

Exotics near thresholds and Hadron interactions in the heavy flavor sector

Yasuhiro Yamaguchi ¹

in collaboration with

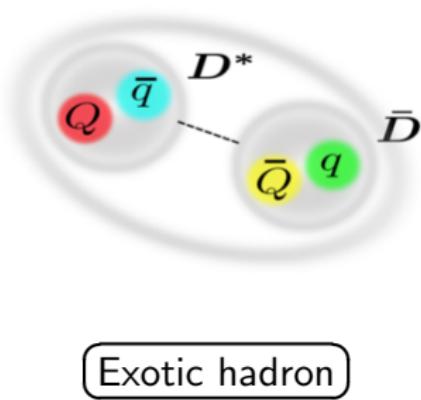
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Seminar@INFN Genova
Genova, Italy, March 2018

Outline

- Introduction
 - Exotic hadrons
 - $Z_c(3900)$
- Meson exchange model
 - Heavy quark symmetry and the meson exchange potential
 - Heavy meson exchange potential
- Quark exchange model
- Summary

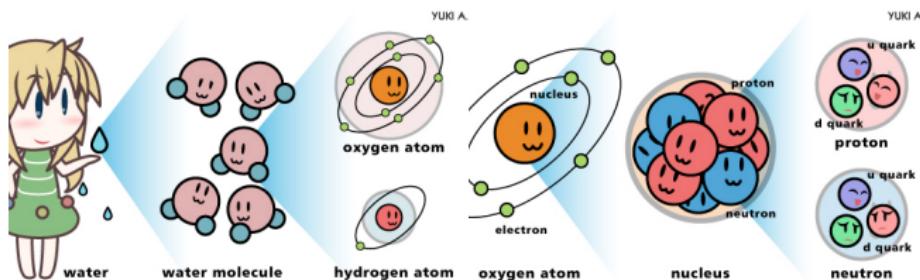


Introduction

- Level structure of “matter”

Introduction

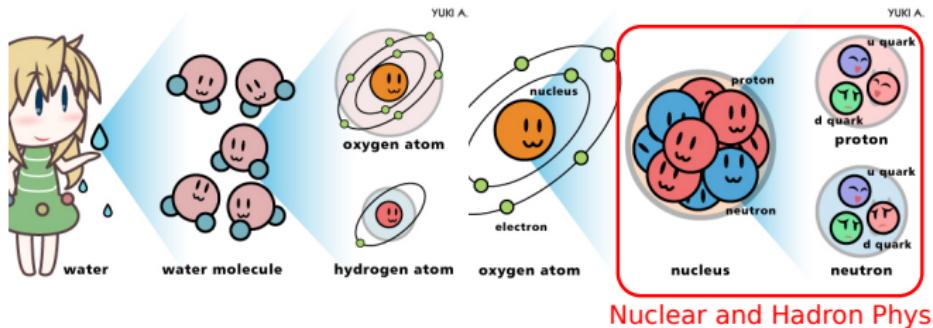
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from HiggsTan.com

Introduction

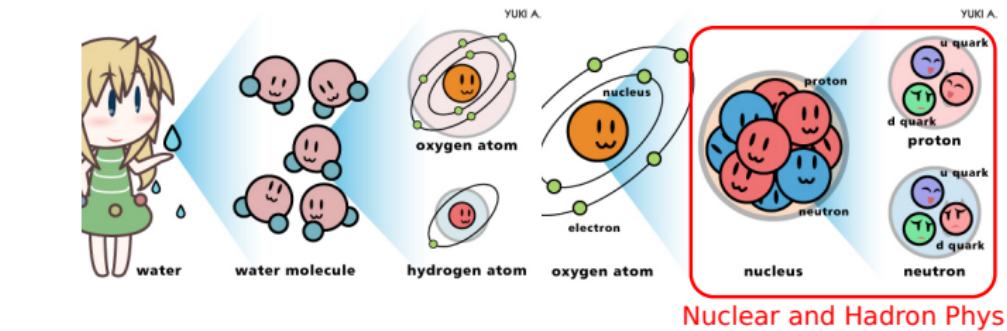
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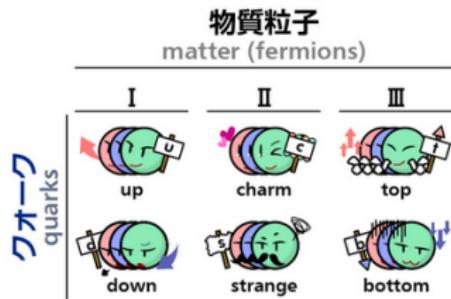
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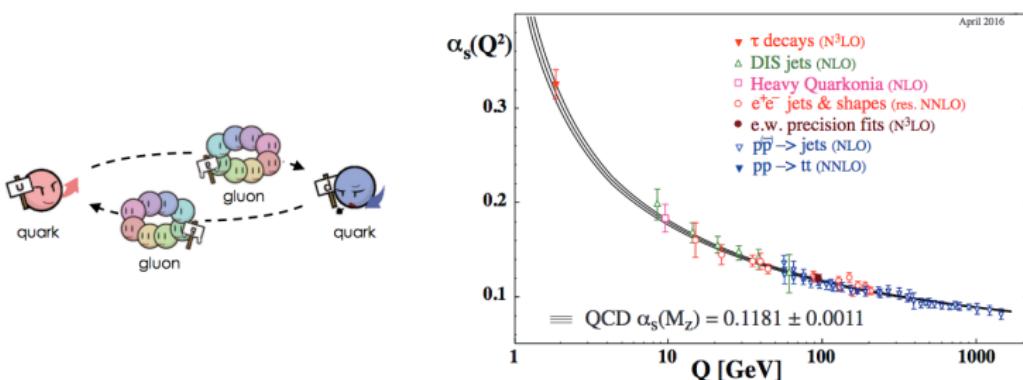
- Elementary particle “Quarks”
- $u, d \rightarrow$ Proton, Nucleon (Ordinary)
- Systems with s, c or b**
beyond ordinary matters

Quark-Gluon dynamics

Introduction

- Quantum ChromoDynamics (QCD)
→ describing the strong interaction of quark and gluon

$$\mathcal{L}_{QCD} = \bar{q}(i\gamma^\mu(\partial_\mu + igA_\mu^a T^a) - m_q)q - \frac{1}{4}G_{\mu\nu}G^{\mu\nu}$$



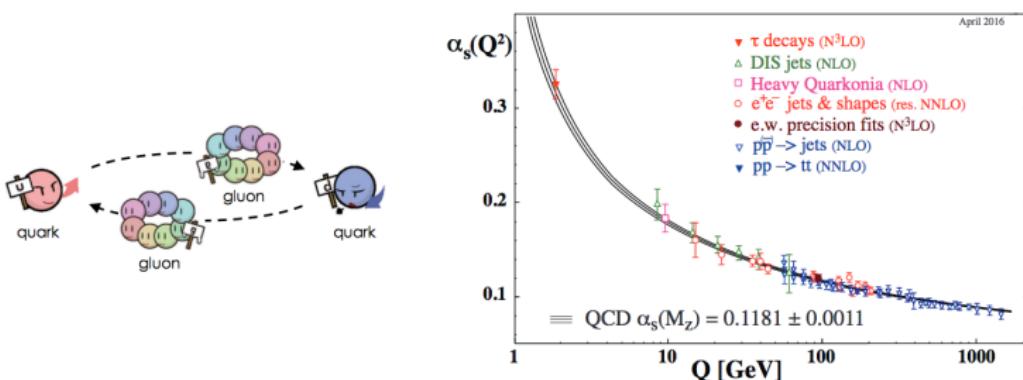
- Running coupling $\alpha_s = g^2/4\pi$
- ⇒ At low-energy, perturbative calculation breaks down
→ **Quark confinement** inside “Hadron” (e.g. proton, nucleon,...)

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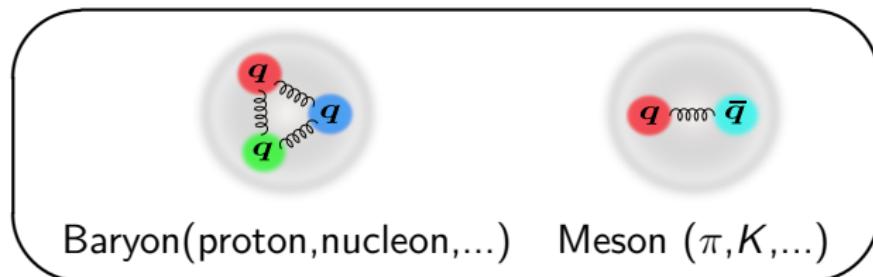
Difficulty of QCD

→ **Model, Lattice**

Constituent quark model: simple but powerful

Introduction

- Simple description of Hadron structure → Constituent quark model
- Constituent quark model (Baryon(qqq) and Meson $q\bar{q}$) has been **successfully applied to the hadron spectra!**



* q : "Constituent quark"

- ▷ Quark potential (One gluon exchange + Linear confinement)

