



# **Gravitational wave as probe of the equation of state of nuclear matter**

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# Main goal

PRL 119, 161101 (2017)

Selected for a **Viewpoint** in *Physics*  
PHYSICAL REVIEW LETTERS

week ending  
20 OCTOBER 2017

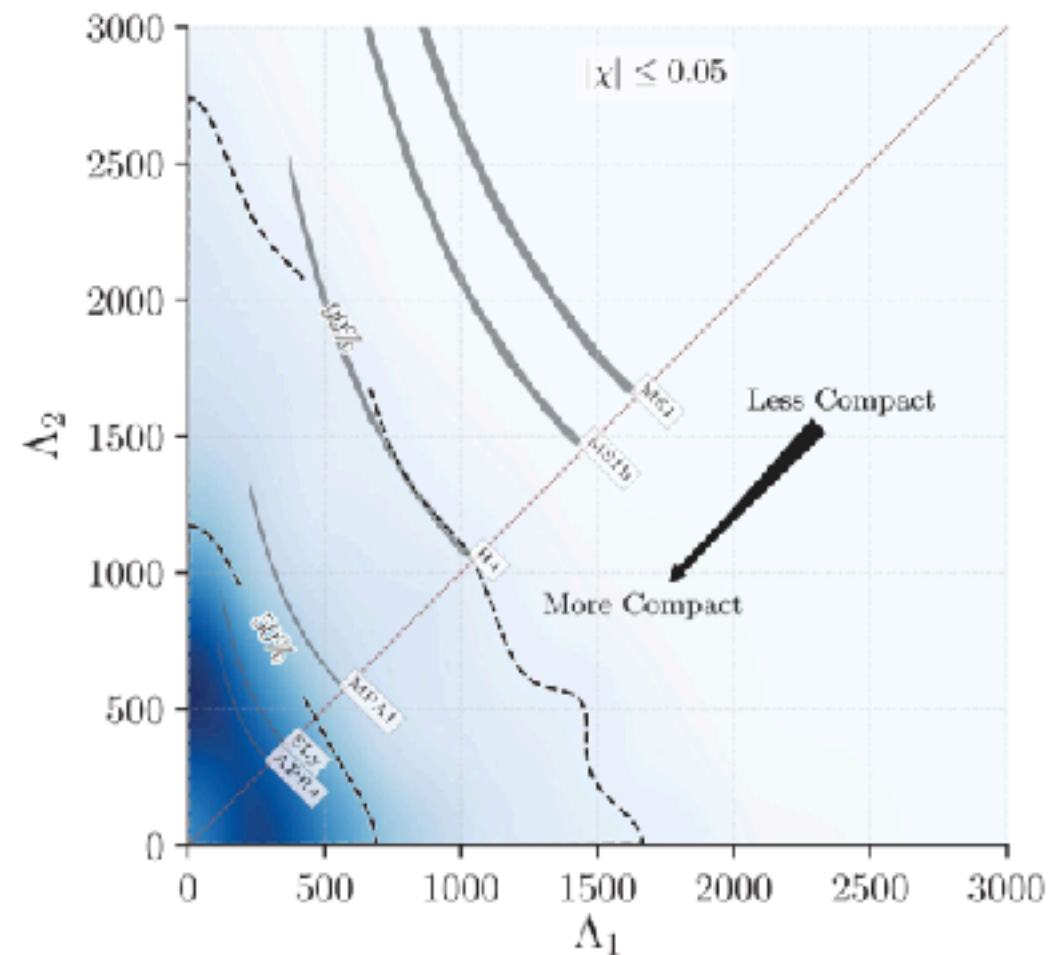
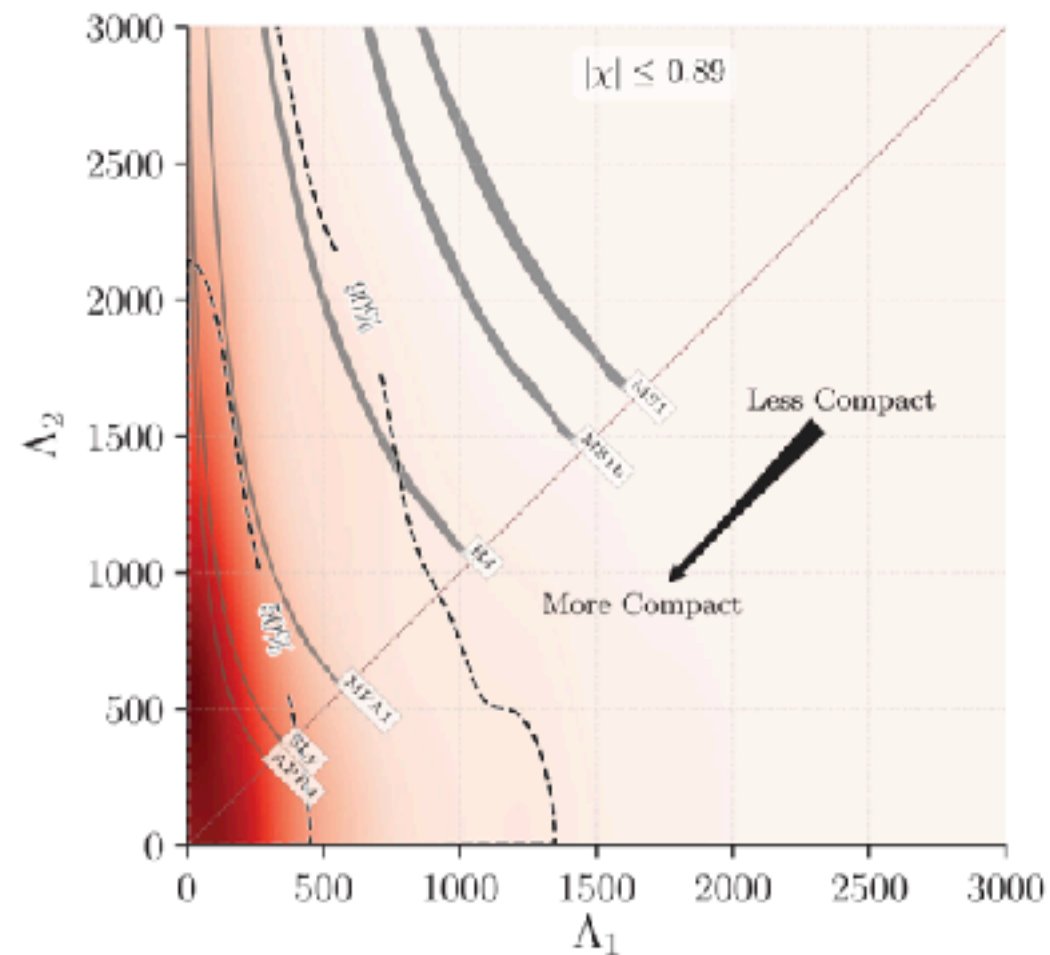


## GW170817: Observation of Gravitational Waves from a Binary Neutron Star Inspiral

B. P. Abbott *et al.*\*

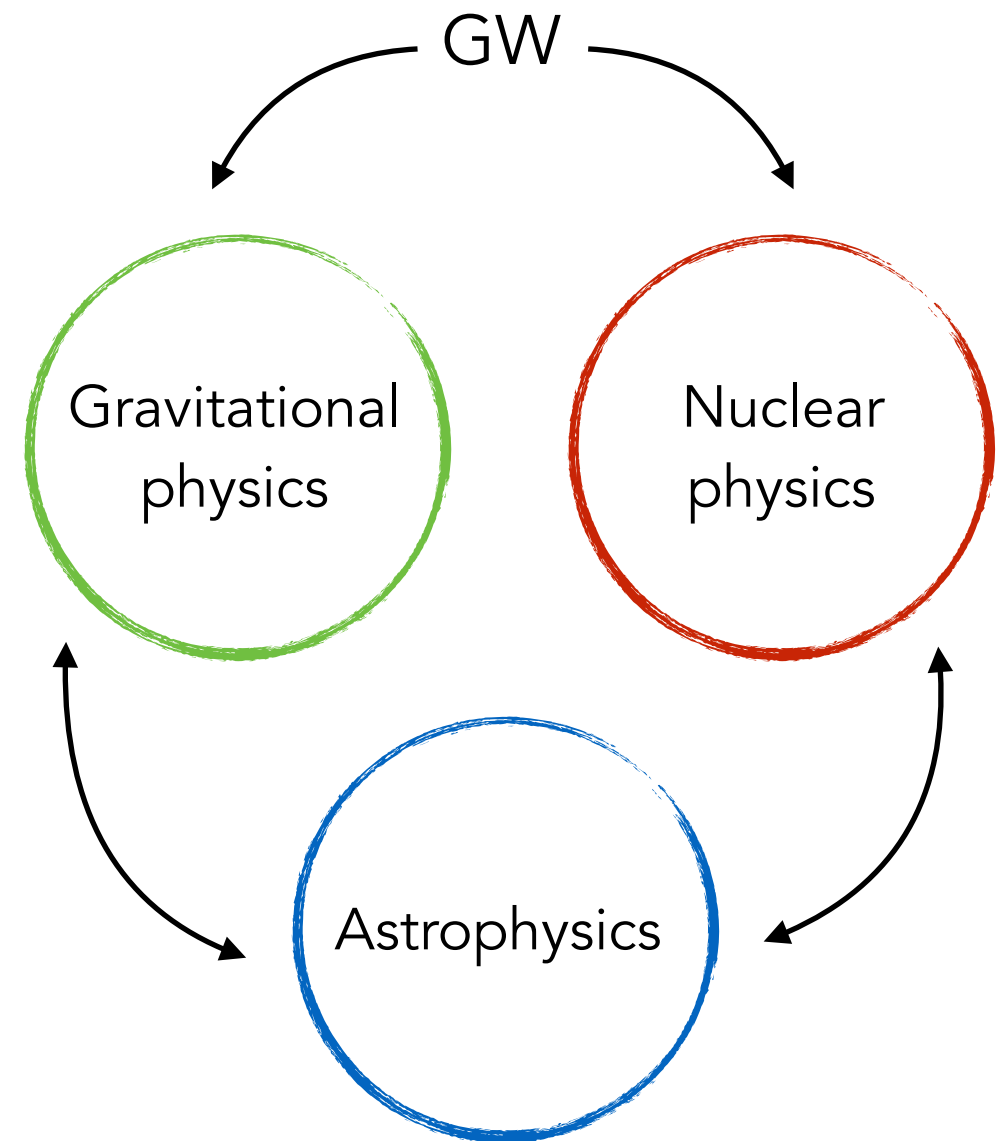
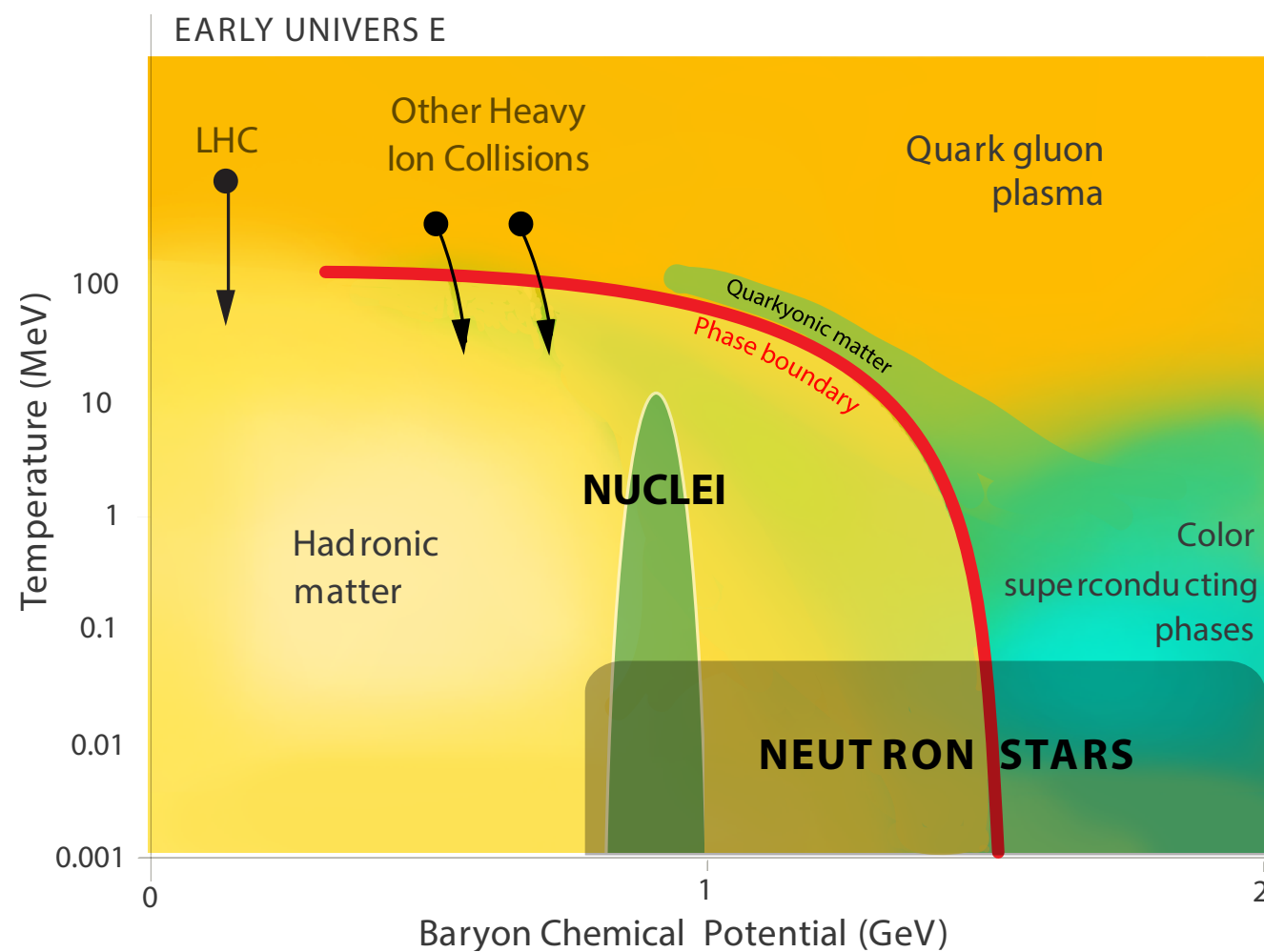
(LIGO Scientific Collaboration and Virgo Collaboration)

(Received 26 September 2017; revised manuscript received 2 October 2017; published 16 October 2017)



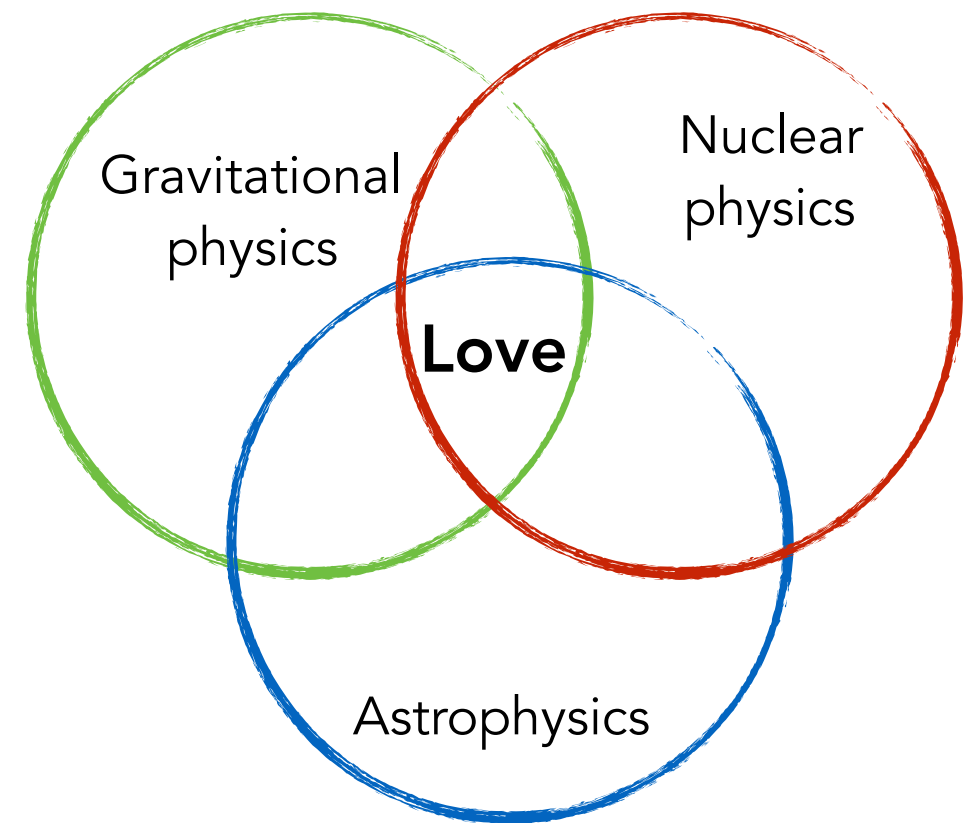
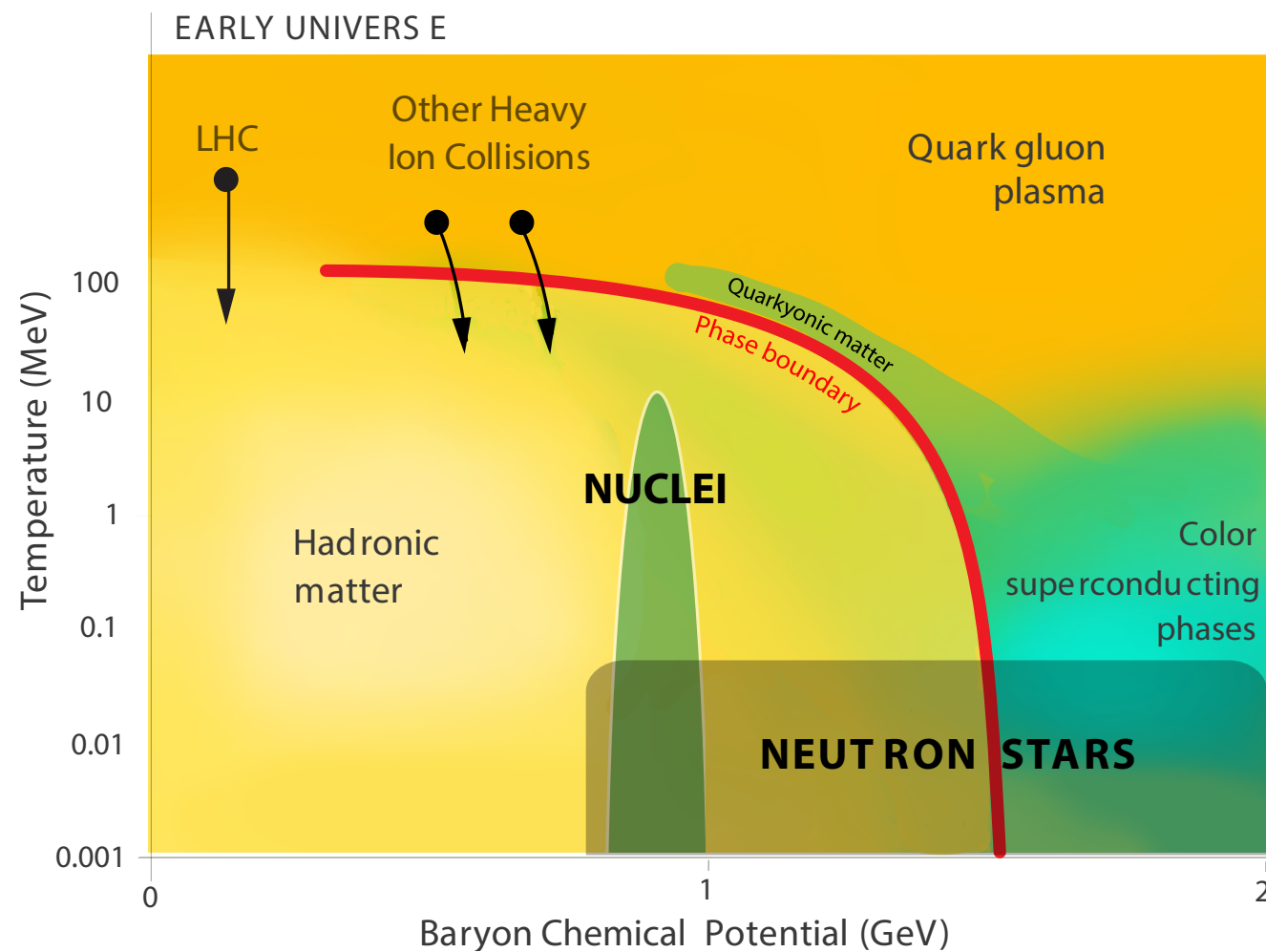
# Fundamental physics with NS

✓ How nuclear matter behaves under extreme conditions?



