#### Recent results from $\varUpsilon$ decays at BABAR

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Searches for

- gluonium in radiative  $\Upsilon(1S)$  decays
- stable six-quark states in  $\Upsilon(2S)$  and  $\Upsilon(3S)$  decays

### The BABAR experiment



#### Data samples

As of 2008/04/11 00:00





# Search for gluonium in quarkonia radiative decays

 $J^{PC}=0^{++}$  glueball expected in the region  $M\approx 1.5\div 2~{\rm GeV}$ 

Searched extensively in  $J/\psi$  radiative decays



in radiative  $\Upsilon(1S)$  decays  $\approx 25$  reduction (quark mass and width differences) see e.g. Ochs, J. Phys. G 40, 043001 (2013)



First study of  $\Upsilon(1S)$  radiative decays from CLEO PRD 73, (2006) 032001 large backgrounds from  $e^+e^- \rightarrow \gamma$  Vector





 $\Upsilon(1S) 
ightarrow \gamma h^+ h^-$  from  $\Upsilon(2S), \Upsilon(3S)$ 

Exploit  $\Upsilon(2S)$  and  $\Upsilon(3S)$  tagging: fully reconstruct

 $\Upsilon(3S), \Upsilon(2S) \to \pi_s^+ \pi_s^- \Upsilon(1S) \to \pi_s^+ \pi_s^- (\gamma h^+ h^-)$ 



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### $\Upsilon(1S) \rightarrow \gamma h^+ h^-$ : hadronic spectrum A rich $\pi^+\pi^-$ and $K^+K^-$ spectrum



"simple" fit with S-wave plus interfering BW (+  $\rho^0(770)$  background)

some of the resonances clearly visible!

S-wave =  $|BW_{f_0(500)}(m) + cBW_{f_0(980)}(m)e^{i\phi}|^2$ 

, The fraction of S-wave associated with the  $f_0(500)$  is  $(27.7\pm3.1)\%$ 

#### PRD 97 (2018) 112006

# $\Upsilon(1S) ightarrow \gamma h^+ h^-$ : Legendre moments



Assuming only S and D wave 
$$\begin{split} &\sqrt{4\pi}\langle Y^0_0\rangle = S^2 + D^2, \\ &\sqrt{4\pi}\langle Y^0_2\rangle = 2SD\cos\phi_{SD} + 0.639D^2, \\ &\sqrt{4\pi}\langle Y^0_4\rangle = 0.857D^2, \end{split}$$

extract amplitudes as a function of mass (model independent!)





### $\Upsilon(nS) \rightarrow \pi_s^+ \pi_s^- \Upsilon(1S) \rightarrow \pi_s^+ \pi_s^- \gamma h^+ h^-$ angular analysis PRD 97 (2018) 112006

Full angular analysis of the decay chain in the helicity formalism (see backup) in three different mass ranges



#### Fit results for the yields and helicities

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Resonance	Mass range (GeV/ $c^2$ )	Events	Spin	$ A_{00} ^2/ A_{01} ^2$		
$\pi\pi S$ -wave	0.6–1.0	104	0	$0.09\pm0.33$		
$\begin{array}{l} f_2(1270) \to \pi^+\pi^- \\ f_2'(1525) \to K^+K^- \\ f_0(1500) \to K^+K^- \end{array}$	1.092–1.460 1.424–1.620	280 36 40	2 2 0	$\begin{aligned} & A_{01} ^2 /  A_{00} ^2 \\ &1.07 \pm 0.31 \\ &47.9 \pm 10.8 \\ &0.04 \pm 0.07 \end{aligned}$	$\begin{array}{c}  C_{11} ^2/ C_{10} ^2 \\ 0.00 \pm 0.03 \\ 0.42 \pm 0.36 \end{array}$	$\begin{array}{c}  C_{12} ^2 /  C_{10} ^2 \\ 0.29 \pm 0.08 \\ 1.43 \pm 0.35 \end{array}$
Reweight yields by efficiency to obtain $\mathcal{B}$ 's: $w_R = \frac{\sum_{i=1}^{N_R} 1/\epsilon_i (\cos \theta_H, \cos \theta_R)}{N_R}$						

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# $\mathcal{B}(\Upsilon(1S) \to \gamma R)$

