

Physical interpretation and asymptotic analysis of particle's energy losses on the polarization radiation process

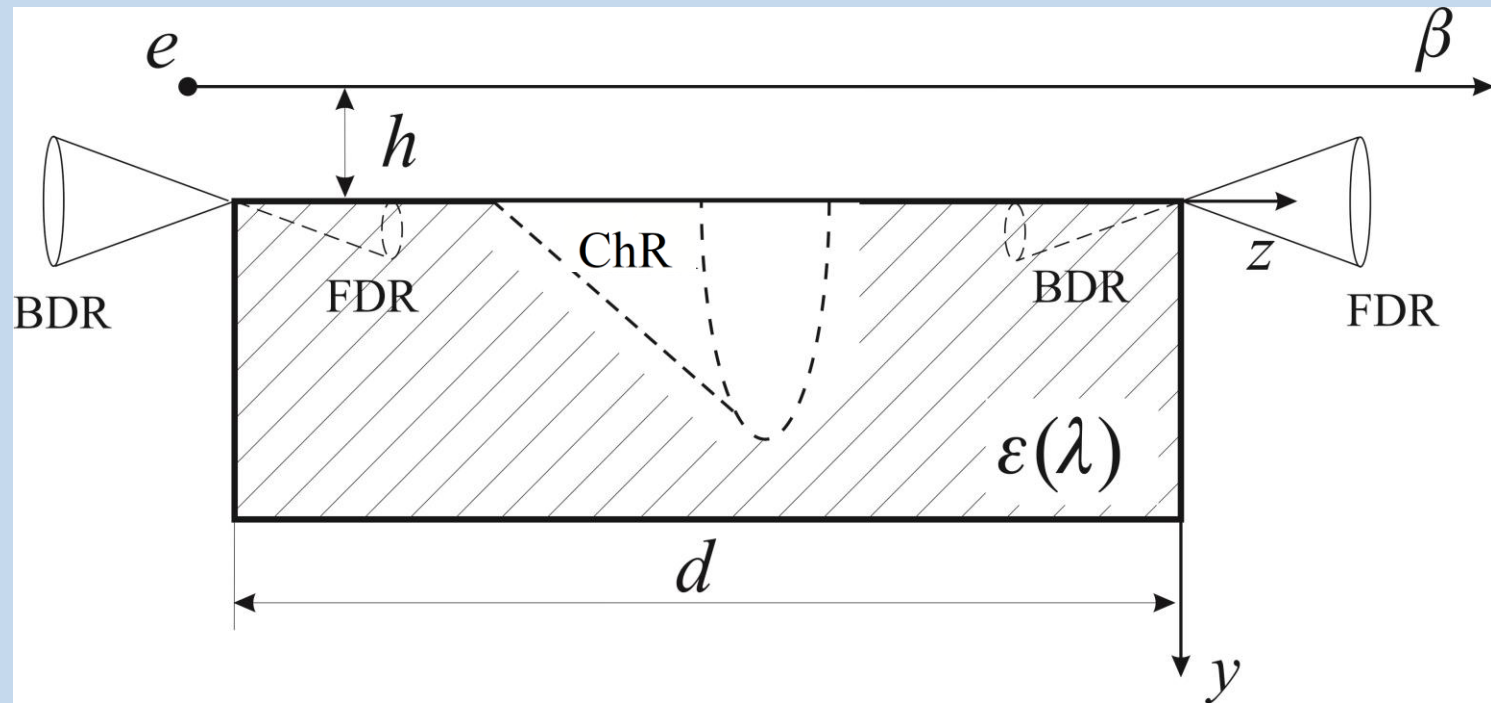
Anatoly S. Konkov

Channeling - 24 Sep. 2018

Motivation

- Beam instrumentation and beam dynamic purposes
- Novel acceleration techniques using dielectric material
- Sources of electromagnetic radiation

Geometry



D.V. Karlovets and A.P. Potylitsyn, JETP Lett. **90** (2009) 326.

