

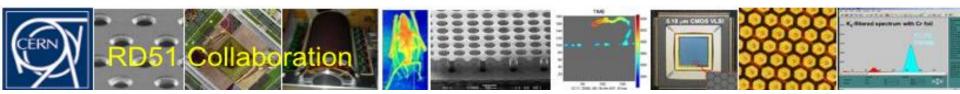
The micro-RWELL

M. Poli Lener¹

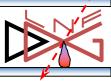
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- 2. CERN
- 3. Kobe University



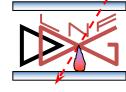
Muon Detectors for future colliders



- □ The future colliders (CepC, SppC and FCC hh) requires for extremely large muon detectors :
 - $\hfill\square$ ~10000 m² in the barrel
 - $\hfill\square$ 3-5000 m^2 in the endcap
 - □ 300 m² in the very forward region
- □ The detectors have to be operated in high background (very large uncertainties depending on shielding, actual structure, etc.):
 - □ O(1 10 kHz/cm²) in the barrel
 - O(10 100 kHz/cm²) in the end-cap
 - **O(1 MHz/cm²)** in the forward region
- Taking into account the surface and the expected rates gaseous detectors and in particular MPGDs is the natural solution (straight-forward for CepC, requiring an R&D for the harsher conditions of SppC & FCC-hh)
- **R&D for HL-LHC** (LHCb phase-2 muon upgrade) is clearly a good starting point
- Last but not least, cylindrical or planar MPDG, could be adopted to equip tracking devices



The **µ-RWELL**



The R&D on μ -RWELL is mainly motivated by the wish of improving the

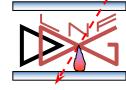
- stability under heavy irradiation (discharge suppression)
- construction technology (simplifying the assembly)
- Technology Transfer to industry (mass production)

a MUST for **very large scale applications** in fundamental research at the future colliders as well as for technology dissemination beyond HEP

The original idea was conceived in 2009 @ LNF during the construction of the CGEM, to try to find a way to simplifying as much as possible the construction of the CGEM and its toolings. Only in the 2014 we really started a systematic study of this new technology in collaboration with Rui de Oliveira

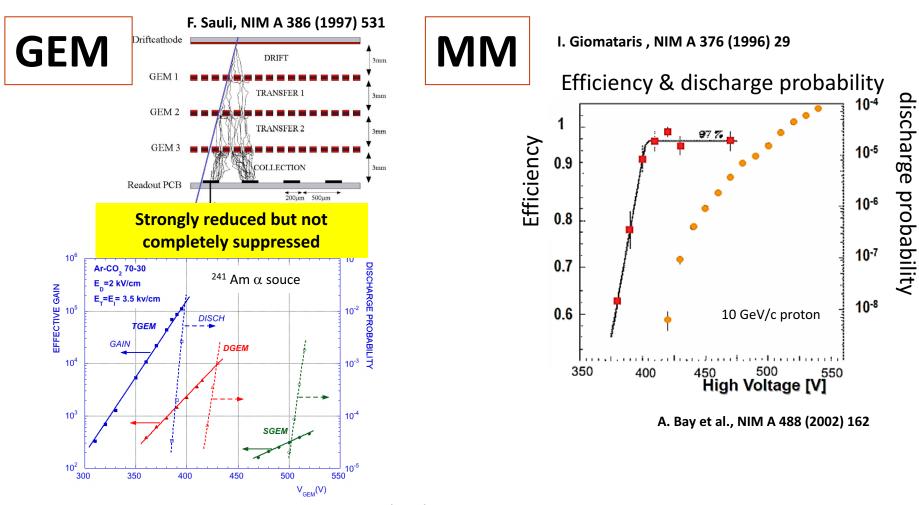


MPGDs: stability



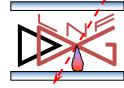
The **biggest "enemy"** of MPGDs are the **discharges.**

Due to the **fine structure** and the **typical micrometric distance of their electrodes**. MPGDs generally suffer from **spark occurrence** that can eventually **damage the detector and the related FEE**.

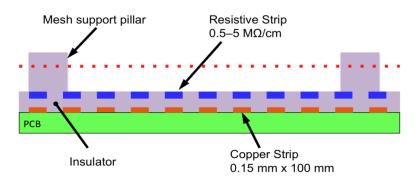




Technology improvements for MicroMegas

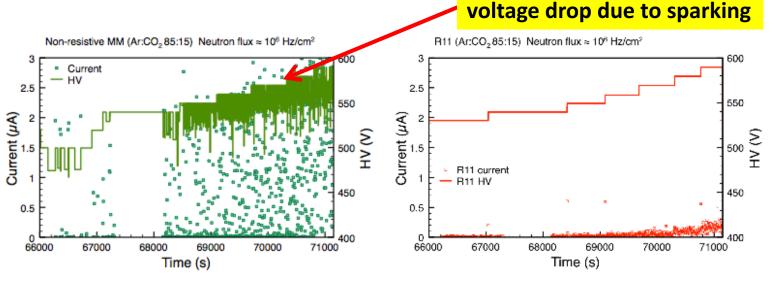


For MM, the spark occurrence between the metallic mesh and the readout PCB has been overcome with the implementation of a "resistive layer" on top of the readout. The principle is the same as the resistive electrode used in the RPCs: the transition from streamer to spark is strongly suppressed by a local voltage drop.



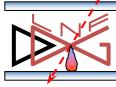
by R.de Oliveira TE MPE CERN Workshop

The resistive layer is realized as resistive strips capacitive coupled with the copper readout strips.



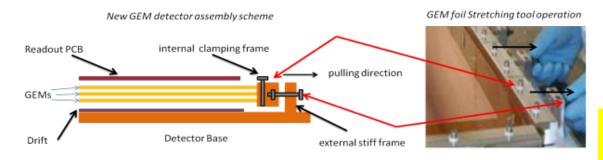


MPGDs: construction issues (I)



An **important limitation of such MPGDs** is correlated with the **complexity of their assembly procedure**, particularly evident in case of **large area devices**.

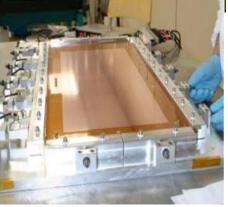
□ The construction of a GEM chamber requires time-consuming assembly steps such as the stretching (with quite large mechanical tension to cope with – 1 kg/cm) and the gluing of the GEM foil on frames





NS2(CERN – R. de Oliveira): no gluing, but still stretching

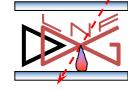
- A 2 m long detector requires a ~200 kg mechanical tension that must be sustained by stiff mechanical structures (large frames, rigid panels ...). While the max width of the raw material is about 60 cm.
- ❑ The splicing/joining of smaller detectors in order to realize large surfaces (as used for silicon detectors) is difficult unless introducing not negligible dead zones.



