Channeling of Radiations: from Crystal Undulators to Capillary Waveguides

Russian Academy of Sciences



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@ Channeling: Orientational Effects of Transmission

• 1962-63:

Robinson &

Oen:

Piercy &

Lutz:

• 1965:

Lindhard:

• • • •

Andersen

Uggerhoj

Kagan

Kononez

Firsov

Tsyganov

Gibson

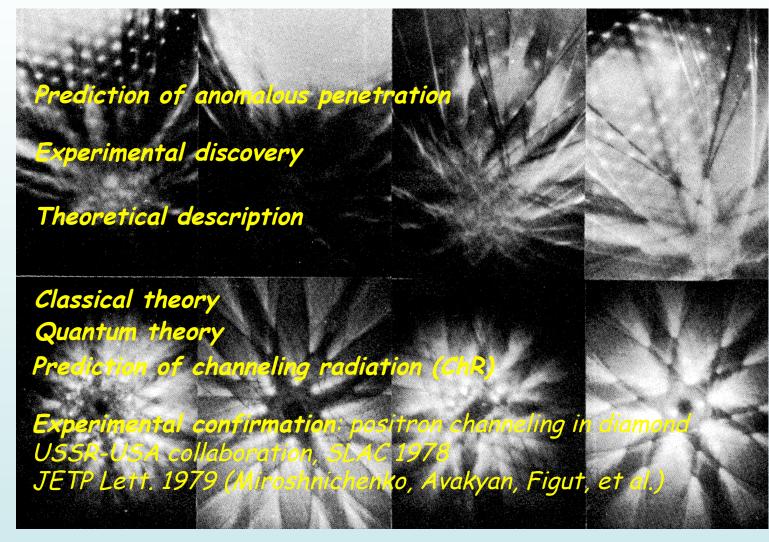
Kumakhov

Beloshitsky

Gemmel

Appleton

.

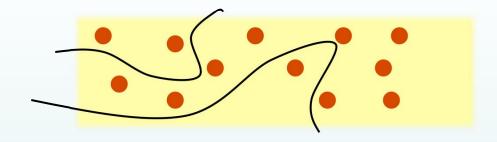


Rev. Mod. Phys. 1974: 1 MeV e⁻ @ Cu crystal

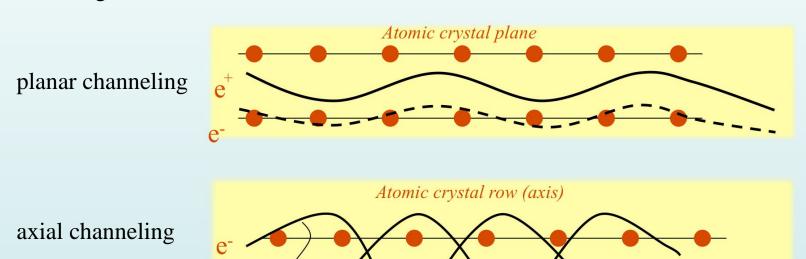
more than 1000 articles + a number of monographs

@ Channeling of Charged Particles

@ Amorphous:



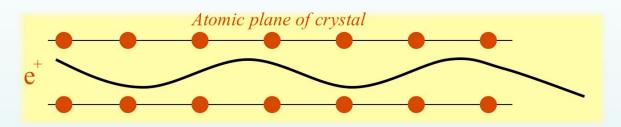
@ Channeling:



$$\varphi <<1$$
 $(\varphi < \varphi_L \sim \sqrt{U/E})$ - the Lindhard angle is the critical angle for the channeling

@ Channeling of Charged Particles & Channeling Radiation

@ Channeling:

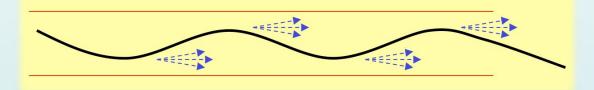


$$\varphi << 1$$
 $(\varphi < \varphi_L \sim \sqrt{U/E})$

- the Lindhard angle is the critical angle for the channeling

@ Channeling Radiation:

$$\omega = \omega(\theta) = \frac{\omega_{fi}}{1 - \beta_{\parallel} \cos \theta}$$



$$\omega_{fi}$$
 - optical frequency

$$\longrightarrow \omega_0 \gamma^{3/2} - \omega_0 \gamma^2$$

Powerful radiation source of X-rays and γ-rays:

$$(\gamma >> 1)$$

@ Channeling: Continuum model

$$V(r) = \frac{Z_1 Z_2 e^2}{r} \left(\varphi(r/a) \right)$$

screening function of Thomas-Fermi type

$$a = .8853a_0 \left(Z_1^{1/2} + Z_2^{1/2}\right)^{-2/3}$$

screening length

$$\varphi(r/a)$$
: $\sum_{i=1}^{3} \alpha_i \exp(-\beta_i r/a)$ Molier's potential
$$1 - \left[1 + \frac{Ca}{r^2}\right]^{-1/2} \qquad C^2 \approx 3 \quad \text{Lindhard potential}$$

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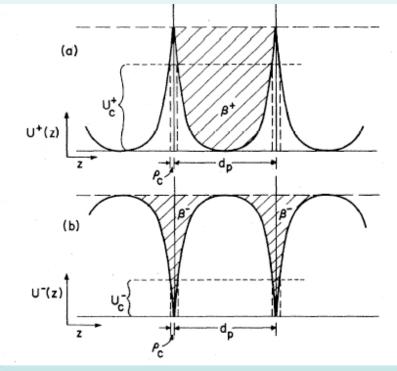
..... Firsov, Doyle-Turner, etc.

Lindhard:

Continuum model –

continuum atomic plane/axis potential

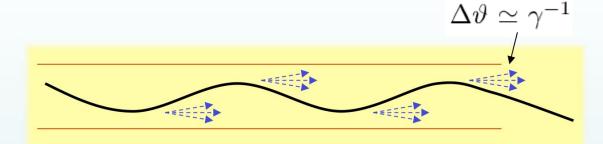
$$V_{RS}(\rho) = \frac{1}{d} \int_{-\infty}^{+\infty} V(\sqrt{\rho^2 + x^2}) dx$$



@ Channeling Radiation

@ Channeling Radiation:

$$\omega = \omega(\theta) = \frac{\omega_{fi}}{1 - \beta_{\parallel} \cos \theta}$$



$$\omega_{\it fi}$$
 - optical frequency

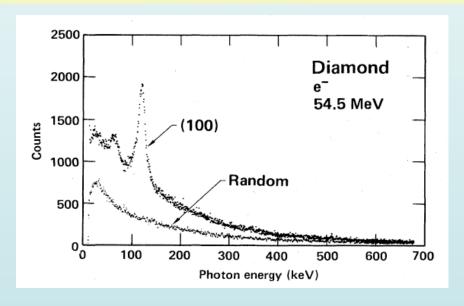
$$\longrightarrow$$
 Doppler effect \longrightarrow $\omega_0 \gamma^{3/2} \omega_0 \gamma^2$

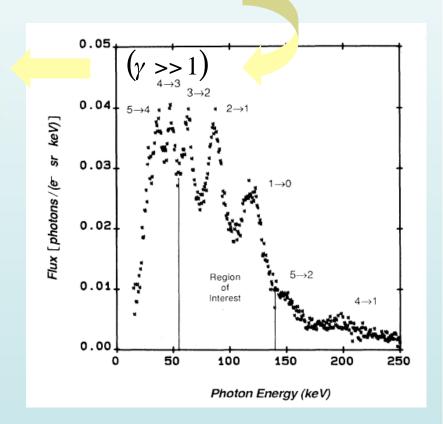
Powerful radiation source of X-rays and γ-rays:

polarized

•tunable

narrow forwarded





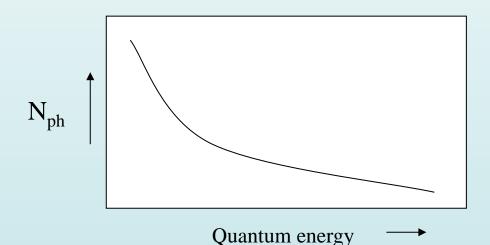
@ Bremsstrahlung & Coherent Bremsstrahlung vs Channeling Radiation

- @ amorphous electron:
 - •Radiation as sum of independent impacts with atoms
 - •Effective radius of interaction a_{TF}
 - •Coherent radiation length $l_{coh}>>a_{TF}$
 - •Deviations in trajectory less than effective radiation angles:

$$\Delta\theta \propto a_{TF}/p$$

$$\Delta \vartheta \simeq \gamma^{-1}$$

$$\left(\frac{d^2I}{d\omega\Omega}\right)_{BR}\simeq (\pi L_R)^{-1}\gamma^2\frac{1+\gamma^4\theta^4}{(1+\gamma^2\theta^2)^4} \longrightarrow \left(\frac{dI}{d\omega}\right)_{BR}\simeq \frac{4}{3}L_R^{-1}$$



@ Bremsstrahlung & Coherent Bremsstrahlung vs Channeling Radiation

@ interference of consequent radiation events:

phase of radiation wave
$$\longrightarrow$$
 $(\omega t - \mathbf{kr}(t))$

Radiation field as interference of radiated waves:

$$l_{coh} \approx \frac{v}{\omega - \mathbf{k}\dot{\mathbf{r}}} = \frac{\lambda\beta}{1 - \beta\cos\theta} \longrightarrow \boxed{l_{coh} \propto \gamma^2\lambda}$$

Coherent radiation length can be rather large even for short wavelength

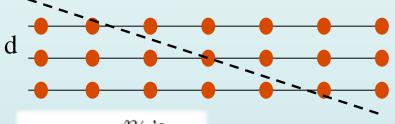
@ crystal:

$$l_1 = nl_{coh}$$

$$l_1 = n l_{coh}$$
 $l = d / \sin \alpha$

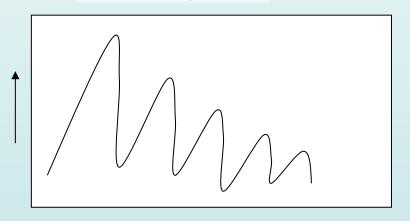
$$l_1 = \frac{n\lambda\beta}{1 - \beta\cos\theta} \quad \omega_0 \equiv \beta/l_1$$

$$\omega_0 \equiv \beta/l_1$$



$$\omega = \frac{n\omega_0}{1 - \beta\cos\theta}$$

$$\left(\frac{d^2I}{d\omega\Omega}\right)_{CBR} \propto \delta\left(\omega(1-\beta\cos\theta)-n\omega_0\right)$$



Quantum energy

@ Bremsstrahlung & Coherent Bremsstrahlung vs Channeling Radiation

@ crystal: $\alpha \to 0$ channeling

$$\omega = \frac{\omega_{fi}}{1 - \beta_{||} \cos \theta}$$

$$\left(\frac{dI}{d\omega}\right)_{CR} \propto \omega \left[1 - 2\left(\frac{\omega}{\omega_m}\right) + 2\left(\frac{\omega}{\omega_m}\right)^2\right], \ \omega \leq \omega_m \simeq 2\gamma^2 \omega_{fi}$$

