

# Massive stars as progenitors of merging black hole binaries

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# Introduction

## Open question:

- What are the **formation mechanisms** of **binary black holes**?



## Purpose:

- to study the **demography** of **compact object binaries** in **different environments**.

## Issue:

- most of current **population-synthesis codes** **do not use recent stellar evolution models**.

# MOBSE (Massive Objects in Binary Stellar Evolution)



**Updated version** of the **most popular** and **used population synthesis code** (Hurley+ 2002).

**BSE**: includes **obsolete stellar-evolution models**:

- **Tout+ 1997** for the **stellar winds**;
- **Hurley+ 2000** for the **supernova explosions** (SNe).

**MOBSE**: major updates:

- **recent stellar winds** Vink+ 2001 and Gränerfer+ 2011;
- **new SNe** Fryer+ 2012, **Pulsation-Pair-Instability** (PPISN) and **Pair-Instability** (PISN) Woosley 2017.

# Upgrades: stellar winds

The **main differences** with respect to the **old recipes** for the stellar winds are:

**Dependence on metallicity  $Z$**   
during **Wolf-Rayet** phase and  
**Luminous Blue Variable** stars:

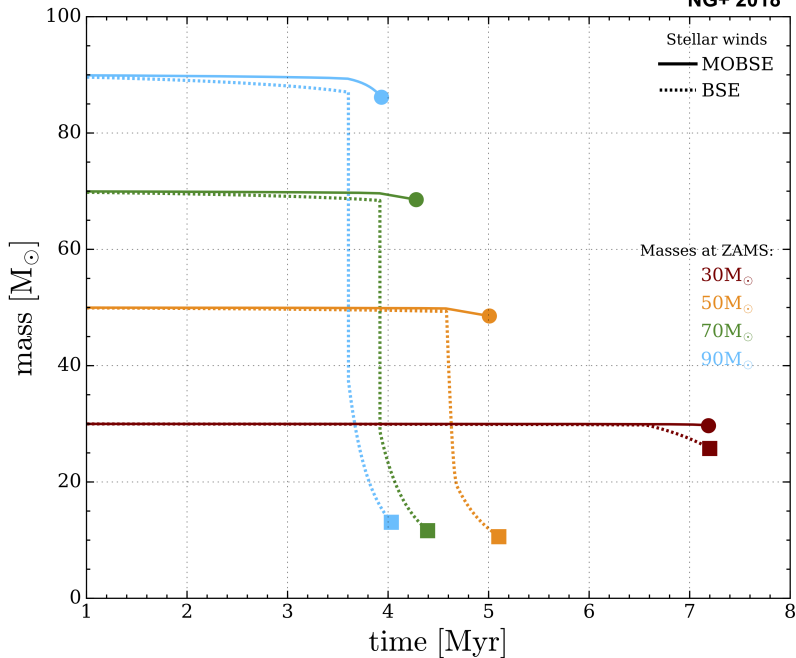


$$\dot{M} \propto Z^{\alpha} M_{\odot} \text{ yr}^{-1}$$

$$\begin{cases} \alpha = 0.85 & \Gamma_e < \frac{2}{3} \\ \alpha = 2.45 - 2.4 \Gamma_e & \frac{2}{3} \leq \Gamma_e \leq 1 \end{cases}$$

Effect of the electron - scattering  
**Eddington factor** on mass loss:  
(Chen+ 2015)





# Upgrades: SNe models

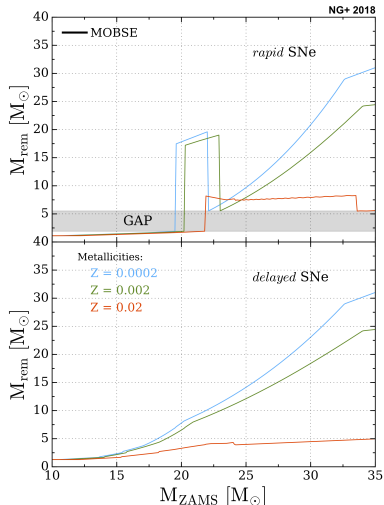
Implementation of two new SNe models described in Fryer+ (2012).

## Rapid SNe:

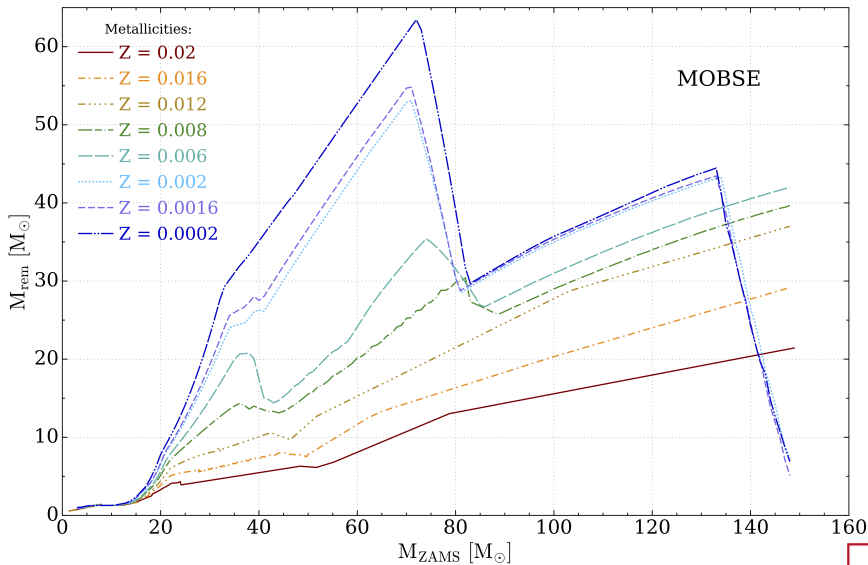
explosion occurs at  $t \lesssim 250$  ms after the bounce.

