# **Di-hadron Fragmentation Functions (DiFF)**





 $P_h \mu = P_1 \mu + P_2 \mu$  $R \mu = (P_1 \mu - P_2 \mu) / 2$ 

$$R_T^2 = \frac{z_1 z_2}{z^2} M_h^2 - \frac{z_2}{z} M_1^2 - \frac{z_1}{z} M_2^2$$

Bacchetta & Radici, P.R.D**67** (03) 094002

		Quark polarization		
		Unpolarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)
Hadron polarization	U	$D_1$	$G_{1} \perp \mathbf{s_{L}} \cdot \mathbf{P_{hT}} \times \mathbf{R_{T}}$ $\left( \bullet \bullet \rightarrow \bigcirc \right) - \left( \bullet \bullet \rightarrow \bigcirc \right)$	$H_{1}^{*} \hat{z} \cdot s_{T} \times R_{T}$ $H_{1} \perp \hat{z} \cdot s_{T} \times P_{hT}$ $\left( \hat{z} \stackrel{\circ}{\longrightarrow} \right) - \left( \hat{z} \stackrel{\circ}{\longrightarrow} \right)$

# **Di-hadron Fragmentation Functions (DiFF)**

Bianconi et al., P.R. D**62** (00) 034008



 $P_h \mu = P_1 \mu + P_2 \mu$  $R \mu = (P_1 \mu - P_2 \mu) / 2$ 

 $\int d\mathbf{P}_{hT}$  collinear framework

leading twist,  $R_T^2 \ll Q^2$ 



#### Access to transversity via DiFF

Collins, Heppelman, Ladinsky, N.P. B420 (94)



correlation between **quark polarization** and  $\mathbf{R}_T = (z_2 \mathbf{P}_{1T} - z_1 \mathbf{P}_{2T}) / z$ or, equivalently, azimuthal orientation of  $(h_1, h_2)$  plane

(only if  $h_1 \neq h_2$ )

effect encoded in  $h_1(x) H_1^{\triangleleft}(z, M_h^2)$ 

 $z=z_1+z_2$  $P_h^2=M_h^2 <=> R_T^2$ 

alternative to Collins effect

Radici, Jakob, Bianconi, P.R.D65 (02) 074031 Bacchetta & Radici, P.R.D67 (03) 094002

Bacchetta et al., JHEP **1303** (13) 119 Radici et al., JHEP **1505** (15) 123

#### extraction of DiFF from e+e-



back-to-back hadron pairs  $\rightarrow \cos(\Phi_R + \overline{\Phi}_R)$  modulation

Artru & Collins, Z.Ph. C69 (96) 277

Boer, Jakob, Radici,  
*P.R.D67 (03) 094003*  

$$A^{\cos(\phi_R + \bar{\phi}_R)} = \frac{\sin^2 \theta_2}{1 + \cos^2 \theta_2} \frac{|\mathbf{R}_T|}{M_h} \frac{|\bar{\mathbf{R}}_T|}{\bar{M}_h} \frac{\sum_q e_q^2 H_1^{\triangleleft q}(z, M_h^2)}{\sum_q e_q^2 D_1^q(z, M_h^2)} \bar{D}_1^{\bar{q}}(\bar{z}, \bar{M}_h^2)$$
same as in SIDIS

### extraction of DiFF from e+e-



Boer, Jakob, Radici, P.R.D67 (03) 094003  $A^{\cos(\phi_{R}+\bar{\phi}_{R})} = \frac{\sin^{2}\theta_{2}}{1+\cos^{2}\theta_{2}} \frac{|\mathbf{R}_{T}|}{1+\cos^{2}\theta_{2}} \frac{|\mathbf{R}_{T}|}{M_{h}} \frac{|\bar{\mathbf{R}}_{T}|}{\bar{M}_{h}} \frac{\sum_{q} e_{q}^{2} H_{1}^{\triangleleft q}(z, M_{h}^{2})}{\sum_{q} e_{q}^{2} D_{1}^{q}(z, M_{h}^{2})} \frac{\bar{H}_{1}^{\triangleleft \bar{q}}(\bar{z}, \bar{M}_{h}^{2})}{\ln \mathrm{SIDIS}}$ Same as in SIDIS For D<sub>1</sub> for D<sub>1</sub> for D<sub>1</sub> first extraction of DiFF, but using PYTHIA Courtoy et al., P.R.D85 (12) 114023 Radici et al., JHEP 1505 (15) 123

new Belle data for unpolarized cross section

Artru & Collins, Z.Ph. C69 (96) 277

### e<sup>+</sup>e<sup>-</sup> cross section for ( $\pi\pi$ ) in same hemisphere



# universality of $H_1^*$





Boer, Jakob, Radici, P.R.D67 (03) 094003



Matevosyan, Bacchetta, Boer, Courtoy, Kotzinian, Radici, Thomas, arXiv:1802.01578

in the cross section, all terms  $\perp \hat{\mathbf{h}}$  change sign  $\rightarrow H_1^{\triangleleft}$  in  $A^{\cos(\phi_R + \bar{\phi}_{\bar{R}})}$  same as in SIDIS  $A_{UT}$ 