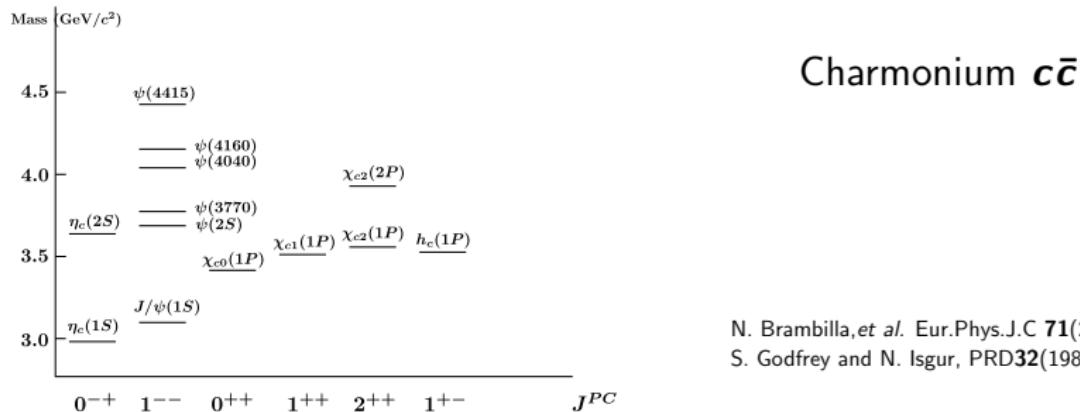
$$V_q(r) = -\frac{a}{r} + br + c + d \frac{\vec{S}_1 \cdot \vec{S}_2}{m_1 m_2} \delta^{(3)}(r) + \dots$$

(Parameters are fixed to reproduce grand state hadrons)

Constituent quark picture and beyond

Introduction

► Charmonium ($c\bar{c}$)

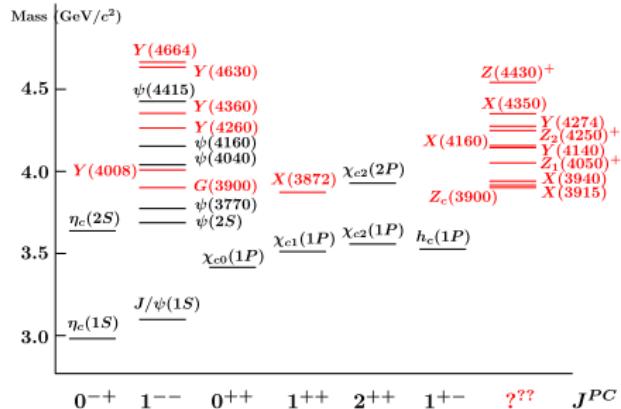


- Constituent quark model is powerful.

Constituent quark picture and beyond

Introduction

► Charmonium ($c\bar{c}$) and New Exotic hadrons X, Y, Z



Charmonium $c\bar{c}$
and
Exotic hadrons ($\neq c\bar{c}$)
 X, Y, Z

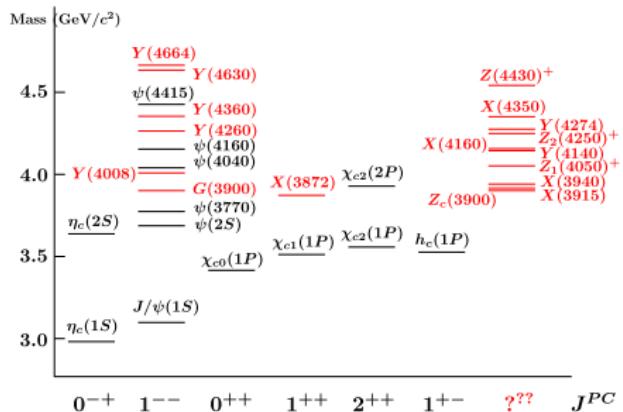
N. Brambilla, et al. Eur.Phys.J.C 71(2011)1534
S. Godfrey and N. Isgur, PRD 32(1985)189

- Constituent quark model is powerful. But, **many exotic states** have found in the heavy quark (c, b) sector!

Constituent quark picture and beyond

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Charm $X(3872)$, Charged Z_c

Belle PRL **91** (2003) 262001. A. Hosaka *et al.* PTEP **2016** (2016) no.6, 062C01

Bottom $Z_b(10610)$ and $Z_b(10650)$ Belle PRL **108** (2012) 122001

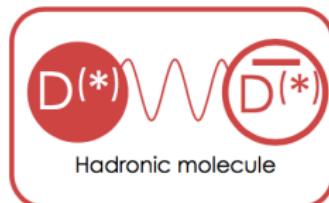
Structures of Exotic hadrons

Introduction

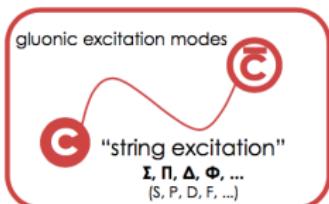
- What is the structure of Exotic hadrons? Candidates:



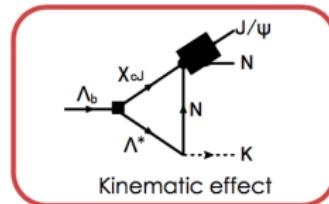
Compact multiquark



Hadronic molecule



gluonic excitation modes
"string excitation"
 $\Sigma, \Pi, \Delta, \Phi, \dots$
 (S, P, D, F, \dots)



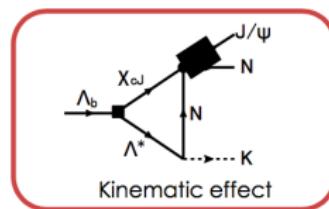
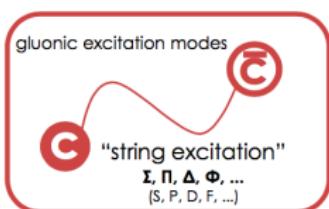
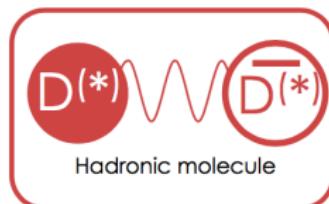
Kinematic effect

- ▷ Multiquark states?:
Compact multiquark and/or Hadronic molecule
- ▷ Gluon excitation?: Hybrid state
- ▷ Kinematical effect?: It is not bound state!

Structures of Exotic hadrons

Introduction

- What is the structure of Exotic hadrons? Candidates:



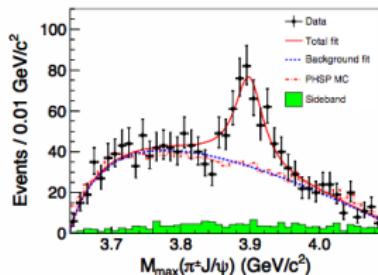
- ▷ Multiquark states?:
Compact multiquark and/or **Hadronic molecule**
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- ▷ **Kinematical effect?**: It is not bound state!

Charged Charmonium: $Z_c(3900)$

Introduction

- Charged Charmonium??
- $Y(4260) \rightarrow Z_c(3900)\pi \rightarrow J/\psi\pi\pi$

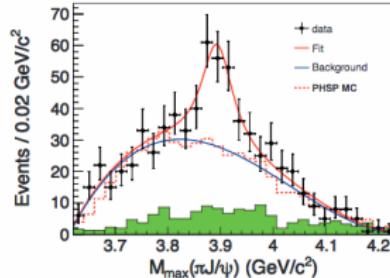
BESIII, PRL **110**(2013)252001



$$M = 3899.0 \pm 3.6_{sta} \pm 4.9_{sys} \text{ MeV}$$

$$\Gamma = 46 \pm 10_{sta} \pm 20_{sys} \text{ MeV}$$

Belle, PRL **110**(2013)252002



$$M = 3894.5 \pm 6.6_{sta} \pm 4.5_{sys} \text{ MeV}$$

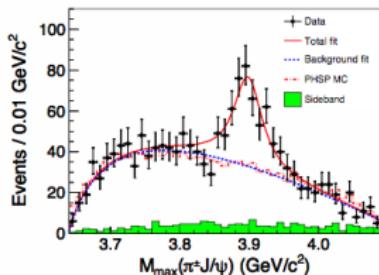
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Charged Charmonium: $Z_c(3900)$

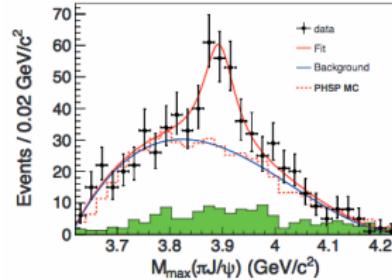
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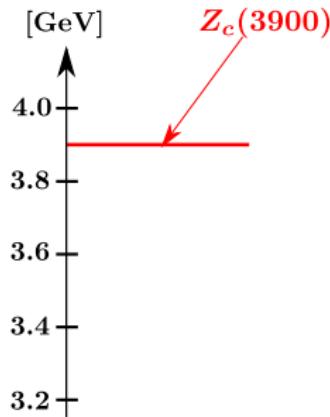
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- ▷ Ordinal Charmonium $c\bar{c}$: no electric charge.
⇒ $Z_c^+(3900)$: **Genuine Exotic State!**? $c\bar{c}u\bar{d}$

$Z_c(3900)$: Hadronic Molecules?

Introduction

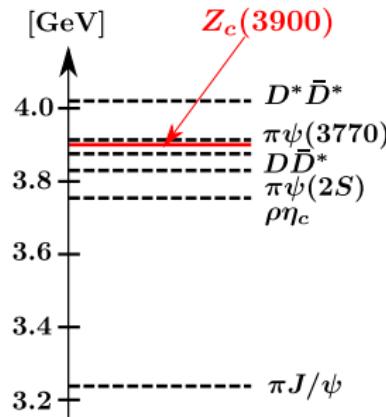
- $Z_c(3900)$ close to the hadron-hadron thresholds



$Z_c(3900)$: Hadronic Molecules?