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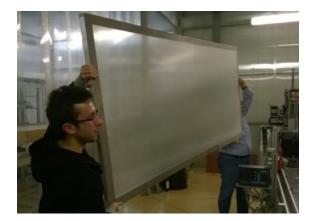


MPGDs: construction issues (II)



Similar considerations hold for MM:

- ❑ the splicing/joining of smaller PCBs is possible, opening the way towards the large area detection covering
- □ the **fine metallic mesh**, that defines the amplification gap, is a "*floating component*", because it is stretched on the cathode (@ 1 kg/cm) and electrostatically attracted toward the PCB ($P = \epsilon_0 \times (\Delta V/_d)^2$).





- □ this could be a source of instability because a "not well defined" amplifying gap could generate gain non-uniformity.
- In addition: handling of large meshes is clearly "not trivial" (of course for large area) cleaning of the amplification stage could be an issue

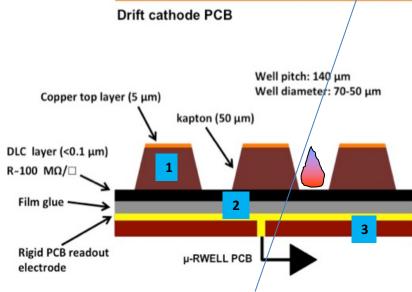
\mathbf{V}_{INF} The μ -RWELL: the detector architecture

The μ-RWELL is composed of only two elements: the μ-RWELL_PCB and the cathode

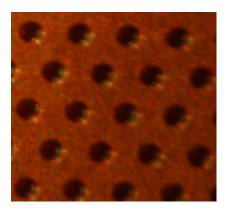
The **µ-RWELL_PCB**, the core of the detector, is realized by coupling:

- 1. a WELL patterned kapton foil as amplification stage
- 2. a **resistive layer**^(*) for discharge suppression:
 - Single resistive layer (SRL) <100 kHz/cm²: single resistive layer → surface resistivity ~100 MΩ/□ (SHiP; CepC, Novosisbirsk, EIC, HIEPA)
 - ii. Double resistive layer (DRL) >1 MHz/cm²: for LHCb-Muon upgrade
- 3. a standard readout PCB

(*) DLC = Diamond Like Carbon highly mechanical & chemical resistant



G. Bencivenni et al., 2015_JINST_10_P02008



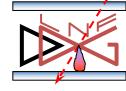




Principle of operation

top copper

layer



HV

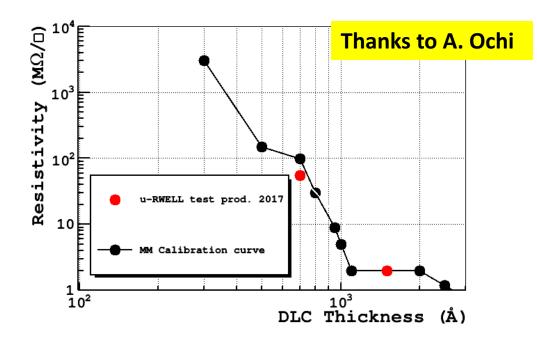
Applying a suitable voltage between **top copper layer and DLC** the "WELL" acts as multiplication channel for the ionization.

The charge induced on the resistive foil is dispersed with a *time constant*, $\tau = \rho C$, determined by kapton resistive stage Insulating medium Pad/strip r/out Not in scale

- the *surface resistivity,* ρ
- the capacitance per unit area, which depends on the distance between the resistive foil and the pad/strip readout plane, t
- the *dielectric constant* of the insulating medium, \mathcal{E}_r [M.S. Dixit et al., NIMA 566 (2006) 281]
- The main effect of the introduction of the resistive stage is the suppression of the transition from streamer to spark
- As a drawback, the capability to stand high particle fluxes is reduced, but an appropriate grounding of the resistive layer with a suitable pitch solves this problem (see High Rate scheme)



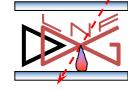
- The **kapton foils (copper etched on one side)** are at moment sputtered with DLC, provided by Be-Sputter Co., Ltd. in Japan. They are able to sputter 6 foils (1.2x0.6 m²) per production batch.
- The resistivity depends on several manufacturing conditions, but can be parametrized as function of the DLC thickness. The resistivity uniformity is at level of 10-20%



Recently we are starting a profitable collaboration with Zhou Yi and Jianbei Liu from USTC – Hefei for the manufacturing of improved DLC foil



Main detector features



The **μ-RWELL** is a **single-amplification stage**, intrinsically **spark protected** MPGD characterized by:

- simple assembly procedure:
 - only two components $\rightarrow \mu$ -RWELL_PCB + cathode
 - no critical & time consuming assembly steps:
 - no gluing
 - *no stretching* (→ no stiff & large frames needed)
 - easy handling
 - suitable for large area with PCB splicing technique w/small dead zone

cost effective:

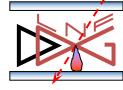
• 1 PCB r/o, 1 μ-RWELL foil, 1 DLC, 1 cathode and very low man-power

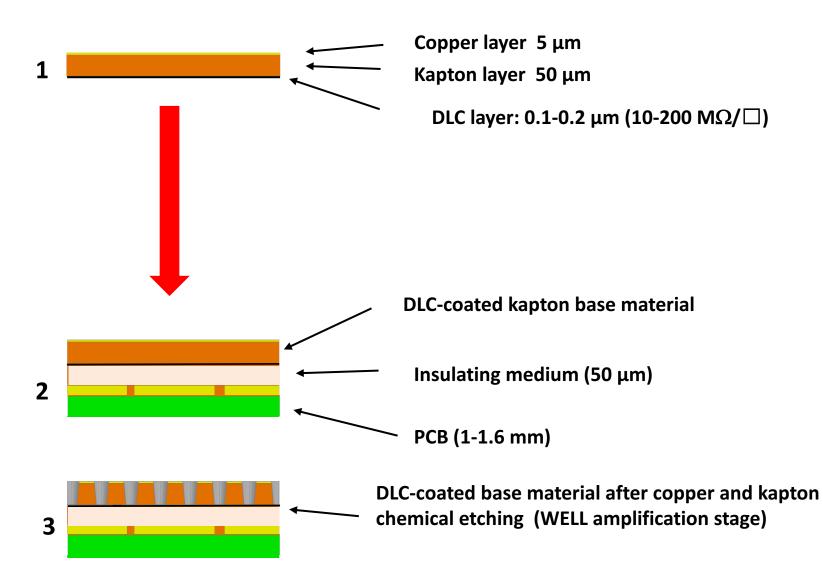
easy to operate:

very simple HV supply → only 2 independent HV channels or a trivial passive divider (while 3GEM detector → 7 HV floating/channels)