 $\frac{ChR}{R} \propto \gamma^{1/2} Z^{-2/3}$ at definite conditions channeling radiation can be significantly powerful than bremsstrahlung

B: CB: ChR:
$$N \iff l_{coh} \propto \gamma^2 / \omega$$
 N_{eff} $\propto NZ^2$ $\propto (NZ)^2$ $\propto (N_{eff}Z)^2$

@ Channeling Radiation & Thomson Scattering

$$\omega_{lab}^{ChR} \approx \frac{2\gamma^2}{1 + \theta^2 \gamma^2} \omega_0^{ChR}$$
 $\propto \gamma^{3/2}$

$$\omega_{lab}^{ChR} \approx \frac{2\gamma^2}{1 + \theta^2 \gamma^2} \omega_0^{ChR} - \text{radiation frequency} - \omega_{lab}^{TS} \left\{ \begin{array}{l} \vartheta = 0 \\ \vartheta = \pi/2 \\ \vartheta = \pi \end{array} \right\} \simeq \left\{ \begin{array}{l} 1 \\ 2 \\ 4 \end{array} \right\} \frac{\gamma^2}{1 + \vartheta^2 \gamma^2} \omega_0^{TS}$$

$$\left(\frac{dN_{ph}}{dt}\right)_{ChR} \propto \gamma^{1/2}$$

 $\left(\frac{dN_{ph}}{dt}\right)_{ChR} \propto \gamma^{1/2}$ - number of photons per unit of time -

$$\left(\frac{dN_{ph}}{dt}\right)_{TS} \propto Const$$

$$P \propto \gamma^2$$

- radiation power -

$$P \propto \gamma^2$$

@ comparison factor:

$$f \simeq \frac{\mathbf{A}_{Ch}^2}{\mathbf{A}_{TS}^2} \frac{L_{Ch}}{L_{TS}}$$

 $f\simeqrac{\mathbf{A}_{Ch}^{2}}{\mathbf{A}_{TS}^{2}}rac{L_{Ch}}{L_{TS}}$ $L_{Ch}\left(z
ight)\simeq\int_{0}^{z}\ N_{ch}\left(z
ight)dz$

Laser beam size & mutual orientation

@ strength parameters – crystal & field:

 \mathbf{A}_{TS}^2 ~ 700 eV/Å³ for the 10 TW laser with a beam diameter of 0.1 mm

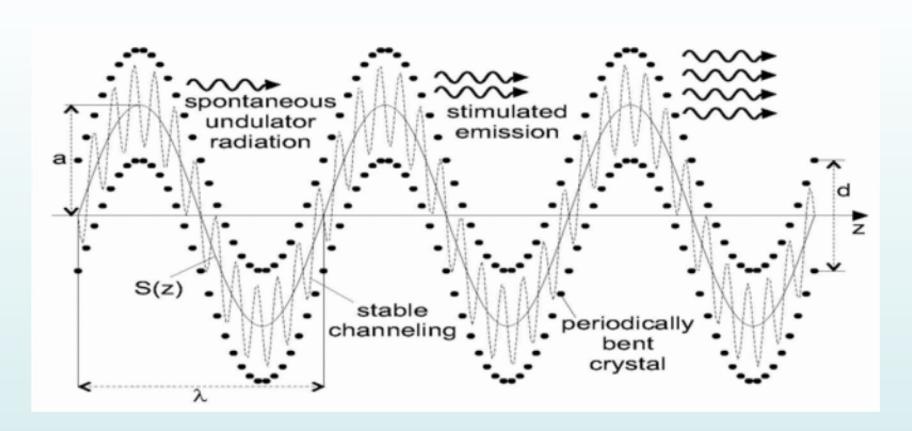
@ Channeling Radiation & Thomson Scattering

For X-ray frequencies: 100 MeV electrons channeled in 105 μ m Si (110) emit ~ 10⁻³ ph/e-corresponding to a Photon Flux ~ 10⁸ ph/sec

ChR – effective source of photons in very wide frequency range:

- in x-ray range higher than B, CB, and TS
- however, TS provides a higher degree of monochromatization and TS is not undergone incoherent background, which always takes place at ChR

@ Bent Crystal as Microundulator & Lasing Effect

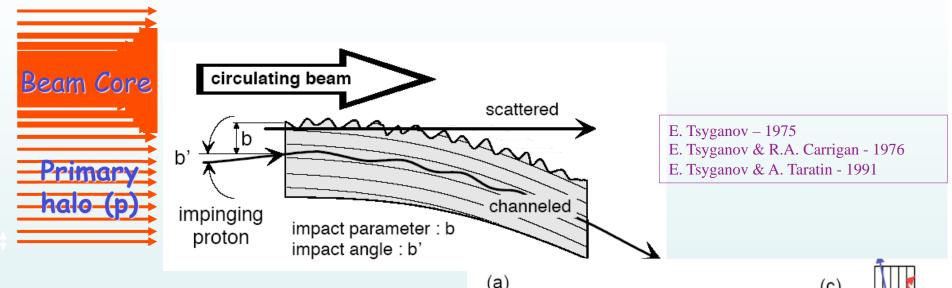


The important and novel idea consists in **realizing a microundulator**, that is a suitable periodically bent crystal, that allows to achieve much higher energy range for the emitted photons respect to the conventional free electron laser (up to 100 keV). A crystalline undulator can be realized either dynamically, using high amplitude transverse acoustic wave, or statically using graded composition of strained layers. Both methods are easily applicable by means of modern technology. The theoretical results establish the feasibility of a crystalline undulator, but in order to make the project to advance it is mandatory to carry out experiments to test the idea and to characterize the emitted radiation as a function of several parameters.

@ BTF Layout

 BTF as unique European facility to deliver positron beams in the range of the energy required for CUP strong photon peak with energy from 20 keV up to 1,5 MeV should appear according well accepted DHSTB02 channeling and undulator theories for 400-600 MeV WCMpositrons SLTB04 Crysral SLTB03 undulator High background/noise DHST 01 DHSTB101 to accumulator 18 81 11 18 8 1 1008 8 D 1 18 80 from Linac DHSTP001 SLTB02 SLTB01 TARGET DHTB001

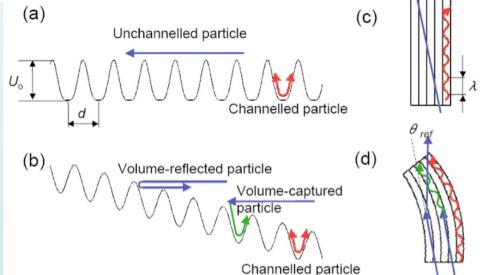
@ Crystal collimation



Possible processes:

- multiple scattering
- channeling
- volume capture
- de-channeling
- **♦** *volume reflection*
- ◆ Primary halo directly extracted!
- ♦ Much less secondary and tertiary halos!?

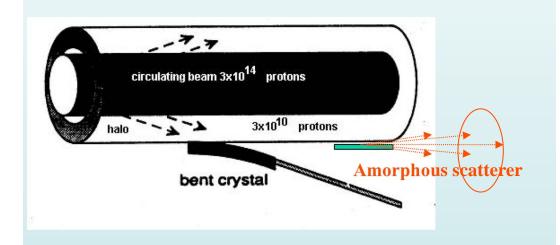
..but no enough data available to substantiate the idea...

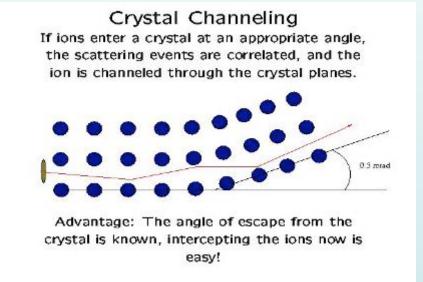


@ UA9: Channeling experiments at SPS

A bent crystal should efficiently deflected halo particles away from the beam core toward a downstream massive absorber

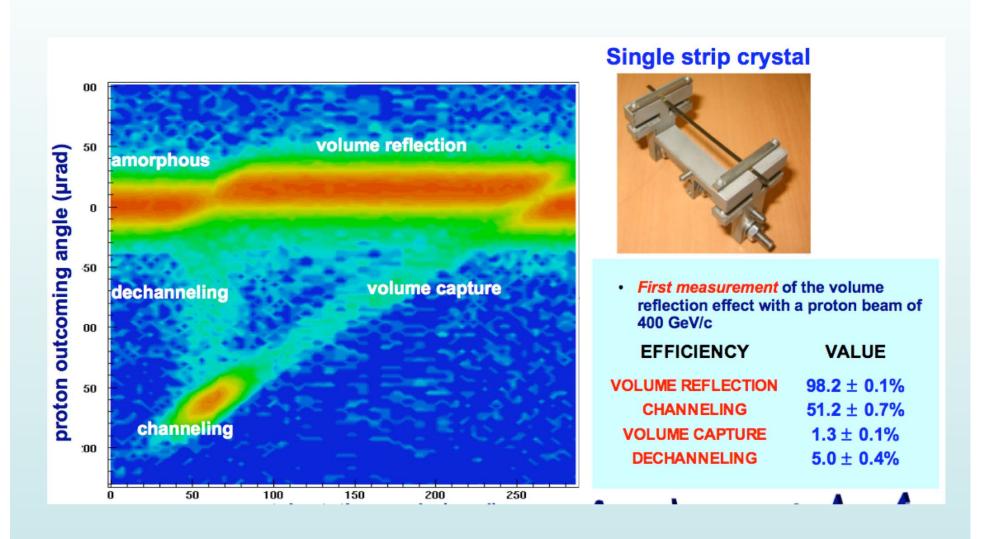
The selective and coherent scattering on atomic planes of an aligned Sicrystal replaces the random scattering process on single atoms of an amorphous target



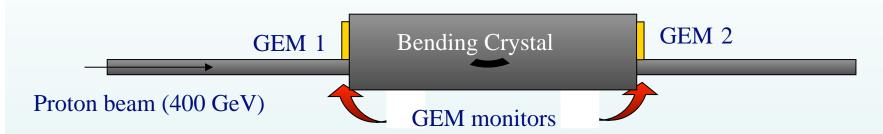


The possible use of crystal for phase II collimation at LHC can be extensively tested at SPS now.

