K^{*0} e e$	20k	2% on A_T
$\Lambda_b \rightarrow \Lambda^0 (p \pi) \gamma$	10k	$\alpha_\gamma < 2\%$
$\Xi_b \rightarrow \Xi^- (\Lambda^0 \pi) \gamma$		$\alpha_\gamma < 10\%$

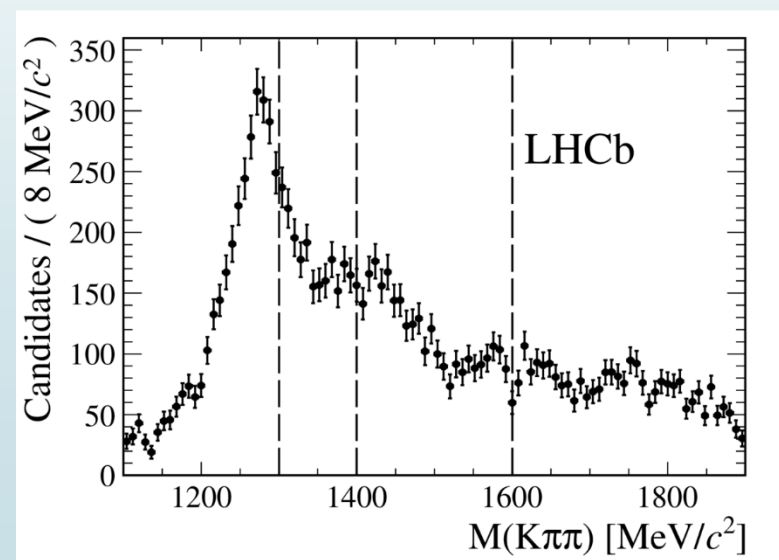
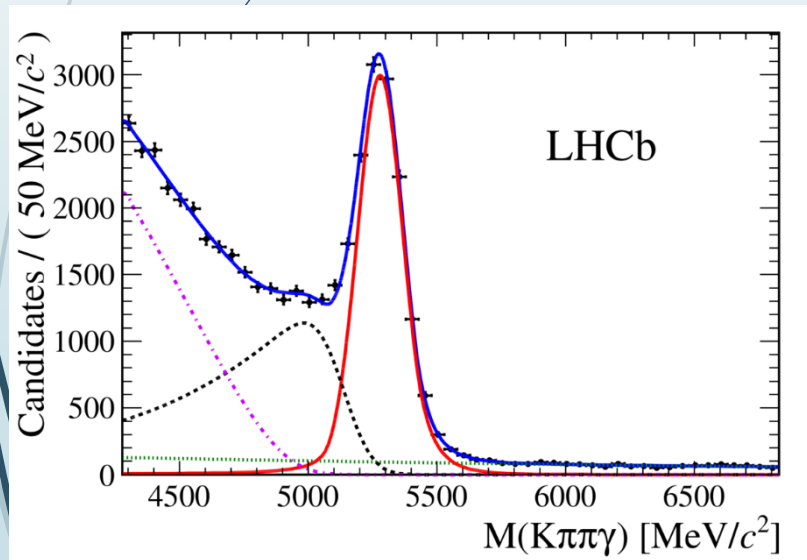
Previous published measurements

- Measurement of the ratio of branching fractions $\mathcal{B}(B^0 \rightarrow K^{*0}\gamma)/\mathcal{B}(B_s^0 \rightarrow \phi\gamma)$. [Phys. Rev. D 85 \(2012\) 112013](#)
- Measurement of the ratio of branching fractions $\mathcal{B}(B^0 \rightarrow K^{*0}\gamma)/\mathcal{B}(B_s^0 \rightarrow \phi\gamma)$ and the direct CP asymmetry in $B^0 \rightarrow K^{*0}\gamma$. [Nucl. Phys. B 867 \(2013\) 1](#)
- Observation of photon polarization in the $b \rightarrow s\gamma$ transition. [Phys. Rev. Lett. 112 \(2014\) 161801](#)
- Angular analysis of the $B^0 \rightarrow K^{*0}e^+e^-$ decay in the low- q^2 region. [JHEP 04 \(2015\) 064](#)
- Search for the rare decays $B^0 \rightarrow J/\psi\gamma$ and $B_s^0 \rightarrow J/\psi\gamma$. [Phys. Rev. D 92 \(2015\) 112002](#)
- First experimental study of photon polarization in radiative B_s^0 decays. [Phys. Rev. Lett. 118 \(2017\) 021801](#)