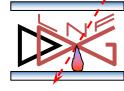


The Low Rate scheme





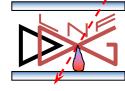




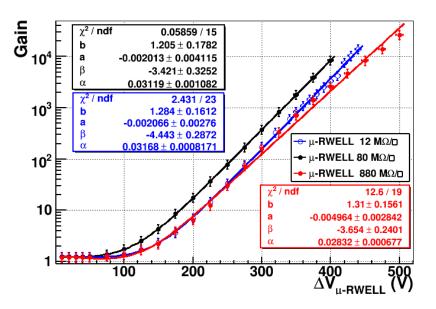
Single Resistive µ-RWELL performance



Detector Gain



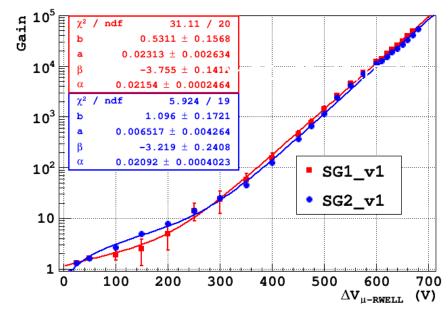
Ar/iC₄H₁₀= 90/10



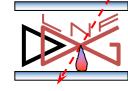
Some recent prototypes achieved a Gain ~ 10^5 in Ar/CO₂/CF₄= 45/15/40

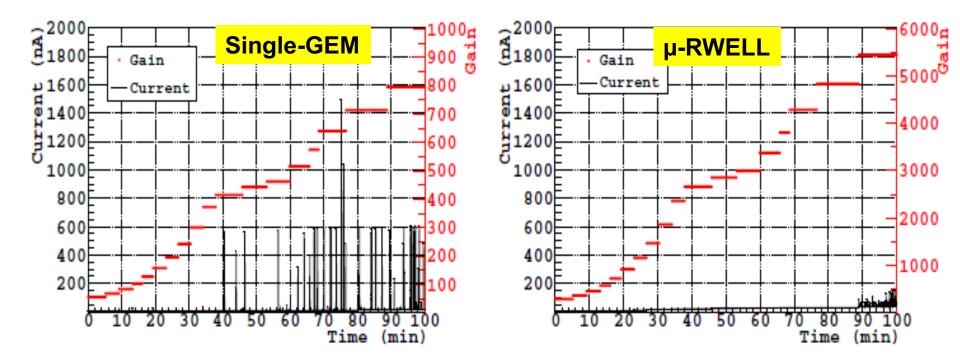
Prototypes with different resistivity have been tested with X-Rays (5.9 keV), with $Ar/iC_4H_{10}=90/10$ gas mixture, and characterized by measuring the gas gain in current mode.

Ar/CO₂/CF₄= 45/15/40







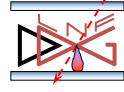


- discharges for <u>μ-RWELL</u> of the order of <u>few tens of nA (<100 nA @ high gain)</u>
- for <u>GEM</u> discharges the order of $\underline{1\mu A}$ are observed at high gas gain

25/05/2018



Beam Tests results



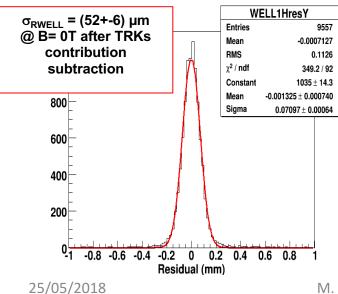
H4 Beam Area (RD51) Muon beam momentum: 150 GeV/c Goliath: B up to 1.4 T

GEMs Trackers

BES III-GEM chambers

μ-RWELL prototype:

- 12-80-880 MΩ /□
- 400 µm pitch strips
- APV25 (CC analysis)
- Ar/iC₄H₁₀ = 90/10

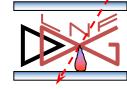


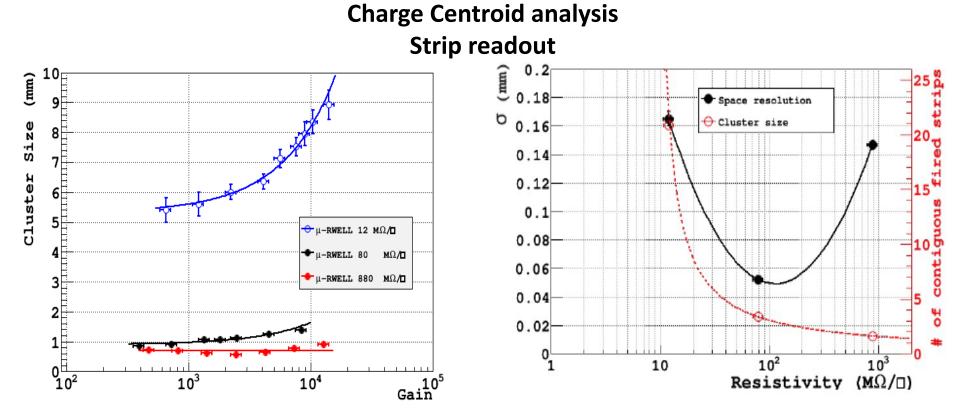
GOLIATH

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Space resolution vs resistivity





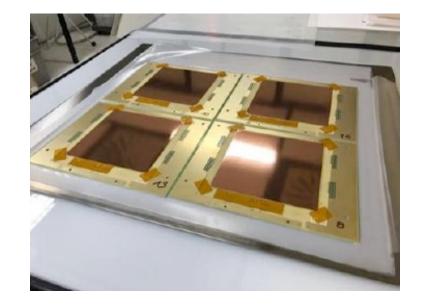
The space resolution exhibits a minimum around 100M Ω / \Box

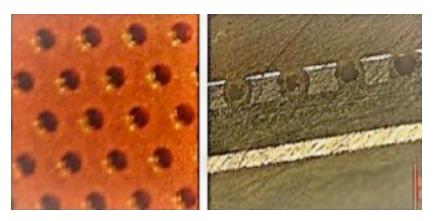
- \rightarrow at **low resistivity** the **charge spread increases** and then σ is worsening
- → at high resistivity the charge spread is too small (Cl_size → 1 fired strip) then the Charge Centroid method becomes no more effective (σ → pitch/ $\sqrt{12}$)

Technology Transfer to Industry (I)

The engineering and industrialization of the μ -RWELL technology is one of the main goal of the project. Transferring the manufacturing process to industry will allow a cost-effective mass production: a must for the construction of muon systems at future HEP Colliders

Manufacturing process of the single resistive layer has been already tested at the ELTOS SpA (http://www.eltos.it)





Production @ ELTOS:

- 8 PCB uRWELL (PAD r/o)
- 16 PCB uRWELL (strip r/o)

coupled with kapton/DLC foils.

