Gamma-ray Flares of FSRQs: 
HE component;

The connection of flaring states with the accretion disk luminosity

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\( \gamma \gamma \) absorption from BLR via \( \gamma \gamma \rightarrow e^+e^- \) interaction

Liu & Bai 2006
Liu, Bai, Ma 2008
\( \gamma \gamma \) absorption from BLR

Liu & Bai 2006; Liu, Bai, Ma 2008

\[ L_{\text{BLR}} = 2.6 \times 10^{44} \text{ erg/s} \] (this is the case of 3C 279, but for 3C 454.3 the BLR is >10 times more luminous)

\( \tau_{\gamma \gamma} \) scales as \( L_{\text{BLR}}^{1/2} \)
KN suppression
(External Compton on BLR photons)


Klein-Nishina suppression for a dissipation region at:
solid curves: $0$, $R_{in}$, $R_{out}$ (from Bottom Up)
solid dashed curve: at the center of the shell
dashed curves: $0.95 \times R_{out}$, $1.25 \times R_{out}$, $1.5 \times R_{out}$, $2 \times R_{out}$
SEARCH within the FERMI-LAT FSRQs sample
We started to search for relevant signal at $E > 20/(1+z)$ GeV in the FERMI-LAT archive from FSRQs and on incoming gamma-ray data (and triggering ToO observations to Swift).

High energy (HE) activity period is defined as the period of time in which the HE photon rate is $> 3 \times$ mean HE rate.
Search within the FERMI-LAT FSRQs sample (I)
Search within the FERMI-LAT FSRQs sample (II)
Search within the FERMI-LAT FSRQs sample (III)
Search within the FERMI-LAT FSRQs sample (IV)

Ghisellini & Tavecchio 2009
SEDs and modeling (i)

**Distances and Properties**

- **dist = 3 pc**
  - B = 25 mG
  - $L_{\text{disk}} = 5.3 \times 10^{45}$ erg/s

- **dist = 0.8 pc**
  - B = 315 mG
  - $L_{\text{disk}} = 24 \times 10^{45}$ erg/s

- **dist = 2.5 pc**
  - B = 35 mG
  - $L_{\text{disk}} = 3.7 \times 10^{45}$ erg/s

- **dist = 2 pc**
  - B = 110 mG
  - $L_{\text{disk}} = 15 \times 10^{45}$ erg/s

**IGM Absorption in UV**
SEDs and modeling (ii)

- **B2 1520+031**
  - $z = 1.49$
  - $d = 1.3$ pc
  - $B = 63$ mG
  - $L_{\text{disk}} = 8 \times 10^{45}$ erg/s

- **B2 1633+38 (4C+83.41)**
  - $z = 1.81$
  - $d = 0.5$ pc
  - $B = 250$ mG
  - $L_{\text{disk}} = 50 \times 10^{45}$ erg/s

- **B2 1846+32A**
  - $z = 0.80$
  - $d = 0.2$ pc
  - $B = 620$ mG
  - $L_{\text{disk}} = 3.4 \times 10^{45}$ erg/s

- **PMN J2345−1555**
  - $z = 0.62$
  - $d = 0.7$ pc
  - $B = 210$ mG
  - $L_{\text{disk}} = 1.5 \times 10^{45}$ erg/s
SEDs and modeling (iii)

For 3C 454.3 (z=0.86),
- Distance: 0.5 pc
- Magnetic Field: B=230 mG
- Disk Luminosity: $L_{\text{disk}} = 33 \times 10^{45}$ erg/s

For CTA 102 (2230+114) (z=1.037),
- Distance: 0.3 pc
- Magnetic Field: B=140 mG
- Disk Luminosity: $L_{\text{disk}} = 41 \times 10^{45}$ erg/s
GB6 J1239+0443 \((z=1.76)\) Multiepoch SED (I)

**AGILE/GRID** and simultaneous data in red

**FERMI-LAT** data (4-day integration around the flare) and simultaneous data in black

*Fermi-LAT* data in green (30-day integration around the flare)

*Fermi-LAT* data in cyan (2FGL catalog)

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**Dissipation region at 0.2 pc** from the SMBH

(Just outside the BLR)

\[ R_{\text{blob}} = 6.7 \times 10^{16} \text{ cm} \]

\[ B = 0.6 \text{ Gauss} \]

**Dissipation region at 7 pc** from the SMBH

\[ R_{\text{blob}} = 2 \times 10^{18} \text{ cm} \]

\[ B = 1 \times 10^{-2} \text{ Gauss} \]

