

### Understanding the x-ray emission spectrum after excitation with a source of x-rays: from theory to experiment

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### Outline

• Description of x-ray emission

- Detector modification
  - Pulse pile-up.
  - Detector response function.



#### MULTIPLE SCATTERING IN X-RAY SPECTROMETRY

- X-rays penetrate deeply into the matter, and, in a thick medium, give place to a phenomenon known as multiple scattering (i.e, multiple collisions).
- Multiple scattering characterizes the radiation field in a thick medium
- Multiple scattering models describe the influence of the prevailing interactions in the x-ray regime to describe the radiation field.
- The emitted X-ray fluorescence spectrum is easily obtained from the radiation field, and is strongly modified by multiple scattering.
- Another important factor to characterize the radiation field is the effect of the polarization.



# Scheme of X-ray interaction mechanisms

#### The full description of the radiation field requires the modeling of coupled photon-electron transport



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## Multiple scattering is usually described using the Boltzmann transport model

## The photon interactions are depicted with the interaction kernels $k_i$

$$\eta \frac{\partial}{\partial z} f(z, \vec{\omega}, E) = -\mu(E) f(z, \vec{\omega}, E)$$

$$+ \sum_{i}^{\text{all interactions}} \int_{0}^{\infty} \left( \int_{4\pi} \cup (z) k_{i} \left( \vec{\omega'}, E', \vec{\omega}, E \right) f\left( z, \vec{\omega'}, E' \right) d\omega' \right) dE'$$

$$+ S(z, \vec{\omega}, E)$$

# Not all the radiative contributions involved in an X-ray transport process are considered by the Boltzmann model and by the used interaction kernels $k_i$

# X-ray production mechanisms from coupling terms

The full description of the radiation field requires the modeling of coupled photon-electron transport



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## MS is better described using the modified Boltzmann transport model

## The Boltzmann transport model has been recently modified to include the electron-photon contributions

 $\begin{aligned} \underbrace{ \boxed{\eta\partial/\partial z \ f\uparrow p \ (z,\omega,E) = -\mu\uparrow p \ (E) \ f\uparrow p \ (z,\omega,E) + \sum i\uparrow all \ photon} \\ \underline{interactions} \\ \underbrace{ \int 0\uparrow \infty } \underbrace{ \left( \int 4\pi\uparrow \\ \hline \cup (z) k\downarrow i \uparrow p \rightarrow p \ (\omega\uparrow ,E\uparrow ,\omega,E) \ f\uparrow p \ (z,\omega\uparrow ,E\uparrow ) \ d\omega\uparrow \right) d} \\ \underline{E\uparrow} + \underbrace{ \sum j\uparrow all \ coupling \ terms} \\ \underbrace{ \int 0\uparrow \infty } \underbrace{ \left( \int 4\pi\uparrow \\ \hline \cup (z) k\downarrow j \uparrow e \rightarrow p \ (\omega\uparrow ,E\uparrow ,\omega,E) \ f\uparrow p \ (z,\omega,E) \right) } \\ \underline{f\uparrow} , \underline{E\uparrow} \ d\omega\uparrow \ dE\uparrow \ + S\uparrow p \ (z,\omega,E) \end{aligned}$ 

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We need to include also the polarization WHY POLARIZATION?



By considering polarization we improve the model of photon diffusion

a good approximation in many cases, but not for phenomena that are influenced by their wave properties



#### PHOTON DIFFUSION IS DESCRIBED BY A "VECTOR" TRANSPORT EQUATION (THE 1-D EQUATION IS SHOWN HERE)

 $\begin{aligned} \frac{[\neg \eta\partial/\partial z \ f\uparrow p \ (z,\omega,E) = -\mu\uparrow p \ (E) \ f\uparrow p \ (z,\omega,E) + \sum i\uparrow all \ photon}{interactions \ [] \int 0\uparrow \infty \ (\int 4\pi\uparrow \ [] \cup (z) \ H\downarrow i\uparrow p \to p \ (\omega\uparrow \ ,E\uparrow \ ,\omega,E) \ f\uparrow p \ (z,\omega\uparrow \ ,E\uparrow \ )d\omega\uparrow \ )d}{E\uparrow \ + \sum i\uparrow all \ coupling \ terms \ [] \int 0\uparrow \infty \ [] (\int 4\pi\uparrow \ [] \cup (z) \ H\downarrow i\uparrow e \to p \ (\omega\uparrow \ ,E\uparrow \ ,\omega,E) \ f\uparrow p \ (z,\omega,E)} \\ \frac{f\uparrow \ ,E\uparrow \ )d\omega\uparrow \ )dE\uparrow \ + S\uparrow p \ (z,\omega,E)}{I \ ,E\uparrow \ ,\omega,E)} \end{aligned}$ 

where

$$\overline{f} = \begin{bmatrix} I(z, \overline{\omega}, \lambda) \\ Q(z, \overline{\omega}, \lambda) \\ U(z, \overline{\omega}, \lambda) \\ U(z, \overline{\omega}, \lambda) \\ V(z, \overline{\omega}, \lambda) \end{bmatrix}$$