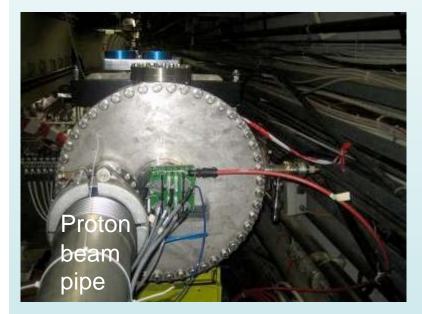
@ First observation: 400 GeV/c CERN, 2006



@ Installation in Crystal Tank at SPS



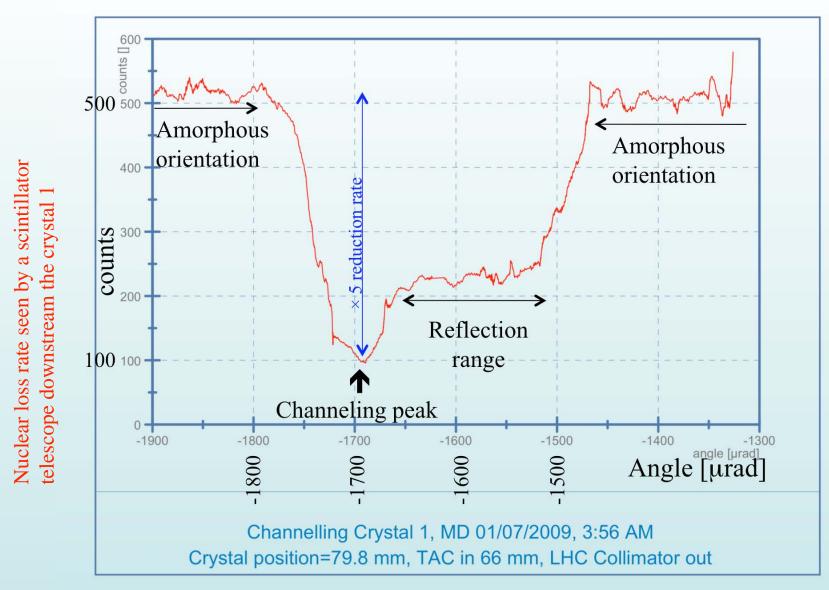
Front view Side view





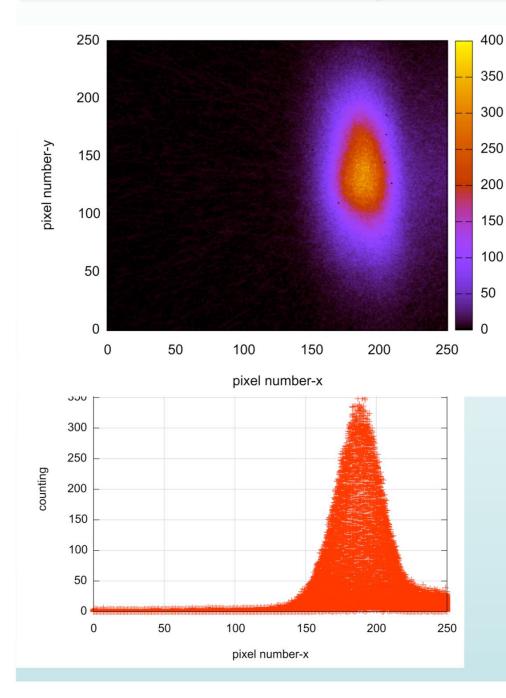
Contribution of Roma1 and Frascati (February 2009)
The GEM monitors will measure the beam halo during the crystal insertion

@ Beam deflection based on channeling process



Nuclear loss rate (including diffractive) strongly depressed

@ Deflected beam profile: Medipix





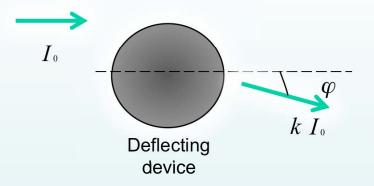
Medipix sensor of the type inserted in the UA9 roman pot, **provided by L. Tlustos (PH/ESE)**

- ◆ 256×256 square pixels
- 1 pixel size = $55 \mu m$
- ◆ 1 frame integration time 1 s

- Pick/valley density ratio = 10
- We observed a ratio of 30 (recording lost for a computer crash)

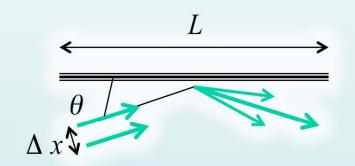
On X-Ray & Neutron Waveguiding: ... from surface down to bulk channeling...

@ X-ray Optics Problem



$$k = k (\varphi) << 1$$

$$k_{\text{max}}(\varphi) \propto 0.1 \div 0.2$$



$$\theta \ll 1$$
 $(\theta \leq \theta_c \equiv \omega_p/\omega)$

$$\begin{array}{ccc} \omega \approx 10 \ keV \\ \Delta x \approx 1 \ mm \end{array} \} \quad \Rightarrow \quad \varphi \approx 1^{\circ} \quad \Rightarrow \quad L \geq \frac{\Delta x \ \varphi}{\theta_{c}^{2}} \approx 2 \ m$$

The system sizes may be significantly diminished (in about two orders of magnitude) by the use of capillary systems.

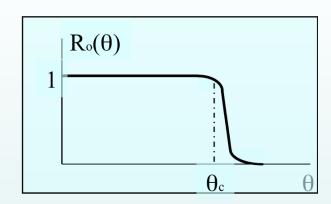
@ Basic principles

$$\Delta \vec{E} + \varepsilon k^2 \vec{E} = 0$$

$$\varepsilon = \begin{cases} 1, z \le z_0; \\ 1 - \theta_c^2, z > z_0 \end{cases}$$

$$\theta_c = \omega_p / \omega \approx 30 eV / \omega$$
, $\omega \rightarrow$ photon energy

$$\theta \le \theta_c \Rightarrow \text{total external reflection (TER)}$$

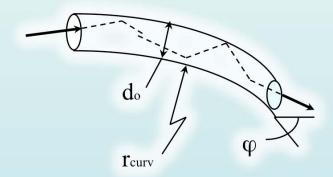




$$\Delta\theta \approx 2\theta_c$$

free space propagation

$$I \propto \int_{\theta_{\min}}^{\theta_c} R(x) \sin(x) dx \propto L^{-1}$$



$$\theta_i \le \theta_c$$
 - multiple TER

$$\frac{r_{curv}\theta_c^2}{2d_0} \ge 1 - \text{effective guide}$$

$$I(\varphi) \propto \int [R_0(x)]^{V(x,\varphi)} dx$$

$$I(\varphi) \propto \int [R_0(x)]^{V(x,\varphi)} dx$$

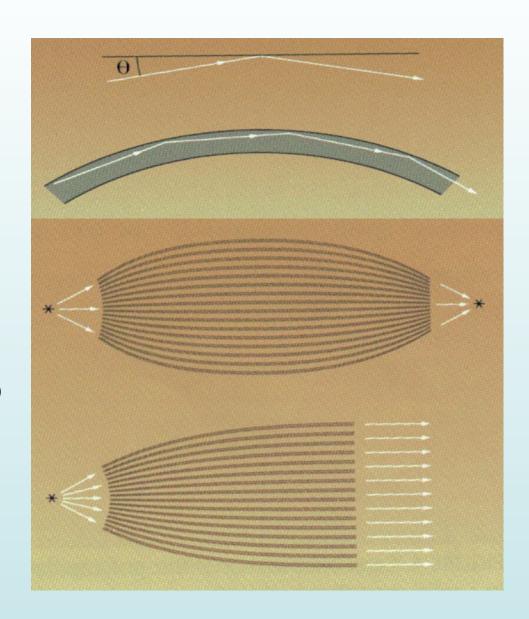
@ X-ray and neutron capillary optics

Basic idea of polycapillary optics is very close to the phenomenon of charged particle channeling

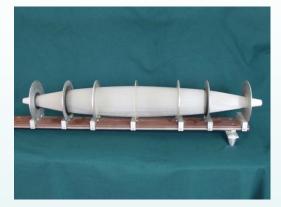
- @ beam bending through large angles
- @ divergent beam to convergent one
- @ divergent to quasiparallel & vv

Number of applications

- @ scientific instrumentation (XRF, XRD)
- @ elemental/structural analysis
- @ medicine (diagnostics, therapy)
- @ astrophysics



@ Samples of capillary optics



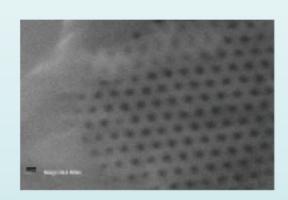
1st generation: [m]



2^d generation: [cm]

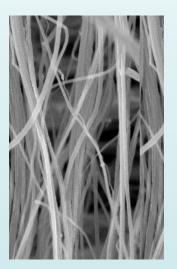


3^d & 4th generations: [mm]



5th generation: [µm]





?n-capillaries?

@ polyCO Evolution: "from micro- down to nano"

Ge	neration	Kind of optics	Sizes: len	gth & channe	l & energy
•	1 st	Assembled lens made of single capillaries	1 m	& 1 mm	& $\leq 10 \ keV$
•	2 nd	Monolithic lens made of single capillaries	10-30 cm	& 0.1-1 mm	& $\leq 10 \ keV$
•	3 rd	Assembled lens made of polycapillaries	10 cm	& 10-50 μm	& $\leq 20 \ keV$
•	4 th	Monolithic lens made of polycapillaries	4-10 cm	& 1-10 μm	& $\leq 50 \ keV$
•	5 th	Monolithic integral micro lens	1-3 cm	& 0.3-1 µm	& $\leq 100 \ keV$



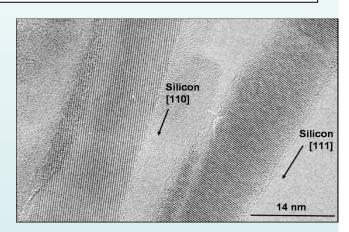
Micro-capillaries



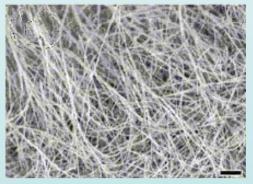
Micro → X-rays



Nano $\rightarrow \gamma$ -rays (?)

