

Monitoring of the plasma generated by a gas-puff target source dedicated for SXR/EUV microscopy

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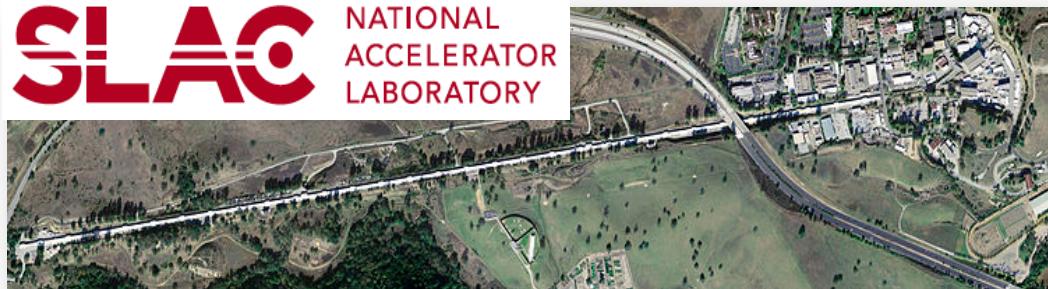
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Outlines

- Motivations on development of laboratory-based SXR and EUV compact sources based on double stream gas-puff target (DSGPT)
- Feasibility of DSGPT for nanoimaging experiments and set-up descriptions
- Plasma generated by DSGPT: diagnostic by Si and SiC detectors (linearity, time evolution, signal intensity)
- Applications of DSGPT
- Conclusions

Large Facilities



Stanford Linear Accelerator (USA)
Source length: 3 km



TOMCAT - X02DA (Switzerland)
Circumference: 288 m

Advantages

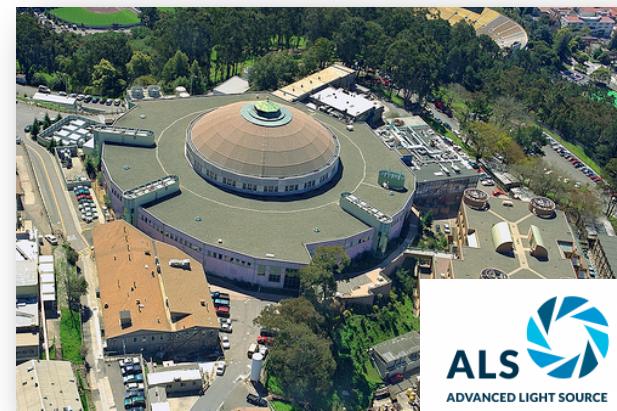
- Bright Sources
- High number of photons
- Tunable

Drawbacks

- Limited Access
- High complexity & Costs
- Expensive maintenance



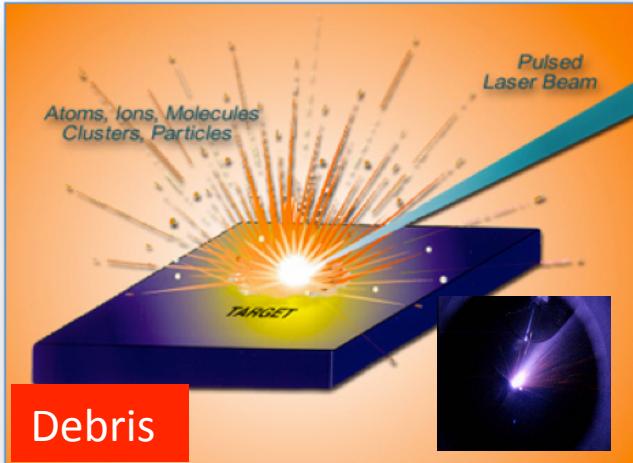
Bessy II
(Berlin, Germany)
Circumference: 240 m



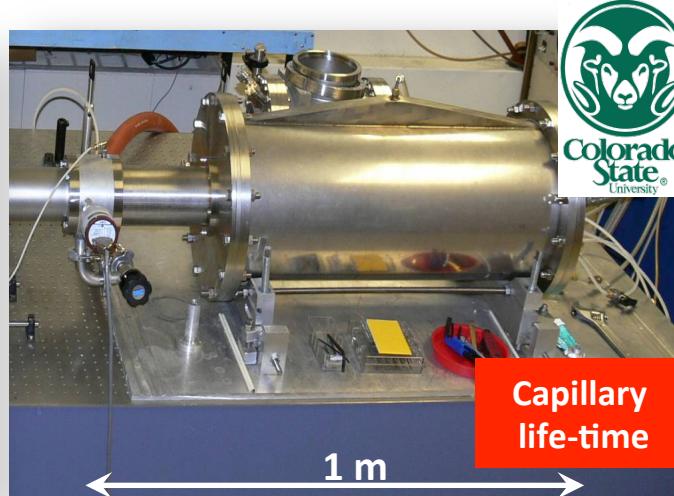
Advanced Light Source
(Berkeley, CA, USA)
Circumference: 200 m

Compact laser-plasma sources

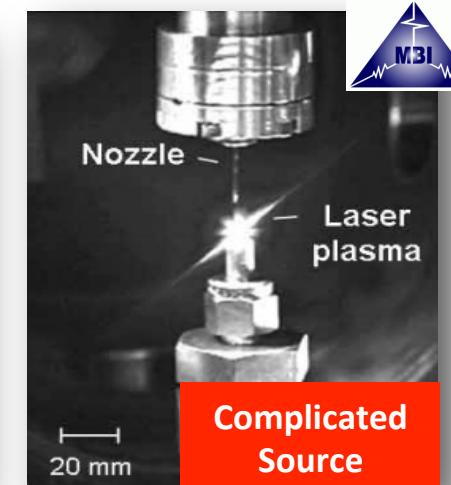
Solid target



Capillary discharge



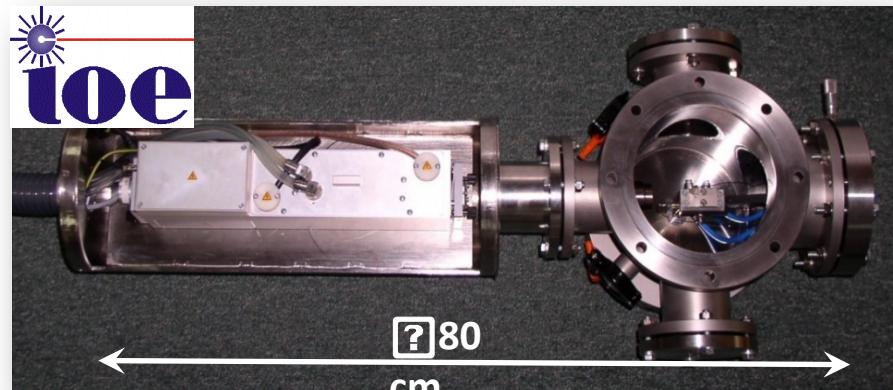
Liquid – Jet Target



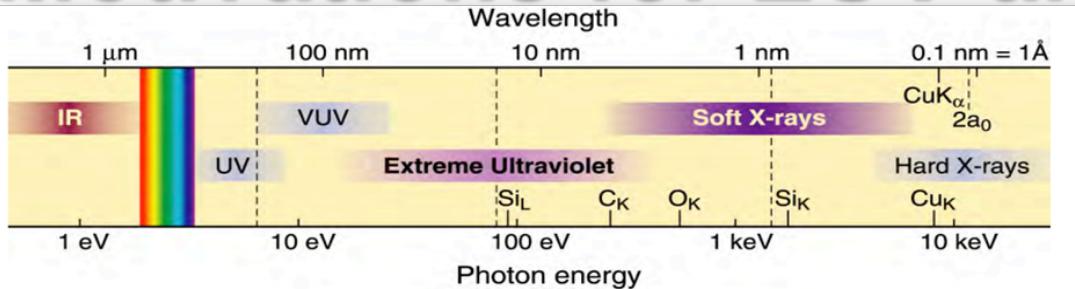
Advantages

- Easily accessible
- User Friendly
- Low cost of operation
- Laboratory environment

