S-matrix Bootstrap revisited

João Penedones



S-matrix Bootstrap I: QFT in AdS [arXiv:1607.06109]

S-matrix Bootstrap II: two-dimensional amplitudes [arXiv:1607.06110]

S-matrix Bootstrap III: higher dimensional amplitudes [arXiv:1708.06765]

S-matrix Bootstrap IV: multiple amplitudes [work in progress]

with M. Paulos, J. Toledo, B. Van Rees, P. Vieira

Motivation

Bootstrap Philosophy: bound the space of theories by imposing consistency conditions on physical observables.

Goal: extend recent success in CFT to massive QFT.

[Rattazzi, Rychkov, Tonni, Vichi '08] + many others

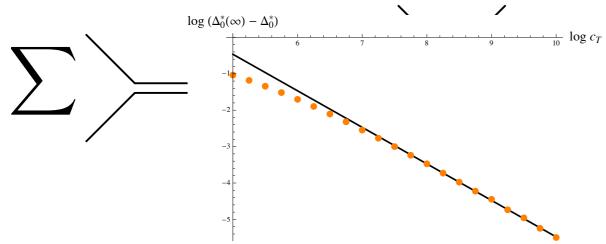
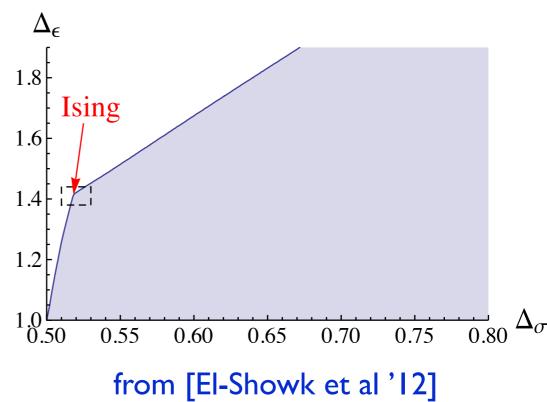


Figure 2: Upper bounds on Δ_0^* (the smallest conformal dimension of a spin-0 long multiplet appearing in the $\mathcal{O}_{35_c} \times \mathcal{O}_{35_c}$ OPE) for large values of c_T . The bounds are computed with $j_{\text{max}} = 20$ and $\Lambda = 19$. The long multiplets of spin j > 0 are only restricted by unitarity. The best fit for the last ten points (shown in black) is $\log(\Delta_0^*(\infty) - \Delta_0^*) = 4.55 - 1.00 \log c_T$.



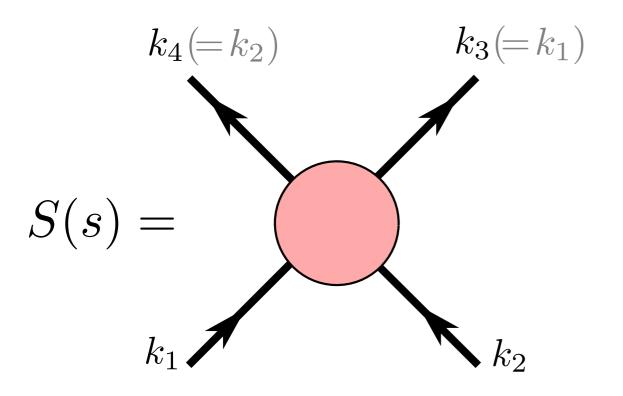
Revisit the $S_{-\infty}$ matrix Bootstrap program of the 60's and 70's.

excluded

Outline

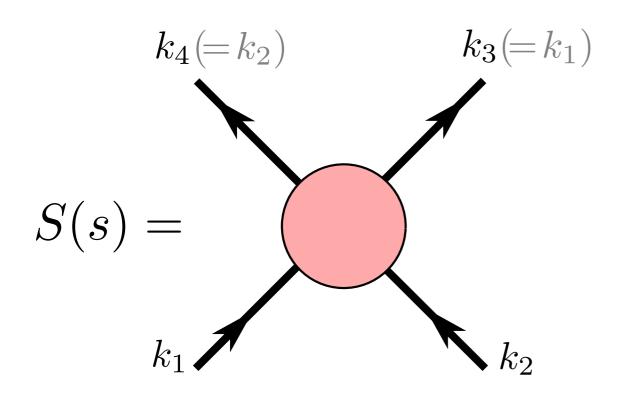
- S-matrix Bootstrap in D=2
- S-matrix Bootstrap in D>2
- Multiple Amplitudes Bootstrap in D=2
- Open questions

S-matrix Bootstrap in 2D QFT



$$k_i^2 = m^2$$

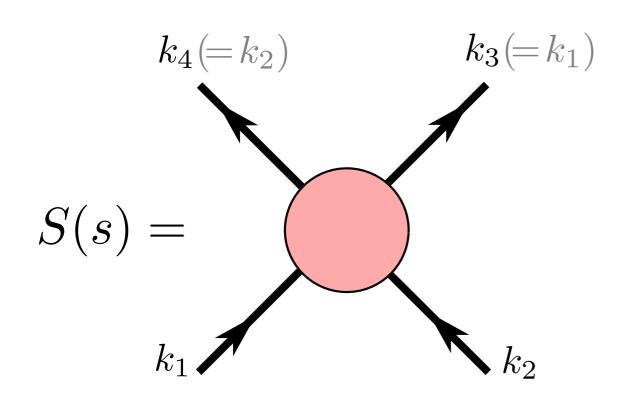
 $s \equiv (k_1 + k_2)^2$
 $t \equiv (k_2 - k_3)^2 = 4m^2 - s$
 $u \equiv (k_3 - k_1)^2 = 0$



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Crossing symmetry: $S(s) = S(4m^2 - s)$



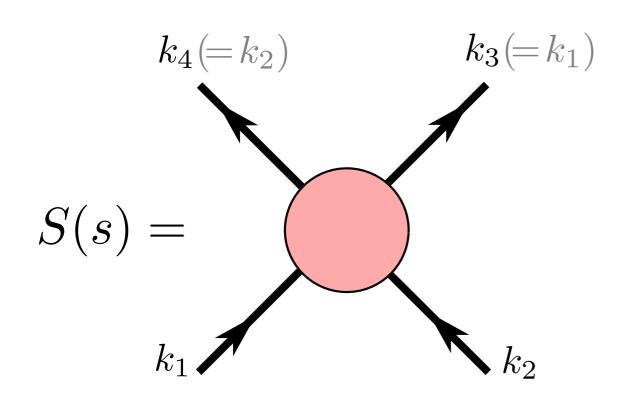
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Analyticity follows from mass spectrum.

 $\frac{1}{m_b}$ $\frac{2m}{m_b}$ $\frac{1}{m_b}$

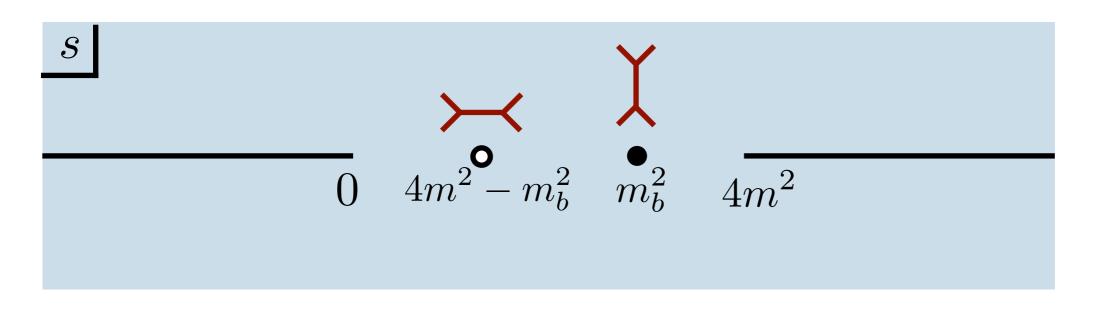


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2m
m_b
m

Crossing:
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Analyticity: $S(s^*) = [S(s)]^*$

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$$S(s) = [S(s)]^*$$

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Unitarity:
$$|S(s)|^2 \le 1$$
, $s > 4m^2$.

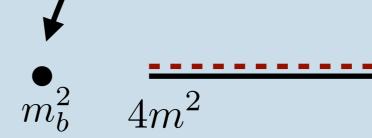
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$$S(s^*) = [S(s)]^*$$

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Unitarity:

$$|S(s)|^2 \le 1$$
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Question: for given spectrum,

$$\max g_b^2 = ?$$

$$S_{opt}(s) = \frac{\sqrt{s(4m^2 - s)} + \sqrt{m_b^2(4m^2 - m_b^2)}}{\sqrt{s(4m^2 - s)} - \sqrt{m_b^2(4m^2 - m_b^2)}} \equiv [m_b](s) \quad \text{[Creutz '72]}$$

$$CDD \text{ factor}$$

$$[Castillejo, Dalitz, Dyson]$$

Pole at
$$s=m_b^2>2$$

No particle production $|S_{opt}(s)|^2 = 1$, $s > 4m^2$.

