

# $\Upsilon$ and $\psi$ decays as probes of solutions to the $R(D^{(*)})$ puzzle

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QWG 2019 - Torino

In collaboration with  
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JHEP 1706 (2017) 019, Arxiv: 1702.07356

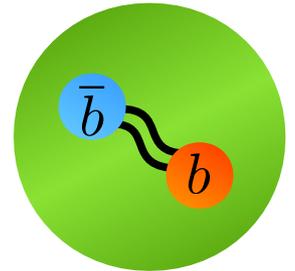


# Motivation

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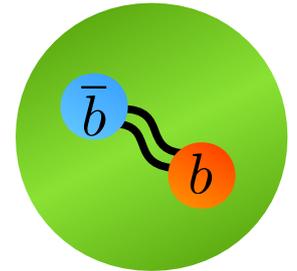
QWG 2019 - The 13th International Workshop on Heavy Quarkonium



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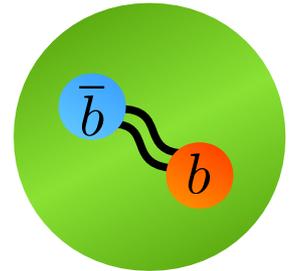


**Torino** derives its name from a Celtic word **tau**, meaning mountain.....

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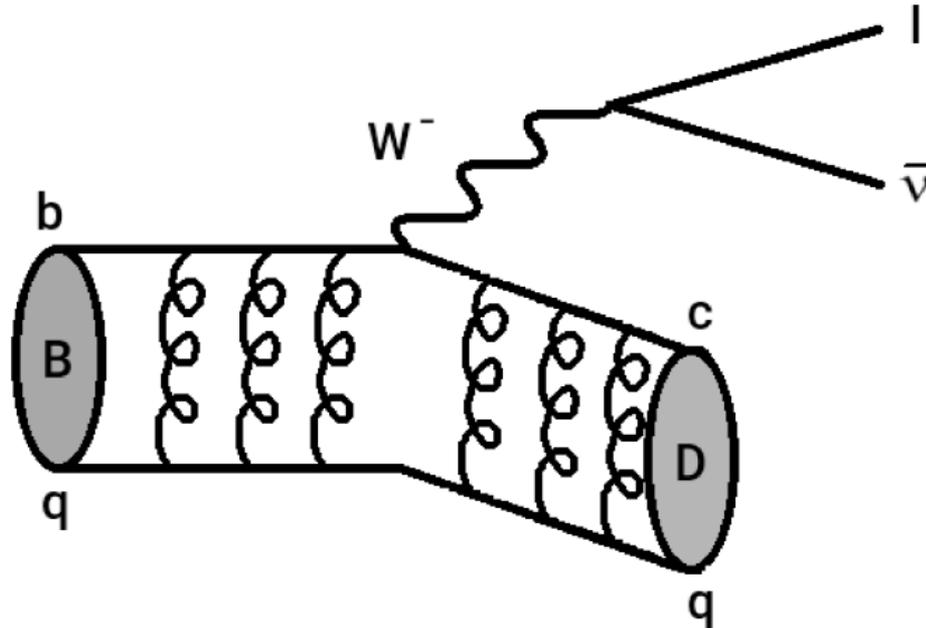


**Torino** derives its name from a Celtic word **tau**, meaning mountain.....

Today:  $\Upsilon$  to tau tau!

# Motivation

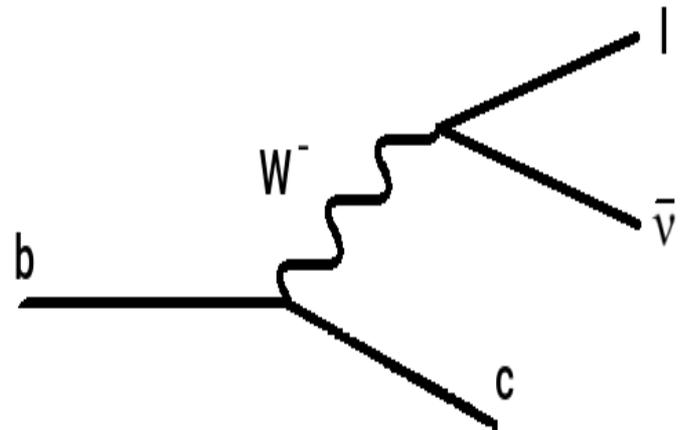
- $R(D^{(*)}) \equiv \frac{BR(B \rightarrow D^{(*)} \tau \bar{\nu})}{BR(B \rightarrow D^{(*)} \ell \bar{\nu})}, \quad \ell = \mu, e$



