We have got an

IDEA

an International Detector for Electron-Positron Accelerator(s)

Workshop on Circular Electron-Positron colliders Roma, May 24-26th, 2018



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Step 1

STEP I: a look at the different technologies, with a different point of view and starting





The essence of the designing and constructing a VERTEX DETECTOR: fit 1 GigaPixel in a Diet Coke can & keep it cool!

Physics First!

ILD DBD 2012



impact parameter resolution

Accelerator	a [μ m]	b $[\mu m \cdot GeV/c]$	
LEP	25	70	
SLC	8	33	
LHC	12	70	
RHIC-II	13	19	
ILC	< 5	< 10	ILD

ILD LOI 2009

The ILC figures apply also when you go beyond the linear approximations

a depends on the single point resolution and the ratio between the innermost radius and the lever arm:

 $\Rightarrow \sigma_{sp} = 3 \mu m$ when $R_{in} = 16 mm$ and $R_{out} = 60 mm$

[The ILD and CePC baseline figures]

b depends on the multiple scattering at the innermost radius: => thickness/layer = 0.15% X₀ [X₀ = 9.37 cm for Silicon]

[140 µm]

The machine comes next; and we have to account for

the time structure of the beams:

at the CepC, collisions are equally spaced (in time) with a frequency depending on the number of bunches. In one of the configurations reported in Beijing-201609, we have:

- 50 bunches at the Higgs factory energy
- 5000 bunches at the Z factory energy [where I estimated 4 kHz event rate]

for a beam Xing every 5 µs (@Higgs) to 50 ns (@Zpole) [3.6 µs is the "official" number]

the expected Beam-induced background:

there is actually NO solid rock number and estimates have a significant dependence on the machine & final focus parameters (HongBo, 2018, Roma).

A rough figure says ≈ 2.5 hits/cm²/Xing (I believe @Higgs energies)

BUT:

- having the spectrum of the bckg particles is important to see if we have "loopers"
- we have to see how it scales with the energy

the expected radiation level: RELAX!

* MAPS have been shown to be able to provide the required resolution with a binary read-out:



M. Winter et al., arXiv: 1203.3750v1 (2012)

The pitch/ $\sqrt{12}$ rule has been violated

Test beam results for the MIMOSA-26 sensor:

- 18.4 µm pitch (5.3 µm binary resolution)
- rolling shutter & end-of-column zero suppression (200 ns/pixel r.o. time)
 250 mW/cm² power consumption



* sophisticated architectures with ON PIXEL sparsification have been designed and qualified:



1 discriminator/pixel + 1 bit memory cell
 analog info locally processed
 the integration time is independent from read-out (r.o.) time
 the r.o. time is dependent from the pixel occupancy
 current power consumption at the level of 50 mW/cm² (ALPIDE)

-NIM A 765 (2014) 177 + A 785 (2015) 61 -pixel 2014 proceedings published on JINST (doi:10.1088/1748-0221/10/03/C03030)

* large systems have been designed and commissioned (or will be, in a short while):



400 sensors
 0.9 Pixel each
 power dissipation 170 mW/cm²

nothing but a toy compared to what is envisaged for the ITS of the ALICE experiment:



	σ_{sp}	t _{r.o.}	Dose	Fluency	T_{op}	Power	Active area
STAR-PXL	$<$ 4 μm	$<$ 200 μs	150 kRad	$3{\cdot}10^{12}~{ m n}_{eq}$ /cm 2	30-35°C	160 mW/cm 2	$0.15\mathrm{m}^2$
ITS-in	\lesssim 5 μm	\lesssim 30 μs	2.7 MRad	1.7 \cdot 10 13 n $_{eq}$ /cm 2	30°C	$<$ 300 mW/cm 2	$0.17 \mathrm{~m}^2$
ITS-out	\lesssim 10 μm	\lesssim 30 μs	100 kRad	$1{\cdot}10^{12}~\mathrm{n}_{eq}/\mathrm{cm}^2$	30°C	$<$ 100 mW/cm 2	\sim 10 m 2

a development based on:

new technologies (Tower-Jazz 180 nm)
 and new design (on pixel sparsification)

* and new technologies based on high resistivity substrates are very appealing:

The INFN SEED (Silicon with Embedded Electronics Development, partnership with LFoundry):



Vertex Detector Conclusions:

The new technologies certainly offer unprecedented opportunities

- Running conditions at the Z shall be carefully considered in designing the detector
- the real CHALLENGE, to me, will be designing an architecture providing the required data evacuation rate with the MINIMUM power dissipation (<20 mW/ cm²), resulting by an optimisation of the ANALOG CELL, the digital architecture, the clock distribution

But I'm confident that fun and excitement will exceed pain & fear!



