

#### Cosmic Rays in the multi-messanger era

#### Giovanni Morlino,

Gran Sasso Science Institute, L'Aquila, ITALY

Vulcano Workshop, 2018 – May 20-26, Vulcano, ITALY



### The path to became a cosmic ray



### The path to became a cosmic ray



### Predictions of diffusive shock acceleration

(1) Spectrum: 
$$f_{CR}(p) \propto p^{-4} \rightarrow f_{CR}(E) \propto E^{-2}$$
  
(2) Acceleration efficiency: ~10%  
(3) Maximum energy:  $E_{max} \simeq 50 \left(\frac{\delta B}{B_0}\right) \left(\frac{B_0}{\mu G}\right) eV$   
Strong dependence on magnetic field

High energies, up to PeV, can be achieved only if  $(\delta B/B_0)^2 >> 1$ 

This condition requires amplification of the magnetic field

G. Morlino, Vulcano Workshop — May 22, 2018

#### Fermi acceleration at work

#### [From Caprioli & Spitkovsky (2013)]



### PIC simulation of particle acceleration



### Non-thermal spectrum from SNRs



### Gamma-rays from SNRs



NOTE: this general trend could be an artifact of the environmental conditions

### Magnetic field amplification: observations

Chandra X-ray map. Data for the green sector are from Cassam-Chenaï et al (2007)



Thin non-thermal X-ray filaments provide evidence for magnetic field amplification

[Hwang el al(2002); Bamba et al (2005)]



$$\begin{cases} D = r_L c/3 \propto E B^{-1} \\ \tau_{syn} = \frac{3 m_e c^2}{4 \sigma_T c \gamma \beta^2 U_B} \propto E B^{-2} \end{cases}$$

### Where is the magnetic field amplified?

**DOWNSTREAM:** MHD instabilities (shear-like)

**UPSTREAM:** only through instabilities driven by CRs (Streaming, Bell)

BUT we need amplification upstream of the shock to reach high energies

Low magnetic field upstream produces a more extended emission NOT OBSERVED!



## Magnetic field amplification: Theory

How is the magnetic field amplified?

#### **Resonant Straming instability**

[e.g. Skilling (1975), Bell & Lucek (2001), Amato & Blasi (2006), Blasi (2014)] Particles amplify Alfvèn waves with wave-number  $k=1/r_{I}(p)$  Fast growth rate but

$$\left(\frac{\delta B/B_0}{2}\right)^2 \simeq few$$

$$E_{max} \approx 10^5 \, GeV$$
A factor 10 below the knee

## Magnetic field amplification: Theory

How is the magnetic field amplified?

#### **Resonant Straming instability**

[e.g. Skilling (1975), Bell & Lucek (2001), Amato & Blasi (2006), Blasi (2014)]



Fast growth rate but



A factor 10 below the knee



## Magnetic field amplification: Theory

How is the magnetic field amplified?



### The path to became a cosmic ray



#### Particle escape from SNRs

If particles are not released all at the same time, in general:

Spectrum injected into the Galaxy  $f_{esc}(p) \neq f_{SNR}(p)$ 

Spectrum inside SNRs

#### Particle escape from SNRs

If particles are not released all at the same time, in general:

Spectrum injected into the Galaxy  $f_{esc}(p) \neq f_{SNR}(p)$  Spectrum inside SNRs

Assume that at time t only particles at maximum momentum  $p_{max}(t)$  can escape

$$4\pi f_{esc}(p) c p p^{2} dp = \xi_{esc}(t) \frac{1}{2} \rho V_{sh}^{3} 4\pi R_{sh}^{2} dt$$
  
Released energy Converted Incoming  
fraction energy flux

### Particle escape from SNRs

If particles are not released all at the same time, in general:

Spectrum injected into the Galaxy

$$\begin{array}{ccc} \text{ed} \\ \text{xy} & f_{esc}(p) \neq f_{SNR}(p) & \begin{array}{c} \text{Spectrum} \\ \text{inside SNRs} \end{array} \end{array}$$

Assume that at time *t* only particles at maximum momentum  $p_{max}(t)$  can escape



## Effect of self-amplification near the CR sources

During the process of escaping, CR can excite magnetic turbulence (via streaming instability) that keep the CR close to the SNR for a long time, up to  $\sim 10^5$  yr

[Malkom et al. (2013) Nava et al. (2015)]

During this time CR spend in the vicinity of sources they can produce diffuse emission via  $\pi^0 \rightarrow \gamma \gamma$ 



## Effect of self-amplification near the CR sources

During the process of escaping, CR can excite magnetic turbulence (via streaming instability) that keep the CR close to the SNR for a long time, up to  $\sim 10^5$  yr

[Malkom et al. (2013) Nava et al. (2015)]

During this time CR spend in the vicinity of sources they can produce diffuse emission via  $\pi^0 \rightarrow \gamma \ \gamma$ 

If a molecular cloud is close enough the enhanced  $\gamma$ -ray emission will be seen for long time

CTA will probably discover tens of SNR-MC associations



### Using molecular clouds as CR barometer

Examples of y-ray emission from clouds close or interacting with SNRs - [*Fermi*-LAT]