EUV/SXR lamp for metrology and microscopy



Motivations for EUV and SXR imaging



SXR ($\lambda=0.1\text{-}10\text{nm}$)

- High-contrast biological imaging

good penetration in micrometer-thick specimens

much larger penetration distances for photons than for electrons

Wavelength (nm)

$$\Delta X = \frac{k \lambda}{NA}$$

$NA = n \sin q$

DX = Spatial Resolution

I = Illumination wavelength

NA = Numerical aperture

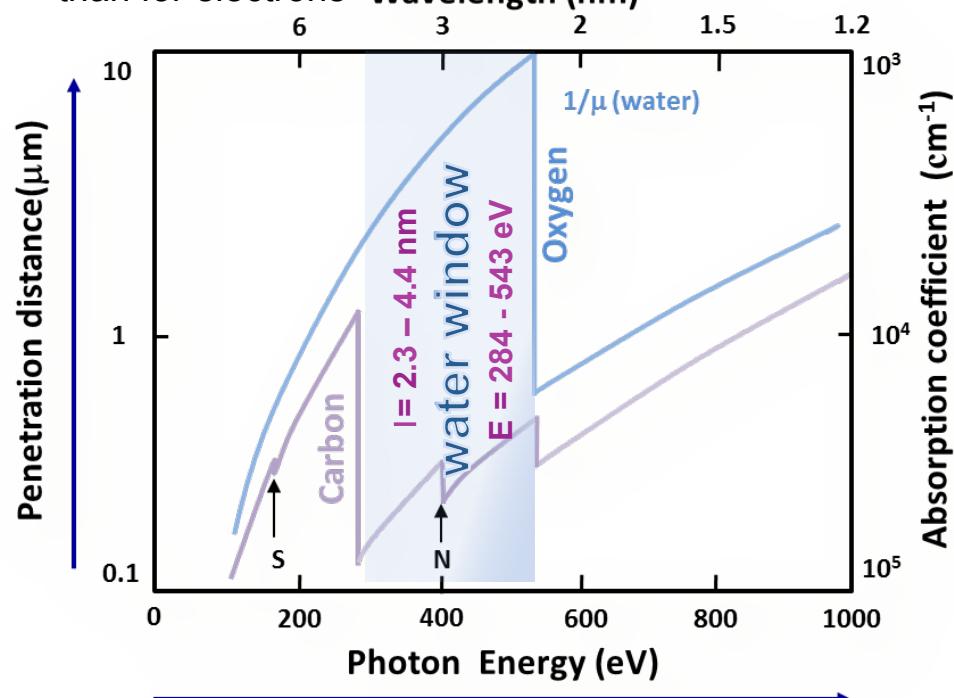
k = 0.61 for incoherent illumination

EUV ($\lambda=10\text{-}120\text{nm}$)

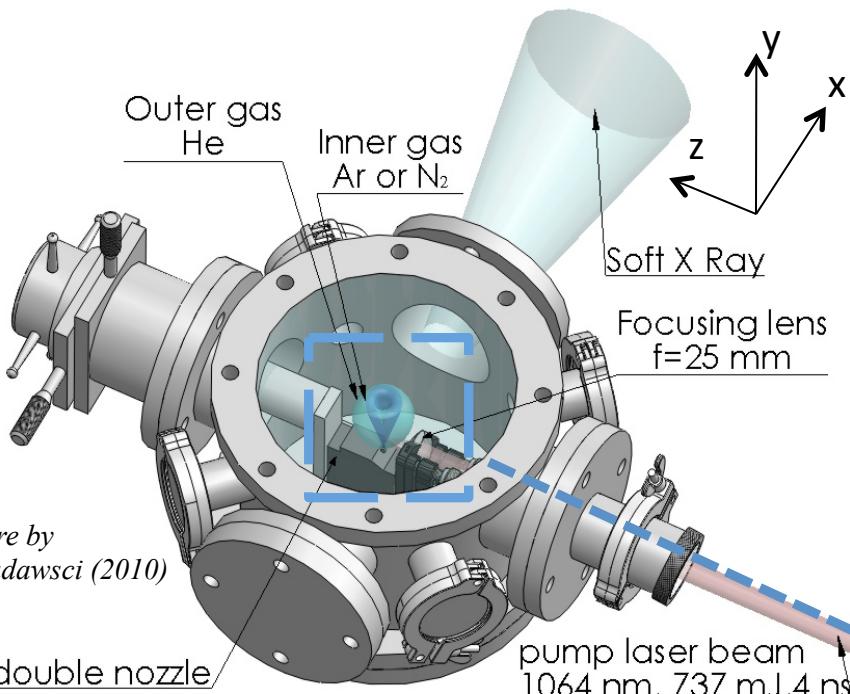
- Optical contrast

due to the atomic resonance frequencies

- High absorption in very thin layers



Gas-puff target EUV and SXR laser-plasma short wavelength source dedicated for imaging



P. W. Wachulak et al., Nucl Instrum Meth B, 268, 10, 1692-1700 (2010)

- Inner nozzle: circular 0.4 mm in diameter
- Outer nozzle: ring 0.7 mm/1.5mm diameters

✓ Compact construction

✓ Debris-free source

✓ High repeatability

✓ Optimal working gas density

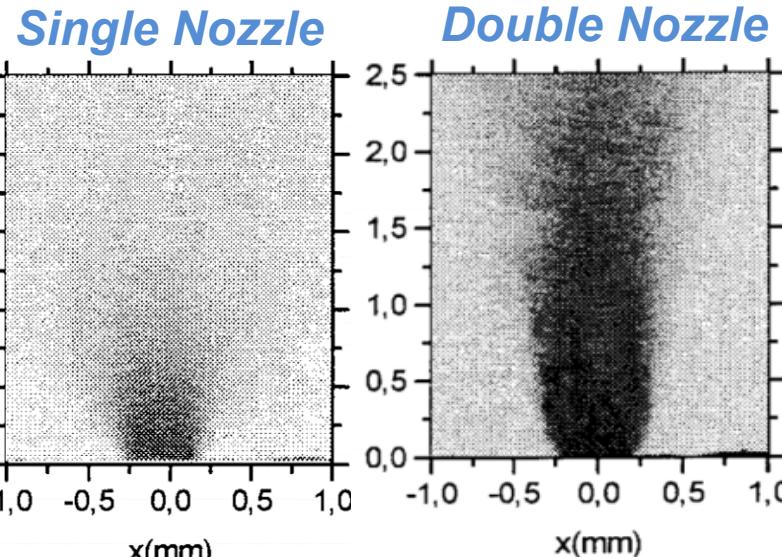
($\sim 10^{18}$ at/cm³)