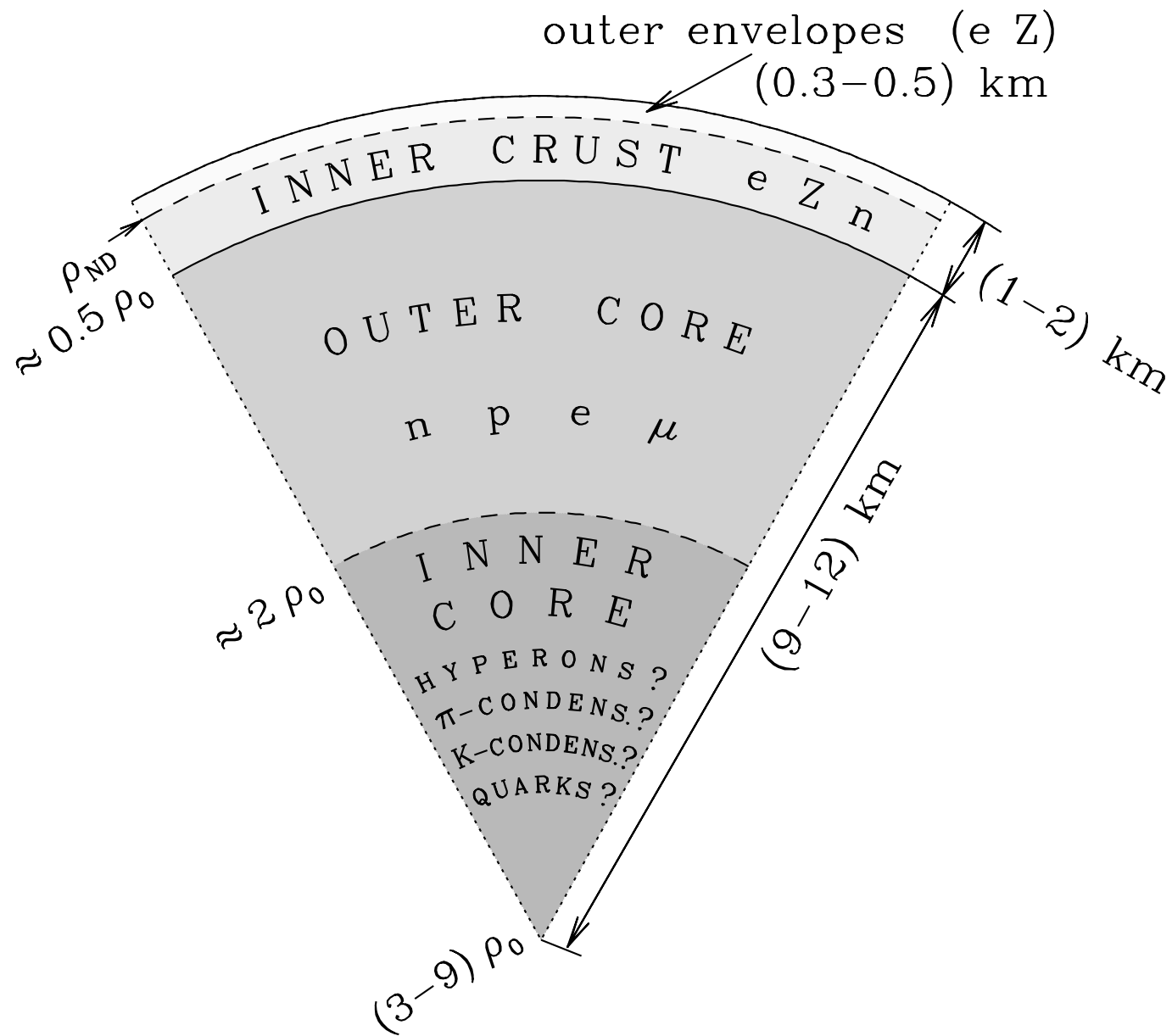
# Fundamental physics with NS

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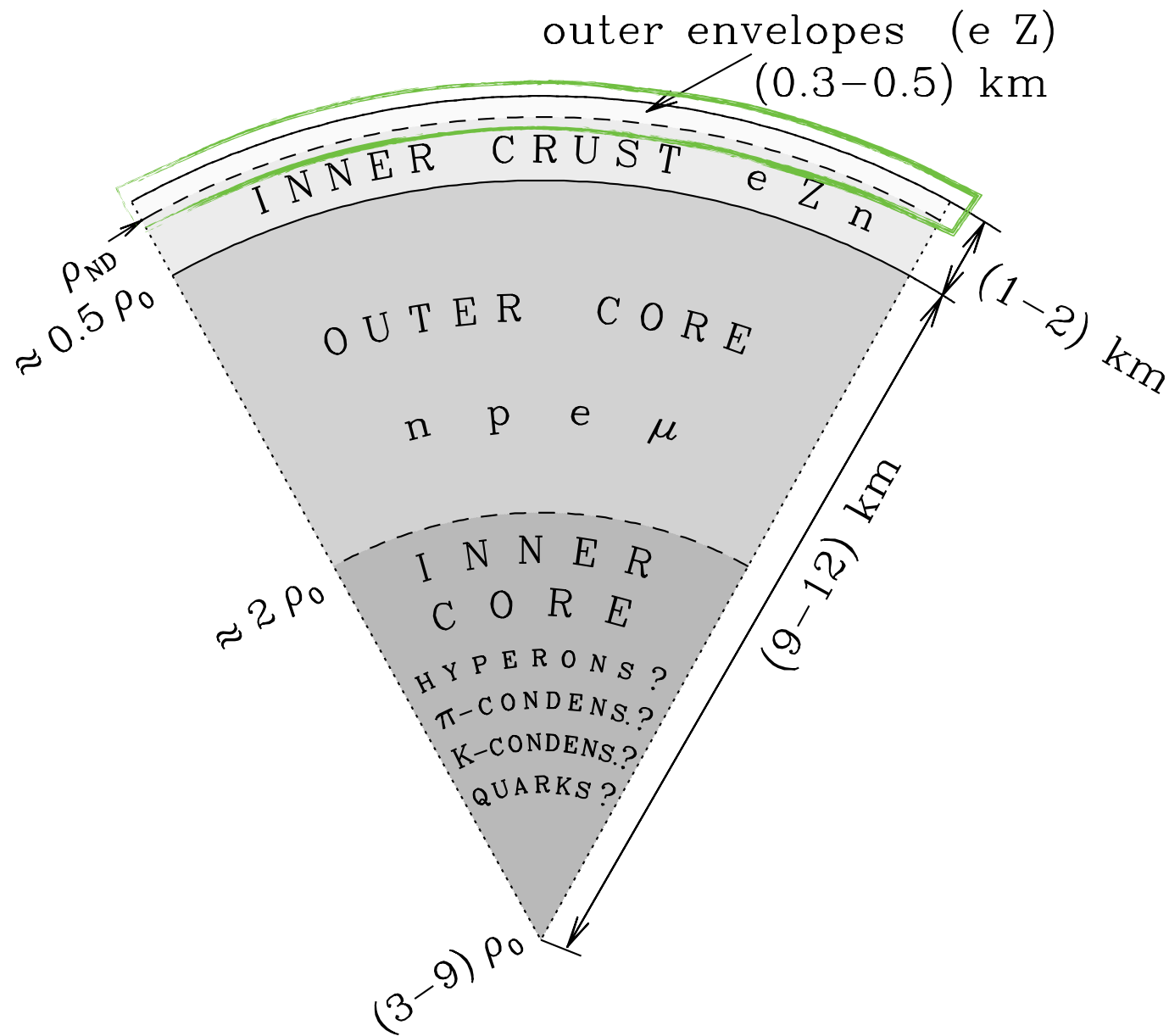


# Inside a NS



$$\rho_0 \simeq 2.67 \times 10^{14} \text{ g cm}^{-3}$$

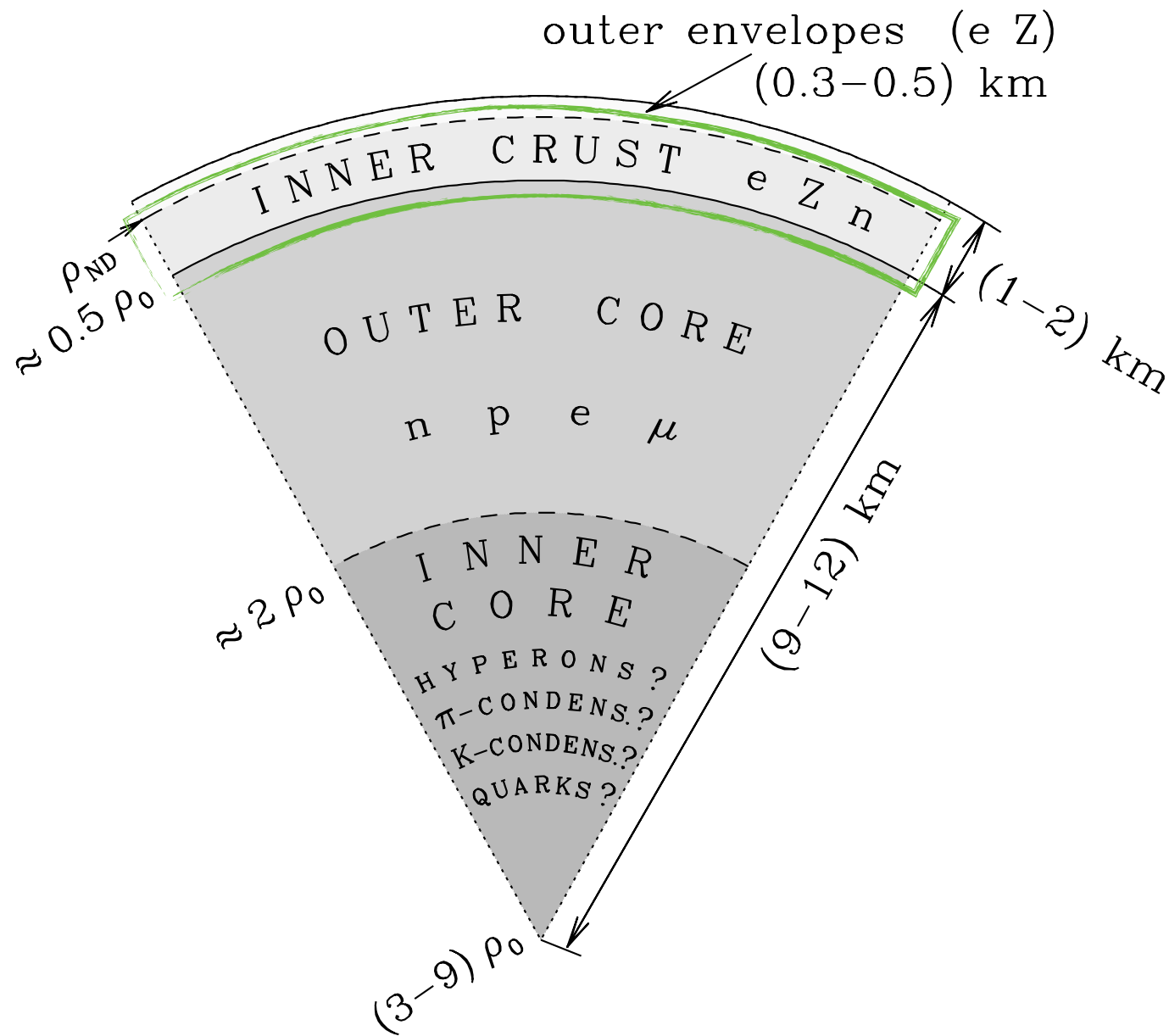
# Inside a NS



Nuclei lattice within  $e^-$  gas

$$\rho_0 \simeq 2.67 \times 10^{14} \text{ g cm}^{-3}$$

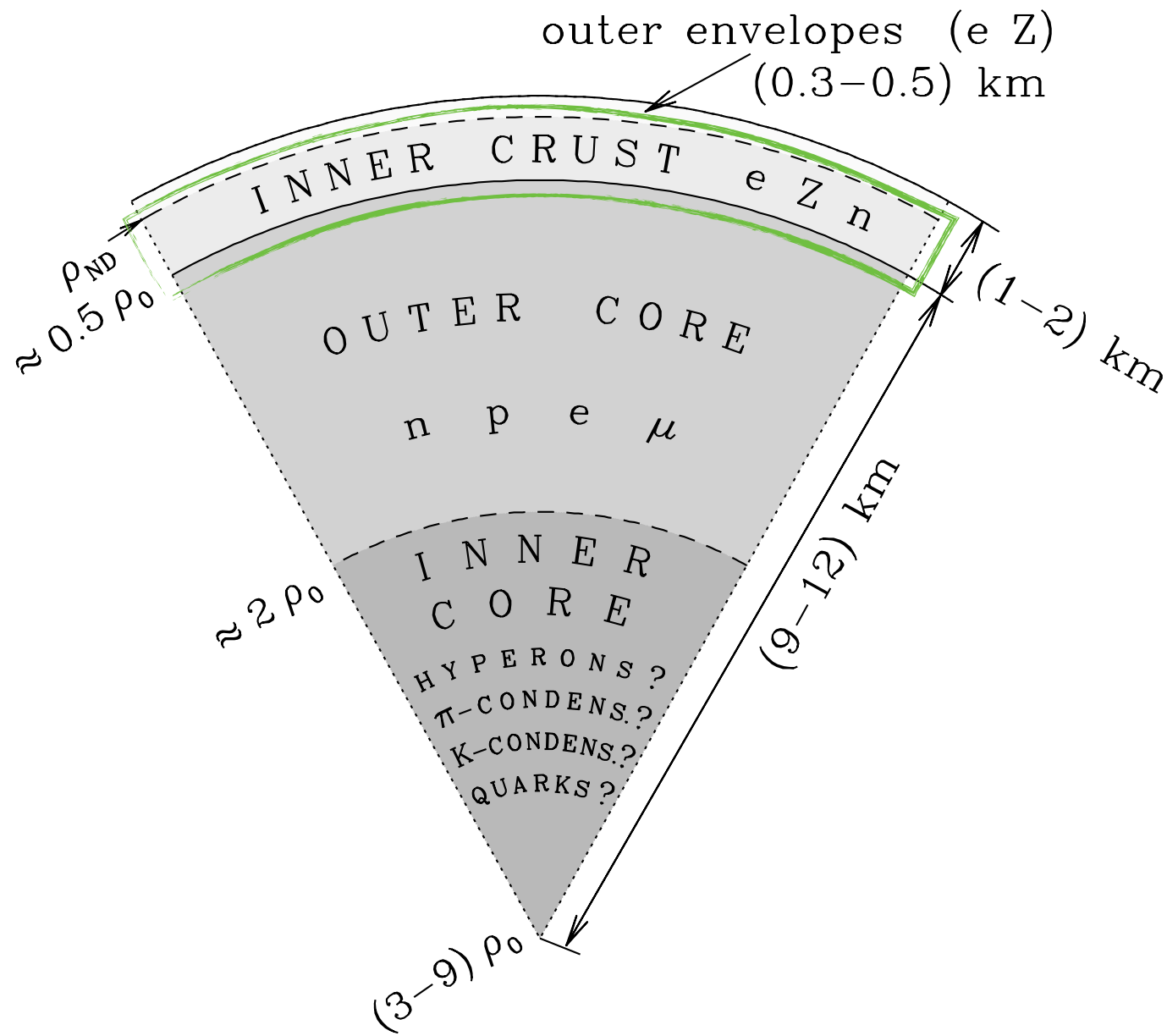
# Inside a NS



Pasta phases

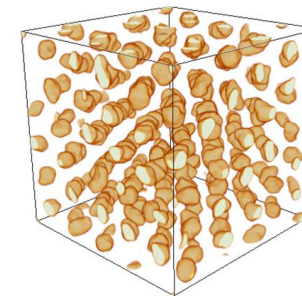
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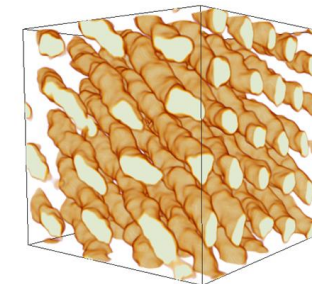