First observation of photon polarisation in $B^+ \rightarrow K^+ \pi^+ \pi^- \gamma$

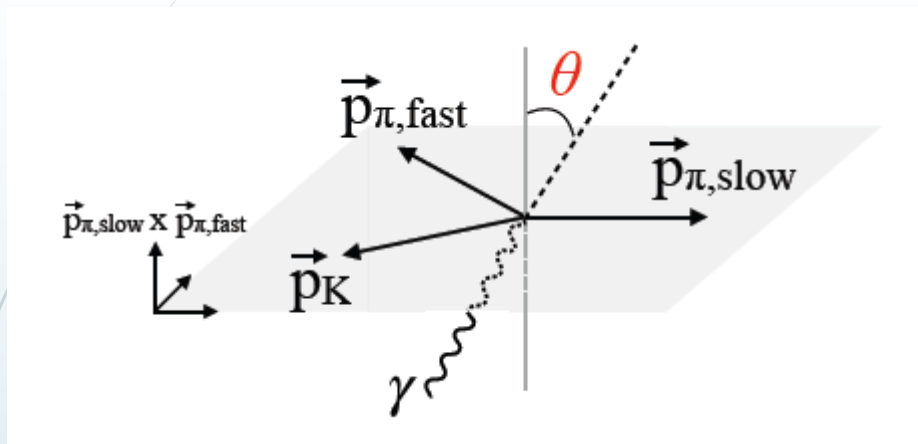
[Phys. Rev. Lett. 112 \(2014\) 161801](#)

- Run-1 data (3 fb^{-1}). $\sim 14\text{k}$ candidates with all intermediate resonances of $K\pi\pi$ in $[1.1, 1.9] \text{ GeV}/c^2$ mass range.



The photon polarisation can be inferred from the polarisation of the K resonance.

- Need three tracks in the final state to form a parity odd triple product.



[Gronau et al.,
PRL 88 (2002) 051802]

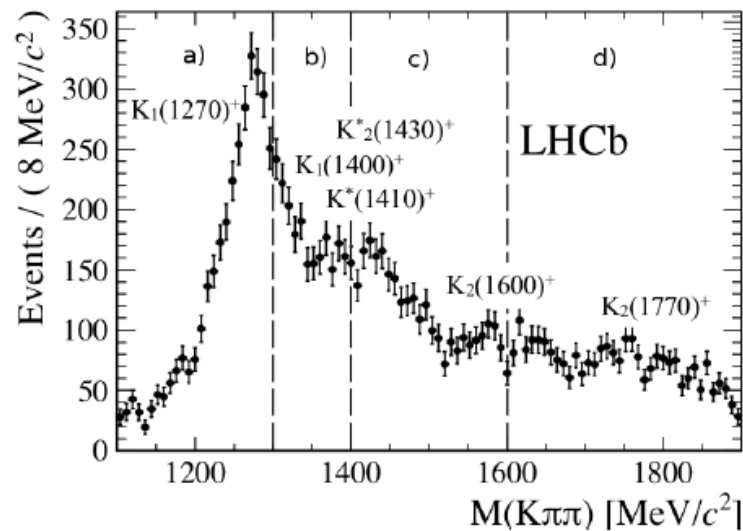
For a mixture of resonances:

$$\frac{d\Gamma}{ds ds_{13} ds_{23} d\cos\theta} \propto \sum_{i=0,2,4} a_i(s, s_{13}, s_{23}) \cos^i\theta + \lambda_\gamma \sum_{j=1,3} a_j(s, s_{13}, s_{23}) \cos^j\theta$$