✓ Energy stability

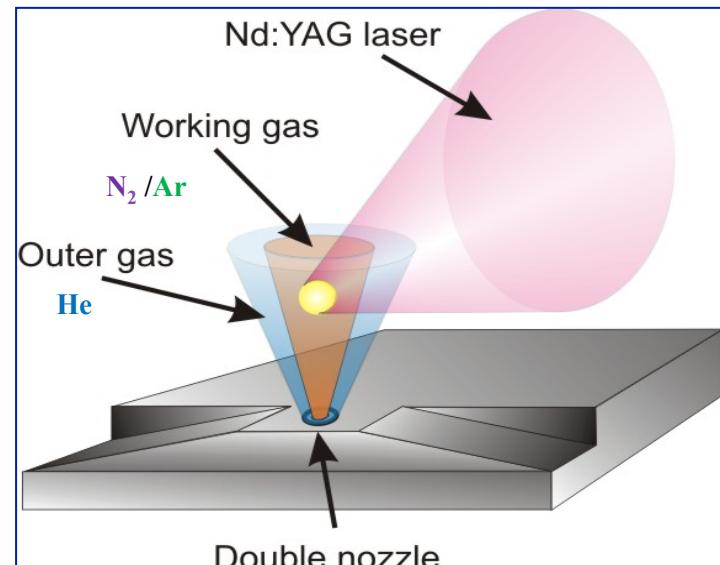
✓ High conversion efficiency

✓ Thousands of shots/day

✓ $n_{\text{plasma}} < n_{\text{critical}}$

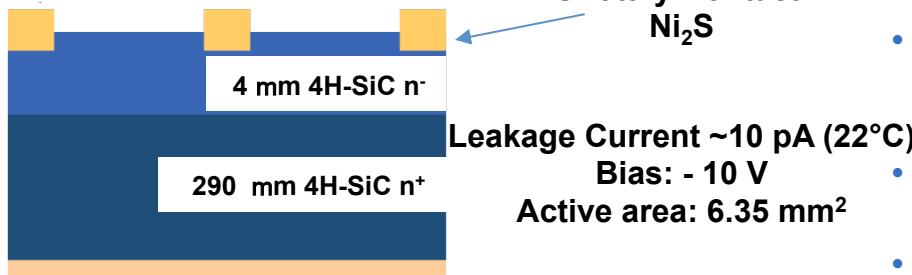
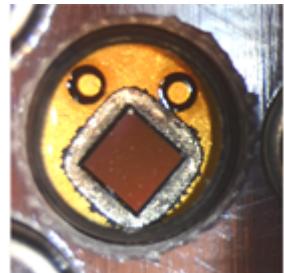


H. Fiedorowicz et al., Appl. Phys. B 70, 305–308 (2000)



Plasma Diagnostic and SiC detectors calibration

$$J_{SiC} = \frac{dq}{dt} = e(n_i E_i / \varepsilon) \mu_{\text{eff}} (V_d/d)$$

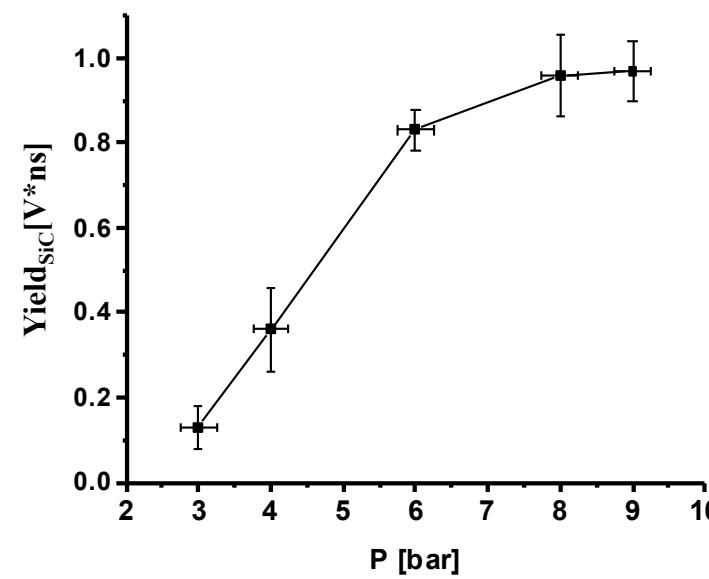
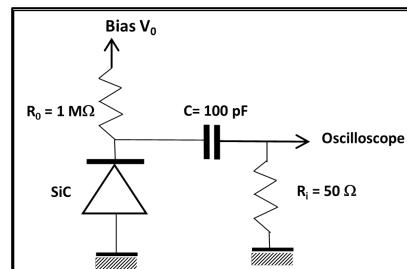
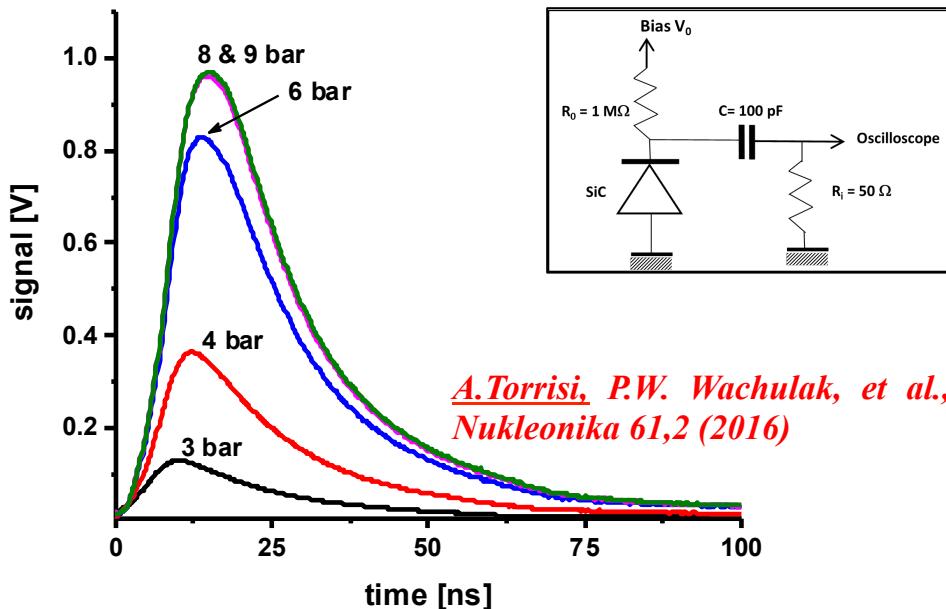


Collaborations with: ST Microelectronics, Catania (Italy)
Messina University (Italy), Dept. of Physics



SiC detectors cut VIS and EUV radiation emitted from plasma, enhancing the sensitivity to very fast ions.

