# Beam Gas background characterization in the IR of FCC-ee

(mostly study of Inelastic Beam Gas for FCCee Z 45.6GeV/beam in this presentation)



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### FCC-ee, the machine

- A 100km tunnel in the CERN area,
- thought to be used for pp collision,

and as a first stage e+e- collisions,

• At 4 different center of mass energies

#### International FCC collaboration (CERN as host lab) to study:

- *pp*-collider (*FCC-hh*)
   → main emphasis, defining infrastructure requirements
- ~16 T  $\Rightarrow$  100 TeV *pp* in 100 km
- 80-100 km infrastructure in Geneva area
- e<sup>+</sup>e<sup>-</sup> collider (FCC-ee) as potential intermediate step / as a possible first step
- *p-e* (*FCC-he*) option, HE-LHC ...





Oide, FCC week 2017 https://indico.cern.ch/event/556692/contributions/2483406/attachments/1466449/2271404/FCCee\_Oide\_170529.pdf

# the BEAM GAS study

Gas in the beam pipe from the initial atmosphera (N2, O2, CO2, ...), or from desorption of the material in the beam pipe surface (H2, CO, CH4, ...) Interact with the beam



It seems that there are two main criteria to evaluate this process : a) Beam life time

b) Background in the experiment

DAFNE Tech Note V-3. Vaccarezza, 1991/JUL/08

"Preliminary Study for the Choice of the DAFNE Vacuum Chamber Material".

I will mostly address beam life time ...

There is another issue about the gas pressure from desorption due to radiation but this presentation DOES NOT address that.

# the BEAM GAS scattering Cross sections

#### Inelastic

e- losses energy in the interaction and is lost because of energy acceptance

$$\sigma(Z,\epsilon)$$

Cross section depends on : Z, the atomic number

 $\boldsymbol{\epsilon}$ , the energy acceptance of the machine

#### Elastic

e- is kicked out of its trajectory and is lost because of large oscillation

 $\sigma(Z, \gamma, \langle \beta \rangle, a)$ 

Cross section depends on : Z, the atomic number γ, the relativistic factor ⟨β⟩, the average beta a, the beam pipe aperture

### Beam Gas Cross Section for part of the ring

Equations give an estimate of the importance (order of magnitud) of the phenomena, mostly valid in the arcs.

However a detailed analisis requires a simulation, usually Monte Carlo.

Lattice Section	Gas	Elastic [barn]	Inelastic [barn]	
		(Ζ, γ, <β>, aperture)	(Ζ, ε)	
Arc	H2	0.0002 From Le Duff	0.3 Burkhardt	
Arc	N2	0.04 From Le Duff	<b>9.4</b> Burkhardt	
MDI	H2	0.2 New formula	-	
MDI	N2	<b>20</b> New formula	-	

# WHAT WE DO WITH MDISim

We interfase optics (MAD-X), with geometry (.gdml) to perform Monte Carlo simulations (in Geant4)

First results of Inelastic Beam-gas scattering latest optics

 $\rightarrow$  Loss map and loss rates are obtained.

Results for:

IR at Z, H, W, H. I am showing plots for the Z

- Note that SR is not considered in the simulation. 6

# Z 45.6GeV/beam MDI Region

<u>at all energies the scattering distribution is similar,</u> <u>BUT, rates differ</u>

# FCCee Z 45.6GeV/beam

Lattice : FCCee\_z\_213\_nosol\_18.seq (ZOOM)



We consider to study from s=-830m to s=370m

Conical Tapers from Kersevan. IR Vacuum Concept. Workshop on the Mechanical optimisation of the FCC-ee MDI https://indico.cern.ch/event/694811/timetable/





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#### FCCee Z geometry scalexy 50 : (scalexy 1 is not displayed correctly, but IS used for the tracking studies) Tracking starts at -830m and ends at 370m = 1200 m in total



#### As reference $\tau$ =100h $\rightarrow$ Scattering Rate = 78.6 Mhz/km/beam Meaning that a beam life time on the order of tens of hours is expected