# Motivation

- $R(D^{(*)}) \equiv \frac{BR(B \rightarrow D^{(*)} \tau \bar{\nu})}{BR(B \rightarrow D^{(*)} \ell \bar{\nu})}$  ,  $\ell = \mu, e$

- At the quark level:  $b \rightarrow c \tau (\ell) \bar{\nu}$



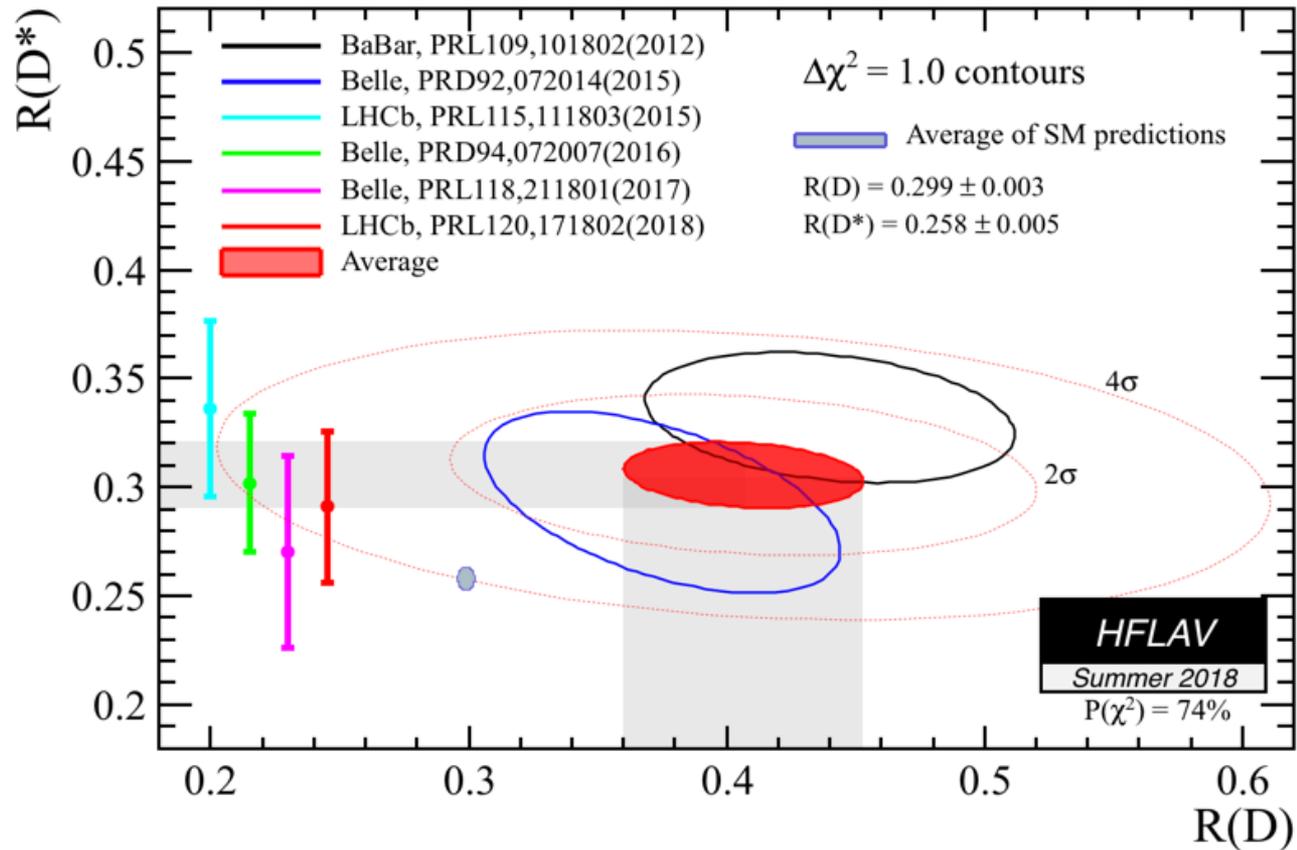
- SM:  $b \rightarrow c \tau (\ell) \bar{\nu}$  transition is mediated by the  $W$  boson