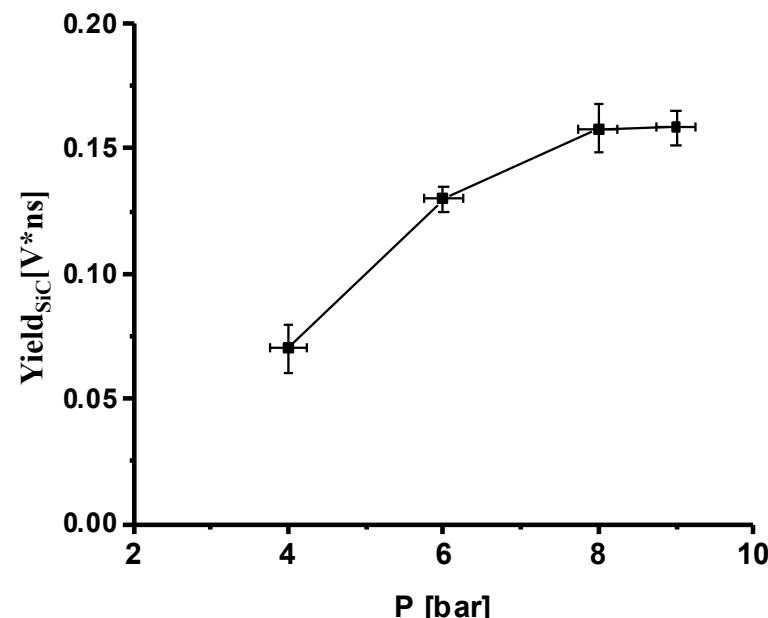
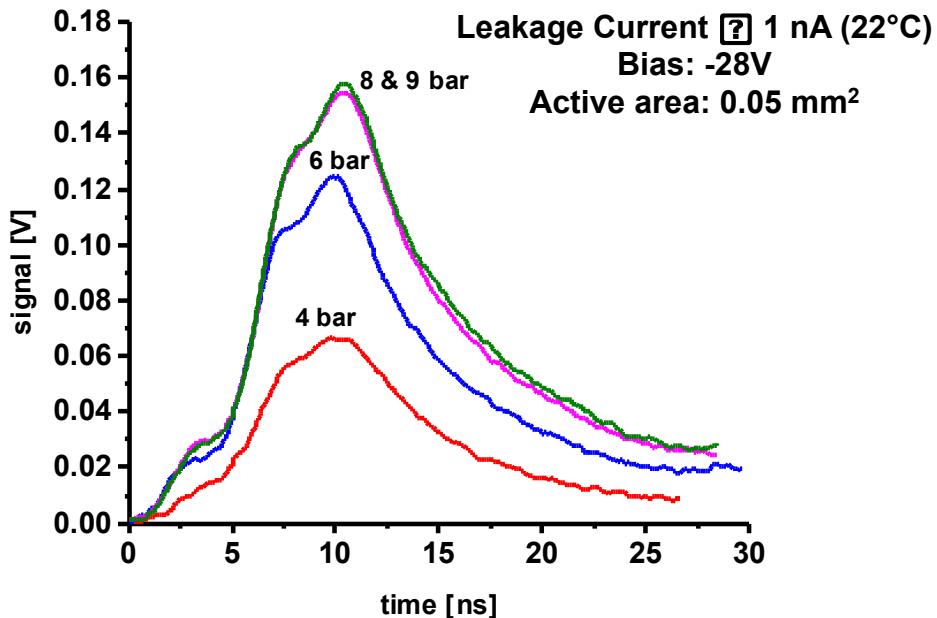
N₂ plasma @ $\lambda = 2.88 \text{ nm}$



Advantages

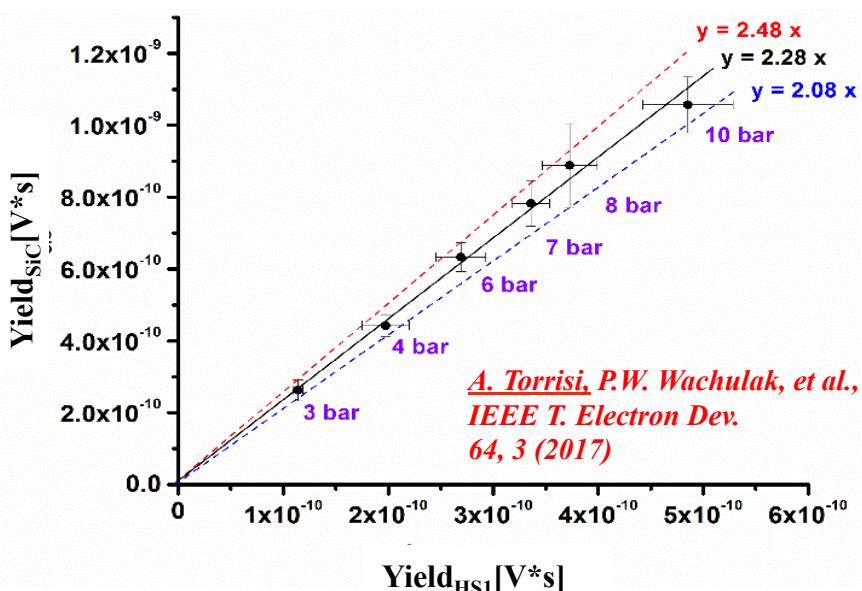
- Higher bandgap [SiC 3.28 eV vs. Si 1.1 eV]
- Noise reduction
- Electrons and fast ions detection
- Low current at room T
- Radiation Hardness
- High electron mobility

