Two level systems and ultrastable glasses

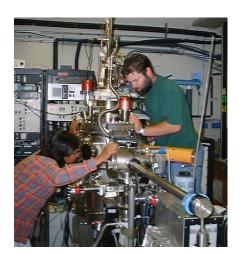
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UHV growth chambers for prototyping materials



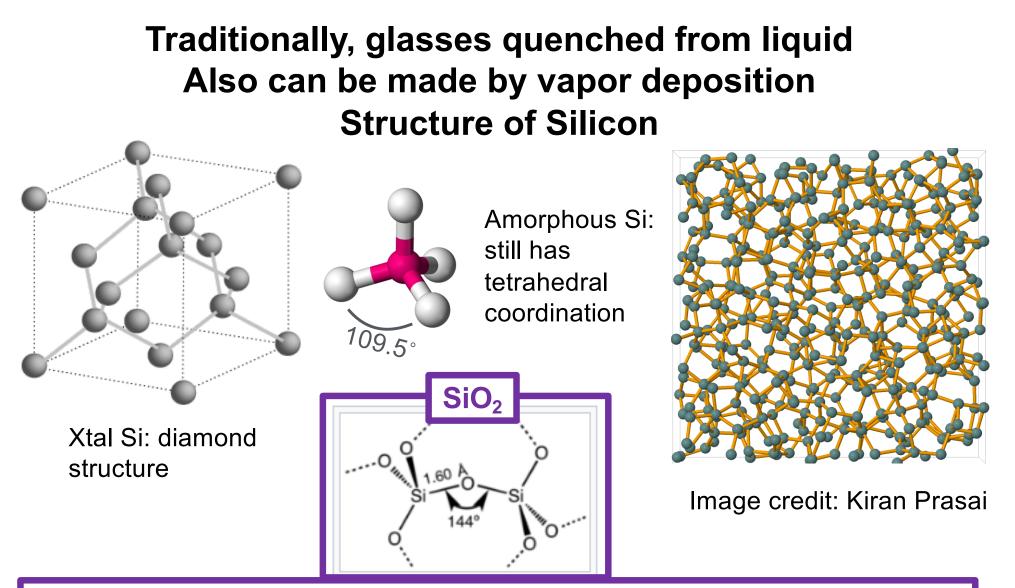






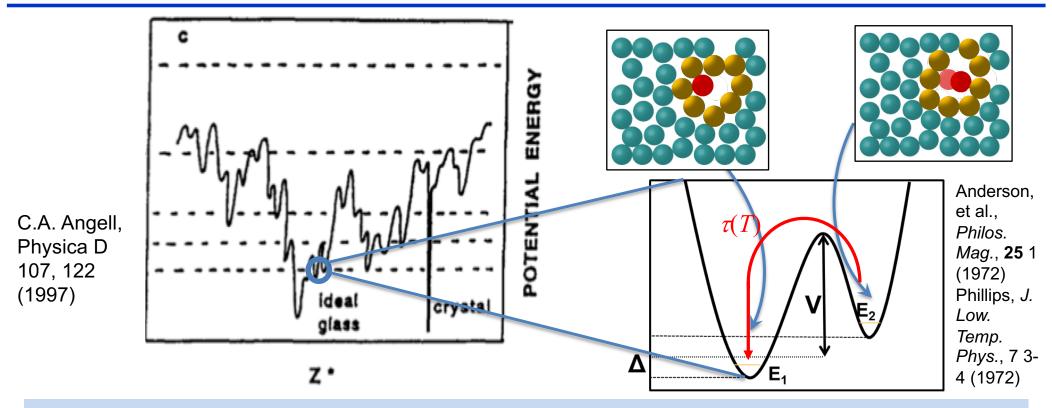


Current mirror coatings are amorphous No long range structural order, but have short range order



O-Si-O bonds are fixed angle but Si-O-Si angle is quite floppy – different energy scales

Energy landscape of configurations: "nearby" minima lead to tunneling *or* thermally-activated motion of groups of atoms



Two-Level Systems from neighboring energy minima in structural landscape:

- At low T, atomic structure **tunnels** between these $\geq \mu eV$ energy splitting $E_{1, 2} \pm \Delta$
- At higher T, atomic motion is thermally activated, requiring k_BT ~ barrier height V

In both cases, atomic motion leads to dissipation (thermal noise)
 For a single V, dissipation at frequency ω will have a peak at temperature T
 > at 1 kHz, 0.5 eV barrier heights has a peak ~ room temperature
 50 meV barrier heights has a peak ~ 30K