# The SM prediction

- $R^{SM}(D) = 0.299 \pm 0.003$  ,  $R^{SM}(D^*) = 0.258 \pm 0.005$  \*
- How do we know that so well?
  - Semileptonic
  - Unknown parameters cancel in the ratio
  - In the heavy quark limit  $m_b, m_c \rightarrow \infty$ , we have only phase space suppression
  - In the degenerate lepton masses limit  $m_\tau \rightarrow m_\ell$ ,  $R(D)=R(D^*)=1$
- We know to expand systematically around this small parameters

\* HFLAV Average

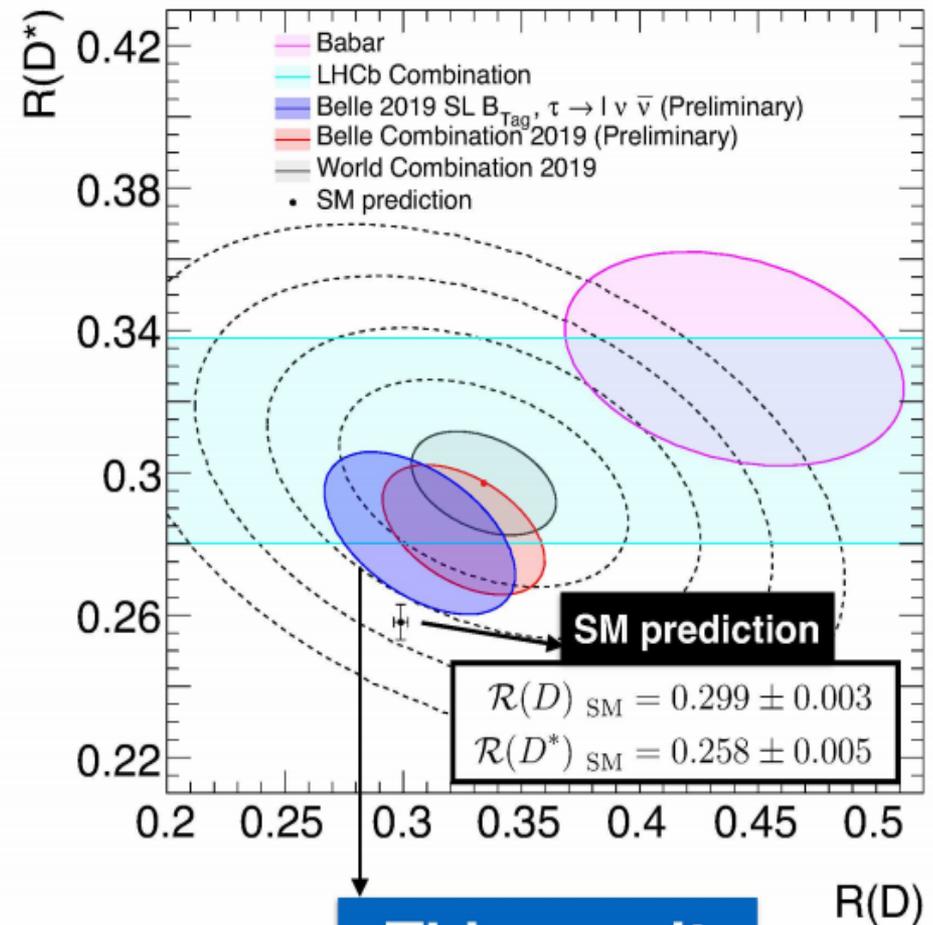
# Measurement



- $\sim 4\sigma$  deviation from SM prediction

# Conclusion / Preliminary $R(D^{(*)})$ averages

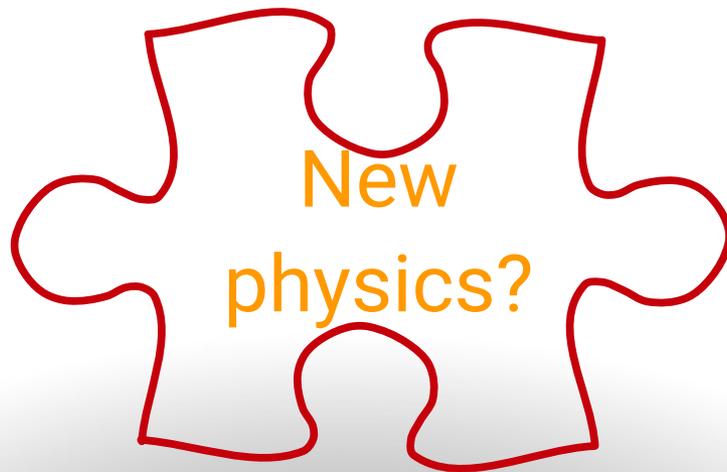
- **Most precise measurement** of  $R(D)$  and  $R(D^*)$  to date
- First  **$R(D)$**  measurement performed with a **semileptonic tag**
- Results **compatible with SM** expectation within  **$1.2\sigma$**
- **$R(D) - R(D^*)$  Belle average** is now within  **$2\sigma$**  of the SM prediction
- **$R(D) - R(D^*)$  exp. world average** tension with SM expectation **decreases from  $3.8\sigma$  to  $3.1\sigma$**



**This result**

$$\begin{aligned} \mathcal{R}(D) &= 0.307 \pm 0.037 \pm 0.016 \\ \mathcal{R}(D^*) &= 0.283 \pm 0.018 \pm 0.014 \end{aligned}$$

This is puzzling



# If new physics

- Central values are enhanced by 30% compared to SM → NP  
amplitude 15%-30% compared to SM
- New physics is non-universal and breaks lepton flavor symmetry
- New physics is probably heavy → Can work with an effective theory