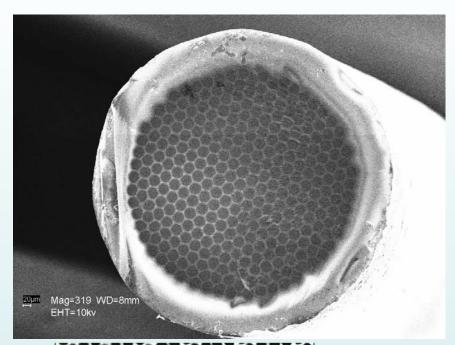


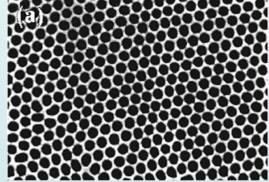
Nanotubes & Nanochannels

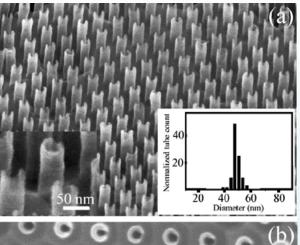


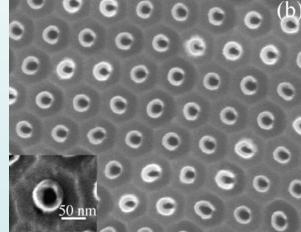
Frascati @ 3 Dec 2009 @ S. Dabagov

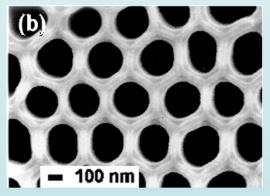
@ Samples of Nanostructures



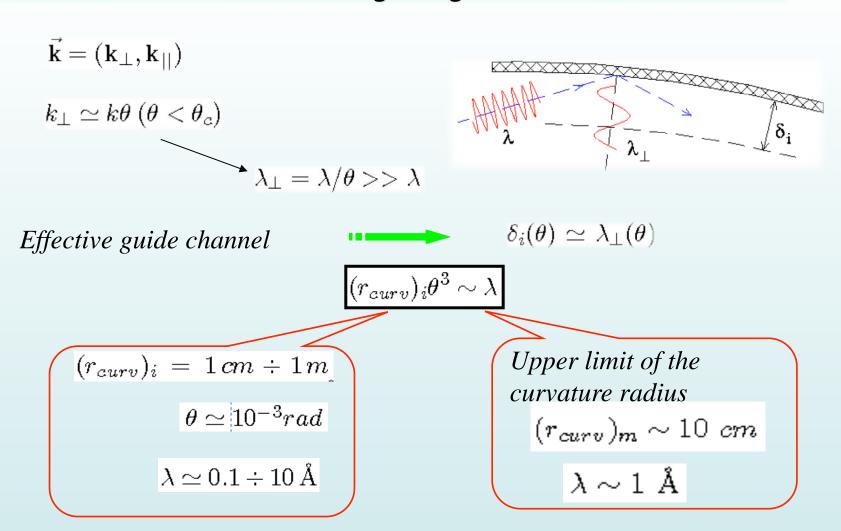






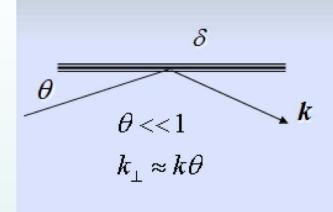


@ Modes of channeling along curved surfaces



Very important expression for the diffraction limit estimation – can be applied for any kind of the optics:

@ Quantum basics



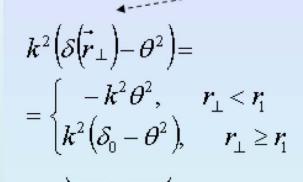
1st order:
$$\Delta \varepsilon(\vec{r}) = 0$$
 -no roughness

Wave equation:

 $k^2\delta$

$$\left(-\nabla^2 + k^2 \delta(\vec{r}_\perp) - k_\perp^2\right) E(\vec{r}_\perp) = 0$$

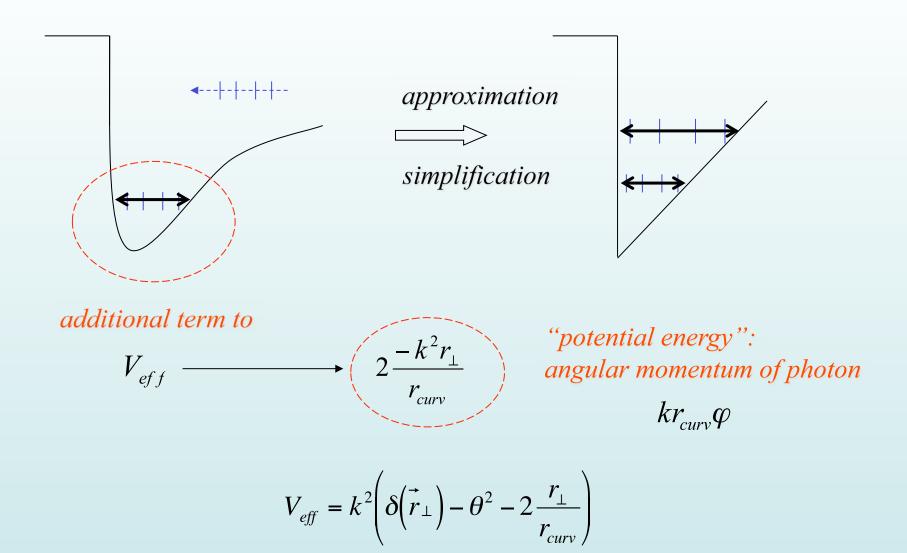
$$V_{eff}$$



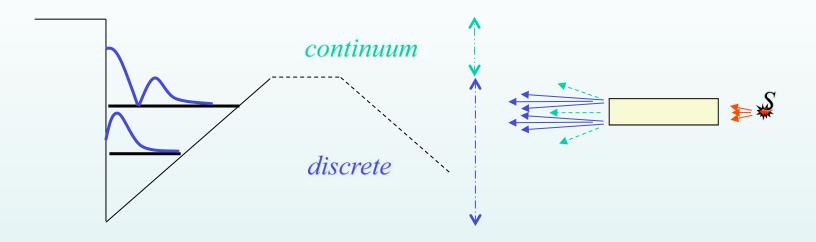
Total external reflection

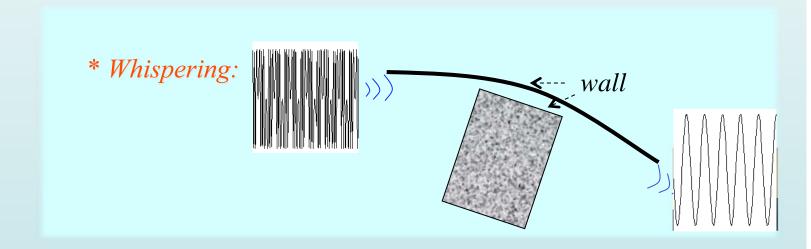
$$V_{\it eff}\equiv 0 \Longrightarrow \theta_c \equiv \theta pprox \sqrt{\delta_0}$$

@ Quantum basics... Curvature



@ Surface channeling - "whispering X gallery"





@ Divergence - surface channeling

APPLIED PHYSICS LETTERS VOLUME 78, NUMBER 19 7 MAY 2001

Divergence behavior due to surface channeling in capillary optics

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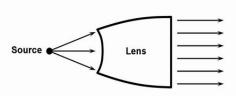
C. Gramaccioni and A. Pifferi

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(Received 21 November 2000; accepted for publication 7 March 2001)

Recent studies on the transmitting and focusing properties of capillary optical systems have shown that several unexpected effects take place during the experiments. One such effect is a dec the beam divergence behind the capillary structures. In this letter, we present results c scattering at grazing angles inside capillaries. During x-ray propagation in capillary channe is a strong angular redistribution of the beam, which has been explained in the framewor wave scattering theory. © 2001 American Institute of Physics. [DOI: 10.1063/1.137011

Today, based on a 10W x-ray tube, developed in-house and combined with optics, we have created an x-ray source producing quasi-parallel monochromatic flux of 10¹⁰ ~ 10¹¹ photons/sec∙mm², which we have used in diffraction studies.

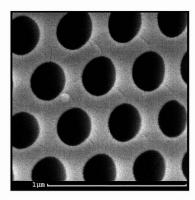


Parallel Beam. Divergence is (2-3)•10⁻³ radian; Parallel nature in both directions Such flux compares well with that usually obtained at synchrotron stations (see Comparison Table 1).

Hence, this system has been named a Laboratory "Synchrotron"

At the same time, in contrast to works carried out at workstations in synchrotron centers where the researchers are strictly limited in time and have other restrains, in our case we provide for a freedom of creativity and possibility of careful and long term investigations.

It should be also taken into consideration that our Laboratory Synchrotron has a reasonable price and is affordable to a great number of scientists, engineers, research laboratories, universities, companies and so on.

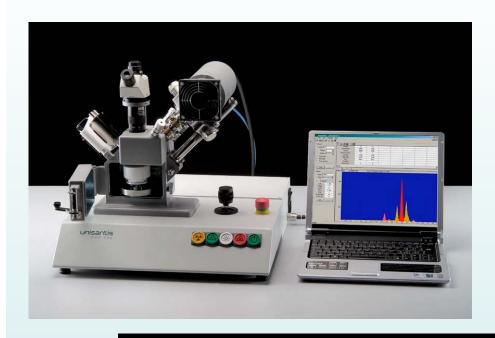


Our Laboratory "Synchrotron"

is a portable instrument and does not need special premises, or radiation protection measures.

In addition, it can be easily incorporated in existent diffractometers, which is another advantage of our system.

@ Capillary optics: applications/new portable instruments





@ Table-top X-Ray Source



"Laboratory Synchrotron"



W		X12-C NLSL	FOP (ESRF)	ANKA	UNISANTIS
	Energy Range	8-13 keV	7-21 keV	4-20 keV	8 keV
	Beam Size	0.3H x 0.5V mm	0.3 x 0.3 mm		20 to 300 micron
	Beam Divergence	2 mrad			2-4 mrad
	Flux	~5 x 10 ¹⁰ photon / sec (0.3x0.5mm ²)	~5 x 10 ¹⁰ photon / sec (0.3x0.3 mm ²)	~1 x 10 ¹² photon / sec (mm ²)	10 ¹⁰ - 10 ¹¹ monochromatic photon / sec (mm ²)
	Source Power	Synchrotron (2.6 GeV, 300 mA)	Synchrotron (2.6 GeV, 300 mA)	Synchrotron (2.6 GeV, 300 mA)	(3÷30) W x-ray tube

@ Comparison of Various X-ray Micro-Focus Sources



Osmic Bruker Unisantis

Maximum power - 50 W Maximum anode voltage - 50 kV Material of the anode - Cu, Mo The size of a tube focal spot \leq 20 μ m

Source Power, W	30	24	(3÷30)
Optics	Multi-layer	Focus Mirror	PolyCO
Flux, Cu Kα ph/s/mm ²	$\approx 1.9 \times 10^9$	109	$10^{10} \div 10^{12}$