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$$h(s) \equiv \frac{S(s)}{[m_b](s)} \Rightarrow \frac{h(s)}{|h(s)| \le 1}$$
 bounded at all boundaries

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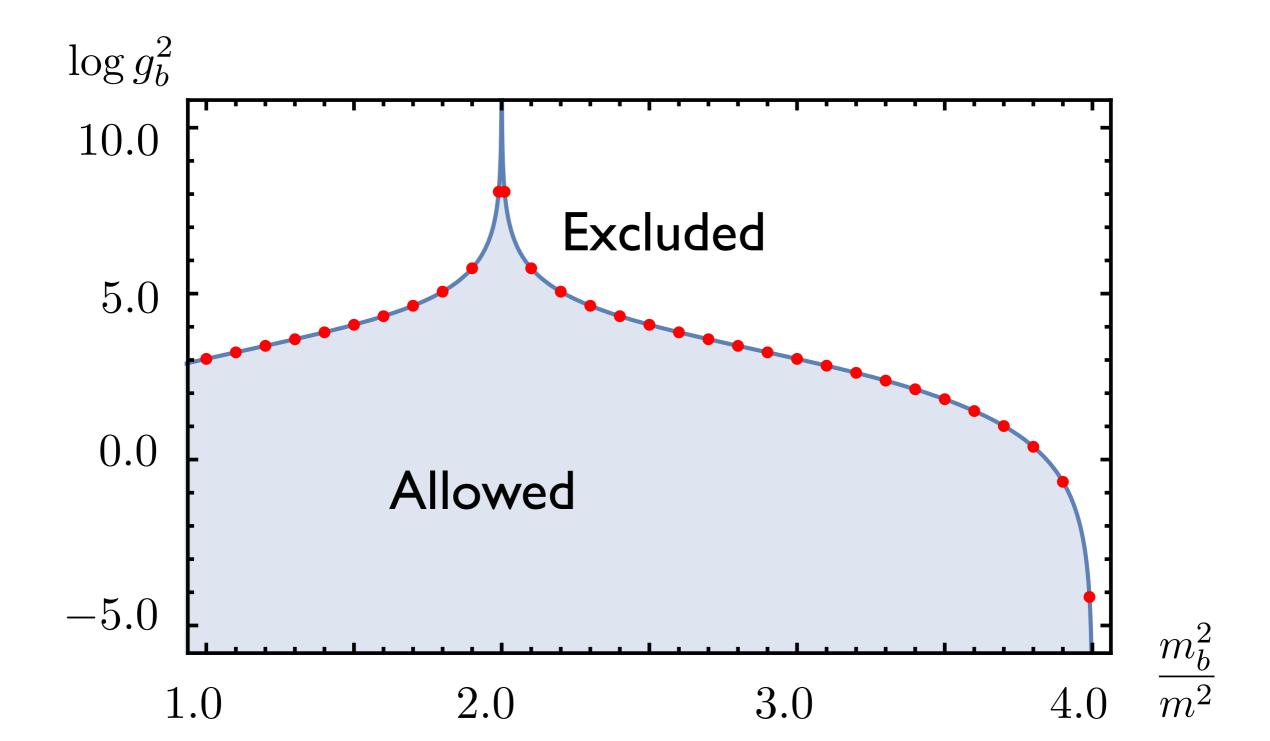
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$$|h(m_b^2)| = \left| \frac{g_b^2}{\text{Res}_{s=m_b^2}[m_b](s)} \right| \le 1$$

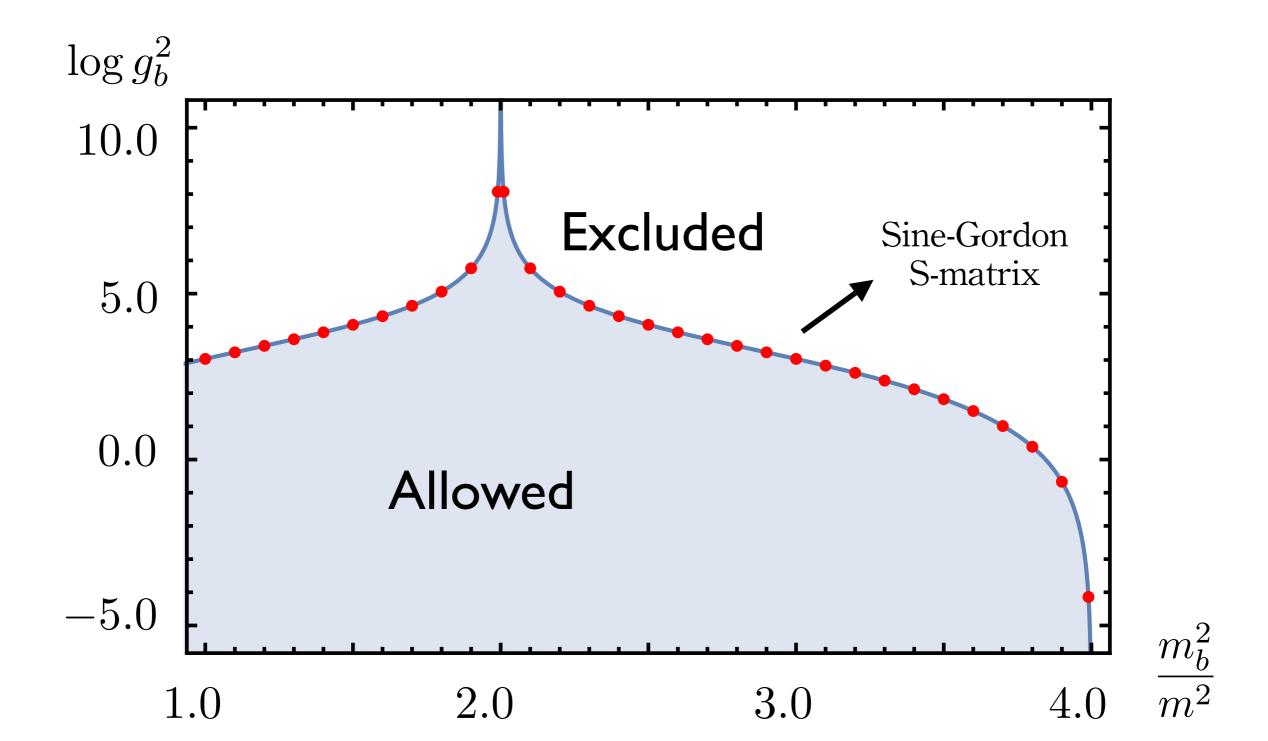
Maximum cubic coupling

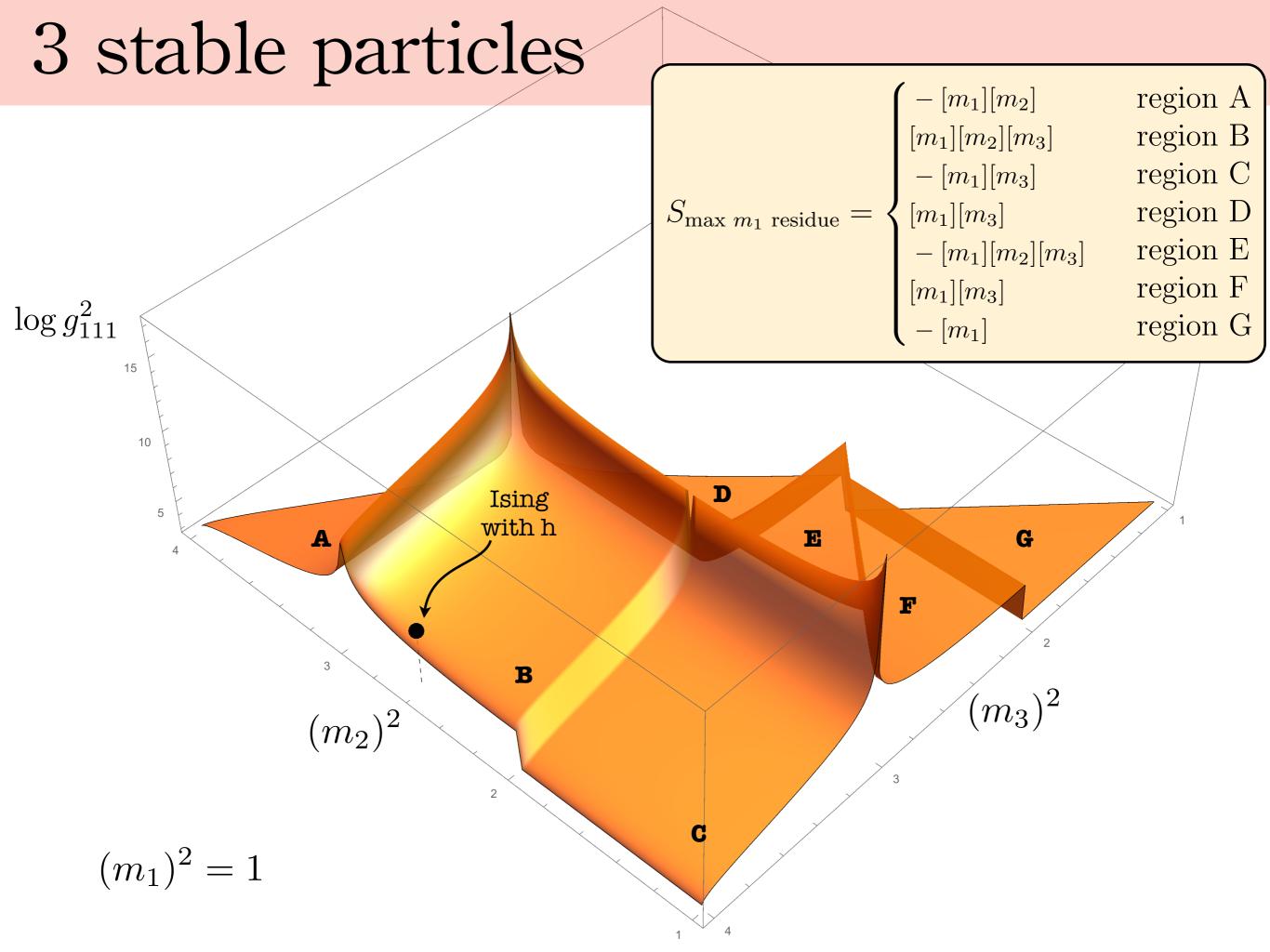
$$S_{opt}(s) = \frac{\sqrt{s(4m^2 - s)} + \sqrt{m_b^2(4m^2 - m_b^2)}}{\sqrt{s(4m^2 - s)} - \sqrt{m_b^2(4m^2 - m_b^2)}}$$