This model gives a satisfactory gamma-ray spectral shape, but the expected variability is \(~10^2\) days.
GB6 J1239+0443 (z=1.76) Multiepoch SED (II)

Model is for a dissipation region at 5 pc from the central BH, a blob radius of $1 \times 10^{17}$ cm, $B = 7 \times 10^{-2}$ Gauss

$R_{\text{blob}} = 0.0067 R_{\text{diss}}$ in agreement within a factor 2 with Bromberg and Levinson 2009 ($R_{\text{blob}} = 10^{-2.5} R_{\text{diss}}$)

inverting $R_{\text{diss}} = 2.5 L_{\text{jet,46}} (R_{\text{BLR}}/0.1 \text{ pc})^{-1}$ and using $R_{\text{diss}} = 5$ pc, we obtain $L_{\text{jet}} = 3.5 \times 10^{46}$ erg/s.

We need to assume that the p/e number ratio is $\sim 0.1$ to accomplish such a luminosity.
From the broad Mg II line: 
$L_{\text{Mg II}}=5.4 \times 10^{43} \text{ erg/s}$ (Stickel 89) 
we derived the BLR luminosity (Celotti 1997) 
and in turn the disk luminosity 
(we assumed a BLR/disk luminosity ratio 0.1) 
$L_{\text{disk}} = 1 \times 10^{46} \text{ erg/s}$

$B = 6 \times 10^{-3} \text{ G}$

$\Gamma = 20$

Dist = 7 pc

R = 1-1.2 pc

Variability time Scale 30 d, 
comparable with long term modulation of the light curve, but not with the daily variability.
PKS 1441+25

$z=0.94$

$L_{\text{disk}} \sim 2 \times 10^{45} \text{ erg/s}$

$R_{\text{in, BLR}} \sim 0.05 \text{ pc}$

$\text{dist} \sim 0.1 \text{ pc}$

A pessimistic evaluation of attenuation ($\gamma \gamma \text{ abs } + \text{ KN}$) at 50 GeV (sat frame) is $< 3$

So the emission region must be at the edge or outside the BLR.
PKS 2023-07 Dec. 2016 flare
(Piano 2018)
Fast HE flares

From the 4 brightest HE flares we searched for fast variability at HE (E> 10 GeV).

For all these 4 sources we found short periods (period A) lasting from 1.5 hours to less than 6 hours of very bright HE emission and hard spectra.

NB: in the following, the gamma-ray photon index of periods A ($\Gamma_{\text{ph}}$) are evaluated in the energy range 0.2-10 GeV (they are not biased by the selection criteria, i.e. the search for bright emission at HE, E>10 GeV)
Fast HE flares and spectral evolution (i)

\[
\Delta t/(1+z) = 0.11 \text{ d} \\
\Gamma_{ph} = 1.99 \pm 0.31
\]

\[
\Delta t/(1+z) = 0.12 \text{ d} \\
\Gamma_{ph} = 1.51 \pm 0.34
\]

\[
\Delta t/(1+z) = 1.4, 2.8 \text{ d}
\]

\[
\Delta t/(1+z) = 2.8, 2.8 \text{ d}
\]
Fast HE flares and spectral evolution (ii)

CTA 102

PERIOD A

\[ \Delta t/(1+z) = 0.076 \text{ d} \]
\[ \Gamma_{ph} = 1.73 \pm 0.14 \]

3C 454.3

PERIOD A

\[ \Delta t/(1+z) = 0.30 \text{ d} \]
\[ \Gamma_{ph} = 1.77 \pm 0.17 \]

CTA 102

PERIOD B, C, D

\[ \Delta t/(1+z) = 1.5, 2.0^*, 2.0 \text{ d} \]

3C 454.3

PERIOD B, C, D

\[ \Delta t/(1+z) = 1.6, 1.6, 1.6 \text{ d} \]

*period C starts 5 d after period B due to a gap in the telemetry.
Fast HE flares and spectral evolution (ii.j)

CTA 102 and 3C 454.3 gamma-ray spectra of period B are consistent with the slow cooling scenario, with:

\[ \text{low energy } \Gamma_{ph} \text{ consistent with } \Gamma_{ph} \text{ of period A,} \]

and

\[ \Delta \Gamma_{ph} = 0.75 \pm 0.32 \text{ (3C 454.3)} \]
\[ \Delta \Gamma_{ph} = 0.72 \pm 0.35 \text{ (CTA 102)} \]