M. Shevelev, A. Konkov, A. Aryshev, Phys. Rev. A. **92** (2015) 053851.

Approximations

- Non-polar, nondispersive mediums with $\epsilon > 2$
- Ultrarelativistic particle's energy
- Small radiation angles

Main equations

$$W_{DR} = \frac{3 e^2}{8 h} \gamma \left[\frac{\sqrt{\varepsilon} \pm 1}{\sqrt{\varepsilon} \mp 1} \right]^2$$

Forward DR – upper sign
Backward DR – lower sign

$$W_{ChR} = \frac{e^2}{32\sqrt{2}h^2} \gamma d \frac{\sqrt{\varepsilon - 1}}{\varepsilon}$$

Non-stationary ChR

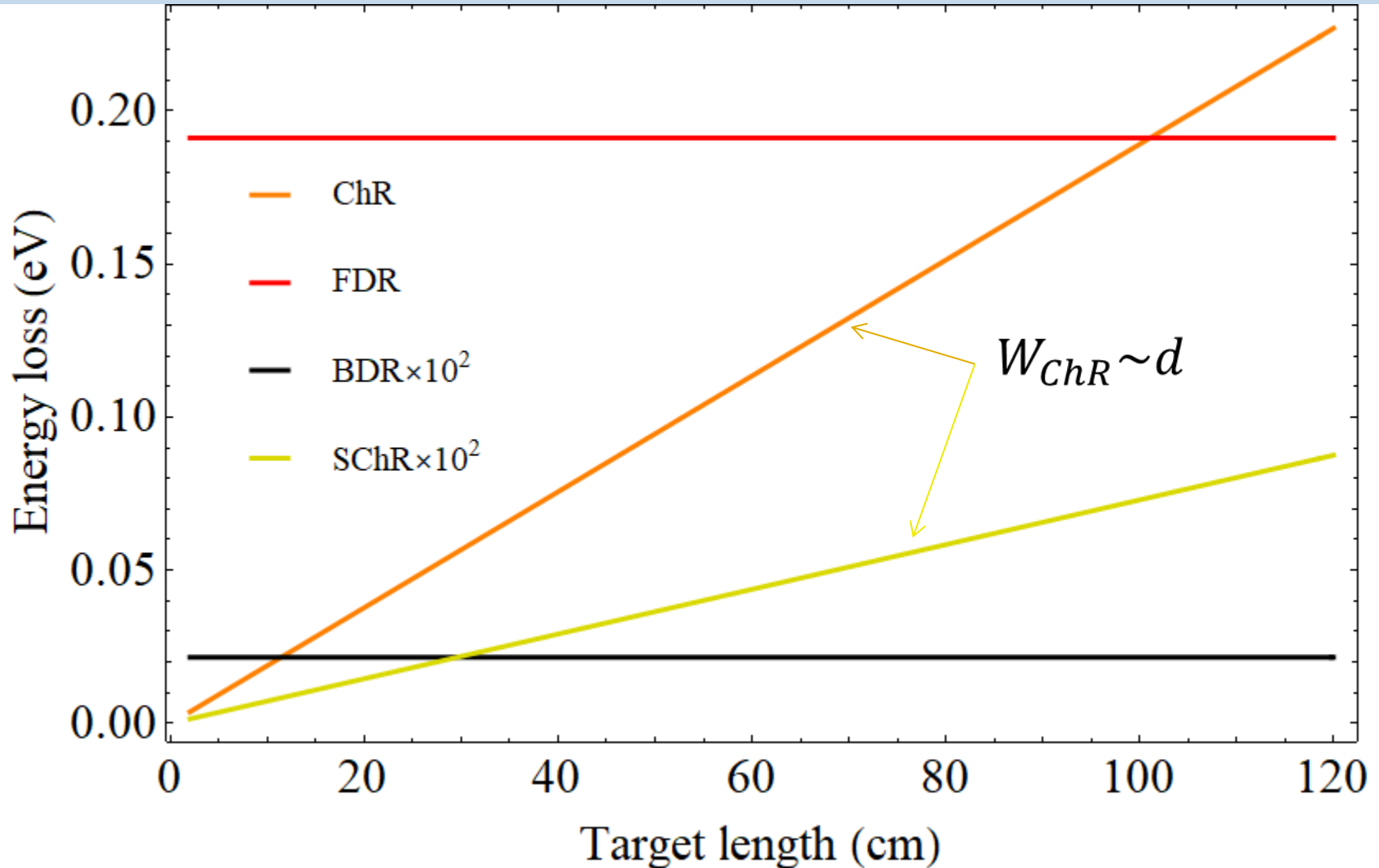
Stationary ChR [A.I. Morozov, JETP **32** (1957) 1260]:

$$W_{SChR} = \frac{e^2}{2\beta h^2} \frac{z}{\varepsilon - 1} \left[\frac{\varepsilon}{\sqrt{1 + \gamma^{-2}\varepsilon}} - \gamma^{-1}\sqrt{\varepsilon - 1} - \beta \right] \approx \frac{e^2}{2h^2} z$$