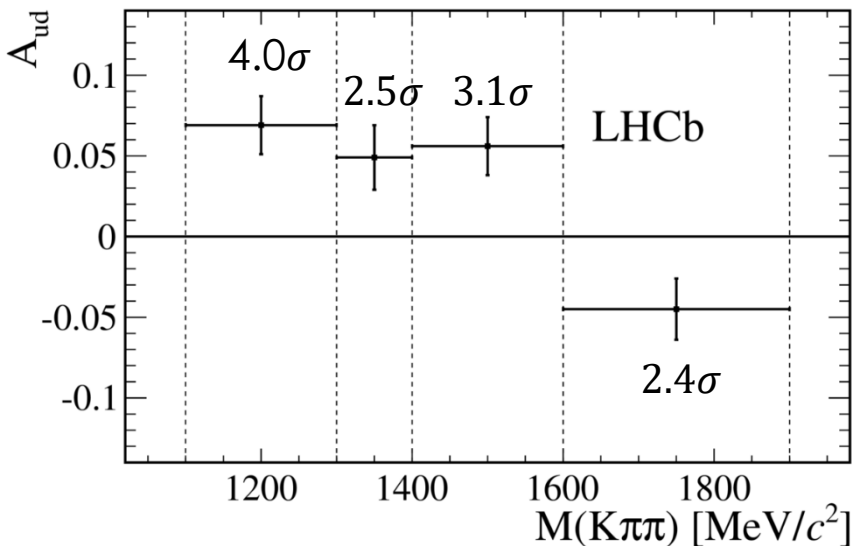
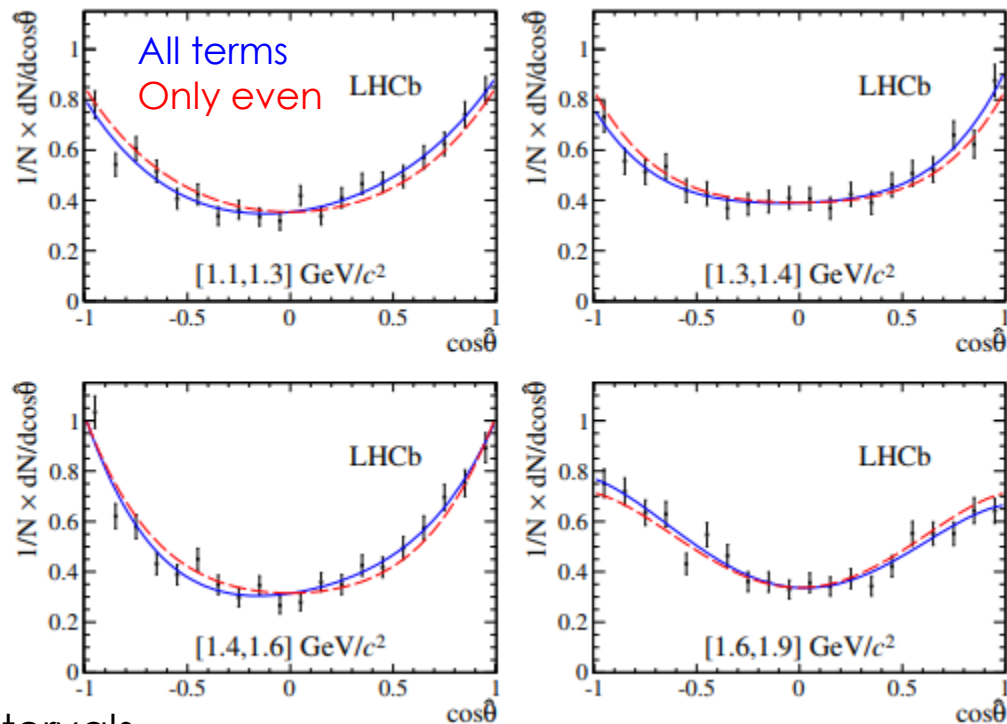
Define an up-down asymmetry, proportional to the photon polarisation parameter, λ_γ

[Gronau and Pirjol,
PRD 96 (2017) 013002]

$$\mathcal{A}_{\text{ud}} \equiv \frac{\int_0^1 d\cos\theta \frac{d\Gamma}{d\cos\theta} - \int_{-1}^0 d\cos\theta \frac{d\Gamma}{d\cos\theta}}{\int_{-1}^1 d\cos\theta \frac{d\Gamma}{d\cos\theta}} = C\lambda_\gamma$$



