

A review on neutrinoless double beta decay



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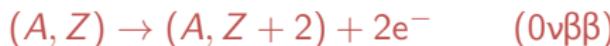
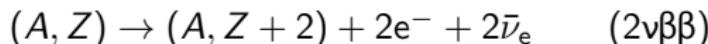
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VULCANO Workshop 2018

Frontier Objects in Astrophysics and Particle Physics

May 20 – 26, 2018 - Vulcano Island, Sicily, Italy

Double Beta Decay: real and virtual neutrinos



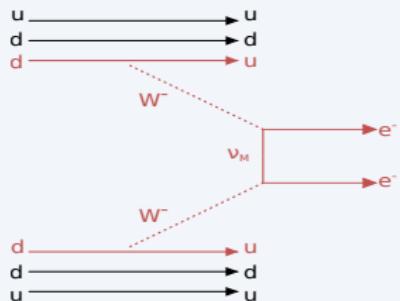
- *L-violation*, i. e. creation of a pair of electrons

- a discovery of $0\nu\beta\beta$ would imply
a violation of the SM global symmetries

- $0\nu\beta\beta$ key tool to *study neutrinos*

- Majorana or Dirac nature
 - ν mass scale and ordering
 - CP -violation in the lepton sector

Feynman diagram



Review: Adv. High En. Phys. 2016, 2162659 (2016)

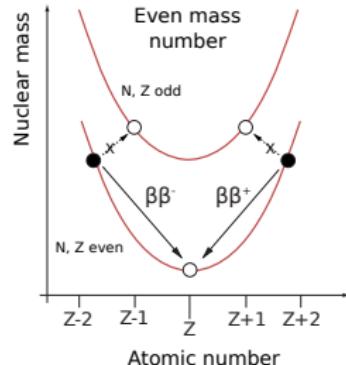
$0\nu\beta\beta$ half-life time

- $0\nu\beta\beta$ is a *nuclear process*
 - 2nd order transition: $(A, Z) \rightarrow (A, Z + 2)$
 - even-even nuclei: β -decay is forbidden
- half-life expression can be factorized as

$$[t_{1/2}^{0\nu}]^{-1} = G_{0\nu} |\mathcal{M}|^2 |f(m_i, U_{ei})|^2$$

- $G_{0\nu}$ = Phase Space Factor (atomic physics)
- \mathcal{M} = Nuclear Matrix Element (nuclear physics)
- $f(m_i, U_{ei})$ = mechanism (particle physics)

$$\rightarrow 3 \text{ light neutrino exchange: } f(m_i, U_{ei}) \equiv \frac{m_{\beta\beta}}{m_e} \quad (\text{Effective Majorana mass})$$

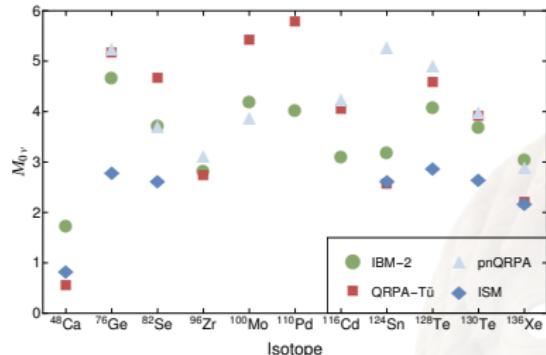


Caveat In order to extract the information on the neutrino mass, it is necessary to pass through *atomic* and *nuclear* physics

Models for the nuclear matrix elements

- calculation of the NMEs for the $0\nu\beta\beta$ is difficult
 - complex nuclear structure with many nucleons
 - ground + multiple excited states to be considered
- different theoretical models
 - Quasiparticle Random Phase Approximation
 - Intermediate Boson Model
 - Interacting Shell Model
 - ...

QRPA / IBM-2 within $\sim 30\%$



QRPA-Tu: F. Šimkovic *et al.*, Phys. Rev. C 87, 045501 (2013)

pnQRPA: J. Hyvärinen *et al.*, Phys. Rev. C 91, 024613 (2015)

IBM-2: J. Barea *et al.*, Phys. Rev. C 91, 034304 (2015)

ISM: J. Menéndez *et al.*, Nucl. Phys. A 818, 139 (2009)

Assessing the uncertainties

- estimate of the uncertainties on PSFs/NMEs is crucial to constrain $m_{\beta\beta}$
- theory of PSFs is known / mostly computational difficulties $\Rightarrow \sim 7\%$
- quite large uncertainties for the NMEs
 - error on single calculations of $\sim 20\%$ / still hard to give an overall value
 - calculations vs. rates discrepancies $\gg 20\%$ for known processes (β , EC, $2\nu\beta\beta$)
- $\mathcal{M} \equiv g_A^2 \mathcal{M}_{0\nu} = g_A^2 \left(\mathcal{M}_{GT}^{(0\nu)} - \left(\frac{g_V}{g_A} \right)^2 \mathcal{M}_F^{(0\nu)} + \mathcal{M}_T^{(0\nu)} \right)$
 - significant effect of g_A : uncertainty on its values \Rightarrow larger uncertainty on \mathcal{M}
 - the value of g_A in the nuclear medium is not reliably known
 - from 1.27 (free nucleon) to < 1 (quenching)

Effective Majorana mass (3 light neutrino exchange)

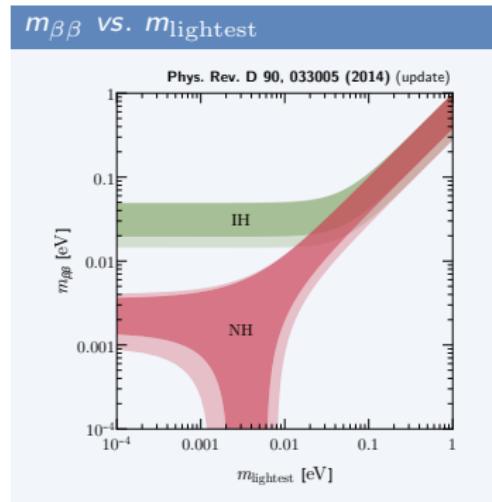
- $m_{\beta\beta}$ is the key quantity in the $0\nu\beta\beta$