Normalize to known  $\Upsilon(3S)$  and  $\Upsilon(2S)$   $\mathcal{B}$ 's

$$\mathcal{B}(R) = \frac{N_R(\Upsilon(nS) \to \pi_s^+ \pi_s^- \Upsilon(1S)(\to R\gamma))}{N(\Upsilon(nS) \to \pi_s^+ \pi_s^- \Upsilon(1S)(\to \mu^+ \mu^-))} \times \mathcal{B}(\Upsilon(1S) \to \mu^+ \mu^-)$$

based on observed yields for the normalization modes

$$N(\Upsilon(2S) \to \pi_s^+ \pi_s^- \Upsilon(1S)(\to \mu^+ \mu^-)) = (4.35 \pm 0.12_{\rm sys}) \times 10^5$$
  
$$N(\Upsilon(3S) \to \pi_s^+ \pi_s^- \Upsilon(1S)(\to \mu^+ \mu^-)) = (1.32 \pm 0.04_{\rm sys}) \times 10^5$$

Good agreement BABAR/CLEO for  $f_J(1500)$  and  $f_2(1270)$ 

Resonance	$\mathcal{B}(10^{-5})~(B\!\!AB\!\!AR)$	CLEO
$\pi\pi$ S-wave	$4.63 \pm 0.56 \pm 0.48$	$(f_0(980)) \ 1.8^{+0.8}_{-0.7} \pm 0.1$
$f_2(1270)$	$10.15 \pm 0.59 \ {}^{+0.54}_{-0.43}$	$10.2 \pm 0.8 \pm 0.7$
$f_0(1710) \rightarrow \pi\pi$	$0.79 \pm 0.26 \pm 0.17$	
$f_J(1500) \rightarrow K\bar{K}$	$3.97 \pm 0.52 \pm 0.55$	$3.7^{+0.9}_{-0.7} \pm 0.8$
$f_2'(1525)$	$2.13 \pm 0.28 \pm 0.72$	
$f_0(1500) \rightarrow K\bar{K}$	$2.08 \pm 0.27 \pm 0.65$	
$f_0(1710) \rightarrow K\bar{K}$	$2.02 \pm 0.51 \pm 0.35$	$0.76 \pm 0.32 \pm 0.08$

First observation (5.7  $\sigma$ ) of  $\Upsilon(1S)$  radiative decays to  $f_0(1710)$ :

$$\frac{\mathcal{B}(f_0(1710) \to \pi\pi)}{\mathcal{B}(f_0(1710) \to K\bar{K})} = 0.64 \pm 0.27_{\rm stat} \pm 0.18_{\rm sys}$$



# Radiative decays $\Upsilon(1S) ightarrow \gamma h^+ h^-$

In the  $\Upsilon(1S)$  radiative decays to  $\pi^+\pi^-$  and  $K^+K^-$  final state we observe

- broad S-wave component in  $J/\psi$  decays impossible to study due to irreducible  $\pi^+\pi^-\pi^0$  background
- $f_0(980)$ ,  $f_2(1270)$ ,  $f_0(1500)$ ,  $f_2'(1525)$  and  $f_0(1710)$  resonances

the scalars include many of the "gluonium" candidates considered by theorists and the  $\mathcal B$ 's we measure will be useful to shed some more light:

- $\mathcal{B}(\Upsilon(1S) 
  ightarrow \gamma f_0(1710))$  predicted  $\mathcal{O}(10^{-4})$  R. Zhu, JHEP 1509, 166 (2015)
  - we have observed it for the first time at a rate that is consistent either with expectations for gluonium or with  $s\bar{s}$  dominance
- $\mathcal{B}(\varUpsilon(1S) o \gamma f_0(1500)) pprox 2 \div 4 imes 10^{-5}$  He et al., PRD 66, 074015 (2002)
  - consistent with our measurement
- $\mathcal{B}(\varUpsilon(15) o \gamma f_0(1370)) pprox 3 imes 10^{-5}$  Zhu, JHEP 1509, 166 (2015)
  - evidence for  $f_0(1370)$ ) is controversial, but prediction in the range for our measurement of  $\pi^+\pi^-$ -S-wave



### Stable six-quark state?

- LQCD calculations suggest tightly bound *uuddss* state could be stable or nearly stable Beane et al, PRD 87, 034506 (2013) Q = 0; B = 2; S = -2
  - $\mathcal{S}:$  scalar, flavour singlet, very compact: r=0.1÷0.4 fm

$$m \approx 1.2 \div 1.86 \text{ GeV}$$

- if  $m < M_A + M_P + m_e$  only doubly weak decays allowed
  - lifetime on the cosmological scale