# EFT – complete basis

- Demand  $G_{SM} = SU(3)_C \times SU(2)_L \times U(1)_Y$  invariance
- Assume no RH neutrinos, *i.e.*  $B \rightarrow D\tau\bar{\nu}_L$
- A complete set for  $b \rightarrow c\tau\bar{\nu}$  transitions contains only four operators
  - ›  $(\bar{e}L)(\bar{u}Q)$
  - ›  $(\bar{e}\sigma^{\mu\nu}L)(\bar{u}\sigma_{\mu\nu}Q)$
  - ›  $(\bar{L}\gamma^\mu\tau_aL)(\bar{Q}\gamma^\mu\tau_aQ)$
  - ›  $(\bar{Q}d)(\bar{e}L)$

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  - ›  $(\bar{L}\gamma^\mu\tau_a L)(\bar{Q}\gamma^\mu\tau_a Q)$
  - ›  $(\bar{Q}d)(\bar{e}L)$

# EFT – unavoidable NC

- All four operators contain also neutral currents (NC)

→ For instance

$$(\bar{e}L)(\bar{u}Q) = \overbrace{(\bar{e}_R\nu_L)(\bar{u}_Rd_L)}^{CC} - \overbrace{(\bar{e}_Re_L)(\bar{u}_Ru_L)}^{NC}$$

- We looked for observables sensitive to those NC
- Neutral currents unavoidably modify  $b\bar{b}$  and/or  $c\bar{c} \rightarrow \tau\bar{\tau}$

# EFT – NC observables

- *D. A. Faroughy, A. Greljo, J. F. Kamenik* \* – High  $P_T$  distribution of  $\tau\bar{\tau}$  signature at the LHC
- We looked on lepton non universality of  $\Upsilon$  and  $\psi$  decays

$$R_{\tau/\ell}^V \equiv \frac{\Gamma(V \rightarrow \tau^+\tau^-)}{\Gamma(V \rightarrow \ell^+\ell^-)}, \quad (V = \Upsilon, \psi(2s); \ell = e, \mu)$$

\*  $\Upsilon = b\bar{b}$  bound state

\*  $\psi = c\bar{c}$  bound state

\* *D. A. Fraoughy, A. Greljo, J. F. Kamenik, **Phys. Lett. B764 (2017) 126-134***