- absolute value of the ee-entry of the neutrino mass matrix

$$\square m_{\beta\beta} = \left| \sum_{i=1,2,3} e^{i\xi_i} |U_{ei}^2| m_i \right|$$

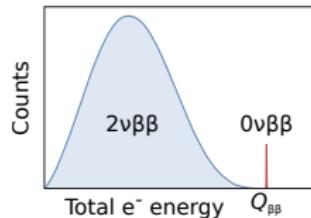
- 1 CP-violating + 3 Majorana phases
 - U can be identified with the mixing matrix of the **oscillation analysis**
 - $U \equiv U|_{\text{osc.}} \cdot \text{diag} \left(e^{-i\xi_1/2}, e^{-i\xi_2/2}, e^{i\phi - i\xi_3/2} \right)$
 - only two phases play a *physical* role

$$\blacksquare m_{\beta\beta} = \left| e^{i\alpha_1} \cos^2 \theta_{12} \cos^2 \theta_{13} m_1 + e^{i\alpha_2} \cos^2 \theta_{13} \sin^2 \theta_{12} m_2 + \sin^2 \theta_{13} m_3 \right|$$



Experimental search for $0\nu\beta\beta$

- **detection:** energy (track) of the 2 emitted e^-
 - monochromatic peak at $Q_{\beta\beta}$
 - smearing due to finite energy resolution
- **observable:** decay half-life of the isotope, $t_{1/2}^{0\nu}$
 - in the case of a peak in the energy spectrum



$$(A, Z) \rightarrow (A, Z + 2) + 2e^-$$

$$t_{1/2}^{0\nu} = \ln 2 \cdot T \cdot \varepsilon \cdot \frac{N_{\beta\beta}}{N_{\text{peak}}} \quad \left(\frac{\delta t_{1/2}^{0\nu}}{t_{1/2}^{0\nu}} = \frac{\delta N_{\text{peak}}}{N_{\text{peak}}} \right)$$

- if no peak is detected, the **sensitivity** corresponds to the maximum signal that can be hidden by the background fluctuations $n_B = \sqrt{M T B \Delta}$

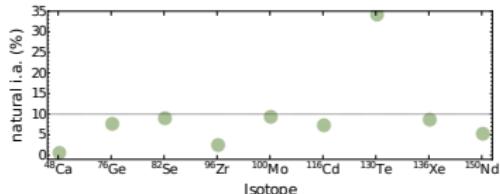
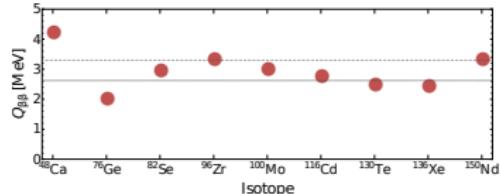
$$S_{1/2}^{0\nu} = \ln 2 \cdot T \cdot \varepsilon \cdot \frac{n_{\beta\beta}}{n_\sigma \cdot n_B} = \ln 2 \cdot \varepsilon \cdot \frac{1}{n_\sigma} \cdot \frac{x \eta N_A}{\mathcal{M}_A} \cdot \sqrt{\frac{M T}{B \Delta}}$$

(M = detector mass T = measuring time B = background level Δ = energy resolution)

Choice of the isotope

- value of $Q_{\beta\beta} \rightarrow$ influences the bkg
 - 2.6 (3.3) MeV end-point of main γ s (β s)
 - avoid radioactivity peak position
 - obs. suitability depends on detector features
- high isotopic abundance
 - ease of material enrichment
(technologically + economically)
- availability of the isotope
 - tonnes required for future 0v $\beta\beta$ experiments
→ high cost + large procurement time
- compatibility with a detection technique

Most suitable **isotope + detector** combination



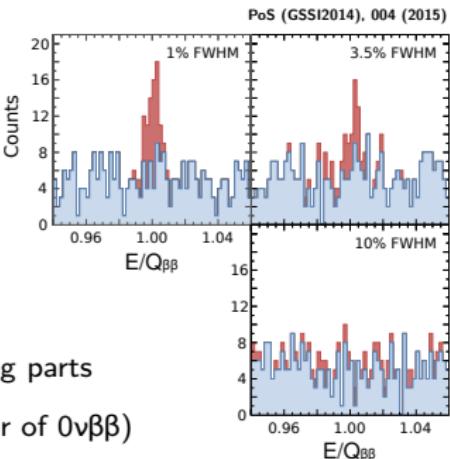
Promising isotopes

^{48}Ca ^{76}Ge ^{82}Se
 ^{96}Zr ^{100}Mo ^{116}Cd
 ^{130}Te ^{136}Xe ^{150}Nd

Detector requirements

- good energy resolution

- only protection against $2\nu\beta\beta$ spectrum tail
 - $R_{0\nu/2\nu} \propto \left(\frac{Q_{\beta\beta}}{\Delta}\right)^6 \frac{t_{1/2}^{2\nu}}{t_{1/2}^{0\nu}}$



- very low background

- underground location + shielding
 - radio-pure materials for detector and surrounding parts
 - (($10^9 - 10^{10}$) yr from natural chains vs. $> 10^{25}$ yr of $0\nu\beta\beta$)
 - analysis rejection techniques

- large isotope mass

- present: some tens up to hundreds of kg
 - tonnes required to cover the IH region

How to get the best sensitivity?
It is up to the experimentalists to
choose **which aspect to privilege**