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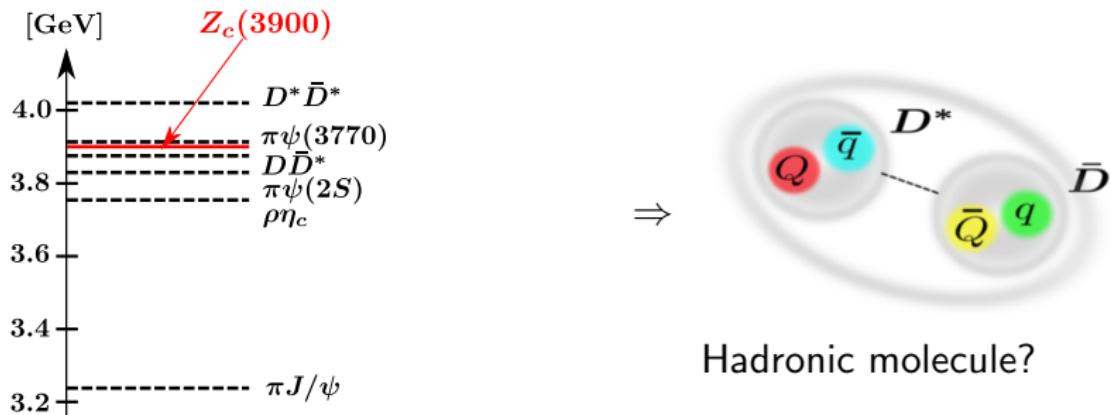
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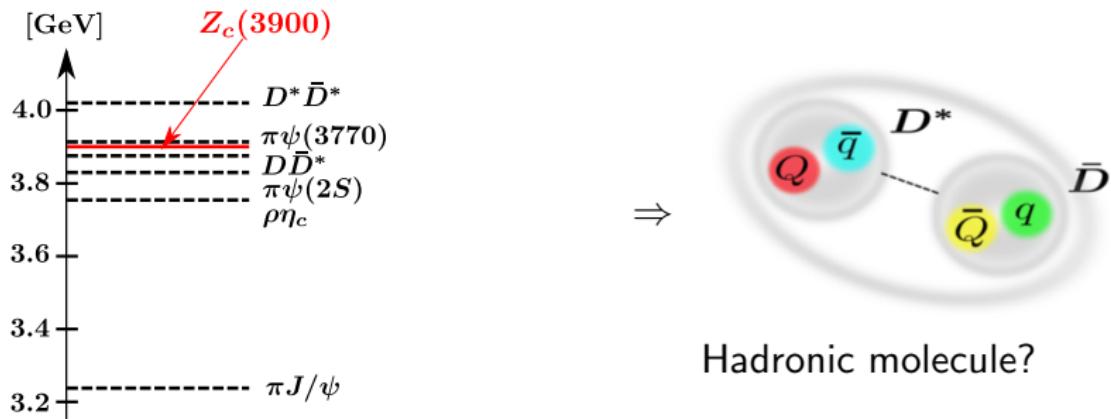


- Exotic state may be a **loosely bound state (resonance)** of the meson-meson.
⇒ Analogous to atomic nuclei (Deuteron: $B \sim 2.2$ MeV)

$Z_c(3900)$: Hadronic Molecules?

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- $Z_c(3900)$ close to the hadron-hadron thresholds



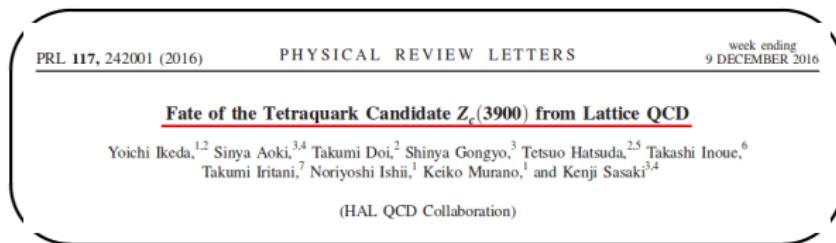
- Exotic state may be **a loosely bound state (resonance)** of the meson-meson.
⇒ Analogous to atomic nuclei (Deuteron: $B \sim 2.2$ MeV)
- $D\bar{D}^*$ molecule? **π exchange**

ref. A. Hosaka et al. PTEP 2016 (2016)062C01, A. Esposito, et al., Phys. Rept. 668 (2016)1, ...

$Z_c(3900)$: Lattice QCD (Numerical Experiments)

Introduction

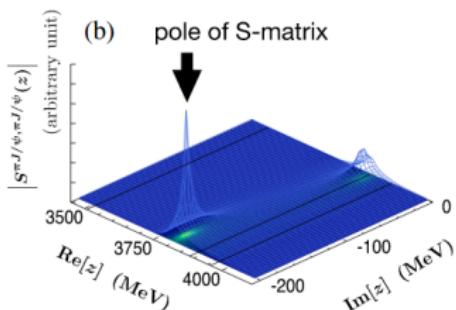
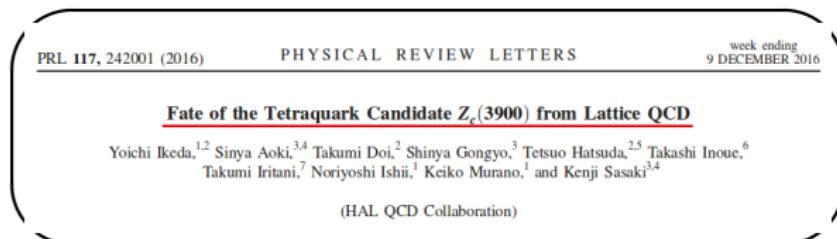
- Lattice QCD simulation by HALQCD at $m_\pi = 410 - 700$ MeV
- ⇒ Coupled-channel $\pi J/\psi - \rho\eta_c - D\bar{D}^*$



$Z_c(3900)$: Lattice QCD (Numerical Experiments)

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⇒ **Virtual state** is obtained.

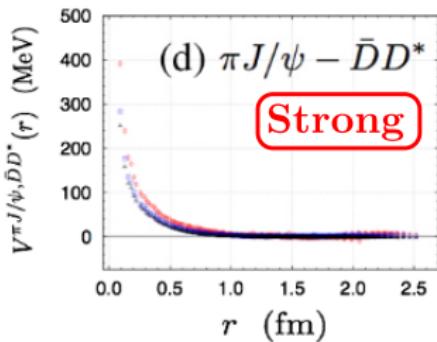
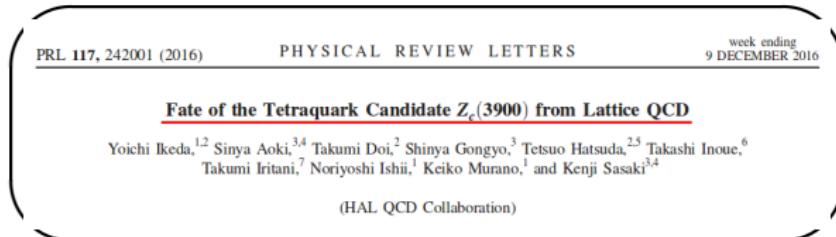
- $Z_c(3900)$ is Threshold cusp

Ikeda, et al., PRL 117(2016)242001

$Z_c(3900)$: Lattice QCD (Numerical Experiments)

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- ⇒ Coupled-channel $\pi J/\psi - \rho\eta_c - D\bar{D}^*$



- ⇒ **Virtual state** is obtained.
- $Z_c(3900)$ is Threshold cusp induced by $\pi J/\psi - \bar{D}D^*$ potential

Charm quark exchange?

Ikeda, et al., PRL 117 (2016) 242001

Bound state? Threshold cusp? → Hadron int.

Introduction

**Exotic structure:
Bound state? Cusp?**

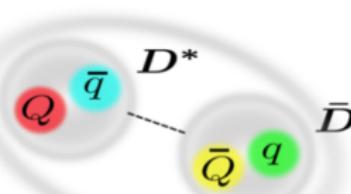
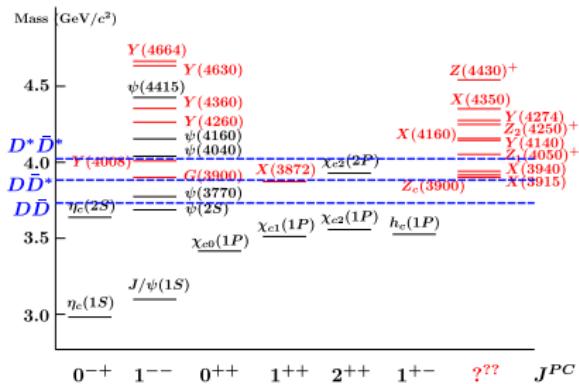
Bound state? Threshold cusp? → Hadron int.

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Bound state? Cusp?

↔ Hadron-hadron interaction

- Hadron-hadron interaction is important to understand the nature of exotic states! not only Z_c but also others.