## Delayed SNe:

explosion occurs after  $t \gtrsim 0.5$  s from the bounce.



# Mass spectrum



# BHBs Demography with MOBSE

## Grid of initial conditions:

$Z \rightarrow$  12 metallicity  $\in [0.02 - 0.0002]$ ;

systems  $\rightarrow 10^7$  for each metallicity;

## Distributions proposed by Sana+ 2012:

$M_1 \rightarrow$  **IMF** of Kroupa+ 2001 in  $M_1 \in [5 - 150]M_{\odot}$ ;

$M_2 \rightarrow$  uniform distribution of  $M_2 \in [0.1 - 1.0]M_1$ ;

$e \rightarrow$  uniform distribution of  $e^{-0.42} \in [0.0 - 1.0]$ ;

$P \rightarrow$  uniform distribution of  $\log_{10}(P/\text{day})^{-0.55} \in [0.15 - 5.5]$ .

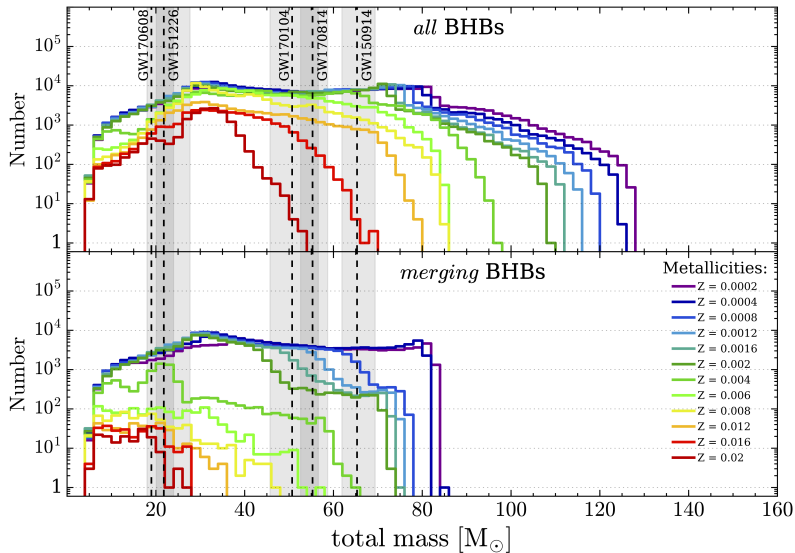
**NG, M. Mapelli & M.Spera, 2018, MNRAS, 474, 2959**





# Mass BHBs

NG+ 2018



# Common-Envelope

CE critical phase for the formation of compact object binaries.

Energy conservation formalism:  $\alpha\lambda$

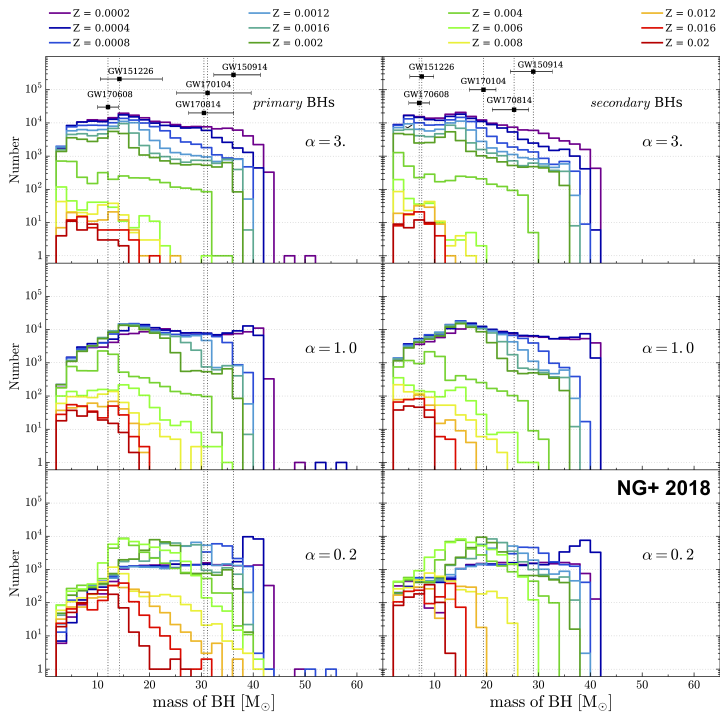
$$\alpha \left( \frac{Gm_1 m_2}{2a_i} - \frac{Gm_{1,\text{core}} m_2}{2a_f} \right) = - \frac{Gm_1 m_{1,\text{env}}}{\lambda R_1}$$