- if  $m < 2 M_p$  is stable
- Experimentally not excluded by earlier searches
- NOT the same as the loosely bound H-dibaryon proposed by Jaffe PRL 38 (1977) 195, PRL 38 (1977) 617 with  $M{\approx}2150~{\rm MeV}$  and typical weak lifetime excluded by many searches
  - even suggested as possible dark matter candidate
    - (G. Farrar, arXiv:1708.08951,arXiv:1805.03723) but not everybody agrees

(eg. Kolb, Turner, PRD 99, 063519 (2019) ,

Gross et al. PRD 98, 063005 (2018) , McDermott et al PRD 99, 035013 (2019) )





# Search for $\Upsilon \to \bar{\Lambda}\bar{\Lambda}\mathcal{S}$ in BABAR

If S is stable it might interact hadronically in the detector, but signature not easy – better to just reconstruct the two  $\Lambda$ 's

 $\Longrightarrow$  Search for peak in missing mass distribution recoiling against the two  $\Lambda{}'{\rm s}$ 

Blind analysis based on 90×10<sup>6</sup>  $\Upsilon(2S)$  and 110×10<sup>6</sup>  $\Upsilon(3S)$  reconstructing  $\Lambda \rightarrow p\pi^-$  ( $\mathcal{B} \approx 0.64$ )

- $\implies$  Select events with two  $\varLambda$  's of same strangeness and nothing else
  - p and  $\pi$  from both  $\Lambda$ 's must be positively identified
  - A's must have a significant flight length and point back to the interaction vertex
  - at most one additional track candidate and little extra energy outside a cone around S direction ( $E_{extra} < 0.5 \text{ GeV}$ )





# Background in $\Upsilon \to \bar{\Lambda}\bar{\Lambda}S$

Dominant source from  $e^+e^- \to \Lambda\Lambda\bar{\Lambda}\bar{\Lambda}(X)$  where the two other  $\Lambda$ 's of same strangeness decay to  $n\pi^0$  ... Cross section not known

B decays do not contribute to background:  $\Upsilon(4S)$  data can be used to model background in continuum

Loose  $\chi^2$  cut (<25) on kinematic fit constraining the two  $\Lambda$ 's to the same vertex and to their known mass

scale  $\Upsilon(2S)$  and  $\Upsilon(3S)$  MC background to match what observed in "sideband" region with  $E_{extra} > 0.5~{\rm GeV}$ 







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# Upper limit on $\mathcal{B}(\Upsilon \to \bar{\Lambda}\bar{\Lambda}\mathcal{S})$

Only 4 events with  $E_{extra} < 0.5 {
m ~GeV}$ 

– None in the  ${\mathcal S}$  mass region

Upper limit from profile likelihood including systematic ( $\approx 10 - 15\%$ ) as a function of mass separately for  $\Upsilon(2S)$  and  $\Upsilon(3S)$  and combined assuming same partial width





$${\cal B}(arphi 
ightarrow ar{\Lambda} ar{\Lambda} {\cal S}) < (1.2 \div 1.4) \cdot 10^{-7}$$



# Conclusions

- Radiative decays  $\Upsilon(1S) \to \gamma h^+ h^-$  measured in tagged  $\Upsilon(1S)$  from  $\Upsilon(3S)$  and  $\Upsilon(2S)$  dipion transitions
  - Clean samples, not affected by  $e^+e^- \to \gamma \ V\!ector$  background
  - historically a good place to search for gluonium
  - numerous resonances observed in  $\pi^+\pi^-$  and  $K^+K^-$  final states
  - $f_0(1710)$  observed for the first time

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- Search for stable exa-quark  $\mathcal{S}$ , *uuddss*, in  $\Upsilon(2S)$  and  $\Upsilon(3S)$  decays
  - $\bullet$  LQCD suggest  ${\mathcal S}$  can be a tightly bound state
  - extremely long lifetime if  $m < M_{\Lambda} + M_{\rho} + M_{e}$ , stable if  $m < 2 N_{\rho}$
  - not excluded by earlier searches whether it's dark matter candidate or not, search for it!
  - Upper limit on  $\mathcal{B}(\Upsilon \to \bar{\Lambda}\bar{\Lambda}S) < (1.2 \div 1.4) \cdot 10^{-7}$ stringent limit, but other modes are possible ...