# Vector meson decay - SM

- Within the SM

$$R_{\tau/\ell}^V \simeq \left[ 1 + \frac{2m_\tau^2}{m_V^2} \right] \left[ 1 - \frac{4m_\tau^2}{m_V^2} \right]^{1/2} = 1 - \mathcal{O} \left( \frac{m_\tau^4}{m_V^4} \right)$$

where  $V = \Upsilon, \psi$

- Dominantly QED – 1 photon mediated

# $R_{\tau/\ell}^V$ - Prediction vs. measurement

$V(nS)$	SM prediction	Exp. value $\pm \sigma_{\text{stat}} \pm \sigma_{\text{syst}}$
$\Upsilon(1S)$	$0.9924 \pm \mathcal{O}(10^{-5})$	$1.005 \pm 0.013 \pm 0.022$
$\Upsilon(2S)$	$0.9940 \pm \mathcal{O}(10^{-5})$	$1.04 \pm 0.04 \pm 0.05$
$\Upsilon(3S)$	$0.9948 \pm \mathcal{O}(10^{-5})$	$1.05 \pm 0.08 \pm 0.05$
$\psi(2S)$	$0.390 \pm \mathcal{O}(10^{-4})$	$0.39 \pm 0.05$

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$$m_{\Upsilon(4S)} > 2m_B \quad \Rightarrow \quad \Upsilon(1S), \Upsilon(2S), \Upsilon(3S)$$

$$m_{\psi(1S)} < 2m_{\tau} \quad \Rightarrow \quad \psi(2S)$$

$$m_{\psi(3S)} > 2m_D$$

# The need for assumptions

- There are four independent CC operators
- There are eight independent NC operators
- CC + Gauge invariance  $\longrightarrow$  NC
  - Not enough measurements to fix the values of the Wilson coefficients of the four CC operators
  - No information on the other four NC operators

# EFT from simplified models

- Recall that
  - Within the SM  $b \rightarrow c\tau\bar{\nu}$  is a tree-level process
  - Central value is  $\sim 30\%$  enhanced compared to prediction
- Huge enhancement of tree-level suggests new bosons which also modify  $b \rightarrow c\tau\bar{\nu}$  at tree-level
- There are eight different possible mediators  
 $W'_\mu \sim (1, 3)_0, U_\mu \sim (3, 1)_{2/3}, X_\mu \sim (3, 3)_{2/3}, S \sim (3, 1)_{-1/3},$   
 $T \sim (3, 3)_{-1/3}, \phi \sim (1, 2)_{1,2}, D \sim (3, 2)_{7/6}, V_\mu \sim (3, 2)_{-5/6}$
- Each boson breaks LFU in  $\Upsilon$  and/or  $\psi$  decays in a different way

# Results

UV field content	$R_{\tau/\ell}^{\Upsilon(1S)}$	$R_{\tau/\ell}^{\psi(2S)}$	Predicted modification to $R_{\tau/\ell}^{\Upsilon(1S)}$
$W'_\mu \sim (1, 3)_0$	0.989-0.991	0.390	Decrease by 0.2% – 0.4%
$U_\mu \sim (3, 1)_{+2/3}$	0.952-0.990	SM	Decrease by 0.3% – 4.0%
$S \sim (3, 1)_{-1/3}$	SM	0.389-0.390	–
$V_\mu \sim (3, 2)_{-5/6}$	0.976-0.987	SM	Decrease by 0.5% – 1.6%
SM	0.992	0.390	
Current measurement	$1.005 \pm 0.025$	$0.39 \pm 0.05$	
Achievable uncertainty (with current data)	$\pm 0.01$	$\pm 0.02$	
Projected uncertainty ( $\mathcal{L}^{\Upsilon(3S)} = 1/\text{ab}$ in Belle II)	$\pm 0.004$	–	

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# Future prospects - $\Upsilon(nS)$

- Current error is  $\sim \sigma_{1S}^{BaBar} \sim 2\%$ ,  $\sigma_{2S}^{CLEO} \sim 5\%$ ,  $\sigma_{3S}^{CLEO} \sim 8\%$
- Babar and Belle have 16-20 times larger data samples than what was analyzed, giving  $\sigma \sim 1\%$
- Running at  $\Upsilon(3S)$  with  $\mathcal{L} \sim 1/ab$  Belle II might reach  $\sigma_{1S} \simeq 0.4\%$

# Summary



- Current error is  $\sigma_{1S}^{BaBar} \sim 2\%$
- Running at  $\Upsilon(3S)$  with  $\mathcal{L} \sim 1/ab$  Belle II might covers most region of parameter space related to  $R(D^{(*)})$
- LFU in  $\Upsilon$  decays provide additional motivation to study  $\Upsilon(3S)$  at Belle II
- Test the SM and Probe NP even if  $R(D^{(*)})$  disappears

Thank you!