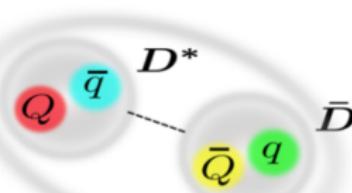
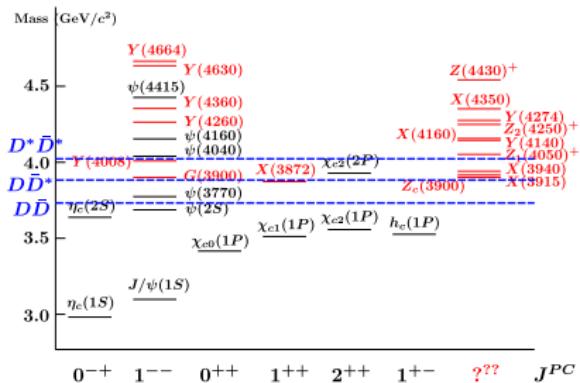
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Exotic structure:
Bound state? Cusp?

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→ Hadron int. is not established yet...

Model of Hadron-hadron interaction

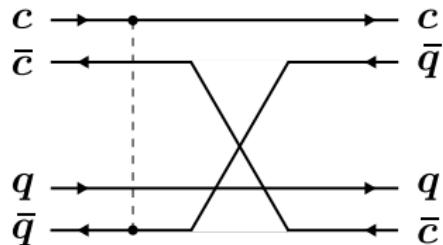
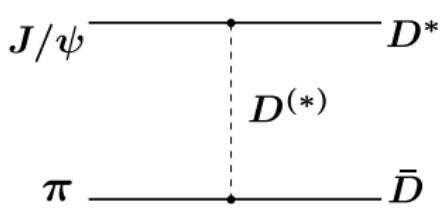
Introduction

- **Long-range force:** one π exchange potential (OPEP)
Lightest meson π , Importance in the nuclear force,
Heavy Quark Symmetry

Model of Hadron-hadron interaction

Introduction

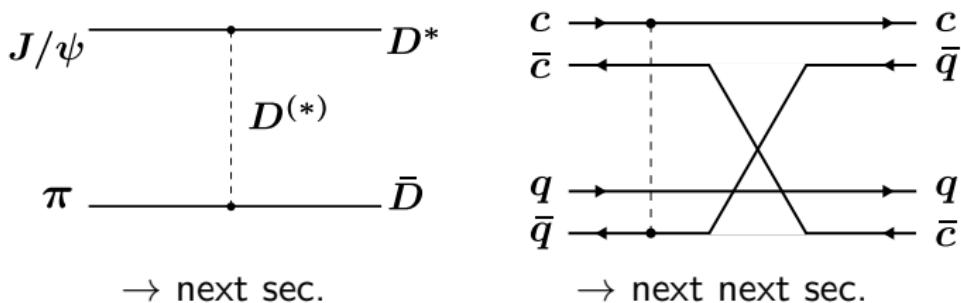
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- **Short-range force:** heavy meson exchange, quark exchange
 - (a) $D^{(*)}$ meson exchange
 - (b) Quark exchange



Model of Hadron-hadron interaction

Introduction

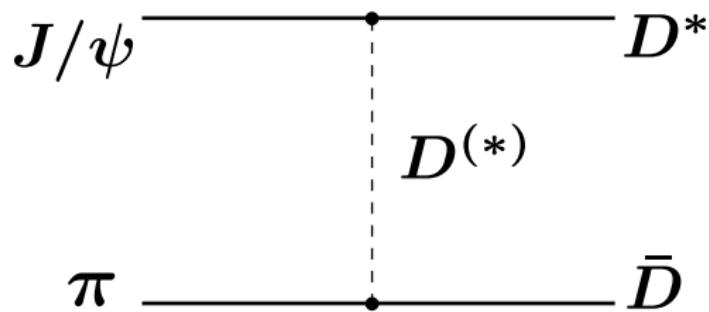
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- How can we understand **strong $\pi J/\psi - D\bar{D}^*$ potential?**
→ Very short range interaction due to large c mass

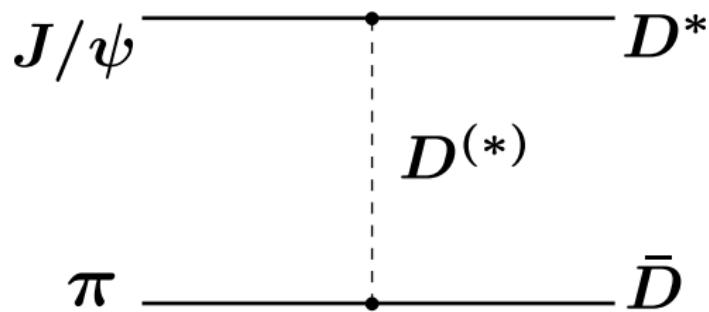
Today: Let us focus on $Z_c(3900)$ as $\pi J/\psi - D\bar{D}^*$

Meson exchange model



- $D^{(*)}\bar{D}^{(*)} - D^{(*)}\bar{D}^{(*)}$: π, ρ, ω
- $\pi J/\psi - D^{(*)}\bar{D}^{(*)}$: D, D^*

Meson exchange model



- $D^{(*)}\bar{D}^{(*)} - D^{(*)}\bar{D}^{(*)}$: π , ρ , ω
- $\pi J/\psi - D^{(*)}\bar{D}^{(*)}$: D , D^*

Heavy quark symmetry and OPEP

Meson exchange model

Heavy Quark Spin Symmetry

Heavy Quark Spin Symmetry

Charm (c), Bottom (b), Top (t)

Heavy Quark Spin Symmetry

Charm (c), Bottom (b), Top (t)



1. Coupled channels
2. Tensor force (OPEP)

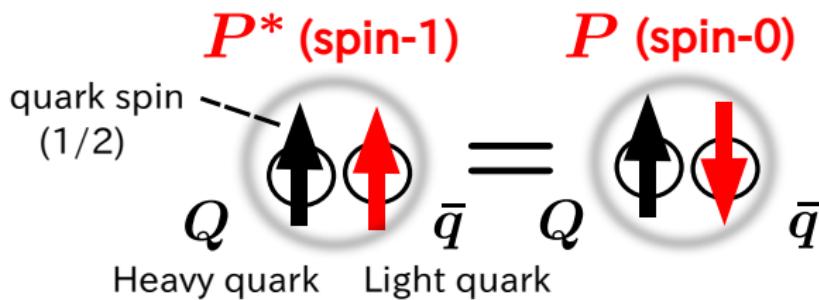
Heavy Quark Spin Symmetry and Mass degeneracy

Meson exchange model

Heavy Quark Spin Symmetry (HQS)

N.Isgur,M.B.Wise,PLB232(1989)113

- **Suppression of Spin-spin force** in $m_Q \rightarrow \infty$.
⇒ **Mass degeneracy** of hadrons with the different J
- e.g. $Q\bar{q}$ meson

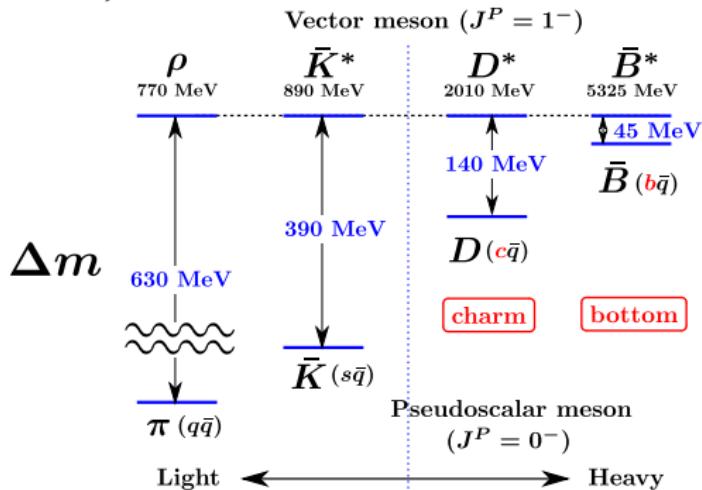


⇒ Mass degeneracy of spin-0 and spin-1 states!

Mass degeneracy of heavy hadrons

Meson exchange model

- Mass difference between vector and pseudoscalar mesons.
 $(Q\bar{q}, q = u, d)$

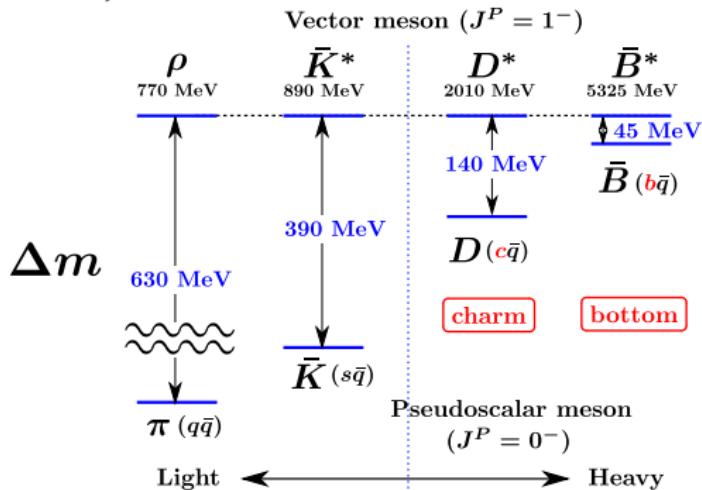


- Δm decreases when the quark mass increases.