Silicon Detector – HS1

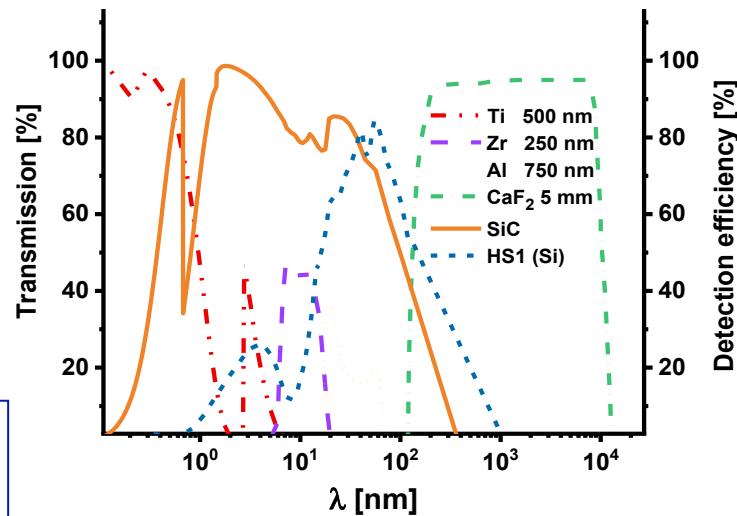
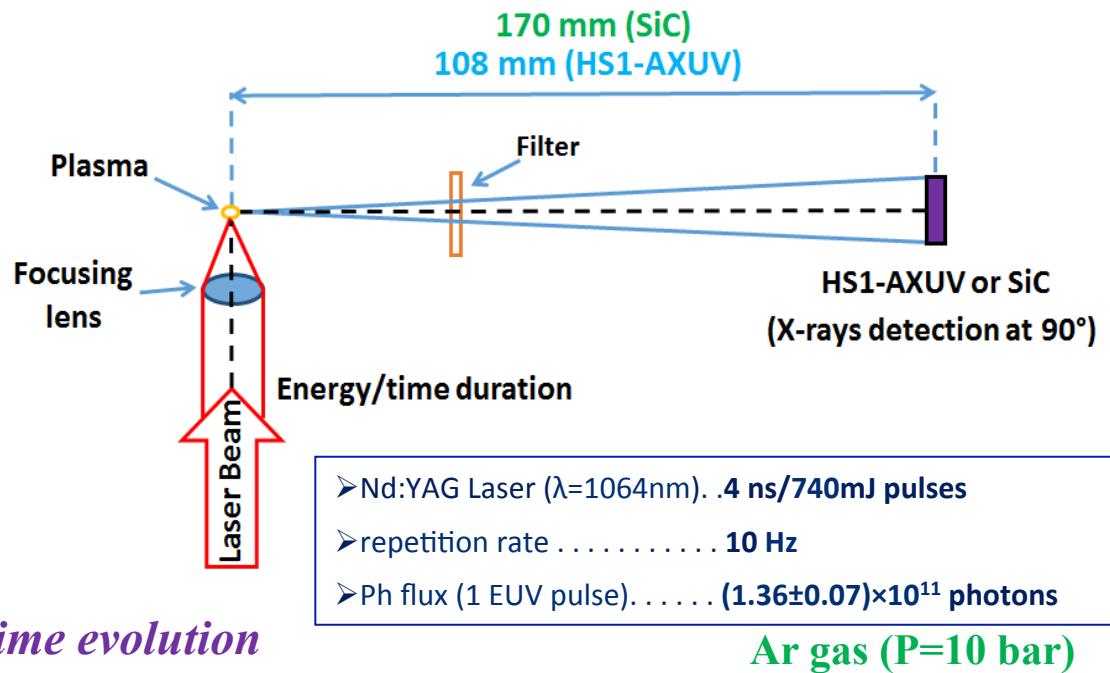


Quantum Efficiency Calculation

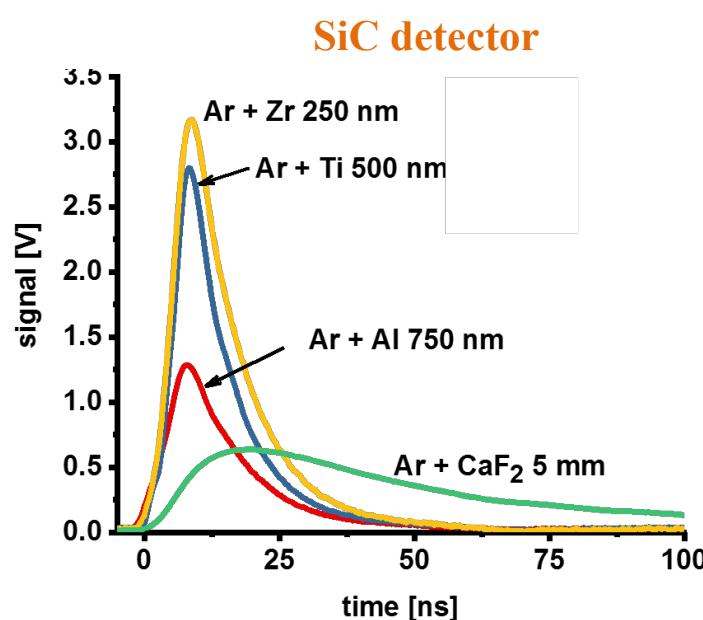
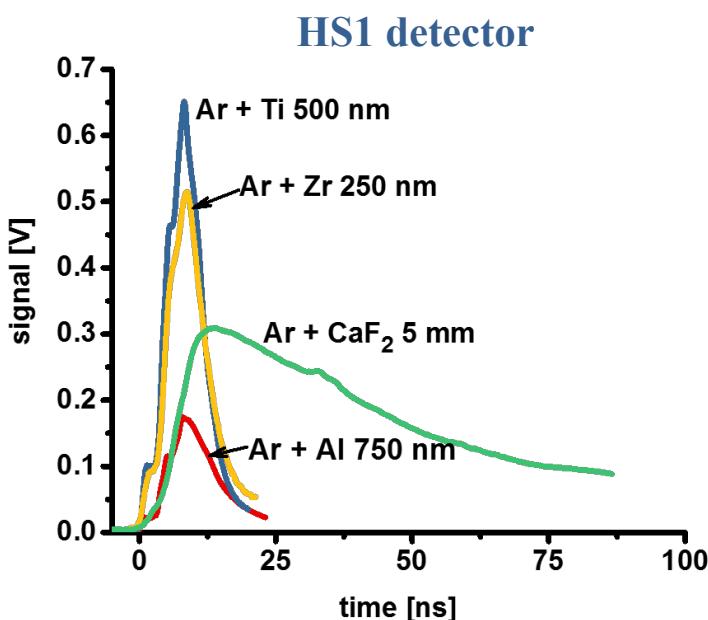
| AXUV/HS1 (manufacturer's data) | | SiC (experimental data) | |
|---|--------------------------|----------------------------|--------------------------|
| QE _λ [e/ph] | R _λ [mA/W] | QE _λ [e/ph] | R _λ [mA/W] |
| $\lambda_{N_2} = 2.88 \text{ nm}$ $E_{N_2} = 430.5 \text{ eV}$ | 119.49 | 277.5 | 3.41 ± 0.13 |
| $\lambda_{Ne} = 1.35 \text{ nm}$ $E_{Ne} = 918.5 \text{ eV}$ | 253 | 275.4 | 7.10 ± 0.16 |



