

NUCLEAR PHYSICS AROUND THE UNITARITY LIMIT

U. van Kolck

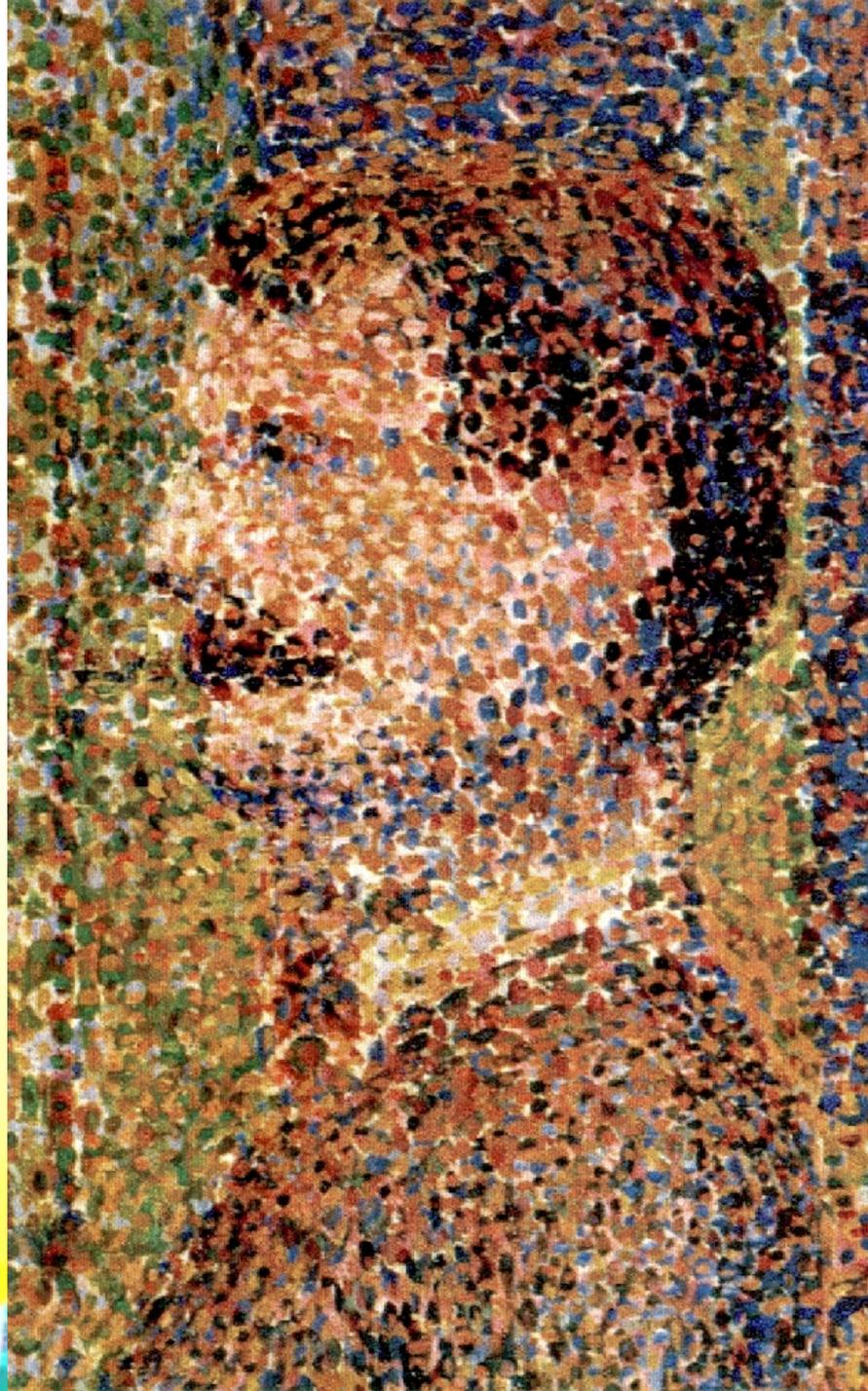
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El Greco,
1586-90

"We must sow the seed, not hoard it."
Santo Domingo (1170-1221)





Seurat,
La Parade
(detail)

Outline

- What is essential?
- Effective Field Theory
- Unitarity: light nuclei
- Unitarity: bosonic clusters
- Unitarity: matter
- Conclusion

with

S. König,
H.W. Griesshammer
& H.-W. Hammer

S. Gandolfi,
J. Carlson,
& S.A. Vitiello

S. Lyu,
& B. Long



What is essential in nuclear physics?

Traditional
approach

- 1) Describe NN precisely to some high energy
- 2) Append 3N forces as needed

→ Lost in details:

e.g., chiral potentials at NLO, N²LO, N³LO, N⁴LO, ...
(apparently no nuclear matter saturation at LO!)

Here

- 1) Describe NN and 3N approximately
- 2) Treat all other interactions in *perturbation theory*
cf. atomic systems in QED

“simplicity emerging
from complexity”

Expansion
around unitarity

König, Grießhammer,
Hammer + v.K. '15 '16
König '16
v.K. '17

A = 2

$$T_2(k \ll R^{-1}) = \frac{4\pi}{m} \left(\underbrace{a_2^{-1}}_{\text{scattering length}} + \underbrace{ik}_{\text{effective range}} - \frac{r_2}{2} k^2 + \frac{P_2}{4} k^4 + \dots \right)^{-1} + (l > 0)$$

unitarity limit

$$a_2^{-1} \approx \sqrt{mB_2} \rightarrow 0$$

$$\frac{r_2}{2} \sim R, \quad \frac{P_2}{4} \sim R^3, \quad \dots \quad \text{typically}$$

nucleons
 $R^{-1} \sim m_\pi$

$$\begin{array}{l}
 |a_{2,I=1,I_3=0} m_\pi|^{-1} \simeq 0.06 \\
 \textcircled{1S_0} \quad |a_{2,I=1,I_3=+1} m_\pi|^{-1} - |a_{2,I=1,I_3=0} m_\pi|^{-1} \simeq 0.12 \\
 |a_{2,I=1,I_3=-1} m_\pi|^{-1} - |a_{2,I=1,I_3=0} m_\pi|^{-1} \simeq 0.02 \\
 \textcircled{3S_1} \quad |a_{2,I=0} m_\pi|^{-1} \simeq 0.26
 \end{array}$$

$$T_2 \left(|a_2^{-1}| \ll k \ll R^{-1} \right) = \frac{4\pi}{m} (ik)^{-1} \left(\textcircled{1} + \mathcal{O} \left(\frac{1}{ka_2}, kR \right) \right)$$

unitarity window

no parameter!

“universality”

A > 2

binding momenta

$$Q_A / m_\pi \sim \left(\frac{2m_N B_A}{m_\pi^2 A} \right)^{1/2} \sim 0.7$$

in unitarity window!

