



Slow positron extraction from DAΦNE

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POSEYDON

Motivations & assumptions

Extracted "long" positron beam motivated by:

- Dark sector searches in positron annihilations (PADME-type experiments)
- Crystal channeling experiments with positrons (parametric radiation, crystal undulators, etc.)
- Test positron beam

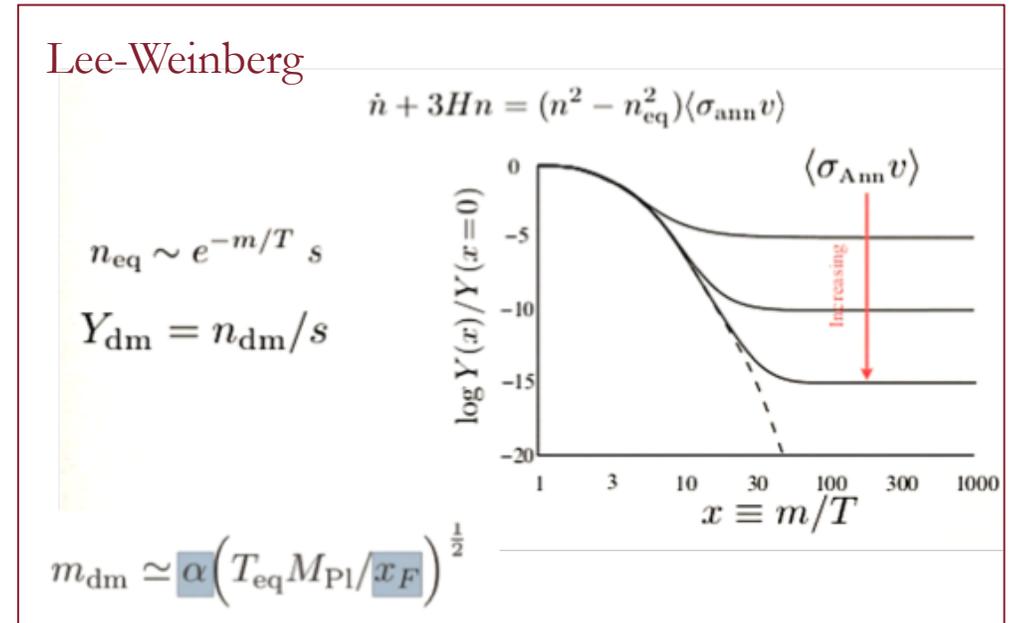
Assume:

- Minimal modifications to the infrastructure
- Compatibility with synchrotron light facility

Why searching for a dark sector

Properties of dark matter:

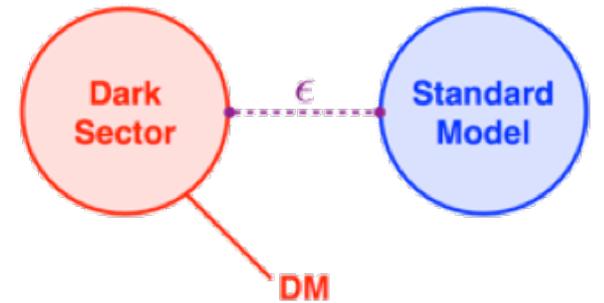
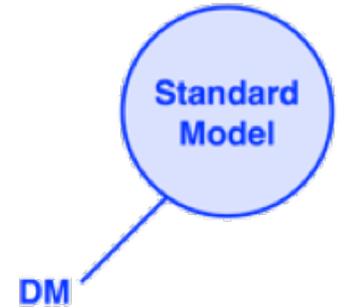
- Dominant with respect to baryonic matter: $\rho_{DM} \sim 5 \rho_b$
- **Suppressed** interactions with QCD and QED
- **Suppressed self-interactions** $\frac{\sigma}{m_{DM}} < 1 \text{ barn/GeV}$
- **At least one stable DM particle**
- **May or may not** been in thermal equilibrium with SM particles.
- If part of the same bath: freeze-out gives relic density
- **Freeze-out:** $m_{DM} \sim \text{TeV}$ if **weakly interacting**



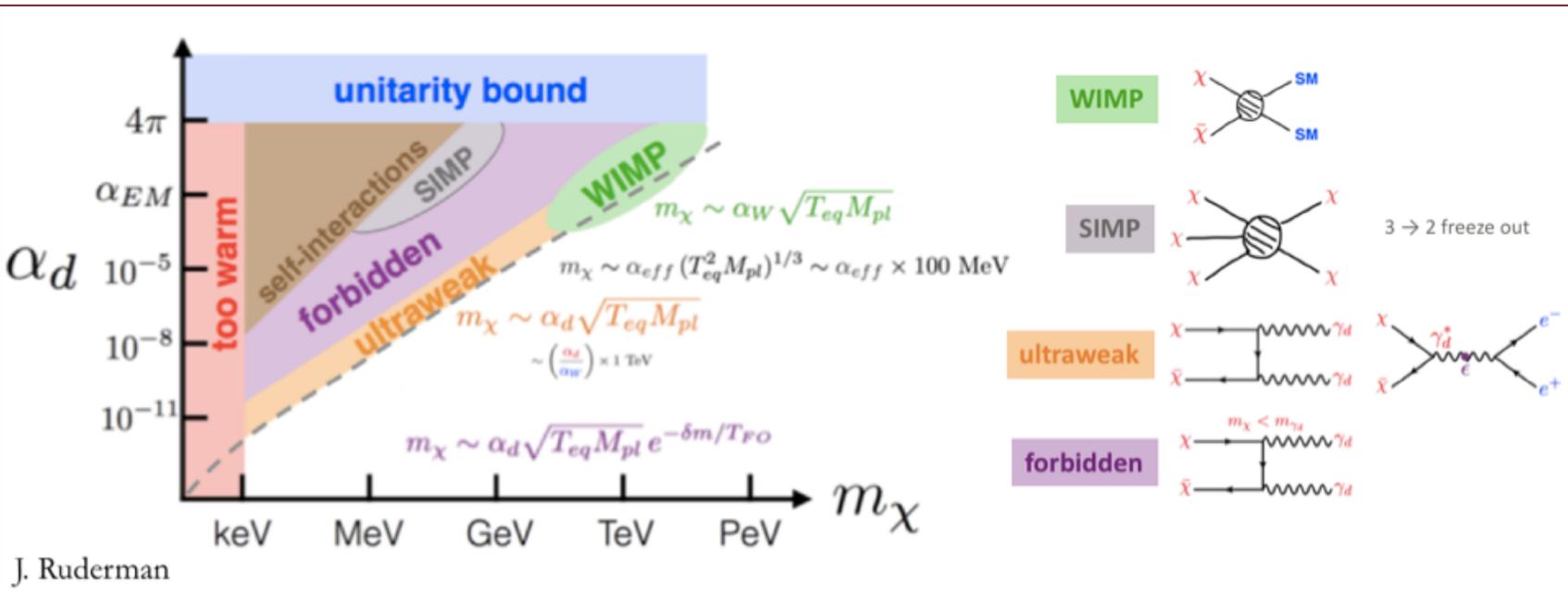
Why searching for a dark sector

Solutions to no WIMP at ~ 1 TeV

- Non-thermal or quasi-thermal models
 - Freeze-in, axions, asymmetric DM, gravitino...
- Dark freeze-out of an **almost secluded** DM sector
 - DM does **not annihilate** into SM particles (at least not directly)
 - The two thermal baths (SM and DM) can have different temperatures



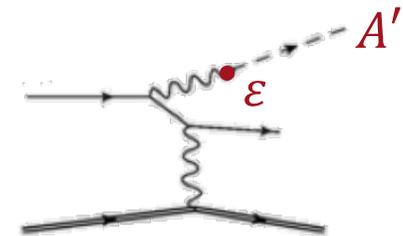
- At least one «**portal**» carrying both DM and SM quantum numbers:
 - Vector (dark photon)
 - Axial vector
 - Scalar (dark Higgs)
 - Fermion



Dark photon production and decay

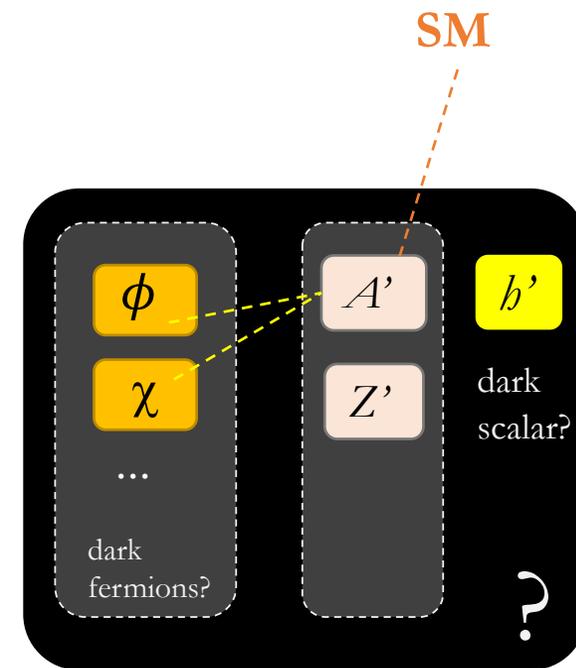
Production:

- In any process with an ordinary photon, we can substitute it with a **dark photon** (A'):
 - A' -strahlung
 - Electron dump and thin-target experiments (visible decays)
 - Missing energy/missing momentum (invisible decays)
 - π^0, η decays
 - e^+e^- annihilations (**missing energy**)
- In the case of **axion-like particles**, we also have the Primakov production mechanism