Plasma Diagnostic and SiC detectors calibration



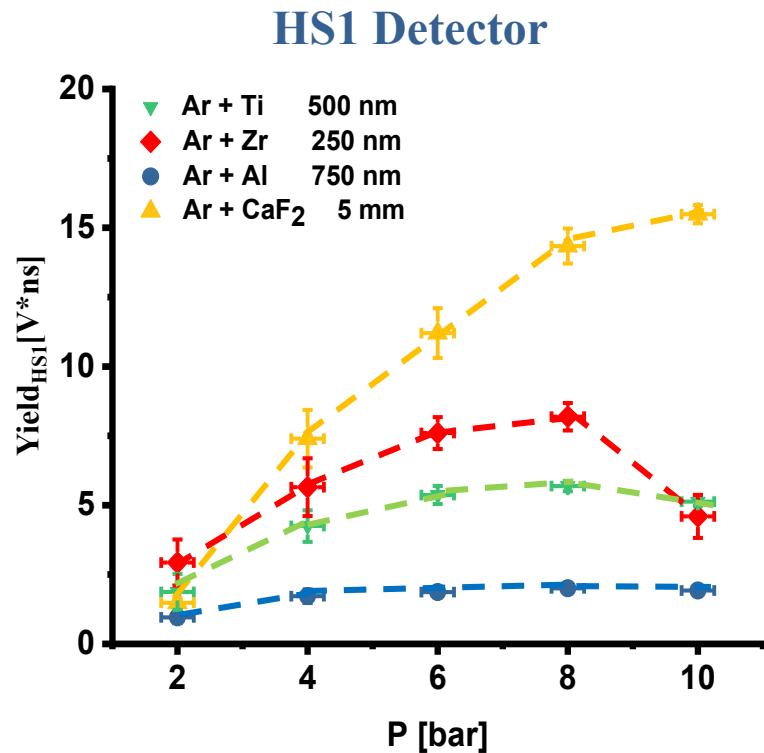
A.Torrisi et al., NIMA (2018) – under review



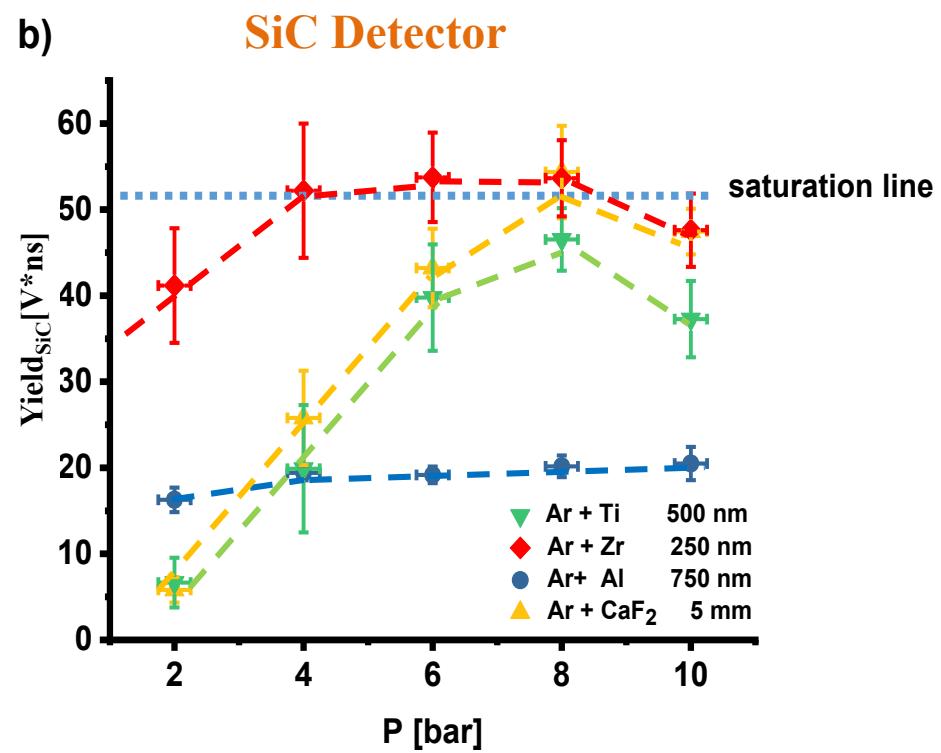
Plasma Diagnostic and SiC detectors calibration

Intensity signal

a)



b)



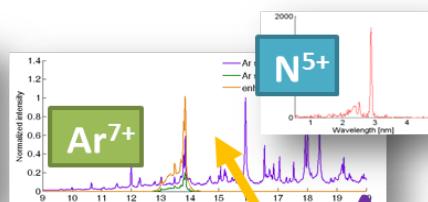
Applications

Double stream gas-puff target laser-plasma EUV/SXR source

Efficient generation of the SXR/EUV high intensity radiation, laser-plasma sources



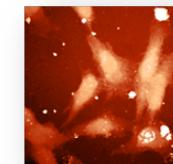
Generation of a monochromatic EUV/SXR radiation from laser-plasma sources



Radiobiology



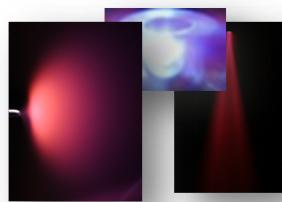
Contact microscopy



SXR/EUV radiation

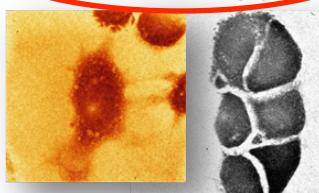
EUV light interactions with gasses

Polymer surface modification using EUV

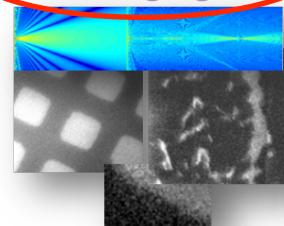


EUV photoionization

SXR "Water window" microscopy



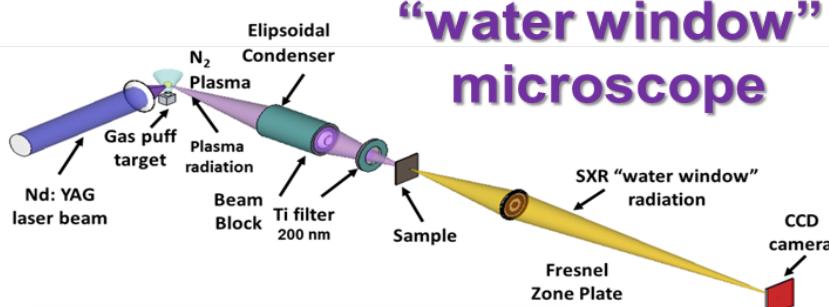
EUV high resolution imaging



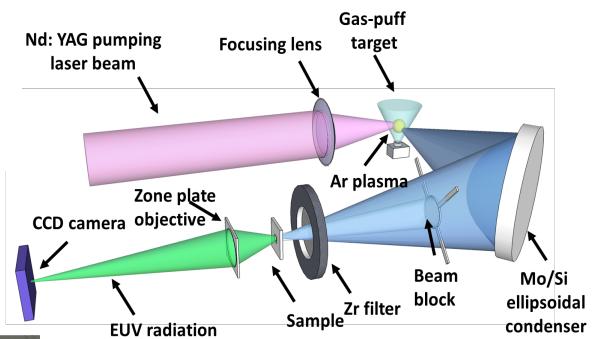
| | |
|-------------------------------|---|
| ⁷ N 14.01 | $\lambda=2.88\text{nm}$ (200nm Ti filter) |
| ¹⁸ Ar 39.95 | $\lambda=2.8\text{-}5\text{nm}$ (200nm Ti filter) |
| ¹⁸ Ar 39.95 | $\lambda=2.7\text{-}4.3\text{nm}$ (200nm Si ₃ N ₄ 200nm Ti) |
| ¹⁸ Ar 39.95 | $\lambda=13.84\text{ nm}$ (250nm Zr filter + Mo/Si) |
| ³⁶ Kr 83.89 | $\lambda=0.7\text{nm}$ (10μm Be, 20μm Si) |
| ⁵⁴ Xe 131.29 | $\lambda\sim0.8\text{-}1.5\text{nm}$ (10μm Al) |