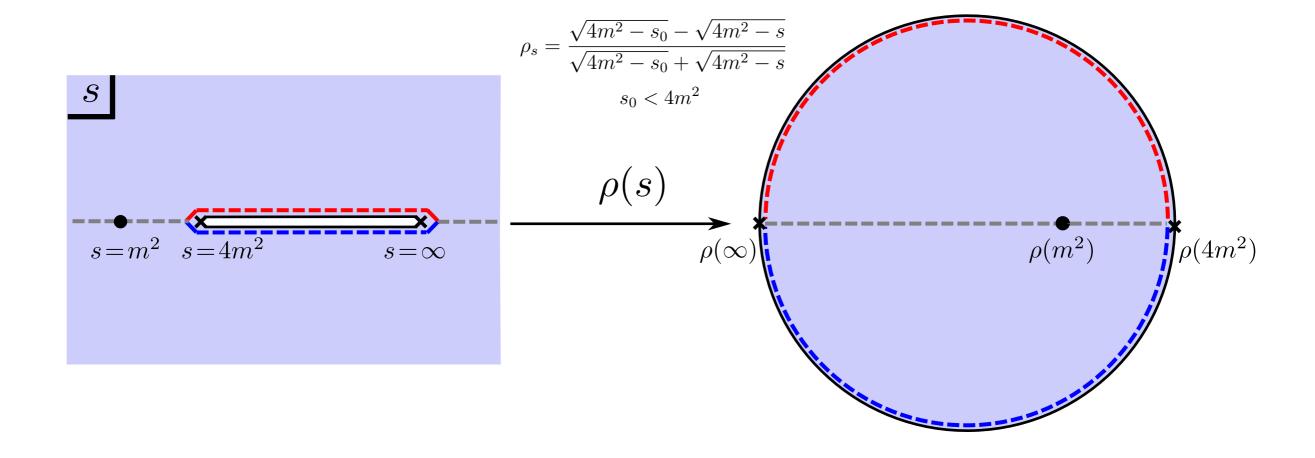


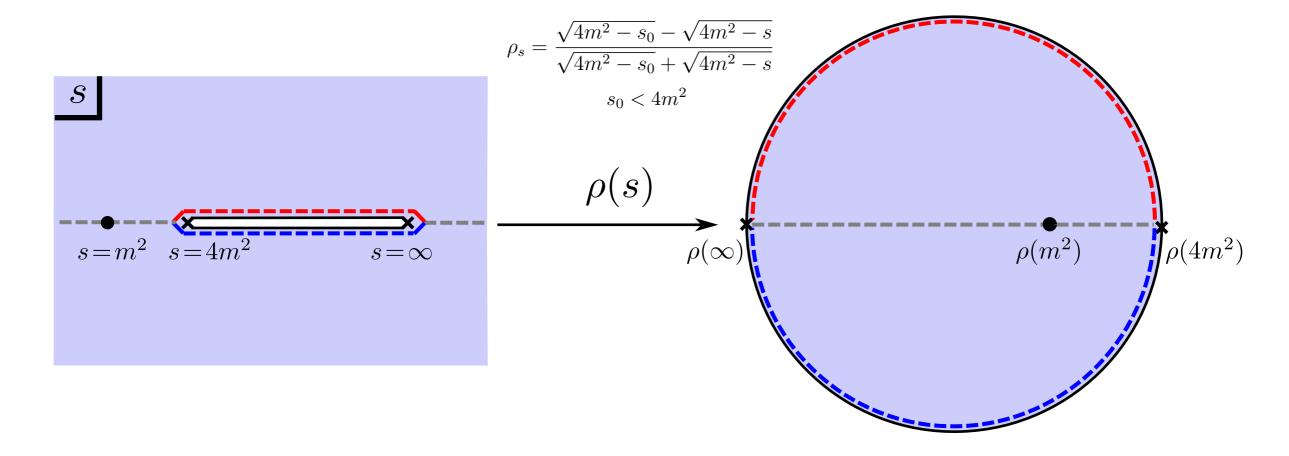
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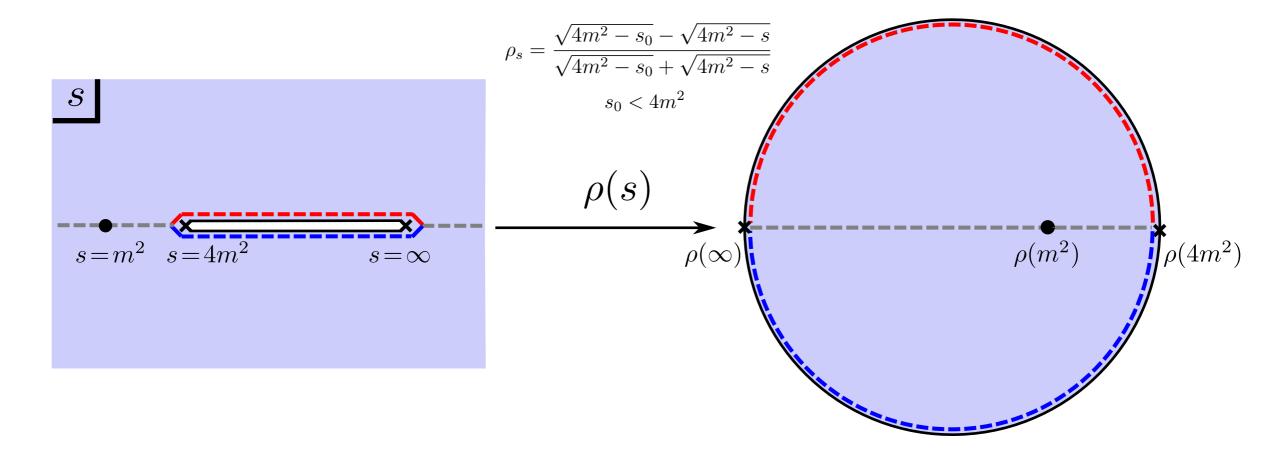








$$S_{ext}(s,t) = \frac{g_b^2}{s - m_b^2} + \frac{g_b^2}{t - m_b^2} + \sum_{a,b=0}^{a} c_{(ab)} \rho_s^a \rho_t^b$$



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Crossing symmetry and analyticity are automatic. Unitarity gives quadratic constraints:

$$|S_{ext}(s, 4m^2 - s)|^2 \le 1, \qquad s > 4m^2$$

Ansatz:

$$S_{ext}(s,t) = \frac{g_b^2}{s - m_b^2} + \frac{g_b^2}{t - m_b^2} + \sum_{a,b=0}^{a} c_{(ab)} \rho_s^a \rho_t^b$$

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Truncate to finite number of variables and quadratic constraints

$$a+b \leq N_{\max}$$

$$\{g_b^2, c_{(ab)}\}$$
at $s = s_1, s_2, \dots, s_M$

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 at $s = s_1, s_2, \dots, s_M$

[Simmons-Duffin '15]

Use semidefinite programming (SDPB) to maximize g_b^2 subject to these constraints. This reproduces the analytic solution as $N_{\rm max} \to \infty$

S-matrix Bootstrap in d+1 QFT

$$\langle \mathbf{p}_3, \mathbf{p}_4 | S | \mathbf{p}_1, \mathbf{p}_2 \rangle = 1 + i(2\pi)^{d+1} \delta^{(d+1)}(p_1 + p_2 - p_3 - p_4) T(s, t, u)$$

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Crossing symmetry & Analyticity:

$$T(s,t,u) = \frac{g_b^2}{s - m_b^2} + \frac{g_b^2}{t - m_b^2} + \frac{g_b^2}{u - m_b^2} + \sum_{a,b,c=0} \alpha_{(abc)} \rho_s^a \rho_t^b \rho_u^c$$

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Partial waves: Gegenbauer polynomial
$$x = \cos \theta$$

$$S_{\ell}(s) = 1 + i \frac{(s - 4m^2)^{\frac{d-2}{2}}}{\sqrt{s}} \int_{-1}^{1} dx (1 - x^2)^{\frac{d-3}{2}} P_{\ell}^{(d)}(x) \left. T(s, t, u) \right|_{\substack{t \to -\frac{1-x}{2}(s - 4m^2) \\ u \to -\frac{1+x}{2}(s - 4m^2)}}$$

Unitarity:
$$|S_{\ell}(s)|^2 \le 1$$
, $s > 4m^2$, $\ell = 0, 2, 4, ...$

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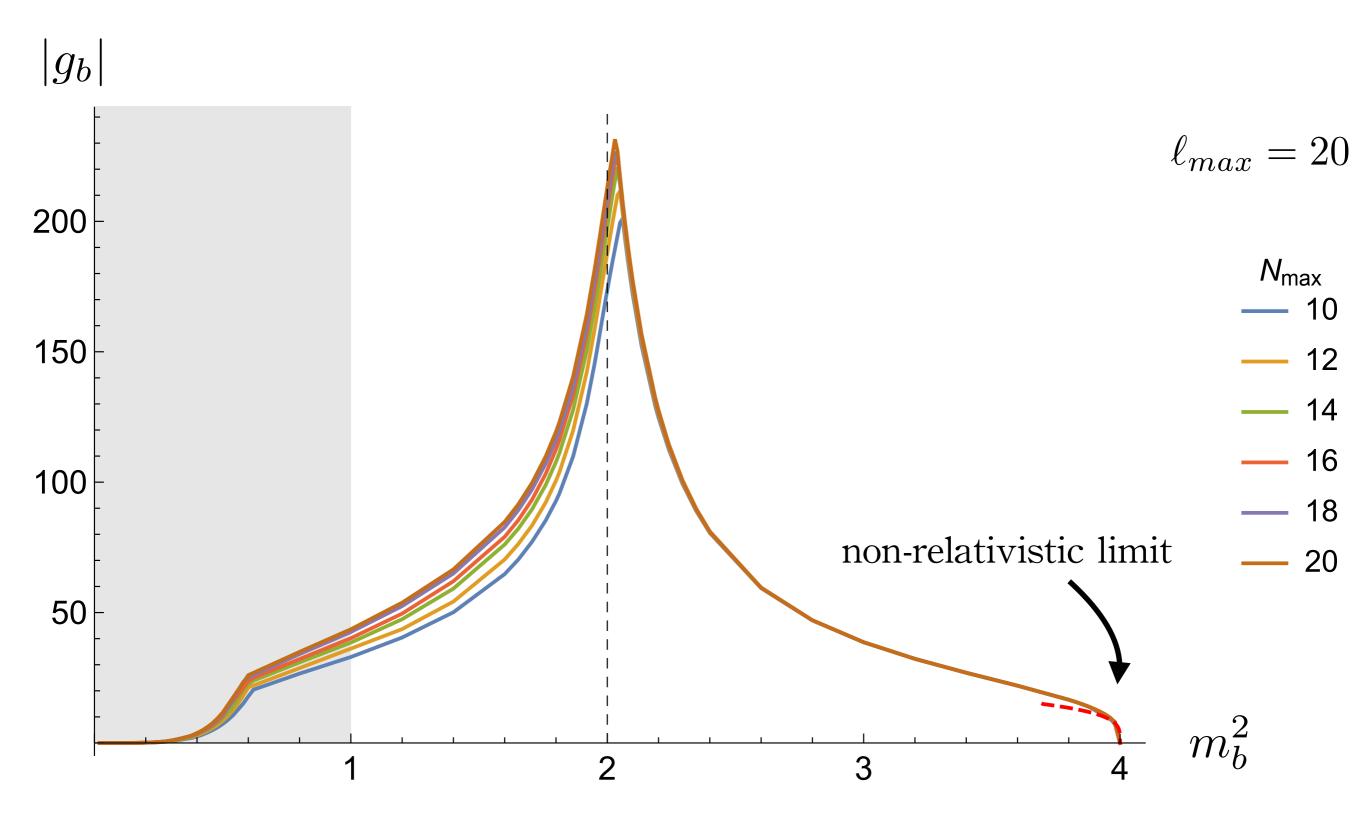
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, $s > 4m^2$, $\ell = 0, 2, 4, \dots \ell_{\text{max}}$