## Pasta phases

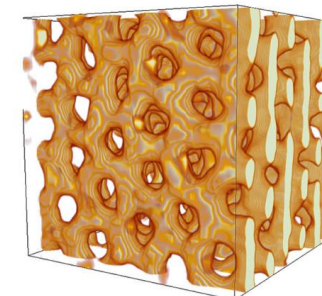
(a) *Gnocchi*



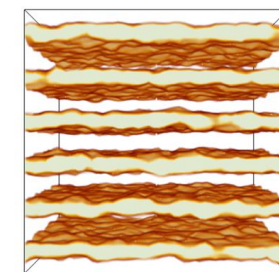
(b) *Spaghetti*



(c) *Waffles*

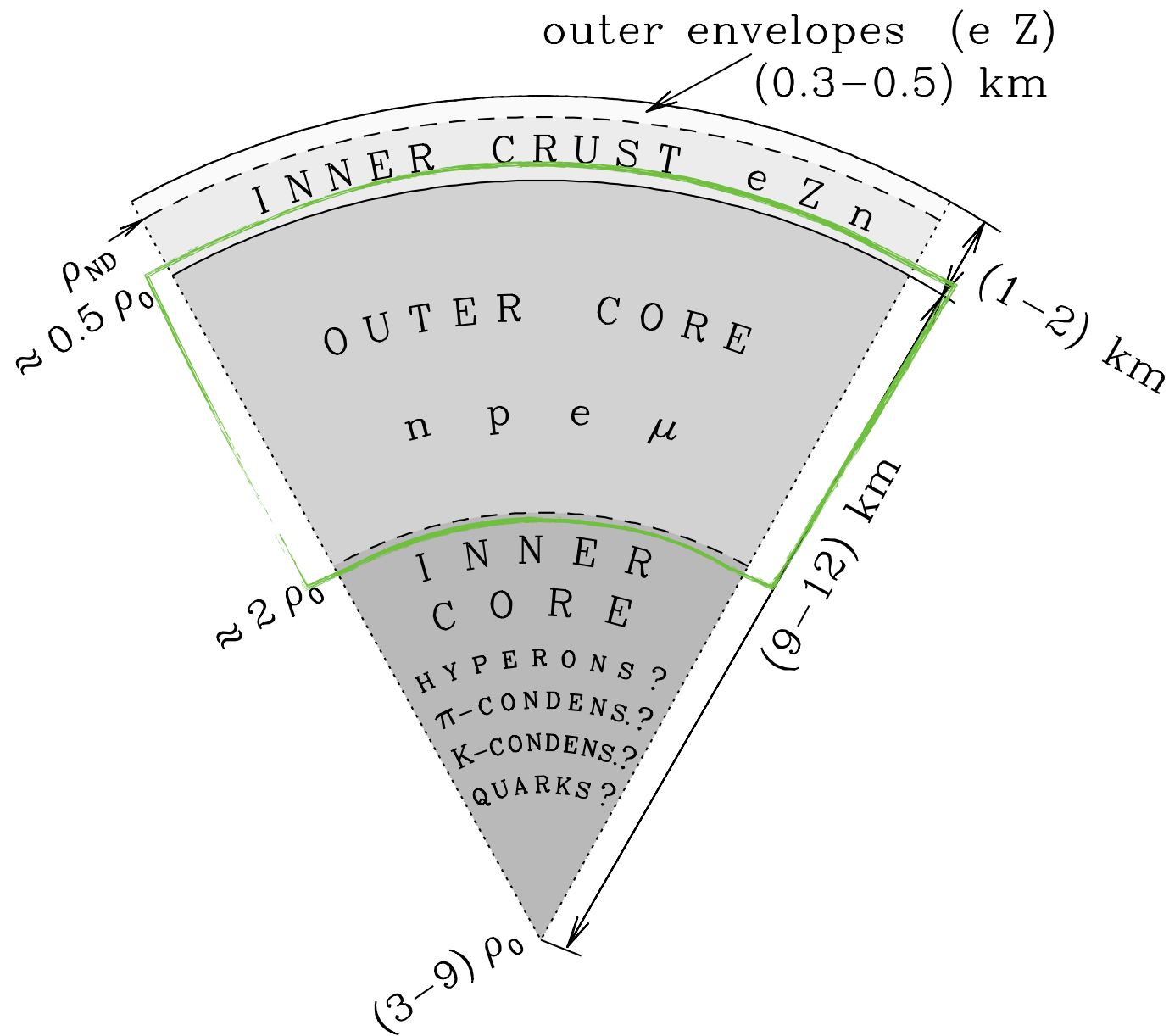


(d) *Lasagna*



$$\rho_0 \simeq 2.67 \times 10^{14} \text{ g cm}^{-3}$$

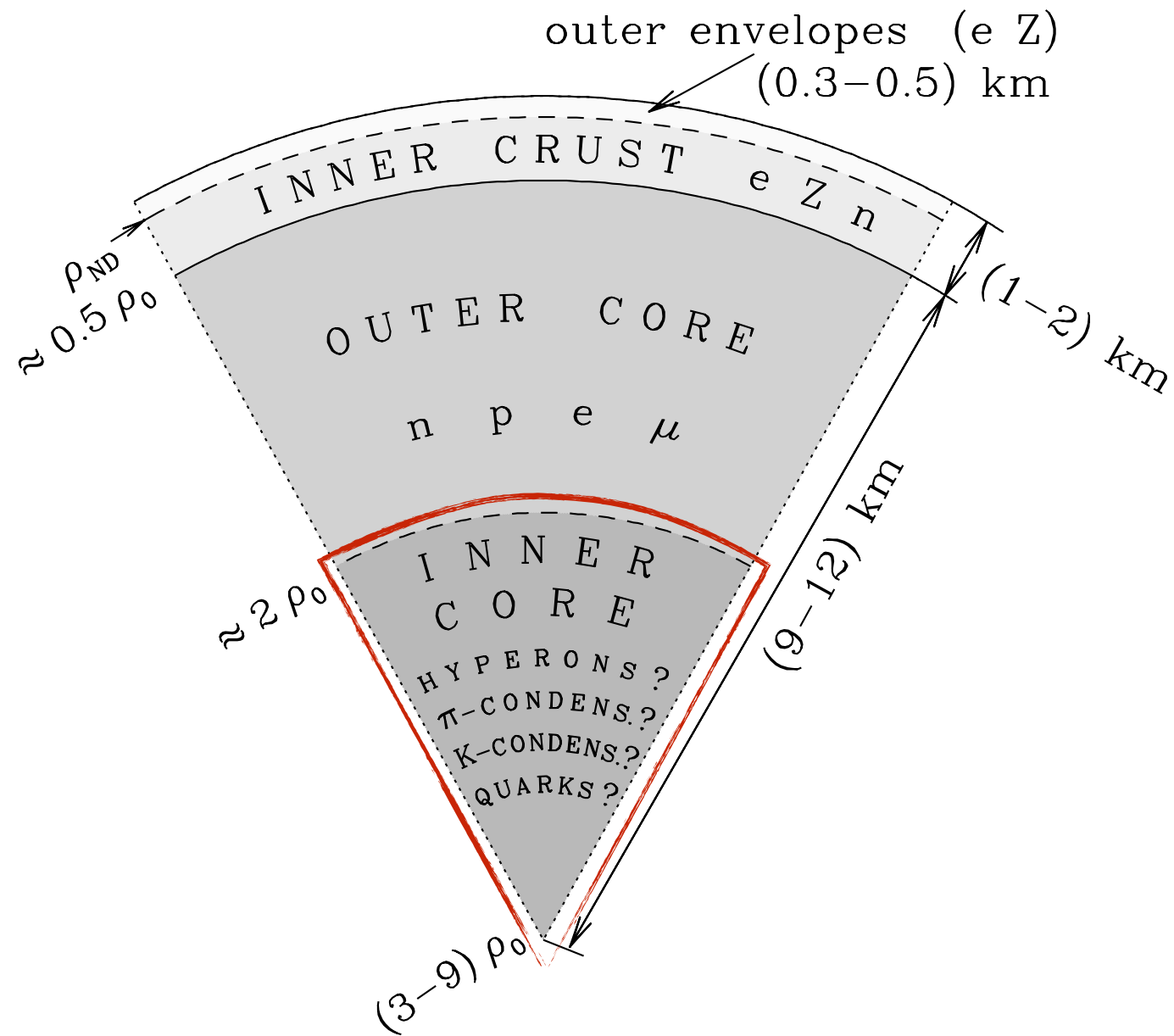
# Inside a NS



Nucleonic matter in  
 $\beta$ – equilibrium

$$\rho_0 \simeq 2.67 \times 10^{14} \text{ g cm}^{-3}$$

# Inside a NS

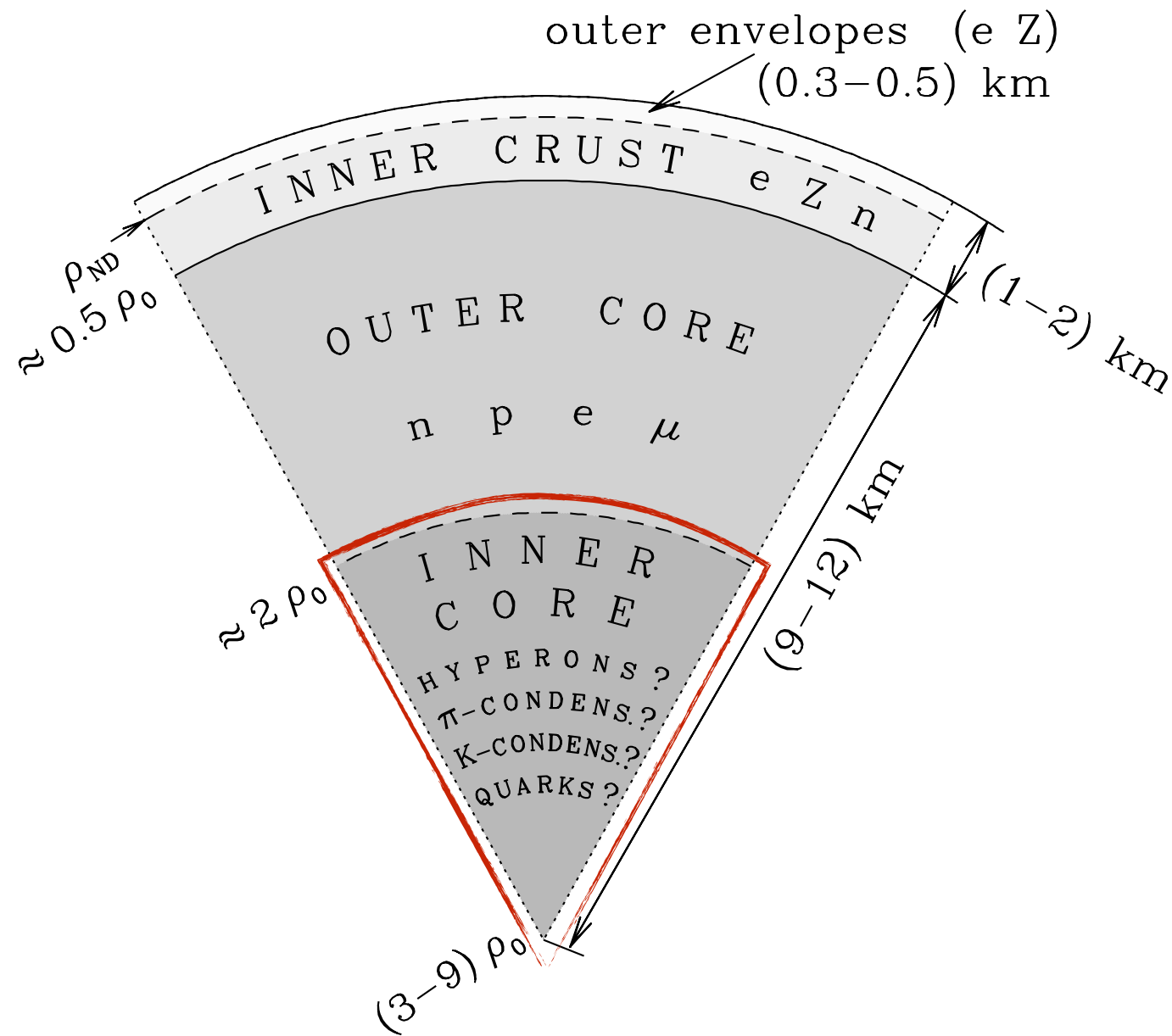


Phase transitions,  
quark-gluon plasma ...

$$\rho_0 \simeq 2.67 \times 10^{14} \text{ g cm}^{-3}$$



# Inside a NS



$$\rho_0 \simeq 2.67 \times 10^{14} \text{ g cm}^{-3}$$

Phase transitions,  
quark-gluon plasma ...

Constraints at  
super-nuclear density

# Inside a NS

---

Different models to be tested by observations

- ✓ Non-Relativistic Many Body theory

$$\mathcal{H} = \sum_i \frac{p_i^2}{2m} + \sum_{j>i} v_{ij} + \sum_{k>i>j} V_{ijk}$$

2-body

3-body

- ✓ Relativistic Mean field theory

$$\mathcal{L} = \mathcal{L}_N + \mathcal{L}_B + \mathcal{L}_{\text{int}}$$

- ✓ Can be generalised to include strange baryons [interactions?]
- ✓ *MIT Bag model* to describe unconfined quarks

# Inside a NS

---

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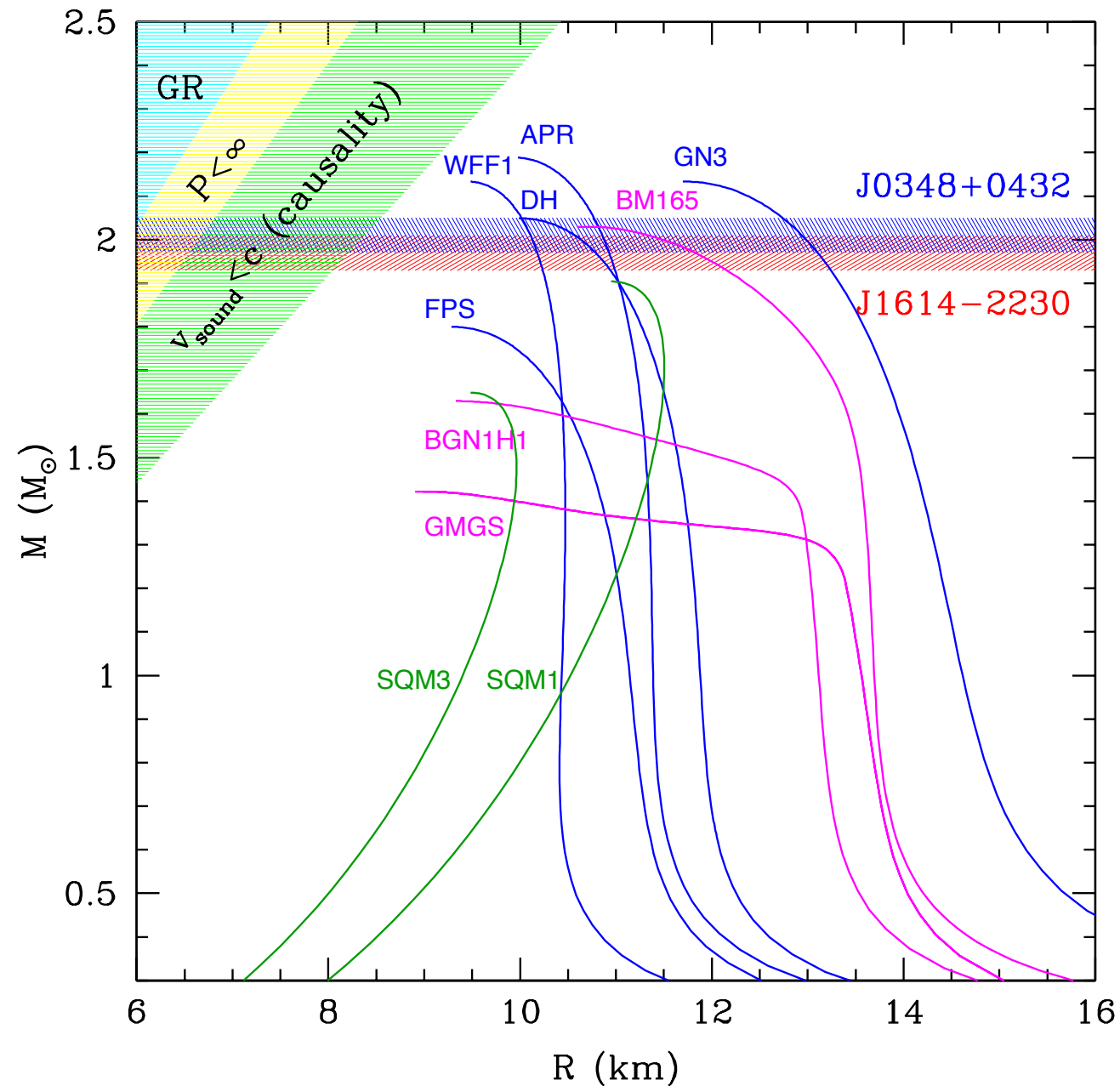
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ask O. Benhar

# From micro to macro

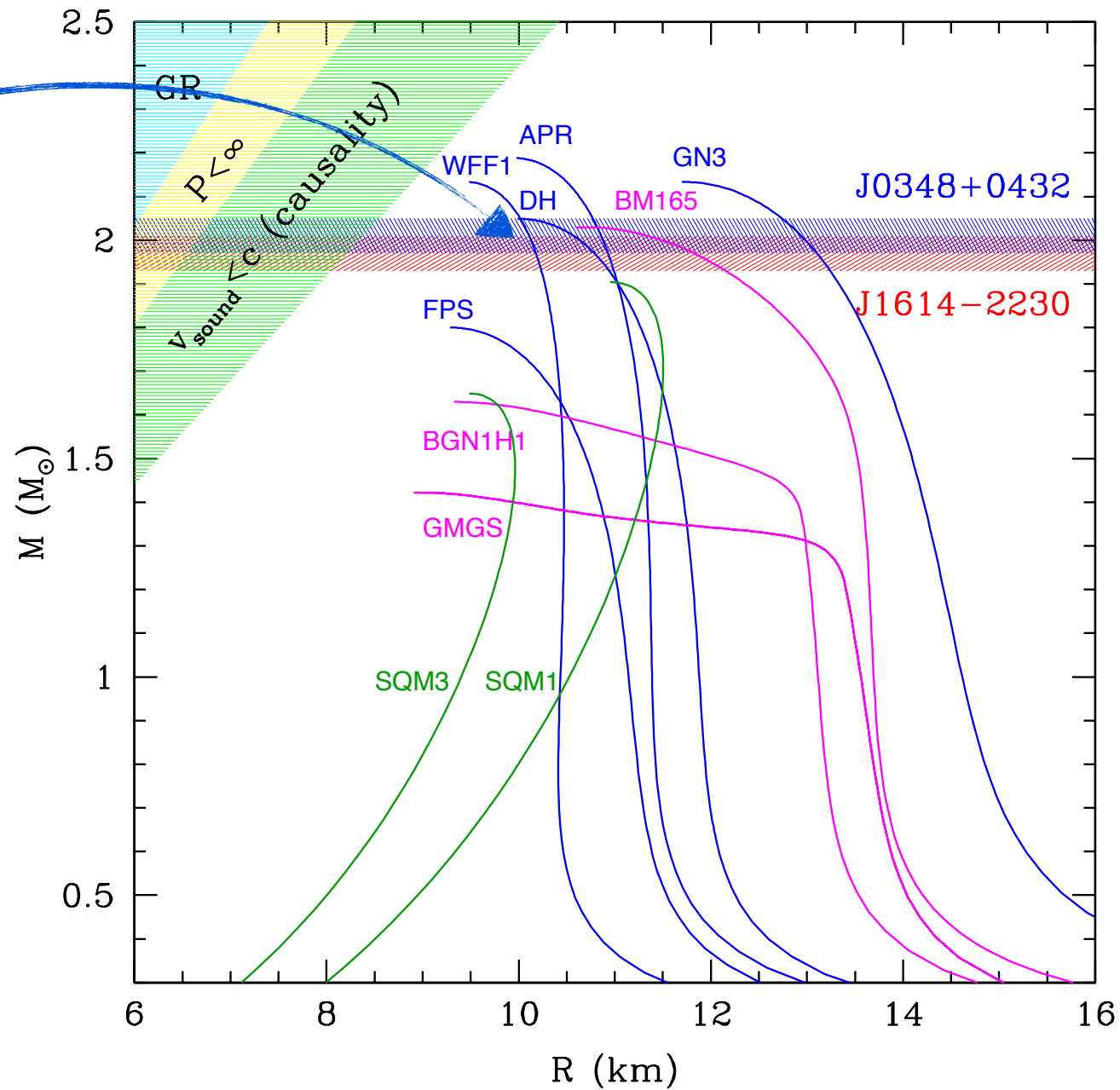
The Equation of State  $p = p(\epsilon)$