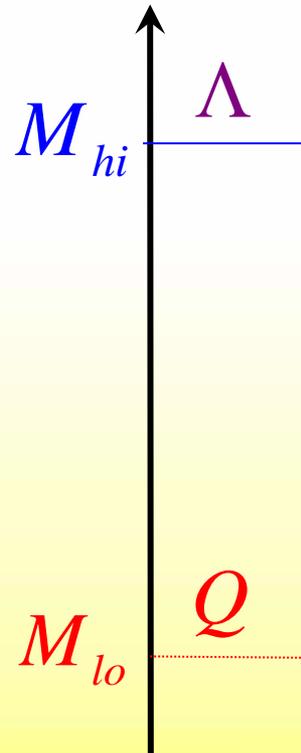


Effective Field Theory

most general action with symmetries



momentum scales



Λ arbitrary UV regulator

model independence

i.e., consistent with QCD!

$$\frac{\partial T}{\partial \Lambda} = 0$$

renormalization-group invariance

"low-energy constants"

$$T(Q \sim M_{lo} \ll M_{hi}) \propto \sum_{\nu=0}^{\infty} \left[\frac{Q}{M_{hi}} \right]^{\nu} F^{(\nu)} \left(\frac{Q}{M_{lo}}, \frac{Q}{\Lambda}; \gamma_i^{(\nu)} \left(\frac{\Lambda}{M_{lo}}, \frac{M_{lo}}{M_{hi}} \right) \right)$$

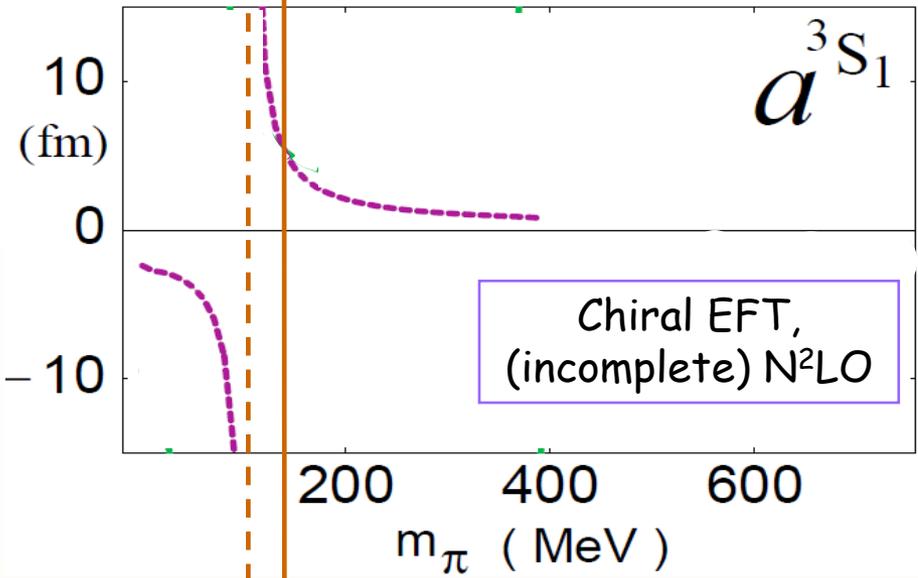
controlled expansion

N^{ν} LO

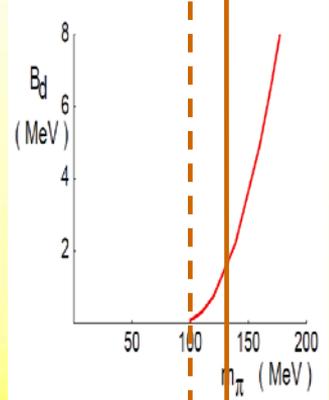
non-analytic functions, from solution of dynamical eq. (e.g. Lippmann-Schwinger)

The quantum generalization of the multipole expansion

unitarity limit



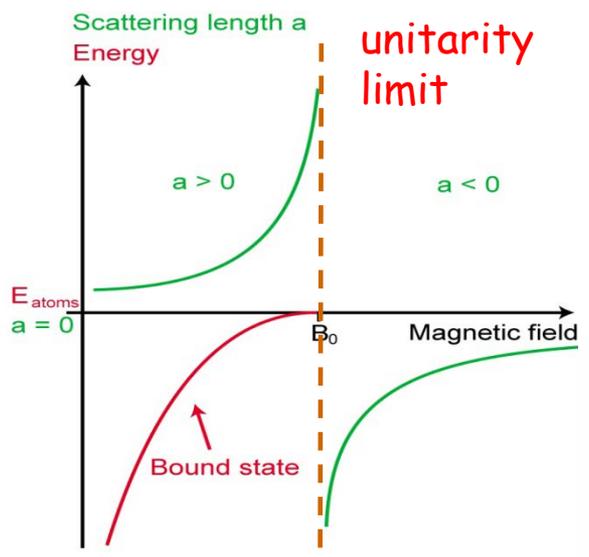
$m_\pi^* (M_{QCD})$ $m_\pi \approx 140 \text{ MeV}$



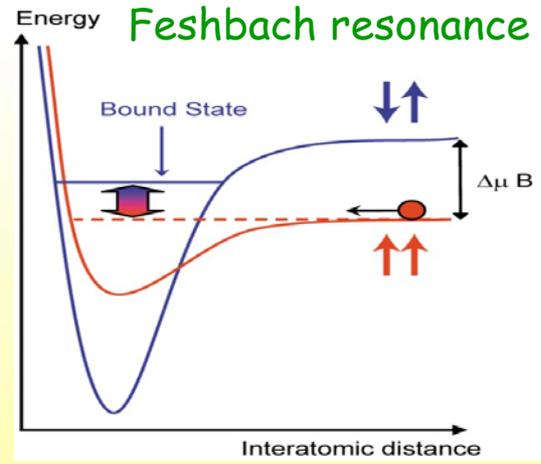
- Beane, Bedaque, Savage + v.K. '02
- Beane + Savage '03'04
- Epelbaum, Glöckle + Meißner '03
- Braaten + Hammer '03

$a_2 \gg m_\pi^{-1} \sim r_2$

[can do nuclear physics at unphysical quark masses as well, but that's another talk]



unitarity limit



Feshbach resonance

or "accidentally", e.g. ⁴He atoms

$a_2 \gg l_{vdW} \sim r_2$

MIT webpage



Pionless EFT

$$Q \sim M_{lo} \ll M_{hi} \sim m_\pi$$

- d.o.f.: nucleons
- symmetries: Lorentz, ~~P~~, ~~T~~, ~~B~~, SU(3)_c, U(1)_{em} (trivial)

$$S_{EFT} = \int \frac{dt}{2m_N} \int d^3r \left\{ \psi^\dagger \left(2im_N \frac{\partial}{\partial t} + \vec{\nabla}^2 \right) \psi \right.$$

projector on isospin /

$$- 4\pi \sum_{I=0,1} C_{0I} \psi^\dagger \psi^\dagger P_I \psi \psi$$

$$- \frac{(4\pi)^2}{3} D_0 \psi^\dagger \psi^\dagger \psi^\dagger \psi \psi \psi$$

more derivatives,
more bodies,
isospin violation,
...