Experimental techniques (I)

■ Ge-diodes

- high-purity enriched crystals
- high energy resolution ($\lesssim 0.2\% @ Q_{\beta\beta}$)
- bkg rejection by pulse shape analysis

Heidelberg-Moscow IGEX

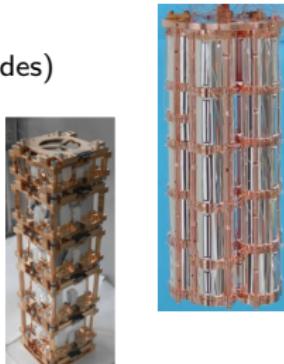


GERDA MAJORANA Demonstrator LEGEND

■ bolometers

- high energy resolution (close to Ge-diodes)
- many compounds with $0\nu\beta\beta$ emitters
- large source masses
- complex cryogenic infrastructure

Cuoricino CUORE-0 CUPID-0



AMoRE CUORE CUPID



Experimental techniques (II)

- Xe liquid . . .

- Xe easily enrichable
 - event topology reconstruction
 - low energy resolution ($\sim 3\%$)



- . . . and gaseous TPCs

- higher energy resolution
 - lower signal efficiency ($\sim 30\%$)



EXO-200 NEXT nEXO PandaX-III

- liquid scintillators loaded with $0\nu\beta\beta$ isotope

- poor energy resolution ($\sim 10\%$)
 - huge amount of material
 - very low background



KamLAND-Zen SNO+

Experimental techniques (III)

- tracker + calorimeter
 - almost no limitations in the choice of the isotope
 - large isotope masses hardly achievable
 - low energy resolution
 - event topology reconstruction



NEMO-3 SuperNEMO

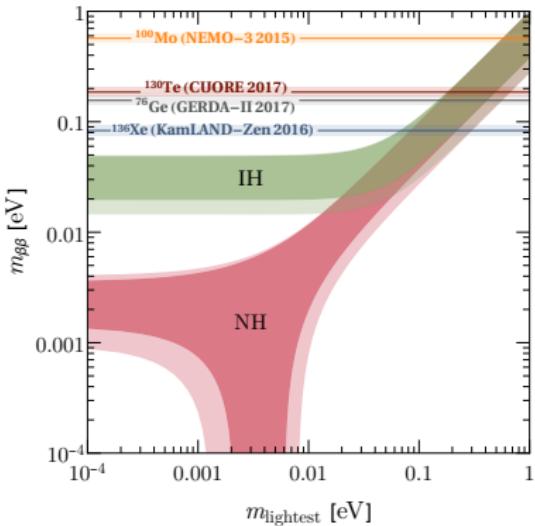
- others
 - variations on the previous, or new techniques
 - numerous running prototypes and R&D projects



CANDLES COBRA ZICOS
DCBA/MTD FLARES ...



Present sensitivity



$$m_{\beta\beta} \leq \frac{m_e}{g_A^2 M_{0\nu} \sqrt{G_{0\nu} t_{1/2}^{0\nu}}}$$

NMEs (IBM-2): *Phys. Rev. C* 91, 034304 (2015)

PSFs: *Phys. Rev. C* 85, 034316 (2012)

$g_A = 1.269$

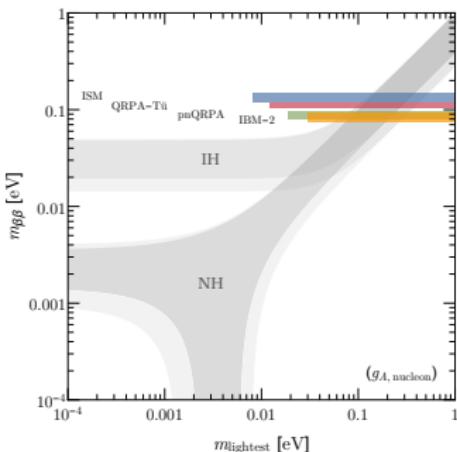
obs. *fixing the NME model allows to compare different experimental sensitivities, but it underestimates the uncertainties*

Isotope	$t_{1/2}^{0\nu}$ (90% C. L.) [yr]	$m_{\beta\beta}^{\min}$ [eV]	Reference
¹⁰⁰ Mo	$1.1 \cdot 10^{24}$	0.569 ± 0.053	<i>Phys. Rev. D</i> 92, 072011 (2015)
¹³⁰ Te	$1.5 \cdot 10^{25}$	0.186 ± 0.017	<i>Phys. Rev. Lett.</i> 120, 132501 (2018)
⁷⁶ Ge	$8.0 \cdot 10^{25}$	0.159 ± 0.014	L. Pandola, <i>results presented @ TAUP 2017</i>
¹³⁶ Xe	$1.1 \cdot 10^{26}$	0.083 ± 0.009	<i>Phys. Rev. Lett.</i> 117, 082503 (2016)