# From micro to macro

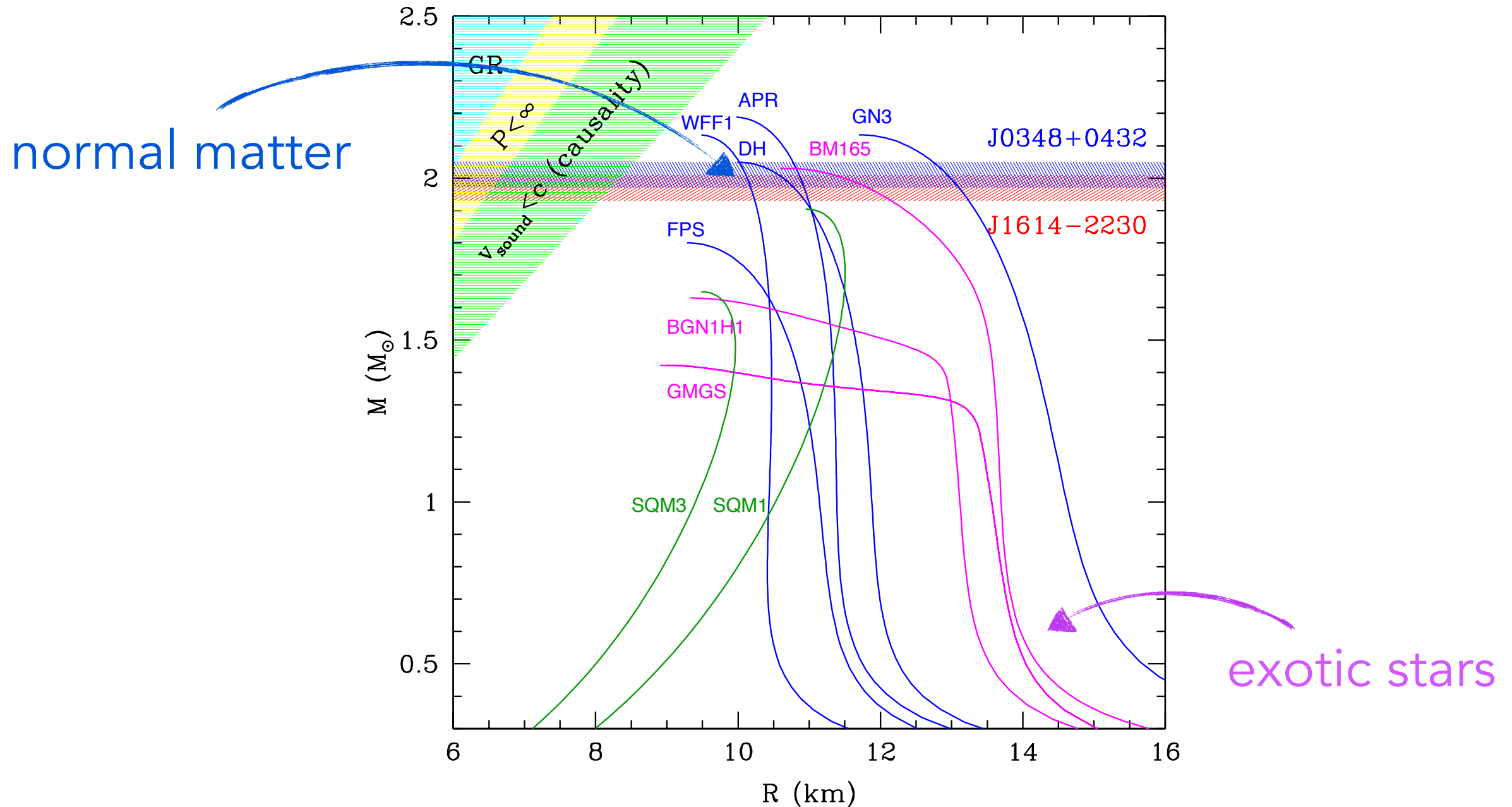
The Equation of State  $p = p(\epsilon)$

normal matter



# From micro to macro

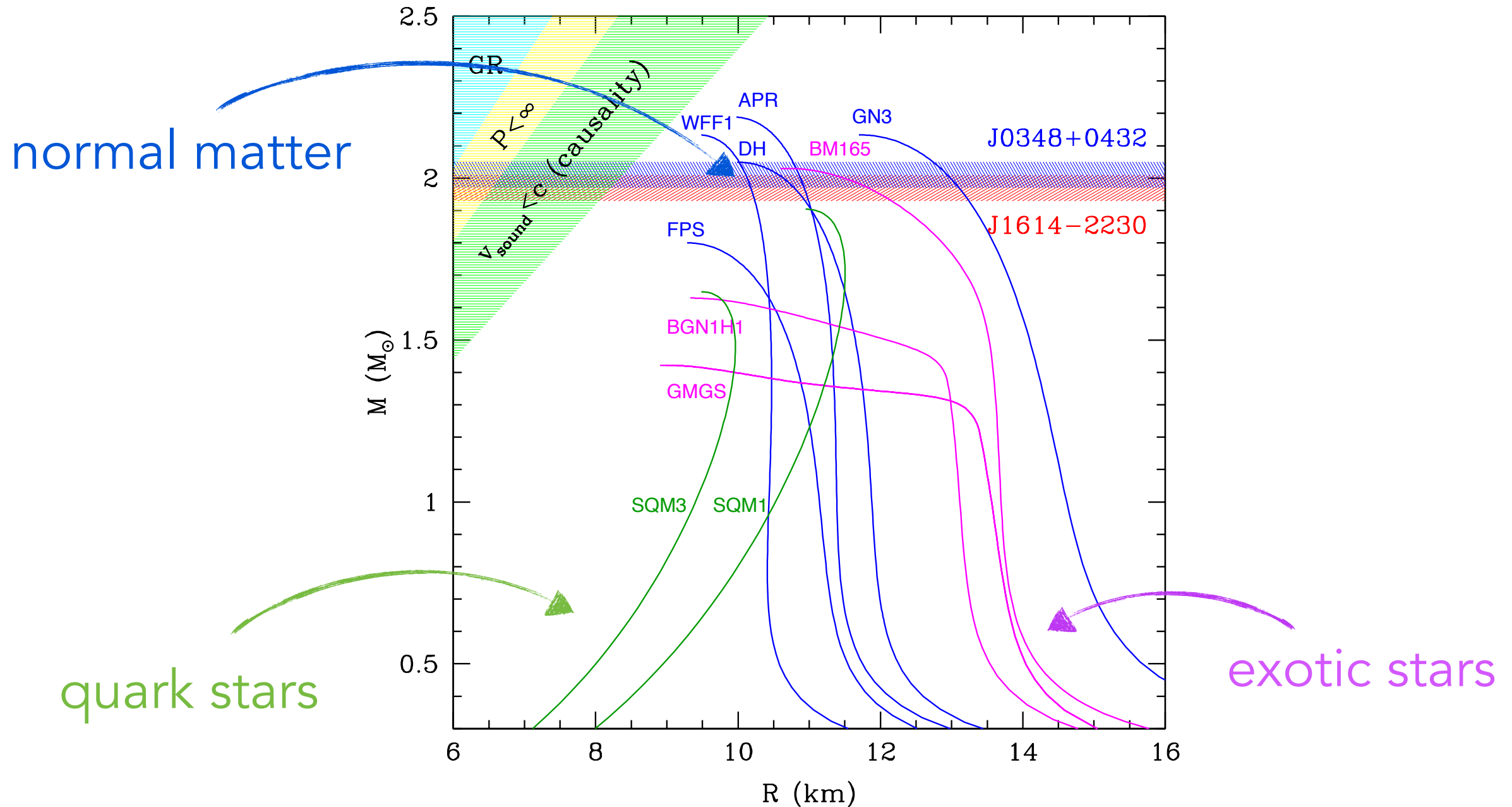
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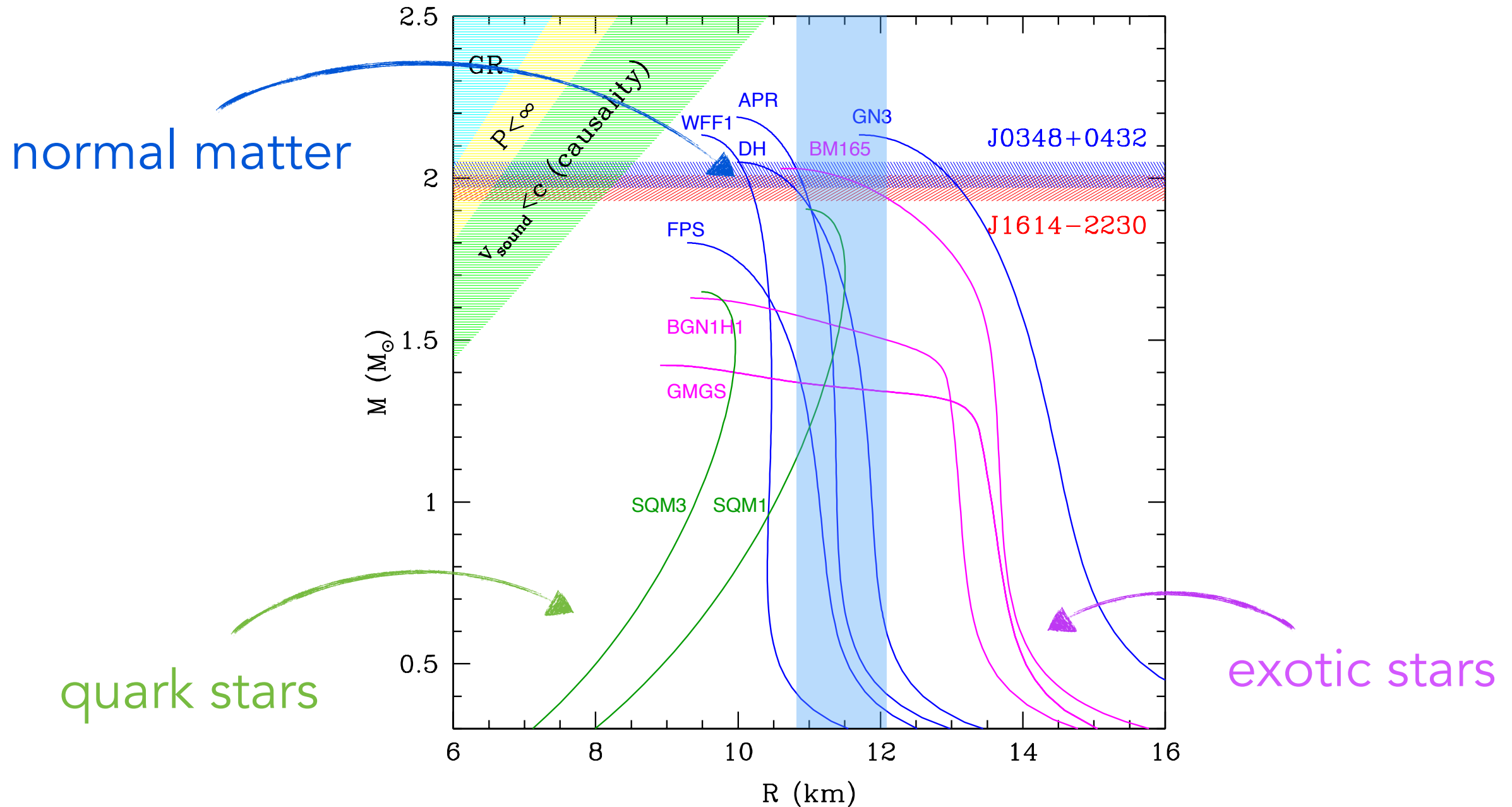
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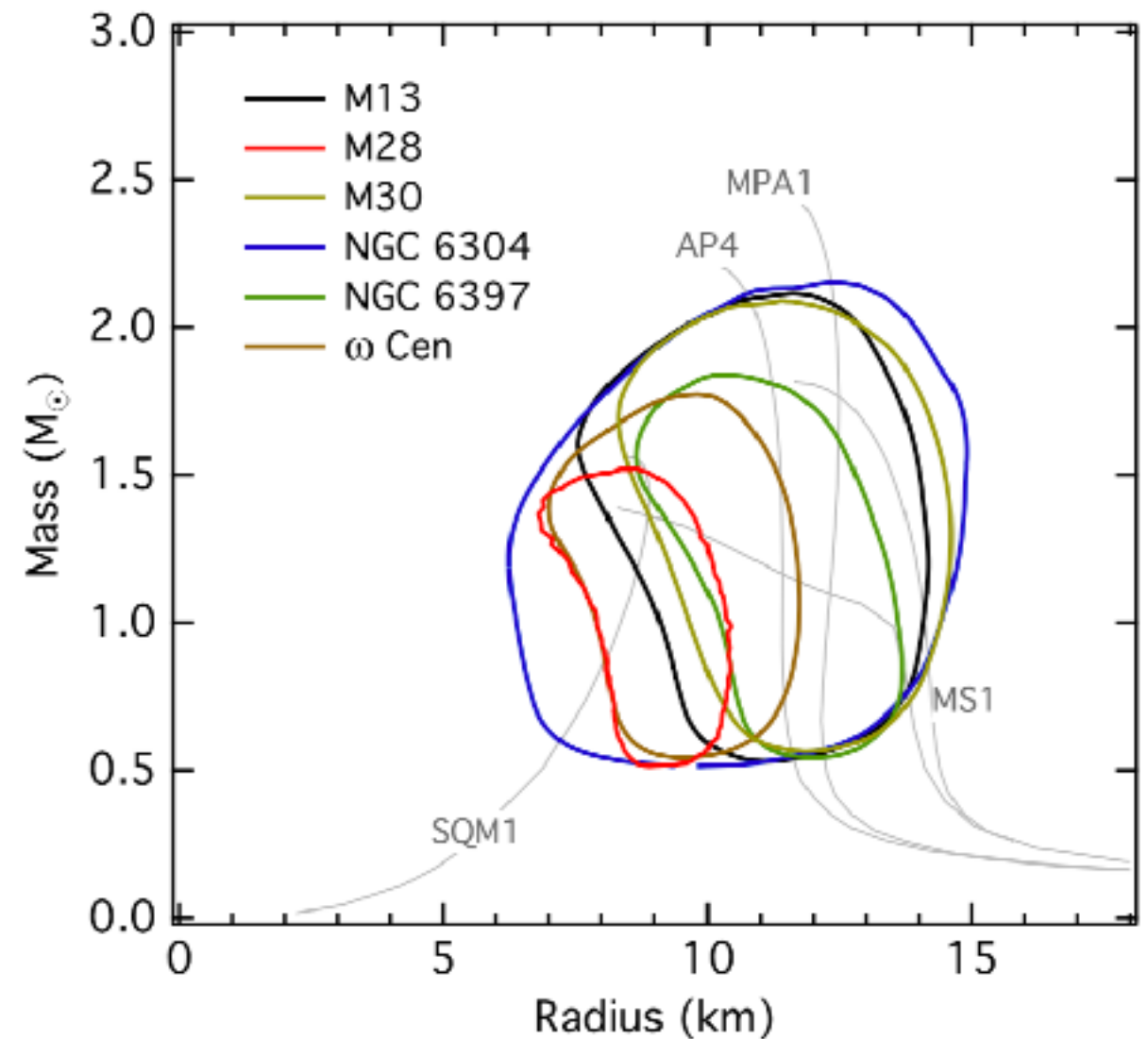


# The radius so far

Observed masses  $[1 \div 2]M_{\odot}$

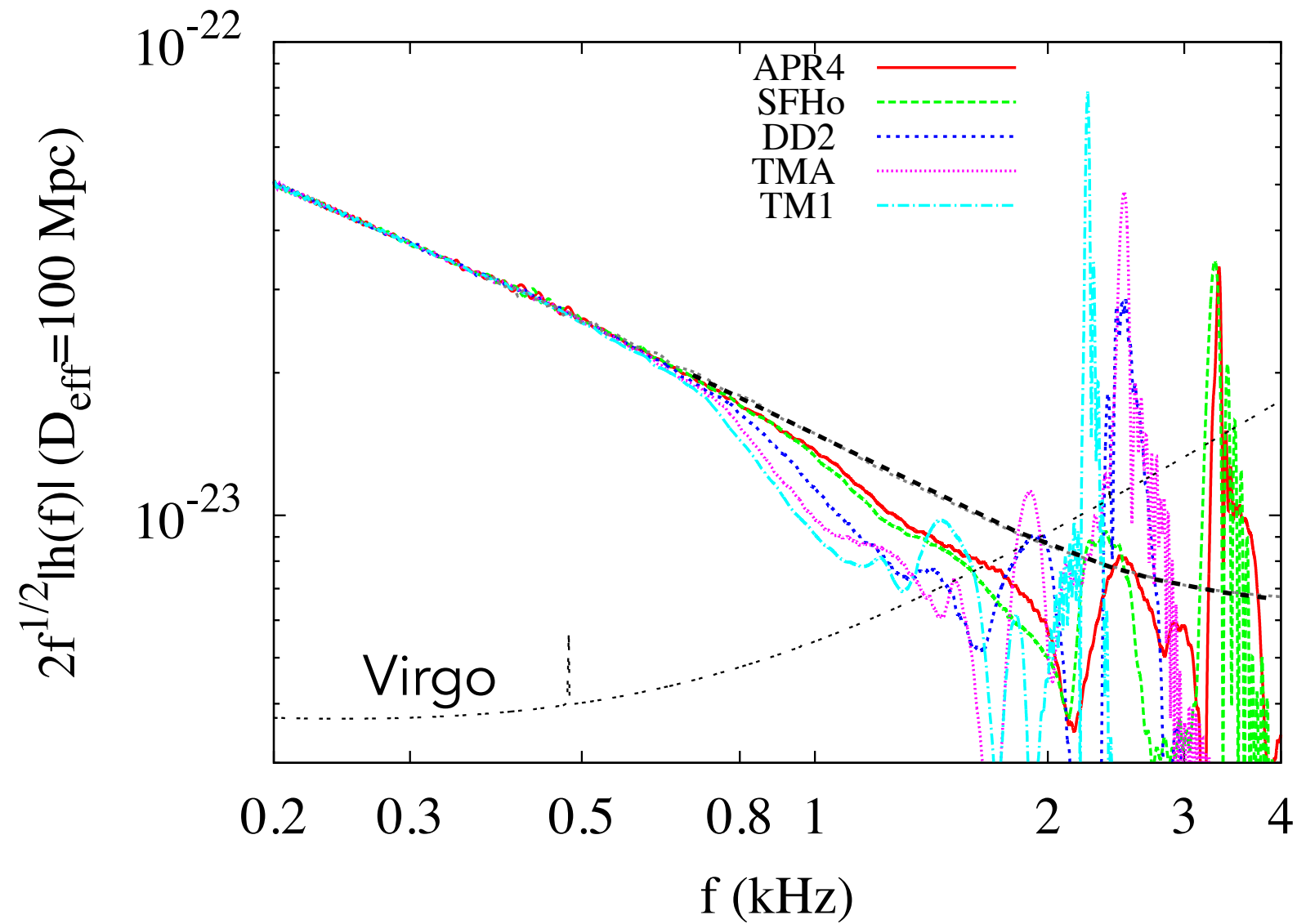
*Ozel et al., 2016*

- ✓ Radius measured with [13-15]% of accuracy
- ✓ Difficult to model emission processes (magnetic fields, atmosphere...)
- ✓ We could look for a **cleaner** framework