Mass degeneracy of heavy hadrons

Meson exchange model

- Mass difference between vector and pseudoscalar mesons.
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- Δm decreases when the quark mass increases.
⇒ **Degeneracy of Heavy hadrons!**

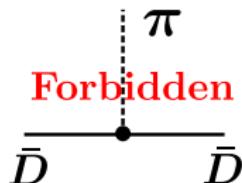
For $Z_c(3900)$, $D - D^*$ mixing

Heavy hadron- π coupling

Meson exchange model

- Effective Lagrangians: Heavy hadron and π

R. Casalbuoni *et al.*, Phys.Rept.**281** (1997)145, T. M. Yan, *et al.*, PRD**46**(1992)1148



- Heavy meson: $\bar{D}^{(*)}\bar{D}^{(*)}\pi$ (**$DD\pi$: Parity violation**)

$$\mathcal{L}_{\pi HH} = -\frac{g_\pi}{2f_\pi} \text{Tr} [H\gamma_\mu\gamma_5\partial^\mu\hat{\pi}\bar{H}], \quad H = \frac{1+\gamma}{2} [D_\mu^*\gamma^\mu - D\gamma_5]$$

- Doublet D and D^* with one coupling const. $g_\pi = 0.59$ (from $D^* \rightarrow D\pi$ decay)
- Form factor (Hadron has **finite size**)

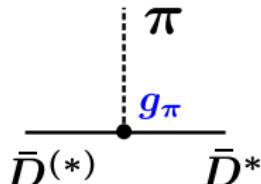
$$F(q^2) = \frac{\Lambda^2 - m_\pi^2}{\Lambda^2 - q^2}, \quad \Lambda_{\bar{D}} \sim 1130 \text{ MeV}$$

Heavy hadron- π coupling

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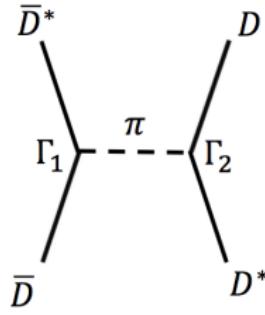
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One pion exchange potential in $D^{(*)}\bar{D}^{(*)}$

Meson exchange model

- One boson exchange potential (OBEP)



Born amplitude

→ Non-relativistic Potential

OPEP

$$V^\pi = -\frac{1}{2} \left(\frac{g_\pi}{f_\pi} \right)^2 \left[\vec{S}_1 \cdot \vec{S}_2 C(r) + S_{12}(\hat{r}) T(r) \right] \vec{\tau}_1 \cdot \vec{\tau}_2$$

$$C(r) = m_\pi^2 \left(\frac{e^{-m_\pi r}}{r} - \frac{e^{-\Lambda r}}{r} - \frac{\Lambda^2 - m_\pi^2}{2\Lambda} e^{-\Lambda r} \right)$$

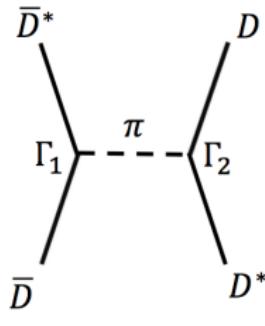
Comments

- HQS induces $D\bar{D}^* - D^*\bar{D}^*$ couplings → OPEP works!

One pion exchange potential in $D^{(*)}\bar{D}^{(*)}$

Meson exchange model

- One boson exchange potential (OBEP) **with Tensor force!**



Born amplitude

→ Non-relativistic Potential

OPEP

$$V^\pi = -\frac{1}{2} \left(\frac{g_\pi}{f_\pi} \right)^2 \left[\vec{S}_1 \cdot \vec{S}_2 C(r) + \mathbf{S}_{12}(\hat{r}) \mathbf{T}(r) \right] \vec{r}_1 \cdot \vec{r}_2$$

$$C(r) = m_\pi^2 \left(\frac{e^{-m_\pi r}}{r} - \frac{e^{-\Lambda r}}{r} - \frac{\Lambda^2 - m_\pi^2}{2\Lambda} e^{-\Lambda r} \right)$$

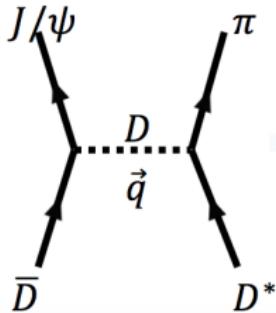
Comments

- HQS induces $D\bar{D}^* - D^*\bar{D}^*$ couplings → OPEP works!
- Tensor force $T(r) \Rightarrow$ the driving force in atomic nuclei
 $S_{12}(\hat{r}) = 3(\vec{S}_1 \cdot \hat{r})(\vec{S}_2 \cdot \hat{r}) - \vec{S}_1 \cdot \vec{S}_2 \rightarrow S-D$ mixing

Heavy meson exchange potential

Meson exchange model

- $D^{(*)}$ meson exchange potential in $\pi J/\psi - D^{(*)}\bar{D}^{(*)}$



D exchange

$$V^D = \frac{2}{3} \frac{g_\psi g_\pi}{f_\pi \sqrt{E_\pi}} [\vec{S}_1 \cdot \vec{S}_2 C(r) + S_{12}(\hat{r}) T(r)]$$

D^* exchange

$$V^{D^*} = \frac{2}{3} \frac{g_\psi g_\pi}{f_\pi \sqrt{E_\pi}} [2\vec{S}_1 \cdot \vec{S}_2 C(r) - S_{12}(\hat{r}) T(r)]$$

$$g_\psi = 8$$

A. Deandrea, G. Nardulli and A. D. Polosa, PRD **68**(2003)034002

Comments

- $D^{(*)}$ meson exchange gives the $\pi J/\psi - D^{(*)}\bar{D}^{(*)}$ potential.
Hidden \leftrightarrow Open-Open
- $D^{(*)}$ mass ~ 2 GeV $\Leftrightarrow 1/m_{D^{(*)}} \sim 0.1$ fm
Does it work?

Numerical results: Phase shift

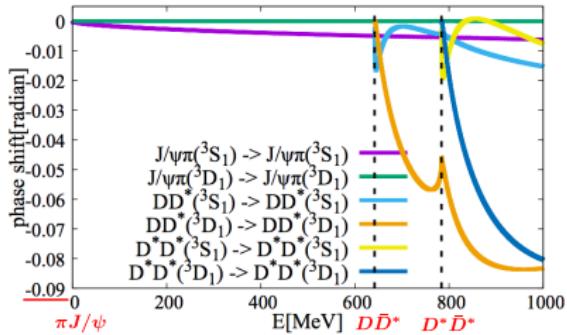
Meson exchange model

- We found...

Numerical results: Phase shift

Meson exchange model

- We found... **No Bound state, No Resonance**
Very Small phase shift $|\delta| < 0.09$ [rad]

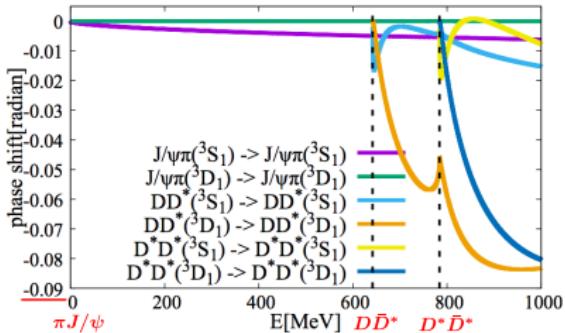


- $D^{(*)}\bar{D}^{(*)}$ channel: **Small** contribution from OPEP
- $\pi J/\psi$ channel: $D^{(*)}$ exchange is **Negligible**

Numerical results: Phase shift

Meson exchange model

- We found... **No Bound state, No Resonance**
Very Small phase shift $|\delta| < 0.09$ [rad]



- $D^{(*)}\bar{D}^{(*)}$ channel: **Small** contribution from OPEP
Why?: Isospin factor $\vec{\tau}_1 \cdot \vec{\tau}_2$, $I = 0: -3$, but $I = 1: +1$
- $\pi J/\psi$ channel: $D^{(*)}$ exchange is **Negligible**
Why?: Suppression by the form factor (finite hadron size)

D meson exchange \rightarrow Quark exchange

Meson exchange model

- No resonance ← agreeing with Lattice QCD result
- ↔ We cannot explain the strong $\pi J/\psi - D\bar{D}^*$ potential.

D meson exchange \rightarrow Quark exchange

Meson exchange model

- No resonance \leftarrow agreeing with Lattice QCD result
- ⇒ We cannot explain the strong $\pi J/\psi - D\bar{D}^*$ potential.

Problems in $D^{(*)}$ exchange potential

- Too large mass, $1/m_{D^{(*)}} \sim 0.1$ fm
- In such short range region, “Hadron” is not good effective d.o.f. ?



Quark exchange interaction!

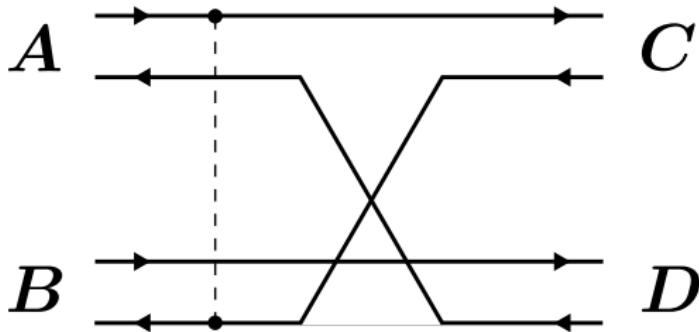
→next section

Quark exchange model (on progress...)

