

LIGHT (ANTI-)NUCLEI AND (ANTI-)HYPERNUCLEI PRODUCTION WITH ALICE AT THE LHC

Stefano Piano on behalf of ALICE Collaboration INFN sez. Trieste



ALICE

MOTIVATION

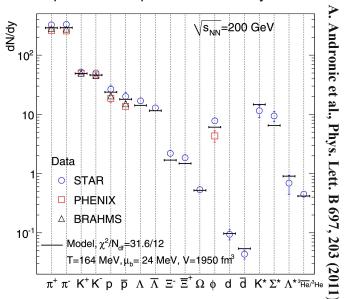
- (anti-)(hyper)nuclei are good probes of coalescence mechanism
- (anti-)(hyper)nuclei yields are sensitive to the freeze-out temperature in heavy-ion collision due to their large mass (e.g. in the Thermal Model yield scales roughly $\propto e^{(-M/Tchem)}$)
- > light (anti-)(hyper)nuclei have small binding energy and small Λ separation energy,
 - e.g. $B_{\Lambda}(^{3}_{\Lambda}H) = 0.13 \pm 0.05$ MeV [H. Bando et al., Int. J. Mod. Phys. A 5 4021 (1990)] :
 - > they should dissociate in a medium with high T_{chem} (~156 MeV) and be suppressed
 - → if their yields equal to thermal model prediction ⇒ sign for adiabatic (isentropic) expansion in the hadronic phase
- \rightarrow A=3 (anti-)(³He, t, ³_{\wedge}H), a simple system of 9 valence quarks:
 - $ightharpoonup ^3 \Lambda H / ^3 He$ and $^3 \Lambda H / t$ (and anti) $\Rightarrow \Lambda$ -nucleon correlation (local baryon-strangeness correlation)
 - ightharpoonup t / ³He (and anti) \Rightarrow local charge-baryon correlation
- Anti-nuclei in nature:
 - matter-antimatter asymmetry
 - [J.Adam et al. (ALICE Collaboration), Nature Phys. 11, no.10, 811 (2015)]
 - ➤ light nuclei measurements in high energy physics can be used in dark matter searches to estimate the background coming from the secondary anti-nuclei [K. Blum, Phys. Rev. D 96 (2017) 103021]



(ANTI-)(HYPER)NUCLEI PRODUCTION IN URHIC

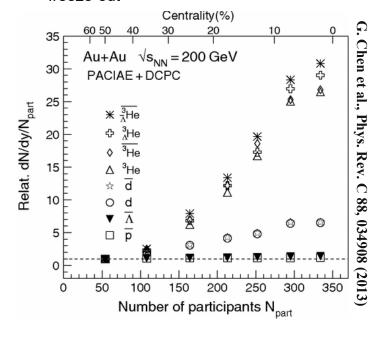
Statistical Thermal model

- Thermodynamic approach to particle production in heavy-ion collisions
- ➤ Abundances fixed at chemical freeze-out (T_{chem}) (hyper)nuclei are very sensitive to T_{chem} because of their large mass (M)
- \triangleright Exponential dependence of the yield $\propto e^{(-M/T \text{chem})}$



Coalescence

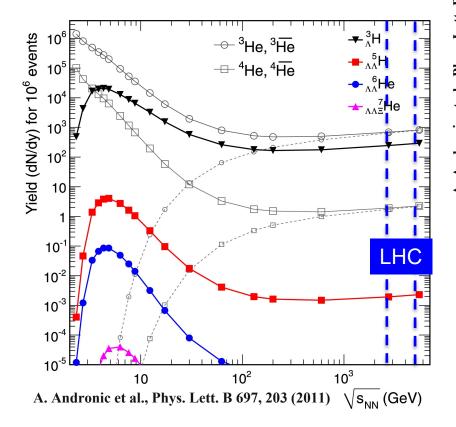
- ➤ If baryons at freeze-out are close enough in Phase Space an (anti-)(hyper)nucleus can be formed
- (Hyper)nuclei are formed by protons (Λ) and neutrons which have similar velocities after the freeze-out





(ANTI-)(HYPER)NUCLEI PRODUCTION AT LHC

Production yield estimate of (anti-)(hyper)nuclei in central heavy-ion collisions at LHC energy based on thermal model:

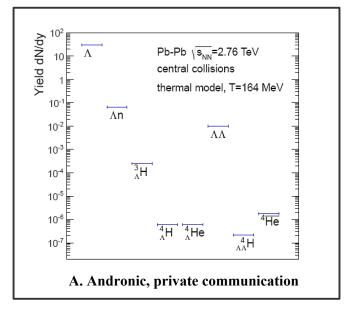


203 (2011)		Yield/event at mid-rapidity and central collisions
onic et al., Fnys. Lett. B 697, 203 (2011)	π	~800
	р	~40
	Λ	~30
	d	~0.17
	³ He	~0.01
	$^3\Lambda$ H	~0.003

Light nuclei

Hypertriton

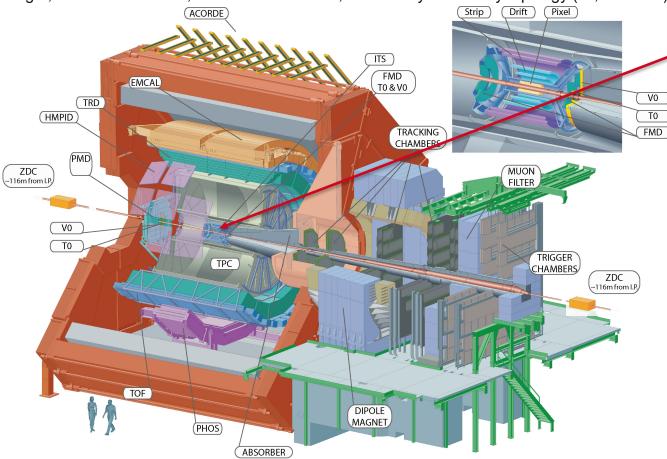
Search for: Λn, ΛΛ dibaryons





ALICE particle identification capabilities are unique. Almost all known techniques are exploited: dE/dx, time-of-

flight, transition radiation, Cherenkov radiation, calorimetry and decay topology (Vo, cascade)



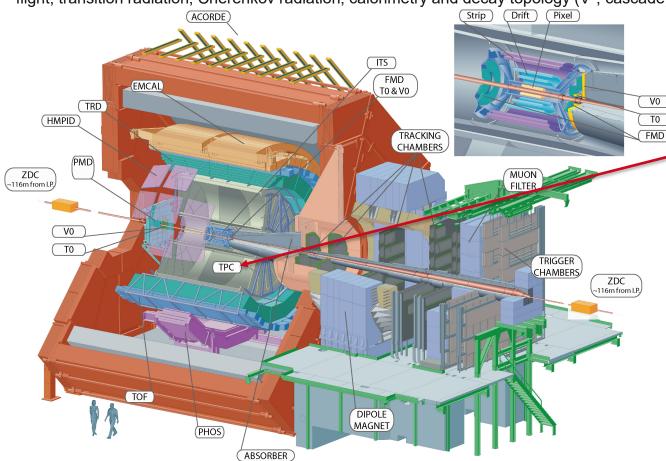
ITS: precise separation of primary particles and those from weak decays (hypernuclei) or knock-out from material

K. Aamodt et al. (ALICE Collaboration), JINST 3 (2008) S08002



ALICE particle identification capabilities are unique. Almost all known techniques are exploited: dE/dx, time-of-

flight, transition radiation, Cherenkov radiation, calorimetry and decay topology (V⁰, cascade)



ITS: precise separation of primary particles and those from weak decays (hypernuclei) or knock-out from material

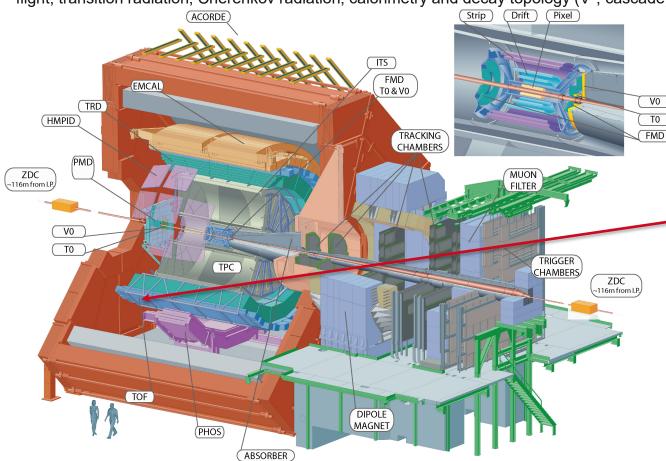
TPC: particle identification via d*E*/d*x* (allows also separation of charges).

K. Aamodt et al. (ALICE Collaboration), JINST 3 (2008) S08002



ALICE particle identification capabilities are unique. Almost all known techniques are exploited: dE/dx, time-of-

flight, transition radiation, Cherenkov radiation, calorimetry and decay topology (Vo, cascade)



ITS: precise separation of primary particles and those from weak decays (hypernuclei) or knock-out from material

TPC: particle identification via d*E*/d*x* (allows also separation of charges).