#### VECTOR TRANSPORT EQUATION (CONT.)

#### where

 $\mathrm{H}^{(S)}(\vec{\varpi},\lambda,\vec{\varpi}',\lambda') {=} \mathrm{L}^{(S)}(\pi - \Psi) \mathrm{K}^{(S)}(\vec{\varpi},\lambda,\vec{\varpi}',\lambda') \mathrm{L}^{(S)}(-\Psi')$ 

## $H^{(S)}$ = kernel matrix in the meridian plane of reference

 $K^{(S)}$  = scattering matrix in the scattering plane of reference



#### IMPORTANT PROPERTIES OF THE "VECTOR" TRANSPORT EQUATION

- Describes the evolution of the full polarization state (not only the intensity of the beam)
- Is linear (for the Stokes representation)
- Requires the simultaneous solution of the whole set of transport equations
- Cannot be transformed in a scalar equation !! (due to the coupling in the scattering term)



## TWO RELEVANT ASPECTS

- A collision always changes the polarization state
- The angular distribution for scattered unpolarized and polarized photons is very different



### MCSHAPE

- MCSHAPE is a Monte Carlo code developed at the University of Bologna which can simulate the diffusion of photons with arbitrary polarization state and has the unique feature of describing the evolution of the polarization state along the interactions with the atoms.
- The adopted transport model is derived from the so called Boltzmann-Chandrasekhar 'vector' transport equation. The polarization state of the photons is described by using the Stokes parameters I, Q, U and V, having the dimension of intensities and containing the physical information about the polarization state.
- This code simulates the propagation in heterogeneous media of photons injected by either polarized (i.e., synchrotron) or unpolarized sources (x-ray tubes).
- Website: http://shape.ing.unibo.it



## Differences between the computational structures of scalar and vector MC models





Schematic diagram of a simulation with MCSHAPE, compared with the experimental steps for a spectrum measurement.



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### Outline

#### **Detector modification**

- Pulse pile-up.
- Detector response function.



First principles pulse pile-up balance equation and fast deterministic solution

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L. Sabbatucci, Viviana Scot, J.E. Fernandez: Multi shape pulse pile-up correction: the MCPPU code. Radiat. Phys. Chem. 104,(2014) 372-375

L. Sabbatucci, J.E. Fernandez: First principles Pulse Pile-Up balance equation and fast deterministic solution, Radiat. Phys. Chem. 137 (2017) 12-17

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### Pulse Pile-Up problem



• Radiation emission is a process randomly spaced in time. At high counts many pulses have time gap much smaller then their width, then pile-up effects occur.

Pulse pile-up has three consequences:

- Loss of counts: the sum of the pulses is detected instead of the separate pulses
- Piled-up pulses are distributed at wrong energies
- The whole spectrum is distorted since the lost pulses are not collected at the proper energies
- Modern detectors usually comprise built-in electronics which partially reduce pile-up effects either by limiting the pulse tail (i.e. Rise Time Discrimination, RTD) and/or by using additional rejection circuits.



### Monte Carlo Approach: The code MCPPU

- MCPPU (Monte Carlo Pulse Pile-Up) is a general purpose code to perform post-processing pile-up correction on spectra obtained with different detectors.
- The code can correct pile-up distortions for spectra collected with or without electronic reduction circuitry.
- It is based on a modified version of the Monte Carlo algorithm developed originally by Gardner and Lee, Adv. X-ray Anal. 41 (1999) 941–950 and by Guo, Lee and Gardner, NIMA 531 (2004) 520-529.
- All orders of pile-up are taken into account.
- It allows the use of the pulse shape of the detector introduced by the user by means of an external text file.
- MCPPU automatically identify the dead time of the counting system to use in the pile-up recovery.
- MCPPU presents an user-friendly graphical interface.

L. Sabbatucci, V. Scot, J. E. Fernandez: Multi-shape pulse pile-up correction: the MCPPU code. Radiat. Phys. Chem. 104 (2014) 372–375.