most general action

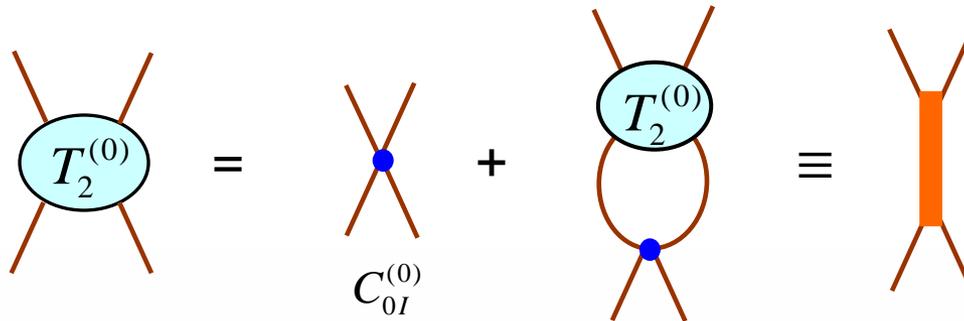
Universality:
first orders
apply also to
neutral atoms

$$m_\pi \rightarrow 1/l_{vdW} \quad \text{where} \quad V(r) = -\frac{l_{vdW}^4}{2mr^6} + \dots$$

Bedaque, Hammer
+ v.K. '99'00
Bedaque, Braaten
+ Hammer '01

A = 2

LO



singular potential!

$$V_2^{(0)} = \frac{4\pi}{m_N} C_{0I}^{(0)} \delta^{(3)}(\vec{r}) \rightarrow \frac{4\pi}{m_N} C_{0I}^{(0)}(\Lambda) \delta_{\Lambda}^{(3)}(\vec{r})$$

~~$$C_{0I}^{(0)}(\Lambda) = C_{0I}^{(0)} \Rightarrow B_d \propto \Lambda^2 / m_N$$~~

renormalization: $C_{0I}^{(0)}(\Lambda) = -\frac{1}{\theta_0 \Lambda} \Leftrightarrow B_d^{(0)} = 0$

regulator-dependent number



$$x \rightarrow \alpha x$$

$$\Lambda \rightarrow \alpha^{-1} \Lambda$$

$$\frac{t}{m} \rightarrow \alpha^2 \frac{t}{m}$$

$$\psi \rightarrow \alpha^{-3/2} \psi$$

$$S_{EFT}^{(0)} \rightarrow S_{EFT}^{(0)}$$

NLO

- scattering lengths
- effective ranges
- Coulomb

in distorted-wave Born approximation

etc.

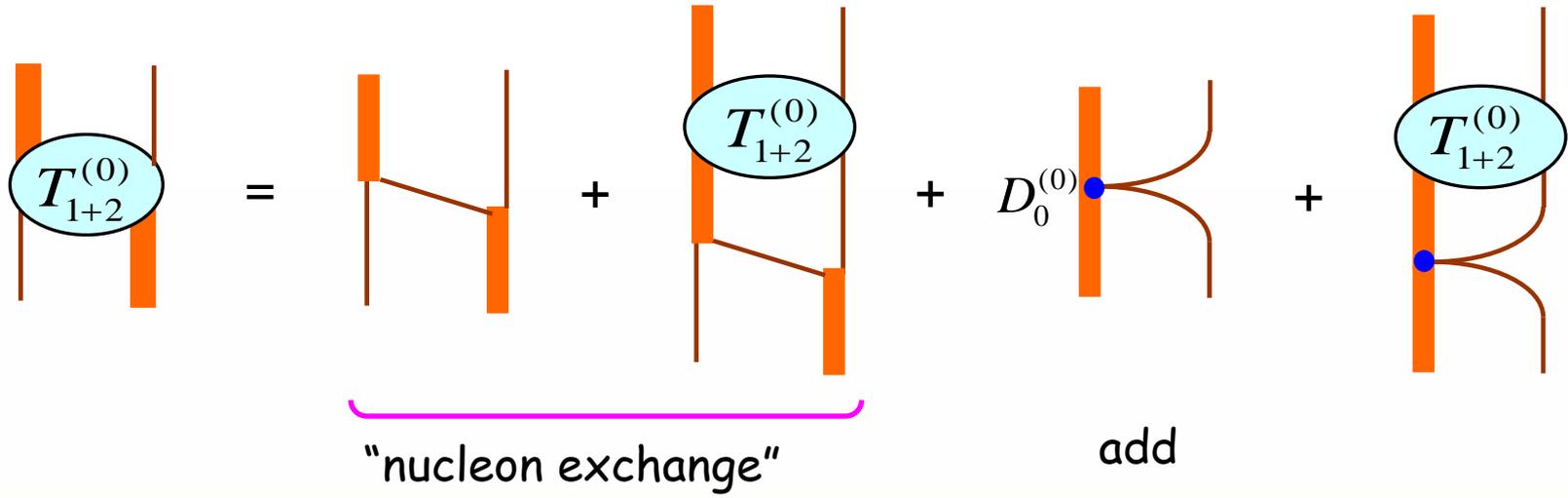
scale invariance
broken explicitly

$$e.g., B_d^{(2)} = 1/(m_N a_{2,I=0}^2)$$



LO

$A = 3$



Thomas '35

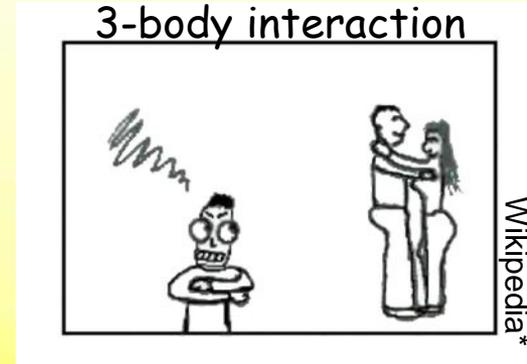
~~$B_t \propto \Lambda^2 / m_N$~~

$$V_3^{(0)} = \frac{(4\pi)^2}{m_N} D_0^{(0)}(\Lambda) \delta_\Lambda^{(3)}(\vec{r}_{12}) \delta_\Lambda^{(3)}(\vec{r}_{23})$$

renormalization:

$$D_0^{(0)}(\Lambda) \propto \frac{1}{\Lambda^4} \frac{\sin(s_0 \ln(\Lambda/\Lambda_*) + \arctan s_0^{-1})}{\sin(s_0 \ln(\Lambda/\Lambda_*) - \arctan s_0^{-1})}$$

$s_0 = 1.00624\dots$



dimensionful parameter (dimensional transmutation)

$\Lambda_* \simeq 49 \text{ MeV} \iff B_t = 8.48 \text{ MeV}$



$$\begin{aligned}
 x &\rightarrow \alpha_n x & \Lambda &\rightarrow \alpha_n^{-1} \Lambda \\
 \frac{t}{m} &\rightarrow \alpha_n^2 \frac{t}{m} & \psi &\rightarrow \alpha_n^{-3/2} \psi
 \end{aligned}$$

$$S_{EFT}^{(0)} \rightarrow S_{EFT}^{(0)}$$

$$\alpha_n = \exp(n\pi/s_0) = (22.7)^n \equiv f^n$$

$$n = \dots, 0, 1, 2, \dots$$

ground state \nearrow



Efimov states

$$B_{3n}^{(0)} = \exp(-2n\pi/s_0) k_*^2 / m_N \lesssim M_{hi}^2 / m_N$$

$$k_* \propto \Lambda_*$$

Efimov '71

...