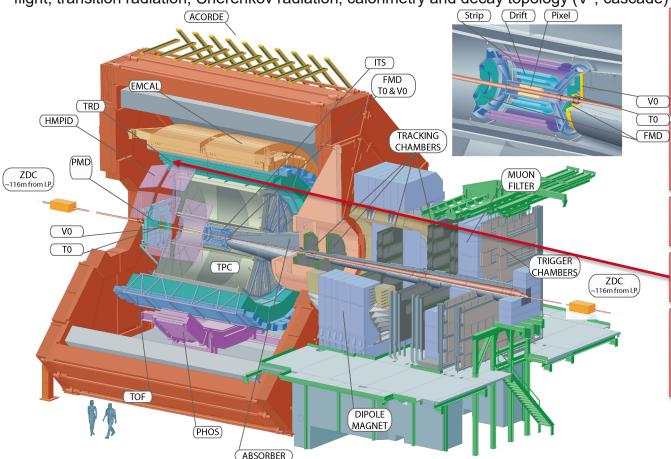
TOF: particle identification via time-of-flight

K. Aamodt et al. (ALICE Collaboration), JINST 3 (2008) S08002



ALICE particle identification capabilities are unique. Almost all known techniques are exploited: dE/dx, time-of-

flight, transition radiation, Cherenkov radiation, calorimetry and decay topology (V⁰, cascade)



ITS: precise separation of primary particles and those from weak decays (hypernuclei) or knock-out from material

TPC: particle identification via d*E*/d*x* (allows also separation of charges).

TOF: particle identification via time-of-flight

TRD: electron identification via transition radiation

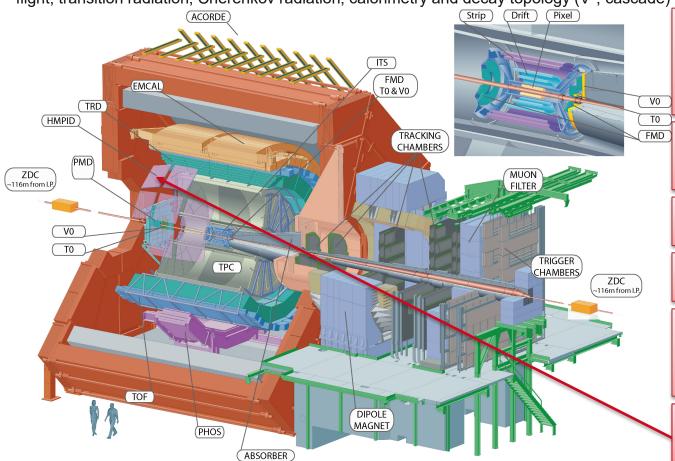
ITS+TPC+TRD: excellent track reconstruction capabilities in a high track density environment

K. Aamodt et al. (ALICE Collaboration), JINST 3 (2008) S08002



ALICE particle identification capabilities are unique. Almost all known techniques are exploited: dE/dx, time-of-

flight, transition radiation, Cherenkov radiation, calorimetry and decay topology (V⁰, cascade)



ITS: precise separation of primary particles and those from weak decays (hypernuclei) or knock-out from material

TPC: particle identification via d*E*/d*x* (allows also separation of charges).

TOF: particle identification via time-of-flight

TRD: electron identification via transition radiation

ITS+TPC+TRD: excellent track reconstruction capabilities in a high track density environment

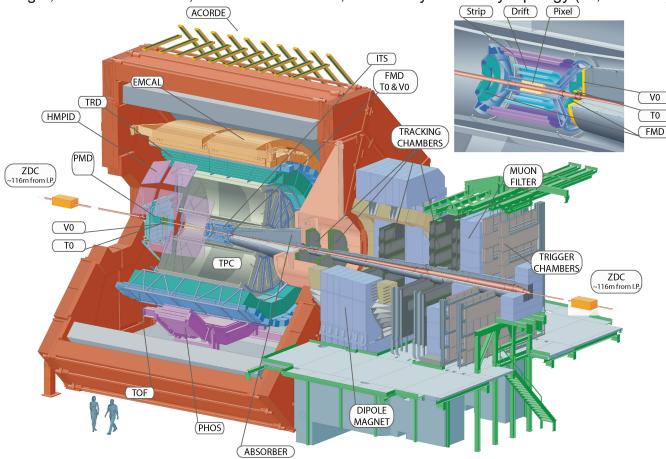
HMPID: particle identification via Cherenkov radiation

K. Aamodt et al. (ALICE Collaboration), JINST 3 (2008) S08002



ALICE particle identification capabilities are unique. Almost all known techniques are exploited: dE/dx, time-of-

flight, transition radiation, Cherenkov radiation, calorimetry and decay topology (V⁰, cascade)



ALICE is ideally suited for the identification of light (anti-)(hyper)nuclei

ITS: precise separation of primary particles and those from weak decays (hypernuclei) or knock-out from material

TPC: particle identification via d*E*/d*x* (allows also separation of charges).