### Deterministic approach: code DRPPU

- It was derived a balance equation based on first principles assuming:
  - 1. an ideal MCA of infinitesimal energy bins width,
  - 2. only second order PPU,
  - 3. a rectangular pulse shape.

$$y(E) = [1 - 2\lambda\tau \exp(-\lambda\tau)]h(E) + N_t \lambda\tau \exp(-\lambda\tau) \int \overline{h}(E_1)\overline{h}(E - E_1)dE_1$$

The solution of this equation is found using the following iterative strategy:

$$\omega(E) = -ax^{(i)}(E) + b\int x^{(i)}(E_1)x^{(i-1)}(E - E_1)dE_1$$
$$x^{(0)}(E) = \omega(E)$$

Where we define:

 $\omega(E) = \overline{y}(E) / N_t$ a = 1 - 2b

 $b = \lambda \tau \exp(-\lambda \tau)$   $N_t = \frac{\int y(E)dE}{1-E}$ 

L. Sabbatucci, J.E. Fernandez: First principles Pulse Pile-Up balance equation and fast deterministic solution, Radiat. Phys. Chem. 137 (2017) 12-17

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#### Pb sample





#### Outline

- Detector modification:
  - Pulse pile-up.
  - Detector response function.



Simulation of the detector response function with the code MCSHAPE

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#### A modeling tool for detector resolution and incomplete charge collection<sup>†</sup>

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J.E. Fernandez, Viviana Scot: Simulation of the detector response function with the code MCSHAPE. Radiat. Phys. Chem. 78,(2009) 882-887

J.E. Fernandez, Viviana Scot, L. Sabbatucci: A modeling tool for detector resolution and incomplete charge collection. X-ray Spectrometry 44 (2015) 177-182



### **Detector Response**



(Detector influence on radiation measures) The measured spectrum is given by the following convolution product:

$$I_{measured}(E) = \int R(E', E) \phi(E') I(E') dE'$$
  
nere

Where

R(E',E)is the response function $\phi(E')$ is the detector efficiencyI(E')is the original spectrum



### Model of detector response

$$R(E_0, E) = \int Q(E'', E_0) G(E'', E) dE''$$

 $Q(E'', E_0)$  is the energy deposition spectrum

G(E'', E) is the detector resolution

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## Energy deposition spectrum

- Is built by computing the escape spectrum distribution
- In a first approximation its integral is normalized (really is not because of the Rayleigh scattering)
- It can be calculated using a MC code

KA MCSHAPE v. 2.61		
Calculation type         Transport in the target         Simulation         Number of hystories         Number of collisions         (MAX 100)         Output energy resolution         E max [keV]         Output energy resolution         E max [keV]         Channel width         0.10000         View run.log         view run.log         View run.log         View run.log         StART	sw sw sw sw sw sw sw sw sw sw sw sw sw s	MCSHAPE Computes the energy deposition spectrum



$$I_{measured}(E) = \int R(E', E) \varphi(E') I(E') dE'$$
  
=  $\int \left( \int Q(E'', E') G(E'', E) dE'' \right) \varphi(E') I(E') dE'$   
=  $\int \left( \int Q(E'', E') \varphi(E') I(E') dE' \right) G(E'', E) dE''$   
=  $\int \left( \int Q(E'', E') \varphi(E') I(E') dE' \right) G(E'', E) dE''$   
=  $\int \left( \int Q(E'', E') \varphi(E') I(E') dE' \right) G(E'', E) dE''$   
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=  $\int \left( \int Q(E'', E') \varphi(E') I(E') dE' \right) G(E'', E) dE''$   
=  $\int \left( \int Q(E'', E') \varphi(E') I(E') dE' \right) dE' \right) G(E'', E) dE''$ 

J..E. Fernandez, V. Scot : Simulation of the detector response function with the code MCSHAPE, Radiation Physics and Chemistry 78 (2009) 882–887.

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 In a first approximation it is described by a normalized Gaussian

$$G(E_0, E) = \frac{0.9395}{FWHM(E_0)} \exp\left\{-2.773 \frac{(E_0 - E)^2}{FWHM^2(E_0)}\right\}$$

• the FWHM is a function of energy



#### **RESOLUTION: CdTe response function**

The new version of RESOLUTION allows also, for a solid state detector, to introduce the effect of the incomplete charge collection.



J. E. Fernández, V. Scot, L. Sabbatucci: A modeling tool for detector resolution and incomplete charge collection, X-ray Spectrom. 44 (2015)177-182.



## Comparison with experimental data (synchrotron experiment)



- Sample: Cu
- Energy: 40 keV
- Linearly polarized source with polarization degree P= 0.885
- Scattering angle: 90°



[1] Fernandez, J.E. and Scot, V. (2009), Simulation of the detector response function with the code MCSHAPE. Rad. Phys. Chem. 78(10):882-887



## Comparison with experimental data (synchrotron experiment)

Comparison between the experimental data and the simulation of the transport in the target performed with the MCSHAPE<sup>[1]</sup> with and without bremsstrahlung contribution



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#### Thank you for your attention!

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