# Large scale separation from mass-split models

## Anna Hasenfratz



Workshop on Bound states in strongly coupled systems Galileo Institute, Florence March 14 2018 Non-perturbative investigations of Composite Higgs BSM systems exhibiting

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This is Part 1: Set the stage - (theoretical) introduction Next talk Part 2: O. Witzel - much more details

Mostly based on R. Brower, A.H., Claudi Rebbi, E. Weinberg, Oliver Witzel, Phys.Rev. D93 (2016) 075028 A.H., Claudi Rebbi, Oliver Witzel, Phys.Let. B773 (2017) 86

## Beyond QCD :

#### it's a wild world out there ...

At fixed  $N_c$ :

- small N<sub>f</sub> : chirally broken, QCD-like
- $N_f^* < N_f < N_f^{(|F|)}$  : conformal
- $N_f^{(|F|)} < N_f$ : IR free



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Start with Higgsless, massless SM  $\longrightarrow$  Full SM

 $\mathcal{L}_{SM0}$  $\rightarrow \mathcal{L}_{SM}$ 

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$$\mathcal{L}_{SD} + \mathcal{L}_{SM0} + \mathcal{L}_{int} \rightarrow \mathcal{L}_{SM} + \dots$$

$$f \\ Full SM + additional \\ states from \\ strong dynamics \mathcal{L}_{SD}$$

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The construction ideally will

- predict the 125GeV Higgs
- give mass to the SM gauge fields
- give mass to the SM fermions : (4-fermion interaction or partial compositness?)
- give mass to  $\mathcal{L}_{SD}$  fermions:  $\mathcal{L}_{UV}$  sector

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$$\begin{array}{cccc} \mathcal{L}_{UV} & \rightarrow & \mathcal{L}_{SD} + \mathcal{L}_{SM0} + \mathcal{L}_{int} & \rightarrow & \mathcal{L}_{SM} + \dots \\ & \uparrow & & \uparrow \\ \text{This could be a UV} \\ \text{complete theory} & & & \text{states from} \\ \text{strong dynamics } \mathcal{L}_{SD} \end{array}$$

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 $\mathcal{L}_{SD}$  :SU(N<sub>c</sub>) gauge, N<sub>f</sub> fermions, chirally broken, coupled to the SM

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- Higgs sector : what keeps the Higgs light ?
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No experimental sign of either scenario (yet):  $\rightarrow$  plausible BSM scenarios require (very) large scale separation between  $\mathcal{L}_{SD}$  and  $\mathcal{L}_{SM}$ 

$$\mathcal{L}_{SD} + \mathcal{L}_{SM0} + \mathcal{L}_{int} \rightarrow \mathcal{L}_{SM} + \dots$$

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#### "Foolproof" realization:

- Take  $N_{\mbox{\scriptsize f}}$  above the conformal window
  - Split the masses:  $N_f = N_{\ell} + N_h$

 $N_h$  flavors are massive,  $m_\ell \ll m_h \ll \Lambda_{cut-off} \rightarrow decouple in the IR$ 

 $N_{\ell}$  (= 2 - 4) flavors are massless,  $m_{\ell} = 0 \rightarrow$  spont. chirally broken

#### A mass-split model

- shows conformal properties in the UV
- chirally broken in the IR
- large scale separation controlled by m<sub>h</sub>

### Phase diagram of mass-split model

Mass-split model,  $N_{\ell}$  (=4) +  $N_{h}$  (=8) ;  $m_{\ell}$  = 0:

N<sub>ℓ</sub> flavors - chirally broken  $m_{h}$  $N_{\ell} + N_h$  flavors - Conformal  $\beta \propto 1/g^2$ 

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## Wilson RG for mass-split models

Use Wilson RG description in conformal systems:

- start with bare parameters at the UV scale
- run RG from  $\Lambda_{\text{cut-off}}$  to low energy  $\mu$
- continuum (infinite cut-off) system: tune bare couplings to criticality while keeping µ fixed



- RG flow runs toward the IRFP, lingers around, then flows along RT
- IR physics is along the RT



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## Running/walking coupling

RG flows predict the running coupling:



3 regions:

• UV :

from cut-off to  $g \sim g^*$ 

- walking: m<sub>h</sub> small, g~g\*
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#### Example: $4\ell + 8h$



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The only free parameter is  $m_h$ :

- sets the lattice scale (like g<sup>2</sup> in QCD)
- tune it with go to control scale separation/walking

Hyperscaling around IRFP restricts the spectrum

In conformal systems Wilson RG predicts hyperscaling in the basin of attraction of the fixed point

If the scale changes as  $\mu \rightarrow \mu' = \mu/b, b > 1$ the couplings run as

> $\hat{m}(\mu) \rightarrow \hat{m}(\mu') = b^{y_m} \hat{m}(\mu)$  (increases)  $g \rightarrow g^*$

and any 2-point correlation function scales as

 $C_H(t;g_i,\hat{m}_i,\mu) \rightarrow b^{-2y_H}C_H(t/b;g^*,b^{y_m}\hat{m}_h,b^{y_m}\hat{m}_\ell,\mu)$ 

$$\equiv b^{-2y_H} C_H(t/b;g^*,b^{y_m}\hat{m}_h,\hat{m}_\ell/\hat{m}_h,\mu)$$

since

$$C_H(t) \propto e^{-M_H t} \longrightarrow aM_H \propto (\hat{m}_h)^{1/y_m} F_H(m_\ell/m_h)$$

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$$aM_H \propto (\hat{m}_h)^{1/y_m} F_H(m_\ell/m_h)$$

Ratios are universal functions of  $m_{\ell}/m_h$ 

$$M_{H_1} / M_{H_2} = \Phi_H (m_\ell / m_h),$$
  
 $M_{H_1} / F_\pi = \tilde{\Phi}_H (m_\ell / m_h)$ 

In terms of  $F_{\pi}$  the spectrum is predictable - no free parameters

- in the  $m_{\ell}=0$  chiral limit the spectrum is independent of  $m_h$
- true for light-light, heavy-light and heavy-heavy spectrum

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# Brief summary of mass-split systems:

#### • predictive like QCD:

- $\bullet m_h$  replaces  $g^2$
- conformal FP replaces the Gaussian FP in the UV
- IR properties are not QCD like:
  - the bound state masses in physical units (like  $F_{\pi}$ ) are independent of  $m_h$  when  $m_\ell$ =0, even for heavy flavors
  - the anomalous dimensions in the UV are controlled by the conformal FP
  - walking can be tuned arbitrarily

## Sample numerical results for $4\ell + 8h$

## Light-light spectrum $(4\ell+8h)$

in terms of  $m_{\ell}/m_h$  at  $m_h = 0.05, 0.06, 0.08, 0.10$ 



- approaches N<sub>f</sub>=12 as  $m_{\ell}/m_{h} \rightarrow 1$  :  $M_{H}/F_{\pi}$  ratios are finite
- universal in  $m_{\ell}/m_h$  at various  $m_h$

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## Hyperscaling

Ratios depend on  $m_{\ell}/m_h$  but not on  $m_{\ell}$ ,  $m_h$  or  $g^2$ 



four different  $m_h = 0.05 - 0.10$ on each panel

similar for other states

## Compare to QCD

QCD:  $m_{\ell} \rightarrow 0$ ,  $m_{h} \rightarrow \infty$ ; not in the basin of attraction of the IRFP (no hyperscaling) —



Connected spectrum is mainly valence fermions

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