Decay: two possibilities

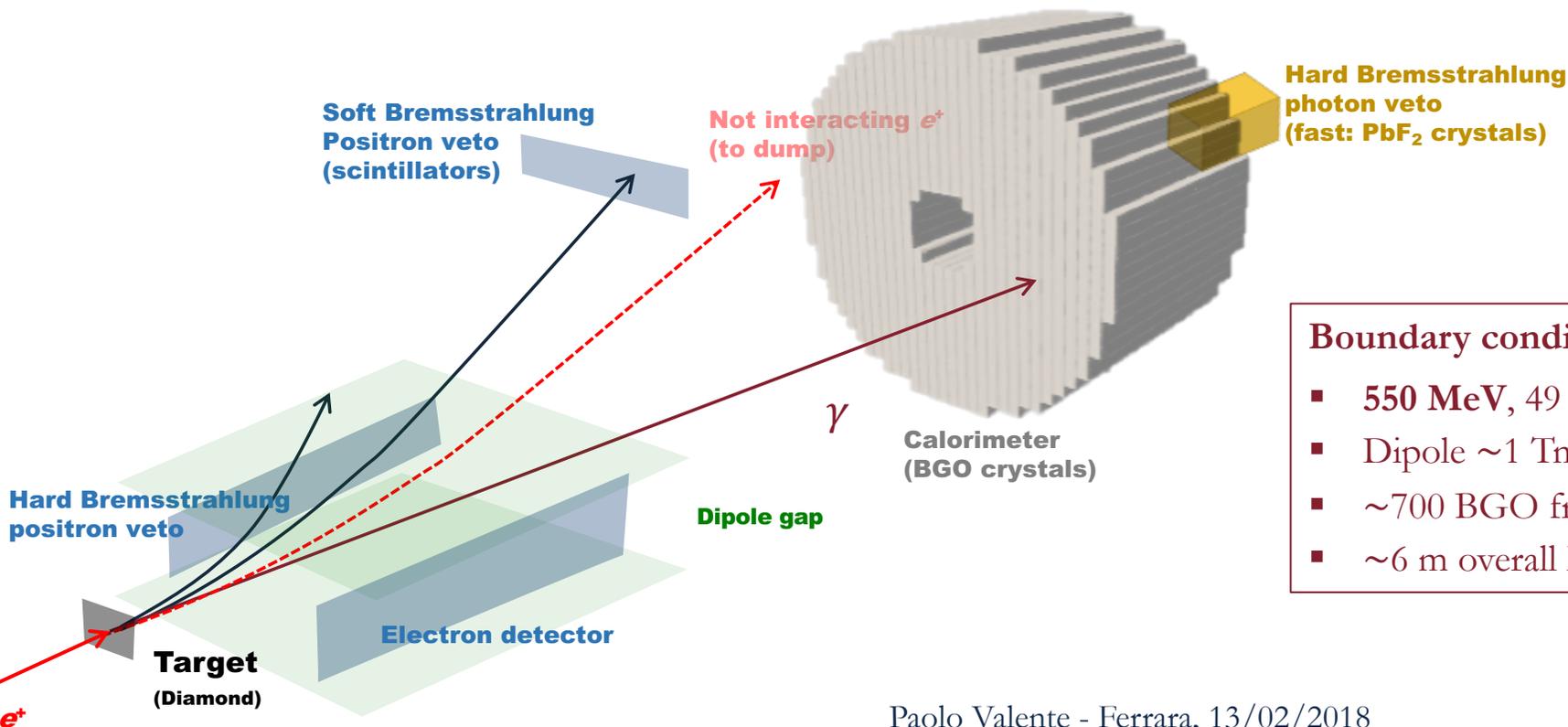
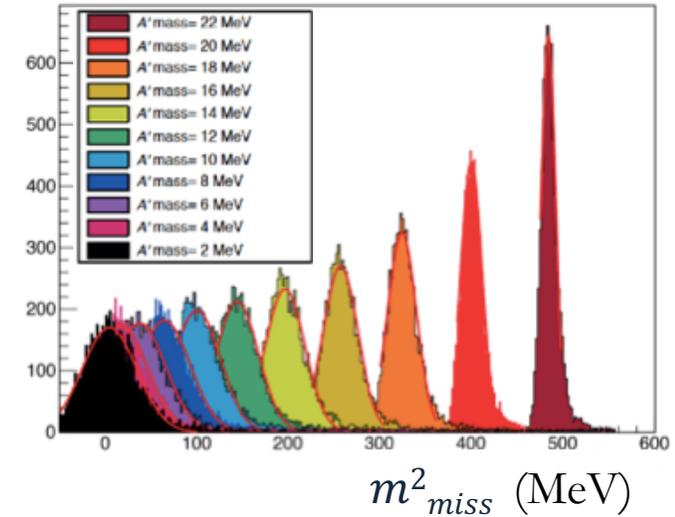
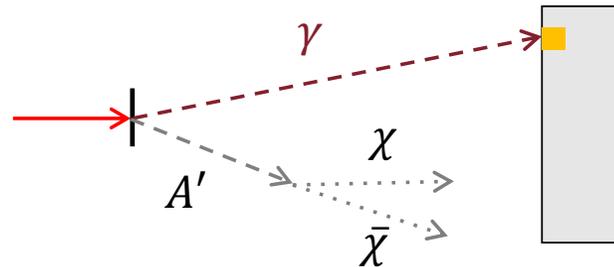
- Looking for decays to SM particles (**lepton pairs** or hadrons if above threshold) or the so-called “**visible**” decays; limits generally rely on the assumption that $A' \rightarrow$ **leptons** is dominant, i.e. the dark photon is the **lightest particle in the dark sector**
- Not looking at the final state, removing the latter assumption, relying on missing energy/momentum or missing mass for identifying “**invisible**” decays $A' \rightarrow \chi\chi$



Dark Sector

PADME experiment

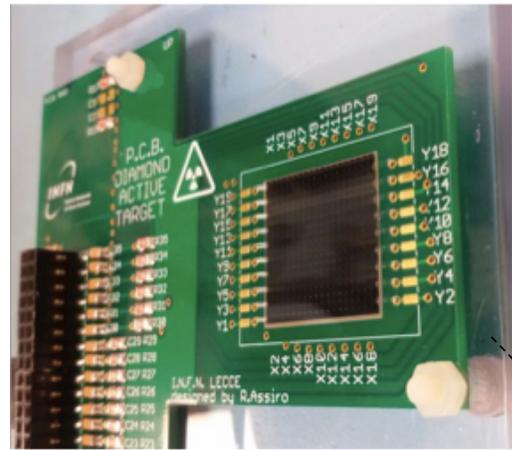
- **Positron beam** with **known 4-momentum**
 1. Small energy and angular spread
 2. Small transverse spot
- Precisely measure the **photon momentum**
- Compute $m^2_{miss} = (P_\gamma - P_{e^+})^2$



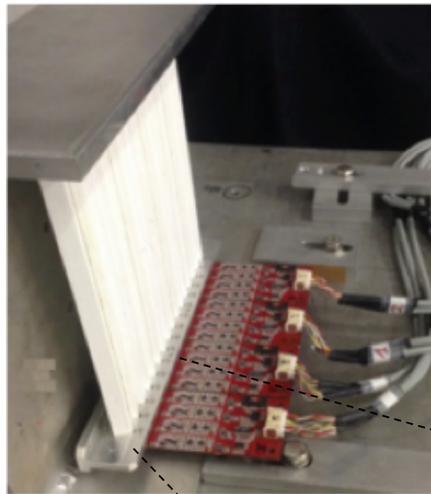
Boundary conditions

- **550 MeV**, 49 pulses/s, **200 ns** long from DAΦNE LINAC
- Dipole ~ 1 Tm, 23 cm gap **available** from CERN SPS
- ~ 700 BGO from L3 calorimeter endcap **available**
- ~ 6 m overall length available in BTF hall

PADME experiment



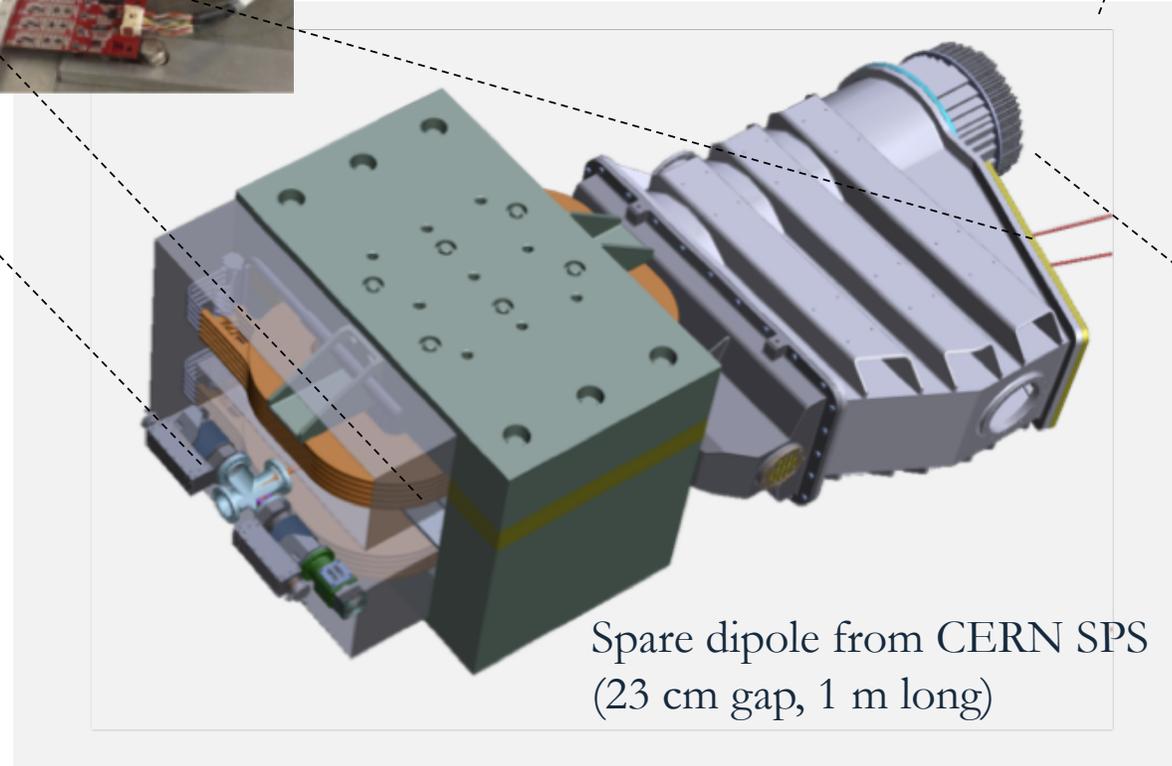
Diamond with grafite strips:
All-Carbon active **target**, beam
position & size and luminosity
monitor; custom electronics