Effect of the nuclear uncertainties: Xe case

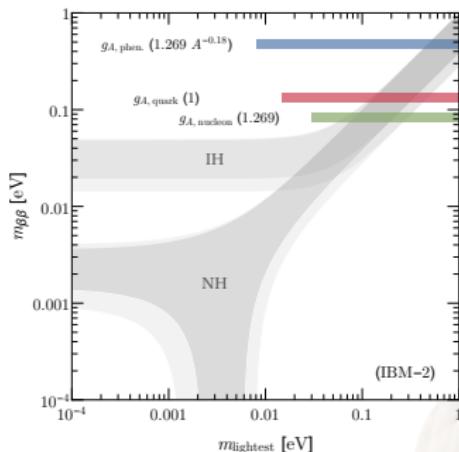
- different NMEs / fixed g_A

□ $73 \text{ meV} < m_{\beta\beta} < 149 \text{ meV}$



- different g_A / fixed NMEs

□ $73 \text{ meV} < m_{\beta\beta} < (149) 542 \text{ meV}$



$$t_{1/2}^{0\nu} \propto g_A^4 \mathcal{M}_{0\nu}^2$$

Towards the next generation of $0\nu\beta\beta$ experiments

- attaining a sensitivity of the order of **10^{28} yr or more** is important or even essential, allowing to test the type of **neutrino mass spectrum**
- this goal imposes **stringent requirements to the detectors**
 - the **feasibility and effectiveness** of the proposed technique have **to be tested** by means of demonstrators or extensive R&D programmes
 - **money**, i. e. $\$/\text{mole}$ of *detectable* isotope, must be included in sensitivity studies
 - technological costs: procurement, enrichment/purification, infrastructures...
 - efficiency of detector/setup scalability to tonne-size
 - **merging of experiments** sharing the same technology and studying **different nuclei** with a specific detector should be considered
- a **large effort** has to be put **in the nuclear studies** (NMEs and effective value of g_A) in order to maximize the information that can be extracted from the experimental searches

A crucial issue: background suppression

- recall: $S_{1/2}^{0\nu} = \ln 2 \cdot T \cdot \varepsilon \cdot \frac{n_{\beta\beta}}{n_\sigma \cdot n_B} = \ln 2 \cdot \varepsilon \cdot \frac{1}{n_\sigma} \cdot \frac{x \eta N_A}{\mathcal{M}_A} \cdot \sqrt{\frac{MT}{B\Delta}}$
- when B is sufficiently low \rightarrow **zero background condition**
 - transition region in between: $MTB\Delta = \mathcal{O}(1)$ (no expected events in the ROI)
- $S_{1/2,0B}^{0\nu} = \ln 2 \cdot T \cdot \varepsilon \cdot \frac{N_{\beta\beta}}{n_\sigma \cdot n_B} = \ln 2 \cdot \varepsilon \cdot \frac{x \eta N_A}{\mathcal{M}_A} \cdot \frac{MT}{N_S}$
 - the sensitivities scales linearly with the exposure!

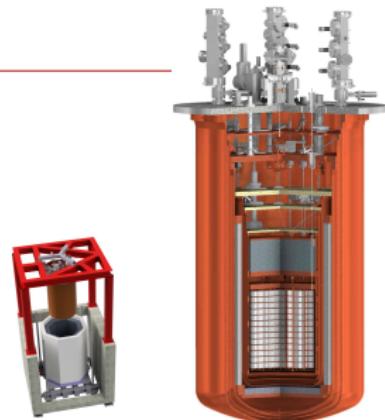
The zero bkg condition depends on M : **the larger the detector mass, the more strict the request on the background**

- the same bkg level can suffice for a kg-size experiment, but not for a tonne-size one

Future players: CUPID

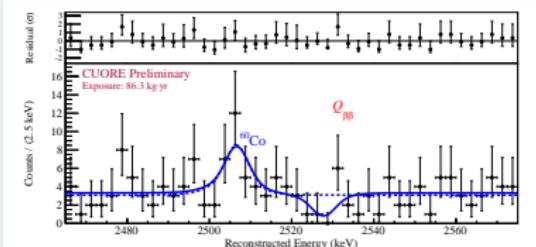
■ CUORE Upgrade with Particle IDentification (Cryogenic Underground Observatory for Rare Events)

- 750 kg bolometric array of enriched crystals
- TeO_2 or Mo/Se compounds
- background in ROI: $0.1 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
- energy resolution: 5 keV FWHM @ $Q_{\beta\beta}$
- sensitivity goal on $0\nu\beta\beta$ half-life:
 $5 \cdot 10^{27} \text{ yr}$



CUORE first results (Oct 2017)

Phys. Rev. Lett. 120, 132501 (2018)



$$t_{1/2}^{0\nu}(\text{Te}) > 1.5 \cdot 10^{25} \text{ yr} @ 90\% \text{ C. L.}$$

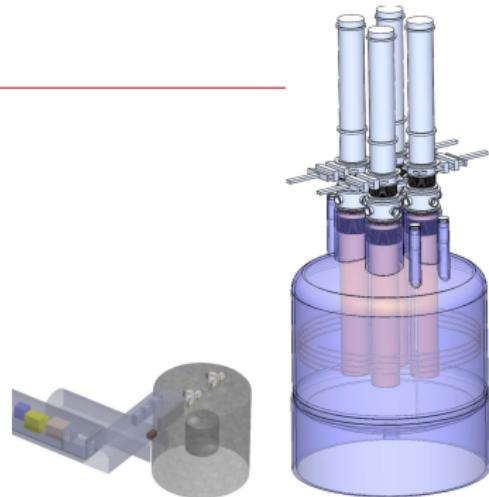
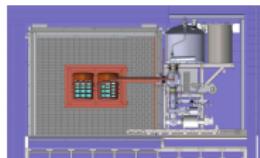
Future players: LEGEND

- Large Enriched Germanium Experiment for Neutrinoless $\beta\beta$ Decay

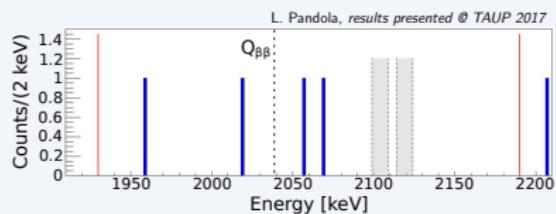
- 200 kg \rightarrow 1 t of ^{76}Ge enriched HPGe-diodes
 - background in ROI: $0.1 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
 - energy resolution: 2.5 keV FWHM @ $Q_{\beta\beta}$
 - sensitivity goal on $0\nu\beta\beta$ half-life: 10^{28} yr

- best of GERDA & MJD

- water + LAr for low-A shielding (G)
 - LAr active veto (G)
 - radio-pure material, especially Cu (M)
 - low-noise electronics (M)



GERDA-II new results (Jul 2017)

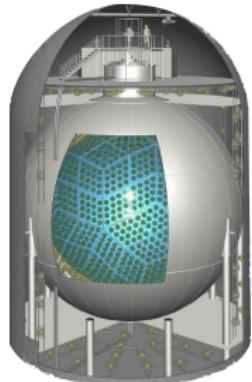


$t_{1/2}^{0\nu} > 8.0 \cdot 10^{25} \text{ yr}$ @ 90% C. L.