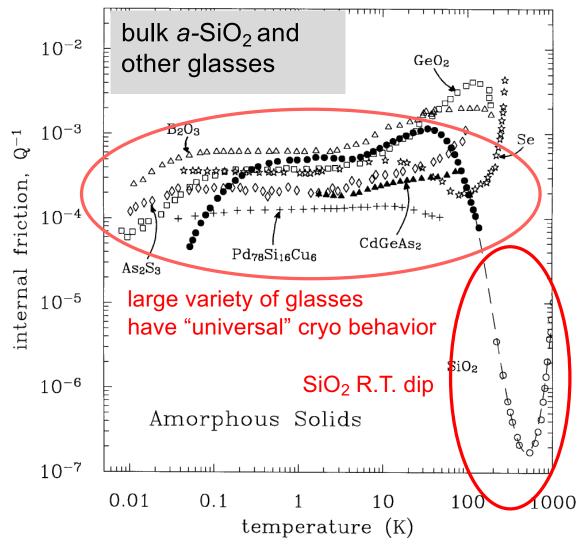
A distribution of µeV tunneling-induced energy splitting leads to T-independent losses 3

"Universal" mechanical losses at low T (tunneling) Higher T more variable, including peaks (thermally activated)

Internal friction $Q^{-1}(T)$ **low T plateau** due to *tunneling two level systems* (TLS)-phonon interactions Q_o^{-1} proportional to \overline{P} (density of TLS) with *poorly understood TLS* – phonon coupling parameter γ

Lower T data (below plateau): TLS model predicts drop off

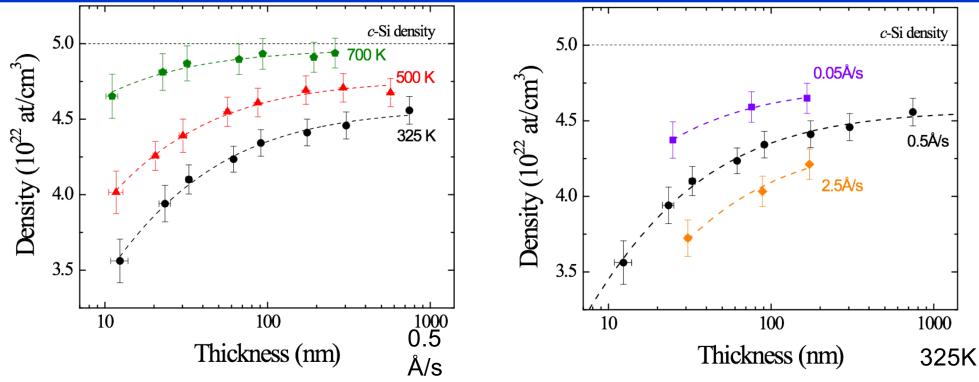
- Higher temp thermally activated regime: a peak implies a single barrier V (or narrow distribution)
- a-SiO₂ dip implies **a** gap in the distribution of barriers, followed by α relaxation approaching melting



K.A. Topp, Z. Physik B Condensed Matter 101 235-45 (1996)

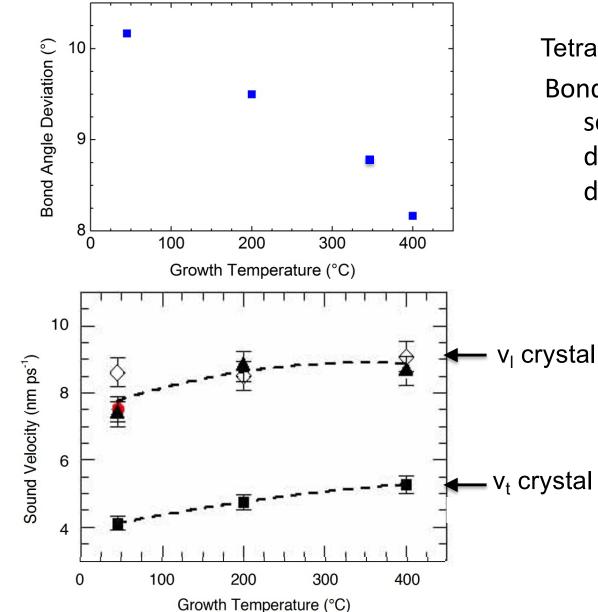
* Berret & Meissner Z.Phys 1988

Growth parameters substantially modify amorphous Si film density and some measures of structure



- Thickness, growth temperature T_g, and growth rate affect film density and roughness; room T growth flattest for all thicknesses; higher T_g thin is flat, roughens with thickness (1.5 nm RMS at 300 nm)
- Thinner, low growth T, high growth rate films are less dense
- On what length scale(s) do density changes occur? Little variation in dangling bond density, sound velocity, or macroscale structure, particularly with thickness
- Variations in bond angle disorder, medium range order, nanovoid size and number (Raman, Fluctuation Electron Microscopy, positron doppler broadening spectroscopy)
 Which of these matter to TLS? Appears that nanovoids may be most important

Amorphous Si: Disorder decreases with increasing growth T



Tetrahedral bond angle $109.5^{\circ} \pm \delta$ Bond angle disorder δ (from Raman scattering width of TO-like peak) decreases with increasing T_s. No dependence on film thickness