Overview of double stream gas-puff target microscopes

“water window” microscope



EUV microscope



Advantages

- Compactness
- Spatial resolution: ~50 – 60 nm
- Short time acquisition 3 ns – 1 min
- Applications in biology, material sciences and nanotechnology
- Possibility of commercialization

Drawbacks

- Limited number of photons and low rep. Rate compared to synchrotrons
- Possible improvement of SNR

Summary

- The presented results demonstrated that the SiC detectors have sensitivity and performance comparable to Si detectors in the diagnostics of laser-generated plasma emitting SXR up to UV radiation.
- The **low reverse current** (of about three orders of magnitude lower than in silicon), allows to employ the SiC detectors **at room temperature**, giving **higher sensitivity** with respect to silicon detectors.
- The higher energy gap of 3.3 eV in SiC, with respect to the 1.1 eV of Si, gives to the detector **insensitivity to the visible light** produced by the laser-generated plasma so that it **does not need filters** to reduce the spurious VIS radiation arriving on the detector.
- The **higher displacement energy** of the SiC (25 eV) with respect to the Si-Si crystalline structure (15 eV), permits to **reduce significantly the crystal damage** to the detector under high radiation doses and high deposited energies.
- The higher melting point and effective atomic number of the SiC, with respect to the Si, allows to use the detector with thinner active regions and to have **high efficiency for X-rays, electrons, and ions at high energy**.
- *The obtained results are very promising to be extended to a micro-pixel CCD camera to acquire an X-ray map from applications of SiC detectors to X-ray microscopy or to detect images in the UV region.*



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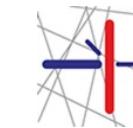
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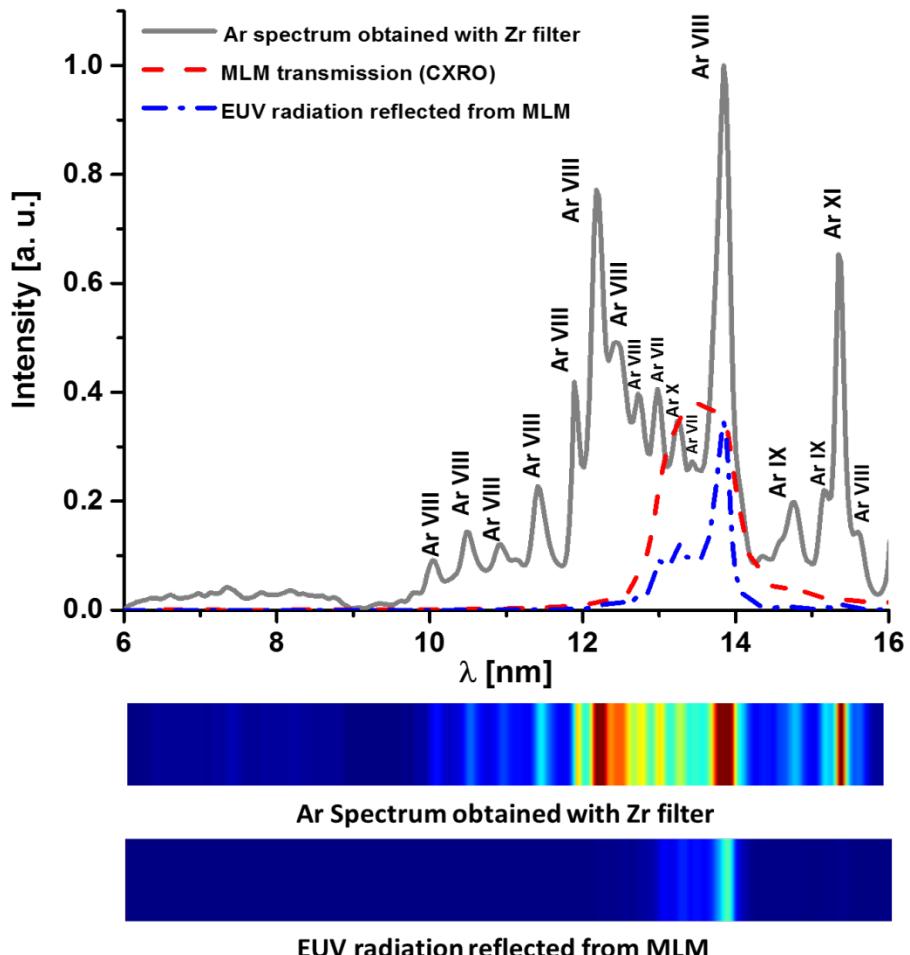
Center of Accelerators
and Nuclear Analytical Methods
(CANAM)