Similar for bosons

Bedaque, Hammer + v.K. '99 '00

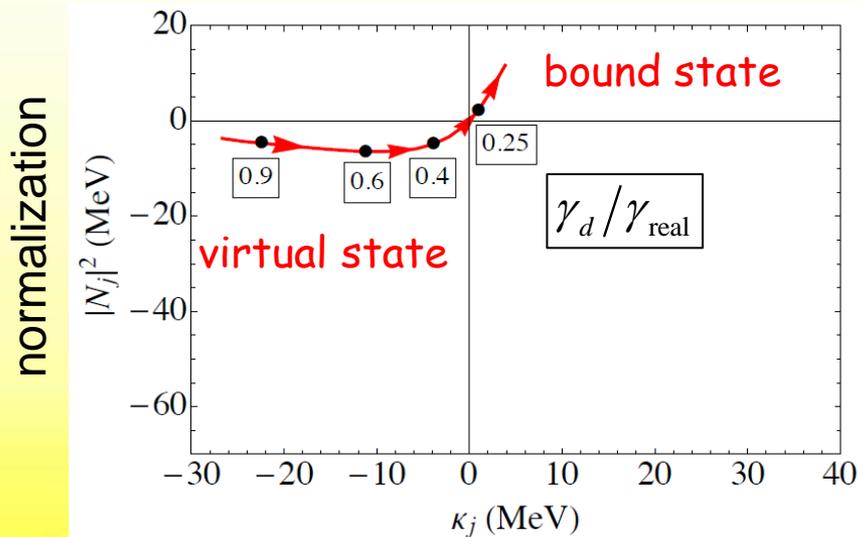
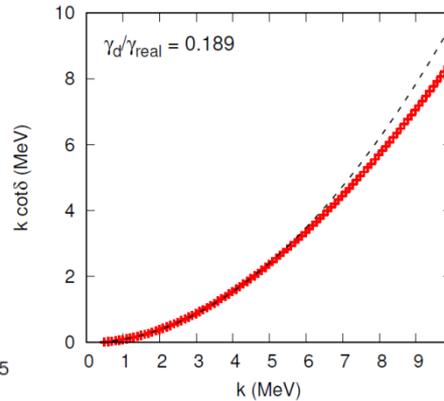
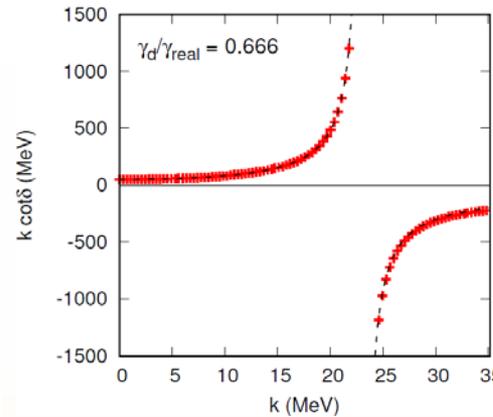
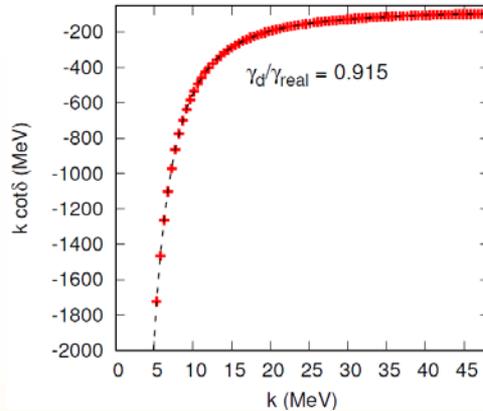
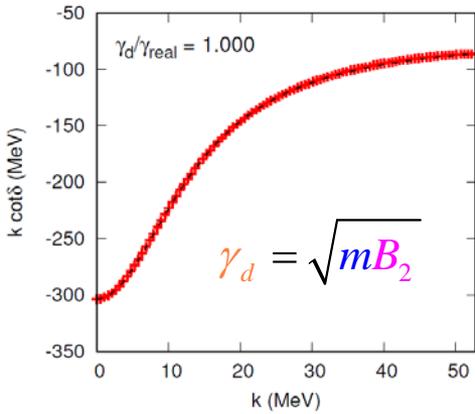
...