- Born-order quark-exchange diagram

T. Barnes and E. S. Swanson, PRD **46**(1992)131. Swanson, Ann. Phys. **220**(1992)73.

- $AB \rightarrow CD$ scattering $\mathcal{M}_{fi} \propto \langle C, D | H_I | A, B \rangle$



- Ingredients: Meson Wavefunctions(A, B, C, D)
Quark interaction (Quark Model)
- Born amplitude \Rightarrow Meson-meson Potential can be obtained

Quark Model

Model Setup

- Quark Hamiltonian (One gluon exchange + Linear potentials)

Barnes and Swanson, PRD**46**(1992)131.; Swanson, Ann. Phys. **220**(1992)73.

$$H_{ij}^q = K_q + \left(-\frac{3}{4} br + \frac{\alpha_s}{r} - C \right) \vec{F}_i \cdot \vec{F}_j \\ - \frac{8\pi\alpha_h}{3m_i m_j} \left(\frac{\sigma^3}{\pi^{3/2}} e^{-\sigma^2 r_{ij}^2} \right) \vec{S}_i \cdot \vec{S}_j \vec{F}_i \cdot \vec{F}_j$$

- Parameters are fixed to reproduce the mass of mesons

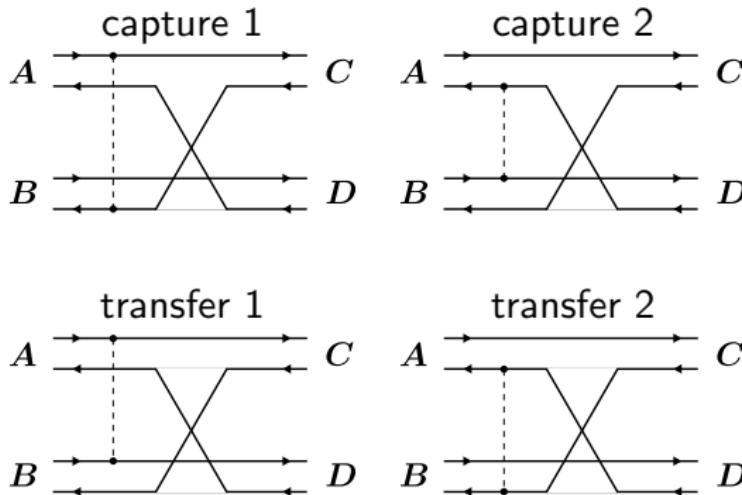
Table: Quark Model Parameters from Ann.Phys.**220**(1992)73.

$m_q = 0.375 \text{ GeV}$	$m_c = 1.9 \text{ GeV}$
$\alpha_s = 0.857$	$\alpha_h = 0.840$
$b = 0.154 \text{ GeV}^{-2}$	$C = -0.4358 \text{ GeV}$
$\sigma = 0.70 \text{ GeV}$	

Scattering Amplitude

Model Setup

- Born quark exchange diagrams T. Barnes and E. S. Swanson, PRD **46**, 131 (1992).
Quark interaction between Mesons \Rightarrow Four diagrams

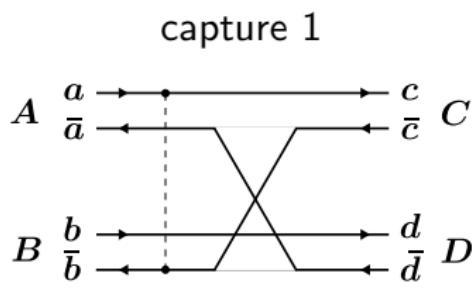


- Scattering Amplitude $\mathcal{M}_{fi} \propto \langle C, D | H^q | A, B \rangle$

$$\mathcal{M}_{fi}^{tot} = \mathcal{M}_{fi}^{capture1} + \mathcal{M}_{fi}^{capture2} + \mathcal{M}_{fi}^{transfer1} + \mathcal{M}_{fi}^{transfer2}$$

Scattering Amplitude

Model Setup



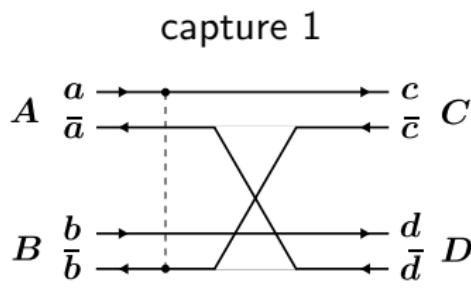
- ▶ Meson momenta: A, B, C, D
- ▶ Quark momenta:
 $a, \bar{a}, b, \bar{b}, c, \bar{c}, d, \bar{d}$
- ▶ Conservation:
 $A + B = C + D,$
 $\bar{a} = \bar{d}, b = d$

• Amplitude

$$\rightarrow \int \int d^3 a d^3 c \phi_C^*(2\vec{c} - \vec{C}) \phi_D^*(2\vec{a} - 2\vec{A} - \vec{C}) V(\vec{a} - \vec{c}) \phi_A(2\vec{a} - \vec{A}) \phi_B(2\vec{a} - \vec{A} - 2\vec{C})$$

Scattering Amplitude

Model Setup



- ▶ Meson momenta: A, B, C, D
- ▶ Quark momenta:
 $a, \bar{a}, b, \bar{b}, c, \bar{c}, d, \bar{d}$
- ▶ Conservation:
 $A + B = C + D,$
 $\bar{a} = \bar{d}, b = d$

- Amplitude

$$\rightarrow \int \int d^3 a d^3 c \phi_C^*(2\vec{c} - \vec{C}) \phi_D^*(2\vec{a} - 2\vec{A} - \vec{C}) V(\vec{a} - \vec{c}) \phi_A(2\vec{a} - \vec{A}) \phi_B(2\vec{a} - \vec{A} - 2\vec{C})$$

- Potentials (momentum space)

Coulomb: $V^{Coul}(q) = -\frac{\alpha_s}{2\pi^2} \frac{1}{\vec{q}^2}$, **Hyperfine:** $V^{Hyp}(q) = -\frac{8\pi\alpha_h}{3m_i m_j} e^{-\vec{q}^2/4\sigma^2}$

Linear (Regularized):

$$V^{Lin}(r) = br \times e^{-\varepsilon r} \rightarrow V^{Lin}(q) = b \left[\frac{-8\pi}{(\vec{q}^2 + \varepsilon^2)^2} + \frac{32\pi\varepsilon^2}{(\vec{q}^2 + \varepsilon^2)^3} \right]$$

Numerical Results

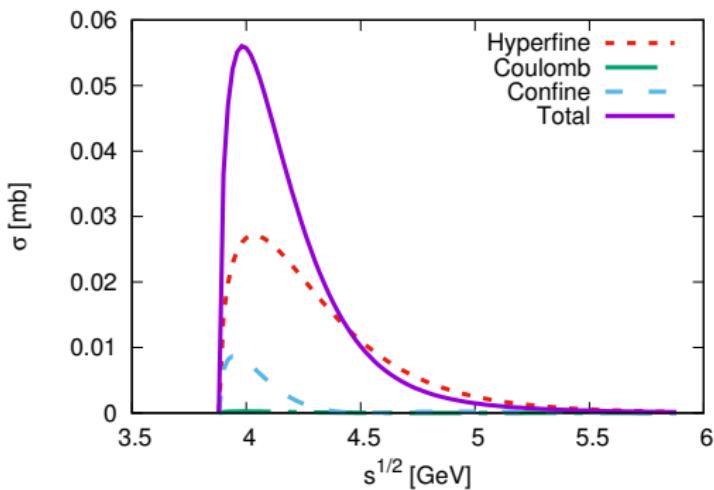
(on progress...)

- $\pi J/\psi - D^{(*)} \bar{D}^{(*)}$ Cross Section (Born-term)

Cross Section (Born term): $\pi J/\psi - D\bar{D}^*$

Numerical Result

- $\pi J/\psi - D\bar{D}^*$: Amplitude \rightarrow Cross section



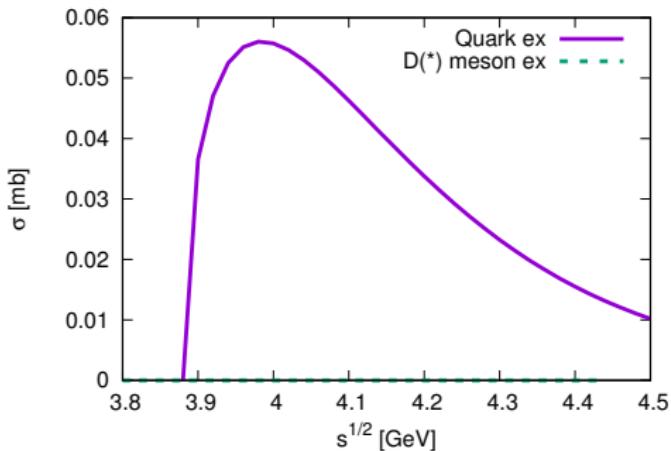
- Cross section $\Rightarrow |(\text{Coulomb}) + (\text{Confine}) + (\text{Hyperfine})|^2$
- Dominant role of the Hyperfine (Spin-spin) term
↔ Minor role of the Coulomb one.