# Backup

# Vector meson decay – SM corrections

- Within the SM

$$R_{\tau/\ell}^V \simeq (1 + 2x_\tau^2)(1 - 4x_\tau^2)^{1/2}$$

where  $x_\tau = m_\tau/m_V$ .

- Leading corrections:

- Corrections due to  $m_\mu$  are  $\delta R_{\tau/\ell}^V \lesssim \mathcal{O}(10^{-7})$
- Corrections due to Z exchange are  $\delta R_{\tau/\ell}^V \lesssim \mathcal{O}(10^{-5})$
- QED at 1-loop  $\delta R_{\tau/\ell}^V \simeq 6 \cdot 10^{-3} x_\tau^2 \sim \mathcal{O}(10^{-4})$  (inclusive)

# EFT – fixing all NC

- Motivated by FCNC and LHC constraints
  - We impose  $U(2)_Q$  symmetry of the light left-handed quarks
  - We assume flavor alignment to the down mass basis
- Choosing a simplified model
  - Fixes the 4 CC Wilson coefficients (up to overall normalization)
  - Fixes the 4 corresponding NC, as well as the 4 other NC. For instance a mediator which generates  $(\bar{u}L)(\bar{e}Q)$  will also generate  $(\bar{u}L)(\bar{L}u)$ ,  $(\bar{e}Q)(\bar{Q}e)$

# Vector meson decay – dim(6) operators

- Eight independent dim.(6) operators contribute to leptonic meson decays

$$\begin{aligned} \mathcal{L}_{lq} = & C_{VRR}^{q\ell} (\bar{e}_R \gamma^\mu e_R) (\bar{q}_R \gamma_\mu q_R) + C_{VRL}^{q\ell} (\bar{e}_R \gamma^\mu e_R) (\bar{q}_L \gamma_\mu q_L) \\ & + C_{VLR}^{q\ell} (\bar{e}_L \gamma^\mu e_L) (\bar{q}_R \gamma_\mu q_R) + C_{VLL}^{q\ell} (\bar{e}_L \gamma^\mu e_L) (\bar{q}_L \gamma_\mu q_L) \\ & + C_T^{q\ell} (\bar{e}_L \sigma^{\mu\nu} e_R) (\bar{q} \sigma_{\mu\nu} q) + \text{h.c.} \quad (+\text{scalar operators}) \end{aligned}$$

- A straightforward calculation relates the 8 Wilson coefficient  $C_A^{q\ell}$  to 4 form factors appearing in the ratio

$$\begin{aligned} R_{\tau/\ell}^V \simeq & \frac{\sqrt{1-4x_\tau^2}}{(4\pi\alpha Q_q)^2} \left[ |\tilde{A}_V^{q\tau}|^2 (1+2x_\tau^2) + |\tilde{B}_V^{q\tau}|^2 (1-4x_\tau^2) + \frac{|\tilde{C}_V^{q\tau}|^2}{2} (1-4x_\tau^2)^2 \right. \\ & \left. + \frac{|\tilde{D}_V^{q\tau}|^2}{2} (1-4x_\tau^2) + 2\text{Re} \left[ \tilde{A}_V^{q\tau} \tilde{C}_V^{*q\tau} \right] x_\tau (1-4x_\tau^2) \right] \end{aligned}$$

# Vector meson decay – Dim(6) operators

Relation between Wilson coefficient and Form-factors:

$$A_V^{q\ell} = -4\pi\alpha Q_q + \frac{m_V^2}{4} \left[ \left( C_{VLL}^{q\ell} + C_{VRR}^{q\ell} + C_{VLR}^{q\ell} + C_{VRL}^{q\ell} \right) + 16x_\ell \frac{f_V^T}{f_V} \text{Re} \left[ C_T^{q\ell} \right] \right],$$

$$B_V^{q\ell} = \frac{m_V^2}{4} \left( C_{VRR}^{q\ell} + C_{VRL}^{q\ell} - C_{VLR}^{q\ell} - C_{VLL}^{q\ell} \right),$$