Future players: KamLAND2-Zen

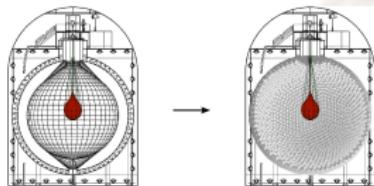
- **Kamioka Liquid scintillator Anti-Neutrino Detector 2 - Zero neutrino**

- 1 t of ^{136}Xe enriched xenon dissolved in LS
 - background in ROI: $0.01 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
 - energy resolution: 50 keV FWHM @ $Q_{\beta\beta}$
 - sensitivity goal on $0\nu\beta\beta$ half-life: 10^{27} yr



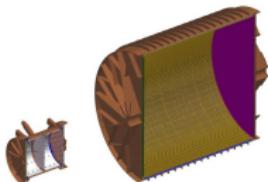
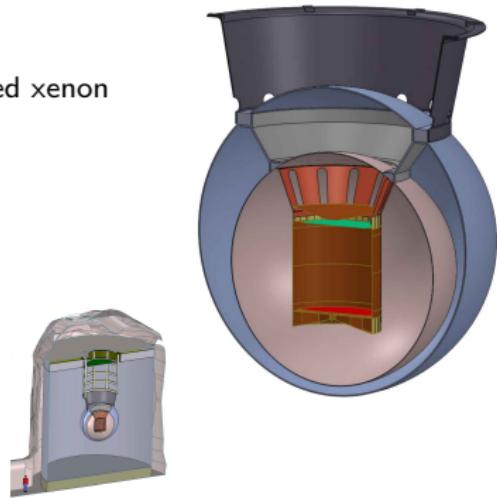
- from the experience of KamLAND-Zen

- improved light collection (new LS & PMTs + collectors)
 - scintillating balloon (^{214}Bi tagging)
 - new method for LS purification
 - pressurized Xe-LS



Future players: nEXO

- next Enriched Xenon Observatory
 - liquid TPC with 4.7 t of active ^{136}Xe enriched xenon
 - background in ROI: $0.01 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
 - energy resolution: 60 keV FWHM @ $Q_{\beta\beta}$
 - sensitivity goal on $0\nu\beta\beta$ half-life: $9 \cdot 10^{27} \text{ yr}$
- from the experience of EXO-200
 - 3x larger size \Rightarrow 30x mass/volume
 - improved design & components
 - increased light collection
(larger coverage + APDs \rightarrow SiPMs)