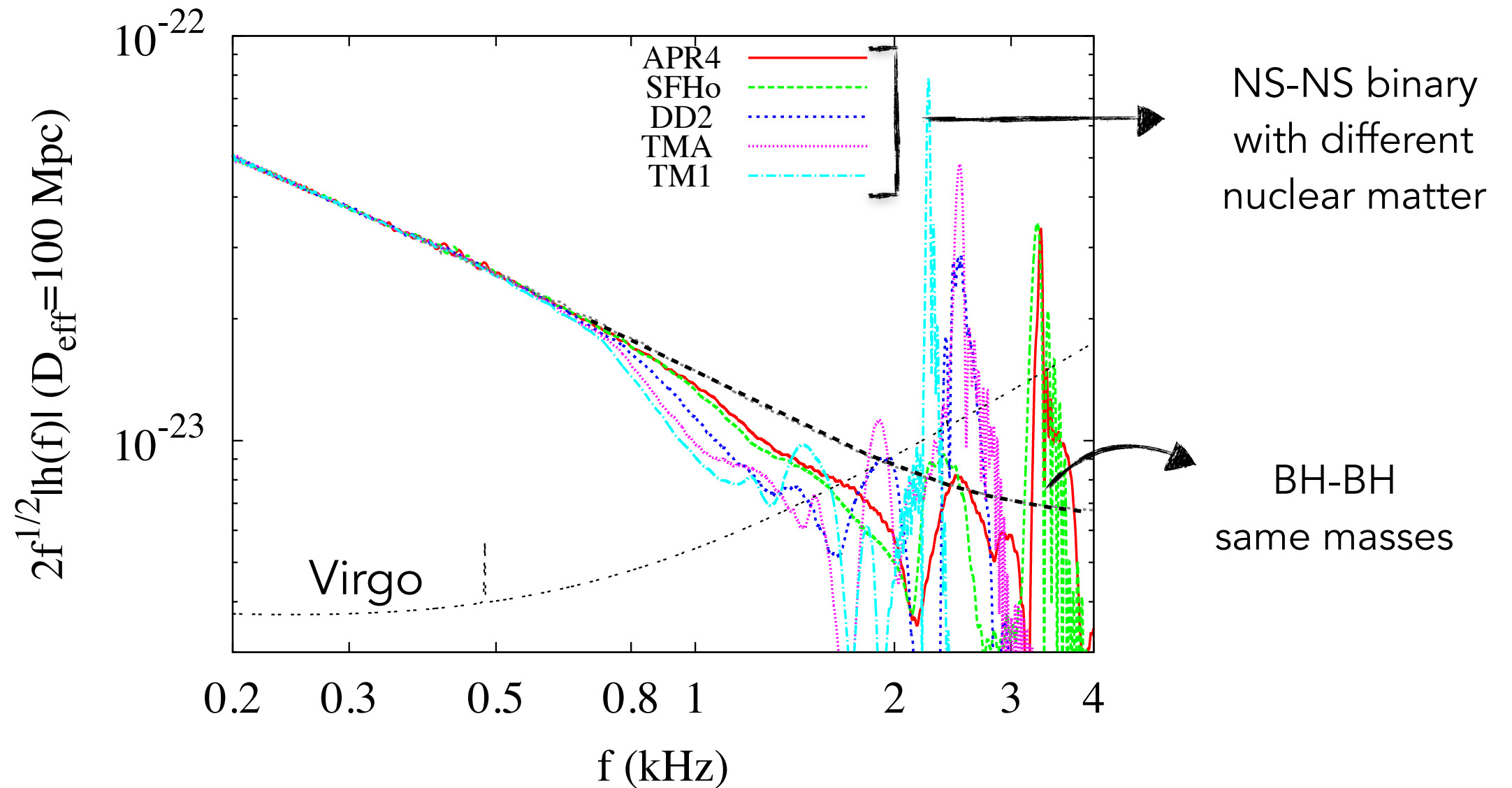
# What do we look for?

Hotokezaka, Kyutoku, Sekiguchi, Shibata, 2016



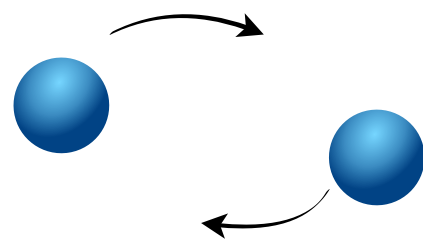
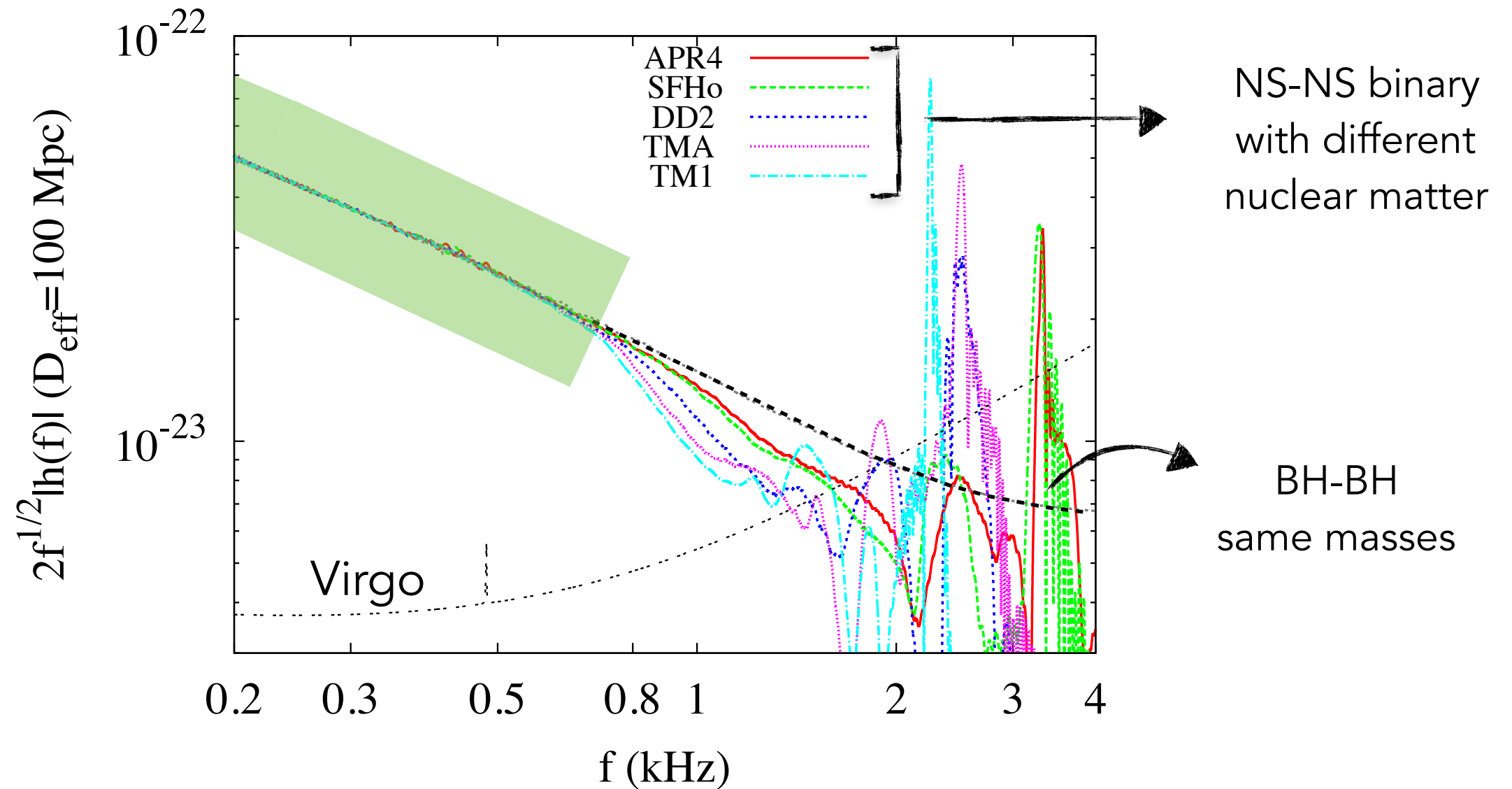
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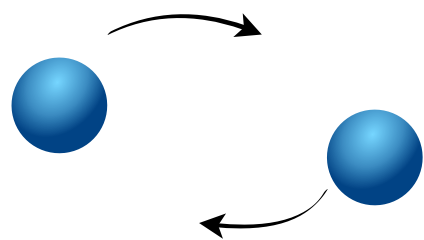
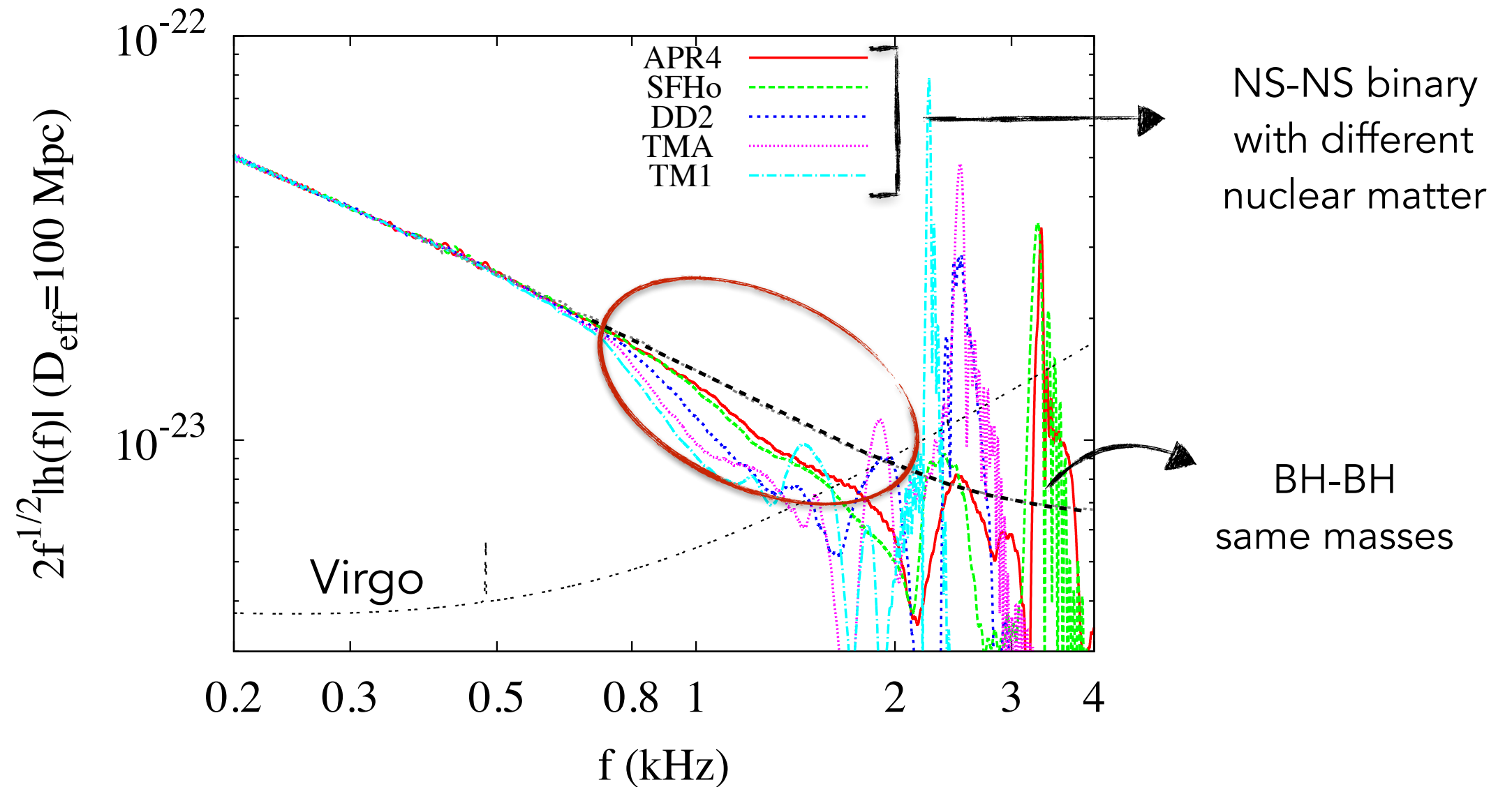


point masses

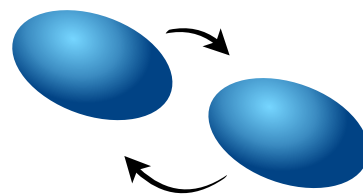


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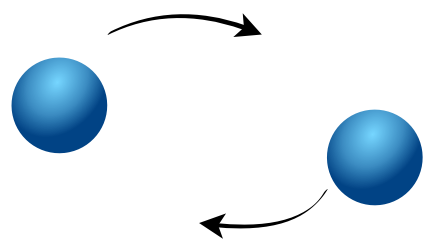
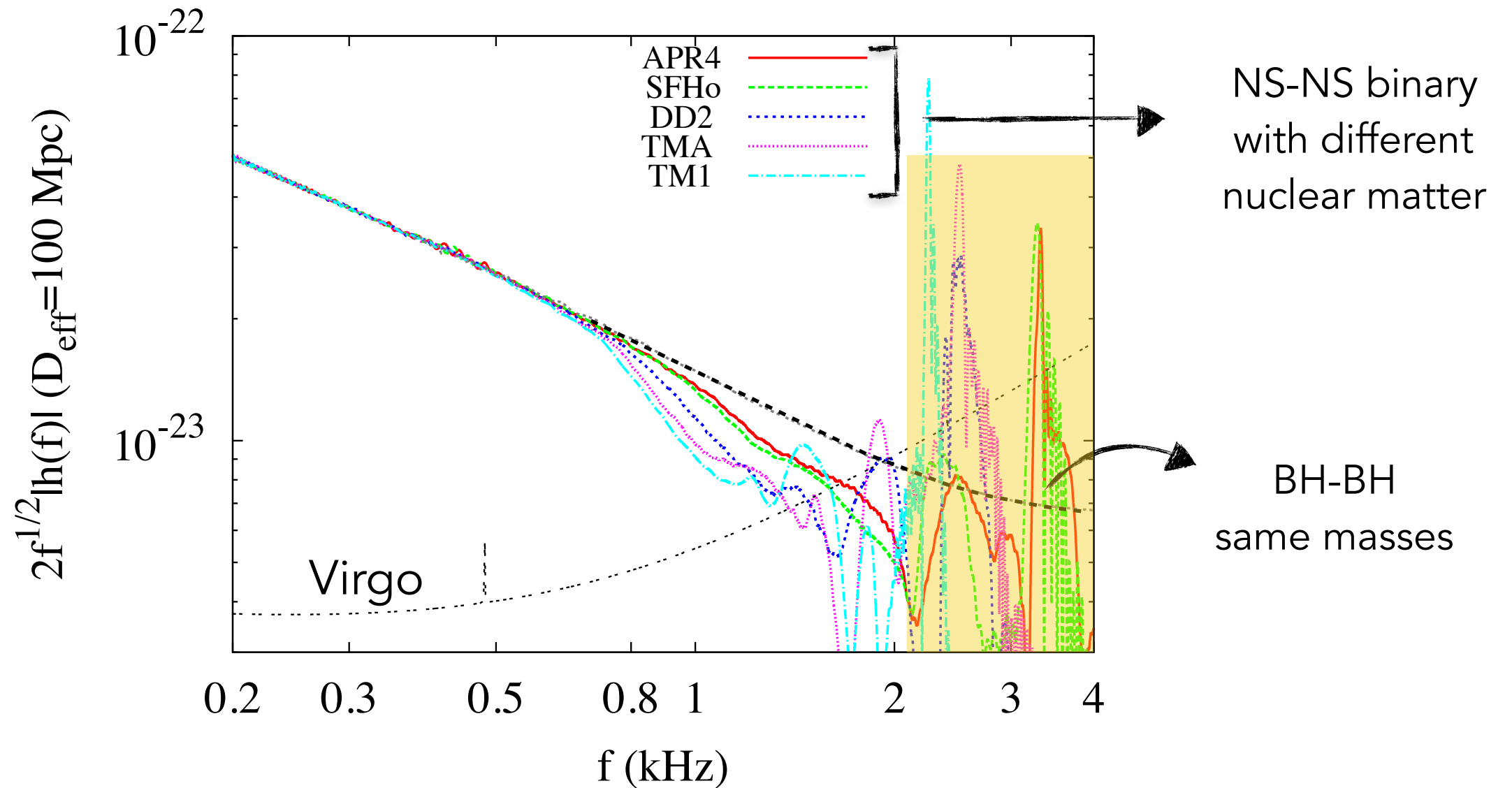
point masses



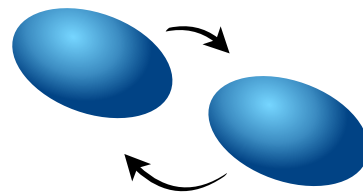
tidal interactions

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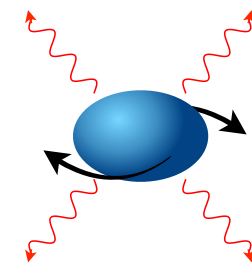
Hotokezaka, Kyutoku, Sekiguchi, Shibata, 2016



point masses



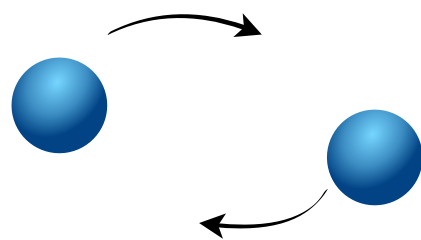
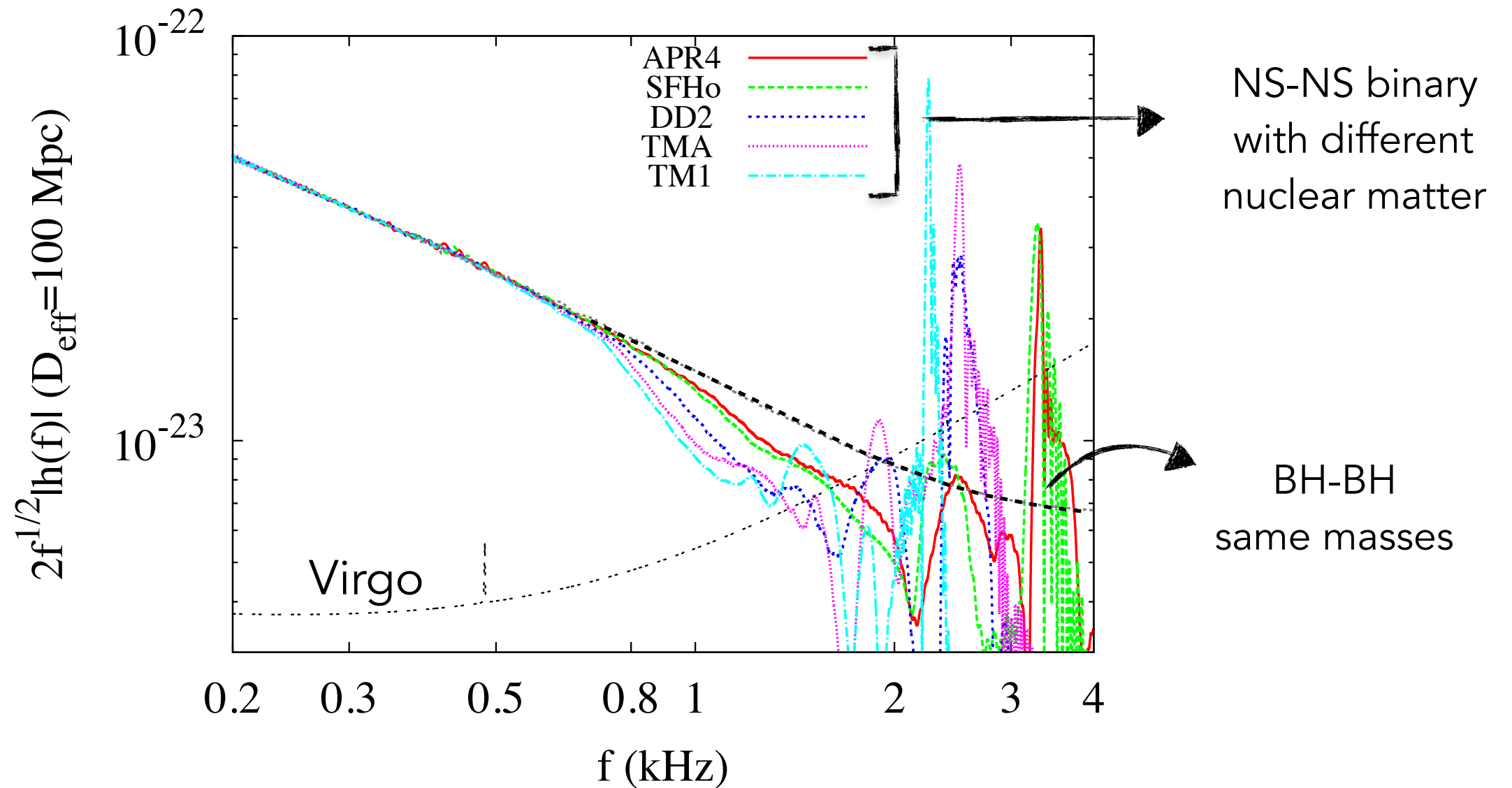
tidal interactions