Asymmetry measured in four K_{res} mass intervals



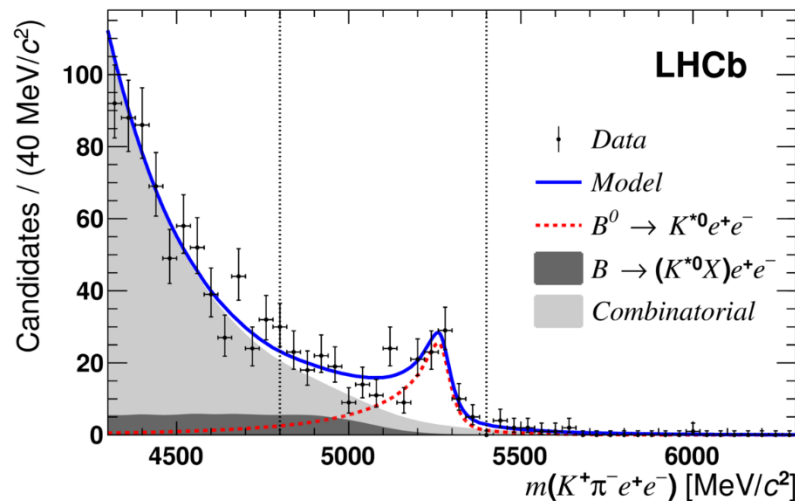
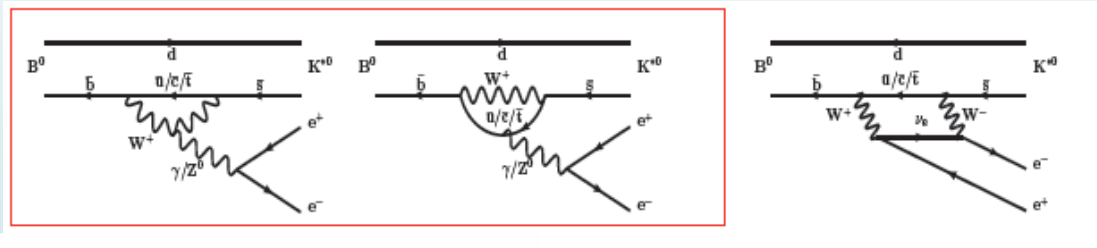
First direct observation of a non-zero photon polarisation in $b \rightarrow s\gamma$, with 5.2 σ significance.

Theory input and full amplitude analysis (ongoing) are needed to determine the exact value of the polarisation.

$B^0 \rightarrow K^{*0} ee$ in the low q^2 region

[JHEP 04 \(2015\) 064](#)

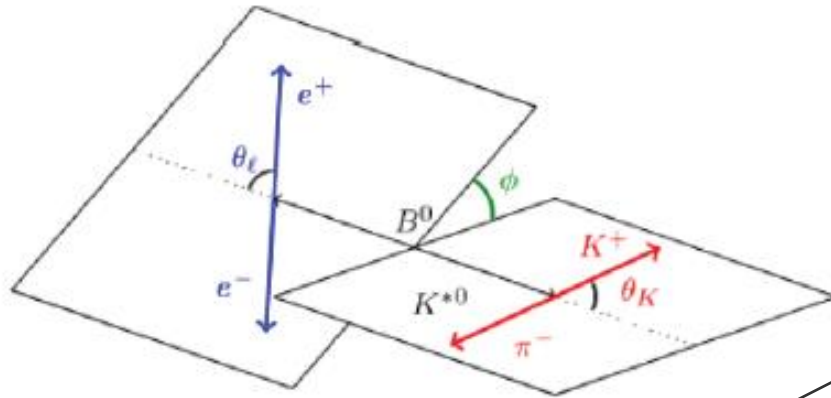
- Run-1 data (3 fb^{-1}). 150 $B^0 \rightarrow K^{*0} e^+ e^-$ candidates.
- Angular analysis at very low $m(ee)$. Sensitive to photon polarisation in the limit $m(ee) \rightarrow 0$.