Longitudinal and transverse sound velocity v increases with increasing T_s

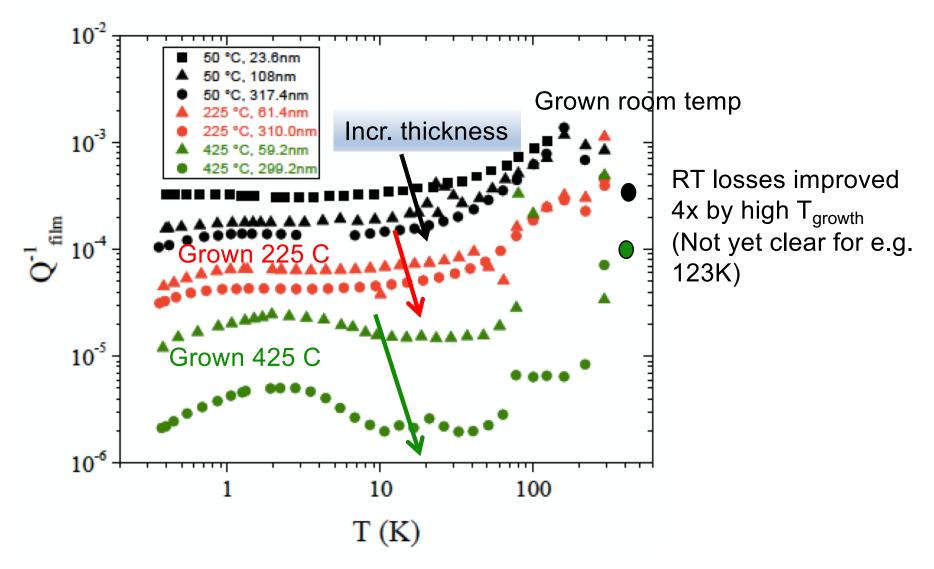
Elastic properties (shear modulus, sound velocity) soften with disorder in amorphous network

Independent of film thickness

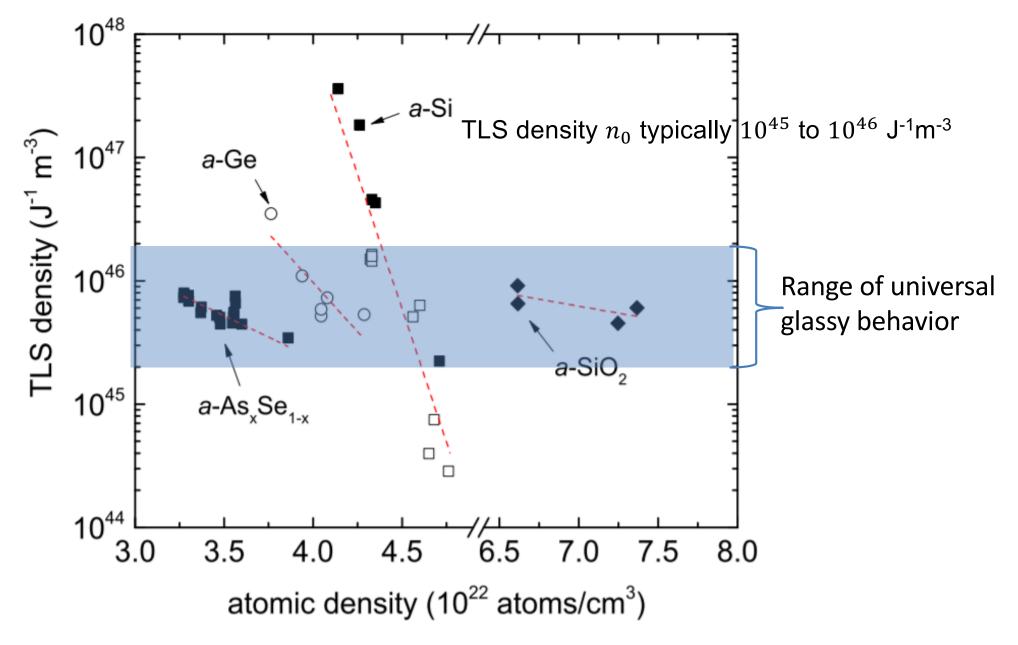
Open symbols: ~100nm films Closed symbols: ~300nm films **Amorphous Silicon losses:** growth T reduces low T losses Thin films are more lossy than thick films; correlates with atomic density

differences (thick films are denser)