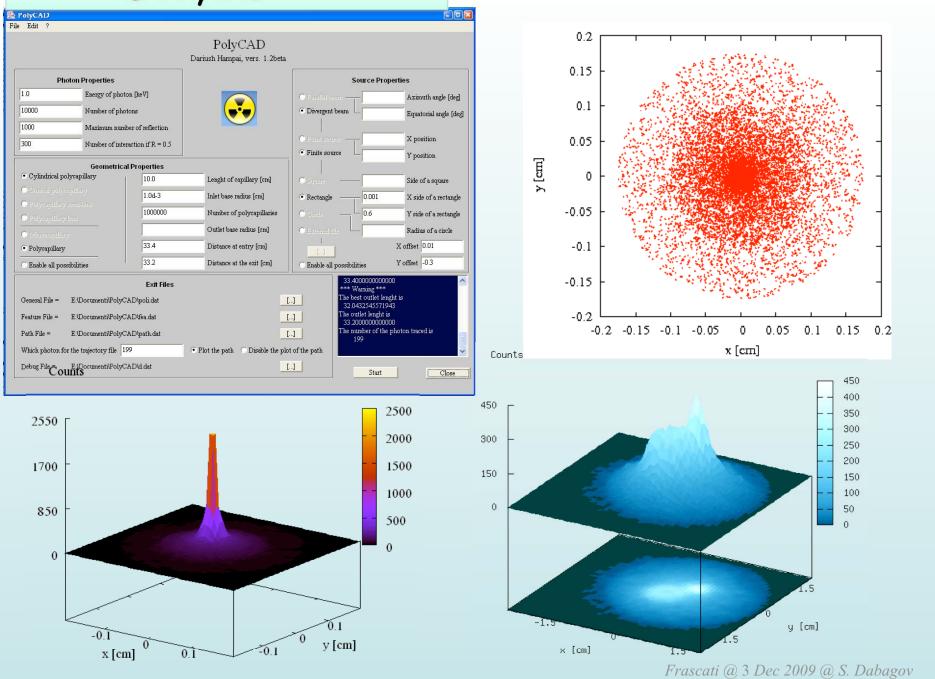
Beams parameters: semilens

Diameter of a beam - 0,1-1 mm Radiation flux - 10¹⁰ ph/s/mm² Divergence of a beam - 4 mrad Cu; 2 mrad Mo

Beams parameters: full lens

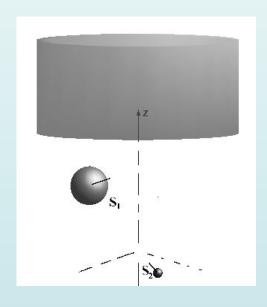
Diameter of a focal spot - 10-50 μm Radiation flux - 10¹² ph/s/mm² Focus length of a lens - 2-10 mm.

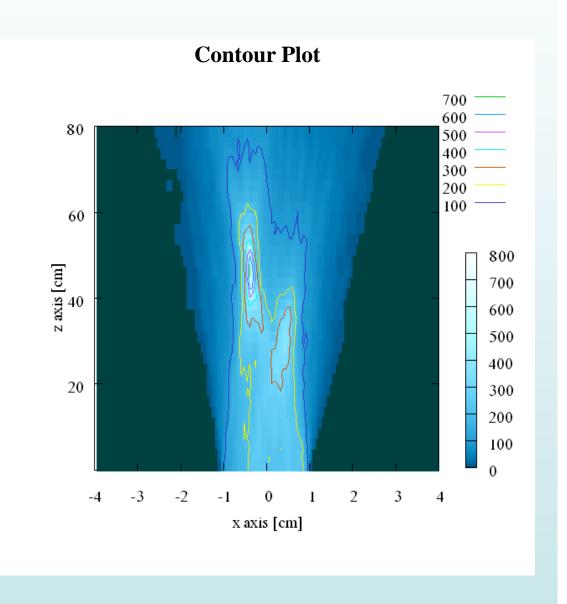
@ PolyCAD



@ 3D Source - PolyCAD

- Reconstruction of the 3D intensity distributions beyond a cylindrical polycapillary due to a couple of Xray spherical source
- The main peaks are located at the conjugate points respect to each sphere position





@ Rendering - PolyCAD

_ B × PolyCAD – Monte-Carlo simulation code for any kind of capillary optical systems angolo_nocubi

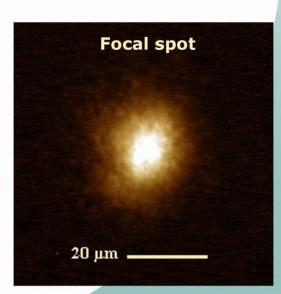
Frascati @ 3 Dec 2009 @ S. Dabagov

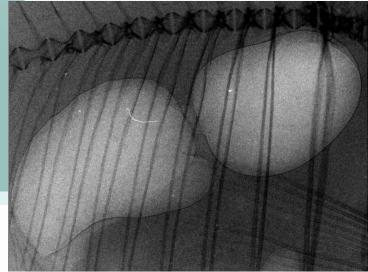
@ X-Lab: Collaboration LNF - UNISANTIS

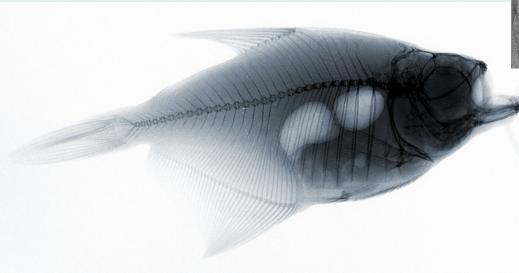
There is a capability of compressing the electron flux to (1-3) micron at less power.

In case of reflection anode, the power of (40-50) watt can be focused to a (10-15) micron focal spot.

In this instance, anode would be at a 1 mm distance from the window.







@ X-ray μ -scope: 1st option

X-ray Cu Source
Power=50 Watt
Intensity=1mA
Voltage=50kV

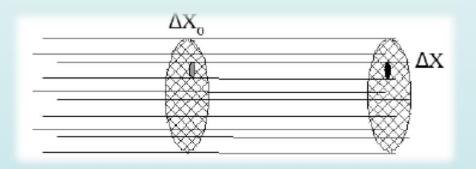
Quasi-Parallel Beam

CCD Detector
res. (3.5x3.5)µm

Sample

Image on CCD

Polycapillary Semi-Lens

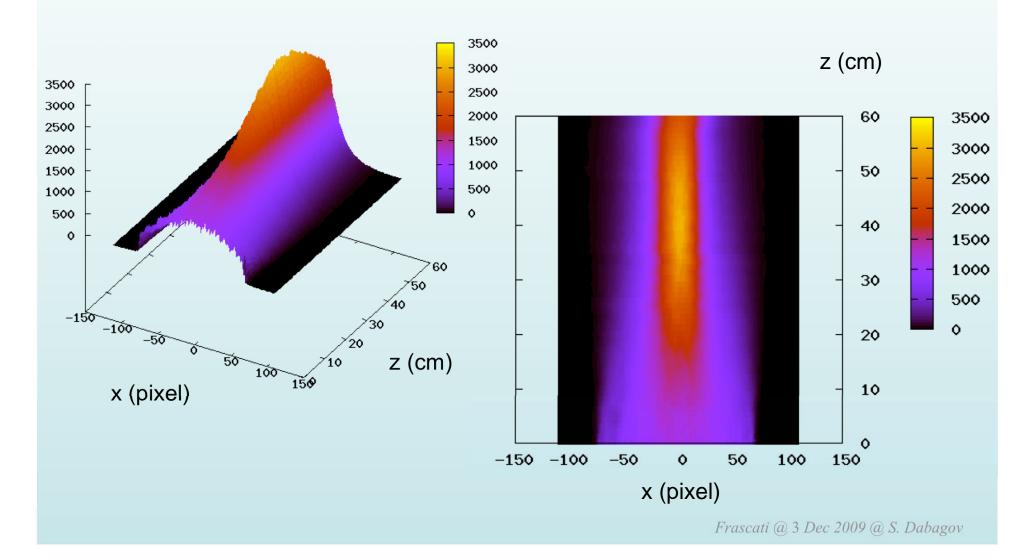


$$|\Delta X - \Delta X_0| \propto l \cdot \Delta \theta$$

1- "sample-detector" distance; $\Delta\theta$ - residual divergence behind the optics

@ Polycapillary Semilens

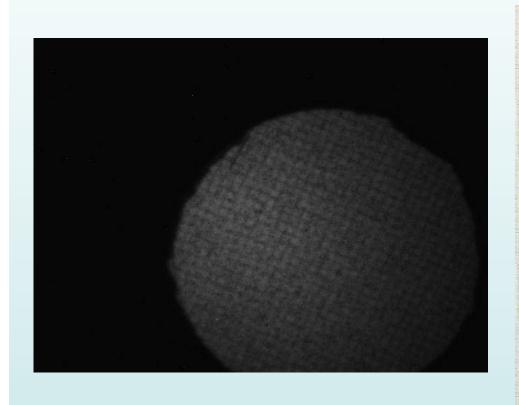
- Input focal distance (mm) = 52.7
- Residual Divergence (mrad) ~ 1.
- Transmission $\sim 60\%$