$$C_V^{q\ell} = 2m_V^2 \frac{f_V^T}{f_V} \text{Re} \left[ C_T^{q\ell} \right],$$

$$D_V^{q\ell} = 2m_V^2 \frac{f_V^T}{f_V} \text{Im} \left[ C_T^{q\ell} \right].$$

where

$$\langle 0 | \bar{q} \gamma^\mu | V(p) \rangle = f_V m_V \epsilon^\mu(p), \quad \langle 0 | \bar{q} \sigma^{\mu\nu} | V(p) \rangle = f_V^T [\epsilon^\mu(p) p^\nu - \epsilon^\nu(p) p^\mu]$$

# Numerical results – example $U_\mu \sim (3, 1)_{2/3}$

- UV

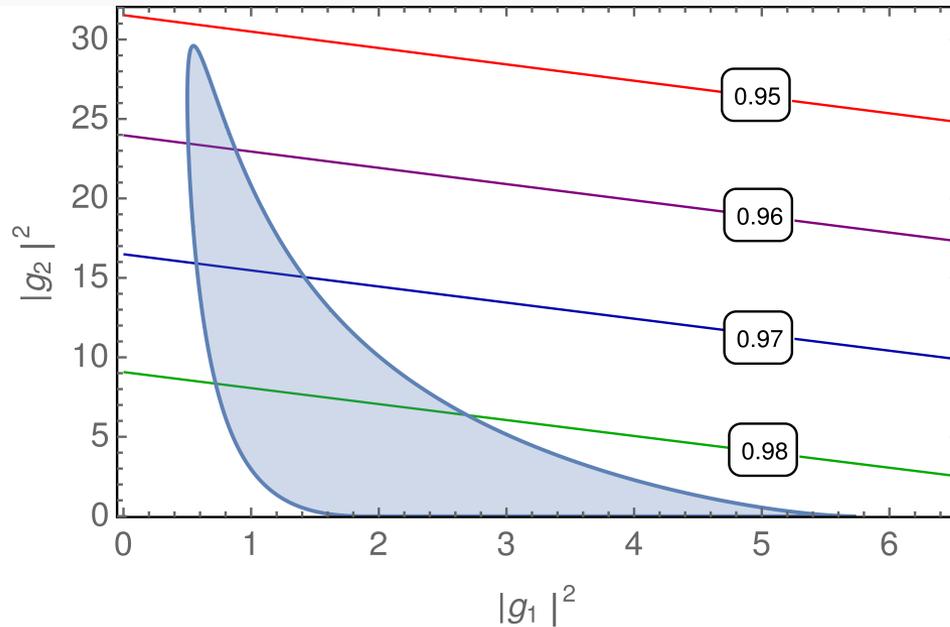
$$\mathcal{L}_U = g_1 \bar{Q}_3 \not{U} L_3 + g_2 \bar{d}_3 \not{U} e_3 + h.c.$$

- EFT

$$\begin{aligned} \mathcal{L}_U^{EFT} = & -\frac{|g_1|^2}{2M_U^2} \overbrace{(\bar{L}\gamma^\mu L)(\bar{Q}\gamma_\mu Q)}^{N.C.} - \frac{2|g_1|^2}{M_U^2} \overbrace{(\bar{L}\gamma^\mu \tau_a L)(\bar{Q}\gamma_\mu \tau_a Q)}^{N.C.+C.C.} \\ & - \frac{|g_2|^2}{M_U^2} \overbrace{(\bar{e}\gamma^\mu e)(\bar{d}\gamma_\mu d)}^{N.C.} + \left[ \frac{g_1 g_2^*}{M_U^2} \overbrace{(\bar{Q}d)(\bar{e}L)}^{N.C.+C.C.} + h.c. \right] \end{aligned}$$