The etching of the kapton done by Rui

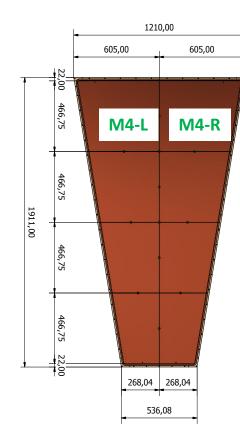
Technology Transfer to Industry (II)

In the framework of the **CMS-phase2 muon upgrade** different prototypes of **large size singleresisitive layer μ-RWELLs** has been built at ELTOS:

- 1.2x0.5m² μ-RWELL

- 1.9x1.2m² μ-RWELL



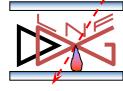


In collaboration with CMS-Muon: L. Benussi, L. Borgonovi, P. Giacomelli, A. Ranieri, M. Ressegotti, I. Vai, V. Valentino





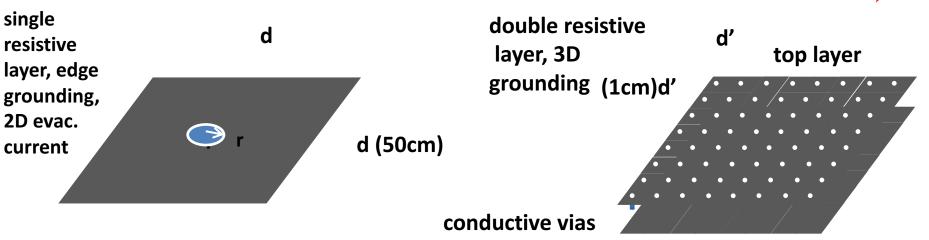




From Low Rate μ-RWELL to High Rate version



Towards the High Rate



bottom layer

(*) point-like irradiation, $r \ll d$ Ω is the resistance seen by the current generated by a radiation incident the center of the detector cell

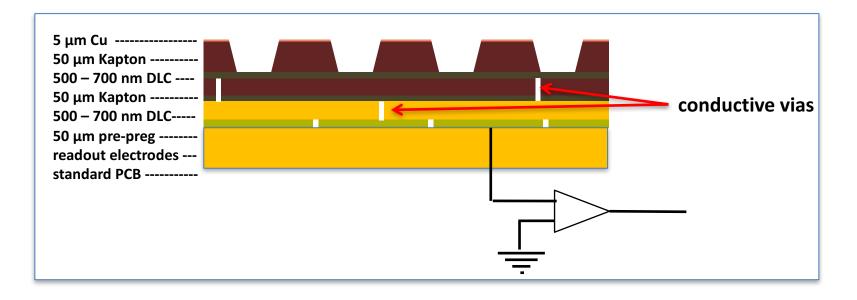
$$\begin{split} \Omega &\sim \rho_{s} \ge d/2\pi r & \Omega' &\sim \rho_{s}' \ge 3d'/2\pi r \\ \Omega &\cap \Omega' &\cap (\rho_{s} / \rho_{s}') \ge d/3d' \\ & \text{If } \rho_{s} = \rho_{s}' \xrightarrow{>} \Omega / \Omega' &\sim \rho_{s} / \rho_{s}' \ge d/3d' = 50/3 = 16.7 \end{split}$$

(*) Morello's model: appendix A-B (G. Bencivenni et al., 2015_JINST_10_P02008)

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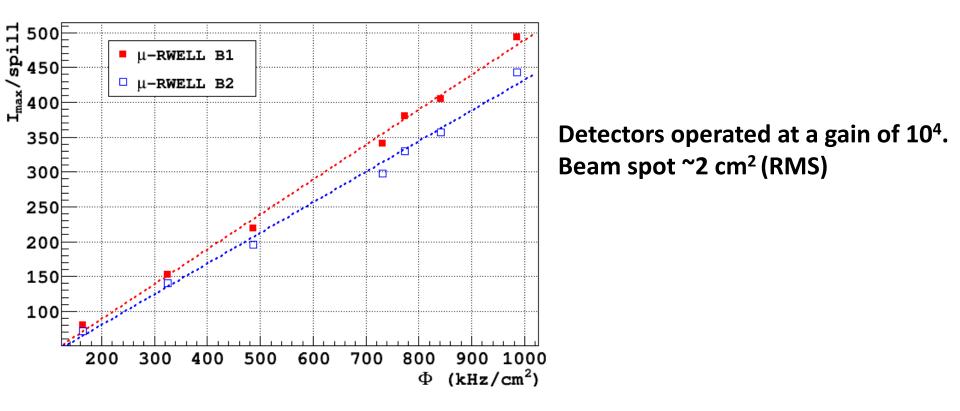


The idea is to reduce the path of the current on the DLC, a matrix of conductive vias connecting the two resistive layers is introduced. Another matrix of vias connects the second resistive layer to ground through the readout





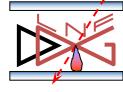
Rate capability as a function of the pion beam (H4-SpS CERN) intensity



WARNING: The engineering/industrialization of the double-resistive layer is difficult due to the manufacturing of the conductive vias



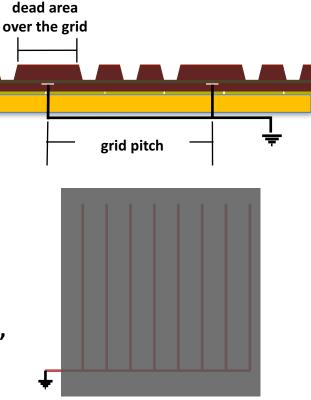
New ideas for the HR version



The aim is **to maintain a very short path for current moving on the resistive layer**, while **simplifying the construction process.**

Two ideas are now under development: **silver grid and resistive grid on the bottom of the DLC**

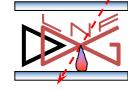
High Rate scheme	Resistivity [MΩ/]	Dead Area over grid	Grid Pitch	Geometrical efficiency [%]	Туре
Silver Grid 1 (SG1)	60-70	2 mm	6 mm	66	conductive grid
Silver Grid 2 (SG2)	60-70	1,2 mm	12 mm	90	conductive grid
Resistive Grid (RG)	60-70	-	6 mm	Full	resistive grid



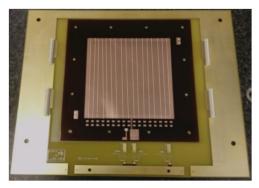
The conductive grid on the bottom of the amplification stage can induce instabilities due to discharges over the DLC surface, requiring for the introduction of a dead zone on the amplification stage. This is not the case for the resistive grid scheme.