Scintillating bars with **SiPM** readout
for rejecting Bremsstrahlung events
(tagging positrons); inside vacuum vessel



Very fast **PbF₂ Cherenkov**
calorimeter for rejecting 2
and 3 γ backgrounds and
withstand Bremsstrahlung,
fast PMT readout

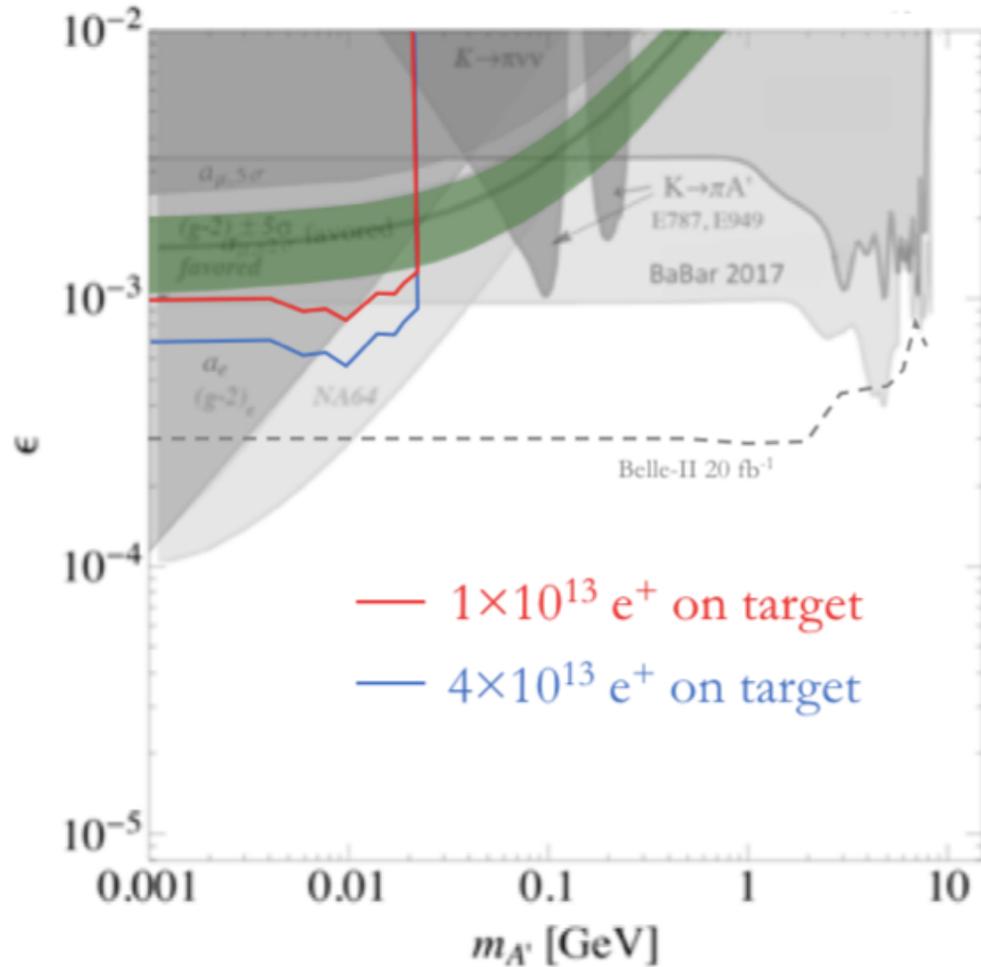


Spare dipole from CERN SPS
(23 cm gap, 1 m long)

BGO calorimeter,
19 mm PMT readout



PADME sensitivity (dark photon invisible decays)



PADME Run 1

- 6 months of data taking in 2018, starting in April

PADME Run 2 at BTF?

- Possible run after Siddharta data taking in late 2019?

Nota bene:

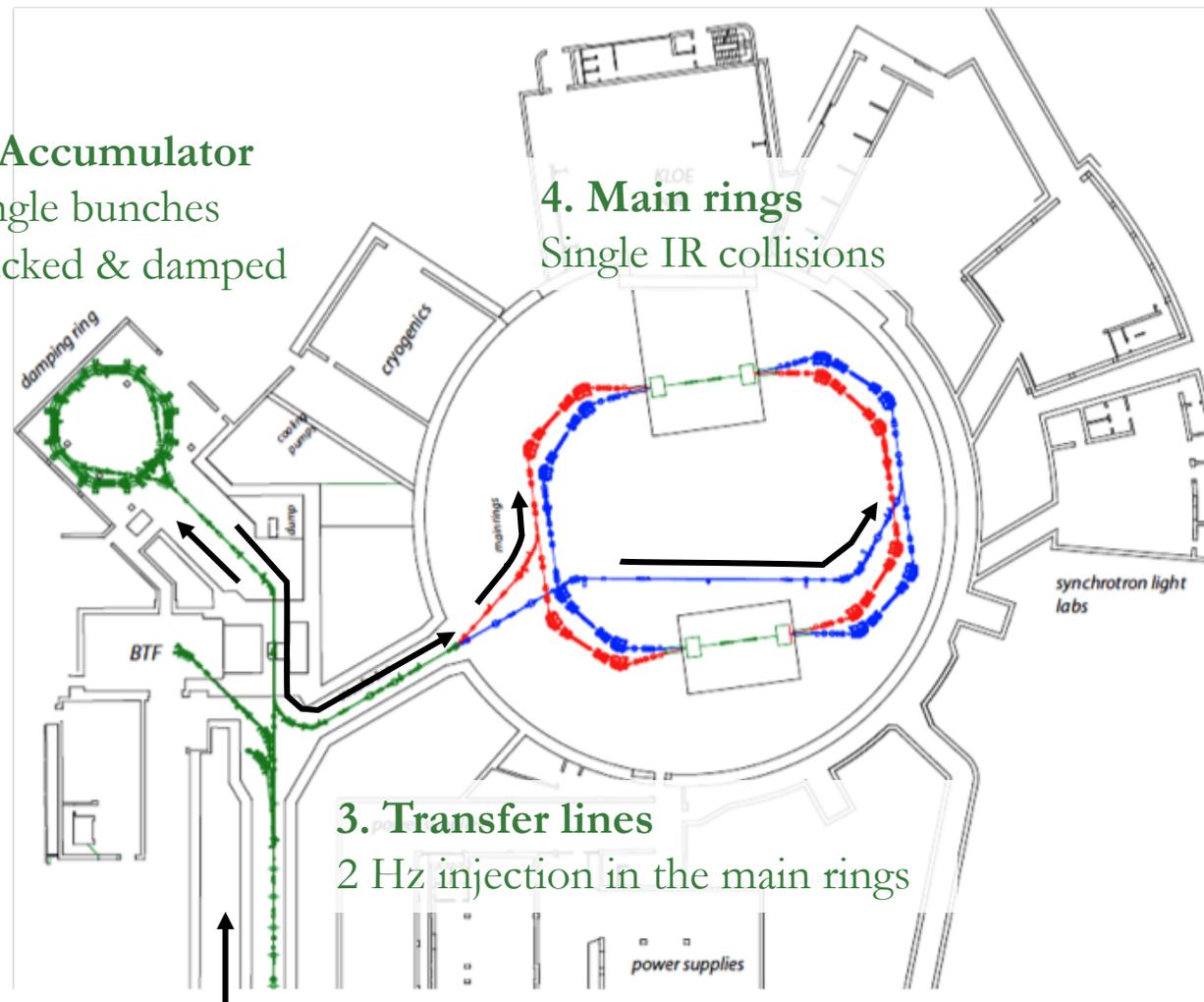
PADME requires long ($\gg 10$ ns) LINAC pulses at $E > 510$ MeV, which is not compatible with continuous injections for DAΦNE topping-up

DAΦNE collider in a nutshell

Parameter	Value
LINAC RF	2857 MHz
LINAC typical charge	1 nC (e^-) 0.5 nC (e^+)
Energy per beam	510 MeV
Machine length	97 m = 324 ns
N. of colliding bunches	100-111
Betatron tune	3.12 (hor.) 1.14 (vert.)
Emittance	1.0 (hor.) 0.01 (vert.) mm mrad
Max. beam current (KLOE run)	2.5 A (e^-) 1.4 A (e^+)
RF frequency	368.67 MHz (main rings) 74 MHz (damping ring)
RF voltage	100-250 kV
Harmonic number	120
Bunch spacing	2.7 ns
Max. Luminosity	
SIDDHARTA run	$4.5 \cdot 10^{32} \text{ cm}^{-2}\text{s}^{-1}$
KLOE run	$2.1 \cdot 10^{32} \text{ cm}^{-2}\text{s}^{-1}$