TOF: particle identification via time-of-flight

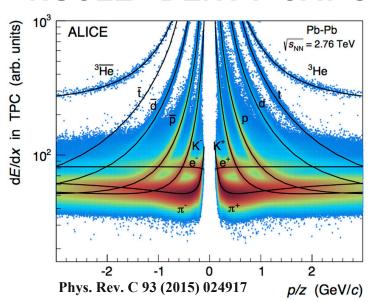
TRD: electron identification via transition radiation

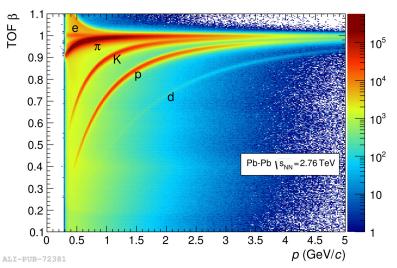
ITS+TPC+TRD: excellent track reconstruction capabilities in a high track density environment

HMPID: particle identification via Cherenkov radiation

ALICE

NUCLEI IDENTIFICATION





Low momenta

Nuclei identification via dE/dx measurement in the TPC:

- ➤ dE/dx resolution in central Pb-Pb collisions: ~6.5%
- Excellent separation of (anti-)nuclei from other particles over a wide momentum range
- ➤ About **10 anti-alpha candidates** identified out of 23x10⁶ events by combining TPC and TOF particle identification

Higher momenta

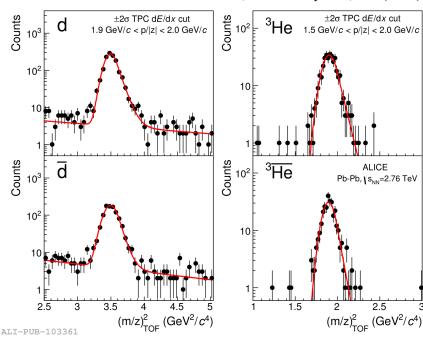
Excellent TOF performance:

- \succ σ_{TOF} ≈ 85 ps in Pb-Pb collisions allows identification of light nuclei over a wide momentum range
- \blacktriangleright Velocity measurement with the TOF detector is used to evaluate the m^2 distribution and to subtract background from the signal in each p_T -bin by fitting the m^2 distribution



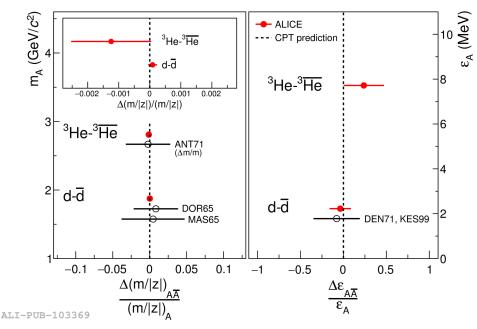
PRECISE MASS MEASUREMENT

ALICE Collaboration, Nature Phys. 11, 811 (2015)



- Masses and binding energies of nuclei and anti-nuclei are compatible within uncertainties
- Measurement confirms the CPT invariance for light nuclei

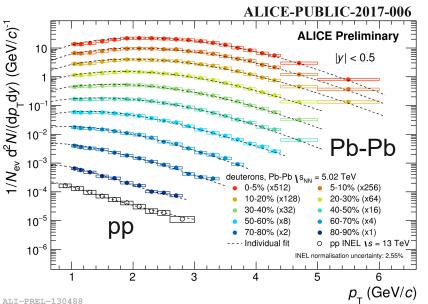
- ➤ The precise measurement of the mass difference between nuclei and their anti-counterparts allows one to probe any difference in the interaction between nucleons and anti-nucleons.
- Looking at the mass difference between nuclei and their anti-nuclei it is possible to test the CPT invariance of the residual QCD "nuclear force"