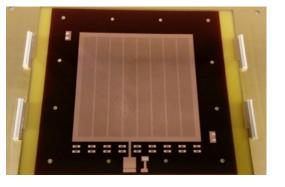
New ideas for the HR version



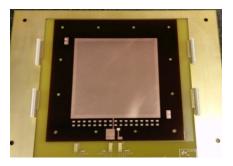
Silver Grid v1

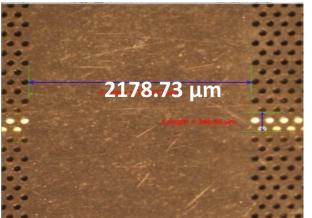


Silver Grid v2

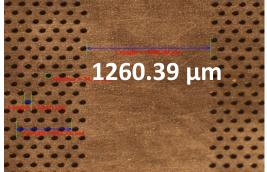


Resistive Grid





grid-pitch 6 mm dead area 2 mm → Geometrical effi.: 66%



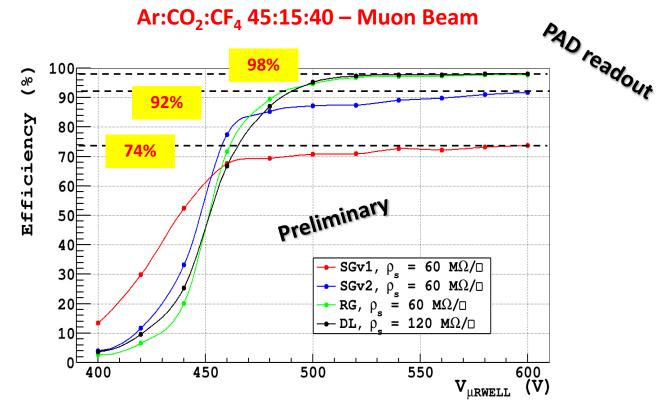
grid-pitch 12 mm dead area 1 mm → Geometrical effi.: 90%

grid-pitch 6 mm NO dead area → Geometrical effi.: full



HR layouts performance: efficiency





As expected, the **RG & DL prototypes** reach full tracking efficiency ~ **98%** (NO DEAD AREA in the amplification stage is present).

The **SG1 & SG2** show lower efficiency (**74% - 92%**) BUT higher than their geometrical acceptance (**66% and 90% respectively**), thanks to the **efficient electron collection mechanism** of the combined drift/amplification electric field, **reducing the effective dead zone**.





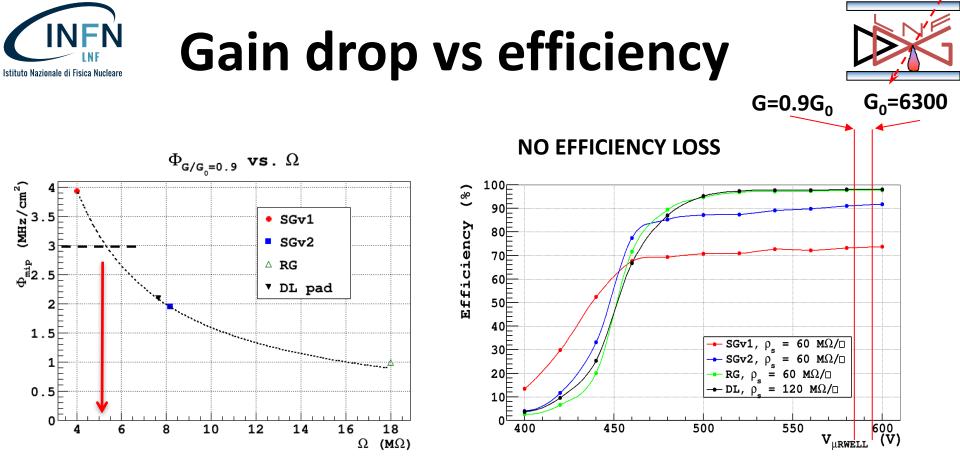
Ar:CO₂:CF₄ 45:15:40, G₀=6300, Spot size = 38.5 mm G/G 1.05 0.95 10% 0.9 SGv1 #1, ρ_=60 MΩ/⊡ 0.85 SGv2 #1, ρ_=60 MΩ/□ 0.8 RG #1, ρ_=60 MΩ/□ 0.75 DL pad #1, $\rho = 120 M\Omega/\Box$ 0.7 10⁵ 10⁴ (Hz/cm^2) Φ

The gain drop is due to the Ohmic effect on the resistive layer: charges collected on the DLC drift towards the ground facing an effective resistance Ω , depending on the evacuation scheme geometry and DLC surface resistivity. Ω is computed by the parameter p_0 coming from

the fit of the Gain curve.

$$rac{G}{G_0} = rac{-1 + \sqrt{1 + 4p_0 \Phi}}{2p_0 \Phi}$$

where $p_0 \propto \Omega$ (effective resistance)



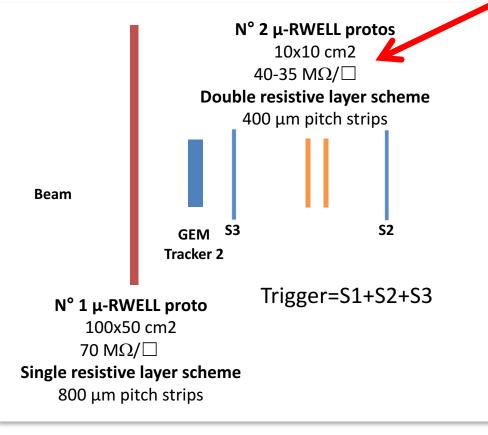
The primary ionization of 5.9 keV is ~7 times larger than the one created by a m.i.p. In order to stand a 3 MHz/cm² m.i.p. fluence, accepting a 10% gain drop, the effective resistance Ω must be <3 MΩ.

It must be stressed that **10% drop of G₀=6300** allows **still to operate the detector at full efficiency**.

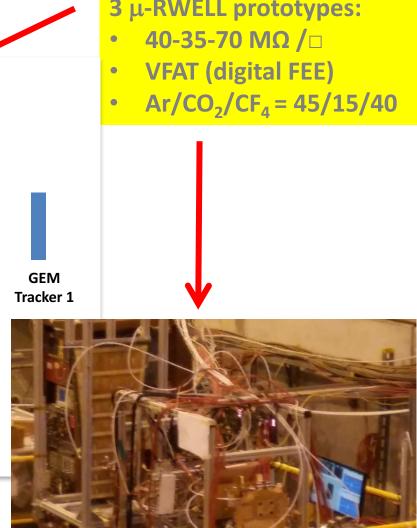


S1

Muon/Pion beam: 150 GeV/c



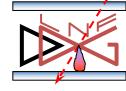
GOAL: <u>time resolution measurement</u> (never done before)



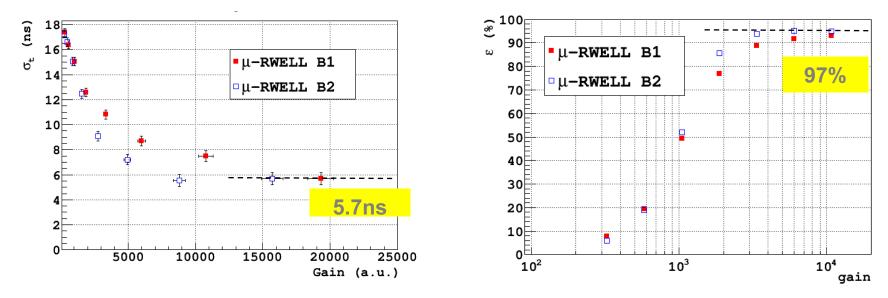
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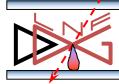
Ar:CO₂:CF₄ 45:15:40 – Strip readout