2. Accumulator
Single bunches
stacked & damped

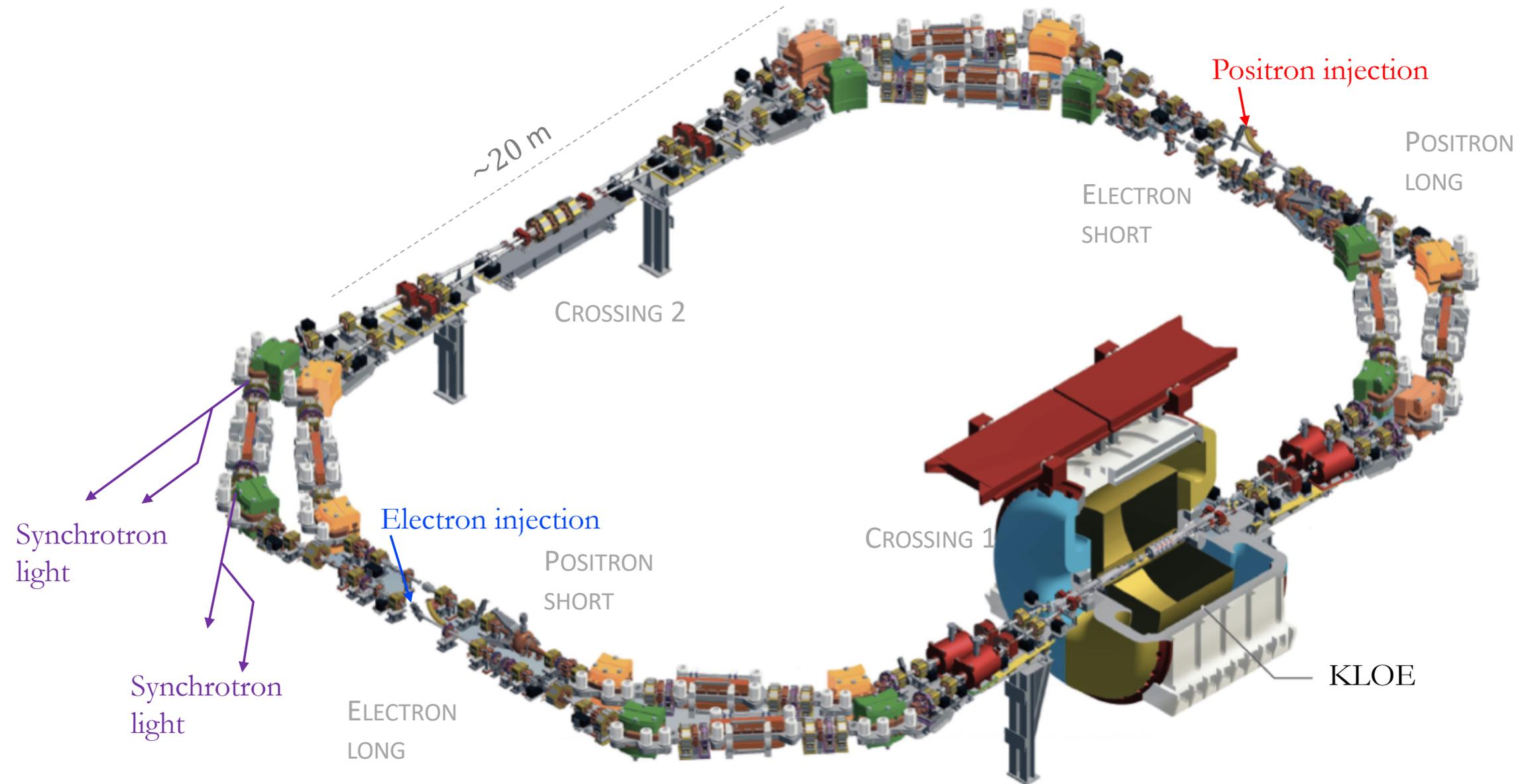
4. Main rings
Single IR collisions



3. Transfer lines
2 Hz injection in the main rings

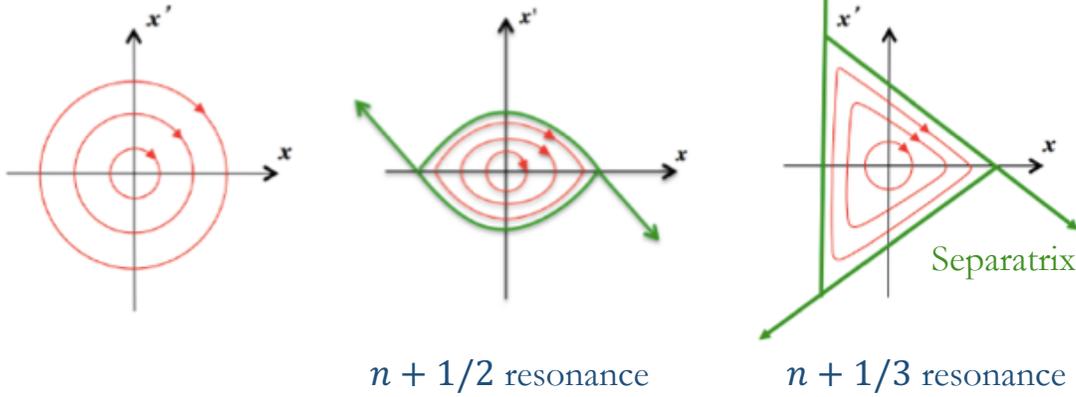
1. High-current LINAC
50 Hz, 10 ns macro-bunches

DAΦNE main rings



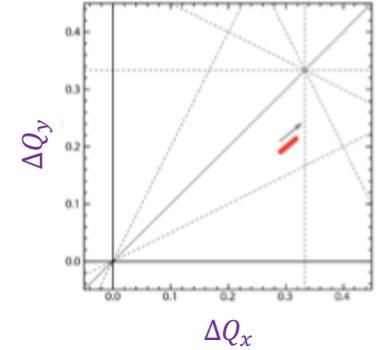
Slow extraction

1. Create an **instability** with one (or more) stable region in the phase space



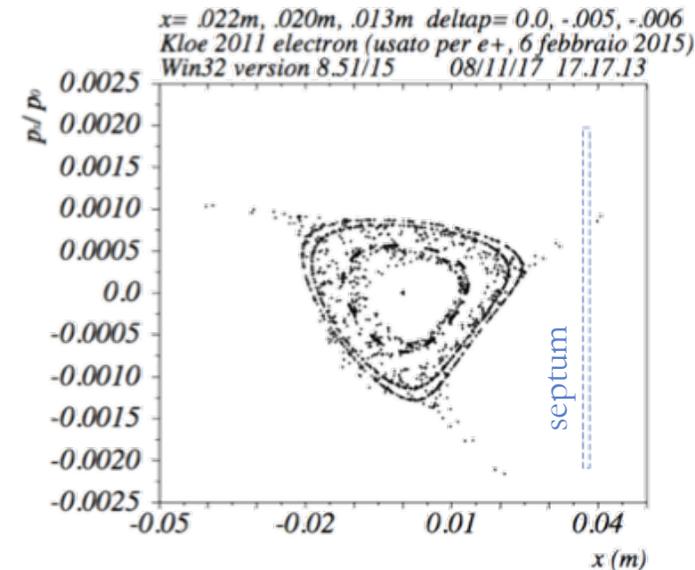
2. Push (at least part of) the beam **out of the stable region**:
 - Shrink the stable region
 - Drive the beam towards a separatrix
3. Extraction septum in a dispersion-free region kicking the beam out (generally in the horizontal plane)

For **electrons** we can use the energy lost by synchrotron radiation to change the tune, slowly approaching the resonance, if the chromaticity is $\neq 0$ ($\Delta Q = C\Delta E$).



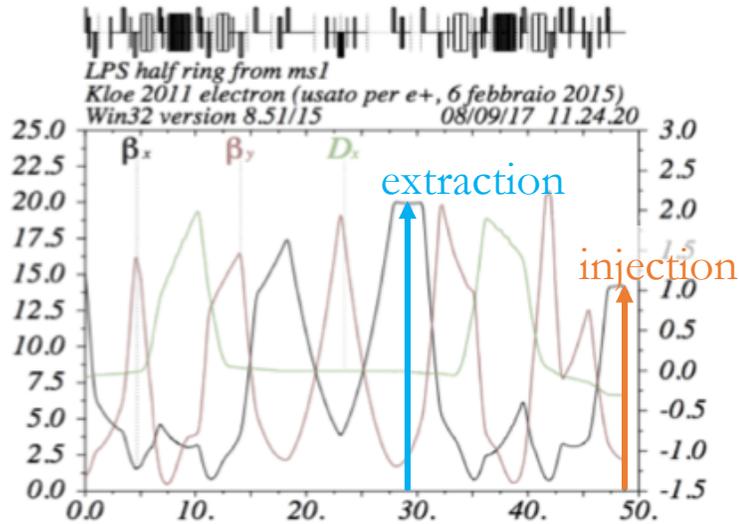
Also called “**monochromatic**” extraction, since the energy at extraction is practically constant. When the energy loss is equal to the spread, all particles are extracted.

Beam injected off-orbit (horizontally), fills an hollow circle in phase space, defined by minimum and maximum emittance.