Used $B^0 \rightarrow J/\psi(e^+ e^-) K^{*0}$ as control channel

Simplified angular distribution with 4 observables.

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Angular analysis with $\cos \theta_l$, $\cos \theta_K$, $\tilde{\phi}$

$$\frac{1}{d(\Gamma + \bar{\Gamma})/dq^2} \frac{d^4(\Gamma + \bar{\Gamma})}{dq^2 d\cos \theta_\ell d\cos \theta_K d\tilde{\phi}} = \frac{9}{16\pi} \left[\frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \left(\frac{1}{4}(1 - F_L) \sin^2 \theta_K - F_L \cos^2 \theta_K \right) \cos 2\theta_\ell + \frac{1}{2}(1 - F_L) A_T^{(2)} \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\tilde{\phi} + (1 - F_L) A_T^{\text{Re}} \sin^2 \theta_K \cos \theta_\ell + \frac{1}{2}(1 - F_L) A_T^{\text{Im}} \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\tilde{\phi} \right].$$

Sensitive to photon polarisation

[Krüger, Matias,
Phys. Rev. D 71 (2005) 094009]
[Becirevic, Schneider,
J. Nucl. Phys. B 09 (2011) 004]

$$F_L = \frac{|A_0|^2}{|A_0|^2 + |A_{||}|^2 + |A_{\perp}|^2}$$

$$A_T^{(2)} = \frac{|A_{\perp}|^2 - |A_{||}|^2}{|A_{\perp}|^2 + |A_{||}|^2}$$

$$A_T^{\text{Re}} = \frac{2\text{Re}(A_{||L} A_{\perp L}^* + A_{||R} A_{\perp R}^*)}{|A_{||}|^2 + |A_{\perp}|^2}$$

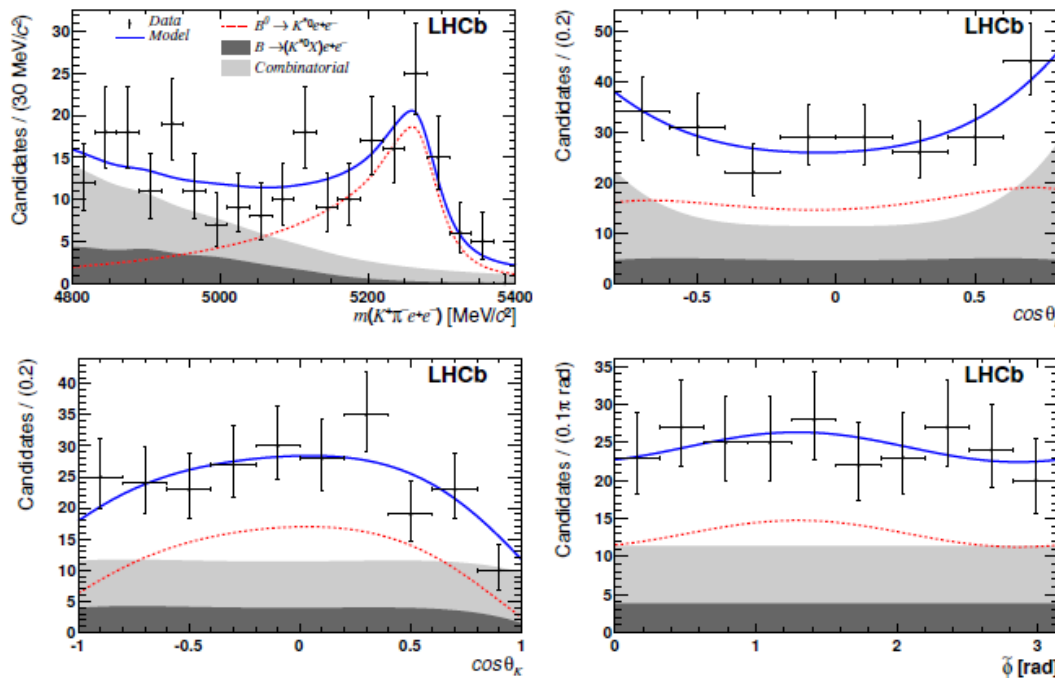
$$A_T^{\text{Im}} = \frac{2\text{Im}(A_{||L} A_{\perp L}^* + A_{||R} A_{\perp R}^*)}{|A_{||}|^2 + |A_{\perp}|^2},$$