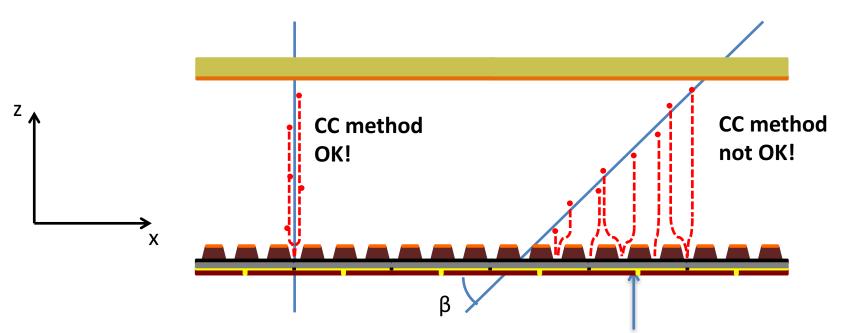
A time resolution of 5.7 ns has been measured with VFAT2. The saturation at 5.7 ns is dominated by the FEE. To be compared with past measurements done by our LHCb with GEM: σ_t = 4.5 ns with VTX chip and CF discriminator [G. Bencivenni et al., NIM A 494 (2002) 156]

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The use of an **analogic front-end** allows to associate a hit to a track using the charge centroid (CC) method. The uncertainty associated to the hit with this algorithm is dependent on the track angle: minimum for orthogonal tracks and larger as the angle increases

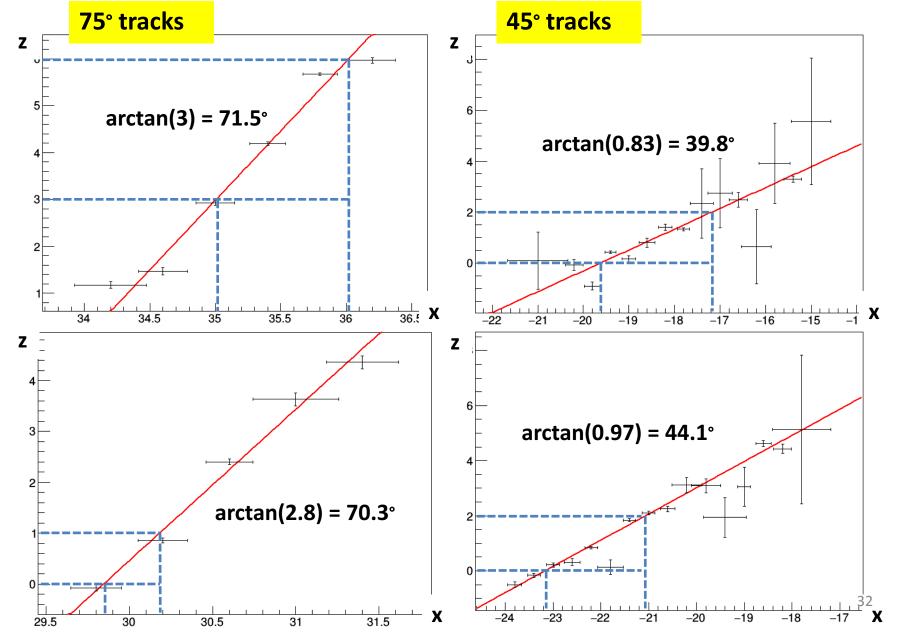


To improve the space resolution the **u-TPC algorithm** combined with the CC method has been implemented

Example of µ-TPC reconstruction

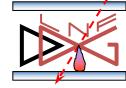
INFN

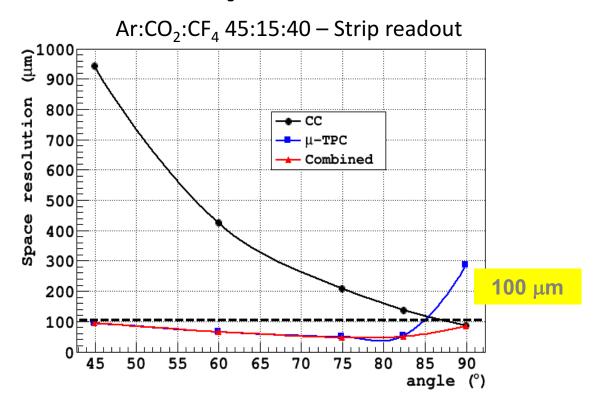
Istitute Nazionale di Fisica Nucleare Some examples where the tracks have an angle w.r.t. the readout plane





Improving space resolution: the µ-TPC mode





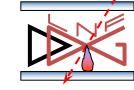
The combination of the CC and the μ -TPC mode with $E_d = 1$ kV/cm The spatial resolution is flattened for a wide range of angles.

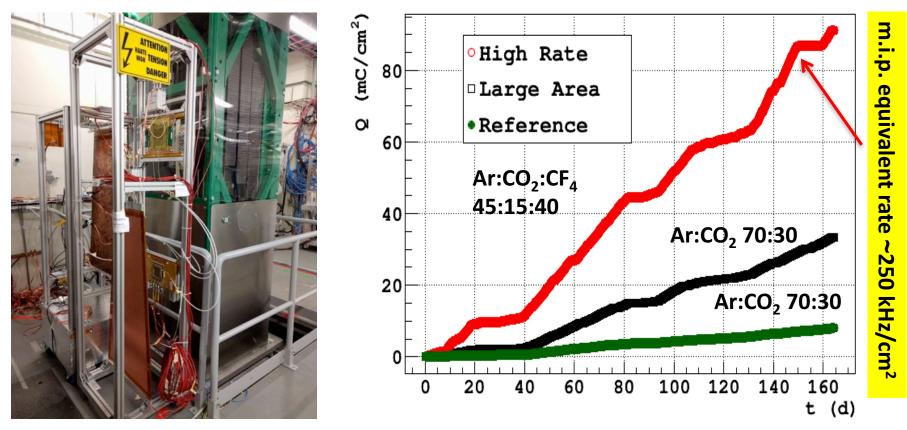
$$x_{merge} = \frac{x_{cc} \cdot w_{cc} + x_{tpc} \cdot w_{tpc}}{w_{cc} + w_{tpc}} \qquad \qquad w_{cc} \propto (clsize)^{-2} \quad w_{\mu-TPC} \propto (clsize)^{2}$$

M. Poli Lener, LNF-INFN - CepC Workshop



Ageing test at GIF⁺⁺ (CERN)





The ageing effects on DLC is under study at the GIF++ by irradiating different μ -RWELL prototypes operated at a gain of 4000 .