LO

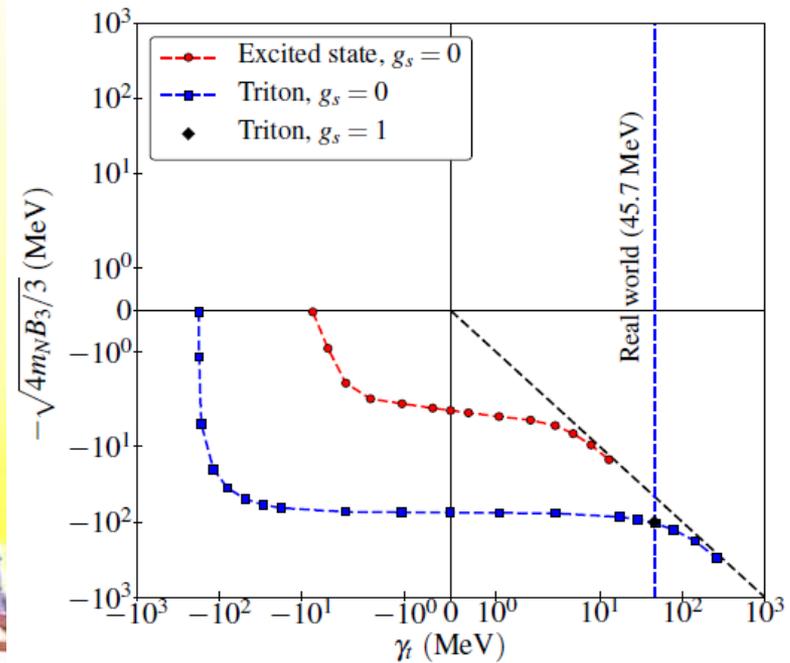
The first excited state of triton

nd scattering below deuteron break-up

→ unitarity

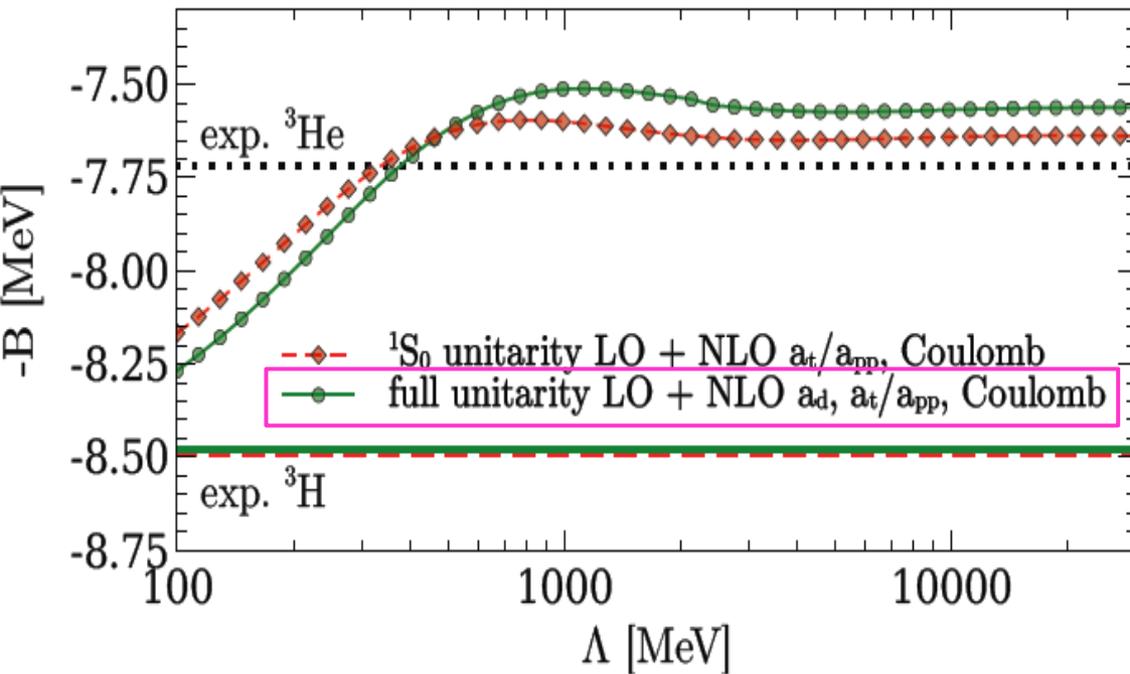


binding momentum



NLO

no new three-body force



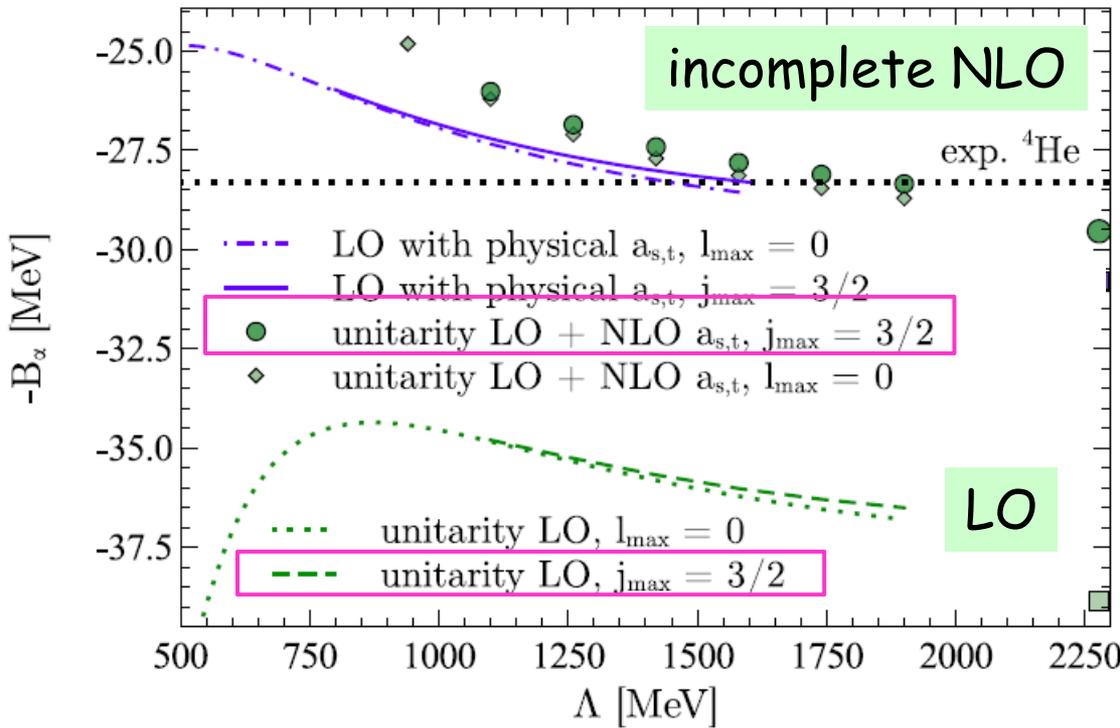
$$B_h^{(1)} - B_t \approx -(0.92 \pm 0.18) \text{ MeV}$$

vs.

$$-0.764 \text{ MeV (exp)}$$

(Fadeev-Yakubovski)

König, Griebhammer,
Hammer + v.K. '16



no ranges
no Coulomb

A = 4

infinite-cutoff extrapolation

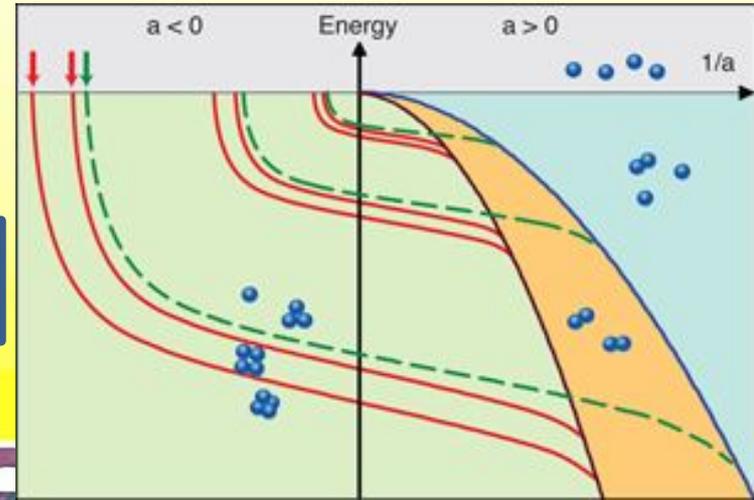
no 4BF up to this order

Hammer, Meißner + Platter '05
Stetcu, Barrett + v.K. '07
...