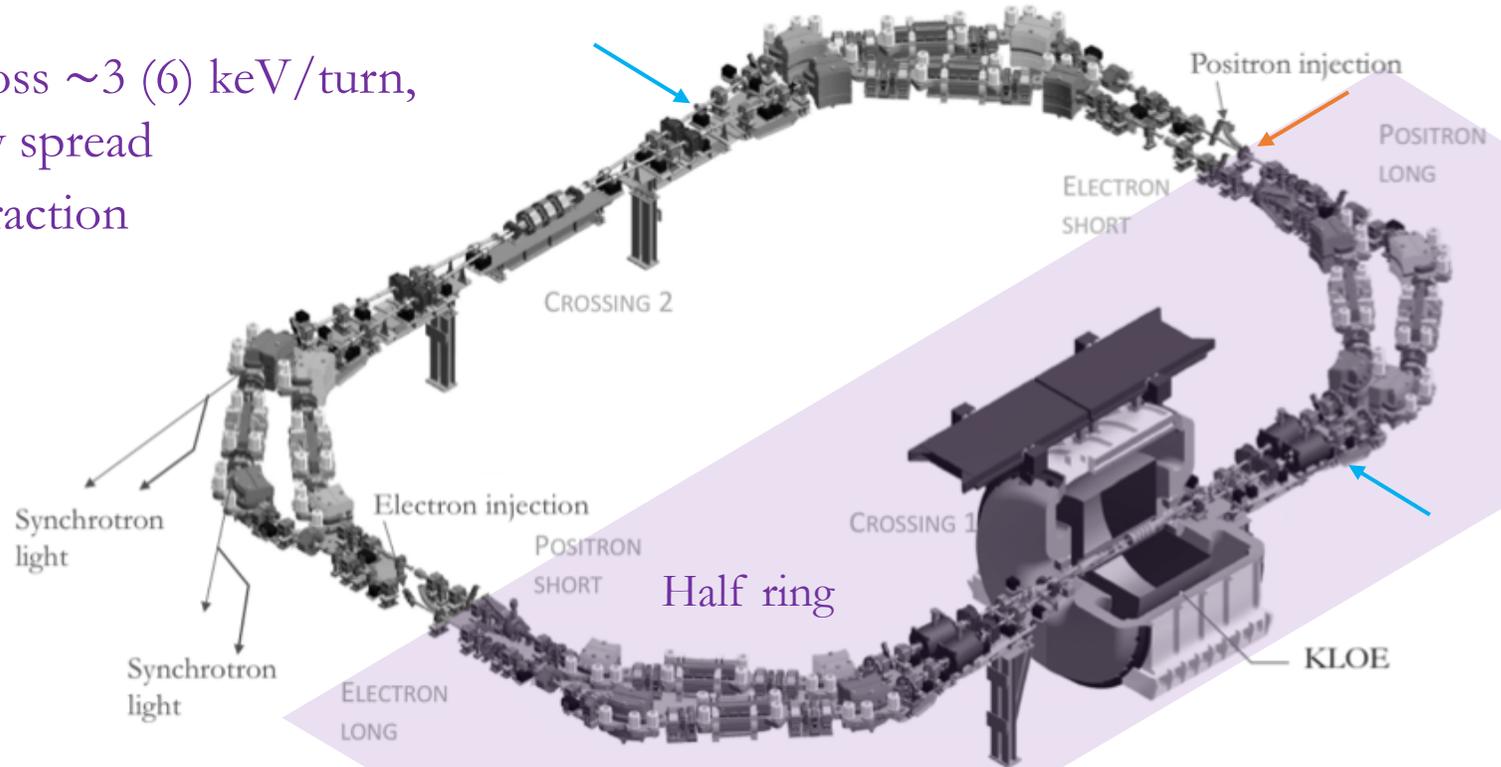


DAΦNE optics

- Inject long beam (up to ~ 324 ns filling the ring)
- Keep ring RF off
 - Or **RF knock-out** (used for protons at CNAO, to be studied for DAΦNE)
- With wigglers off (on), synchrotron loss ~ 3 (6) keV/turn, spill length 0.4 (0.2) ms for 1% energy spread
- $D \sim 0$ and $\alpha \sim 0$, $\beta \sim$ maximum at extraction



S. Guiducci

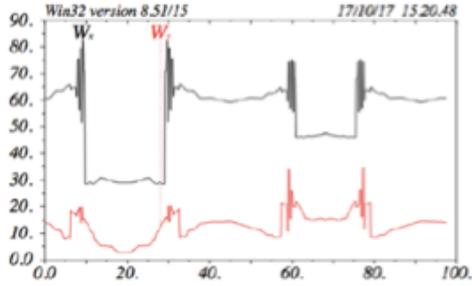
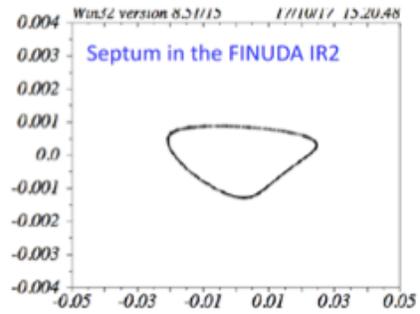
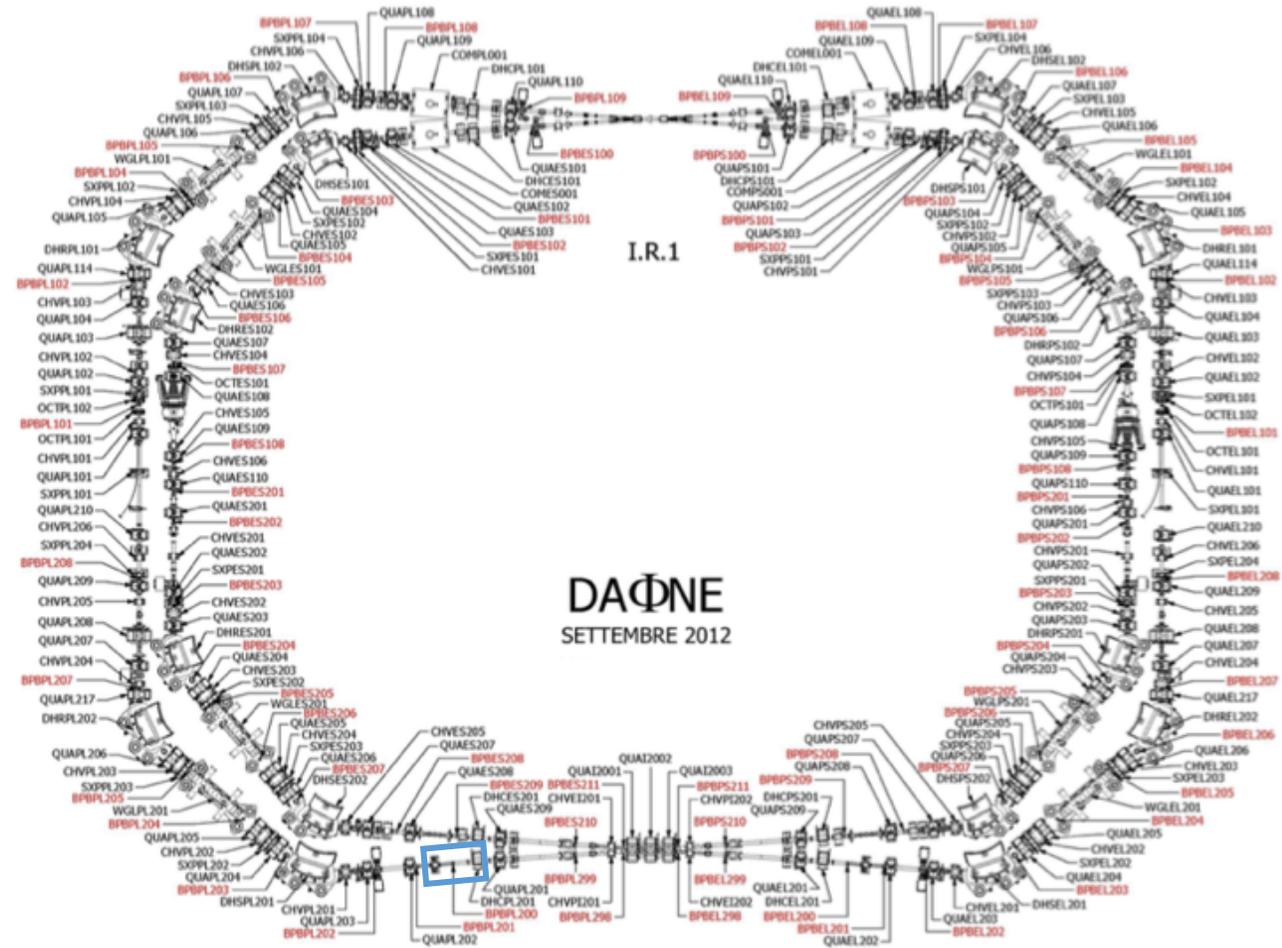


DAΦNE optics

- Chromaticity has to be chosen to get the correct orientation of the stability triangle ($\alpha' > 0$); Strength of the sextupoles.
- Size of injected beam impacts final emittance.
- Adjust the distance of the tune from the resonance: **negative** (positive) chromaticity, approach resonance from **below** (above).