Up to now on the most irradiated detector (~250 kHz/cm² m.ip. equivalent) a charge of about 90 mC/cm² has been integrated

In collaboration with CMS-Muon: L. Borgonovi, P. Giacomelli, A. Ranieri



Summary



Recent R&D on a novel MPGD architectures lead to the introduction of the μ -RWELL in the MPGDs world.

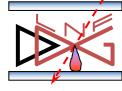
The μ -RWELL is a very promising technology showing important advantages for large area applications in harsh environment: the detector is compact, simple to assemble and intrinsically spark-protected

- gas gain > 10⁴
- rate capability > 1 MHz/cm² (HR version)
- space resolution < 100μm (over a large incidence angle of tracks)
- time resolution ~ 5.7 ns

R&D/engineering in progress:

- Low rate (<100kHz/cm²) :
 - small and large area prototypes built and extensively tested
 - <u>Technological Transfer to industry is ongoing with good achievements</u>
- High rate (>1 MHz/cm²):
 - R&D well advanced, completed by end of 2018
 - prototypes show very good performance

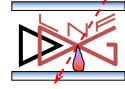




Thanks for your attention



Reference for the μ -RWELL detector:



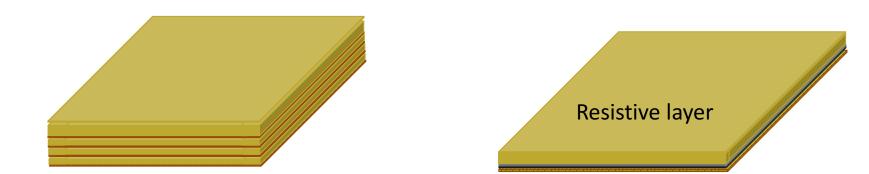
Low rate version:

- G. Bencivenni et al., "The micro-Resistive WELL detector: a compact spark-protected single amplification-stage MPGD", 2015_JINST_10_P02008
- G.Bencivenni et al., "The Resistive-WELL detector: a compact spark-protected single", PoS (BORMIO2015) 024
- G. Bencivenni et al., "The μ-RWELL: a compact, spark protected, single amplificationstage MPGD", NIM A 824 (2016) 565
- G.Bencivenni et al., "Advances on micro-RWELL gaseous detector", PoS (BORMIO2017) 002
- G.Bencivenni et al., "The μ-RWELL detector", 2017_JINST_114P_0517
- G.Bencivenni et al., "Performance of u-RWELL detector vs resistivity of the resistive stage", NIM A 886 (2018) 36

High rate version:

- G.Bencivenni et al., "Recent results of u-RWELL detector", PoS(MPGD2017)019
- G.Bencivenni et al., "The u-RWELL technology: status and perspective", to be submitted to Pos

backup



GEM µ-RWELL MicroMegas

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The µ-RWELL vs single-GEM

μ -RWELL is expected to exhibit a gas gain larger than a single-GEM

Single-GEM

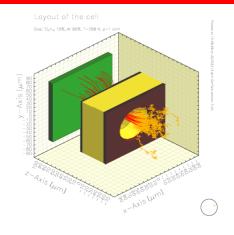
- ~50% of the electron charge produced into the hole contributes to the signal, the rest of the electron charge is collected by the bottom side of the GEM foil
- the signal is mainly due to the electron motion, the ion component is largely shielded by the GEM foil itself

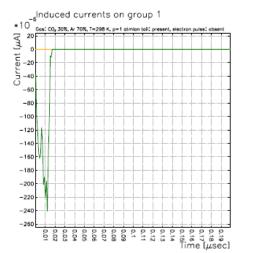
μ -RWELL

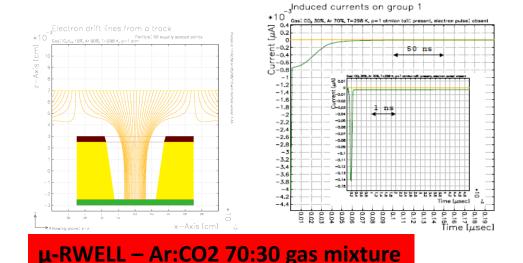
- 100% electron charge produced into the amplification channel is promptly collected on the resistive layer
- the ionic component, apart ballistic effects, contributes to the formation of the signal
- further increase of the gain achieved thanks to the resistive electrode which, quenching the discharges, allows to reach higher amplification field inside the channel

The µ-RWELL vs GEM (Garfield)

GEM – Ar:CO2 70:30 gas mixture







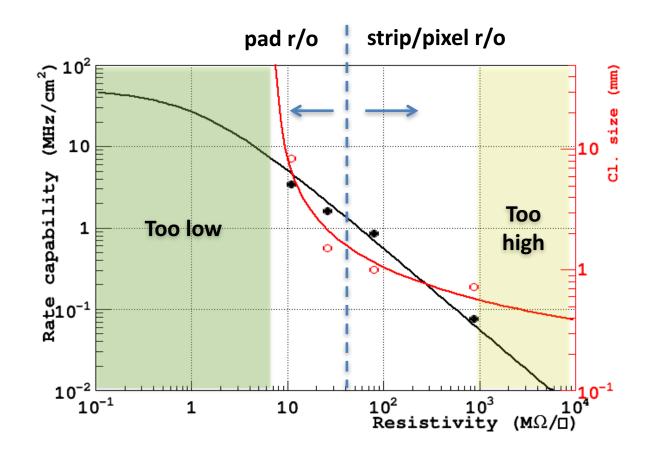
Signal from a single ionization electron in a GEM. The duration of the signal, about 20

ns, depends on the induction gap thickness, drift velocity and electric field in the gap.

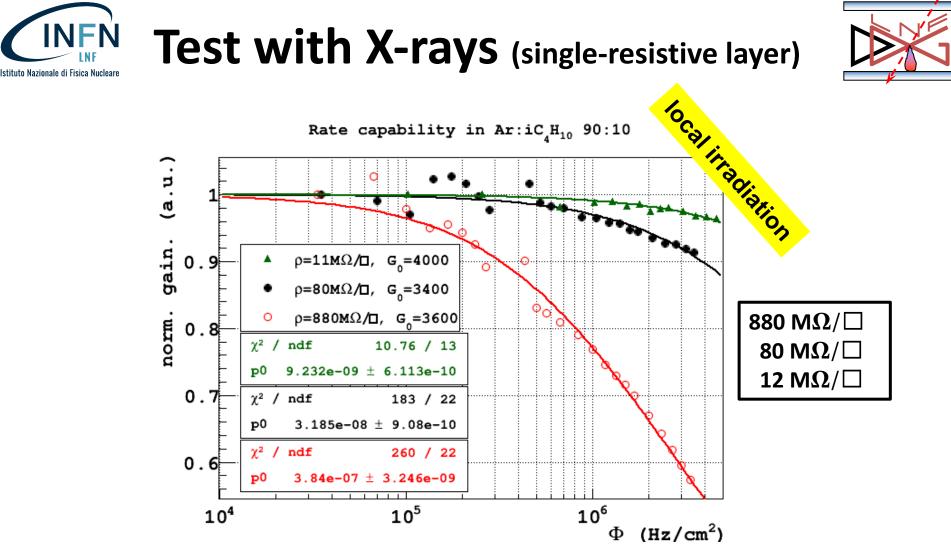
Signal from a single ionization electron in a μ -RWELL.