# Local pressure variations

- Up to now, in the study it was assumed that the pressure is constant all over the beam pipe
- However, this is not the case:
  - In fact, we expect to have a locally variating pressure profile due to gas desorption caused by:
    - Interaction of SR photons with the beam pipe
    - Pump effect gradient due to pipe conductance
- It is possible to study this pressure profile using SynRad + MolFlow
- This study was performed @t (182.5GeV) configuration (expected worst case)

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## Local pressure variations



<u>Kersevan, Vaccum Systems for the FCC-ee</u> <u>7th Low Emittance Workshop, Jan-2018, CERN</u> <u>https://indico.cern.ch/event/671745/contributions/2788843/attachments/</u> 1582988/2502801/7thLowEmittanceWorkshop\_RKersevan\_Vacuum.pdf





Collamatti, FCC Week 2018 - Amsterdam Beam-gas background characterization in the FCC-ee IR <u>MOPMF085 Boscolo. BEAM-GAS BACKGROUND CHARACTERIZATION IN THE FCC-ee IR. IPAC18</u>

# Track display



Mogens Dam, FCC Week 2018 - Amsterdam LumiCal for FCC-ee and beam-background impact

### Beam-gas background

O. Blanco, F. Collamati

- At LEP, off-momentum particles from inelastic beam-gas scattering was the main background process to the luminosity measurement
- FCC-ee simulation of beam-gas scattering at Z-pole energy has been performed
   Loss rate inside region of ± 2.1 m around IP of

**2 MHz/beam** (a)  $10^{-9}$  mbar of N<sub>2</sub> at 300 K



# Sum-up and perspectives

- MDISim enables to import the whole machine configuration (beam + geometry) in GEANT4 starting from MAD-X file
  - Initially thought for Synchrotron Radiation studies
  - Once in GEANT4, an accurate study of the inelastic beam-gas background induced in the IR can be done (loss maps, scattering points, total loss rate, particle tracks)
- If assuming a vacuum pressure of 10-9 mbar of N<sub>2</sub>, results suggest a loss rate of the order of 100 MHz per Ampere of beam current in the region around IP
  - Tapering between different pipe apertures are critical points in which most BG scattered particles are lost
  - If local pressure variations due to gas desorption are taken into account, the total loss rate is increased of about 40%
- Tracks of particles exiting the pipe can be fed into detector simulations to evaluate their impact

• Initial study of background from Beam Gas scattering indicates negligible effect in the LumiCal

# BACK UP

### ENERGY OF THE LOST PARTICLES

Energy lost by the primary particle in the interaction with the gas molecule that led to particle loss

Energy of the particles that get lost due to BG when hitting the pipe



Lattices available in afs is X Initial studies have been done with the ones marked with X Latest lattice 213 for the t is now available

### All these plots and numbers are available for all these optics/energies

	fcc_ee_208	fcc_ee_213
Z	Х	Х
W	Х	Х
н	Х	Х
Т	Х	

Z (Euclidean) is NOT equal to S (C-S coordinates) ZBG is the location along Z where a Beam Gas Interaction occurs.

It looks pretty flat up to 3000m, I cut at 3000m For N :113812/1e7 = 1.13% are lost in 1km



### WHAT IS THE LOSS RATE FROM BEAM GAS IN THE ARC?

As calculated by Francesco :

Neloss = NelossMC/NeprimMC . Nebunch . Nbunches . Preal/Pmc

<u>For H the rate is 2.35 MHz per km of arc at 10<sup>-9</sup> mbar</u> Neloss = 68724/1e7 . 1.7e11 . 16640 . 1e-9/24.8 = 0.783e3 Rate\_eloss = Neloss/Trev = 18.1e3/0.333ms = 2.35 Mhz