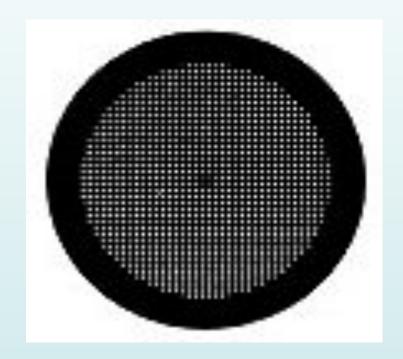


@ High Resolution Image: Cu 400 mesh

- Hole Width 37μm
- Bar Width 25μm

polycapillary pillar

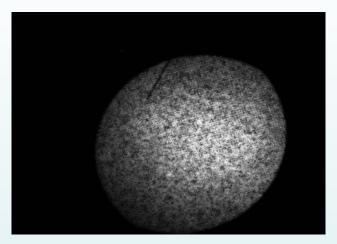




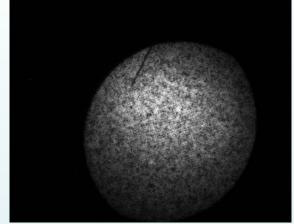
23 cm Polycapillary – CCD Detector

@ High Resolution Image: Au 1000 mesh

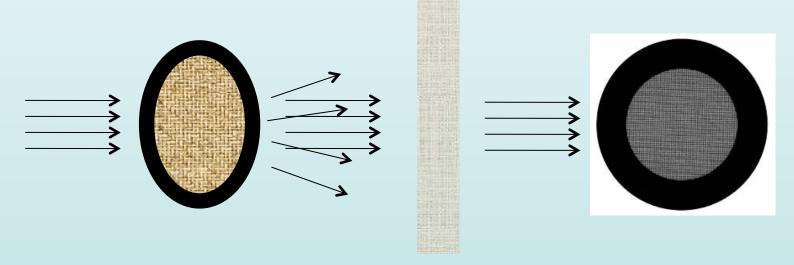
- Hole Width 19μm
- Bar Width 6μm



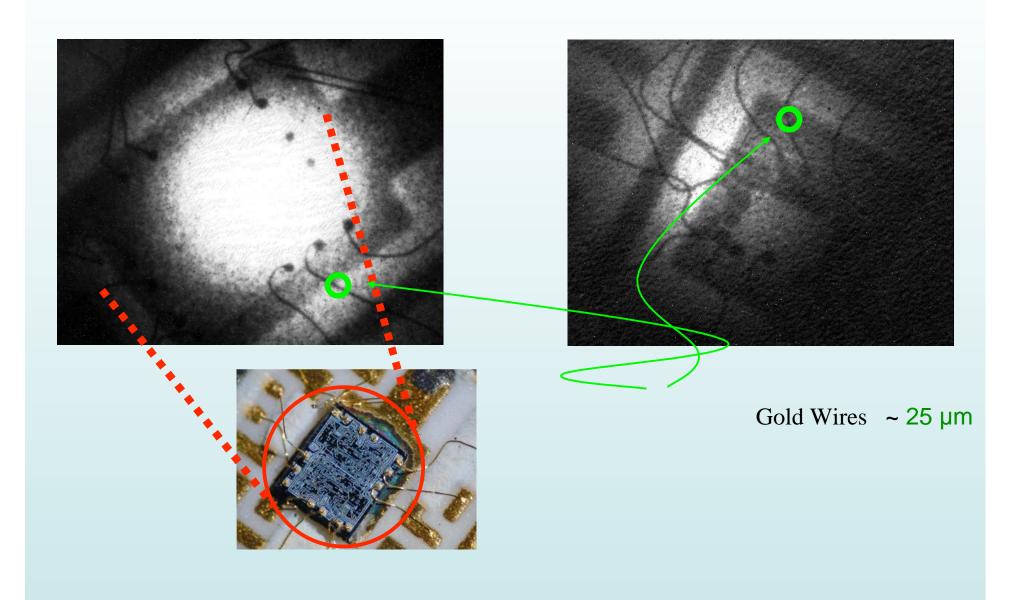
23 cm optics – CCD



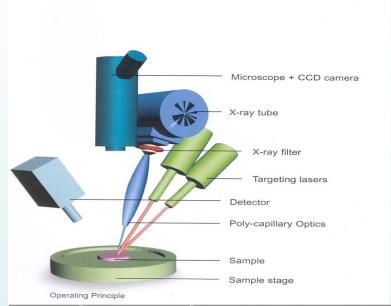
44 cm optics – CCD



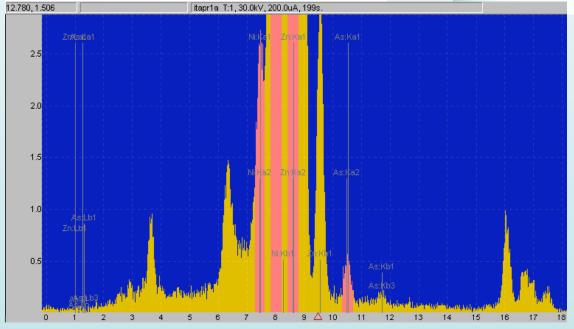
@ High Resolution Image: Chips



@ Collaboration: LNF - UNISANTIS

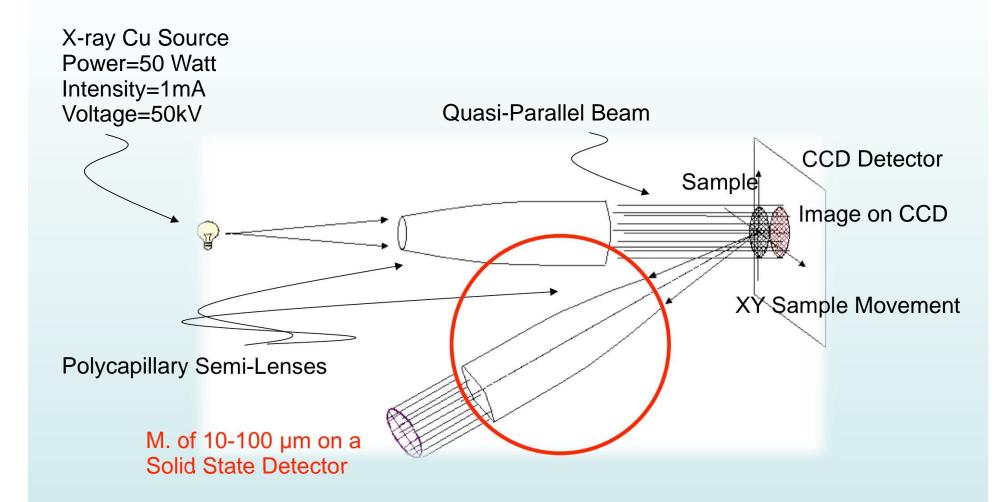




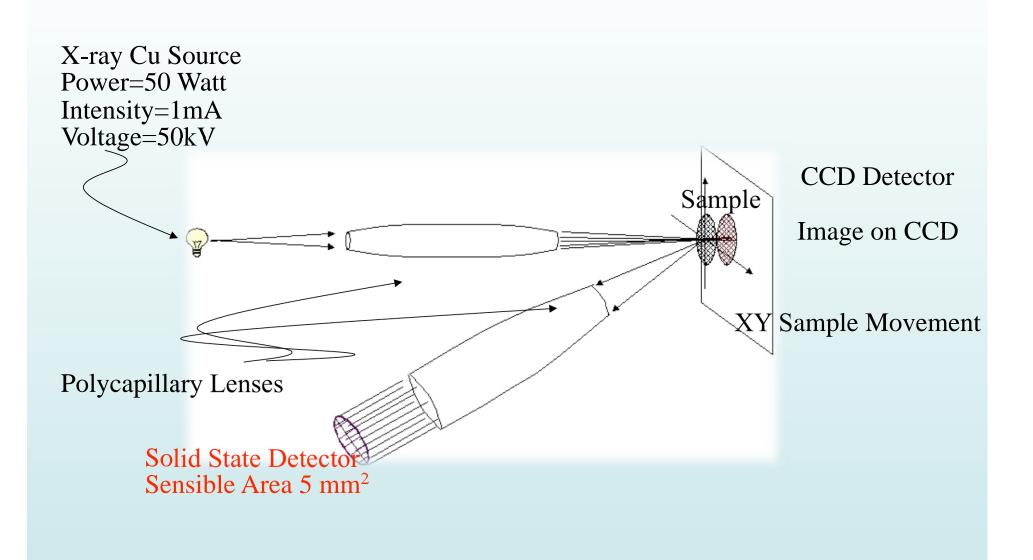


Alloy from Sample copper ancient archeological dig and spectrum of this alloy obtained using this micro analyzer (13 Century, Italy).

@ μ XRF mapping (i)

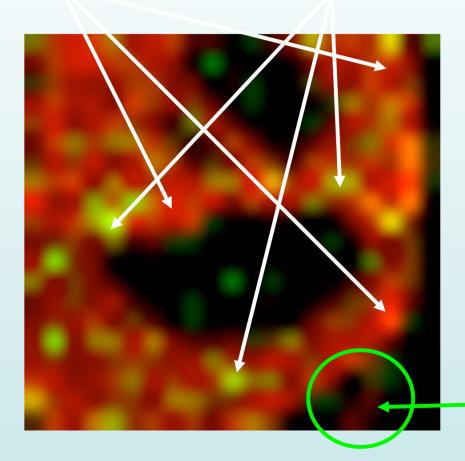


@ μXRF mapping (ii)

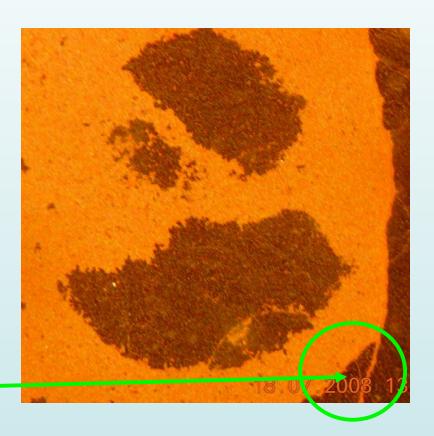


@ µXRF mapping on Fe₂O₃

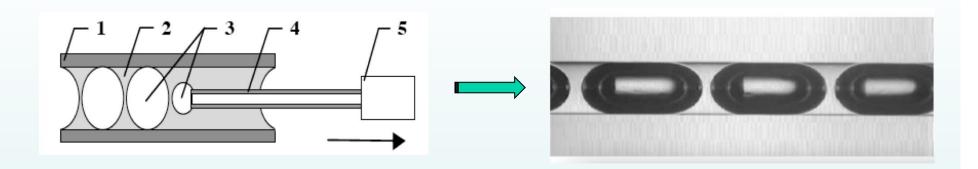
Fe Mn



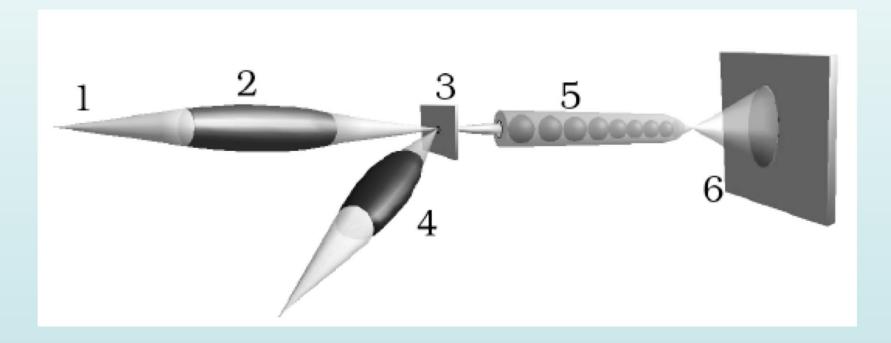
µXRF Mapping – 200x200 μm² spot 4x4 mm² measured area (40x40 steps) 2 minutes for each spot CuKα Tube; V=24 KV, I=0.750 mA



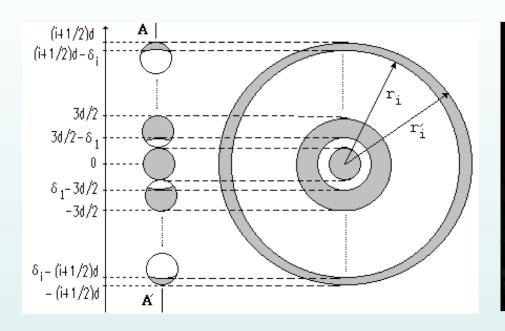
@ polyCO + Spherical CRL

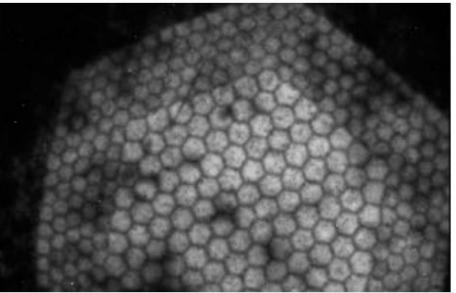


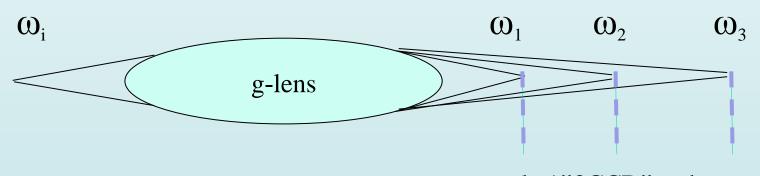
 μ XRF + μ X-imaging = polyCO confocal + spherical ("buble") compound refractive lens