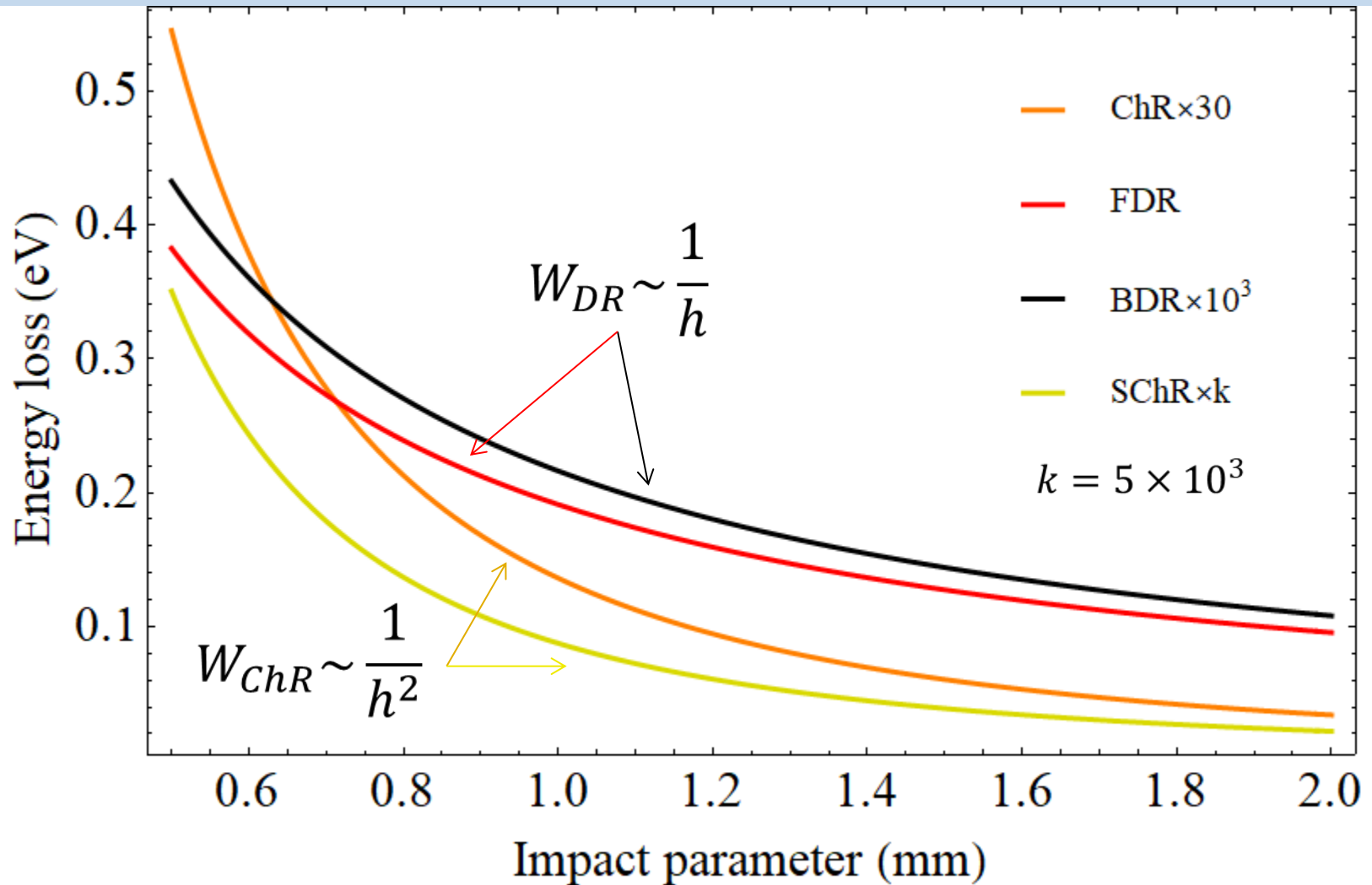
$$\gamma = \sqrt{1 - \beta^2}$$

Particle's Lorentz factor

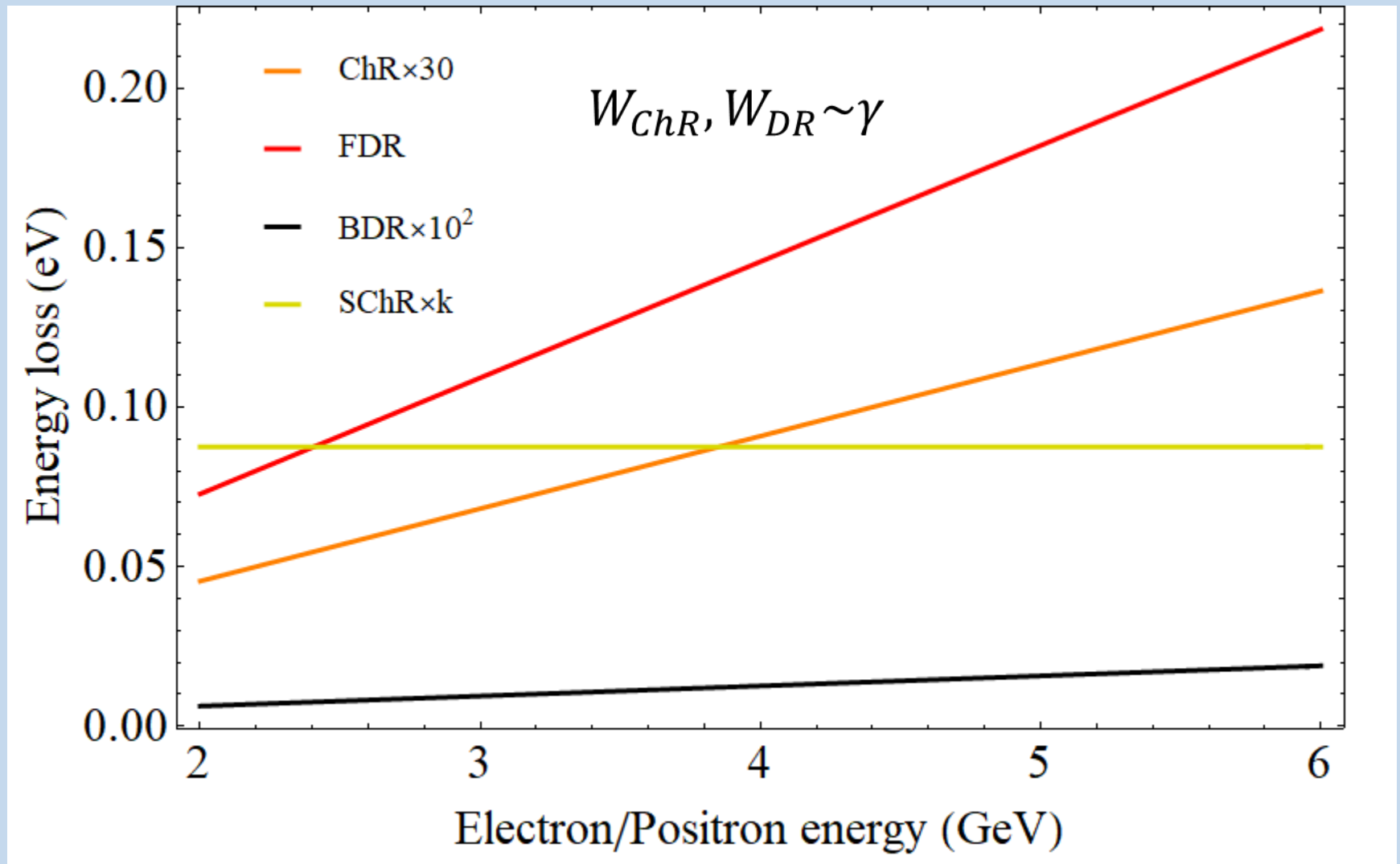
Target length influence