$$Q_x = 4.366, Q_y = 4.27$$

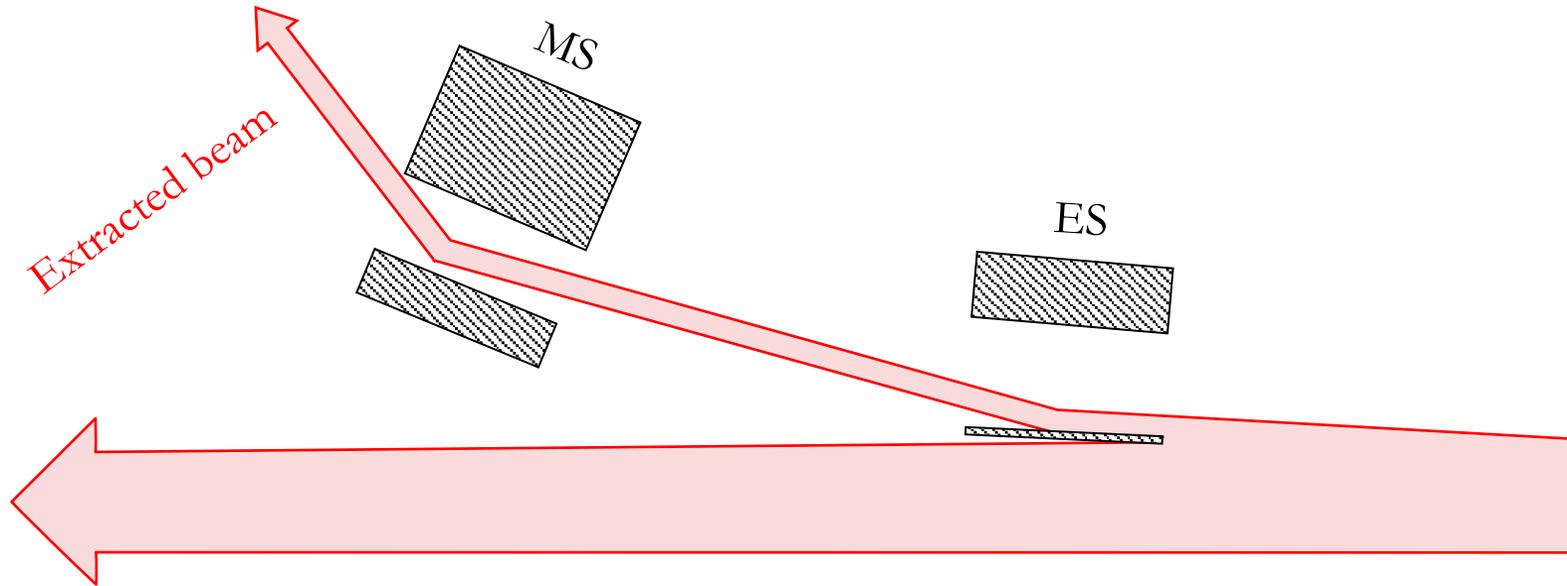
$$C_x = +3.42, C_y = -0.58$$



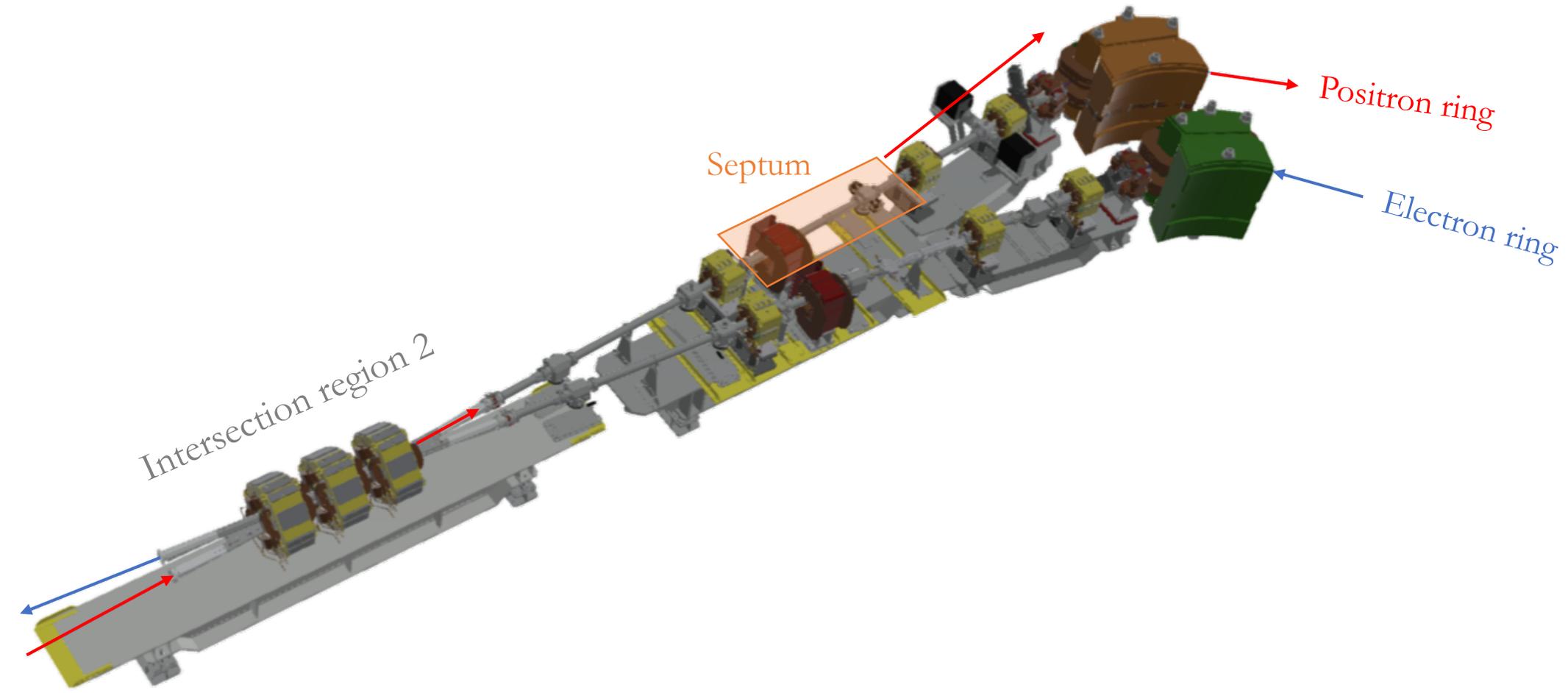
S. Guiducci

Septa

- In the typical configuration an **electrostatic septum** is followed by a **magnetic septum**

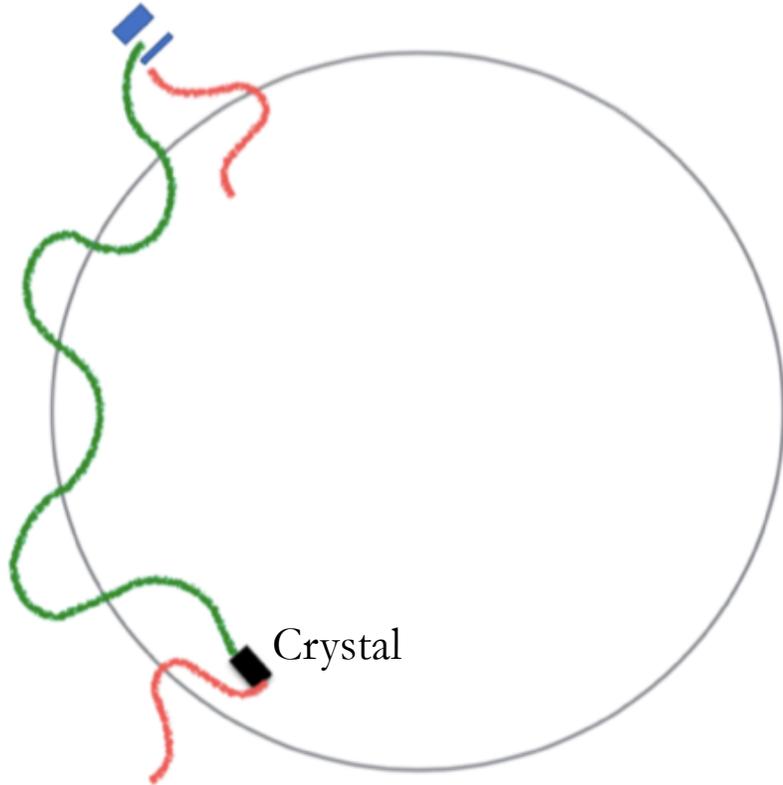


Possible septum location



Why crystals?

Magnetic septum



Beam envelope

- Monochromatic extraction does not allow controlling the **extraction time** (driven by the synchrotron emission)
- RF knock-out can in principle control the process of spilling positrons out of the RF bucket
- An alternative approach: use **crystal channeling**

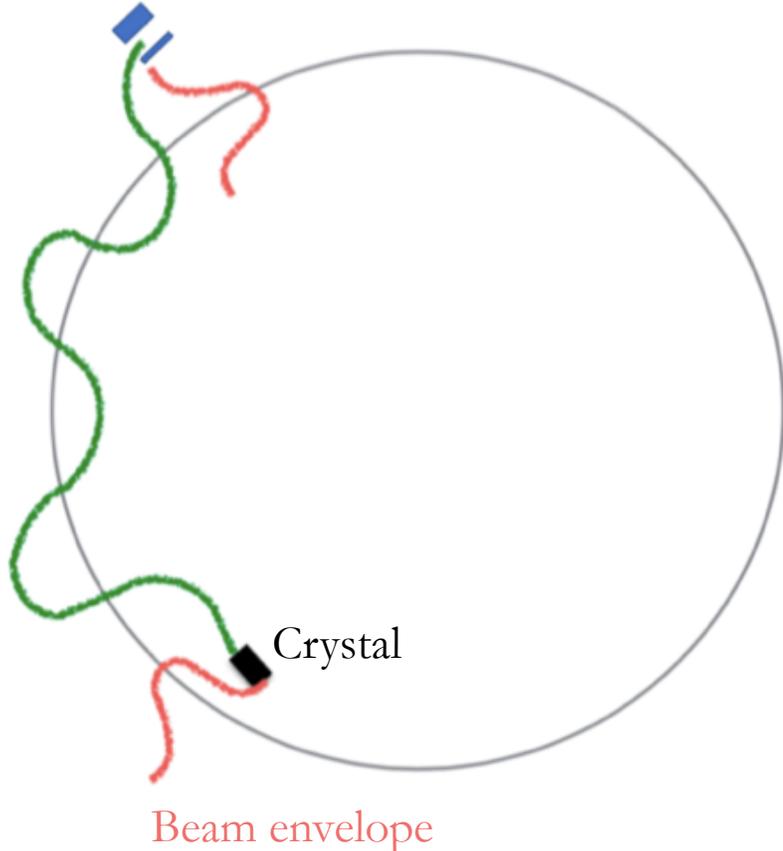
The PADME experiment requirements are basically:

1. Dilute positrons in the **longest possible beam pulse**
2. Maximum particle density determined by the **pile-up** probability in the calorimeter: **$\sim 10^2$ particles/ns**
 - PADME@BTF: $2 \cdot 10^4$ positrons/200 ns (10^6 /s)
 - PADME@Poseydon: $4 \cdot 10^7$ positrons/0.4 ms ($2 \cdot 10^9$ /s)

In both cases, limited by **50 Hz repetition rate**

Why crystals?