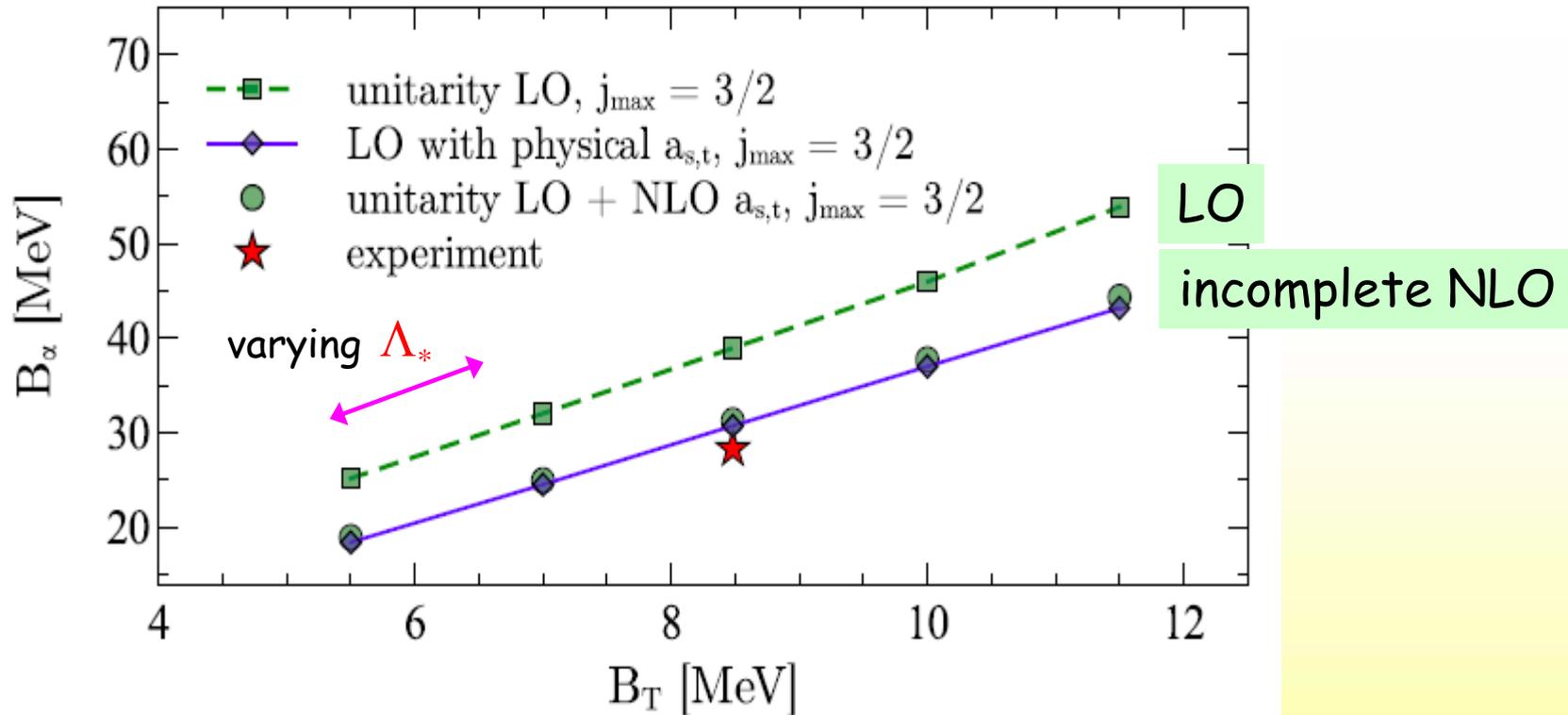
$B_{\alpha}^{(1,inc)} \simeq (29.5 \pm 8.7) \text{ MeV}$ vs. 28.4 MeV (exp)
 $B_{\alpha^*}^{(0)} - B_t \approx 0$ vs. -0.3 MeV (exp)

Hammer + Platter '07
...