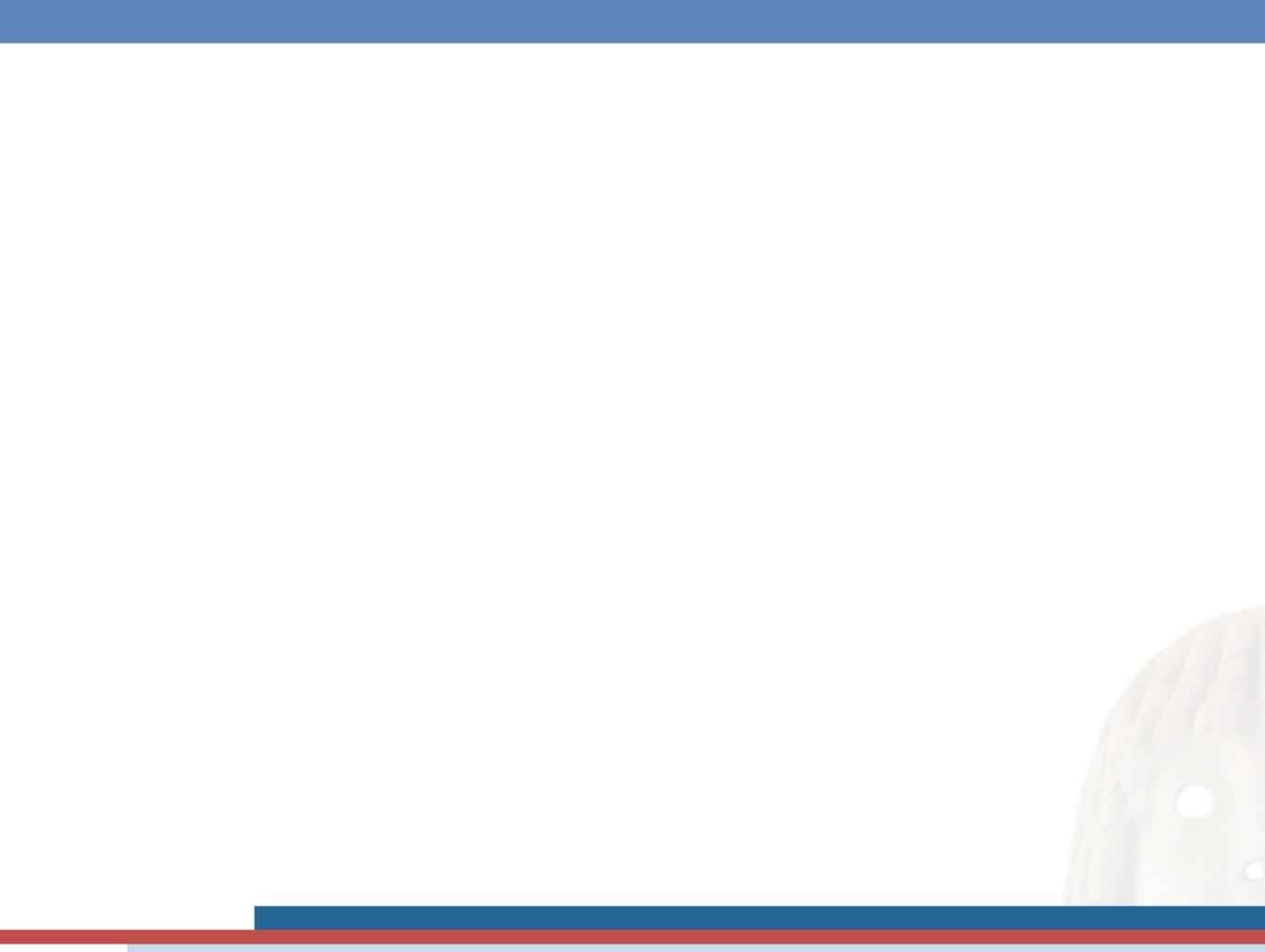
A major challenge: Ba tagging

- $^{136}\text{Xe} \rightarrow ^{136}\text{Ba}^{++} + 2 e^-$
- complete background elimination
- 40x in expected sensitivity: $4 \cdot 10^{28} \text{ yr}$

Summary

- $0\nu\beta\beta$ is a unique tool to study lepton number violation and ν masses
- today, sensitivities of the order of $(10^{25} - 10^{26})$ yr on the decay half-life time have been reached
- the next generation of detectors aims at $(10^{27} - 10^{28})$ yr sensitivities
 - money will represent a critical issue for the experiments
 - reaching the zero background condition will be of crucial importance
- a better understanding of the nuclear physics is needed (\rightarrow probing IH)
- the field is very active, with ambitious experimental proposals and numerous R&D programmes

Thank you!



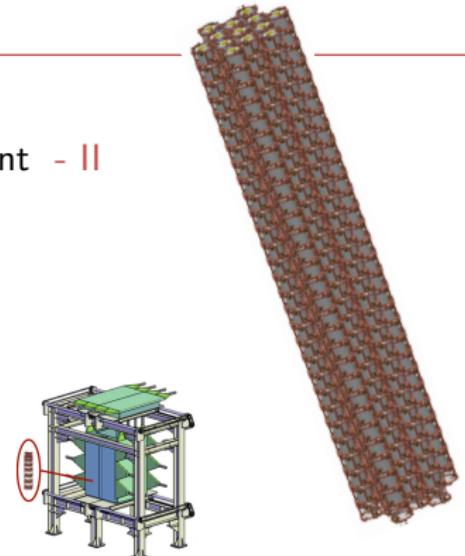
Future players: AMoRE-II

■ Advanced Mo-based Rare process Experiment - II

- 200 kg bolometric array of ^{enr}Mo -based crystals
- background in ROI: $0.1 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
- energy resolution: 10 keV FWHM @ $Q_{\beta\beta}$
- sensitivity goal on $0\nu\beta\beta$ half-life: $5 \cdot 10^{26} \text{ yr}$

■ result of the AMoRE programme

- new underground laboratory: ARF
@ Handeok Iron Mine (1100 m overburden)
- new cryogenic infrastructure
- improved detector performance

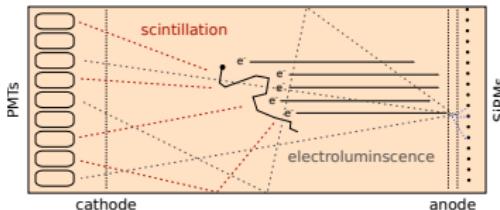
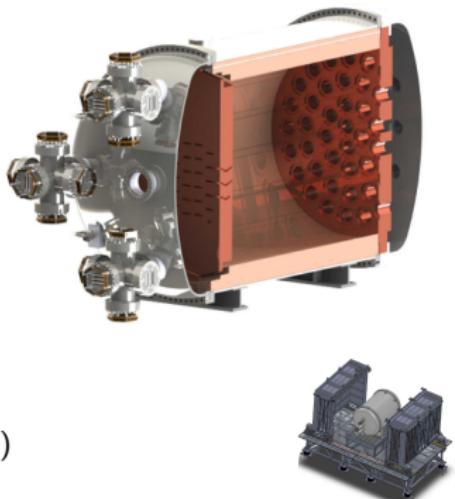


AMoRE pilot

- 6 crystals of $^{40}\text{Ca}^{100}\text{MoO}_4$ ($\sim 1.8 \text{ kg}$)
- several upgrades in detector/system design
- FWHM @ ^{208}Tl : $43 \text{ keV} \rightarrow 10 \text{ keV}$

Future players: NEXT-100

- Neutrino Experiment with a Xenon TPC
 - gas TPC with 100 kg of ^{136}Xe enriched xenon
 - background in ROI: $0.1 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
 - energy resolution: 15 keV FWHM @ $Q_{\beta\beta}$
 - sensitivity goal on $0\nu\beta\beta$ half-life: $5 \cdot 10^{25} \text{ yr}$
- result of a strong R&D programme
 - NEXT-WHITE: final validation prototype
 - signal amplification by electroluminescence
 - tracking plane (SiPMs) + energy plane (PMTs)



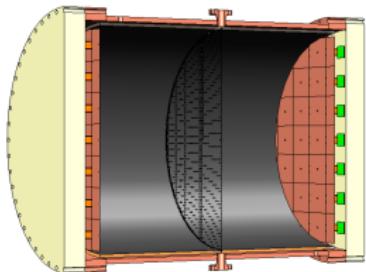
The NEXT program

- NEXT-100 → NEXT-250 → NEXT-ton
- background estimate: $0.05 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
- final expected sensitivity: 10^{27} yr