@ g-lens \rightarrow graded lens

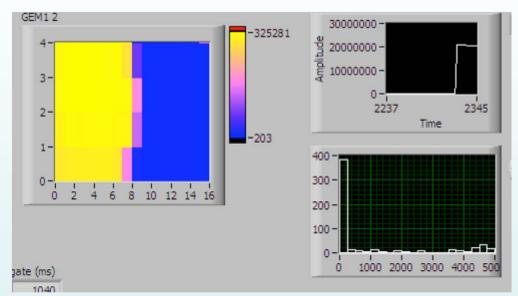


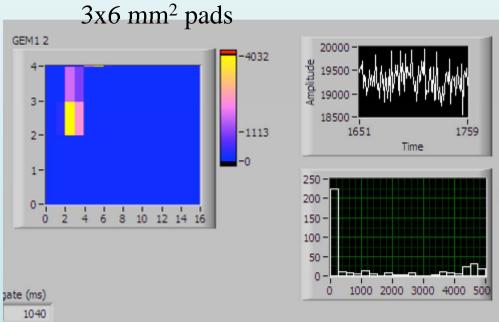


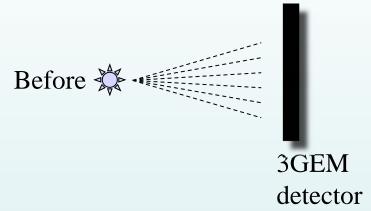


sample / "3CCD" x-detector

@ GEM X-Ray Spot with Polycapillary: GEMINI

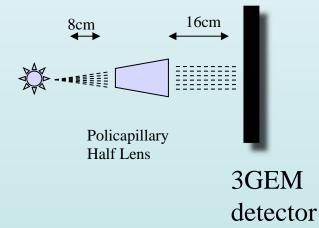




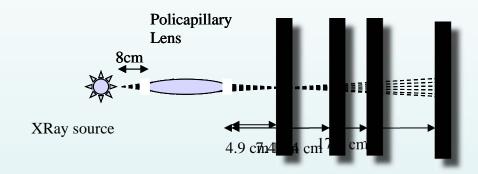


First measurements with polycapillary and GEM detector

After



@ Measurements with 3GEM



GEM12 11372.00 15000 -10000 mean 5000 -11530.2 3-1160 1268 sigma 2-319.3 Time sample lengt^{gth}ngth 1000 -10 800 -0 2 4 6 8 10 12 14 16 600 -400 -200 -2500 te (ms)

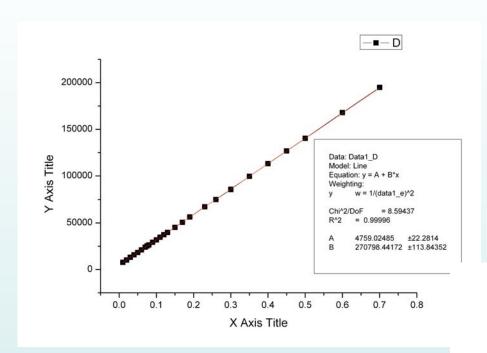
3GEM

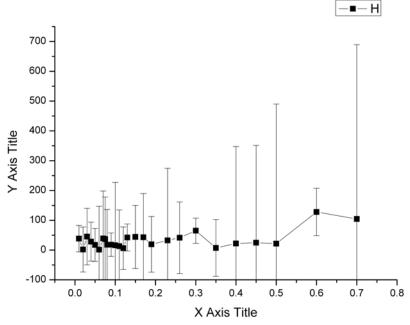
detector

X Ray Energy 10.6 kV Current 0.01 mA Power 0.1 Watt 3GEM Plateau @ 980 Volt (AR CO₂ CF₄) GEM spot counting rate 380 Hz Efficiency X Ray detection 7 %



@ Linearity: GEM detector

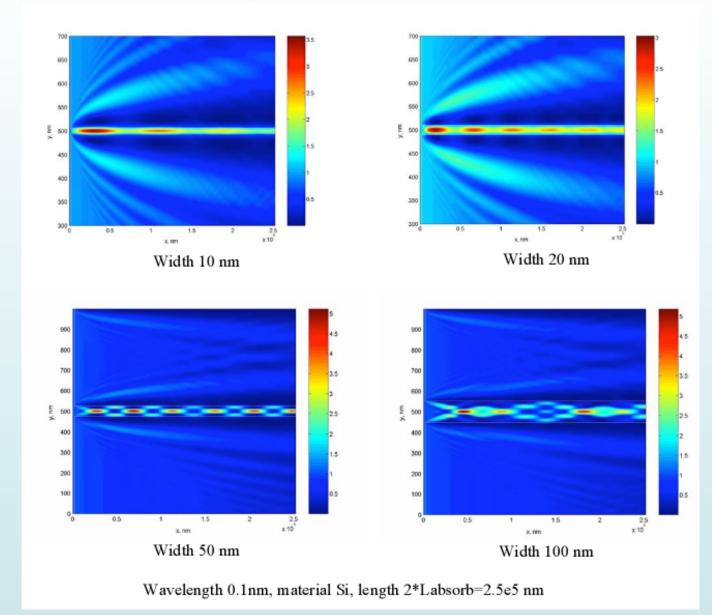




@ Down to bulk X-ray channeling

* Channeling: charged particles \oplus crystals Example: e^- captured by the string potential (smeared atomic) averaged potential $\varphi <<1$ $(\varphi < \varphi_L \sim \sqrt{U/E}$ - Lindhard angle)

@ Nanocapillaries



$$S = \frac{2\pi\theta_c a}{\lambda} \equiv \frac{2\pi a}{\lambda_{\perp c}}$$

- number of modes

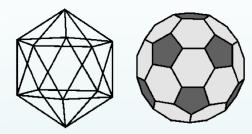
$$\lambda_{\perp c} \cong \lambda/\theta_c$$

~ 40 nm for glass

Tunneling length

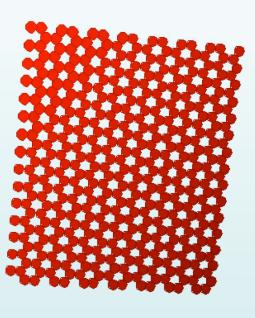
$$\propto \frac{c}{\omega_p} \approx \frac{137e^2}{\omega_p} \approx 8 \ nm$$

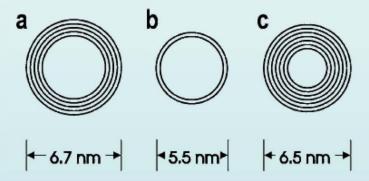
@ Nanotube simulations



Base: fullerene molecule C_{60} sphere of $d \sim 0.7 \text{ nm}$





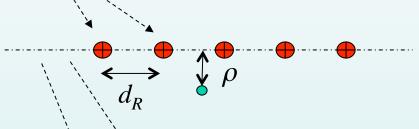


Roled graphite sheets: nested nanotubes

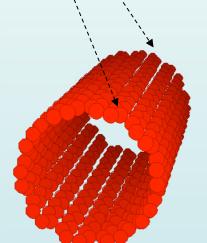
@ Potentials: Doyle-Turner approximation



$$f(\mathbf{k}) = 4\pi Ze \sum_{j=1}^{N} a_j \exp(-k^2/4b_j^2)$$
 - form-factor for the separate fullerene



$$V_R(\rho) = (4Ze^2/d_R) \sum_{j=1}^{N} a_j b_j^2 \exp(-b_j^2 \rho^2)$$



$$U(\mathbf{r}) = \sum_{i} V_{R}(||\mathbf{r} - \mathbf{r}_{i}||)$$
 continuum potential as sum of row potenti

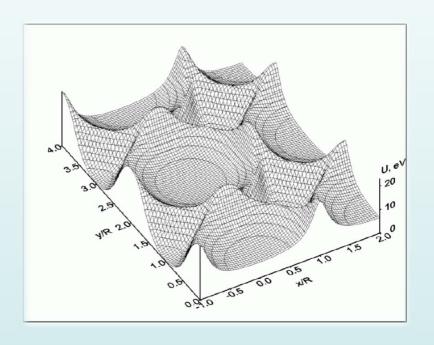
as sum of row potentials

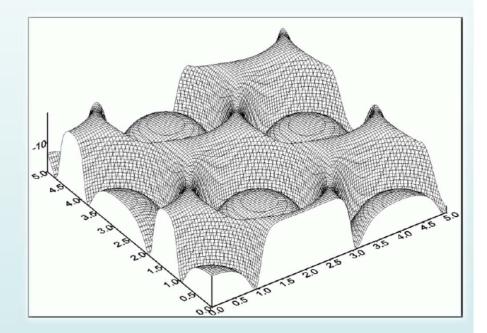
$$U(r) = (16\pi dZe^2/3\sqrt{3}l^2)\sum_{j=1}^{N} a_j b_j^2 \exp\{-b_j^2[r^2 + (d/2)^2]\}I_0(b_j^2rd)$$

r – distance from the tube $I_0(x)$ – mod. Bessel function

@ Potentials: Doyle-Turner approximation

Continuum potential in C60 fullerite: [100] and [110] after averaging of the wall potential





@ Potential for neutral particles: Moliere approximation

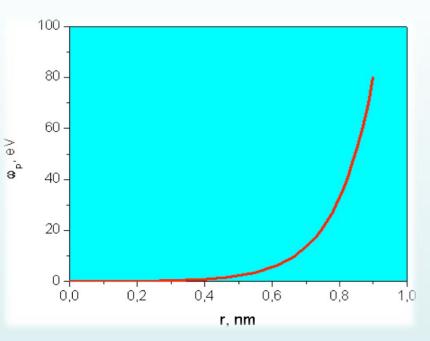
$$N_e(r) = \frac{Z}{4\pi a^2 r} \sum_{i=1}^{3} \alpha_i \beta_i^2 \exp\left(-\frac{\beta_i r}{a}\right)$$

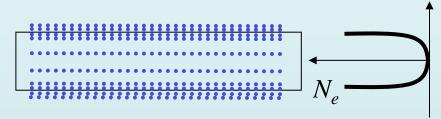
$$C: Z=6$$

$$a \approx 0.05Z^{-1/3}$$
 – screening length

$$\overline{N}_{e}(r) \approx \frac{r_{curv} n_{a} Z}{\pi a^{2}} \sum_{i} \alpha_{i} \beta_{i}^{2} \int_{0}^{\pi} d\theta K_{0} \left(\frac{\beta_{i} \rho}{a}\right)$$

$$\rho = \left(r^{2} + r_{curv}^{2} - 2r r_{curv} \cos \theta\right)^{1/2}$$





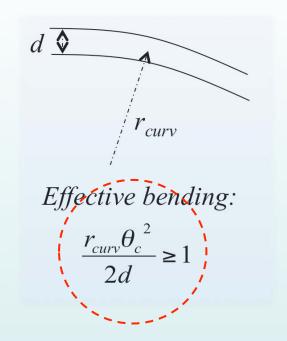
"Continuous filtration by energy"