Magnetic septum

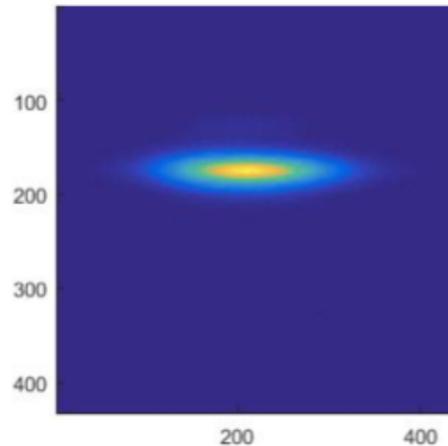


DAΦNE circulating positrons: 1 A current in 120 bunches

$$N_+ \sim 2 \cdot 10^{12}$$

Assuming an average of 1 extracted particle/turn $\rightarrow 1 e^+ / 324 \text{ ns}$
 $\cong 3 \cdot 10^8$ positrons/s:

- 300× with respect to PADME@BTF
 - 1/20 with respect to PADME@Poseydon
- ... but with a continuous time structure of the beam, i.e. “single particle” mode on the target, virtually **no background** in the detector



Transverse size: at SLM: $1 \text{ mm} \times 200 \mu\text{m}$
at IP: $260 \mu\text{m} \times 3 \mu\text{m}$

Emittance: $\epsilon_x = 0.26 \text{ mm mrad}$
 $\epsilon_y = 0.03 \text{ mm mrad}$

Coupling: 0.5%

Energy spread: $0.4 \cdot 10^{-3}$

A few numbers

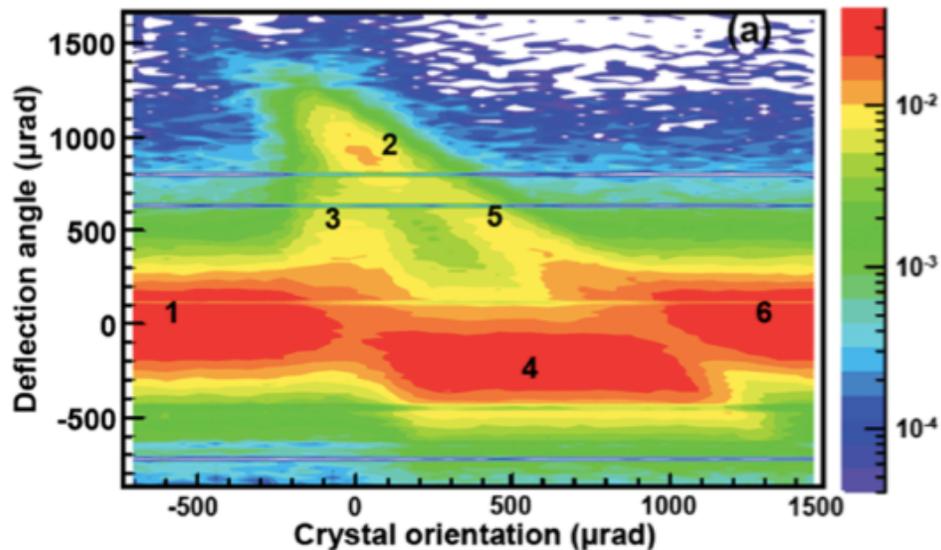
Critical angle for planar channeling, Si(110) at 510 MeV: $\theta_c \sim 210 \mu\text{rad}$

Multiple scattering angle:

510 MeV, 30 μm of Si \rightarrow 480 μrad

855 MeV, 30 μm of Si \rightarrow 280 μrad

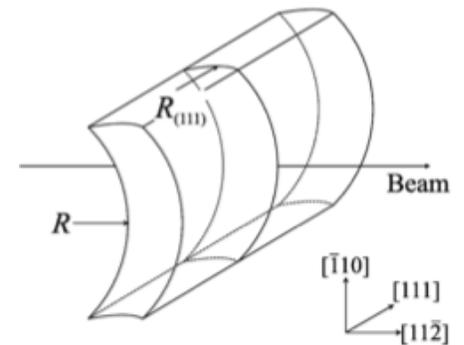
De-channeling length for positrons at 510 MeV $\sim 400 \mu\text{m}$



MAMI experiment

910 μrad **bending** along Si(111) planes,
 $L = 30.5 \mu\text{m}$, $R = 33.5 \text{ mm}$
855 MeV electrons

Si(111) at 855 MeV $R_c \sim 0.4 \text{ mm}$



A few numbers

Problem:

Use thin crystals to reduce **multiple scattering** implies getting close to breaking limit for bending radius for having a sufficient deflection angle $\vartheta = L/R$

$\sim 1 \div 2$ mrad deflection angle: $2 \div 4$ cm from beam at **20 m distance**

Material	x_0 (mm)
Si	93.7
Ge	23
W	3.5

Critical radius of W(110) is $\sim 1/7$ of Si(110), but $x_0 \sim 1/27$
Unless multiple scattering is NOT an issue...

Cut (σ)	Fraction
3.1	$1 \cdot 10^{-3}$
3.7	$1 \cdot 10^{-4}$
4.3	$1 \cdot 10^{-5}$
4.75	$1 \cdot 10^{-6}$
5.2	$1 \cdot 10^{-7}$
5.6	$1 \cdot 10^{-8}$
6.0	$1 \cdot 10^{-9}$
7.0	$1.3 \cdot 10^{-12}$

$2 \cdot 10^4$ e^+ / s at 5.6σ

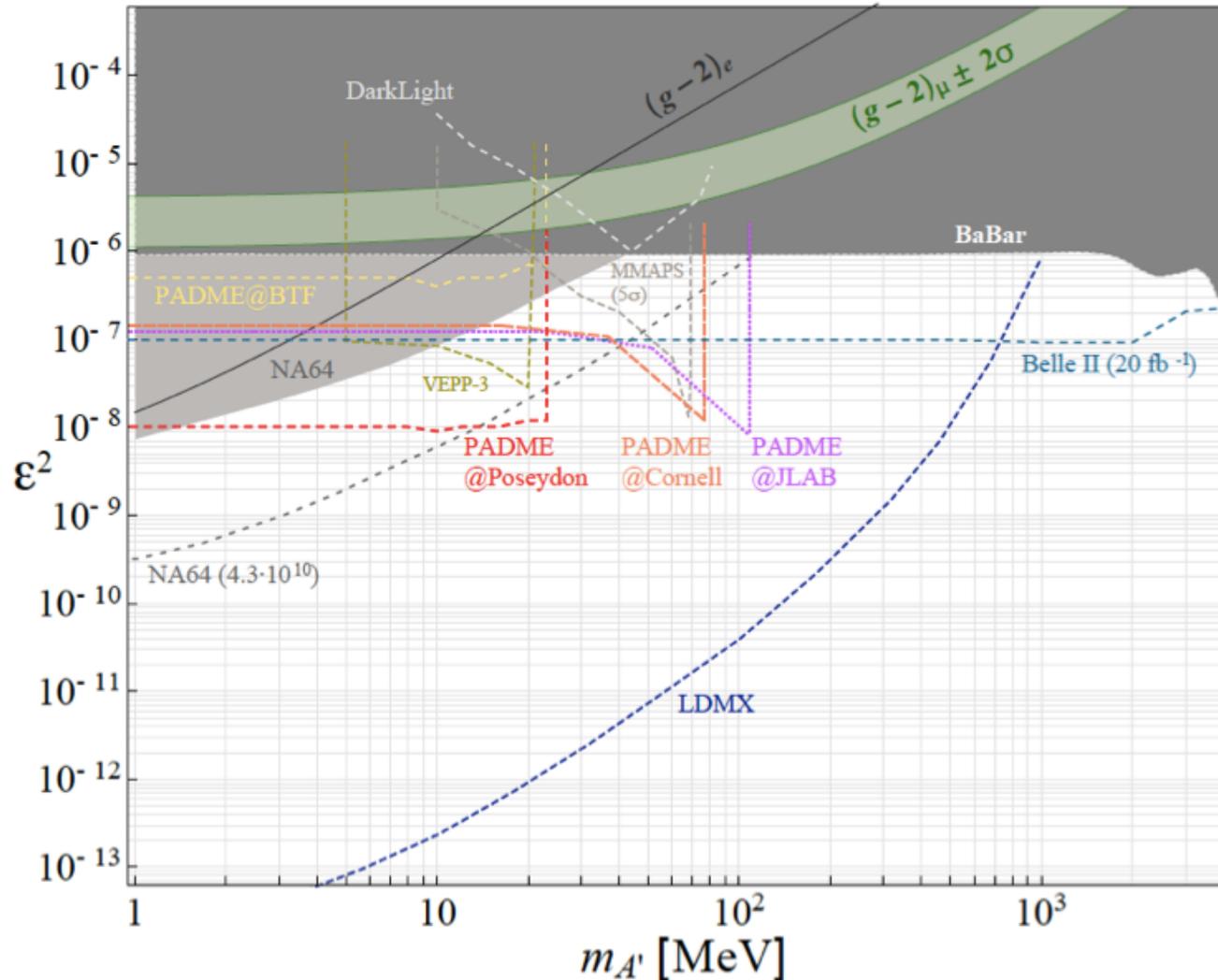
Assuming **all of them** are lost,
lose entire beam in 10^8 turns, i.e. 0.3 s

Problem:

How to send particles to the halo?