Future players: PandaX-III

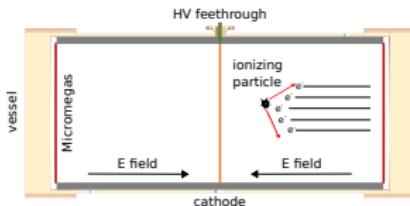
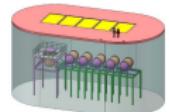
■ Particle and astrophysical Xenon Detector - III

- 5 gas TPCs with 200 kg of ^{136}Xe enriched xenon
- background in ROI: $0.01 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
- energy resolution: $75 \text{ keV FWHM} @ Q_{\beta\beta}$
- sensitivity goal on $0\nu\beta\beta$ half-life: 10^{27} yr



■ $0\nu\beta\beta$ search with the PandaX programme

- symmetric TPC instrumented with Microbulk MicroMegas
- extensive material screening campaign
- commissioning of prototype TPC (10 kg Xe) ongoing



First 200 kg module

- energy resolution: $225 \text{ keV FWHM} @ Q_{\beta\beta}$
- background in ROI: $0.1 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
- expected sensitivity: 10^{26} yr

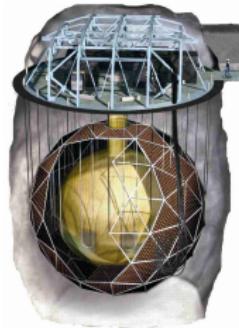
Future players: SNO+

■ Sudbury Neutrino Observatory +

- 3.9 t of tellurium dissolved in LS
- background in ROI: $0.1 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
- energy resolution: 270 keV FWHM @ $Q_{\beta\beta}$
- sensitivity goal on $0\nu\beta\beta$ half-life: $2 \cdot 10^{26} \text{ yr}$

■ commissioning ongoing

- tellurium stored underground
- purification system under construction
- calibration system ready
- loading of LS forthcoming



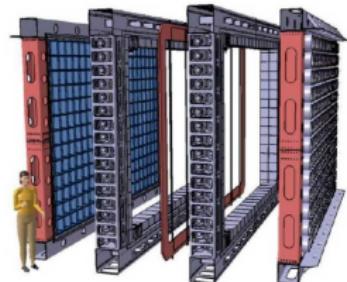
SNO+ $0\nu\beta\beta$ programme

- Te concentration in LS: 0.5% → 5%
- 13.3 t of isotope mass
- expected sensitivity: $> 10^{27} \text{ yr}$

Future players: SuperNEMO

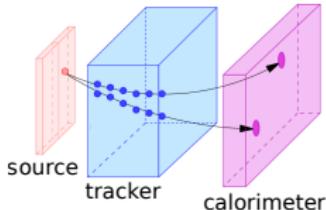
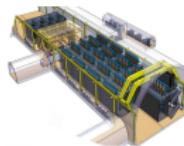
■ Super Neutrino Ettore Majorana Observatory

- tracker + calorimeter with 100 kg of ^{82}Se $\beta\beta$ source
- background in ROI: $5 \text{ c keV}^{-1} \text{ t}^{-1} \text{ yr}^{-1}$
- energy resolution: 120 keV FWHM @ $Q_{\beta\beta}$
- sensitivity goal on $0\nu\beta\beta$ half-life: 10^{26} yr



■ from the experience of NEMO-3

- improved detector design + modularity
- increased detector radio-purity
- increased source radio-purity

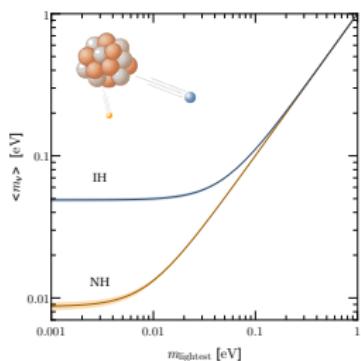


superNEMO demonstrator

- first produced module (1 of 20)
- 7 kg of isotope mass
- expected sensitivity: $6 \cdot 10^{24} \text{ yr}$

Assessing the neutrino masses

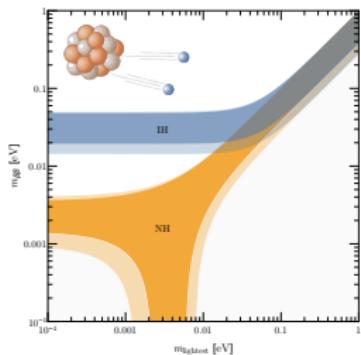
Direct measurement



- effect of electron neutrino mass in β decay
- model independent: relies on pure kinematic considerations
- sensitive to effective electron neutrino mass:

$$\langle m_\nu \rangle = \sqrt{\sum_{i=1}^3 |U_{ei}^2| m_i}$$

Search for $0\nu\beta\beta$

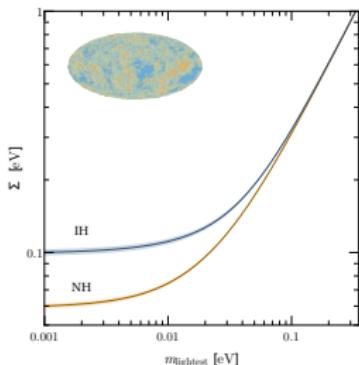


- requires Majorana neutrinos
- large theoretical uncertainties
- sensitive to Majorana mass:

$$m_{\beta\beta} = \left| \sum_{i=1}^3 U_{ei}^2 m_i \right|$$

- bands due to Majorana phases

Cosmology



- strong model dependence
- very stringent bounds
- sensitive to sum of neutrino masses:

$$\Sigma \equiv \sum_{i=1}^3 m_i$$

Implication for the $0\nu\beta\beta$ search (I)