#### **OBSERVATIONS of MCs in γ-RAYS:**

- CRs interact inside MCs  $pp \rightarrow \pi^0 \rightarrow \gamma \gamma$
- strong emission in GeV range
- $\gamma$ -emission sensible to CR energy E > 280 MeV

#### **DETECTION OF IONIZATION**

• The ionization rate of several molecules depends on the CR flux  $(H_2, H_3^+, CH, OH, C_2, DCO^+, HCO^+,....)$ 

• Ionization sensible to CR energy E > 0.1 MeV

Is it possible to use combined information from ionization and y-ray emission to infer the CR spectrum from ~MeV up to ~TeV

#### Using molecular clouds as CR barometer



### Diffusion in the Galactic Halo



### Basic Halo model

In the basic picture of CR propagation model:

- CRs diffuse in a magnetic halo larger than the Galactic disc
- The CR distribution vanish at z = H ( $H \sim 3-4$  kpc from diffuse synchrotron emission)
- The diffusion coefficient D(E) is assumed constant everywhere in the halo



Suggesting Kolmogorov turbulence

This picture is unsatisfactory for at least two reasons:

- Which is the physical meaning of *H* ?
- What generates the diffusion?



#### Some observed anomalies suggest a more complex propagation model

### Anomaly: spectral hardening



Recent measurements by PAMELA and AMS-02 revealed the existance of a fine structure:

At rigidity of ~300 GV all spectra show a spectral hardening

$$N_{p}(E) = Q_{inj}(E) \frac{H^{2}}{2D(E)}$$

Either the injected spectrum or the diffusion present a break at ~300 GV

See talk by M. Vecchi

#### Self-generated magnetic Halo

(Evoli, Blasi, GM, Aloisio, 2018, submitted]

Large scale turbulence generated inside the Galaxctic disc by SN explosion advected and decaying through Kolmogorov cascade.

Dimention of the Halo determined by Advection time = decaying time



#### Self-generated magnetic Halo

(Evoli, Blasi, GM, Aloisio, 2018, submitted]



#### Self-generated magnetic Halo

[Evoli, Blasi, GM, Aloisio, 2018, submitted]

Large scale turbulence generated inside the Galaxctic disc by SN explosion advected and decaying through Kolmogorov cascade.

#### Dimention of the Halo determined by Advection time = decaying time

[see also: Blasi, Amato Serpico (2012) Aloisio, Blasi, Serpico (2015)]



#### **Predicted proton spectrum**

G. Morlino, Vulcano Workshop — May 22, 2018

 $10^{4}$ 

 $10^{5}$ 

## The problem of the cosmic ray gradient in the Galactic plane seen by Fermi-LAT



Recent results from FermiLAT collaboration on the CR distribution in the Galactic plane

[Acero et al. arXiv:1602.07246]

- In the outer region (R > 8kpc) the CR density at ~20 GeV is flat (i.e. decreases much slower than the source distribution)

- In the inner region the CR density has a peak at  $\sim 3$  kpc

## The problem of the cosmic ray gradient in the Galactic plane seen by Fermi-LAT



Recent results from FermiLAT collaboration on the CR distribution in the Galactic plane

[Acero et al. arXiv:1602.07246]

- In the outer region (R > 8kpc) the CR density at ~20 GeV is flat (i.e. decreases much slower than the source distribution)

- In the inner region the CR density has a peak at  $\sim 3$  kpc

## The problem of the cosmic ray gradient in the Galactic plane seen by Fermi-LAT



Recent results from FermiLAT collaboration on the CR distribution in the Galactic plane

[Acero et al. arXiv:1602.07246]

- In the outer region (R > 8kpc) the CR density at ~20 GeV is flat (i.e. decreases much slower than the source distribution)

- In the inner region the CR density has a peak at  $\sim 3$  kpc

- The slope @ 20 GeV is not constant

This scenario is difficult to accommodate in a standard diffusion model

### **Possible solutions**

In the context of leaky-box model several solution have been proposed:

- Extended halo, H > 4 kpc (Dogiel, Uryson, 1988; Strong et al.,1988; Bloemen, 1993, Ackerman et al., 2011)
  - ^ predices a flat spectrum (but not flat enough)
  - ^ cannot explain the denity bump in the inner Galaxy
- Flatter distribution of SNR in the outer Galaxy (Ackerman et al., 2011)
- Enhancement of CO/H<sub>2</sub> density ratio (X<sub>CO</sub>) in the outer Galaxy (Strong et al., 2004)
- Injection dependence on the ISM temperature (Erlykin et al., 2015)
- Advection effects due to the Galactic wind (Bloemen, 1993; Breitschwerdt, Dogiel, Voelk, 2002)

#### CAN SELF-GENERATED DIFFUSION EXPLAIN THE OBSERVATIONS?

None of these ideas can simoultaneously account for all signatures

- flatness R > 8 kpc,
- peak at R~3-4 kpc,
- variation in the slope

#### Self-generated turbulence and the gradient problem (Recchia, Blasi, GM, MNRAS 462, 2016)

Self-generated turbulence could explain the gradient and the spectral index changes because it is more effective where *B* is smaller

^ less effective in the inner Galaxy more effective in the outer Galaxy Λ

Strength of large scale magnetic field in the Galaxy





#### Why a Galactic magnetic Halo?

#### **Evidences for the Galactic magnetic halo:**

Detection of magnetic field around other galaxies
 Detection of synchrotron emission around the Milky Way
 Evidence of hot plasma from X-ray emission and absorption lines

#### What is the origin of the magnetic Halo?

Sometimes the X-shaped magnetic field structure in the halo is accompanied by strong vertical fields above and below the central region of the disk.