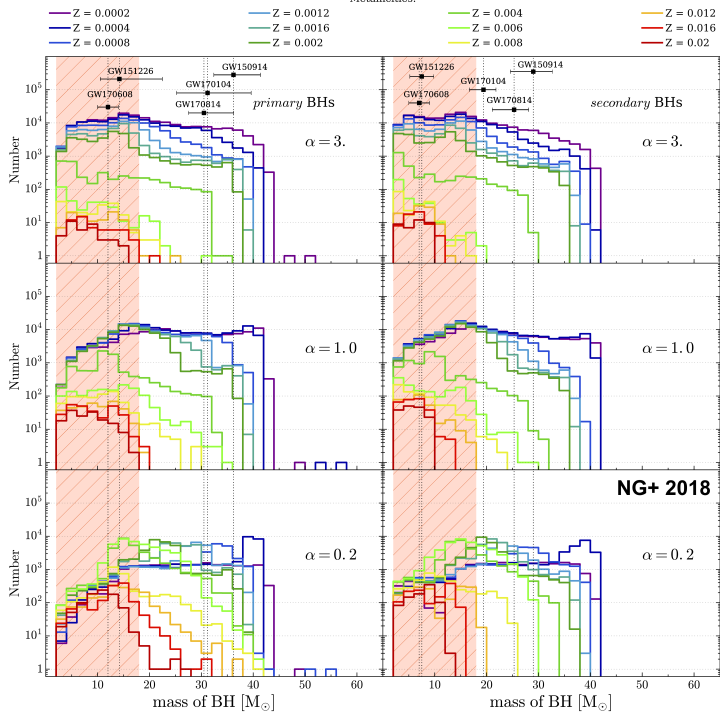


$\alpha \rightarrow$  efficiency transfer of orbital energy to the common envelope.



$\lambda \rightarrow$  describes the binding energy of the common envelope.





# Merger per unit mass

$$R = \frac{N_{\text{merger}}}{M_{\text{tot,sim}}}$$

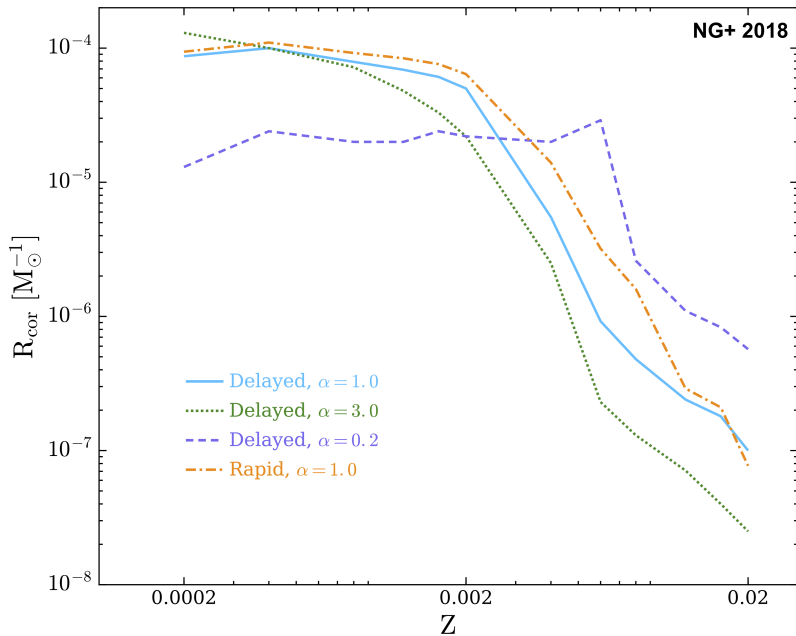


$$R_{\text{cor}} = f_{\text{bin}} f_{\text{IMF}} R$$

$f_{\text{bin}} = 0.5$  → we assume 50 per cent of binary

$f_{\text{IMF}} = 0.285$  → we simulate only  $M_1 \geq 5 M_{\odot}$

# Merger per unit mass



# Conclusions

1. → the **heaviest BHs** ( $\sim 60 M_{\odot}$ ) formed at  $Z \lesssim 0.002$ ;
2. → the most **massive BHBs** ( $\gtrsim 85 M_{\odot}$ ) **do not merge**;
3. → the masses of **our merging BHBs match** those of the five reported **GW events**;
4. → **merging BHBs** form much **more efficiently** from **metal-poor** ( $R_{\text{cor}} \sim 10^{-4} M_{\odot}^{-1}$ ) than from metal-rich ( $R_{\text{cor}} \sim 10^{-7} M_{\odot}^{-1}$ ) binaries.