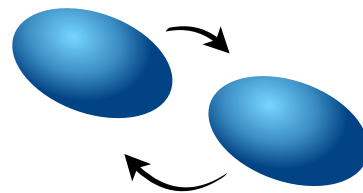
merger/post-merger

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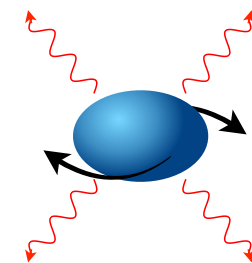
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point masses



tidal interactions



merger/post-merger

# The Love number

---

Tidal interactions leave the footprint of the NS structure on the signal

*Hinderer, '08; Binnington & Poisson '09  
Damour & Nagar '10*

- ✓ Deformation properties encoded within the **Love numbers**

$$Q_{ij} = \frac{2}{3} k_2 R^5 \mathcal{E}_{ij} = \lambda \mathcal{E}_{ij}$$

- ✓  $\lambda$  depends on the EoS only, for a given compactness  $M_{\text{NS}}/R_{\text{NS}}$
- ✓  $\lambda$  enters within the gravitational waveform

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Tidal interactions leave the footprint of the NS structure on the signal

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The diagram shows a central equation  $Q_{ij} = \frac{2}{3} k_2 R^5 \mathcal{E}_{ij} = \lambda \mathcal{E}_{ij}$ . A horizontal line above the equation has two vertical lines extending downwards from its ends. From each vertical line, an arrow points down to the text 'star's quadrupole' on the left and 'external tidal field' on the right.

$$Q_{ij} = \frac{2}{3} k_2 R^5 \mathcal{E}_{ij} = \lambda \mathcal{E}_{ij}$$

star's quadrupole                      external tidal field

- ✓  $\lambda$  depends on the EoS only, for a given compactness  $M_{\text{NS}}/R_{\text{NS}}$
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# The recipe of Love

supplementary  
material



## Polar-electric-type perturbation of background metric

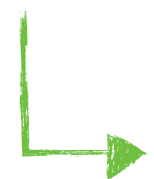
Regge & Wheeler '57

$$g_{\mu\nu} = g_{\mu\nu}^{(0)} + h_{\mu\nu}$$

$$\begin{pmatrix} -e^{\nu(r)} & 0 & 0 & 0 \\ 0 & e^{\lambda(r)} & 0 & 0 \\ 0 & 0 & r^2 & 0 \\ 0 & 0 & 0 & r^2 \sin^2 \theta \end{pmatrix} + \begin{pmatrix} -e^{\nu(r)} H_0(r) & 0 & 0 & 0 \\ 0 & e^{\lambda(r)} H_2(r) & 0 & 0 \\ 0 & 0 & r^2 K(r) & 0 \\ 0 & 0 & 0 & r^2 K(r) \sin^2 \theta \end{pmatrix} Y_{lm}(\theta, \phi)$$

✓ Cook everything within Einstein's equations  $G_{\mu\nu} = \kappa T_{\mu\nu}$

✓ Solve at linear order in the perturbations  $H_0, H_2, K$



set of ODEs sourced by star's density perturbations

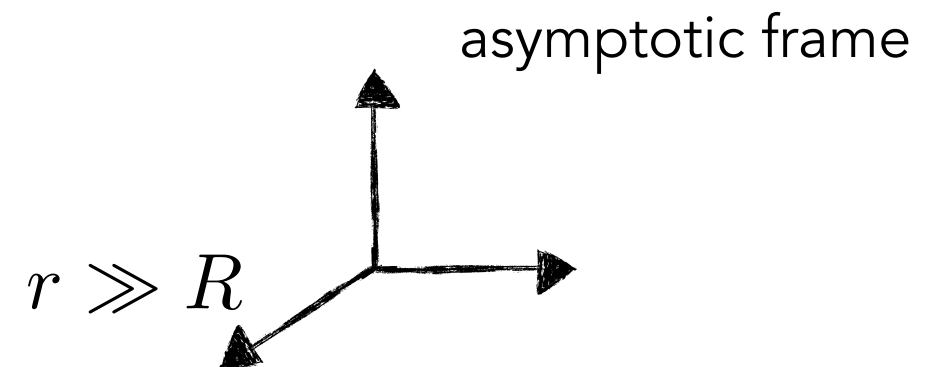
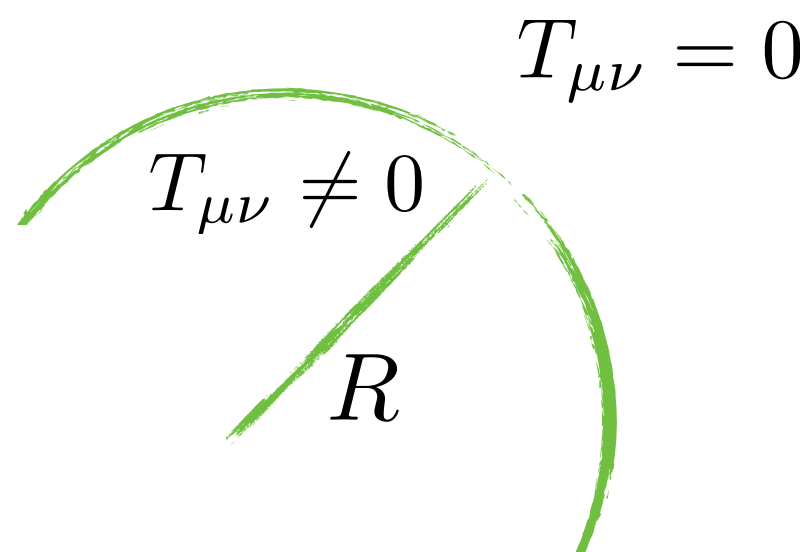


# The recipe of Love

supplementary  
material



- ✓ Numerically integrate from center to outside with appropriate boundary conditions



- ✓ Matching interior & exterior solution

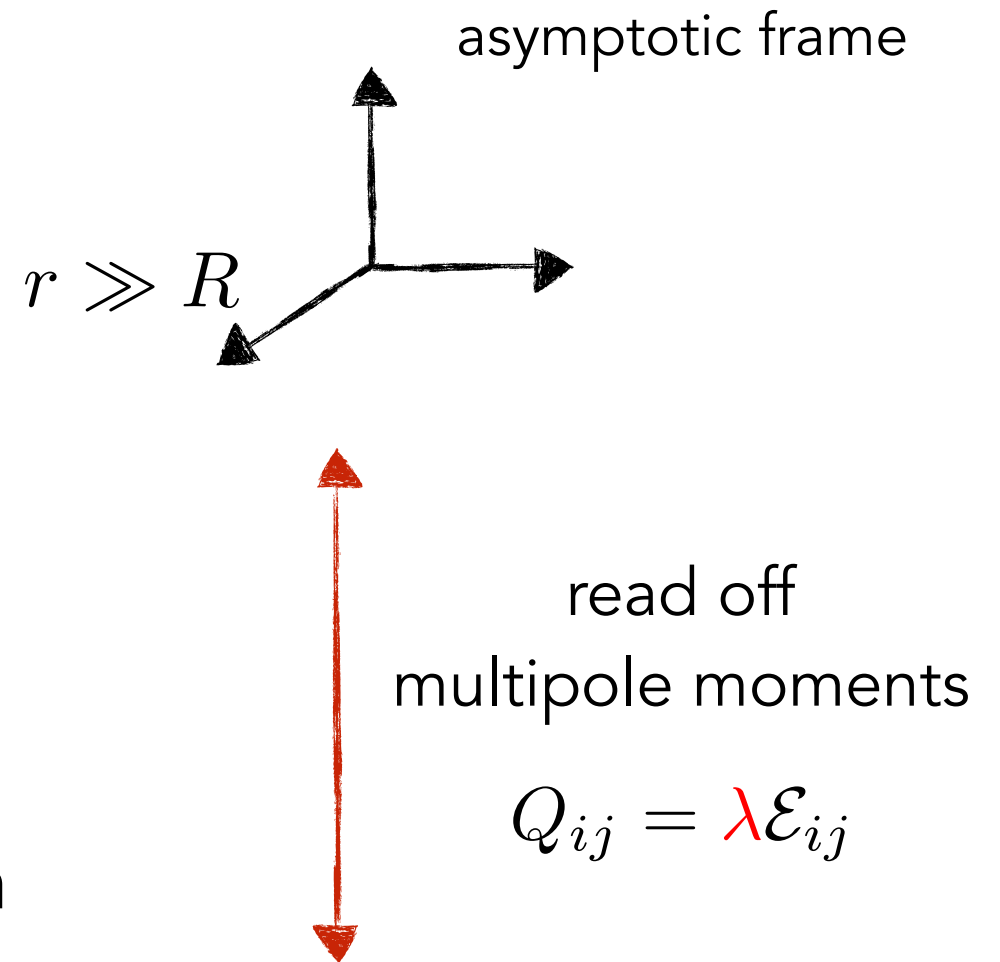
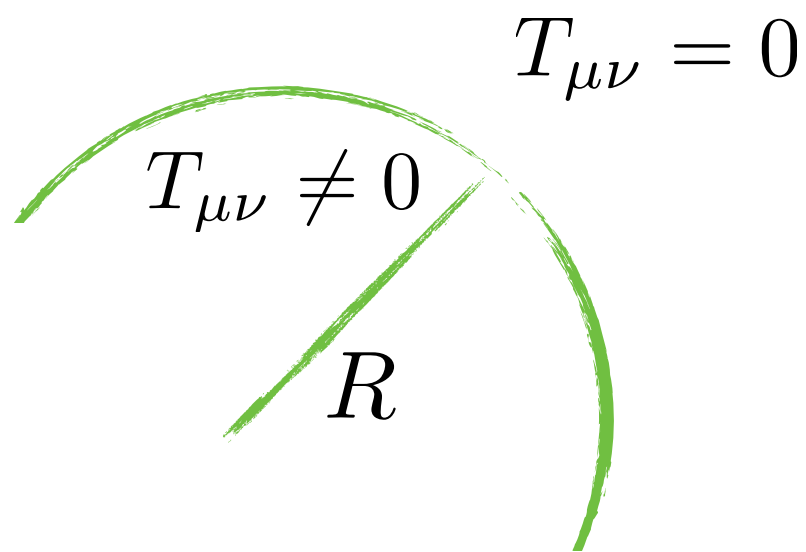
$$g_{tt}^{(\text{int})} = g_{tt}^{(\text{ext})} = -1 + \frac{2M}{r} + \frac{3Q_{ij}\hat{n}^i\hat{n}^j}{r^3} - \mathcal{E}_{ij}x^ix^j$$

# The recipe of Love

supplementary  
material



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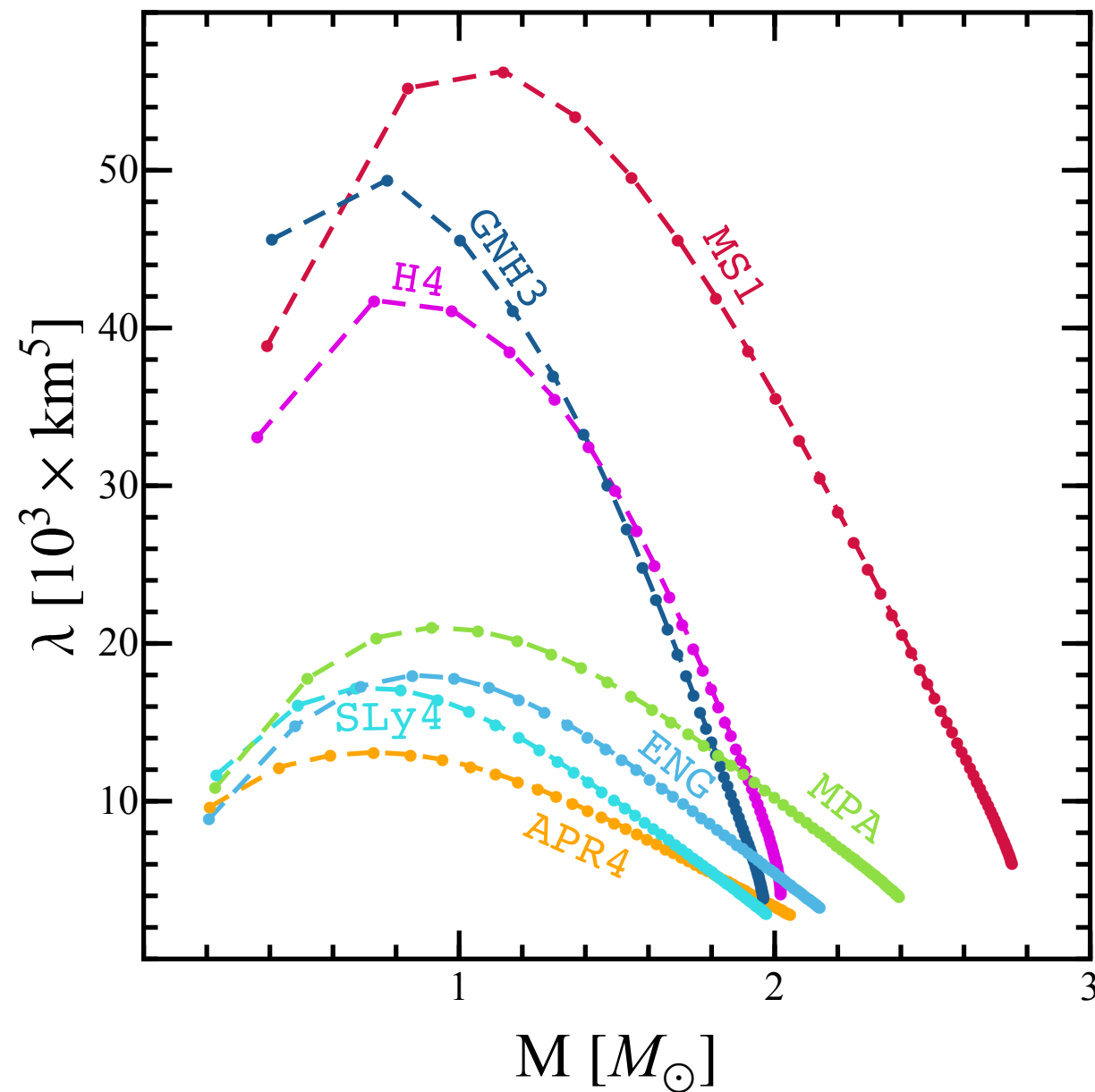