<u>For H2 the rate is 6.22 MHz per km of arc at 10<sup>-9</sup> mbar</u> Neloss = 90764/1e7 . 1.7e11 . 16640 . 1e-9/12.4 = 2.070e3 Rate eloss = Neloss/Trev = 18.1e3/0.333ms = 6.22 Mhz

<u>For N the rate is 54.3 MHz per km of arc at 10<sup>-9</sup> mbar</u> Neloss = 113812/1e7 . 1.7e11 . 16640 . 1e-9/1.78 = 18.1e3 Rate\_eloss = Neloss/Trev = 18.1e3/0.333ms = 54.3 MHz

<u>For N2 the rate is 189.1 MHz per km of arc at 10<sup>-9</sup> mbar</u> Neloss = 99034/1e7 . 1.7e11 . 16640 . 1e-9/0.445 = 62.95e3 Rate\_eloss = Neloss/Trev = 63.0e3/0.333ms = 189.1 MHz

### THE PARTICLE LOSS IS NOT UNIFORMLY DISTRIBUTED



### Cumulative distribution of losses

#### **ELASTIC SCATTERING**

<u>Average cross section from average beta.</u> <u>J Le Duff. Current and current density limitations in existing electron storage rings</u> <u>Laboratoire de l'accélérateur Linéaire, LAL, NIM in Physics Research A239 (1985) 83-101.</u>

#### Average beta is invalid in the MDI region.

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J. Le Duff / Current and current density limitations

#### 2.1.1. The elastic scattering on nuclei

Elastic scattering on nuclei leads to an angular kick for the betatron motion. If the induced amplitude exceeds the vacuum chamber aperture the particle gets lost. The total cross section for this process is [1]:

$$\sigma_{\rm t1} = \frac{4r_{\rm e}^2 Z^2}{\gamma^2} \frac{\pi}{2} \left(\frac{\langle \beta \rangle}{a}\right)_z^2,$$

where

 $r_{\rm e}$  = the classical electron radius,

 $\gamma$  = the normalized energy  $(E/m_0c^2)$ ,

Z = the atomic number for residual gas components,

 $\langle \beta \rangle$  = the average betatron envelope function,

a = the half chamber aperture,

and it is assumed that the loss will occur in the vertical plane (z direction).

In the case of a non-smooth optic and assuming the aperture limit  $a_0$  is located at some azimuth where the envelope function is  $\beta_0$  one should replace  $(\langle \beta \rangle / a)^2$  by  $\langle \beta \rangle \beta_0 / a_0^2$  [2]. The elastic scattering effect goes up when the energy goes down.

#### References

- [1] J. Haïssinski, Thèse, Laboratoire de l'Accélérateur Linéaire, Orsay (1965).
- [2] H. Wiedemann, private communication.
- [3] A. Berthelot, Rayonnement de particules atomiques, électrons et photons (Masson, Paris, 1956).
- [4] F.F. Rieke and W. Prepejchal, Phys. Rev. A6 (4) (1972).
- [5] H. Bruck, Accélérateurs circulaires de particules (Presses Universitaires de France, 1966).

### **ELASTIC SCATTERING**

<u>MODIFIED EQUATION</u> <u>Based on Section 4.1 Elastic Scattering. CAS CERN Accelerator School. Fifth General Accelerator</u> <u>Physics Course, Vol 1. Jan/21/1994. Geneva, Switzerland.</u>

$$\overline{\sigma} = \frac{1}{\sum L_k} \frac{1}{\sum L_i} \frac{\tau_2 Z^2 r_e^2}{\gamma^2} \sum_k \sum_i L_k L_i \frac{\beta_k \beta_i}{H_i^2}; \text{ for } k \neq i$$

Sigma = average cross section

k = index of element inside the considered region

i = index of element along the entire ring

Gamma = relativistic gamma factor

Tau\_2 = 6.2832

L = total accelerator length

L\_k = length of kth-element in the chosen region

L\_i = length of the ith-element in the entire accelerator

Z = atomic number

r\_e = the classical electron radius

Beta\_i, Beta\_k = the vertical beta at the kth/ith-sm element

H\_i = the aperture of the ith-element

#### **INELASTIC SCATTERING**

Theoretical cross section from H. Burkhardt. M Brugger, H Burkhardt, and B Goddard. Interactions of Beams With Surroundings. Landolt-Boernstein, 21C:5–1 – 5–17, 2013.