These observations support the idea of a "galactic wind" which is driven by the energy of star formation processes in the disk and transports gas, magnetic fields and cosmic-ray particles into the halo.

The speed of the outflow can be measured from radio observations and it is of the order of 300 km/s.





#### Evidence from X-ray emission and absorption lines

Thermal X-ray emission has been observed from the region around sturbust galaxies.

 In some "normal" galaxies the presence of a hot temperature gax (T~10<sup>6</sup> K) has been inferred from absorption lines in X-rays (especially lines OVI, OVII and OVIII)

• Also the Milky Way presents the same absorption lines [e.g. Kalberla & Dedes (2008), Miller & Bregman (2013)]

• From those lines the total mass of the halo can be estimated

 $M_{\rm halo} \sim 10^{10} \, M_{\rm sol}$ 

(comparable with the total barionic mass in the disk!!)

• And also the matallicity:  $Z \sim 0.2-0.3$ 

→ The halo has been probably polluted by a Galactic wind
 BUT In the Milky Way thermal pressure alone is unable to drive a Galactic wind.

#### Galactic wind observed in X-rays from starburst galaxy M82



#### Can CRs drive a Galactic wind?

• CRs behave as a relativisic gas  $\Rightarrow$  pressure drops less than thermal pressure  $\Rightarrow$  can drive a wind

• Properties of the wind can explain the hot baryionic halo around Milky Way

- Wind speed increases up to  $\sim few \ 100 \ km/s$ 

- The magnetic turbulence cascade can increase the wind temperature up to  $10^6\,\mathrm{K}$ 

• Wind can produce a mass loss rate  $\sim 0.5 M_{\odot}/yr$ 

 $\Rightarrow$  in agreement with star formation models

 $\Rightarrow$  in agreement with the Halo metallicity > 0

● The CR halo ends where advection dominated over diffusion ⇒ energy dependent halo size

[Breitschwerdt, McKenzie & Vòlk (1993), Everett & Zweibel (2011), Recchia, Blasi, GM (2016) Pakmor, Pfrommer et al. (2016)]





#### Conclusions

#### ACCELERATION

- DSA is a well developed theory but we need to account for different ambient media where SNRs expand
- Strong evidence for magnetic field amplification induced by CRs
- Lack of PeVatrons

#### ESCAPING

- Escaping process is a key issue and is not undestood yet.
- Possible way to study it
  - Emission from molecular cloud close to SNR
  - Contribution of escaping particles to diffuse γ-ray emission

#### **PROPAGATION IN THE GALAXY**

- We still lack of a realstic description of Galactic propagation
- Going beyond the simple view of the leaky-box model is required by data
  - Self-generated turblence produced via streaming instability could play a major role for the propagation of CRs with *E* <~ 100 GeV</p>
  - CR pressure could drive a Galactic wind

### Some anomalies in the CR composition

Low energy CR (~100 MeV/nucleon) give contradicting information

$$\left(\frac{{}^{22}Ne}{{}^{20}Ne}
ight)_{CR} \simeq 5\left(\frac{{}^{22}Ne}{{}^{20}Ne}
ight)_{So}$$

Binns et al., ApJ 634, (2005) ACE-CRIS collaboration



Preferential acceleration of wind material in massive stars Prantzos, A&A 538 (2012)

$${}^{60}Fe\big|_{CR} \simeq 500 \big({}^{60}Fe\big|_{ISM}$$

Binns et al., Science 352, (2016) ACE-CRIS collaboration Requires acceleration of ejecta material in SNRs (possible role of reverse shock?)

It is not clear how to combine these two constraints

# 1-D slab model with self-generated turbulence (Recchia, Blasi,

(Recchia, Blasi, GM, MNRAS 462, 2016)

• CR escaping from the Galactic plane produce magnetic turbulence through resonant streaming instability

$$\Gamma_{\rm cr} = \frac{16\pi^2}{3} \frac{v_A}{\mathcal{F}(k)B_0^2} \left[ p^4 v(p) \frac{\partial f}{\partial z} \right]_{p=eB_0/kc}$$

• Turbulence scatter CRs (mainly) along large scale mag. field lines with Bohm-like diffusion coefficent

$$D(z,p) = \frac{r_L(p)v(p)}{3} \left[\frac{1}{\mathcal{F}(k)}\right]_{k=1/r_L}$$

• CRs are also advected by the global motion of the waves at the Alfvén speed



### Using the diffuse Galacic y-ray emission



Energy(GeV)

# Contribution of the escaping CRs to the diffuse Galactic emission