Similar for bosons

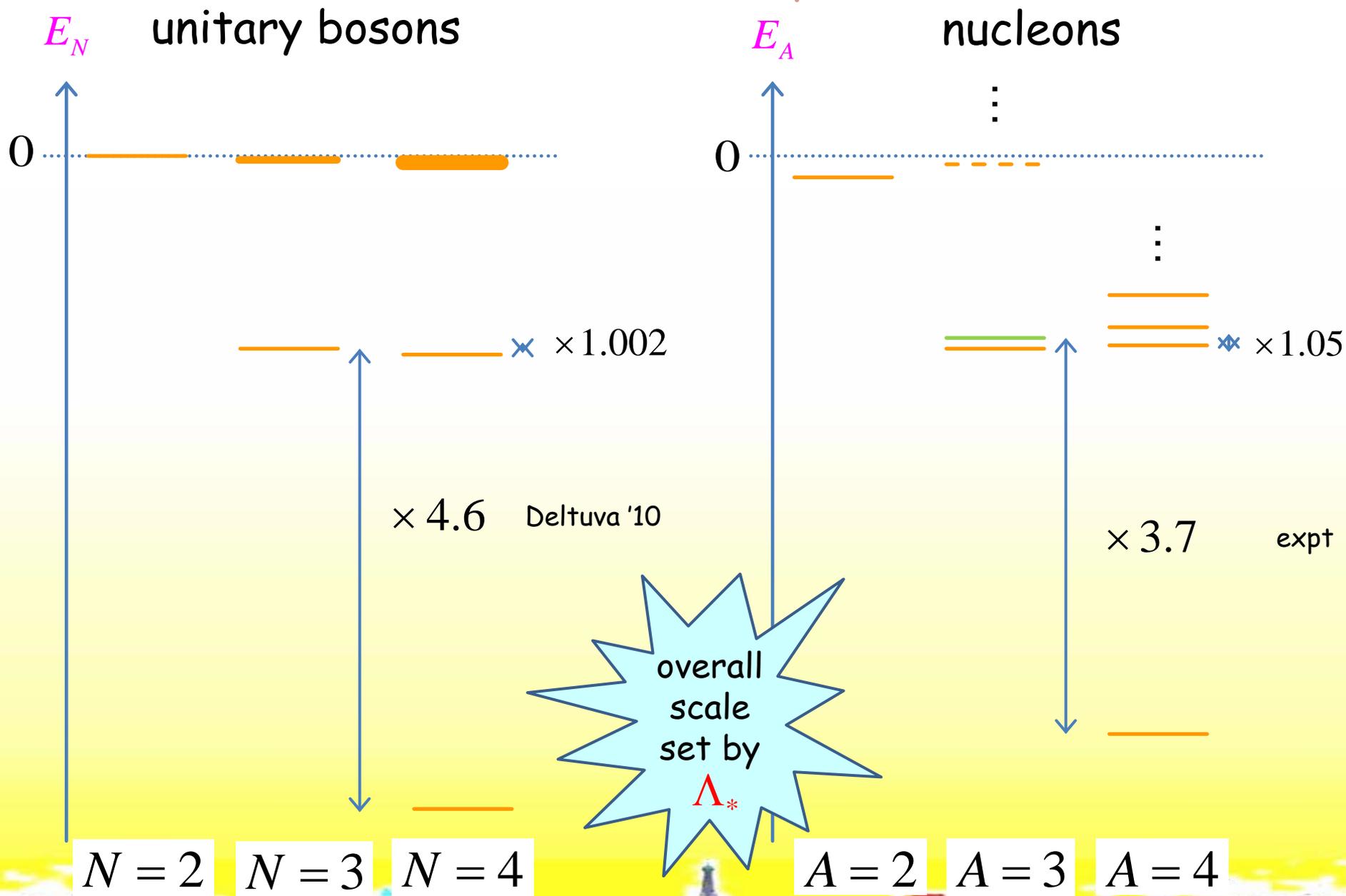


Tjon line



perturbatively close to experimental point

Schematically



Unitary bosons

$$N \geq 4$$

No more-body forces

Bazak, Eliyahu + v.K. '16

LO

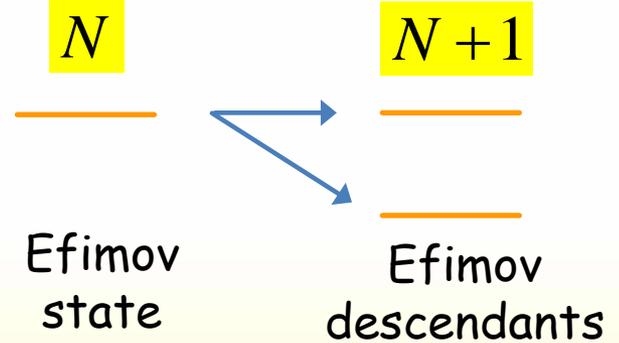


doubling

Hammer + Platter '07

von Stecher '10'11

Gattobigio, Kievsky + Viviani '11'12



Ground states

single scale

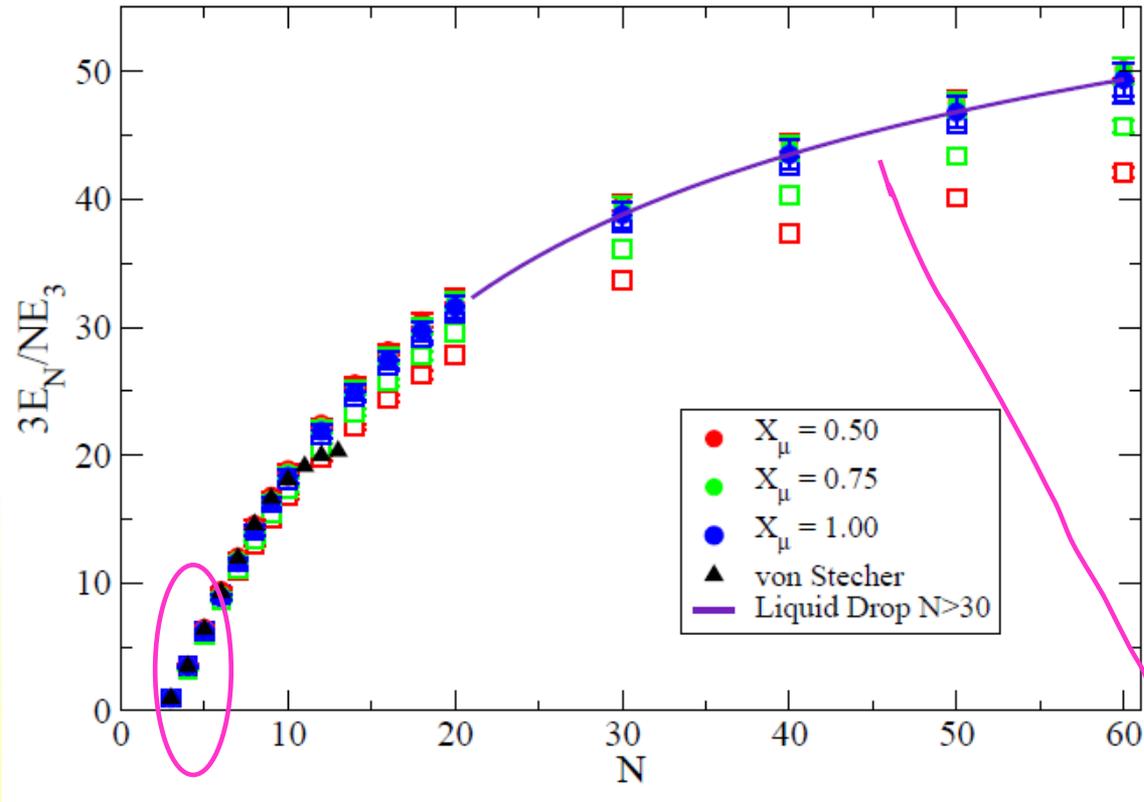


$$\frac{B_N^{(0)}(\Lambda_*)}{N} = \kappa_N \frac{B_3(\Lambda_*)}{3}$$

universal numbers

- $\kappa_2 \equiv 0$
- $\kappa_3 \equiv 1$
- $\kappa_4 \approx 3.5$
- $\kappa_{N \geq 5} \approx ?$

LO



$$\Lambda_2/Q_3 \approx 26 - 38$$

$$\Lambda_3/\Lambda_2 = 0.5 - 1$$

saturation!

cf. Piatecki + Krauth '14
Comparin + Krauth '16

$$\kappa_N \approx \frac{3}{N} (N - 2)^2$$

Bazak, Eliyahu
+ v.K. '16

$$\kappa_N = \kappa_\infty \left[1 + \eta N^{-1/3} + \mathcal{O}(N^{-2/3}) \right]$$

$$\kappa_\infty = 90 \pm 10 \quad \eta = -1.7 \pm 0.3$$



$$A \geq 4$$

Nucleons around unitarity



doubling?

Ground states

single scale

LO



$$\frac{B_A^{(0)}(\Lambda_*)}{A} = \kappa_A \frac{B_3(\Lambda_*)}{3}$$

$$\kappa_2 \equiv 0$$

$$\kappa_3 \equiv 1$$

$$\kappa_4 \approx 3.5$$

$$\kappa_{A \geq 5} \approx ?$$

same as for bosons
should grow slower
than bosons

same
saturation mechanism
as for bosons?