# Numerical results – example $U_\mu \sim (3, 1)_{2/3}$

- Results:  $(g_1^{2 \text{ bfp}}, g_2^{2 \text{ bfp}}) = (3.3, 0.4) \left( \frac{M_U}{\text{TeV}} \right)^2$



$$R_{\tau/\ell}^{\Upsilon(1S)} = 0.952 - 0.990$$

$$R_{\tau/\ell}^{\Upsilon(2S)} = 0.949 - 0.991$$

$$R_{\tau/\ell}^{\Upsilon(3S)} = 0.946 - 0.992$$

- $R_{\tau/\ell}^{\psi(2S)}$  is not modified

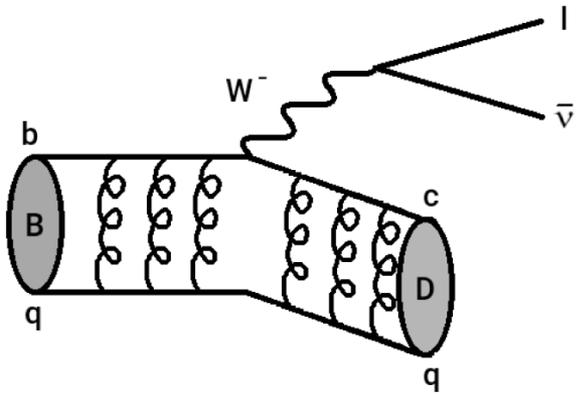
# Future prospects - $\psi(2S)$

- Current error is  $\sigma \sim 13\%$
- Bess II already has 10 times larger data sample than what was analyzed, giving  $\sigma \sim 4\%$
- We predict that  $R_{\tau/\ell}^{\psi}$  is modified by at most few per-mil
- Removing the imposed  $U(2)_Q$  symmetry might lead to a much larger modification (work in progress)
- An order of magnitude improvement is needed to be achieved in Bess III

# In a cartoon

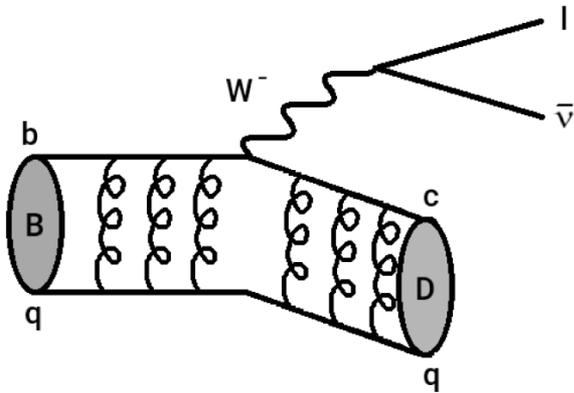
# In a cartoon

## Charge current (CC)

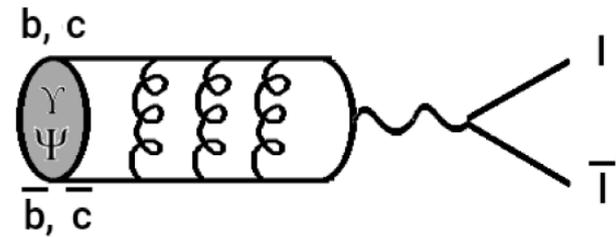


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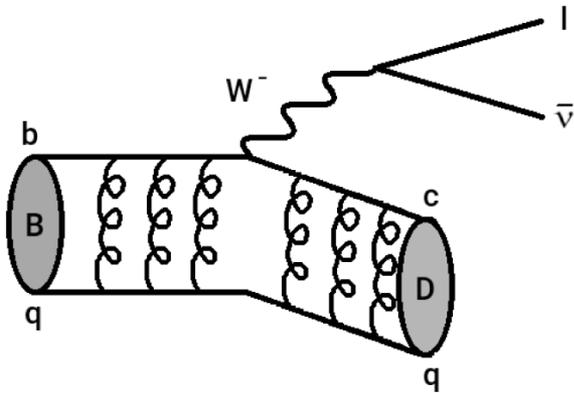


## Neutral current (NC)

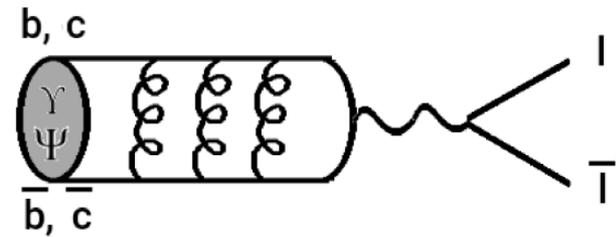


# In a cartoon

## Charge current (CC)



## Neutral current (NC)



$\nu$  is part of a doublet



if we have  $\nu$  we have  $\tau$