In the dusty torus photon field, the expected cooling time is \( \sim 1 \text{ hour} \) for electrons with \( \gamma = 30000 \)
\( \sim 30 \text{ GeV EC photons} \)
The distant scenario

- The bright HE emission witnesses against BLR absorption and Klein-Nishina suppression (for EC on BLR photons)
- The leptonic SED modeling is only consistent with a dissipation region at parsec scale
- The spectral evolution from an hard spectrum is consistent with the slow cooling scenario (chromatic cooling) on Torus seed photons (while the cooling on BLR photons is in Klein-Nishina regime and it is expected to be achromatic).
- **But the CTA 102 light curve** shows a variability pattern which is inconsistent with slow cooling (what is the lower activity period in between two higher activity periods, with a duration of 0.5 days?).
what is the engine?

Reasonable engines acting at large distance from the SMBH are:

- Magnetic reconnection (Giannios 2013)
- Turbulence in the jet (Narayan & Piran 2012, Marscher 2014)
Variability time scale from the SED modeling is \(~30\) d, comparable with long term modulation of the light curve, but we observe also sub-daily variability.

Recent scenario for magnetic reconnections proposed for strongly magnetized jets (Giannios 2013) includes an envelope emission (lasting \(~1\) day) powered by plasmoids, together with fast flares (lasting \(~10\) min) generated by grown “monster plasmoids”.

In low magnetized plasma (such as at several parsec), reconnection time scales are longer and longer flares (days to weeks) could arise (Giannios 2013).

“Monster plasmoids” contain energetic particles freshly injected by the reconnection event (Uzdensky et al. 2010).

Figure 2. A sketch of the envelope-flare structure of the emission from a reconnection layer. The envelope duration corresponds to that of the reconnection event: \(t_{\text{env}} = l' / \Gamma_j c\). Monster plasmoids power fast flares which show exponential rise and last for \(t_{\text{flare}} = 0.1 l' / \delta_p c\). For an envelope of \(~1\) d blazar flaring, the model predicts that monster plasmoids result in \(~10\)-min flares.
Turbulence in the jet

electron acceleration is caused by standing conical recollimation shocks.

Flux and polarization variability originates from turbulence in the flow, approximated as cylindrical cells

(Marscher 2014)
HOW MANY SOURCES?
HOW MANY FLARES?

Work in progress
PowerLaw photon-index distribution for HE flares (I)
(fitting below $E_{THR}$: $200 \text{ MeV} - E_{THR}$)
PowerLaw photon-index distribution for HE flares (II) (fitting below $E_{\text{THR}}$: 200 MeV – $E_{\text{THR}}$)

sources with PowerLaw spectrum in the 2nd FERMI-LAT catalog:
PowerLaw photon-index distribution for HE flares (III)
(fitting below $E_{\text{THR}}$: 200 MeV – $E_{\text{THR}}$)

sources with LogParabolic spectrum in the 2nd FERMI-LAT catalog:
How many flares?
The comparison with the full sample of gamma-ray flares of FSRQs

- With a clustering method applied to gamma-ray collected within a suitable region around the source, we searched for gamma-ray flares of FSRQs in the energy range $0.3\text{-}300\text{ GeV}$.

- **Data set**: $\{X(i)\}$ (gamma-ray events collected within an extraction region) where $X(i)$ is the cumulative exposure (from the start of obs) of the collected event $i$.

- **Clustering law**:

  $$\begin{cases} 
  X_{(i+k)} - X_{(i)} < k*\Delta_{\text{thr}} & (K < N_{\text{tol}}) \\
  l \in [i,i+k]
  \end{cases}$$

- Chance cluster probability is evaluated with a scan statistic related method (maximum score scan statistic, Glaz 2006, conf. level set to $1*10^{-3}$).
How many flares?
The comparison with the full sample of gamma-ray flares of FSRQs

- Solar flares are rejected, asking the source to be at $>15$ deg from the sun during flare.
- Flares from closeby sources are rejected studying the angular distribution of events during flare.
- Peak fluxes evaluated with photometry are compared with the full likelihood analysis, and eventually the flare is validated.
3C 454.3 photometric LC (0.3-300 GeV)

Each horizontal segment represents a cluster. The subtended time period is the cluster length (projected into time domain). For each cluster, the Flux level is the mean Flux within the segment. Each cluster has a chance prob. with respect to the parent one of $< 1.3 \times 10^{-3}$ according to max-score scan statistic. Contrary to usual LC, No typical timescales (time binning) spoil the peak flux estimation.
Applying the method to the whole gamma-ray data set of the Crab Nebula we revealed its fastest flare.