PADME@Poseydon sensitivity

Invisibly Decaying Dark Photon



- Increase in sensitivity: 100× with respect to **PADME@BTf** (1 year of running shown)
- Sensitivity still limited to $m_{A'} \sim 24 \text{ MeV}/c^2$

Possible to extend to higher masses increasing the beam energy **but**:

- (Some of the) DAΦNE dipoles already close to field saturation
- Currently 550 MeV is the maximum positron energy from the LINAC: the ring should **ramp up**
- Significant **cost** and **time**
- Only improves with **square root** of beam energy

Complementary to:

PADME@Cornell

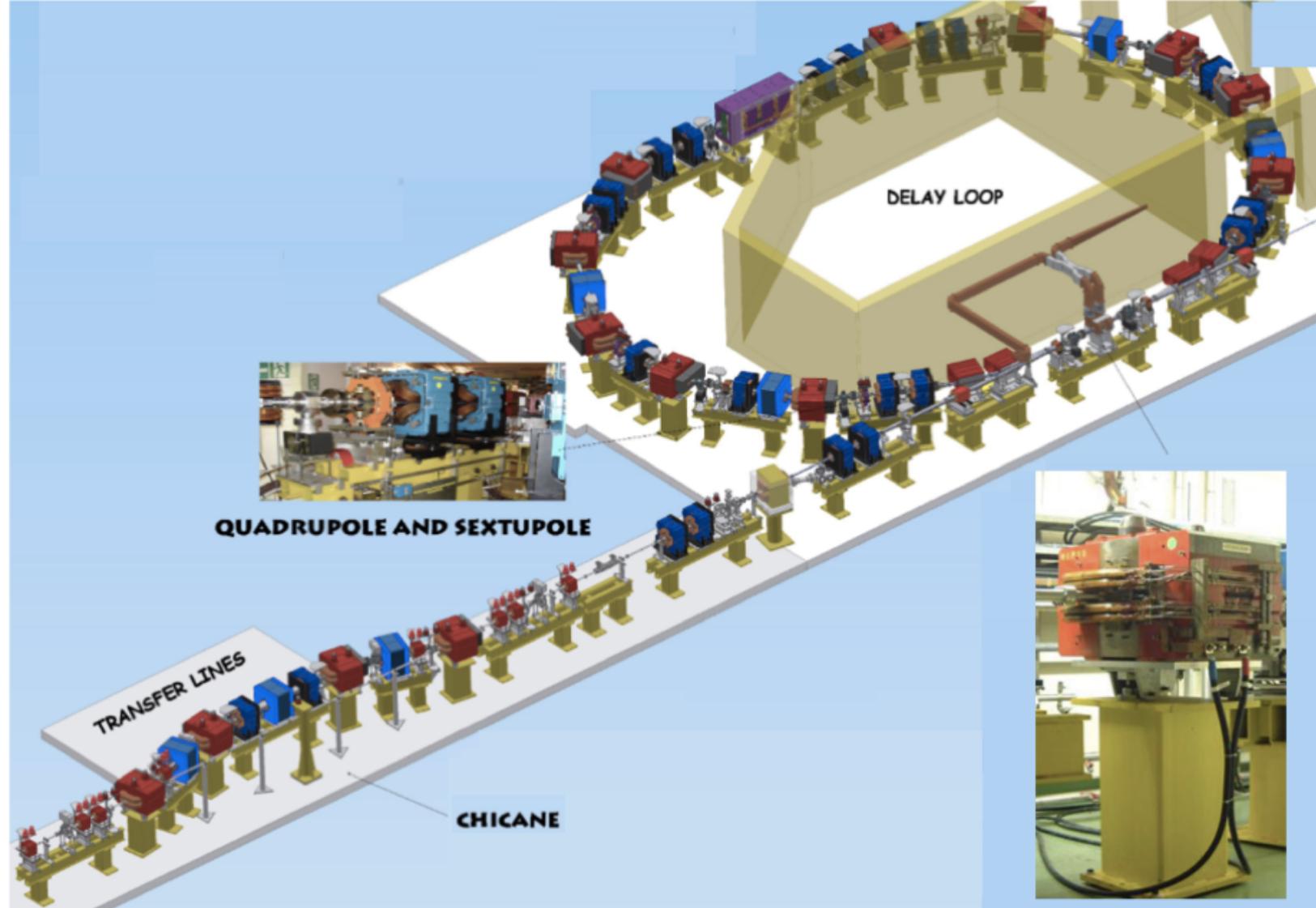
- If NSF MRI proposal is successful ≥ 2022

PADME@JLAB

- No news about positron source at CEBAF

CTF3 magnets

- Official request to CTF3 collaboration for several dipoles, quadrupoles, sextupoles, correctors and diagnostics equipment from machine de-commissioning



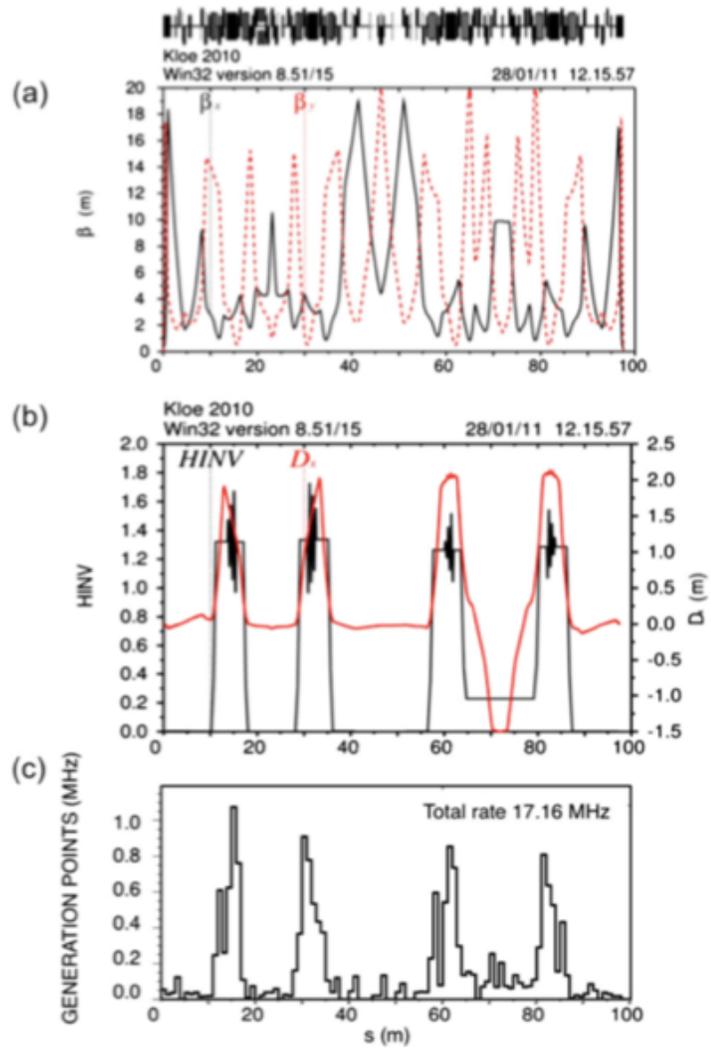
Considerazioni

- Il LINAC di DAΦNE è uno delle pochissime sorgenti di positroni di alta energia (sebbene <0.55 GeV) al mondo
- L'uso di uno dei due anelli di DAΦNE come allungatore dell'impulso del LINAC fornirebbe un fascio di positroni estratto con caratteristiche superiori rispetto a quanto ottenibile alla BTF, e con impulsi lunghi fino a **frazioni di ms/20 ms**
- La **sensitività** di un esperimento di ricerca di mediatori oscuri leggeri come PADME potrebbe essere migliorata di **2 ordini di grandezza** (incremento della statistica di 4 ordini di grandezza)
- L'uso del channeling in luogo dell'estrazione risonante consentirebbe di avere sul setto di estrazione un flusso **praticamente continuo** di positroni, avvicinando le caratteristiche del fascio a quelle ottenibili esclusivamente da un LINAC continuous-wave; un esperimento tipo PADME potrebbe essere condotto in «**single particle**» ovvero praticamente in **assenza di fondo**.
- La disponibilità di un fascio di positroni con queste caratteristiche sarebbe estremamente utile anche per la ricerca nel campo del channeling stesso, per esempio per lo studio della **produzione di radiazione**
- L'uso dell'anello di positroni per questo tipo di attività consentirebbe comunque **l'operazione delle linee di luce di sincrotrone** nell'anello (completamente separato) degli elettroni
- Il costo dell'operazione **senza wigglers**, eventualmente ulteriormente ottimizzato, può essere **$<50\%$** rispetto ad oggi
- **Ad oggi non esiste un piano di medio-lungo termine per il complesso dell'acceleratore DAΦNE**, solo la ragionevole ipotesi che il LINAC e la BTF continueranno ad essere operative anche nel prossimo decennio

P. Valente, "POSEYDON - Converting the DAFNE Collider into a double Positron Facility: a High Duty-Cycle pulse stretcher and a storage ring"

arXiv:1711.06877 [physics.acc-ph], Report INFN-17-15/LNF

Touschek losses



Dispersion

M. Boscolo