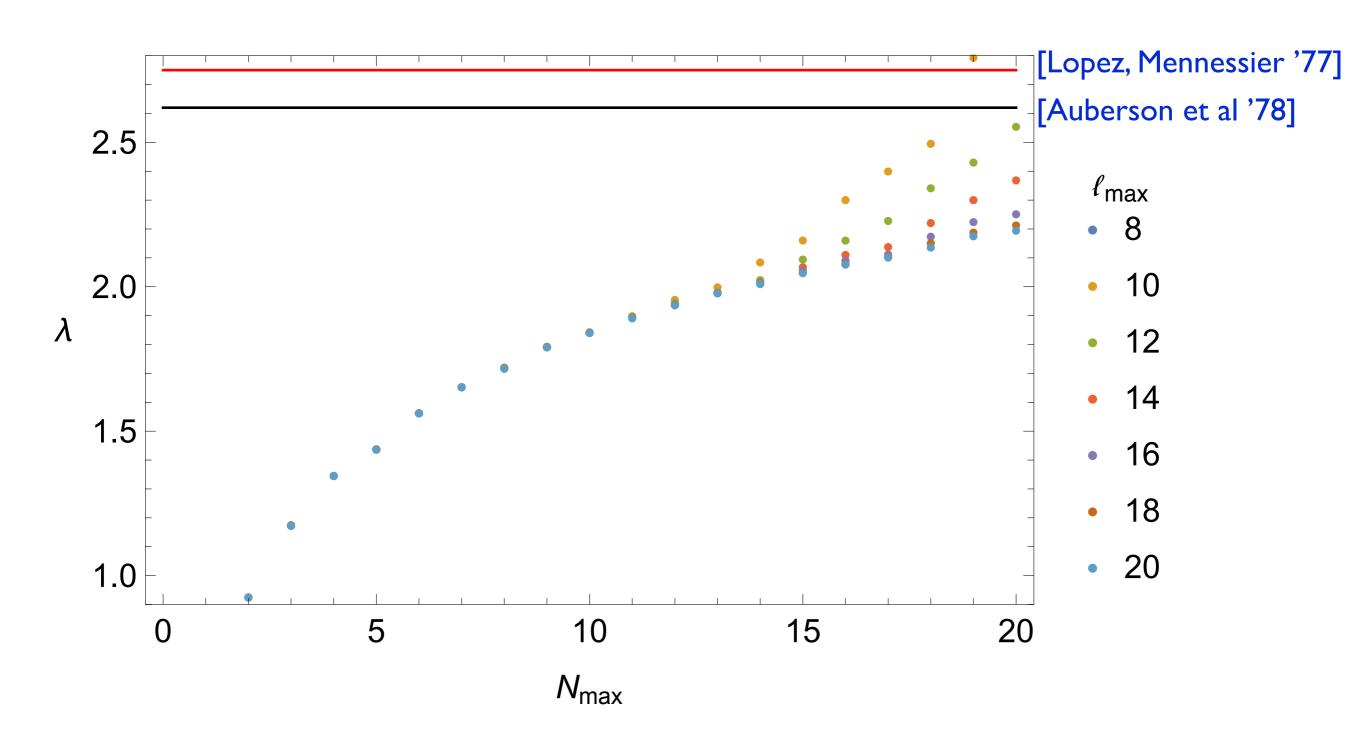
 \Rightarrow Quadratic constraints on the variables $\{g_b^2, \alpha_{(abc)}\}$ $a+b+c \le N_{\text{max}}$

Maximal cubic coupling in 3+1 QFT



Ansatz with no poles. Maximize
$$\lambda=\frac{1}{32\pi}T(s=t=u=\frac{4}{3}m^2)$$
 (e.g. $\pi^0\pi^0\to\pi^0\pi^0$)

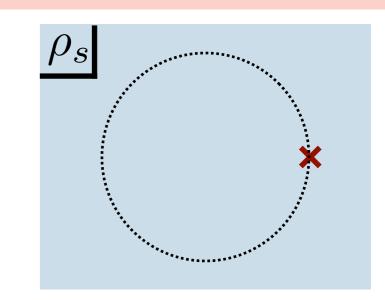
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Improved ansatz with threshold bound state:

$$T(s,t,u) = \beta \left(\frac{1}{\rho_s - 1} + \frac{1}{\rho_t - 1} + \frac{1}{\rho_u - 1} \right) + \sum_{a,b,c=0} \alpha_{(abc)} \rho_s^a \rho_t^b \rho_u^c$$

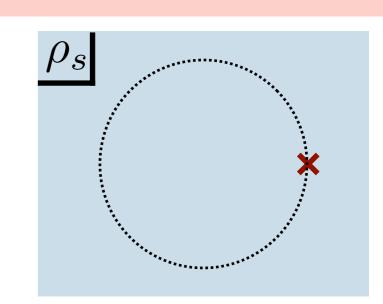
$$a + b + c \le N_{\text{max}}$$

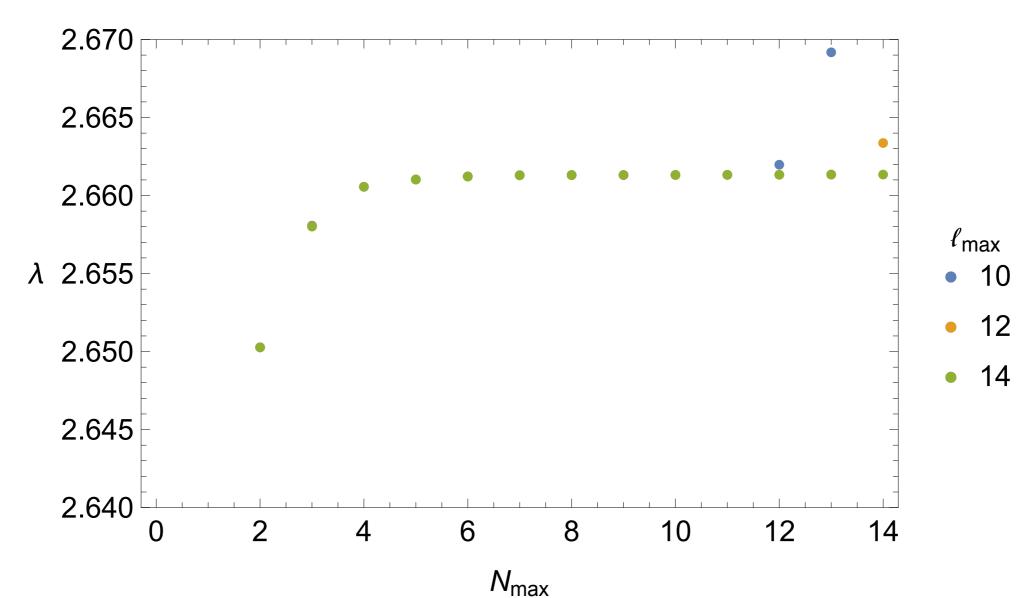


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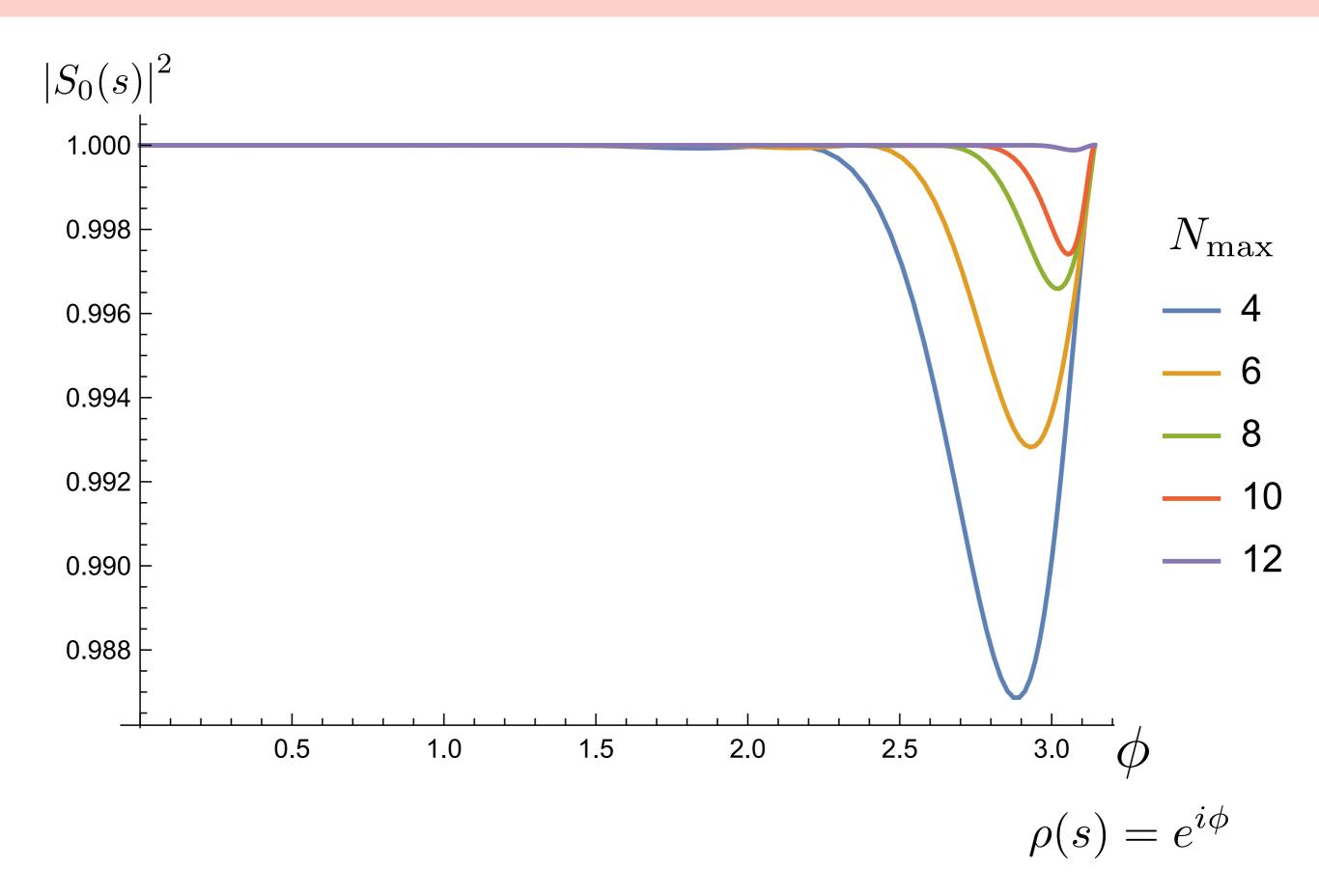
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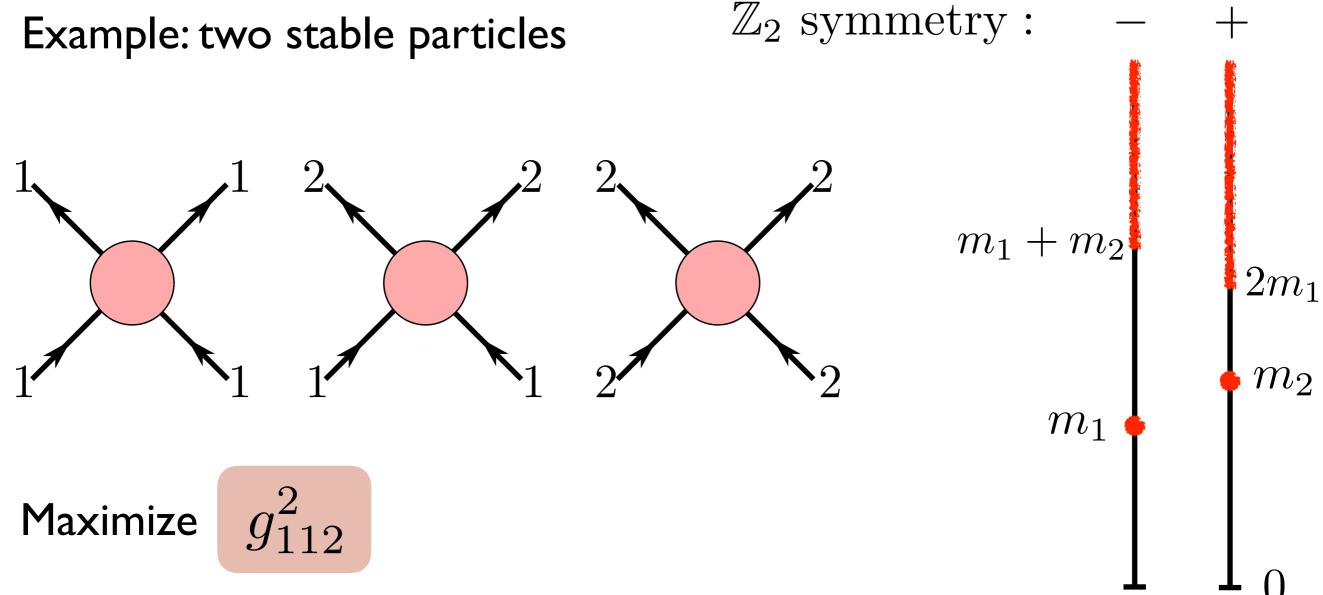
No particle production?