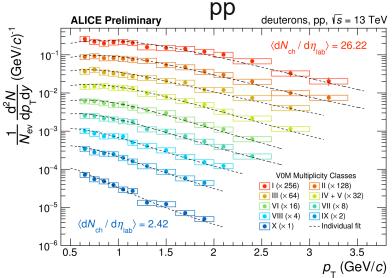


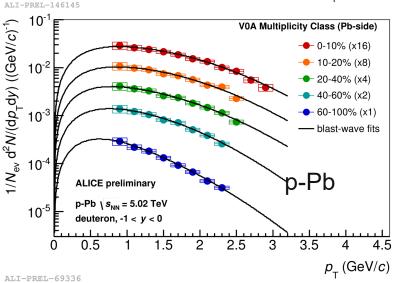


DEUTERON p_T SPECTRA

- Spectra become harder with increasing multiplicity in Pb-Pb and show clear radial flow
- The Blast-Wave fits describe the data well in p-Pb and Pb-Pb
- pp and p-Pb spectrum show no sign of radial flow

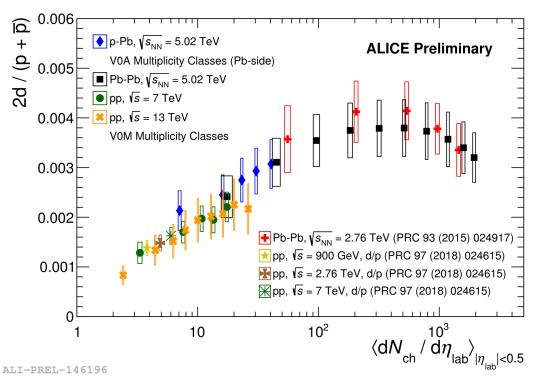








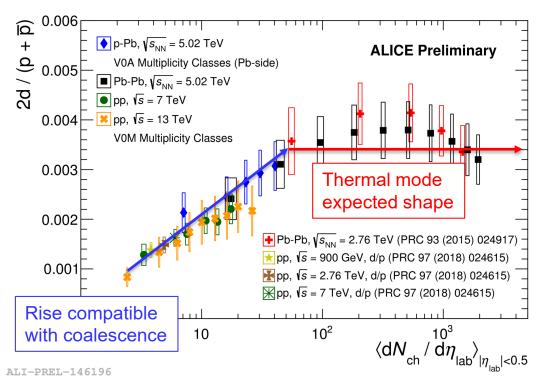
DEUTERON TO PROTON RATIO



- No significant centrality dependence in Pb-Pb
- > Ratio in pp collisions is a factor 2.5 lower than in Pb-Pb collisions
- \triangleright d/p ratio increases when going from pp to p-Pb and peripheral Pb-Pb, until it reaches the grand canonical thermal model value (d/p \sim 3x10⁻³ at T_{ch} = 156 MeV)



DEUTERON TO PROTON RATIO



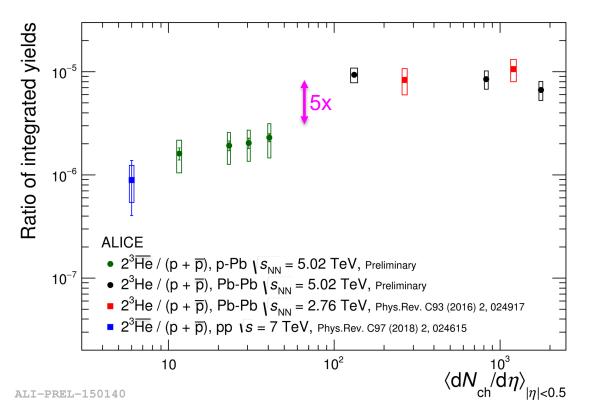
Simple coalescence works in small systems while thermal models describe better the Pb-Pb data

Is this smooth transition suggesting a single description for the nucleosynthesis in HEP?

- No significant centrality dependence in Pb-Pb
- > Ratio in pp collisions is a factor 2.5 lower than in Pb-Pb collisions
- \triangleright d/p ratio increases when going from pp to p-Pb and peripheral Pb-Pb, until it reaches the grand canonical thermal model value (d/p \sim 3x10⁻³ at T_{ch} = 156 MeV)



HELIUM-3 TO PROTON RATIO



Simple coalescence works in small systems while thermal models describe better the Pb-Pb data

Is this smooth transition suggesting a single description for the nucleosynthesis in HEP?