# helicity DiFF $G_1^{\perp}$

Matevosyan, Bacchetta, Boer, Courtoy, Kotzinian, Radici, Thomas, arXiv:1802.01578

$$\langle T \rangle = \int d\xi \int d\bar{\xi} \int d\phi_R \int d\bar{\phi}_{\bar{R}} \int d\mathbf{q}_T \, T \, d\sigma$$

proposed "jet handedness" asymmetry to extract  $G_1^{\perp}$  vanishes  $\langle \cos 2(\phi_R - \phi_{\bar{R}}) \rangle = 0$ Jet A cos(2(Φ<sub>R</sub>,-Φ<sub>R2</sub>)) - 0 0.0 0 - 0 0.0 0 0.012 Jet A<sup>cos(2(Φ<sub>R,</sub>-Φ<sub>R,</sub>)) 000</sup> confirms BELLE findings 0.008 0.004 Abdesselam et al. (BELLE) arXiv:1505.08020 -0.008 -0.008 -0.012 0.3 0.4 0.5 0.6 0.7 0.8 0.9 1 1.1 -0.012 0.3 0.4 0.5 0.6 0.7 0.8 M [GeV/c<sup>2</sup>] which asymmetry to get  $G_1^{\perp}$ ?  $\langle \mathbf{q}_T^2 (2\cos(\phi_R - \bar{\phi}_{\bar{R}}) - \cos(2\phi_1 - \phi_R - \bar{\phi}_{\bar{R}}) \rangle$  $\propto \sum G_1^{\perp q}(z, M_h^2) \, \bar{G}_1^{\perp \bar{q}}(\bar{z}, \bar{M}_h^2)$ Matevosyan et al.,  $G_1^{\perp}(z, M_h^2) \equiv G_1^{\perp, [0], (1)}(z, M_h^2) - G_1^{\perp, [2], (1)}(z, M_h^2)$ arXiv:1712.06384  $G_1^{\perp,[n],(p)}(z, M_h^2) = z^2 \int d\mathbf{k}_T \left(\frac{\mathbf{k}_T^2}{2M_t^2}\right)^p \int d\xi \, \frac{|\mathbf{R}_T|}{M_h} \, G_1^{\perp,[n]}(z,\xi,|\mathbf{k}_T|,|\mathbf{R}_T|)$ 

N.B.  $\langle \cos(\phi_R - \bar{\phi}_{\bar{R}}) \rangle = 0$  but  $\langle \cos(\phi_R - \bar{\phi}_{\bar{R}}) \rangle (\mathbf{q}_T^2) \neq 0 \rightarrow$  there must be a node

# The power of DiFFs

#### collinear framework $\rightarrow$ factorization theorems



# global fit : the SIDIS data



Braun et al., E.P.J. Web Conf. 85 (15) 02018





Adolph et al., P.L. **B713** (12)



90% replicas

# global fit : the STAR data



# $X^2$ of the fit



# theoretical uncertainties

Single-Spin Asymmetry  
in p-p<sup>+</sup> collisions  

$$A_{UT}(\eta, M_h, P_T) = \int_{d\sigma_0}^{d\sigma_{UT}} d\sigma_0$$
typical cross section for  $a+b \rightarrow c+d$  process  

$$d\sigma_0 \propto \sum_{a,b,c,d} \int \frac{dx_a dx_b}{8\pi^2 \bar{z}} f_1^a(x_a) f_1^b(x_b) \frac{d\hat{\sigma}_{ab \rightarrow cd}}{d\hat{t}} D_1^c(\bar{z}, M_h)$$
quark D<sub>1</sub><sup>q</sup> is well constrained by e<sup>+</sup>e<sup>-</sup> (Montecarlo) but  
we don't know anything about the gluon D<sub>1</sub><sup>g</sup> (e<sup>+</sup>e<sup>-</sup> doesn't help..)  
our choice: compute  $d\sigma_0$  with D<sub>1</sub><sup>g</sup> (Q\_0) = 
$$\begin{cases} 0 \\ D_1^u (Q_0) / 4 \\ D_1^u (Q_0) \end{cases}$$
deteriorates our e<sup>+</sup>e<sup>-</sup> fit as  $\chi^2/dof = \begin{cases} 1.69 \\ 1.81 \\ 2.96 \\ 2.01 \end{cases}$ 
background  $\rho$  channels

### comparison with previous fit



#### comparison with previous fit



### comparison with previous fit





# sensitivity to gluon D<sub>1</sub><sup>g</sup>



need dihadron multiplicities from RHIC and better deuteron data from COMPASS

# **MC-based models of FF**



- extended NJL quark-jet model
- string fragmentation model

# extended NJL quark-jet model



quark hadronization chain Nambu-Jona Lasinio model at each vertex  $D_1^{q \rightarrow h}(z, p_T^2) = MC$  statistical average of proper multiplicities results for q=u,d,sand  $h=\pi, K, \rho, K^*, \varphi, N$ 

- p<sub>T</sub>-dependence deviates from Gaussian Matevosyan e

-  $< p_T^2 >$ : flavor- and nonlinear z-dependent

Matevosyan et al., P.R. D**85** (12) 014021 Matevosyan et al., P.R. D**83** (11) 114010 D**86** (12) 059904(E)

- includes transverse momentum Bentz et al., P.R. D94 (16) 034004
   and spin correlations Matevosyan et al., P.R. D95 (17) 014021
- extended to DiFF Matevosyan et al., P.R. D88 (13) 094022
- and polarized DiFF

Matevosyan et al., P.R. D**96** (17) 074010

Matevosyan et al., P.R. D**97** (18) 014019





 $H_1^{\triangleleft}$  as a recursive Collins effect

# string fragmentation model



Kerbizi, Artru, Belghobsi, Bradamante, Martin, arXiv:1802.00962

#### **SIDIS** Collins effect