# Backup









$$\Upsilon(nS) \rightarrow \pi_s^+ \pi_s^- \Upsilon(1S) \rightarrow \pi_s^+ \pi_s^- \gamma h^+ h^-$$
 angular analysis

The slow pions in

$$\Upsilon(nS) o \pi_s^+ \pi_s^- \, \Upsilon(1S)$$



compatible with S-wave (as expected)



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In each reasonance region maximum likelihoood fit to signal+ background (from tail of nearby resonances)

$$\begin{split} \mathcal{L} = & \prod_{n=1}^{N} \bigg[ f_{\text{sig}} \frac{\epsilon(\cos\theta_{H}, \cos\theta_{\gamma}) W_{s}(\theta_{H}, \theta_{\gamma})}{\int W_{s}(\theta_{H}, \theta_{\gamma}) \epsilon(\cos\theta_{H}, \cos\theta_{\gamma}) d\cos\theta_{H} d\cos\theta_{\gamma}} \\ & + (1 - f_{\text{sig}}) \frac{\epsilon(\cos\theta_{H}, \cos\theta_{\gamma}) W_{b}(\theta_{H}, \theta_{\gamma})}{\int W_{b}(\theta_{H}, \theta_{\gamma}) \epsilon(\cos\theta_{H}, \cos\theta_{\gamma}) d\cos\theta_{H} d\cos\theta_{\gamma}} \bigg], \end{split}$$

with angular distribution for J=0 or J=2 from arXiv:1804.04044

$$\begin{split} W_0(\theta_\gamma) = & \frac{dU(\theta_\gamma)}{d\cos\theta_\gamma} = \frac{3}{8} |C_{10}|^2 |E_{00}|^2 (|A_{00}|^2 + 3|A_{01}|^2 \\ & - (|A_{00}|^2 - |A_{01}|^2)\cos 2\theta_\gamma). \end{split}$$

$$\begin{split} V_{2}(\theta_{T},\theta_{H}) &= \frac{dU(\theta_{T},\theta_{H})}{d\cos\theta_{T}} = \frac{1}{1024} |E_{00}|^{2} [6|A_{01}|^{2} (22|C_{10}|^{2} + 8|C_{11}|^{2} + 9|C_{12}|^{2}) \\ &\quad + 2|A_{00}|^{2} (22|C_{10}|^{2} + 24|C_{11}|^{2} + 9|C_{12}|^{2}) + 24(|A_{00}|^{2} + 3|A_{01}|^{2}) (2|C_{10}|^{2} - |C_{12}|^{2}) \cos 2\theta_{H} \\ &\quad + 6(|A_{00}|^{2} (6|C_{10}|^{2} - 8|C_{11}|^{2} + |C_{12}|^{2}) + |A_{0}|^{2} (18|C_{10}|^{2} - 8|C_{11}|^{2} + 3|C_{12}|^{2})) \cos 4\theta_{H} \\ &\quad - 2(|A_{00}|^{2} - |A_{01}|^{2}) \cos 2\theta_{T} (22|C_{10}|^{2} - 24|C_{11}|^{2} + 9|C_{12}|^{2} + 12(2|C_{10}|^{2} - |C_{12}|^{2}) \cos 2\theta_{H} \\ &\quad + 3(6|C_{10}|^{2} + 8|C_{11}|^{2} + |C_{12}|^{2}) \cos 4\theta_{H})]. \end{split}$$

Can be expressed in terms of longl/trans polarization of  $\Upsilon(1S)$  and  $\gamma$  (1 or 3 free parameters for S/D)



### Tightly bound six-quark state?

LQCD calculations for binding energies for various six-quark systems in the limit of SU(3)-flavour,  $m_\pi=m_K\approx 800~{\rm MeV}$ 



Systematic uncertainties on  $\Upsilon \to \bar{\Lambda}\bar{\Lambda}S$ 

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Main sources from

- signal modelling: production amplitudes angular distribution; interaction with detector
- Data/MC differences in reconstruction

S angular distribution	5 – 8%
S particle type	8–11%
$\Lambda$ reconstruction	4% per $\Lambda$
MC statistics	2%
$\mathcal{B}(\Lambda \to \rho \pi^-)$	1.6%
proton PID	1% per proton
Number of $Y$	0.6%