But ... there is a change of factor 5 in the ³He/p between small systems and Pb-Pb

- No significant centrality dependence in Pb-Pb
- > Ratio in pp collisions is a factor 2.5 lower than in Pb-Pb collisions
- \triangleright d/p ratio increases when going from pp to p-Pb and peripheral Pb-Pb, until it reaches the grand canonical thermal model value (d/p \sim 3x10⁻³ at T_{ch} = 156 MeV)

ALICE

COALESCENCE PARAMETER B_A

- ➢ If baryons at freeze-out are close enough in phase space (i.e. geometrically and in momentum) and match spin state, a (anti-)nucleus can be formed
- ➤ Usually, since the nucleus is larger w.r.t. the source, the phase space is reduced to the momentum space
- Assuming that p an n have the same mass and have the same p_T spectra:

$$E_A \frac{\mathrm{d}^3 N_A}{\mathrm{d} p_A^3} = B_A \left(E_\mathrm{p} \frac{\mathrm{d}^3 N_\mathrm{p}}{\mathrm{d} p_\mathrm{p}^3} \right)^A$$

For A=2:
$$B_2 = E_d \frac{\mathrm{d}^3 N_d}{\mathrm{d} p_d^3} \left(E_p \frac{\mathrm{d}^3 N_p}{\mathrm{d} p_p^3} \right)^{-2}$$

Measured deuteron p_{τ} -spectra

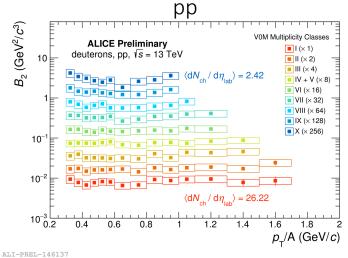
Measured proton p_{T} -spectra

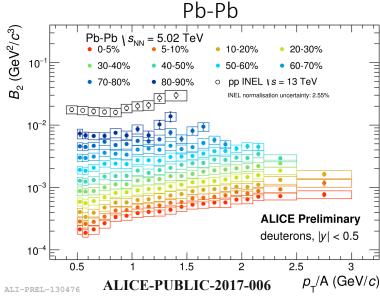


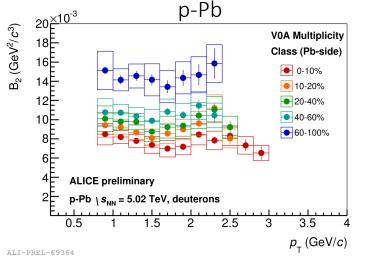
COALESCENCE PARAMETER B₂

$$B_2 = E_d \frac{\mathrm{d}^3 N_d}{\mathrm{d} p_d^3} \left(E_p \frac{\mathrm{d}^3 N_p}{\mathrm{d} p_p^3} \right)^{-2}$$

- ightharpoonup Coalescence parameter B_2 decreases with centrality in Pb-Pb and increases with p_T
- Similar effect seen in p-Pb: decrease with multiplicity, but less pronounced
- Flat behavior
 for pp and no
 dependence
 on multiplicity





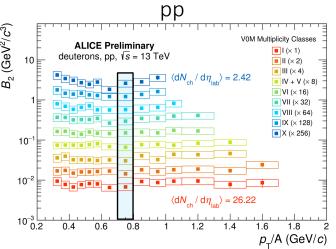


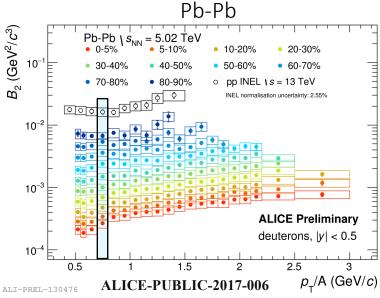


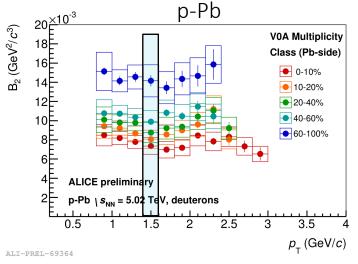
COALESCENCE PARAMETER B₂

$$B_2 = E_d \frac{\mathrm{d}^3 N_d}{\mathrm{d} p_d^3} \left(E_p \frac{\mathrm{d}^3 N_p}{\mathrm{d} p_p^3} \right)^{-2}$$

- \succ Coalescence parameter B_2 decreases with centrality in Pb-Pb and increases with p_{T}
- Similar effect seen in p-Pb: decrease with multiplicity, but less pronounced
- Flat behavior
 for pp and no
 dependence
 on multiplicity