Flare was already found, but its fast temporal shape was not recognized.
How many flares?
The comparison with the full sample of gamma-ray flares of FSRQs

Angular distribution of gamma-ray events during flare
How many flares?

Photometry-likelihood comparison

- The likelihood evaluates sources counts, diffuse bkg within the chosen time interval
- The photometry evaluates the investigate source counts+bkg within the extraction region, bkg is estimated from the long integration, bkg source counts are evaluated from the catalog.
- Photometric source counts and likelihood source counts could not coincide (especially at low source flux)
- If the likelihood and photometric based flux for the investigated source do not correspond, the photometric peak flux could be ascribed to a bkg source, and the peak must be rejected.
Opacity computations at $R_{\text{diss}} = R_{\text{BLR}}/2$

- The opacity becomes low as the dissipation region approaches the internal radius of the BLR. Moreover, the opacity computed value strongly depends on the BLR model.
- At distances of the order of $R_{\text{BLR}}/2$, the opacity is high and the computation do not strongly depends on the details of the modeling.

- We use the Boettcher 2016 model, normalized to the observed Broad Line luminosities.
Opacity computations at $R_{\text{diss}} = R_{\text{BLR}}/2$

- Using the Boettcher 2016 model, normalized to the observed Broad Line luminosities, the opacity at $R_{\text{diss}} = R_{\text{BLR}}/2$ is estimated:
  \[
  \tau_{\gamma\gamma} \sim 4.5 \times (L_{\text{disk}}/1.10^{46})^{0.5}
  \]
  assuming $R_{\text{BLR}} = 10^{17} \times (L_{\text{disk}}/1.10^{45})^{0.5}$ cm (Ghisellini 2009)

With a gamma-ray only sample, the opacity argument can be used to discriminate flares at $R_{\text{diss}} > R_{\text{BLR}}/2$
Temporal FWHM distribution of gamma-ray flares

- $E > 0.3$ GeV
- $E > 35$ GeV

- $E > 0.3$ GeV found at peak
- $E > 35$ GeV expected at peak
The jet to disk correlation during flares

- red: flares with significant emission above $35/(1+z)$ GeV (TS>25)
- blue: flares with some emission above $35/(1+z)$ GeV (TS > 9)

corr coeff=0.498, P chance $\sim 10^{-25}$
How many HE flares were found with respect to the expectation?
The photon index correlation with (disk rescaled) flare Luminosity

Flaring Luminosity evaluated with photon index set to 2.2

corr coeff=-0.193, P chance=6*10^{-5}
corr coeff=-0.399, P chance=10^{-16}
Flaring luminosity does not exceed a $1/\text{FWHM}$ behaviour. Doppler factor is not the main actor to explain the temporal FWHM, because flaring luminosity $\propto \delta^6$ (EC scenario).

We also have prolonged emission above 35 GeV but typical cooling is of order of hour timescales. Flare temporal shape convolved with a long living engine?
Relative Flare Luminosity distribution

\[ Y = A e^{-2x} \]

Sensitivity limit modulates the low Luminosity shape
Conclusions (I)

- From Broad band MWL campaigns on several flares, and the relevant HE gamma-ray emission in HE flares, the gamma-ray dissipating region is placed toward the edges, or outside the Broad Line Region to avoid gamma-gamma absorption (mainly for bright disks $\sim 10^{46} \text{ erg/s}$) and KN suppression.

- From the whole gamma-ray flaring sample, $\sim 40\%$ of gamma-ray flares comes from $R_{\text{diss}} > R_{\text{BLR}}/2$

- But at $dt < 1 \text{ d}$, we observe an HE occurrence of $\sim 100\%$.

- But the remaining $60\%$ come from $R_{\text{diss}} < R_{\text{BLR}}/2$?
The flaring luminosity correlates with disk. This requires that the bulk of the HE flares sample is powered or “catalysed” by accretion (by B?) (Narayan 2003, Tcheckhovskoy 2011, Sbarrato 2014, Ghisellini 2014)

Does this fact rules out reconnection and turbulences?

Resonably, Zamaninasab 2014, showed that for a sample of 76 Radio Loud AGN the $B^2$ field at $10^4 r_s$ correlates with $L_{\text{disk}}$.

Also photon index correlates with rescaled flaring luminosity (rescaled with the disk luminosity)