#### eN scattering relevant for electron rings

The elastic cross section for eN scattering scales strongly with energy (with  $1/\gamma^2$ ) and scattering angle  $1/\theta^4$ . Elastic scattering is mostly relevant as halo production process for lower energy rings and becomes negligible for lifetime estimates for high energy electron rings.

At high energy, the dominating beam-gas process for electron rings is the inelastic scattering or bremsstrahlung in which the incident electron interacts with the field of the residual gas nucleus and radiates a photon.

The high energy cross section for eN scattering can be written in good approximation in dependently of the electron energy as [13]

$$\sigma_{\rm eN} = 4\alpha r_e^2 Z(Z+1) \log(287/\sqrt{Z}) \left( -\frac{4}{3} \log k_{\rm min} - \frac{5}{6} + \frac{4}{3} k_{\rm min} - \frac{k_{\rm min}^2}{2} \right), \tag{5.4}$$

where  $k_{\min}$  is the fractional energy loss or minimum photon energy in units of the electron energy,  $\alpha$  the fine-structure constant (1/137) and Z the atomic number (or number of protons). We can see that the cross section scales with Z(Z+1). Numerical values obtained from Eq. 5.4 for  $k_{\min} = 0.01$ are shown in Tab. 5.4

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· · · · · · · · · · · · · · · · · · ·		$\sigma_{\rm pN}$	$\sigma_{ m eN}$	Gas
-	10 <sup>4</sup>	b	b	
		0.08	0.28	$H_2$
		0.19	0.39	$\mathrm{He}$
_	10 <sup>3</sup>	0.43	3.02	$CH_4$
		0.40	4.38	$H_2O$
		0.56	6.47	$N_2$
		0.56	6.56	CO
	10 <sup>2</sup>	0.87	10.7	$\rm CO_2$
		0.60	17.8	$\operatorname{Ar}$
L <u>.</u> <u></u>			-	

Table 5.2. Numerical values for  $\sigma_{eN}$  for an energy loss of at least 1% and for  $\sigma_{pN}$ , the pN cross section at high energy ( $p_{lab} = 0.01$  to 10 TeV).  $Z^{*}(Z+1)^{*}log(287/sqrt(Z))$ 



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# Z 45.6GeV/beam arc

# FCCee arc Z 45.6GeV/beam

Lattice : FCCee\_z\_213\_nosol\_4.seq (ZOOM)



Constant aperture : 35mm

FCCee Z geometry scalexy 50 : (scalexy 1 is not displayed correctly, but IS used for the tracking studies) Tracking starts at 2000m and ends at 3200m = 1200 m in total



### WHAT IS THE PARTICLE LOSS RATE IN THE ARC?

At 45.6GeV/beam : Npart = 1.7e11 Nbunches = 16640

With an energy acc. of 2% and pressure of 10<sup>-9</sup> mbar

	<u>Scattering Rate/m/beam</u>			
	Expected	Simulation		
<u>H2 :</u> 0.328 barn	6.7 KHz	6.2 KHz		
<u>N2 :</u> 9.386 barn	192.3 KHz	189.1 KHz		
As reference $\tau$ =100h $\rightarrow$ Scattering Rate = 78.6 Khz/m/beam				

Beam gas particles generated in the arc are lost very soon, in the arc dipoles 100m~200m after interaction



### ENERGY EXCHANGE BETWEEN e- and gas?

Energy loss In the Beam Gas event :



Particles loosing less Than 0.95GeV continue In the beam

Energy acceptance 0.95GeV , i.e. 0.95/45.6 = 2.1%

### Particle distribution Along Z : Somehow uniform distribution of the particle loss