Cross Section: Quark exchange vs $D^{(*)}$ exchange

Numerical Result

- $\pi J/\psi - D\bar{D}^*$: Amplitude → Cross section

(i) Quark ex vs $D^{(*)}$ ex



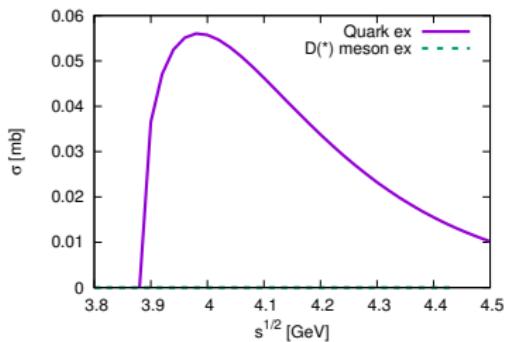
- Comparing results of Quark exchange and $D^{(*)}$ exchange
- Can you see a dashed line ($D^{(*)}$ exchange)?

Cross Section: Quark exchange vs $D^{(*)}$ exchange

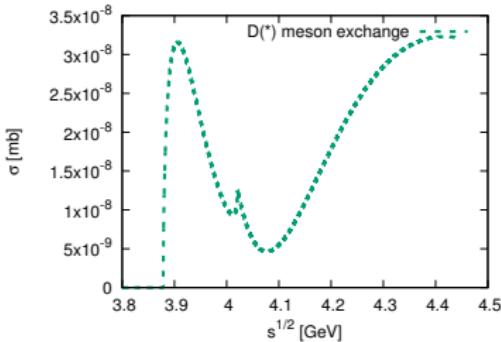
Numerical Result

- $\pi J/\psi - D\bar{D}^*$: Amplitude \rightarrow Cross section

(i) Quark ex. vs $D^{(*)}$ ex.



(ii) $D^{(*)}$ ex. (Zoom)



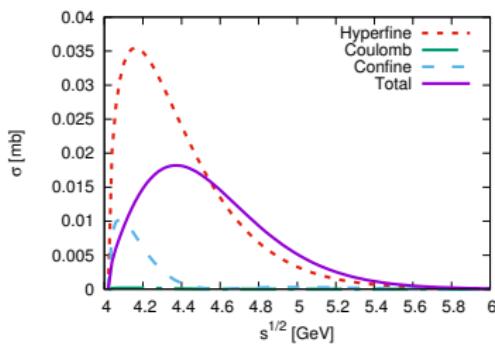
- Comparing results of Quark exchange and $D^{(*)}$ exchange
- Can you see a dashed line ($D^{(*)}$ exchange)? $< 3.5 \times 10^{-8}$ mb
- Large difference between Quark exchange and $D^{(*)}$ exchange!

Cross Section (Born): $\pi J/\psi - D^* \bar{D}^*$

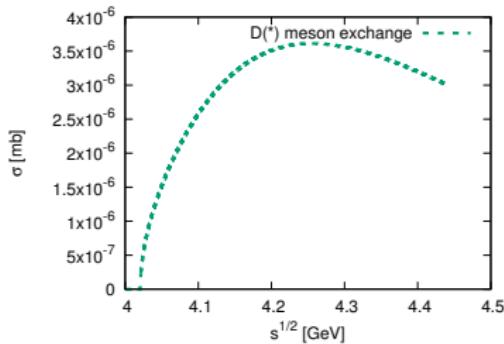
Numerical Result

- $\pi J/\psi - D^* \bar{D}^*$: Amplitude → Cross section

(i) Quark ex vs $D^{(*)}$ ex

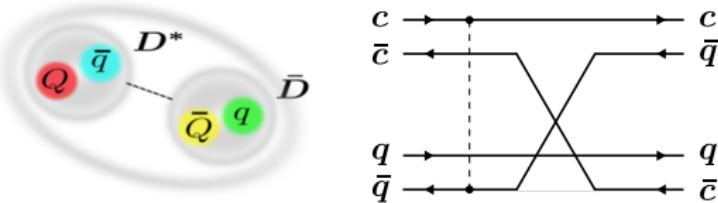


(ii) $D^{(*)}$ exchange



- Cross section $\Rightarrow |(\text{Coulomb}) + (\text{Confine}) + (\text{Hyperfine})|^2$
- **Dominant role of the Spin-spin term.**
Coulomb term \Rightarrow Very small contribution.
- Large difference between Quark exchange and $D^{(*)}$ exchange!

Summary



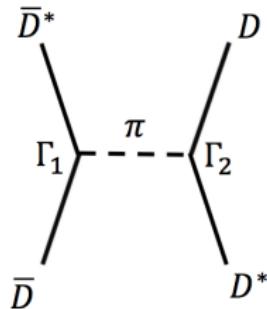
- Many exotic states near the threshold.
→ Understanding **the hadron-hadron interaction** is needed.
- Charged charmonium $Z_c(3900)$ has been discussed as the Hadronic molecules or the threshold cusp.
- OPEP contribution is not strong. $D^{(*)}$ meson exchange is **negligible**.
- Quark exchange interaction is introduced as Short range $\pi J/\psi - D^{(*)}D^{(*)}$ potential.
We find **Large difference** between results from Quark exchange and $D^{(*)}$ meson exchange.

Back Up

One pion exchange potential in $D^{(*)}\bar{D}^{(*)}$

Meson exchange model

- One boson exchange potential (OBEP)



OPEP

$$V^\pi = -\frac{1}{2} \left(\frac{g_\pi}{f_\pi} \right)^2 \left[\vec{S}_1 \cdot \vec{S}_2 C(r) + S_{12}(\hat{r}) T(r) \right] \vec{\tau}_1 \cdot \vec{\tau}_2$$

Vector (ρ, ω) exchange

$$V^\nu = - \left(\frac{\lambda g_V}{\sqrt{3}} \right)^2 \left[2 \vec{S}_1 \cdot \vec{S}_2 C(r) - S_{12}(\hat{r}) T(r) \right] \vec{\tau}_1 \cdot \vec{\tau}_2$$

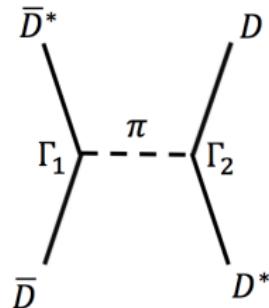
Comments

- Tensor force $T(r) \Rightarrow$ **the driving force** in atomic nuclei
 $S_{12}(\hat{r}) = 3(\vec{S}_1 \cdot \hat{r})(\vec{S}_2 \cdot \hat{r}) - \vec{S}_1 \cdot \vec{S}_2 \rightarrow S\text{-}D$ mixing

One pion exchange potential in $D^{(*)}\bar{D}^{(*)}$

Meson exchange model

- One boson exchange potential (OBEP) **with Tensor force!**



OPEP

$$V^\pi = -\frac{1}{2} \left(\frac{g_\pi}{f_\pi} \right)^2 \left[\vec{S}_1 \cdot \vec{S}_2 C(r) + \mathbf{S}_{12}(\hat{r}) \mathbf{T}(r) \right] \vec{\tau}_1 \cdot \vec{\tau}_2$$

Vector (ρ, ω) exchange

$$V^\nu = - \left(\frac{\lambda g_V}{\sqrt{3}} \right)^2 \left[2 \vec{S}_1 \cdot \vec{S}_2 C(r) - \mathbf{S}_{12}(\hat{r}) \mathbf{T}(r) \right] \vec{\tau}_1 \cdot \vec{\tau}_2$$

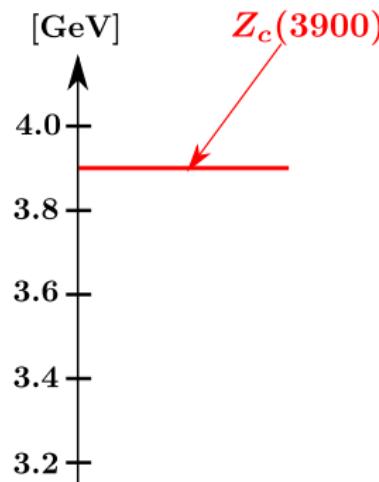
Comments

- Tensor force $T(r) \Rightarrow$ **the driving force** in atomic nuclei
 $S_{12}(\hat{r}) = 3(\vec{S}_1 \cdot \hat{r})(\vec{S}_2 \cdot \hat{r}) - \vec{S}_1 \cdot \vec{S}_2 \rightarrow S-D$ mixing
- G -parity of vector mesons: ρ ($G = -1$), ω ($G = +1$)
 \Rightarrow **Working against each other**, $\rho + \omega$ has a minor role...

Meson-meson thresholds

Model Setup

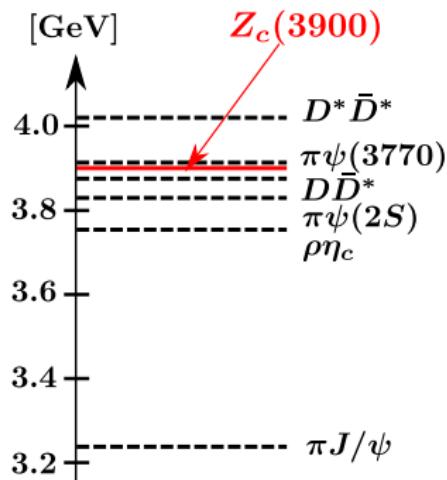
- Meson-meson thresholds,



Meson-meson thresholds

Model Setup

- Meson-meson thresholds, $\pi J/\psi$, $\rho\eta_c$, $\pi\psi(2S)$, $D\bar{D}^*$, $D^*\bar{D}^*$, ...