One step up: the main TRACKING system

In principle, we have two options:









LC events by Graham Wilson, Como workshop 2013

One step up: the main TRACKIN

In practice, we do not since the EMITTANCE PRESERVATION at the IR constraints the B field to be at the 2T level.

600

400

200

So, in order to:

guarantee a good neutral/charged separation at the calorimeters:

d=0.15BR²/pt

the required curvature resolution at the 10⁻⁵ level [about a factor 10 better than LEP]

x-b (mm)

0.2

$$\begin{split} \frac{\Delta p_t}{p_t} &= \frac{8\sqrt{5}\sigma_{xy}}{0.3BR_{out}^2\sqrt{N}} p_t \oplus \frac{0.0523}{\beta BL} \sin\theta \sqrt{\frac{L}{X_0}} \\ \Delta \phi_0 &= \frac{4\sqrt{3}\sigma_{xy}}{R_{out}\sqrt{N}} \oplus \frac{0.0136}{\beta p} \sqrt{\frac{L}{X_0}} \\ \Delta \theta &= \frac{\sqrt{12}\sigma_z}{R_{out}\sqrt{N}} \frac{1 + \tan^2\theta}{\tan^2\theta} \oplus \frac{0.0136}{\beta p} \sqrt{\frac{L}{X_0}} \end{split}$$

We have to be LARGE, precise & LIGHT => there is no other detector than a 3D gas detector imager

Time Projection is not the only way and DRIFT chambers have been shown* to be an interesting alternative:



Dimensions of the MEG II chamber:

★ L = 193 cm

Ex-time

the date

Push the button, it will displa

- ✤ R_{in} = 17 cm
- ✤ R_{out} = 30 cm
- ✤ 10 layers for each 30° azimuthal sector







- M. Adinolfi et al., The tracking detector of the KLOE Eexperiment. NIM. A 488 (2002) 51
 - A. M. Baldini et al., Single-hit resolution measurement with MEG II drift chamber prototypes. 2016 JINST 11 P07011
 - G. Chiarello, The full stereo drift chamber for the MEG II experiment, 2017 JINST 12 C03062

its BIG BROTHER is being proposed as the main tracker of IDEA:



The IDEA drift chamber by numbers:

- ★ L = 400 cm
- ★ R_{in} = 35 cm
- R_{out} = 200 cm
- ✤ 112 layers for each 15° azimuthal sector
- 56 448 squared drift cells of about 12-13.5 mm edge
- max drift time: 350 ns in
 90%He-10%iC₄H₁₀

The "wire cage" and the "gas envelope" are decoupled

The stereo angle α is generated stringing the wire between spokes @ 2 sectors (30°) distance
α ∈ [20 mrad (1.1°); 180 mrad (10.3°)], increasing

with R the electrostatic stability is achieved when the wire tension is about 25g, for a total load of about 7,7 tons!





In cluster counting for improved particle identification: it is essentially based on the well known method of measuring the [truncated] mean dE/dX but it replaces the measurement of an ANALOG information with a DIGITAL one, namely the number of ionisation clusters per unit length:



compliance with the rates at the Z pole [a potential problem with the TPC]:

- a MEG2 prototype has been successfully tested till a rate O(100 kHz/cm²)
- simulation studies at the Z pole and up to 380 GeV indicate that ion drift related problems shall not be an issue

Drift Chamber Conclusions:

it sounds good, as long as you can crunch in real time all of the numbers...

[you do need to digitise the waveform from every sense wire at 2Gs/s on 8 bits => unless you process in real time, the required bandwidth skyrockets to 1 Tb/s...]