Annealing reduces loss, but not much (at low T) compared to growth T effects

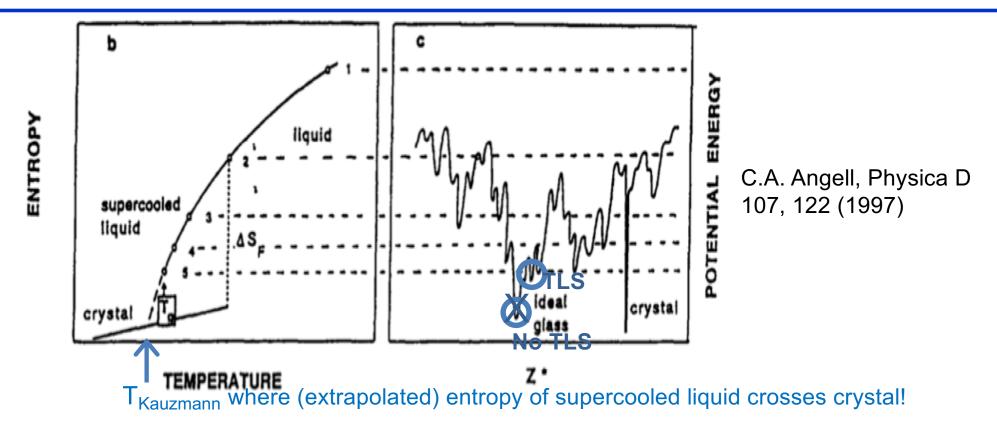


TLS depends on atomic density in *a*-Si and other amorphous materials



From D.R. Queen et al., JNCS 426 (2015), 19-24

Connection between energy landscape, entropy, and TLS

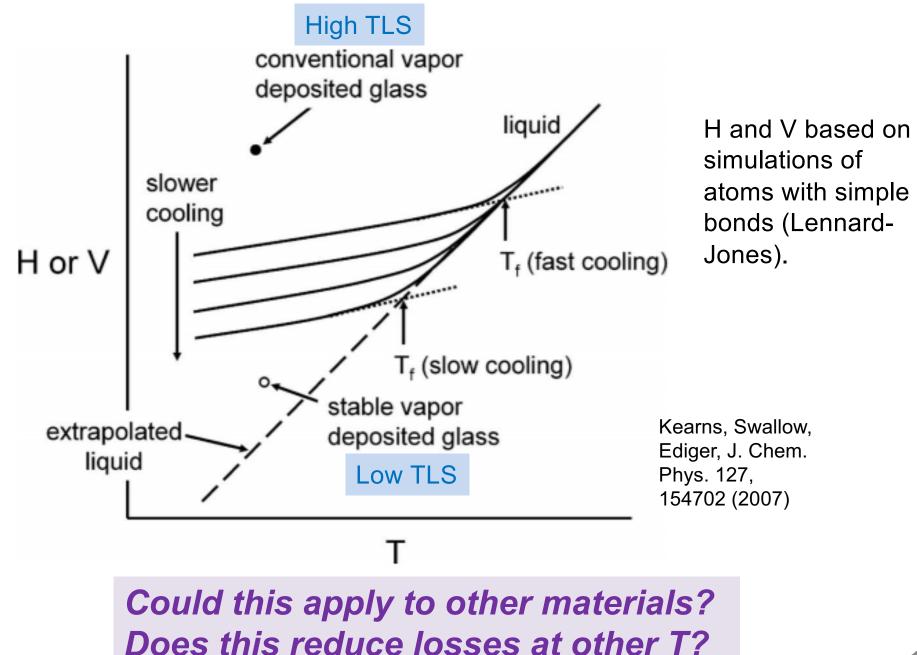


The energy landscape (right) as related to the glass transition of a liquid (left). Glasses falling out of the equilibrium supercooled liquid at a given dashed line correspond to configurations in the energy landscape.

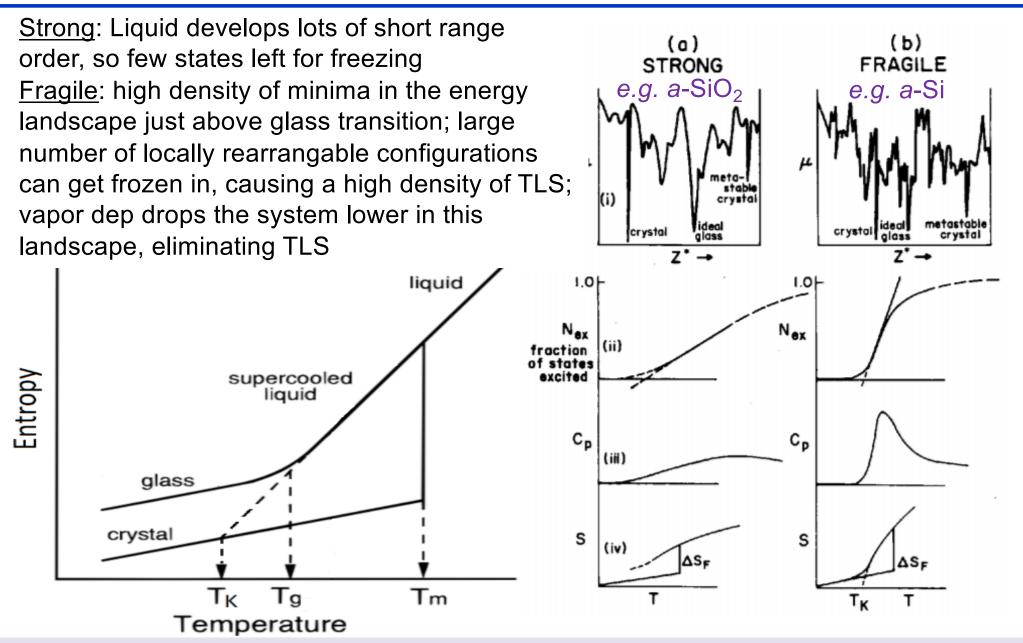
Hypothesis: vapor deposition offers a way to directly access low lying (ideal) glass state Due to high atomic mobility at film growth surface despite being at low T.