Equation of State at Unitarity

$$k_\rho \equiv (6\pi^2 \rho / g)^{1/3} \quad \text{degeneracy}$$

two-state fermions

scale invariance



$$\lim_{N \rightarrow \infty} \frac{E_N}{N} = \frac{3k_\rho^2}{10m} (\xi + \dots)$$

$\xi \approx 0.4$

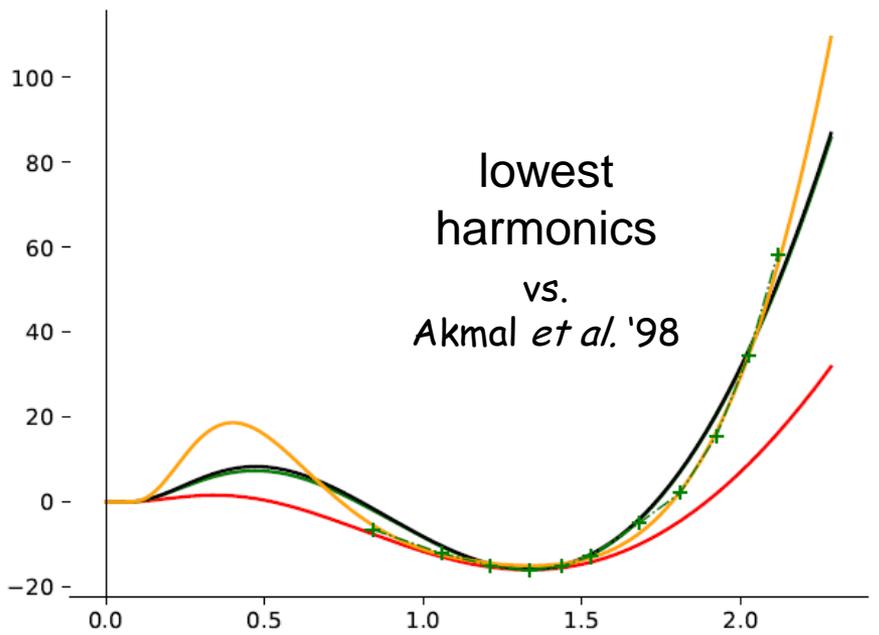
Bertsch '99
Carlson *et al.* '03
...

bosons and nucleons

discrete scale invariance



$$\lim_{N \rightarrow \infty} \frac{E_N}{N} = \frac{3k_\rho^2}{10m} \mathfrak{G}^{(0)} \left[s_0 \ln \left(k_\rho / k_* \right) \right] + \dots$$



$$\mathfrak{G}^{(0)} [\eta + n\pi] = \mathfrak{G}^{(0)} [\eta]$$

$$\equiv \gamma_0 - \sum_{l=1}^{\infty} [\gamma_l \cos(2l(\eta - \eta_0)) + \gamma'_l \sin(2l(\eta - \eta_0))]$$

v.K. '17
Lyu, Long + v.K., in preparation

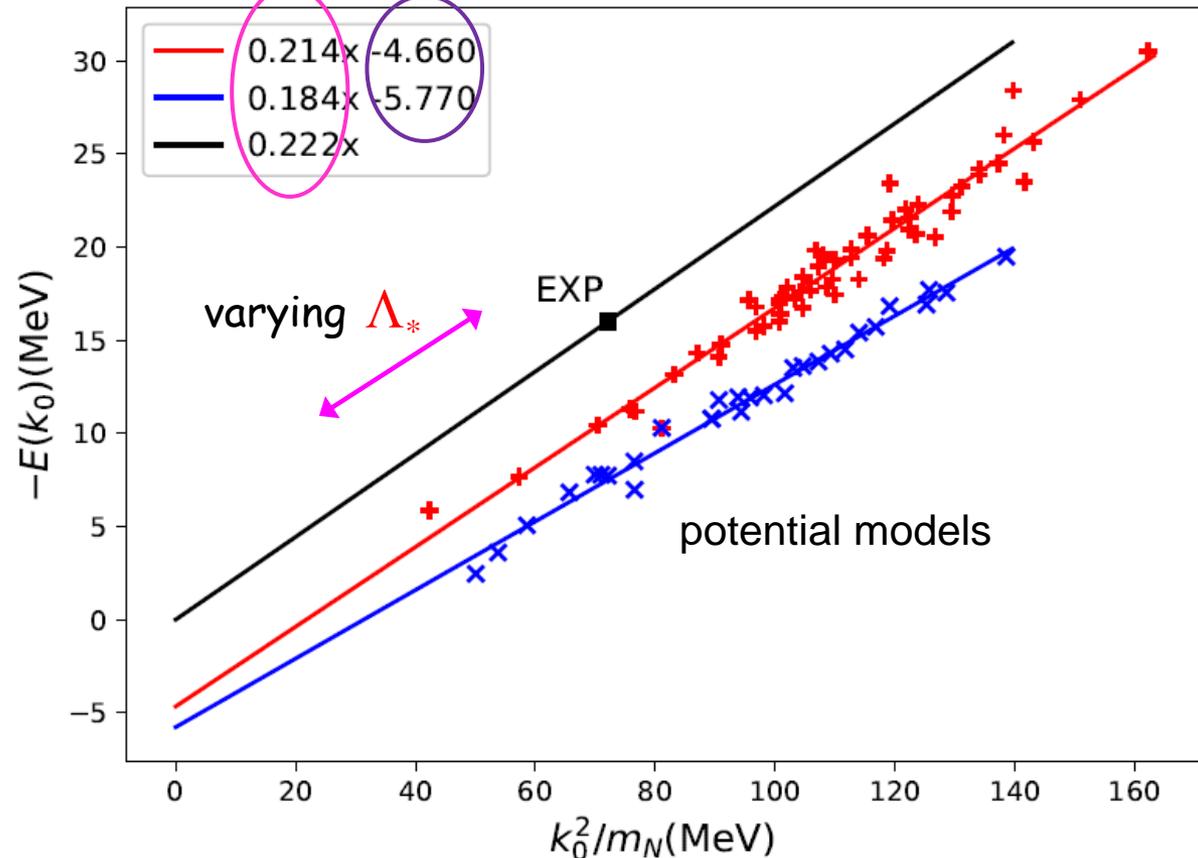


$$\left(\lim_{N \rightarrow \infty} \frac{E_N}{N} \right)_{\text{equil}} = \frac{3}{10} \mathfrak{G}^{(0)}[\eta_0] \frac{k_0^2}{m_N} + \dots \quad \mathfrak{G}^{(0)'}[\eta_0] \approx -\frac{2}{s_0} \mathfrak{G}^{(0)}[\eta_0]$$

v.K. '17
Lyu, Long + v.K.,
in preparation

Coester
line

similar to
Delfino *et al.* '06



Approximate discrete scale invariance
consistent with *ab initio* calculations

Conclusion

- ◆ Systems near unitarity can be described model-independently by Pionless EFT
 - properties given by essentially **one** parameter Λ_*
 - details obtained in perturbation theory
- cf. parallel results from a phenomenological perspective: { Kievsky *et al.*
Grasso *et al.*
- ◆ Light nuclei in a sweet spot:
dilute enough to be insensitive to interaction details,
dense enough for approx discrete scale invariance
 - ◆ A mechanism for nuclear saturation?
 - { more nucleons Gezerlis *et al.*, in progress
 - { nuclear matter Lyu *et al.*, in progress
 - { higher orders Bazak *et al.*, in progress