Approximation:

$$\lim_{q^2 \rightarrow 0} A_T^{(2)} = \frac{2\text{Re}(C_7 C_7')}{|C_7|^2 + |C_7'|^2}$$

$$\lim_{q^2 \rightarrow 0} A_T^{\text{Im}} = \frac{2\text{Im}(C_7 C_7')}{|C_7|^2 + |C_7'|^2}$$

Angular analysis for di-lepton q^2 in $[20, 1120]$ MeV/c^2 .
 Reconstruction challenges with electrons:
 lower reconstruction efficiencies and bremsstrahlung recovery.



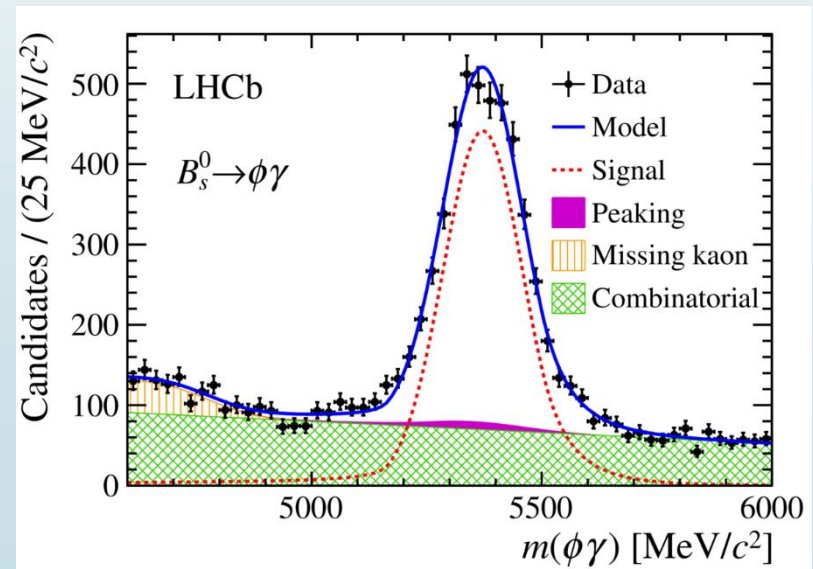
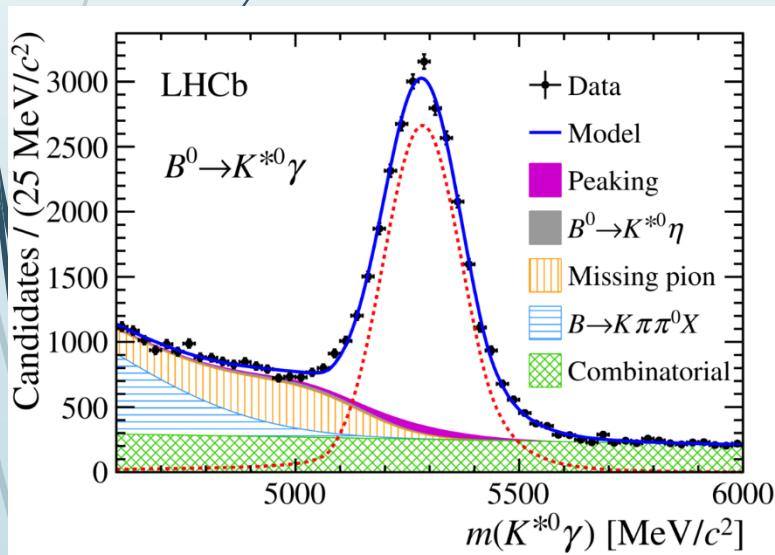
$$\begin{aligned}
 F_L &= 0.16 \pm 0.06 \pm 0.03 \\
 A_T^{(2)} &= -0.23 \pm 0.23 \pm 0.05 \\
 A_T^{\text{Im}} &= +0.14 \pm 0.22 \pm 0.05 \\
 A_T^{\text{Re}} &= +0.10 \pm 0.18 \pm 0.05,
 \end{aligned}$$