Hypothesis: Ideal glass has no nearby energy minima, so no TLS, unlike most other states

Enthalpy or Volume (density) as a function of T starting from liquid Vapor deposited compared to liquid quenching of amorphous material

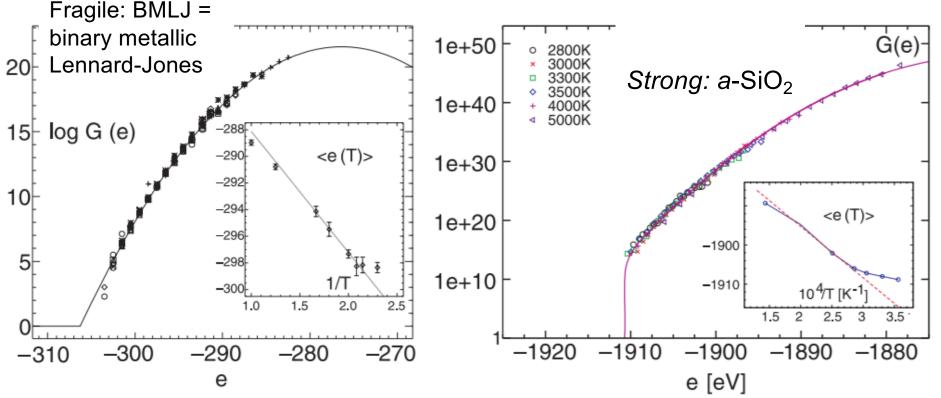


Kauzmann temperature (and paradox)



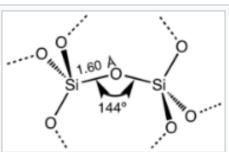
Hypothesis: Ideal glass has no nearby energy minima, so no TLS, unlike most other states Maybe only accessible for fragile glasses where T_K is at a high temperature

Modes of fragile and strong (*a*-SiO₂) glasses



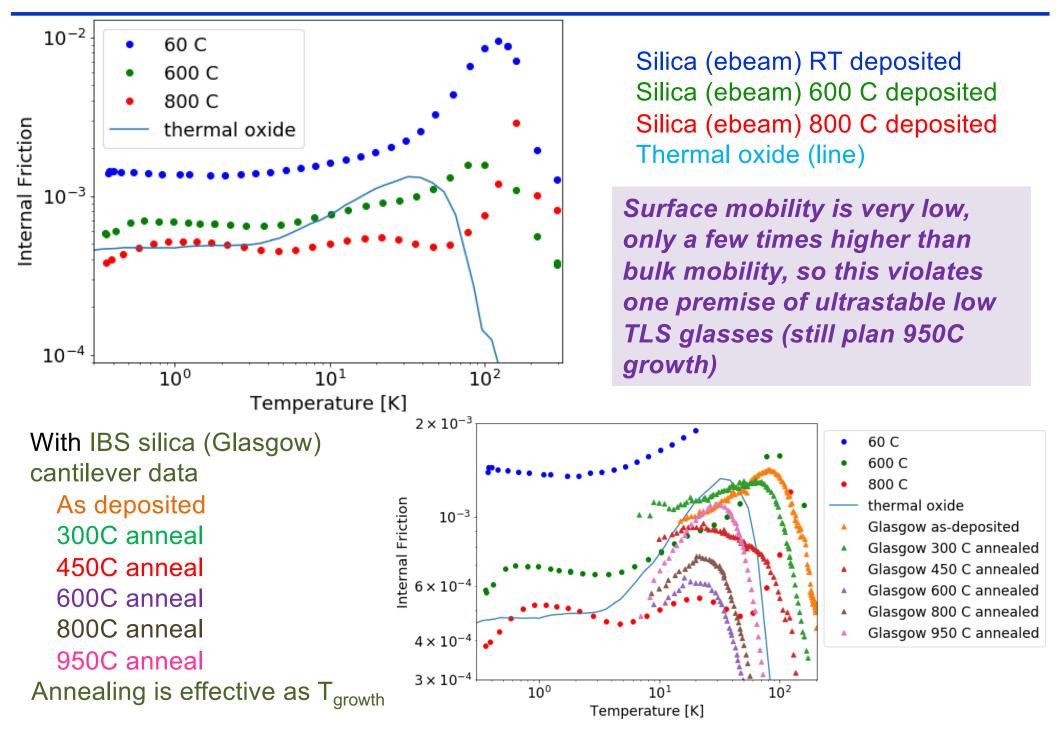
A. Heuer, J. Phys.: Condensed Matter 20, 373101 (2008)