* for further details, see the talk by G.Tassielli



One more step: measuring the ENERGY

We know what we want:

The Blue Plot!



separating hadronically decaying W's from Z's helps for channels where kinematic fits will not work and, obviously, to reduce combinatorics
 in order to compensate from a degraded resolution, 20-40% higher integrated luminosity is required



We know there is one solution:





WARNING! preoptimisation studies! Actual numbers are likely to be different

- 130T of Tungsten (watch the commodity market..)
- ✤ 3000 m² of pixelated Silicon
- ✤ 250 Mpixel (well calibrated and stable...)

Today: reduced to 100 Millions....





We know there is one solution:



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Can I get there in a different way?

We know that:

- Calorimetry is a "fluctuation game" [leakage, sampling, e.m. fraction, invisible energy, noise];
- In hadron initiated showers, the main fluctuations in the event-to-event response are due to:
 - the share between the e.m. and and hadronic component
 - the fluctuations in the "invisible energy"

and the e.m. component is giving a significant contribution, growing with energy:



an example of the improvement that can be expected in the measurement of a sample of 100 GeV π 's if f_{e.m.} is NOT measured (top plot) or if f_{e.m.} bins are singled out

R. Wigmans, NIM A572 (2007) 215-217

We also know that:

if you embed in the same calorimeter a detector responding primarily to the e.m. fraction and detector responding to the total dE/dX, you can single out $f_{e.m.}$.

This was proposed (and successfully demonstrated in a series of different implementations) using Cherenkov light [produced by relativistic particles and dominated by the e.m. shower component] and scintillation => DUAL READOUT CALORIMETRY

Two exemplary results from the **DREAM/RD52 calorimeters**: [NIM A537 (2005) 537-561 - NIM A735 (2014) 130-144 - NIM A732 (2013) 475]

see the talks by: J. Hauptman R. Ferrari R. Santoro







So far, the idea of integrating such a detector concept in a 4π detector turned the DREAM into a nightmare

And it was so until when the Silicon age entered the photonics world and **PMT** were replaced by SiPM:







[sampling fraction 4.5%]

10×10 fibers



more info:
 talk by R. Santoro NOW!
 our NIM paper, available
 on the ArXiv: 1805.03251

S

Results by:

the 2016-2017 beam tests
a preliminary simulation of a 4π geometry





Made us confident this can be a solution worth being considered



When SIZE matters most than RESOLUTION

(e.g. pre-showers, muon detectors and wherever a large area with moderate single point precision is what you need)

=> Micro Pattern Gas Detectors enter the game

In the MPGD world, there is one technology that looks particularly promising for its characteristics and the possibility to be industrialised: the μ -RVVell

It essentially consists of:

- ***** a patterned Kapton foil (amplification stage)
- * a resistive layer sputtered on the back of the Kapton foil to quench the multiplication and avoid sparks (DLC = Diamond Like Carbon)
- * a patterned PCB for readout



Drift cathode PCB

Bencivenni et al., 2017 JINST 12 C06027
Bencivenni et al., 2015 JINST 10 P02008



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Gain vs Voltage



Cluster size and resolution vs Resistivity

Also worth saying that:

- the camera stands rates up to 35 kHz/cm² (simplest process)
- time resolutions at the few ns level have been measured



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a $1.2 \times 0.5 \text{ m}^2$ prototype has been produced in collaboration with an industrial partner (ELTOS s.pa)



Now that we got the pieces, we can try to assemble the detector & complete the exercise:



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▶ Beam Pipe (≈ 1.5 cm radius)



▶ Beam Pipe (≈1.5 cm radius)
▶ Vertex Detector ()





▶ Beam Pipe (≈ 1.5 cm radius)
▶ Vertex Detector (R ∈ [1.7; 34] cm) A BIG Coke can...
▶ Drift Chamber (L = 400 cm, R ∈ [35; 200] cm)





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 ▶ Outer Silicon Layer (strips)





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 ▶ SC Coil (2T, ≈2.1m); THIN! 30 cm (0.74X₀; 0.16 λ @90°)





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Dual Readout Calorimeter (2m, 7 λ)
Yoke & Muon Chambers





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- Yoke & Muon Chambers



Conclusions:

We know it is still a long way to go but:

***** time is on our side

* and we hope to have a good team walking with us!