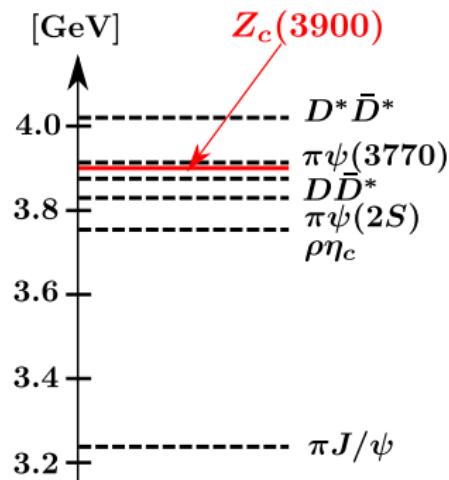
Coupled-Channels

$$\left\{ \begin{array}{l} \pi J/\psi \\ \pi\psi(2S) \\ \rho\eta_c \\ D\bar{D}^* \\ D^*\bar{D}^* \\ \vdots \end{array} \right\} - \left\{ \begin{array}{l} \pi J/\psi \\ \pi\psi(2S) \\ \rho\eta_c \\ D\bar{D}^* \\ D^*\bar{D}^* \\ \vdots \end{array} \right\}$$

Meson-meson thresholds

Model Setup

- Meson-meson thresholds, $\pi J/\psi$, $\rho\eta_c$, $\pi\psi(2S)$, $D\bar{D}^*$, $D^*\bar{D}^*$, ...



Coupled-Channels

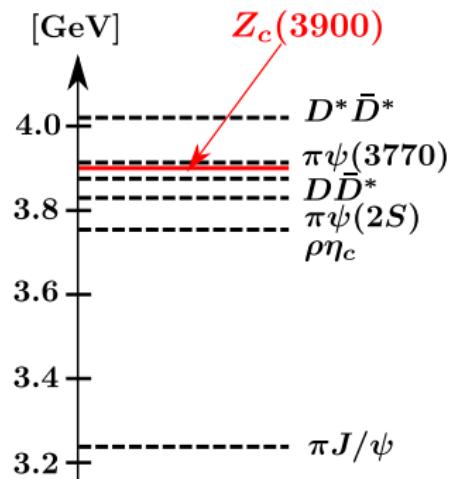
$$\left\{ \begin{array}{l} \pi J/\psi \\ \pi\psi(2S) \\ \rho\eta_c \\ D\bar{D}^* \\ D^*\bar{D}^* \\ \vdots \end{array} \right\} - \left\{ \begin{array}{l} \pi J/\psi \\ \pi\psi(2S) \\ \rho\eta_c \\ D\bar{D}^* \\ D^*\bar{D}^* \\ \vdots \end{array} \right\}$$

- Born-order quark-exchange
⇒ Applicable to charm exchange (**Hidden** ↔ **Open-Open**)

Meson-meson thresholds

Model Setup

- Meson-meson thresholds, $\pi J/\psi$, $\rho\eta_c$, $\pi\psi(2S)$, $D\bar{D}^*$, $D^*\bar{D}^*$, ...



Coupled-Channels

$$\left\{ \begin{array}{l} \textcolor{red}{\pi J/\psi} \\ \textcolor{red}{\pi\psi(2S)} \\ \textcolor{red}{\rho\eta_c} \\ \textcolor{blue}{D\bar{D}^*} \\ \textcolor{blue}{D^*\bar{D}^*} \\ \vdots \end{array} \right\} - \left\{ \begin{array}{l} \textcolor{red}{\pi J/\psi} \\ \textcolor{red}{\pi\psi(2S)} \\ \textcolor{red}{\rho\eta_c} \\ \textcolor{blue}{D\bar{D}^*} \\ \textcolor{blue}{D^*\bar{D}^*} \\ \vdots \end{array} \right\}$$

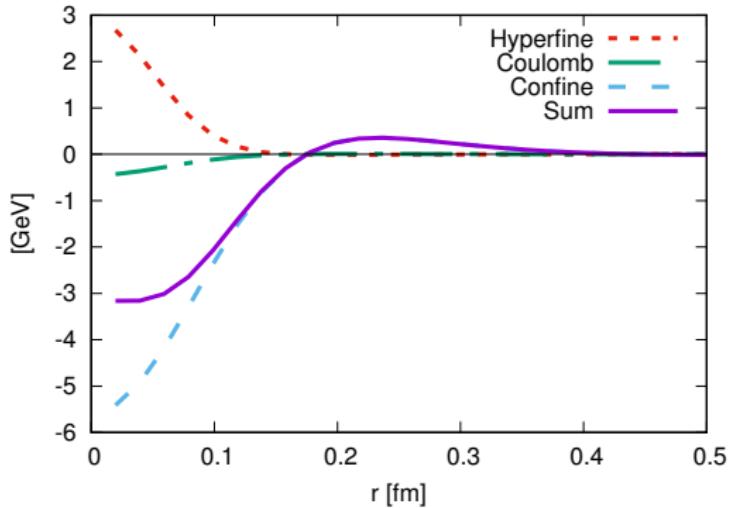
- Born-order quark-exchange
⇒ Applicable to charm exchange (**Hidden** ↔ **Open-Open**)

Today: $\pi J/\psi - D\bar{D}^*$ and $\pi J/\psi - D^*\bar{D}^*$ (*S*-wave)

$\pi J/\psi - D\bar{D}^*$ Potentials (on the threshold)

Numerical Result

- Energy-dependent $\pi J/\psi - D\bar{D}^*$ potential
- Result @ $E = m_D + m_{\bar{D}^*}$ ($D\bar{D}^*$ threshold)

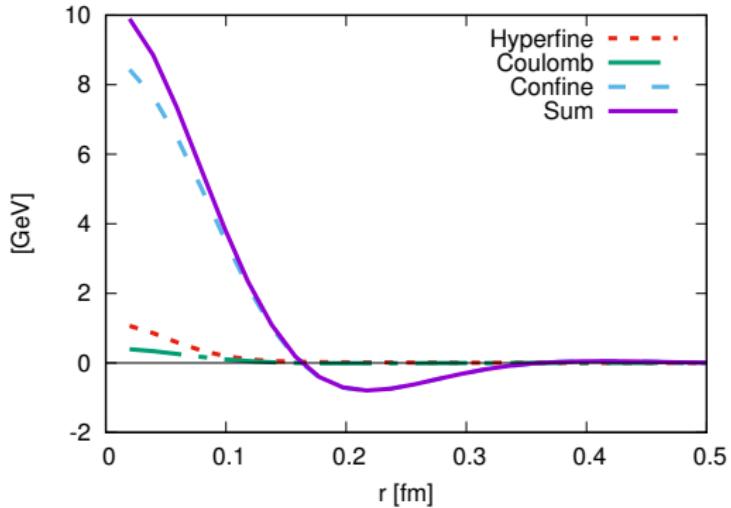


- Hyperfine (Spin-spin) and Confinement potentials
⇒ **Working against each other**

$\pi J/\psi - D^* \bar{D}^*$ Potentials (on the threshold)

Numerical Result

- Energy-dependent $\pi J/\psi - D^* \bar{D}^*$ potential
- Result @ $E = m_{D^*} + m_{\bar{D}^*}$ ($D^* \bar{D}^*$ threshold)

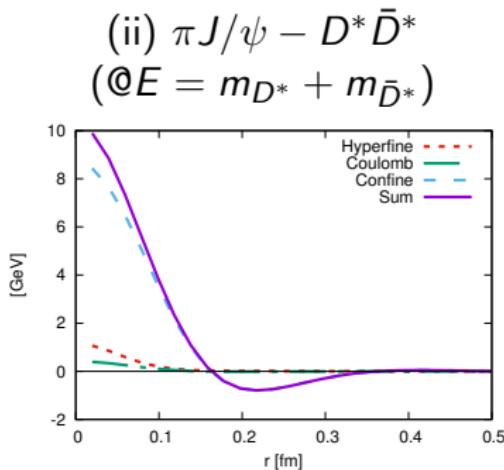
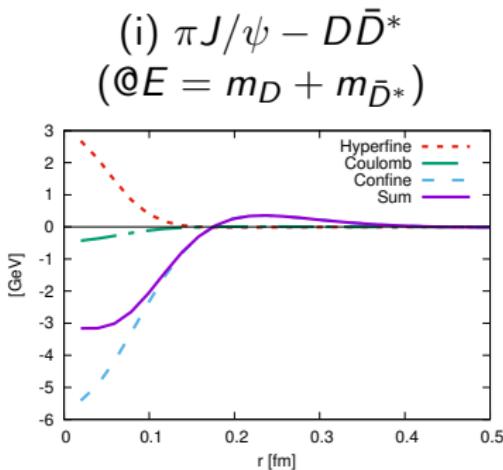


- Hyperfine (Spin-spin) and Confinement potentials
⇒ **Cooperating with each other**

Potentials

Numerical Result

- Comparing $\pi J/\psi - D\bar{D}^*$ with $\pi J/\psi - D^*\bar{D}^*$



- Similar?
- Degenerate in the heavy quark limit?