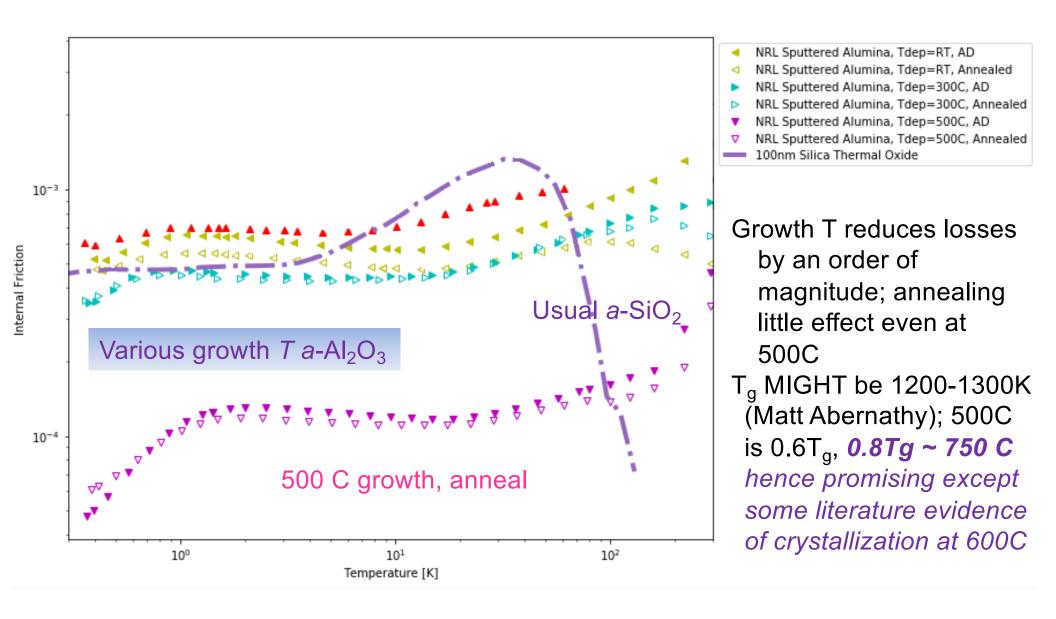
- Low-energy cutoff in G(E) for a-SiO₂ might explain drop in losses at room temperature, yet still allow a highly nearly degenerate ground state where nearly all silicon atoms are fourfold coordinated
- Large degeneracy of such states because Si-O-Si bond angle can vary around 144° without much energetic cost ----> many possible disordered networks of similar energy



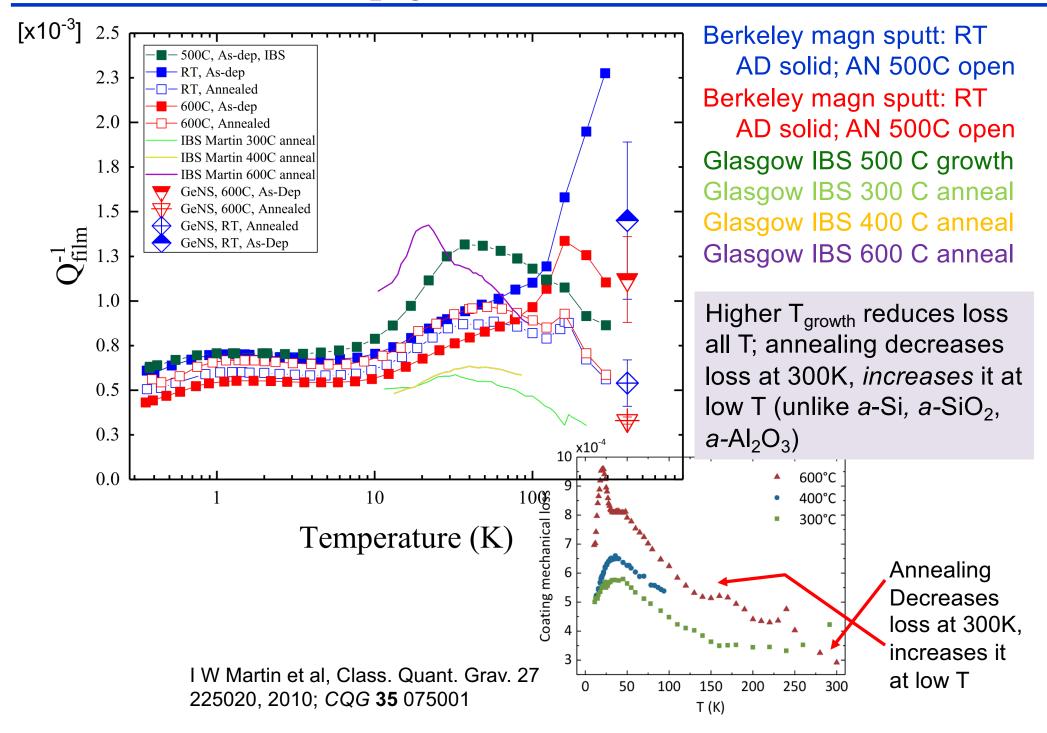
Structural motif found in α-quartz, but also found in almost all forms of silicon dioxide