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# The recipe of Love

✓ Build  $\lambda - M$  profiles



● Neutron stars

$$M/R_{\text{NS}} \in [0.1 \div 0.2]$$



$$\lambda \neq 0$$

● Black holes

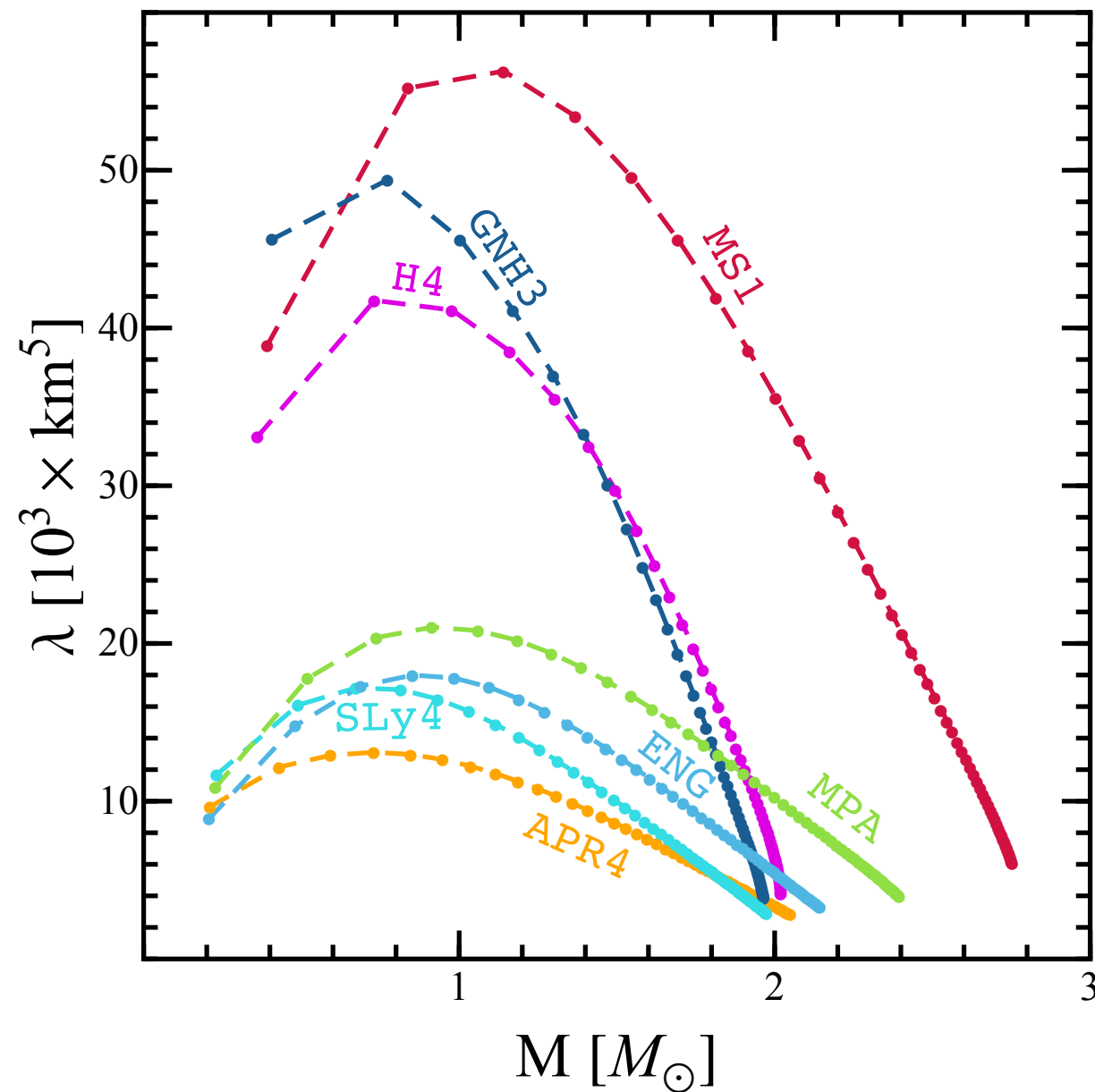
$$M/R_{\text{BH}} = 0.5$$



$$\lambda = 0$$

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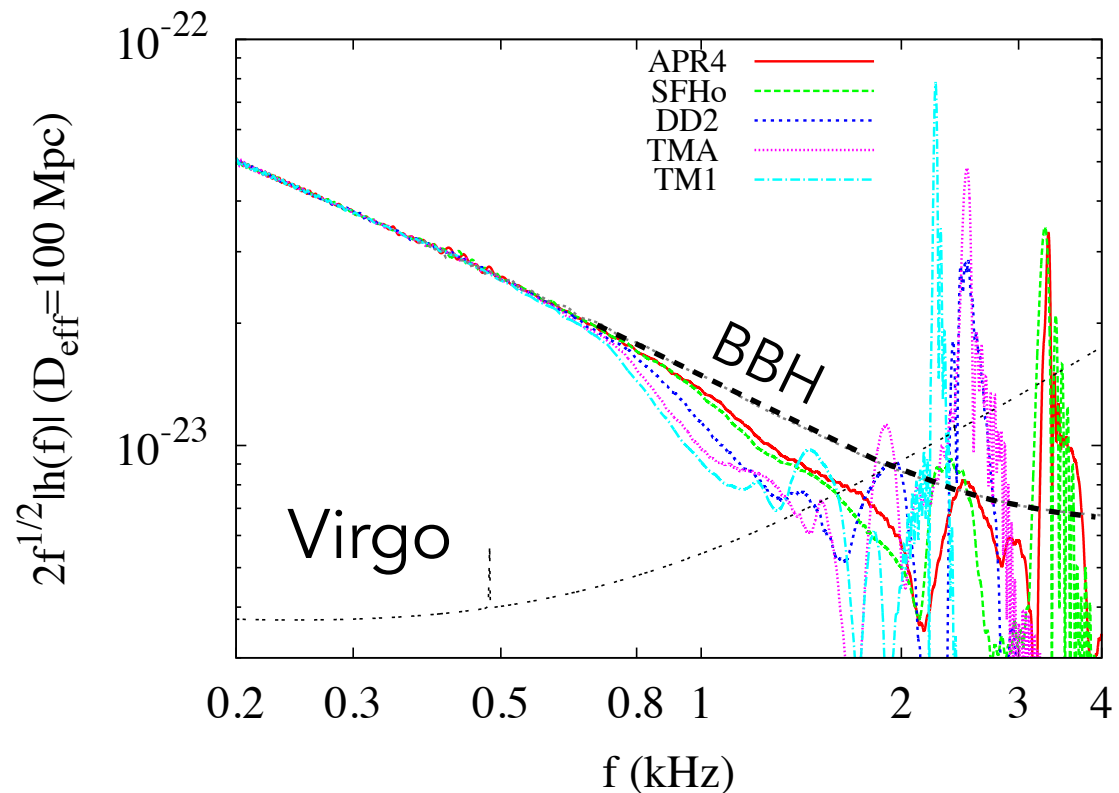


$$\lambda = 0$$

see Paolo's talk!

# How much Love?

We classify EoS for as **stiff** v.s. **soft**



EoS	stiffness	$R_{\text{NS}}$	$10^3 \times \lambda$
<b>APR4</b>	<b>very soft</b>	<b>11.09</b>	<b>10</b>
<b>SFHo</b>	<b>soft</b>	<b>11.91</b>	<b>13</b>
<b>DD2</b>	<b>medium soft</b>	<b>13.20</b>	<b>27</b>
<b>TMA</b>	<b>stiff</b>	<b>13.85</b>	<b>37</b>
<b>TM1</b>	<b>very stiff</b>	<b>14.48</b>	<b>45</b>

## Soft EoS

- ✓ larger densities
- ✓ smaller Love

## Stiff EoS


- ✓ smaller densities
- ✓ more Love
- ✓ larger effect in the signal

# Where do we look for Love?

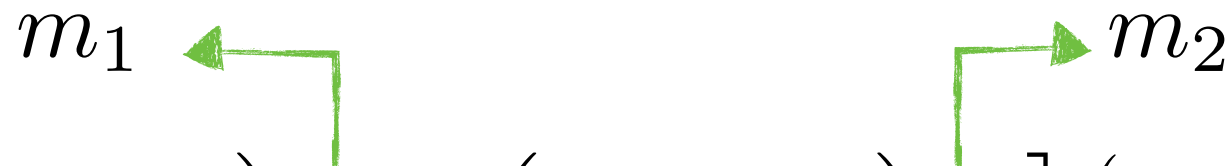
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Tidal effects add linearly  $h(f) = \mathcal{A}e^{i[\psi_{\text{PP}}(f) + \psi_{\text{T}}(f)]}$

$$\psi_{\text{PP}} \propto \left[ 1 - \left( \frac{743}{336} + \frac{11}{4}\eta \right) \frac{(m\pi f)^{2/3}}{c^2} + \frac{\dots}{c^3} + \dots + \frac{\dots}{c^7} \right]$$

  
Newt                      1 PN                      3.5 PN

$$\psi_{\text{T}} \propto \frac{1}{26} \left[ \left( 1 + 12\frac{m_2}{m_1} \right) \lambda_1 + \left( 1 + 12\frac{m_1}{m_2} \right) \lambda_2 \right] \frac{(m\pi f)^{10/3}}{c^{10}} + \frac{\dots}{c^{12}}$$



# Where do we look for Love?

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↑
↑
↑  
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$m_1$  ←
→  $m_2$   
↑
↑  
 5 PN: small term!

# Where do we look for Love?

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 ↑ 3.5 PN
 ⋈

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←  $m_1$ 
←  $m_2$

↑  $\mathcal{O}(10^4)$ 
↑ 5 PN: small term!
 ⋈



# Where do we look for Love?

Tidal effects add linearly  $h(f) = \mathcal{A}e^{i[\psi_{\text{PP}}(f) + \psi_{\text{T}}(f)]}$

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↑ Newt
 ↑ 1 PN
 ↑ 3.5 PN
 ⋈

$$\psi_{\text{T}} \propto \frac{1}{26} \left[ \left( 1 + 12\frac{m_2}{m_1} \right) \lambda_1 + \left( 1 + 12\frac{m_1}{m_2} \right) \lambda_2 \right] \frac{(m\pi f)^{10/3}}{c^{10}} + \frac{\dots}{c^{12}}$$

←  $m_1$ 
→  $m_2$



$\Lambda$   $\mathcal{O}(10^4)$

average Love

# Status of the art

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## Theoretical modelling

- ✓ We know the full form of the  $\frac{1}{c^{10}}$  and  $\frac{1}{c^{12}}$  contribution to the GW phase

*Vines et al, 2011*

- ✓ Rotational corrections to the Love number are under investigation

*Pani et al., 2015*

*Ask Tiziano*

- ✓ First attempts to model time-dependent corrections to the Love number

*Maselli et al., 2013*

*Hinderer et al., 2016*

- ✓ Excitations of neutron star oscillatory modes due to absorption of the tidal energy

*Weinberg, 2015*

# Before GW170817

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Detectability of  $\lambda$  strongly depends on the NS mass & stiffness

Soft EoS

- ✓ smaller deformations
- ✓ deep inspiral to look for
- ✓ weaker bounds

Stiff EoS

- ✓ larger deformations
- ✓ affects a wider frequency range
- ✓ tighter constraints

		EoS	$R_{\text{NS}}^{1.4}$	$\sigma_{\lambda}/\lambda$	
	<b>very soft</b>	<b>APR4</b>	<b>11.30</b>	<b>~60%</b>	[100 Mpc]
LIGO/Virgo $\rightarrow$	<b>medium</b>	<b>H4</b>	<b>13.99</b>	<b>~20%</b>	
	<b>very stiff</b>	<b>MS1</b>	<b>14.95</b>	<b>~10%</b>	

# Before GW170817

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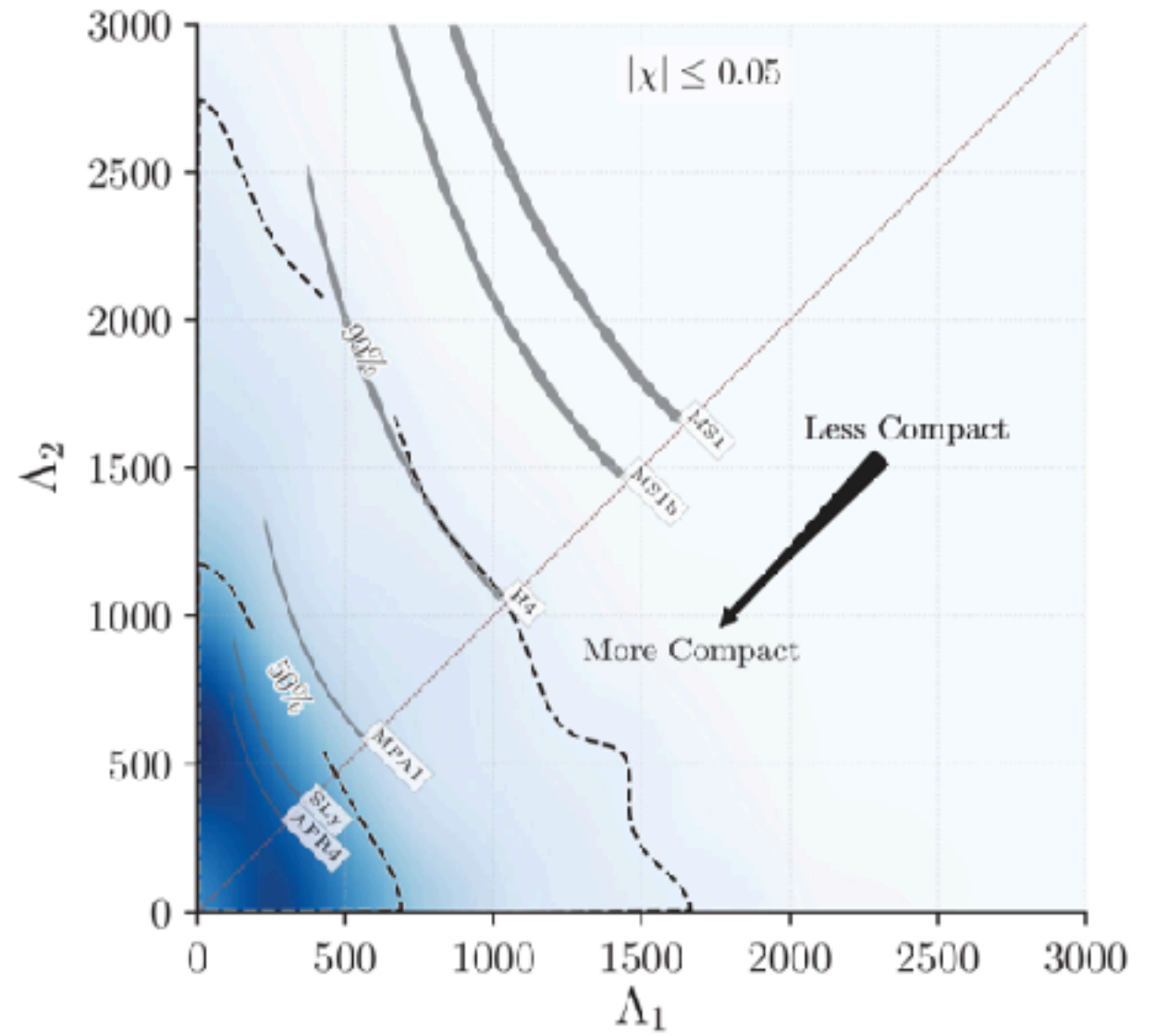
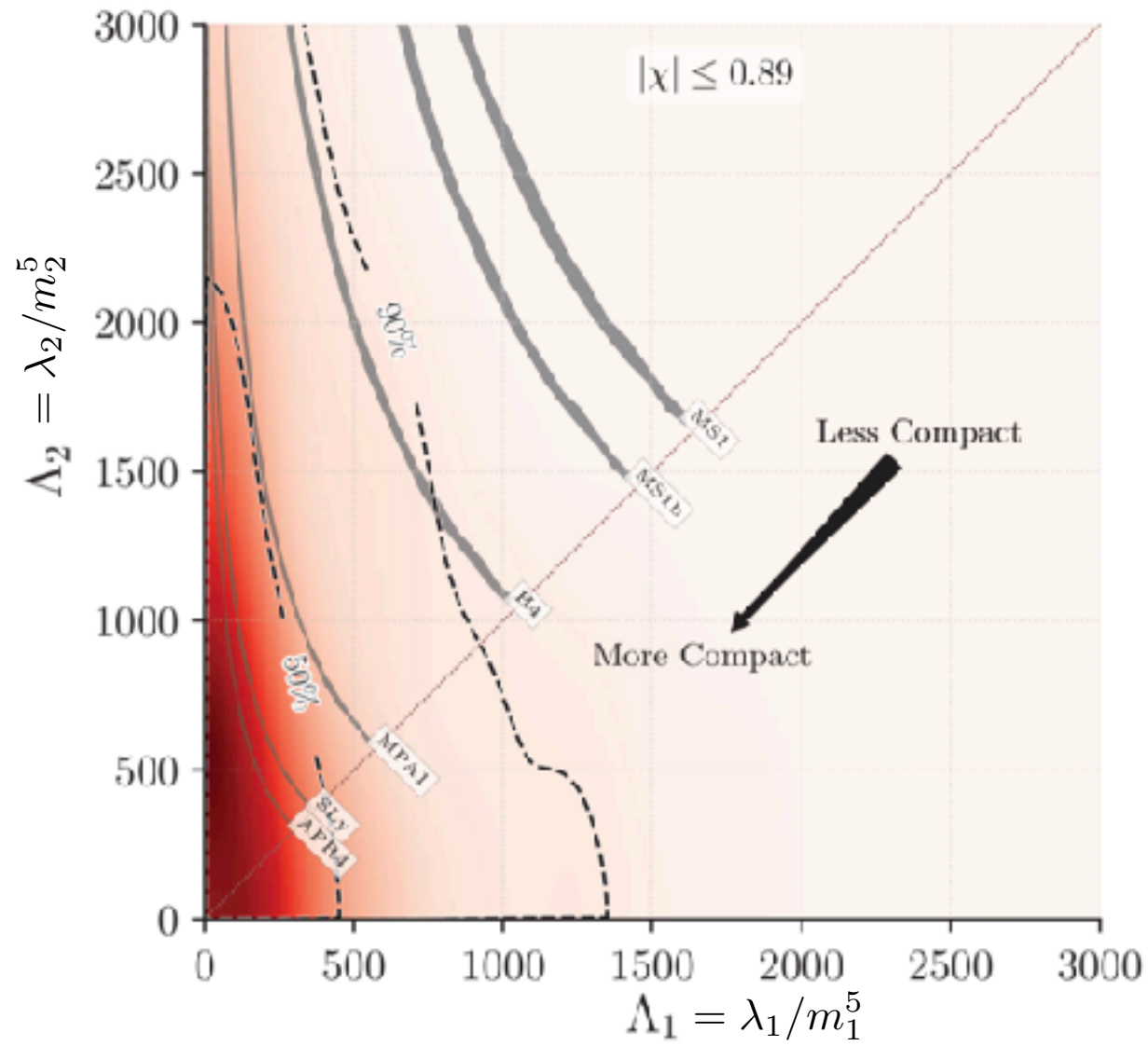
Massive efforts to determine real chances to measure  $\lambda$

*Flanagan et al, 2008; Hinderer et al., 2010; Damour et al, 2012;  
Maselli et al, 2013; Read et al 2013; Favata 2014; Hotokezaka, 2016*

- ✓ For favourable EoS LIGO/Virgo can constrain the Love number and then the NS radius with **1 km** up 200 Mpc
- ✓ A third generation detector, like the Einstein Telescope would improve of an order of magnitude,  $R_{\text{NS}} \sim 100$  m
- ✓ For softer EoS more likely to distinguish from a stiff one, than to actually determine the Love number