The absence of the induction gap is responsible for the **fast initial spike**, **about 200 ps**, induced by the **motion and fast collection of the electrons** then **followed by a ~50 ns ion tail**. **More similar to a MM !!!**

Combining the information



Qualitatively: low resistivity → pad r/out & higher rate high resistivity → strip/pixel r/out & lower rate



The gain decrease is correlated with the voltage drop due to the resistive layer: larger the resistivity higher gain drop

local irradiation ≠ global irradiation (to be compared w/slide # 27)



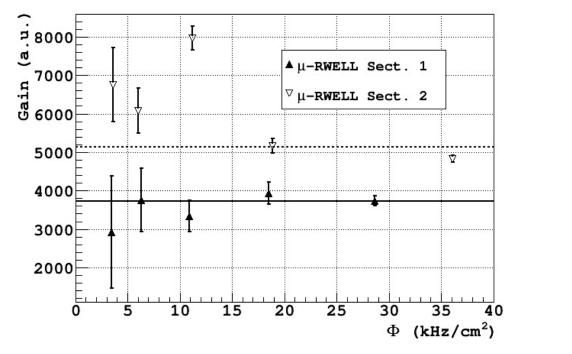
Rate capability of SRL layout



A rate capability measurement of a large single-resistive layer detector $(1,2x0,5 m^2 - 70)$

M Ω **/** \Box **)** has been performed with:

- a local irradiation with a **pion beam @** H8-SPS CERN beam area (**spot area ~3x3 cm²**)
 - → The detector has been **operated up to 35 kHz/cm**²
- a global irradiation @ GIF++ (see slide ...) up to **140 kHz/cm2** without appreciable gain loss (gain ~ 4000)



H8-SPS CERN

In collaboration with CMS-Muon:

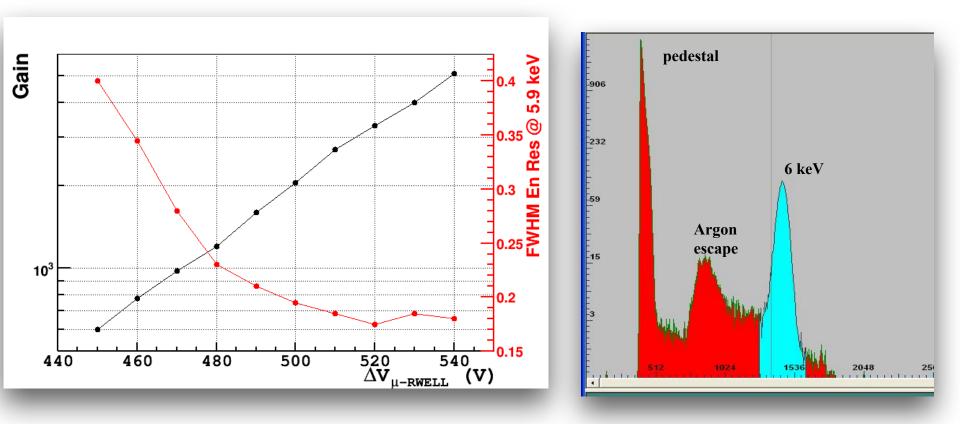
- L. Benussi¹, L. Borgonovi⁴,
- P. Giacomelli⁴, L.Borgonovo⁴, A. Ranieri⁵,
- M. Ressegotti⁶, I. Vai⁶, V. Valentino⁵
- 1. LNF- INFN
- 2. INFN Sezione di Bologna
- 3. INFN Sezione di Bari
- 4. INFN Sezione di Pavia



µ-RWELL: Energy Resolution



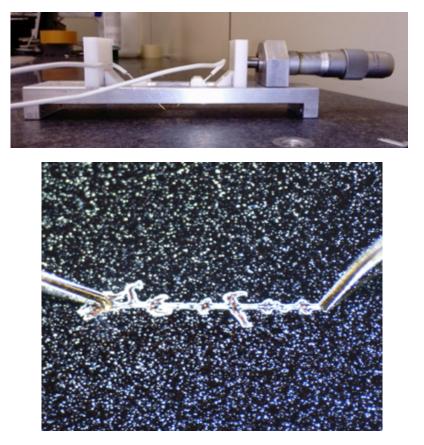
The prototype of μ -RWELL (100 M Ω / \Box) has been tested with X-rays tube (6keV) (Ar/CO2=70/30) & the signal has been readout vith an ORTEC amplifier

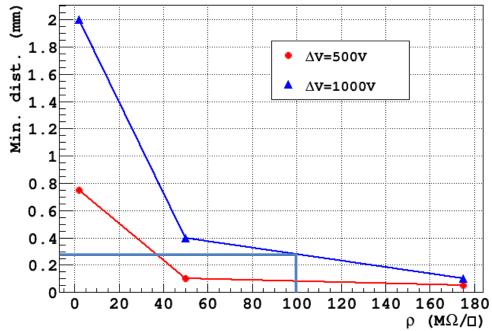


M. Poli Lener - LNF Mini-Workshop Series: Development of novel detectors at LNF, 14th June 2016

Silver Grid: optimization

In order to reduce the dead area, we have studied the Distance Of Closest Approach (*without discharges*) between two tips connected to an HV power supply. We recorded the minimum distance before a discharge on the DLC occurred *vs* the ΔV supplied for foils with different surface resistivity.







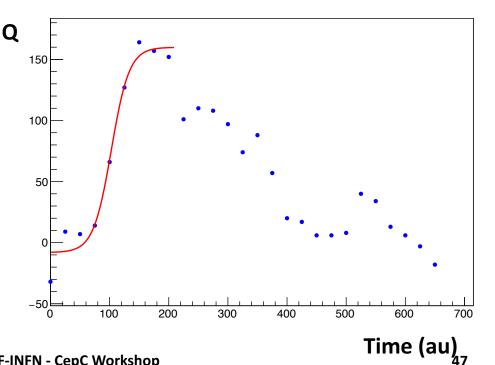
Introduced for **MicroMegas** by **T. Alexopoulos** et al. [NIM **A 617** (2010) 161] it suggests a way to overcome the **poor position reconstruction of the inclined tracks.**

Each hit is projected inside the conversion gap, where the x position is given by each strip and the $z = v_d t$

The drift velocity is provided by the Magboltz libraries.

The drift time is obtained with a fit of the charge sampled every 25 ns (APV25) from each FEE channel associated to the strip.

For each event we obtain a set of projected hits that once fitted provide a track segment



M. Poli Lener, LNF-INFN - CepC Workshop