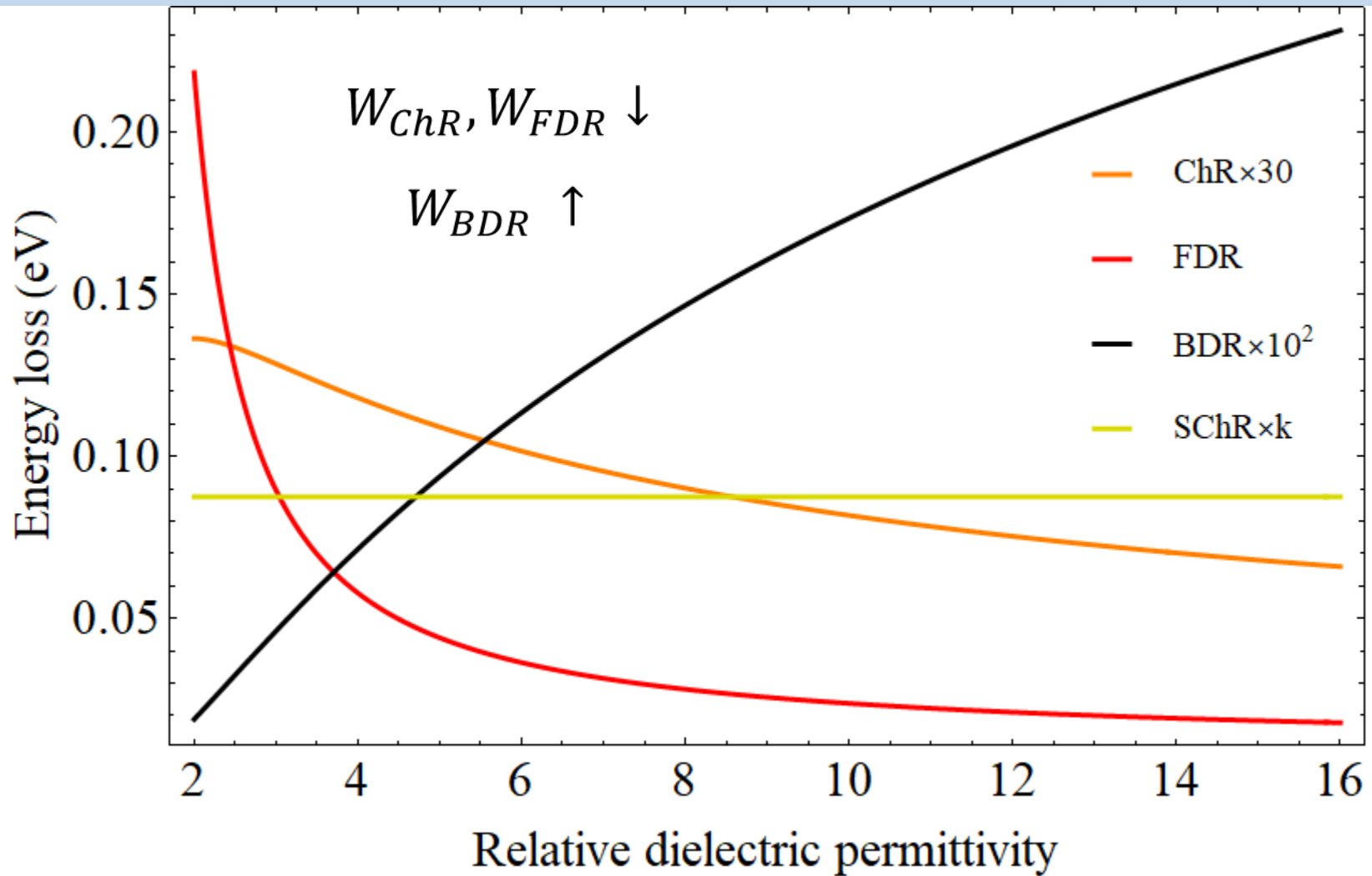
Impact parameter influence



Particle's energy influence



Dielectric properties influence



Conclusion

The devil is in the detail

- Stationary ChR \neq ChR from the finite size target
- The borders are the cause of the field renormalization effect



Thank you for your attention!

Dr. Anatoly S. Konkov
E-mail: Ekwinus@tpu.ru