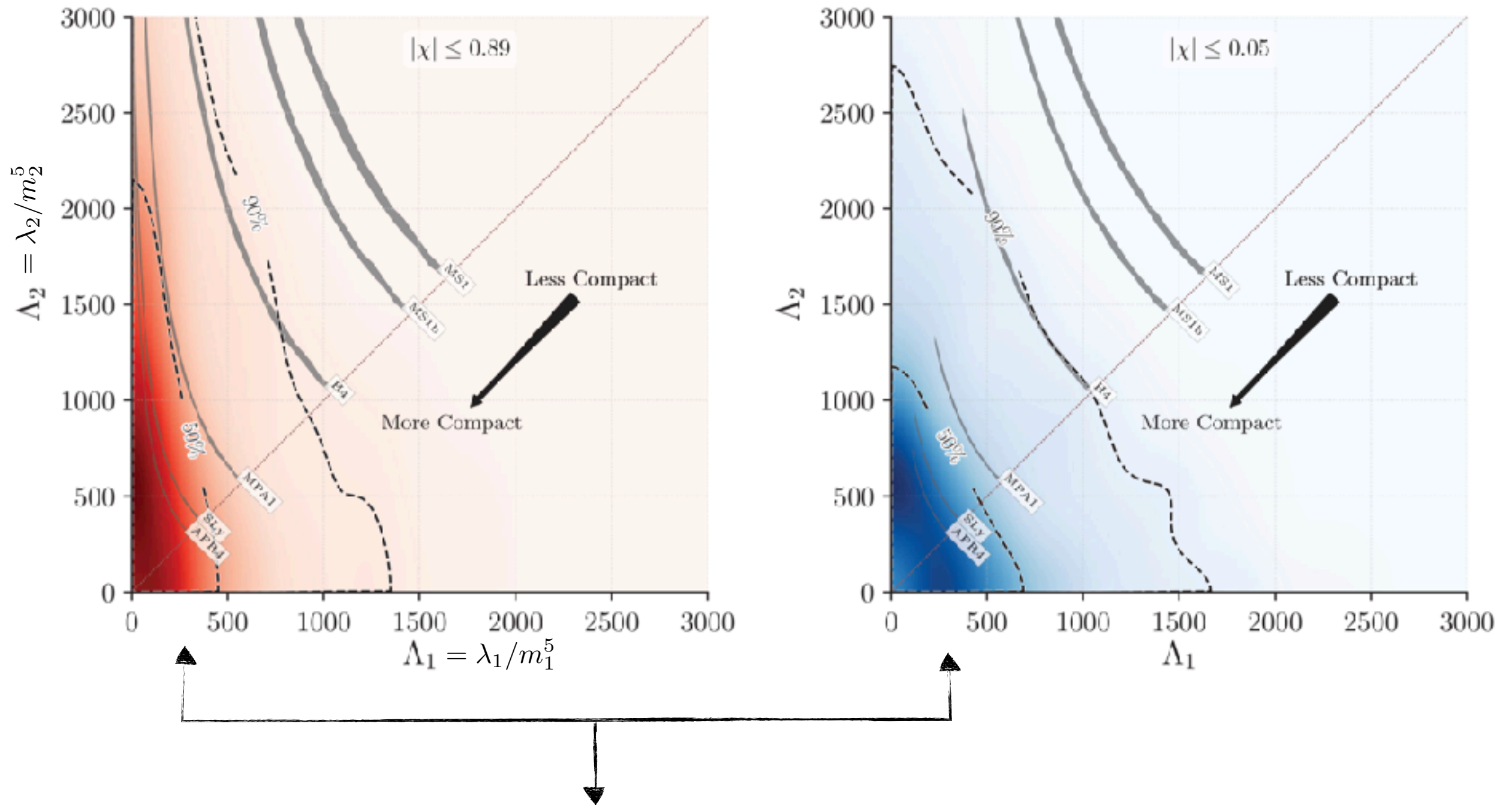
# GW170817

Abbott et al., 2017



# GW170817

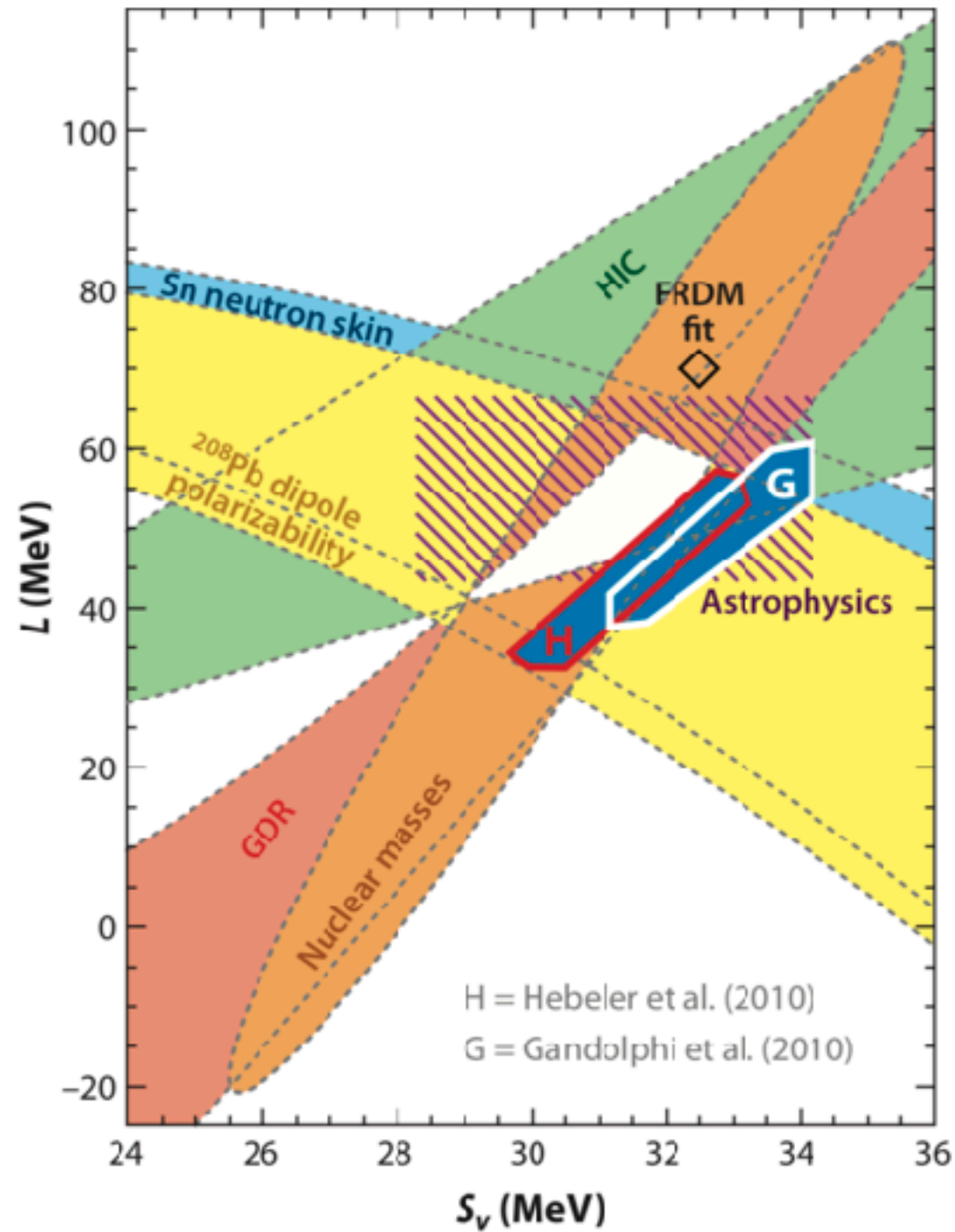
Abbott et al., 2017



First detection seems to favour soft nuclear matter

# From macro to micro

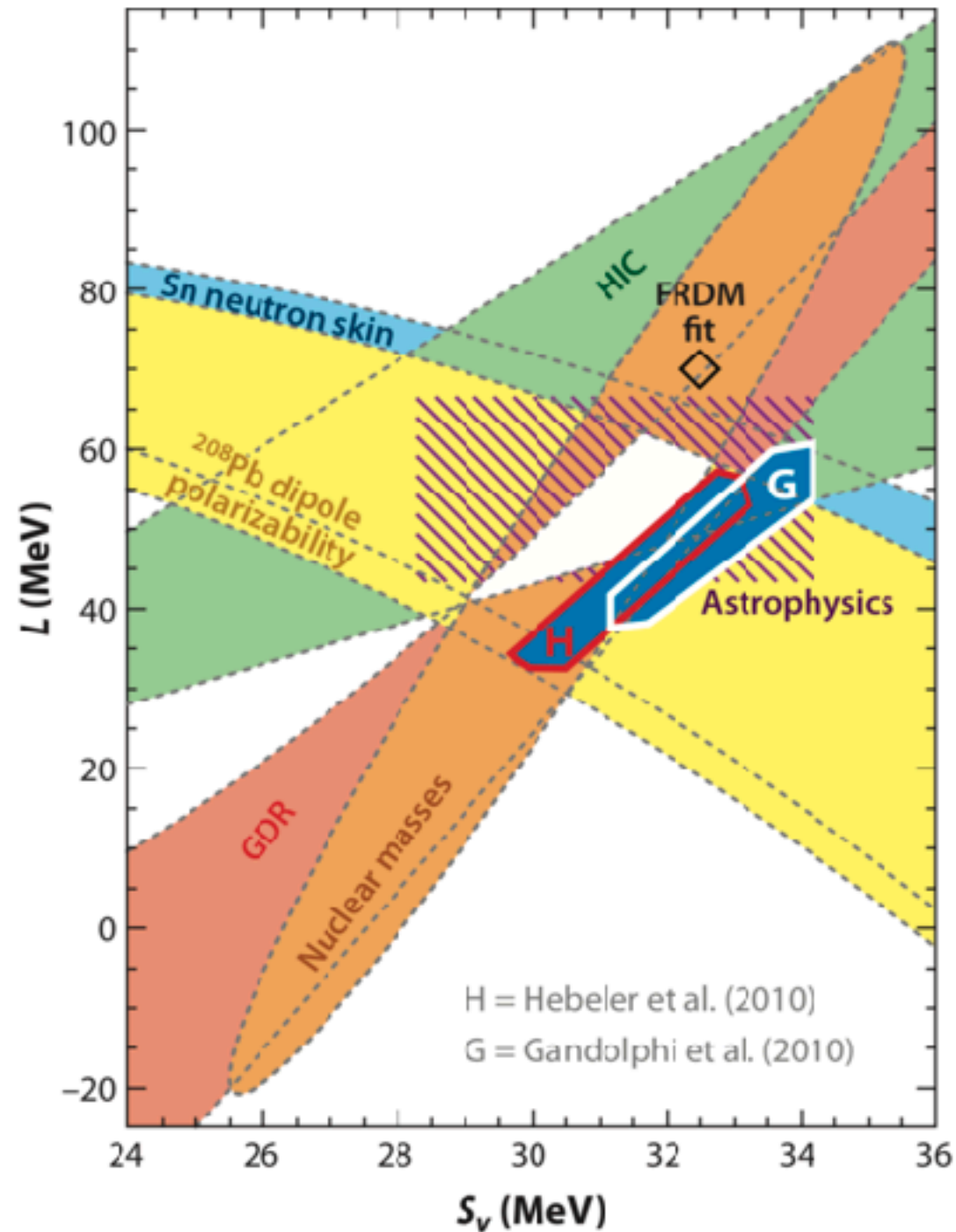
So far for electromagnetic signals





# From macro to micro

So far for electromagnetic signals

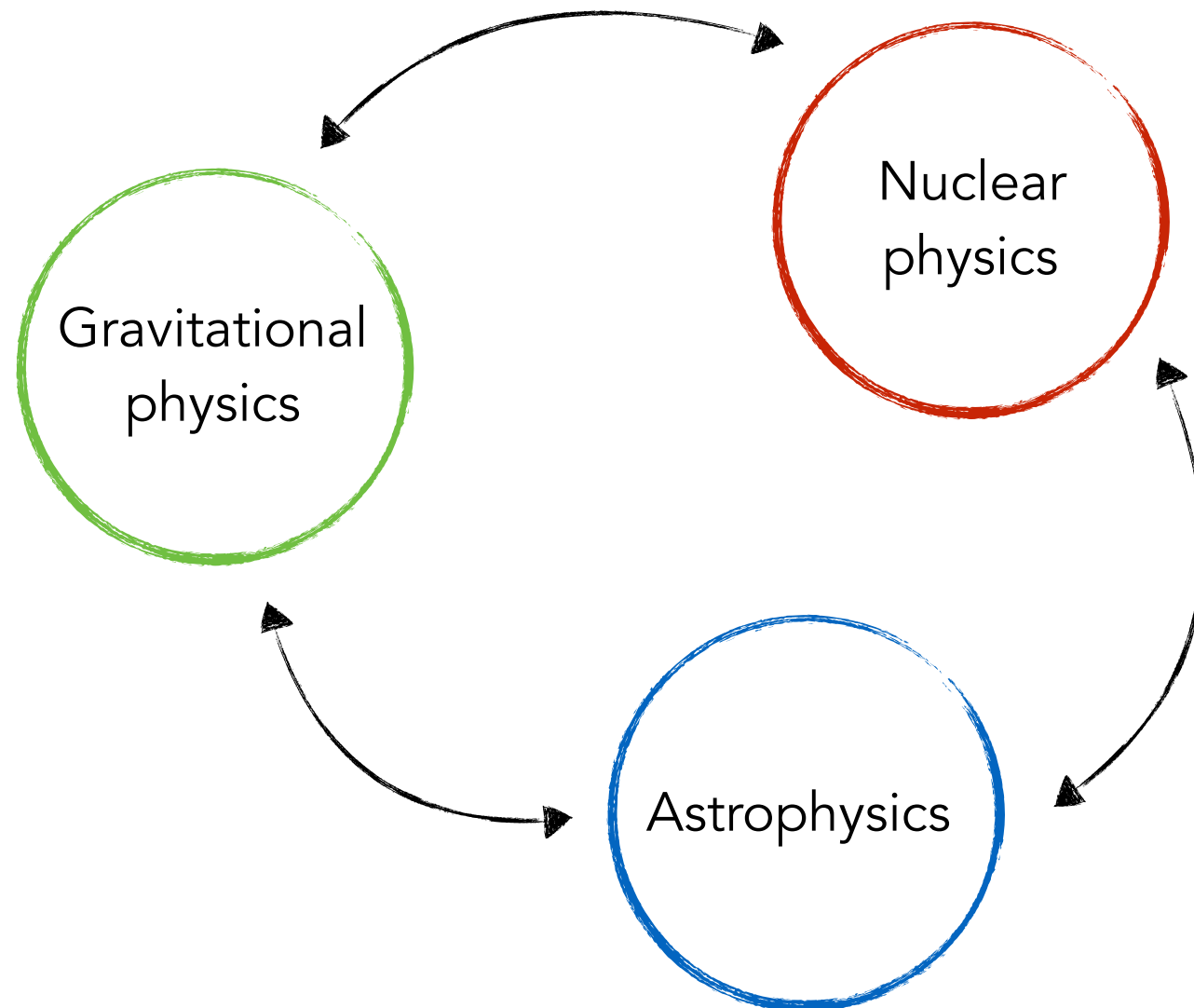


New constraints coming from **GW** and the **Love numbers**



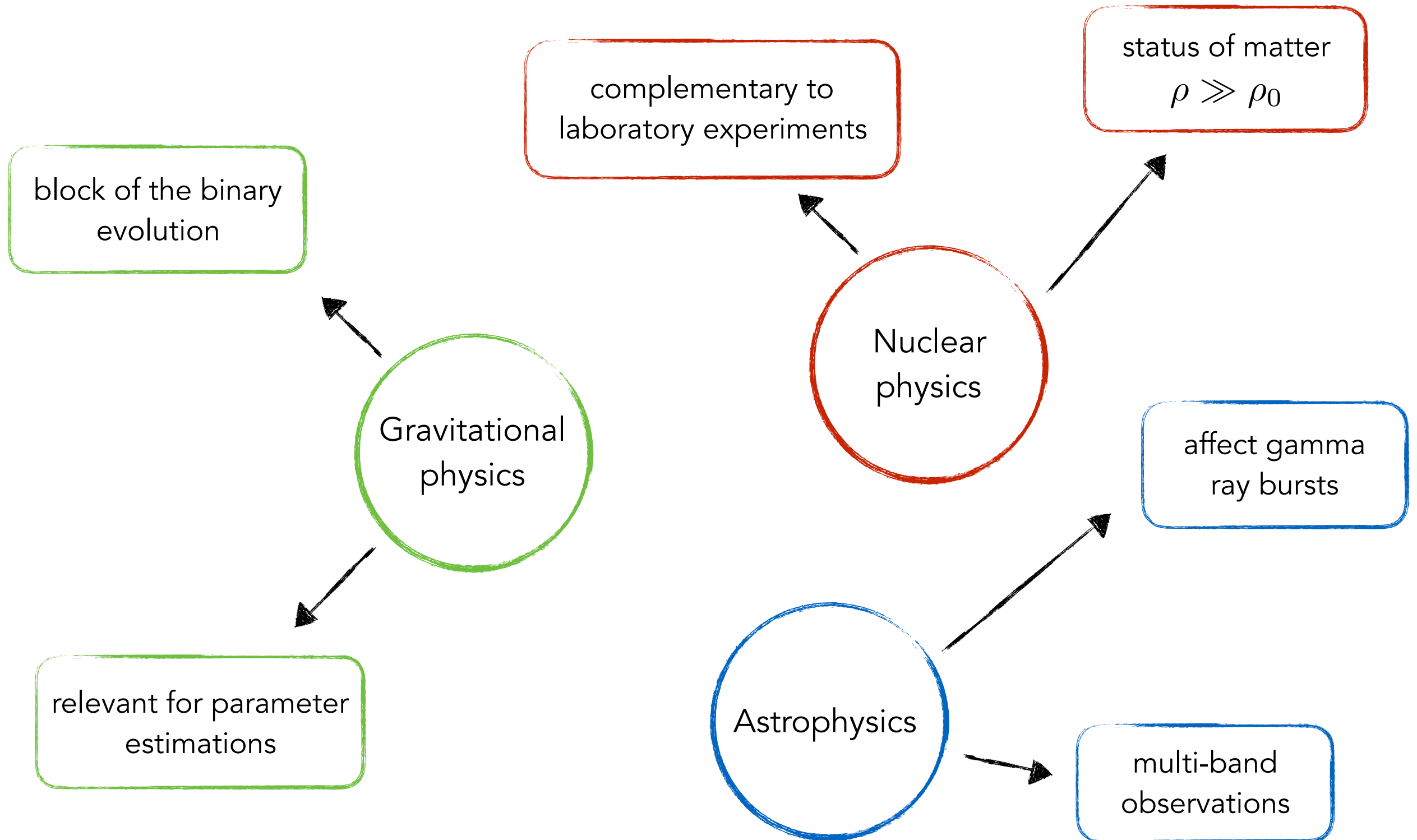
# Summary

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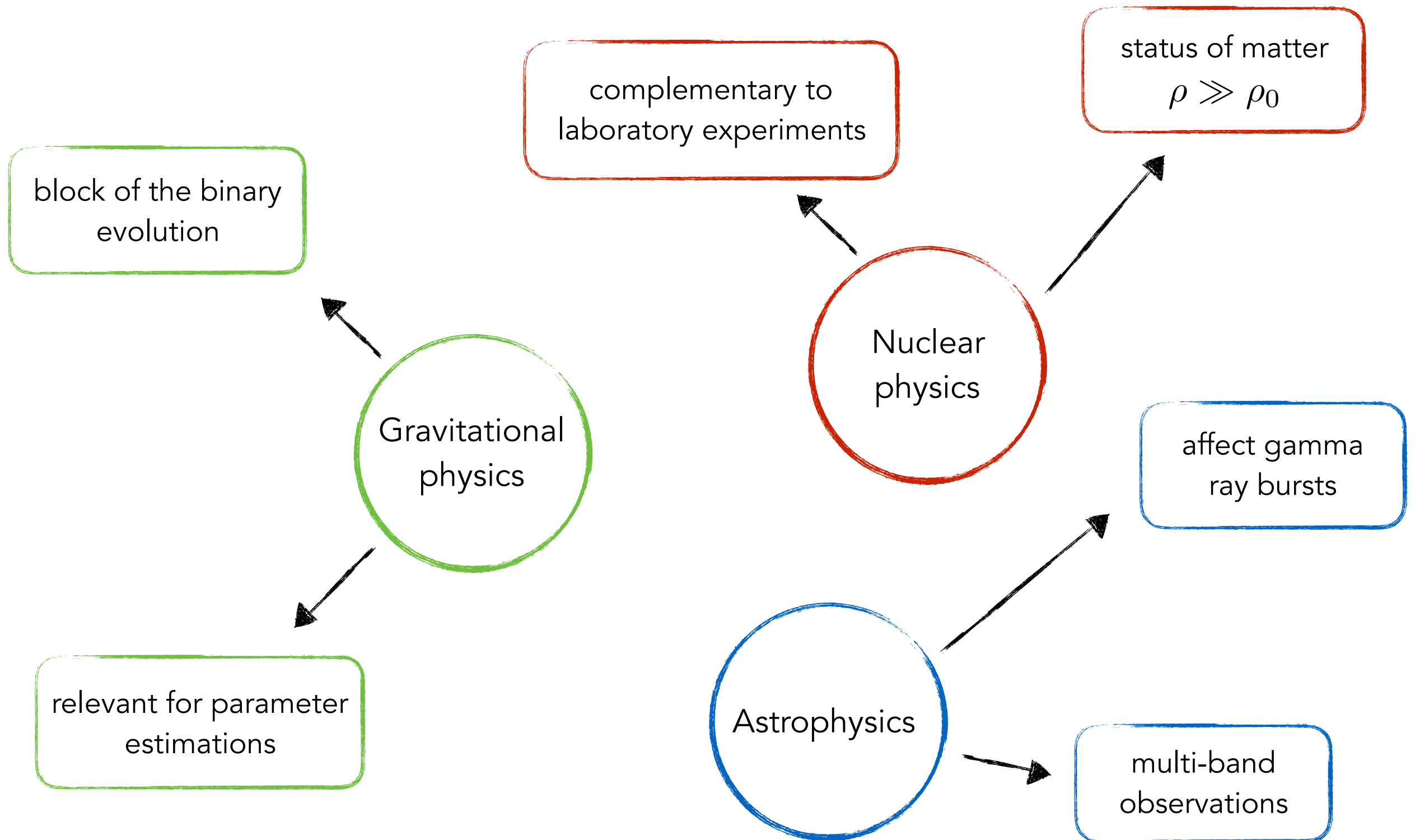


# Summary

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# Summary: take care of Love



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*Backup*

# Observed NS masses

Ozel & Freire, 2017