- $\Sigma = m_1 + m_2 + m_3$

$$= m_I + \sqrt{m_I^2 + a} + \sqrt{m_I^2 + b}$$

- NH: $a = \delta m^2$

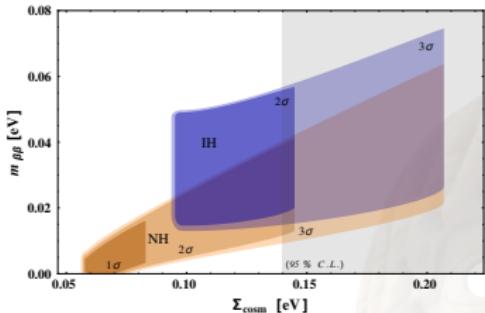
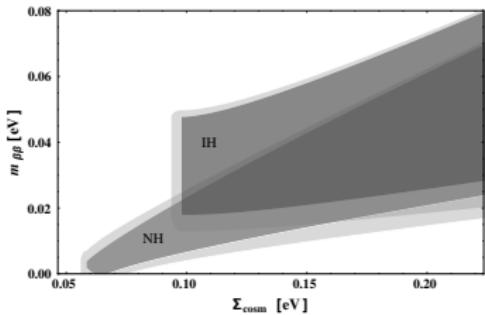
$$b = \Delta m^2 + \delta m^2/2$$

- IH: $a = \Delta m^2 - \delta m^2/2$

$$b = \Delta m^2 + \delta m^2/2$$

- it is possible to include the new constraints on Σ by considering:

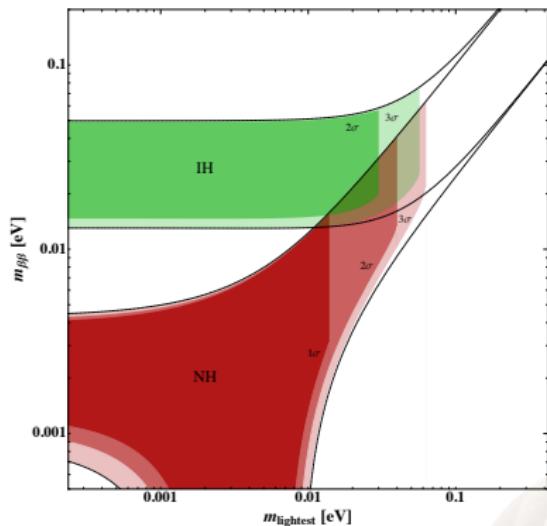
$$\frac{(y - m_{\beta\beta}(\Sigma))^2}{(n\sigma[m_{\beta\beta}(\Sigma)])^2} + \frac{(\Sigma - \Sigma(0))^2}{(\Sigma_n - \Sigma(0))^2} < 1$$



Implication for the $0\nu\beta\beta$ search (II)

$$\frac{(y - m_{\beta\beta}(m))^2}{(n \sigma[m_{\beta\beta}(m)])^2} + \frac{m^2}{m(\Sigma_n)^2} < 1$$

Mass spectrum	$m_{\beta\beta}^{\max}$ [meV] (C. L. on Σ)		
	1σ	2σ	3σ
NH	16	41	64
IH	-	57	75

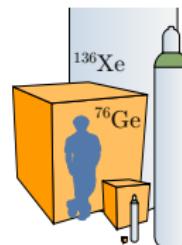
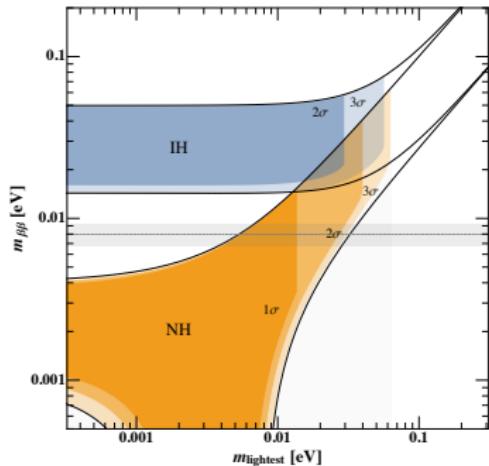


The IH region is excluded at 1σ

Further future challenges

- let us require a sensitivity: $m_{\beta\beta} = 8 \text{ meV}$
 - sign of Δm^2
 - test of cosmology
- $M T B \Delta = \mathcal{O}(1) \Rightarrow S_{1/2,0B}^{0\nu} \propto M T$

Isotope	$S_{1/2,0B}^{0\nu}$ [yr]	Exposure [t yr]	$(B \Delta)_{0B}$ [kg $^{-1}$ · yr $^{-1}$]
<u>g_A, nucleon</u>			
^{76}Ge	$3.0 \cdot 10^{28}$	5.5	$1.8 \cdot 10^{-4}$
^{130}Te	$8.1 \cdot 10^{27}$	2.5	$4.0 \cdot 10^{-4}$
^{136}Xe	$1.2 \cdot 10^{28}$	3.8	$2.7 \cdot 10^{-4}$
<u>g_A, quark</u>			
^{76}Ge	$7.9 \cdot 10^{28}$	14	$7.0 \cdot 10^{-5}$
^{130}Te	$2.1 \cdot 10^{28}$	6.5	$1.5 \cdot 10^{-4}$
^{136}Xe	$3.0 \cdot 10^{28}$	9.8	$1.0 \cdot 10^{-4}$
<u>g_A, phen.</u>			
^{76}Ge	$6.9 \cdot 10^{29}$	125	$8.0 \cdot 10^{-6}$
^{130}Te	$2.7 \cdot 10^{29}$	84	$1.2 \cdot 10^{-5}$
^{136}Xe	$4.0 \cdot 10^{29}$	130	$7.7 \cdot 10^{-6}$



(Many-)tonne
scale detectors
needed