Multiple Amplitudes Bootstrap in 2D QFT

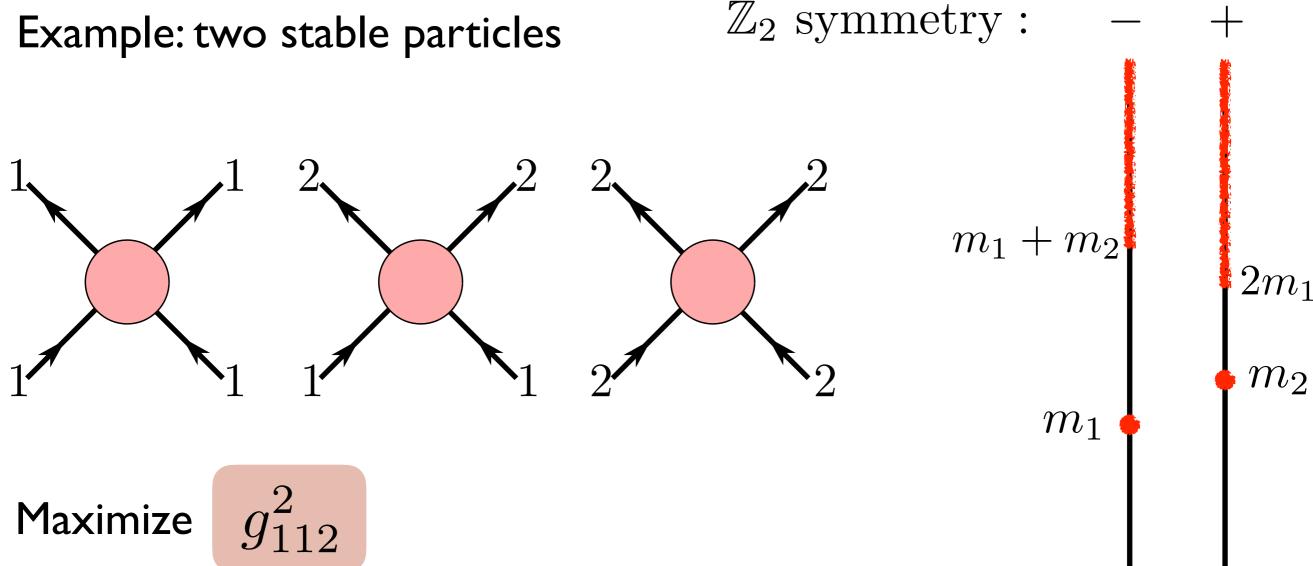
2 to 2 Scattering Amplitudes

Example: two stable particles



2 to 2 Scattering Amplitudes

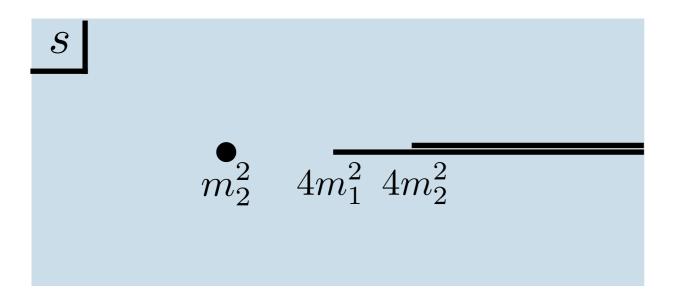
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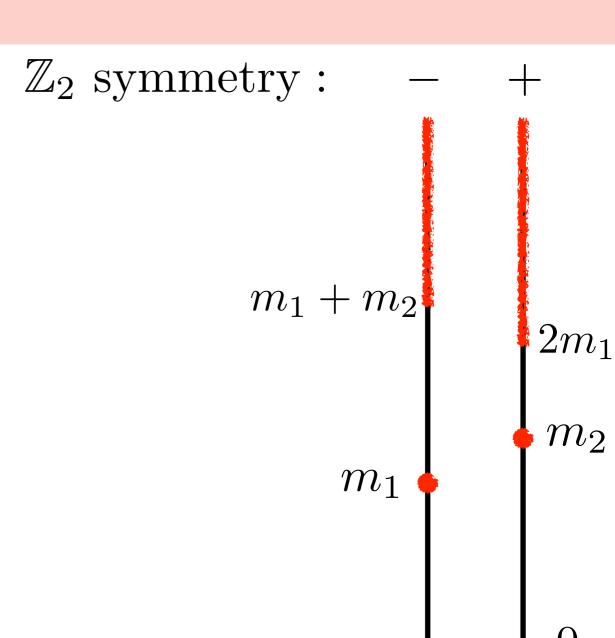


Unitarity: $|S_{11\to 11}|^2 + |S_{11\to 22}|^2 \le 1$

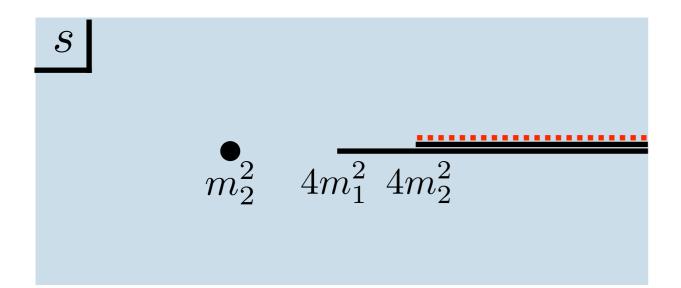
Not zero in optimal solution

Analyticity:

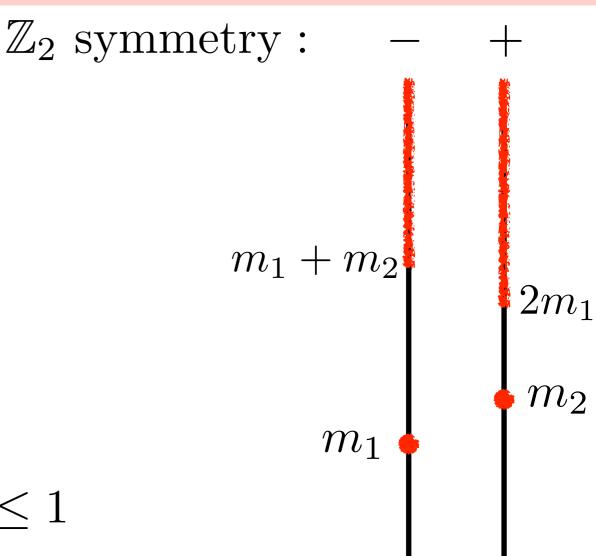




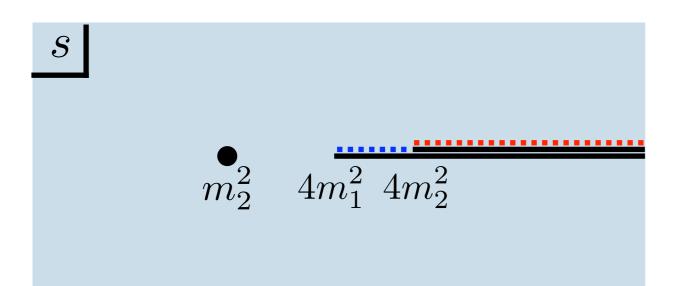
Analyticity:



Unitarity:
$$|S_{22\to 11}|^2 + |S_{22\to 22}|^2 \le 1$$
 $s \ge 4m_2^2$



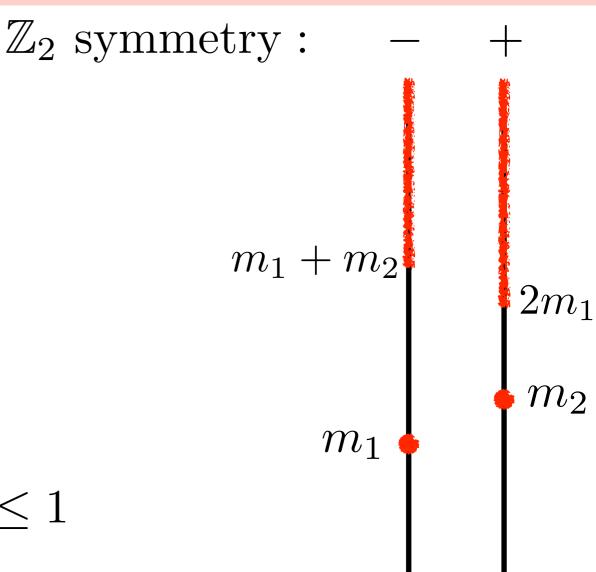
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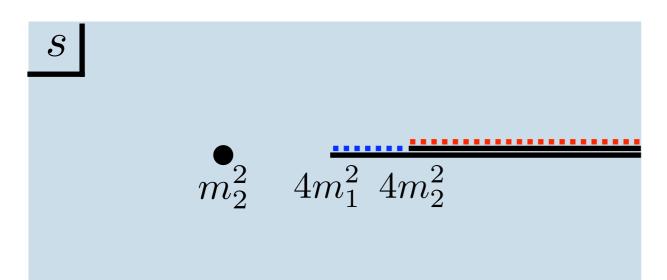
 $s > 4m_2^2$

$$4m_1^2 \le s \le 4m_2^2$$



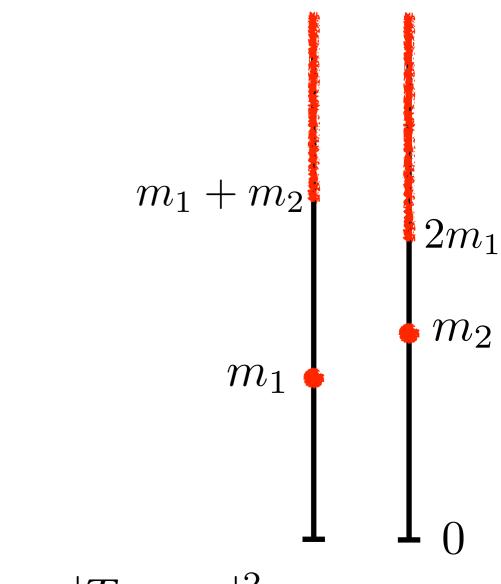
Extended Unitarity:
$$2 \text{Im} T_{22 \to 22} = \frac{|T_{22 \to 11}|^2}{2\sqrt{s(s-4m_1^2)}}$$

Analyticity:



Unitarity:
$$|S_{22\to 11}|^2 + |S_{22\to 22}|^2 \le 1$$

 $s > 4m_2^2$



 \mathbb{Z}_2 symmetry:

Extended Unitarity:
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3-state Potts model saturates the bound for $m_2=m_1$ and $\frac{g_{222}}{g_{112}}=-1$

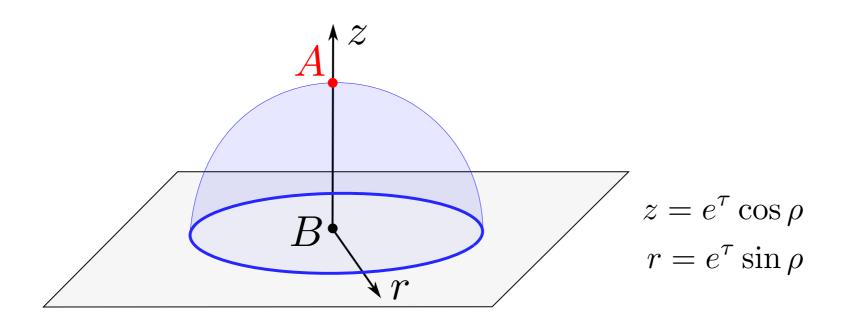
Open questions

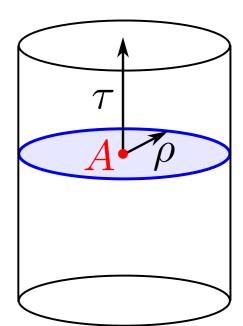
Future work

- Anomalous thresholds (Landau diagrams)
- Particles with spin (internal and external)
- Particles with flavour (global symmetries)
- Use analyticity beyond the physical sheet
- Connect with conformal bootstrap for D>2
- Other interesting questions? Maximize particle production? Resonances?
- Can we input UV data about the QFT? Hard scattering?
 Form factors?

Thank you!

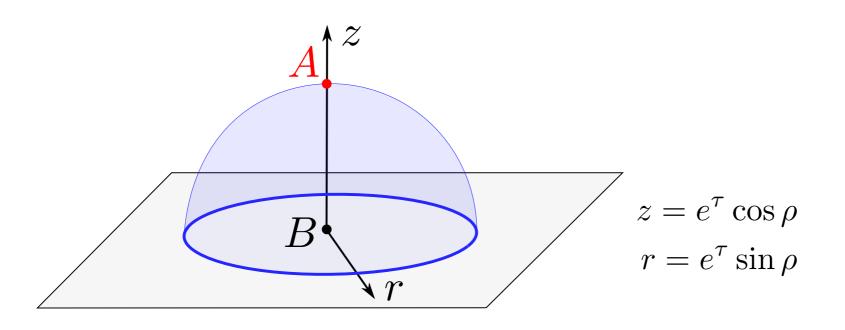
S-matrix from the Conformal Bootstrap

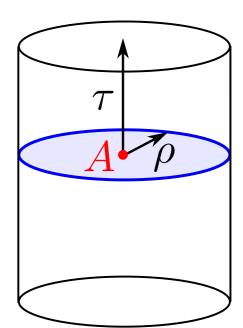




Correlation functions of boundary operators

$$\langle \mathcal{O}(x) \dots \rangle = \lim_{z \to 0} z^{-\Delta} \dots \langle \phi(z,x) \dots \rangle$$
 bulk operator

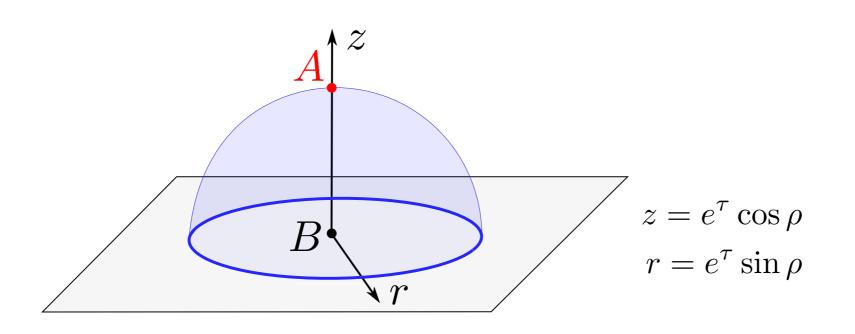


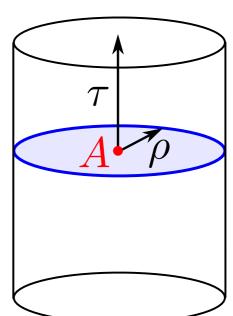


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Isometry group of AdS = SO(d+1,1) = Conformal group

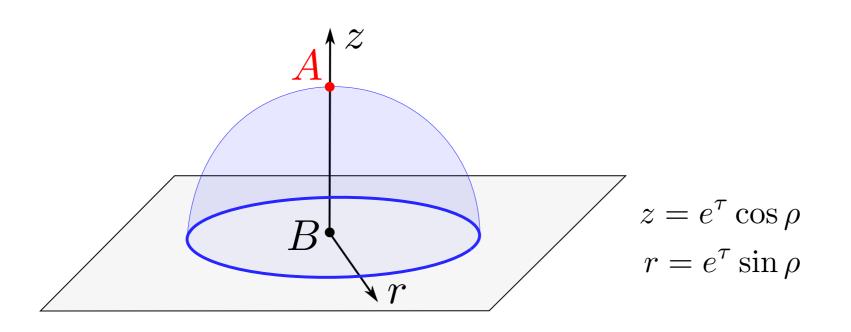


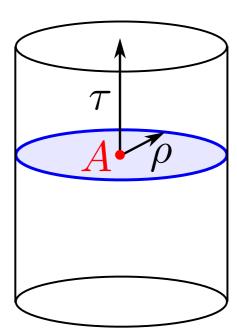


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Isometry group of AdS = SO(d+I,I) = Conformal group Convergent OPE for boundary operators





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Isometry group of AdS = SO(d+I,I) = Conformal group Convergent OPE for boundary operators

 \Rightarrow Use conformal bootstrap to study $\langle \mathcal{O}_1(x_1) \dots \mathcal{O}_n(x_n) \rangle$

AdS radius $R o \infty$

AdS radius $R \to \infty$

Mass spectrum:

$$\Delta_i \sim m_i R$$

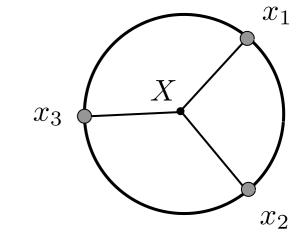
$$\frac{m_i}{m_1} = \lim_{\Delta_i \to \infty} \frac{\Delta_i}{\Delta_1}$$

AdS radius $R \to \infty$

Mass spectrum:

$$\Delta_i \sim m_i R$$

$$\frac{m_i}{m_1} = \lim_{\Delta_i \to \infty} \frac{\Delta_i}{\Delta_1}$$



Cubic couplings:

OPE coefficient

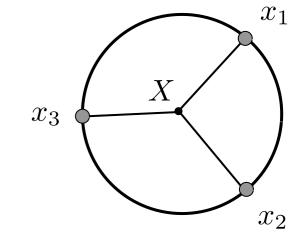
$$g_{123} = \lim_{\Delta_i \to \infty} \lambda_{123}^{4} \times \frac{2(\Delta_1)^{\frac{d-5}{2}}}{\pi^{\frac{d}{2}} \Gamma(\frac{1}{2} \sum_{i=1}^{3} \Delta_i - \frac{d}{2})} \prod_{i=1}^{3} \frac{\Gamma(\Delta_i)}{\Gamma(\frac{1}{2} \sum_{i=1}^{3} \Delta_i - \Delta_i) \sqrt{C_{\Delta_i}}}$$