Compatible with SM predictions

Time-dependent decay rate of $B_s^0 \rightarrow \phi\gamma$

[Phys. Rev. Lett. 118 \(2017\) 021801](#)

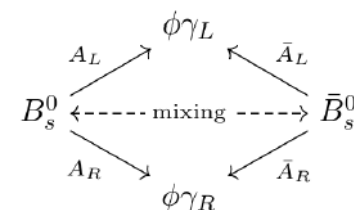
- Run-1 data (3 fb^{-1}). 4.2k $B_s^0 \rightarrow \phi\gamma$ candidates.
- 25.8k $B^0 \rightarrow K^{*0}\gamma$ used as control sample for decay time acceptance (simultaneous fit).



Direct access to the photon polarisation through the time-dependent decay rate of $B_s^0 \rightarrow \phi\gamma$.

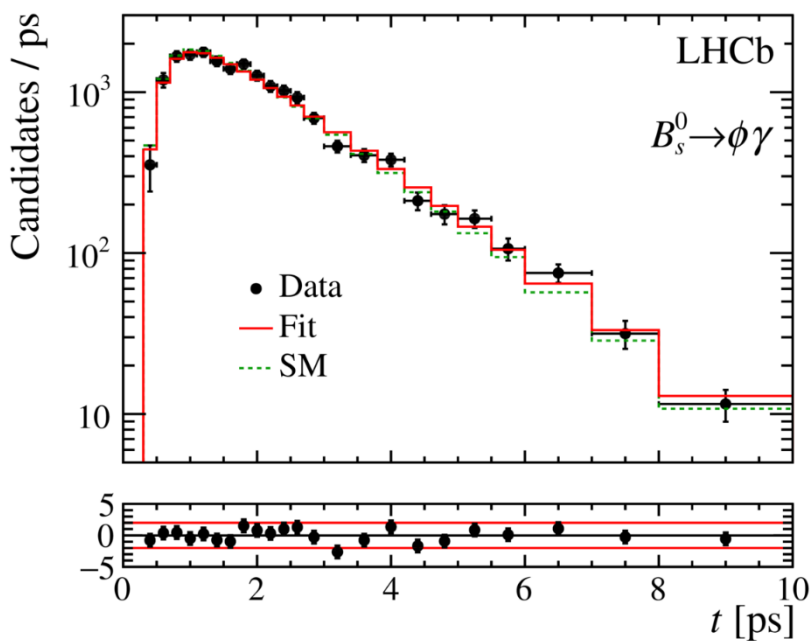
$$\Gamma(B_s^0 \rightarrow f^{CP}\gamma) \sim e^{-\Gamma(s)t} \left\{ \cosh \frac{\Delta\Gamma(s)t}{2} - \mathcal{A}^\Delta \sinh \frac{\Delta\Gamma(s)t}{2} \pm \mathcal{C} \cos \Delta m_{(s)}t \mp \mathcal{S} \sin \Delta m_{(s)}t \right\}.$$

Mixing/decay interference



$$\mathcal{A}_{\phi\gamma}^\Delta \simeq \frac{\text{Re}(e^{-i\phi_s} C_7 C_7')}{|C_7|^2 + |C_7'|^2} \quad S_{\phi\gamma} \simeq \frac{\text{Im}(e^{-i\phi_s} C_7 C_7')}{|C_7|^2 + |C_7'|^2}$$

An untagged measurement was performed, so only \mathcal{A}^Δ was measured:



$$\mathcal{A}^\Delta = -0.98^{+0.46}_{-0.52}(\text{stat.})^{+0.23}_{-0.20}(\text{syst.})$$

Compatible with SM at 2.6σ

SM prediction: $\mathcal{A}^\Delta = 0.047^{+0.029}_{-0.025}$

[Muheim et al., PRB 664 (2008) 174]

In Left-Right Symmetric model $\mathcal{A}^\Delta \sim 0.7$