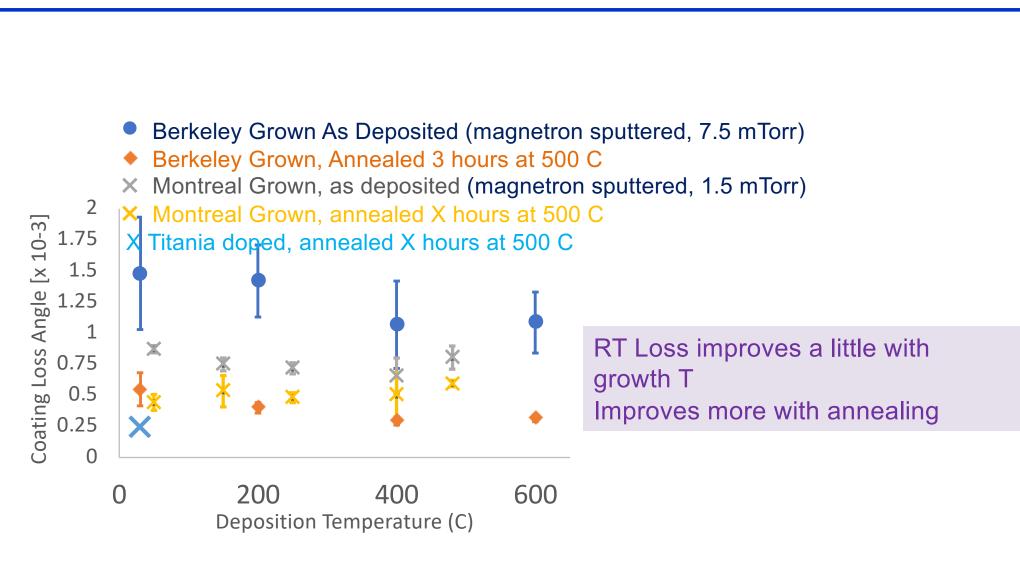
DPO Losses in silica ebeam: RT, 600 C , 800 C growth T





Losses in *a*-Ta₂O₅: DPOs, cantilevers, and GeNS





R Flaminio *et al* 2010 *Class. Quantum Grav.* **27** 084030

Amorphous Si: Absorption is high, even at 1.5 or 2 μm, due to dangling bonds (~2x10¹⁸ cm⁻³)
 Hydrogenate 500 nm thick *a*-Si, 3-12 hours at 425 ° C in forming gas (5 at.%H) to reduce absorption; measure mechanical losses

Hydrogen Forward Scattering (HFS)

- Black 50 °C: relatively homogeneous 1.45 at.%H
- red 425 °C: inhomogeneous and only 0.55 at.%H

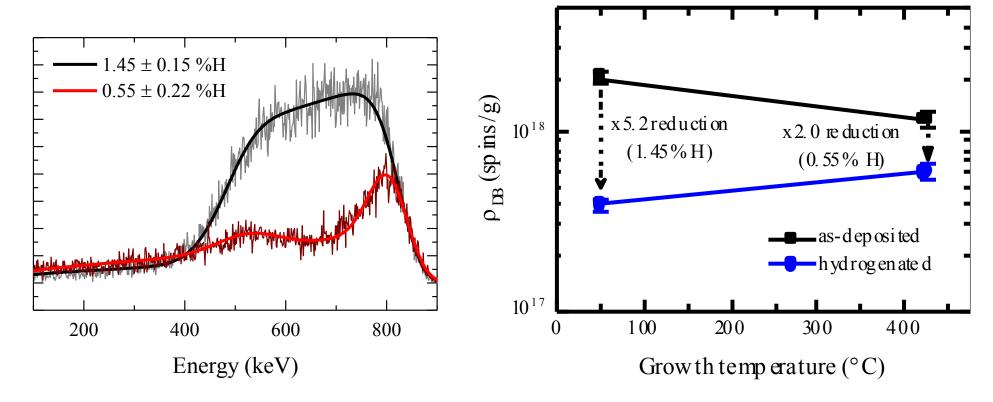
Intensity (a. u.)