AdS radius $R \to \infty$

Mass spectrum:

$$\Delta_i \sim m_i R$$

$$\frac{m_i}{m_1} = \lim_{\Delta_i \to \infty} \frac{\Delta_i}{\Delta_1}$$



Cubic couplings:

$$g_{123} = \lim_{\Delta_i \to \infty} \lambda_{123}^{\bullet} \times \frac{2(\Delta_1)^{\frac{d-5}{2}}}{\pi^{\frac{d}{2}} \Gamma(\frac{1}{2} \sum_{i=1}^{3} \Delta_i - \frac{d}{2})} \prod_{i=1}^{3} \frac{\Gamma(\Delta_i)}{\Gamma(\frac{1}{2} \sum_{i=1}^{3} \Delta_i - \Delta_i) \sqrt{C_{\Delta_i}}}.$$

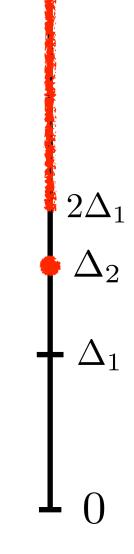
Scattering amplitudes:

Mellin amplitude

$$(m_1)^a T(k_i) = \lim_{\Delta_i \to \infty} \frac{(\Delta_1)^a}{\mathcal{N}} M \left(\gamma_{ij} = \frac{\Delta_i \Delta_j}{\Delta_1 + \dots + \Delta_n} \left(1 + \frac{k_i \cdot k_j}{m_i m_j} \right) \right)$$

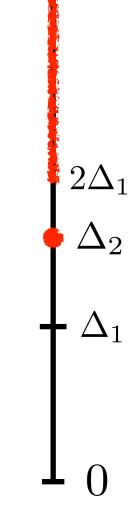
$$a = n(d-1)/2 - d - 1 \qquad \mathcal{N} = \frac{1}{2}\pi^{\frac{d}{2}}\Gamma\left(\frac{\sum \Delta_i - d}{2}\right) \prod_{i=1}^n \frac{\sqrt{\mathcal{C}_{\Delta_i}}}{\Gamma(\Delta_i)}, \qquad \mathcal{C}_{\Delta} \equiv \frac{\Gamma(\Delta)}{2\pi^{\frac{d}{2}}\Gamma\left(\Delta - \frac{d}{2} + 1\right)}.$$

$$\mathcal{O}_1 \times \mathcal{O}_1 = 1 + \lambda_{112}\mathcal{O}_2 + \dots$$
 (operators with $\Delta > 2\Delta_1$)...



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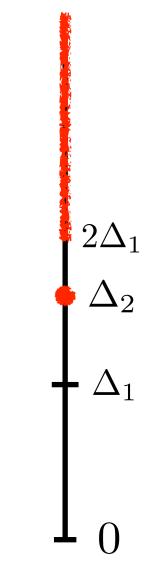
$$\langle \mathcal{O}_1(0)\mathcal{O}_1(z)\mathcal{O}_1(1)\mathcal{O}_1(\infty)\rangle = \frac{1}{z^{2\Delta_1}} \sum_k \lambda_{11k}^2 G_{\Delta_k}(z)$$



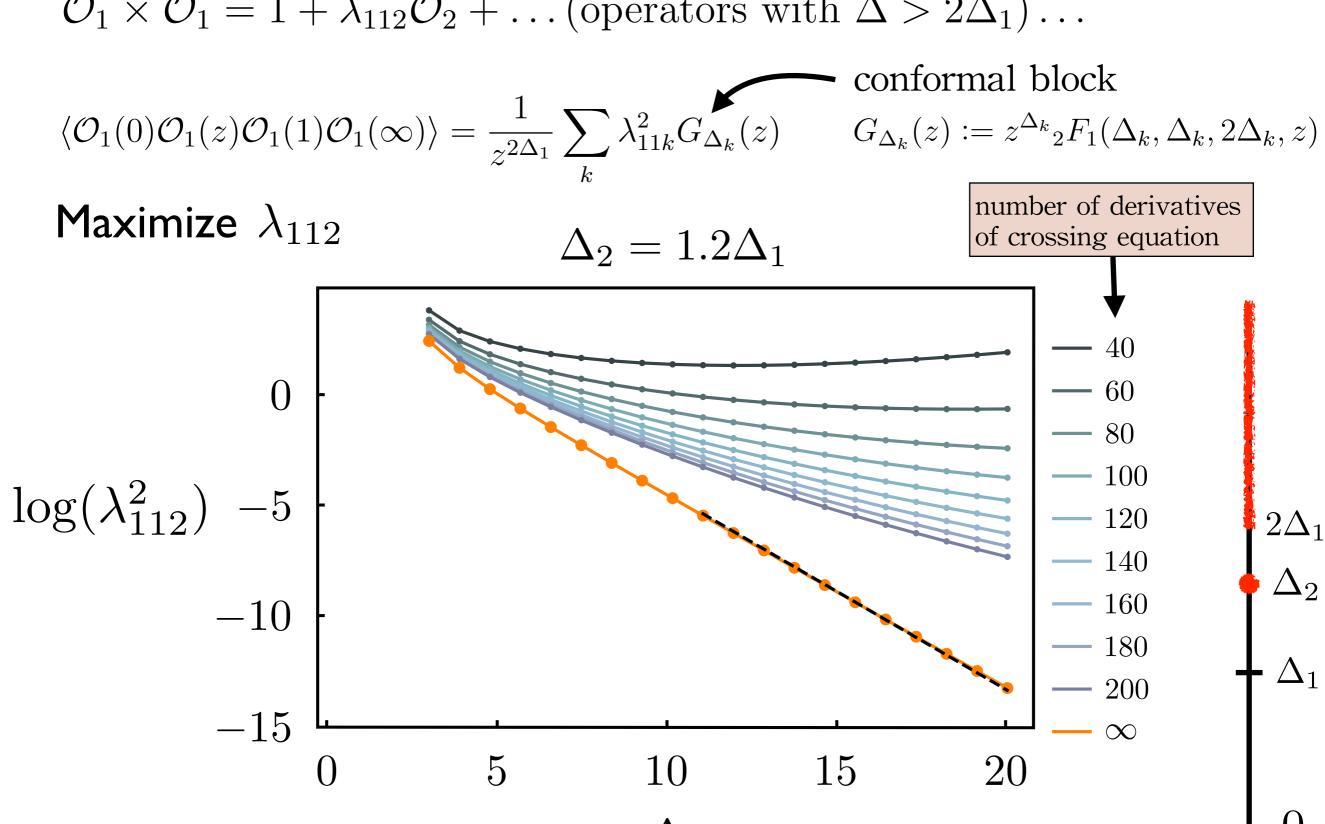
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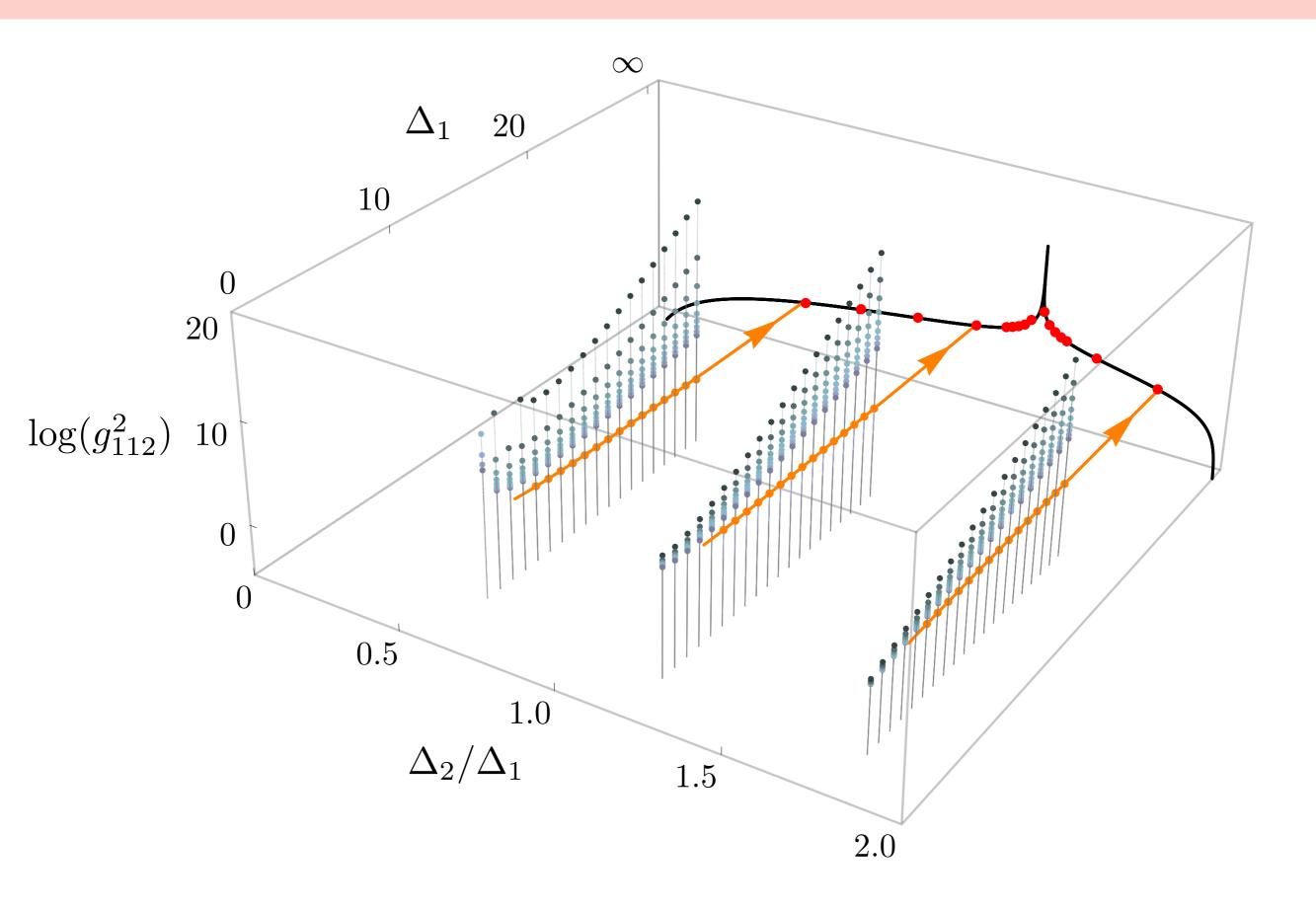
$$G_{\Delta_k}(z) := z^{\Delta_k} {}_2F_1(\Delta_k, \Delta_k, 2\Delta_k, z)$$