@ Nanocapillary: Bending efficiency

$$n = \sqrt{1 - \theta_c^2} \approx \sqrt{1 - \frac{\omega_p^2}{\omega^2}}$$

$$\omega_p = \sqrt{\frac{4\pi N_e e^2}{m}} - plasma \quad frequency$$

$$\omega - photon \quad frequency$$



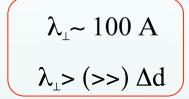
 μ -capillary: 10^{0} - 30^{0} through 10-20cm

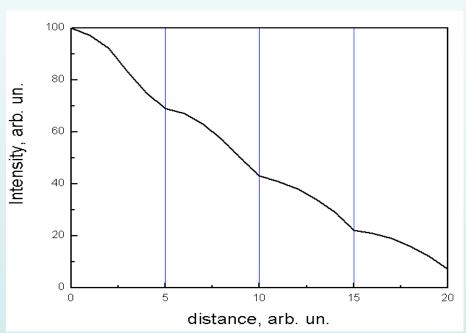
n-capillary: the reduce of the dimensions by several orders with much higher efficiency of the bending

@ Bulk channeling: tunneling & diffraction

The even modes exists at any ratio between the channel size and layer distance







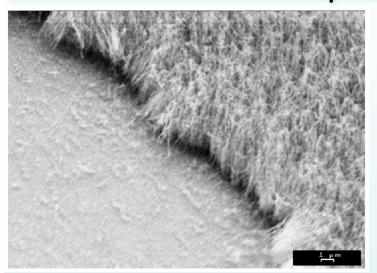
Spatial distribution of radiation through the various narrow channels

Evidence of high leakage through the nanotube wall barriers

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The necessity of the cladding - multiwall nanotube

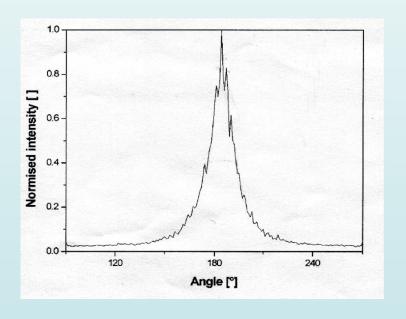
@ Nanotube sample - forest of nanotube bundles

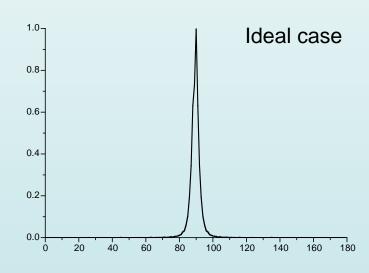


SEM image of aligned carbon nanotube bundles deposited perpendicularly to the substrate surface. <L> \sim 10 μ m.

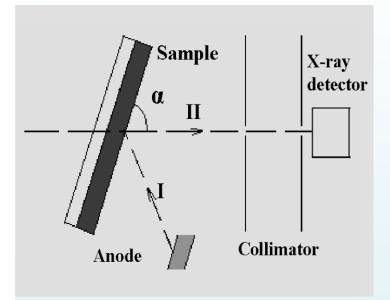
By TEM investigations, multilayered carbon nanotubes, making film structure, have from 5 up to 25 layers, external diameter of 100÷200 Å, and diameter of an internal cavity of 50÷70 Å.

Angular distribution for nanotube forest

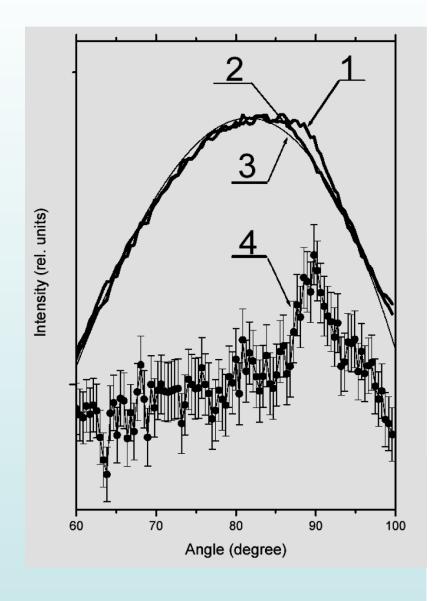




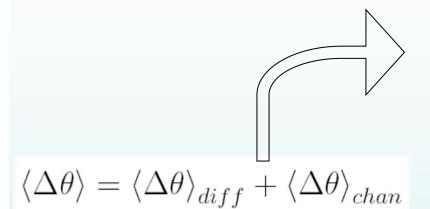
@ X-ray channeling in nanotube



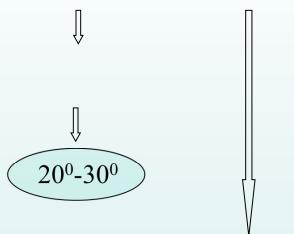
The sample of 8 mm in length and 2 mm in width was fitted to the rotatable axis provided the change of angle α between the sample surface and the optical axis of spectrometer from 60° to 100° . The angle β between the radiation source (anode surface of 6 mm in width) and the rotating axis was 20° .



@ X-ray channeling in nanotube: estimations



Diffuse scattering +Inner cavity

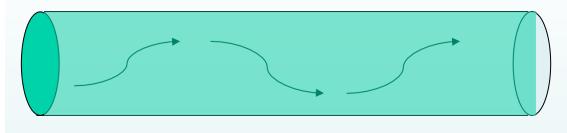


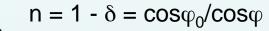
$$\langle \Delta \theta \rangle_{text} + \langle \Delta \theta \rangle_{pass}$$

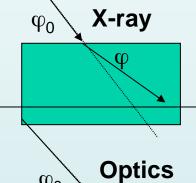
$$\langle \Delta \theta \rangle_{chan} \sim \frac{1}{r_0 N_0} \int_{\omega_{\min}}^{\omega_{\max}} d\omega \, \frac{N(\omega)}{\omega} \int_0^{r_0} dx \, \omega_0(x)$$

$$\langle \Delta \theta \rangle_{chan} \lesssim \omega_{0 \, \text{max}} r_0 / 10 \approx 6^0 \div 10^0 \text{ at the given values of } \omega_{0 \, \text{max}} \simeq 100 \div 150 \text{ eV}$$

$fearsfar{2}{c}$ Features of $fearsfar{2}{c}$ K α propagation in nanotubes



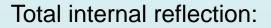




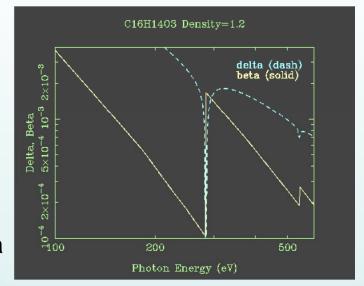
 ϕ_0

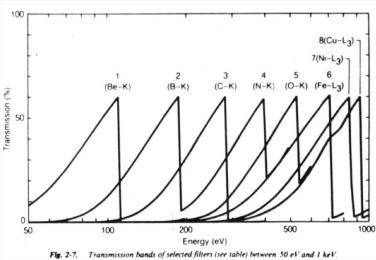
 δ ~10⁻⁶ λ ~ 0.1 nm for C media

for
$$\omega >> \omega_1$$
 $\delta = r_0 N \lambda^2 / 2\pi$

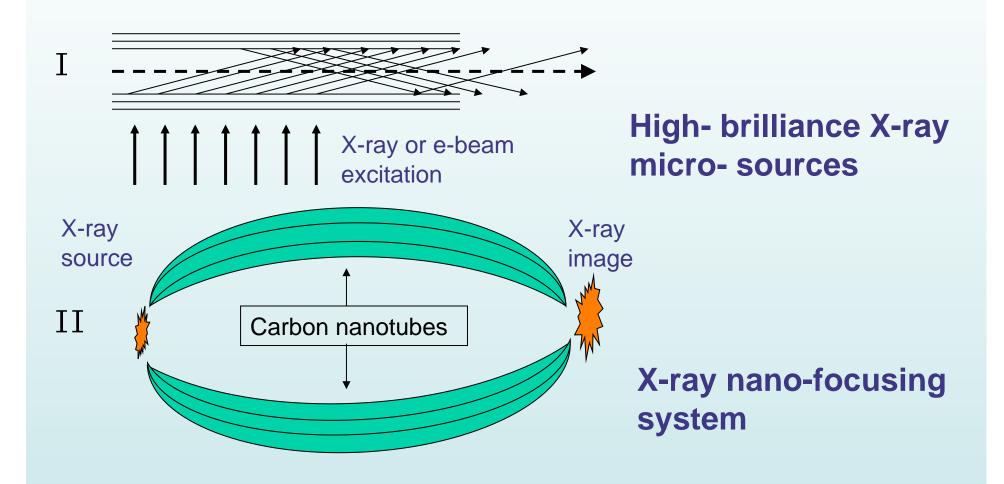


$$\phi_0 \sim \sqrt{2}\delta \sim 2.5^\circ$$
 for $\delta \sim 10^{\text{-}3}$





@ X-ray nanotube channeling: possible applications



@ Capillary Optics as Instrument

@ focusing efficiency:

condensing radiation in a spot of submicron & nanometer diameter with flux gain of 10³-10⁵

- x-ray spectroscopy (XRF analysis min. detect. limit up to 10⁻¹⁵ g)
- x-ray microscopy (source image study nm scale: µ-CHANDRA)
- µ-diffraction (XRD analysis analysis of micro-materials: "table synchrotron")

@ bending efficiency:

deflection through the angle up to 30° over the distance 10-15 cm

- special angle-tuning bender (space problem for new experiments)
- @ new technological ideas:

combination of various types of optics

@ X-ray waveguiding for basic research

surface x-ray and neutron channeling

- @ bound states of x-ray propagation through capillary channels
 - quantum effects in curved space (simple model of gravity)
 - coherent/incoherent x-ray & n scattering study (whispering gallery for x-rays)
 - phase distortion study ("hidden coherence")

bulk channeling of neutral particles & modal propagation

- @ nanotubes as nanocapillaries (?):
 - new instrument to study bulk parameters of nanotubes (bulk potentials etc.)
 - novel x-ray & n optical devices (nanofocusing, nanocollimation etc.)
- @ nanochannels as nanocapillaries (!):
 - X-ray & n "hollow fibers" inverse to optical waveguide

. . .

Thanks to
the international collaborations
in
Channeling and X-ray Optics
Projects,
particularly:

D. Hampai, G. Cappuccio, G. Cibin, F. Murtas, I. Bukreeva, A. Popov ...