Thank you
for
your attention

Compact EUV microscope based on Ar plasma

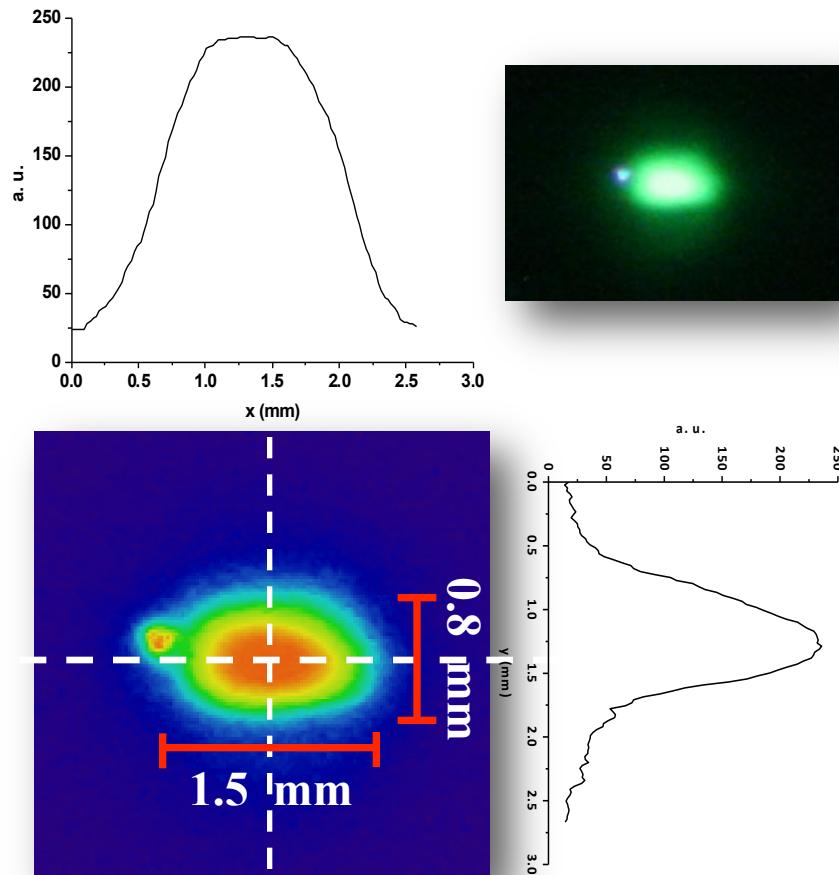
Source Parameters



Inverse rel. bandwidth (FWHM) of emission

$\lambda/\Delta\lambda \sim 60$ @ $\lambda = 13.84$ nm

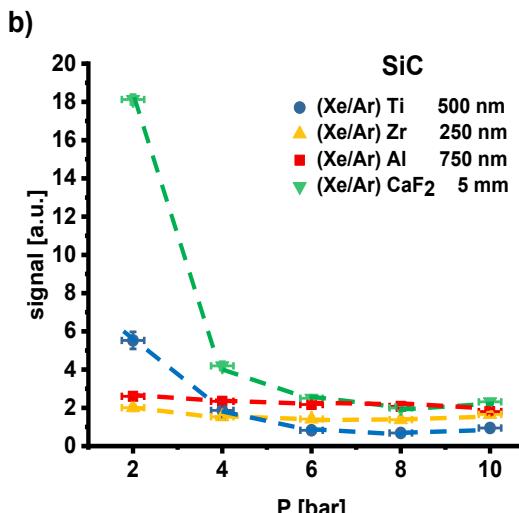
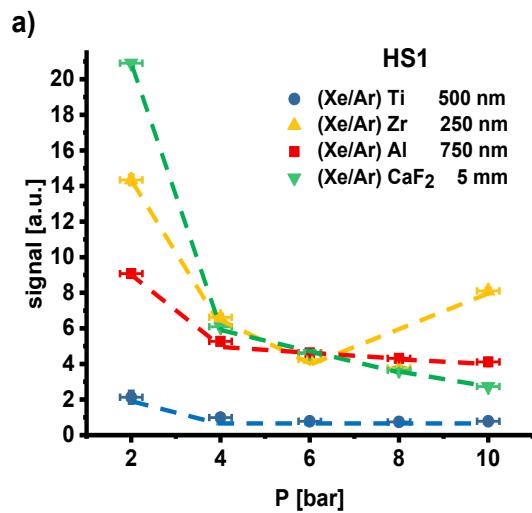
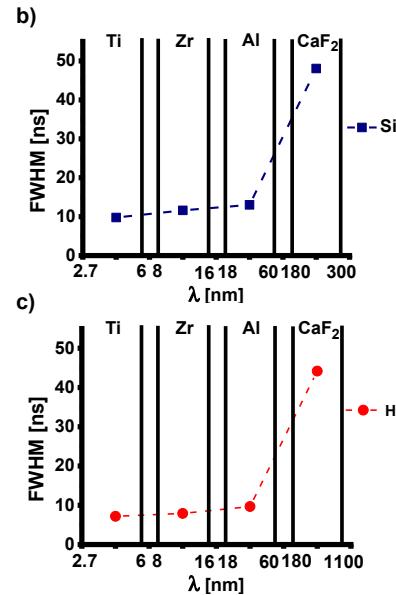
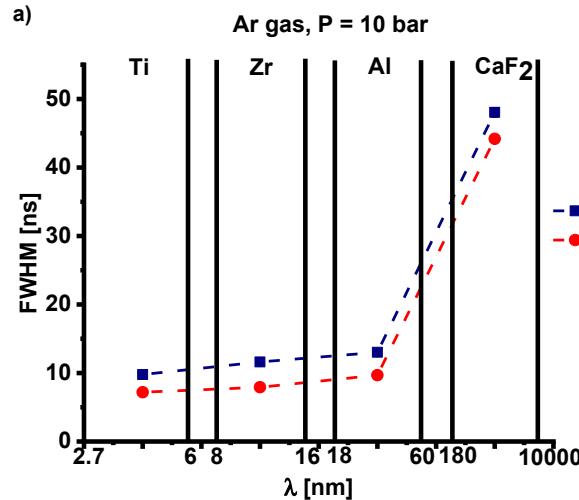
A.Torrisi et al., – Journal of Microscopy, Vol. 265, Issue 2 (2017)



CCD camera

iKon-M (Andor), 1k x 1k pix, 13x13 μm^2

Plasma Diagnostic and SiC detectors calibration



HS1 – Si Detector

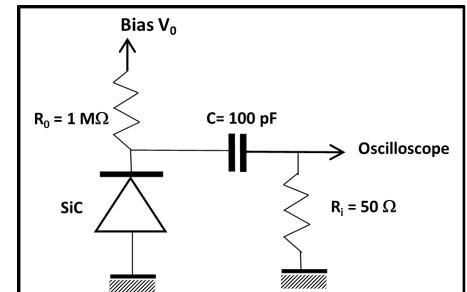
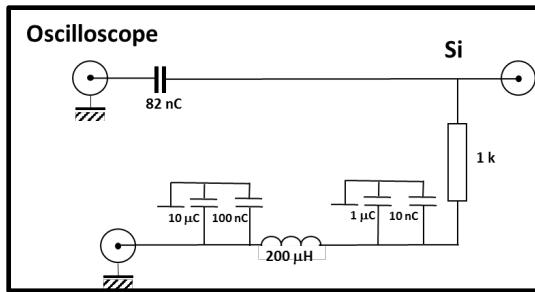
Detector

Electro-optical characteristics of a HS1 - Si detector at 25°C

| PARAMETERS | TEST CONDITIONS | TYPICAL VALUES | UNITS |
|-----------------|---------------------------------|----------------|-----------------|
| Active Area | 0.02 mm x 0.02 mm | 100 | mm ² |
| Responsivity | @2.88 nm, V _R = 0 V | 277.5 | mA/W |
| Reverse Current | I _R = 1 mA | 10 | Volts |
| Capacitance | V _R = 0 V | 10 | nF |
| Rise Time | V _R = 0 V, RL = 50 W | 10 | msec |

Electro-optical characteristics of a SiC detector at 25°C

| PARAMETERS | TEST CONDITIONS | TYPICAL VALUES | UNITS |
|-----------------|-----------------------------------|----------------|-----------------|
| Active Area | 2.52 mm x 2.52 mm | 16 | mm ² |
| Responsivity | @2.88 nm, V _R = - 20 V | 7.92 | mA/W |
| Reverse Current | I _R = 6 pA | 1- 10 | Volts |
| Capacitance | V _R = - 20 V | 100 | pF |
| Rise Time | < 1 | < 1 | nsec |



$$\psi_{\text{ph}} = \frac{\int V(t)dt}{\text{QE}_\lambda \cdot e \cdot R_i}$$

$$R_\lambda = \frac{\text{QE}_\lambda \cdot \lambda \cdot e}{\hbar \cdot c}$$