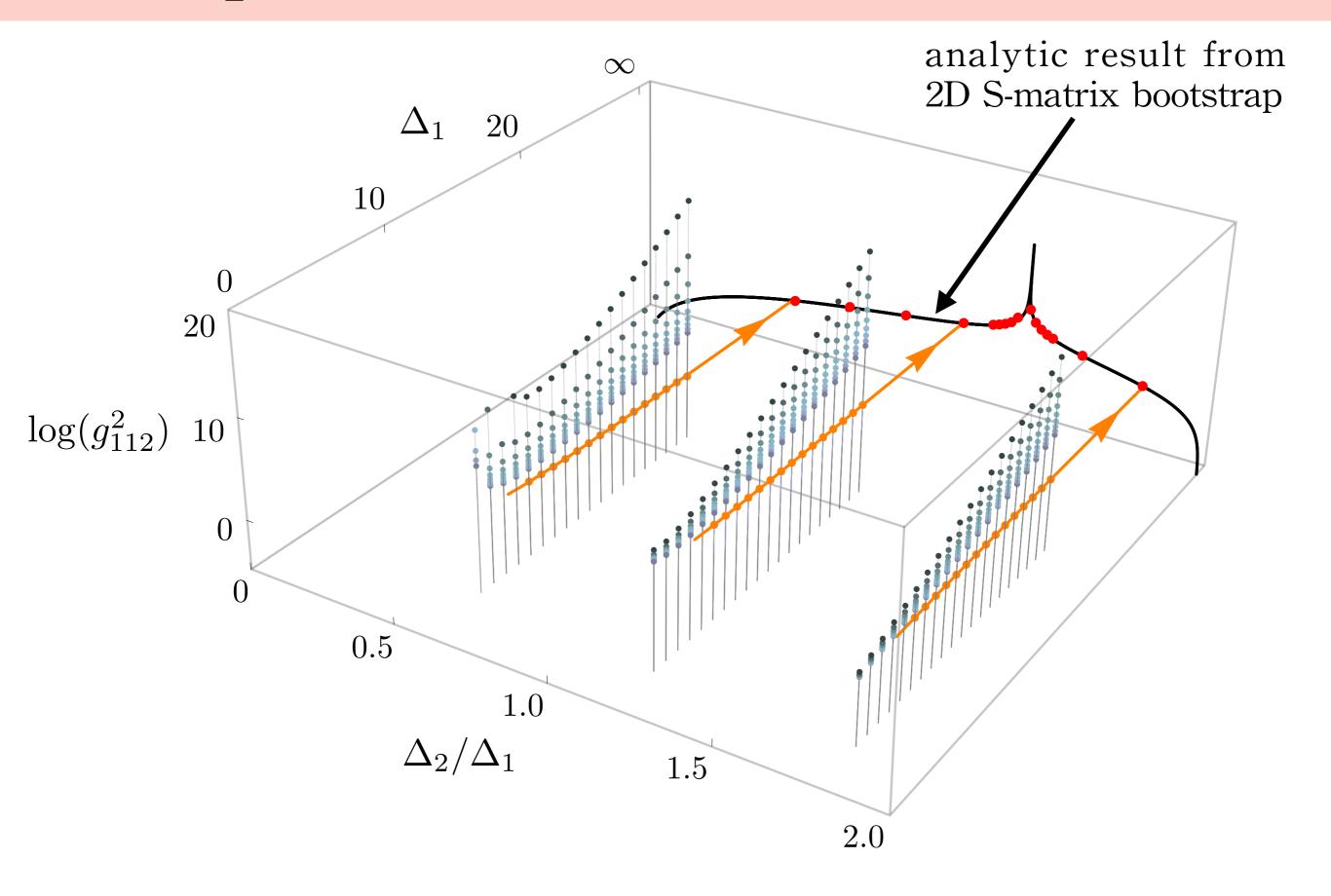
$$\mathcal{O}_1 \times \mathcal{O}_1 = 1 + \lambda_{112}\mathcal{O}_2 + \dots$$
 (operators with $\Delta > 2\Delta_1$)...



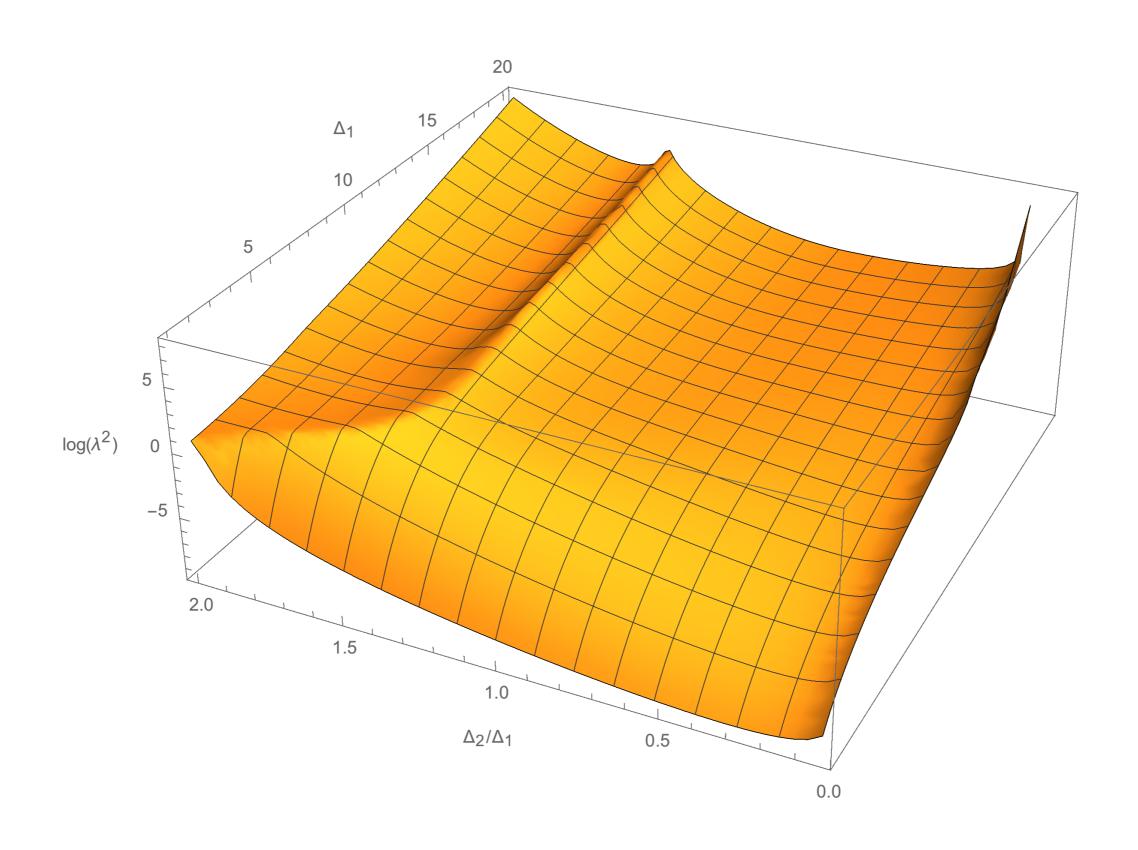
Extrapolation²



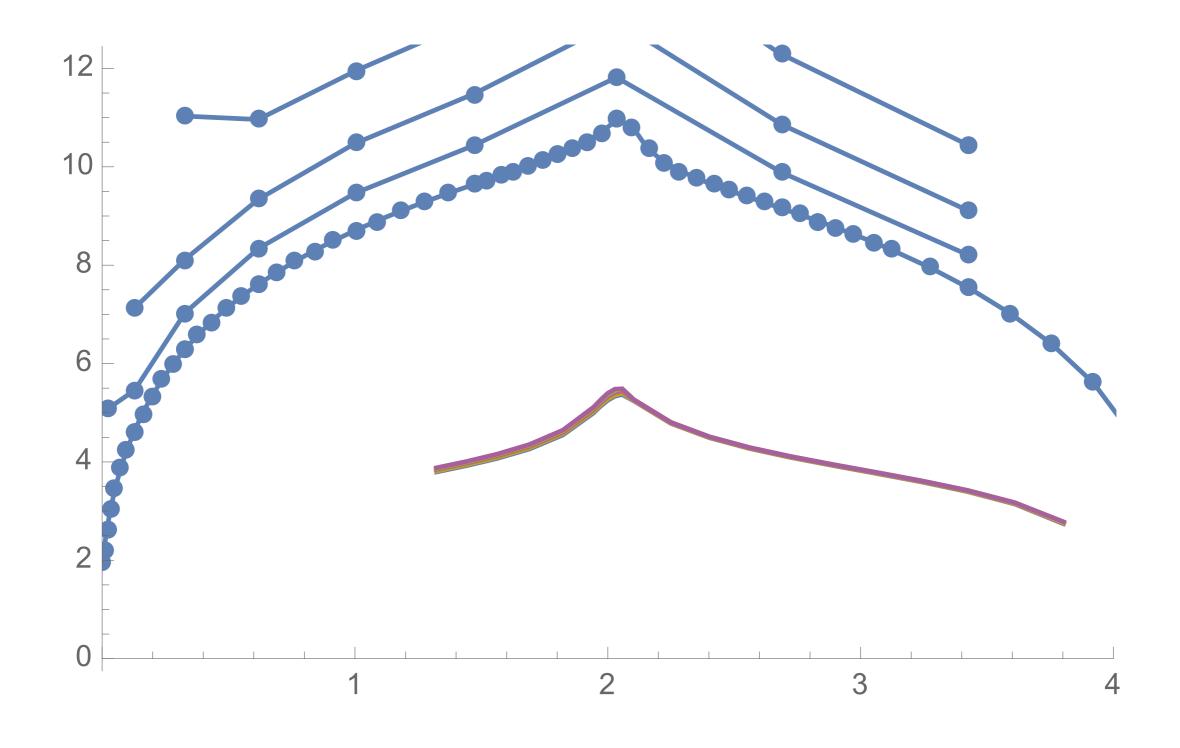
Extrapolation²



2D Conformal bootstrap - preliminary



2D Conformal bootstrap - preliminary



RG flows from QFT in AdS