Electron Paramagnetic Resonance – dangling bond density

RT growth: 5x reduction after Hydr.

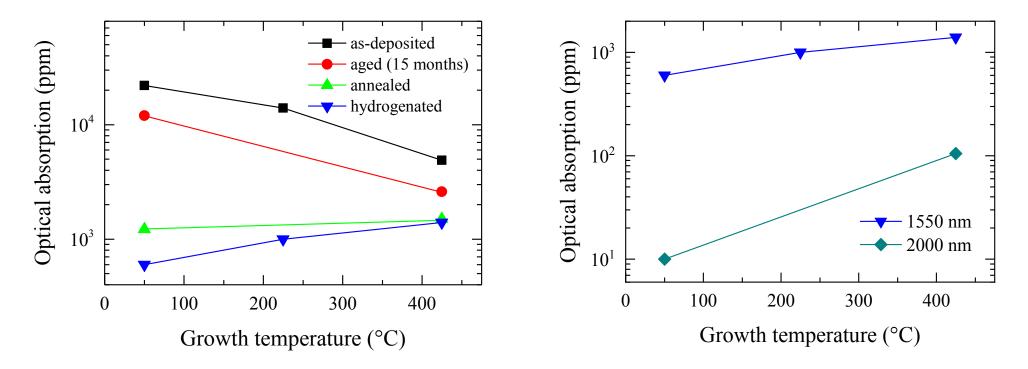
425C growth: 2x reduction after Hydr.

225 C growth pending



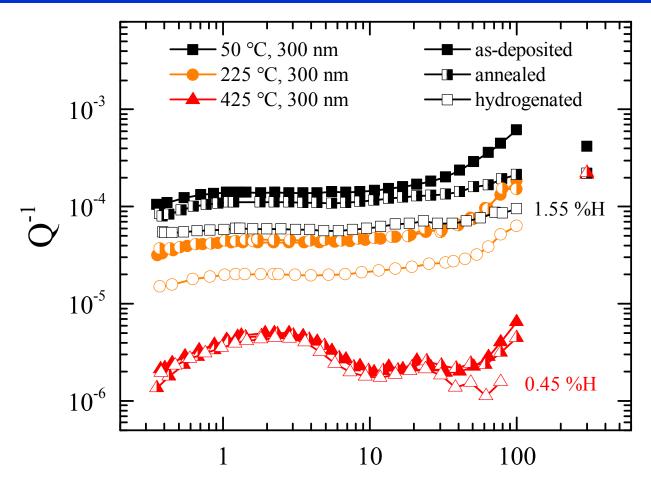
Optical absorption

- 1. as-deposited (500 nm) at 50, 225 and 425 $^\circ\,$ C,
- 2. aged (15 months),
- 3. annealed (12 hours at 425 °C) or hydrogenated (12 hours at 425 °C and 5 at.% H FG).



Optical absorption improves with growth T, and somewhat with annealing, but improves far more with hydrogenation. Best result at 2 mm RT growth, <10 ppm

Mechanical losses in *a*-Si for various growth T, annealing, hydrogenation



T (K)

As before, low T mechanical losses are improved 100x by increased growth T, improve only slightly on annealing at 425 C, and improve significantly for RT and 225 C growth on hydrogenation. RT losses less affected by any of these.
A trade-off of absorption and losses; likely sweet spot at e.g. 250C growth where H is incorporated, losses 10-5, absorption may be low

Conclusions and Open Questions

- Are low TLS in ultrastable a-Si (and IMC) the "exception that proves the rule" of universal low T glass properties? Or, is there a new rule – "universal glass properties" at low T are perhaps due to the universal nature of liquid quenching and domain growth/correlation length growth/boundaries?
- Is low TLS related to growth near T_K? (If (and only if) surface mobility during growth is high). Fragile glasses have T_K near T_g, where mobility is high, so low TLS would be correlated with fragility
- Or is low TLS related to nature of bonding: overconstrained (tetrahedral Si) versus underconstrained e.g. Si-O-Si bonds in *a*-SiO₂ and TLS in *a*-Si due to nanovoids
- Silica, alumina show increased density and reduced loss at low T with increased T_{growth}; not as much as a-Si, but not yet at T_{growth}= 0.8T_g.
- Tantala shows reduced losses at low T with increased growth T; not as much as a-Si, and likely at T_{growth} = 0.8T_g; annealing big effects, T_{growth} not stabilizing structure.
- Low losses at room temperature in all are not well correlated with low losses at low T
- Route to low room T losses is to find a material like *a*-SiO₂ with strong well formed bonds in liquid state (i.e. strong glass) and moderately high T_g
- Route to low low T losses is fragile glass with moderate T_g and suppress crystallization