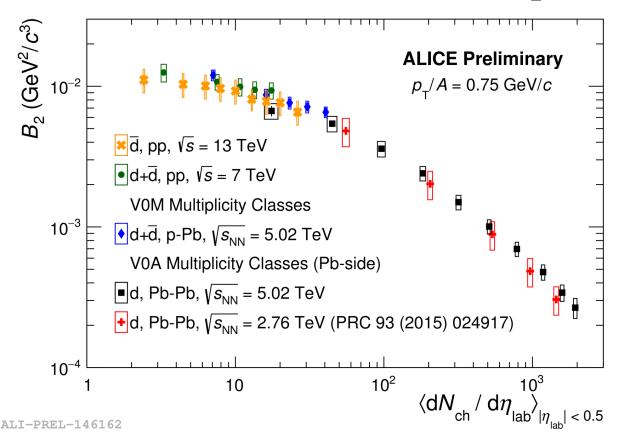


 $p_{\rm T}/A = 0.75 \; {\rm GeV}/c$

2018 European Nuclear Physics Conference | 04-09-2018 | Stefano Piano



COALESCENCE PARAMETER B₂



The coalescence parameter evolves smoothly as a function of multiplicity with no discontinuity between different colliding systems

(*) R. Scheibl and U. Heinz, Phys.Rev. C59, 1585 (1999)

Simple coalescence model

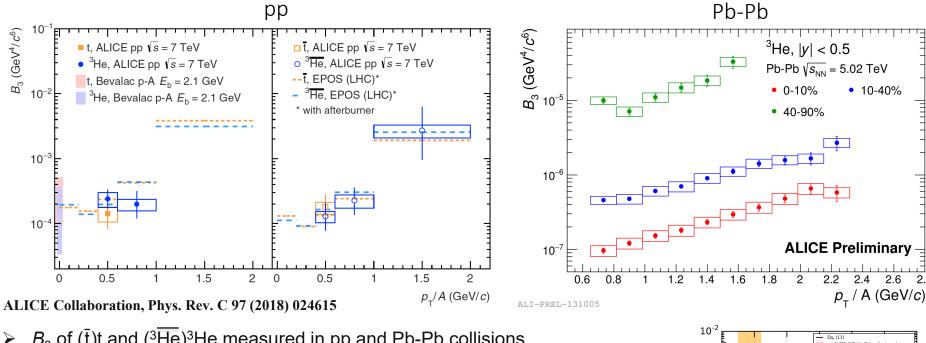
- Flat B₂ vs p⊤ and no dependence on multiplicity/centrality
- Observed "small systems": pp, p-Pb and peripheral Pb-Pb

More elaborated coalescence model takes into account the volume of the source:

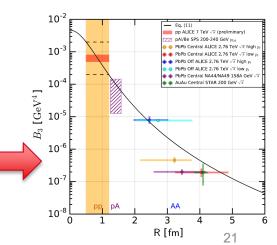
- ➢ B₂ scales like HBT radii (*)
- Decrease with centrality in Pb-Pb is explained as an increase in the source volume
- Increase with p_T in central Pb-Pb reflects the k_T-dependence of the homogeneity volume in HBT
- Qualitative agreement in central Pb-Pb collisions



COALESCENCE PARAMETER B₃

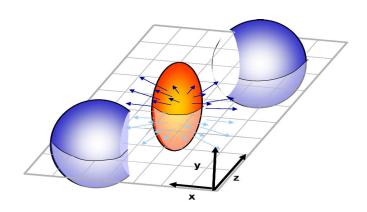


- B_3 of $(\bar{t})t$ and $(^3\overline{He})^3He$ measured in pp and Pb-Pb collisions
- First ever measurements of the B_3 of \bar{t} and ${}^3\overline{\text{He}}$ in pp collisions
- Increasing trend with p_{T} and centrality observed in Pb-Pb collision
- Present data allow to make a new estimate to assess the astrophysical secondary flux of antihelium [K. Blum, Phys. Rev. D 96 (2017) 103021]



ALICE

DEUTERON V₂ FOR DIFFERENT CENTRALITIES



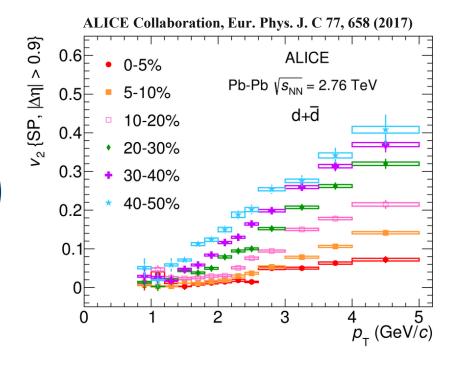
Angular distribution of reconstructed charged particles can be expanded into a Fourier series w.r.t. symmetry plane:

$$E\frac{\mathrm{d}^3 N}{\mathrm{d}p^3} = \frac{1}{2\pi} \frac{\mathrm{d}^2 N}{p_{\mathrm{T}} \mathrm{d}p_{\mathrm{T}} \mathrm{d}y} \left(1 + \sum_{n=1}^{\infty} 2v_n \cos\left(n\left(\varphi - \Psi_n\right)\right) \right)$$

$$v_n = \langle \cos (n(\varphi - \Psi_n)) \rangle$$

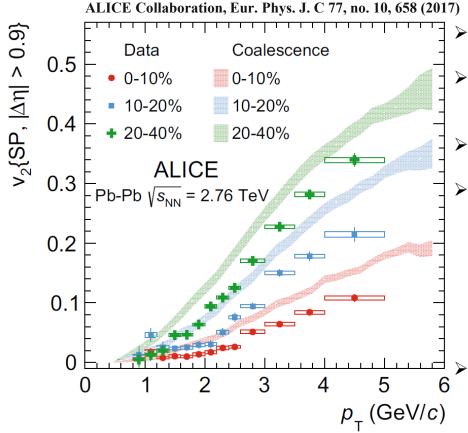
Elliptic flow (v₂) is sensitive to the system evolution:
 It probes initial conditions and constrains particle production mechanisms

- \triangleright A significant v_2 is observed for deuterons
- The value of $v_2(p_T)$ increases progressively from central to semi-central collisions





DEUTERON v₂ FOR DIFFERENT CENTRALITIES



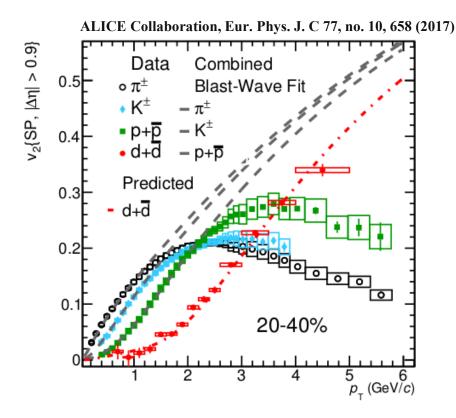
- A significant v_2 is observed for deuterons
- The value of v_2 (p_T) increases progressively from central to semi-central collisions
- If protons have only elliptical flow
- And if the light nucleus is formed by simple coalescence, v_2 can be expressed as

$$v_{2,d}(p_{\rm T}) = \frac{2v_{2,p}(p_{\rm T}/2)}{1 + 2v_{2,p}^2(p_{\rm T}/2)}$$

Such a simple coalescence model is not able to reproduce the measured elliptic flow of deuterons



DEUTERON v₂ FOR DIFFERENT CENTRALITIES



Simplified hydro model (Blast-Wave) is able to describe spectra and flow at the same time, suggesting an early "freeze out"

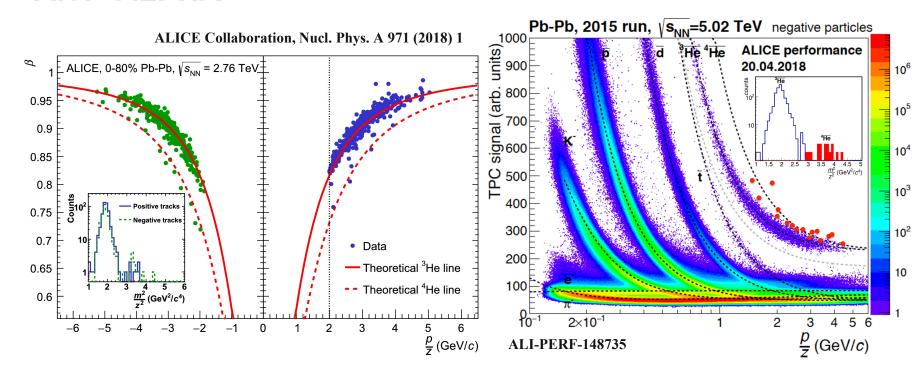
- \triangleright A significant v_2 is observed for deuterons
- The value of v_2 (p_T) increases progressively from central to semi-central collisions
- If protons have only elliptical flow
- ➤ And if the light nucleus is formed by simple coalescence, *v*₂ can be expressed as

$$v_{2,d}(p_{\rm T}) = \frac{2v_{2,p}(p_{\rm T}/2)}{1 + 2v_{2,p}^2(p_{\rm T}/2)}$$

Such a simple coalescence model is not able to reproduce the measured elliptic flow of deuterons



ANTI-ALPHA

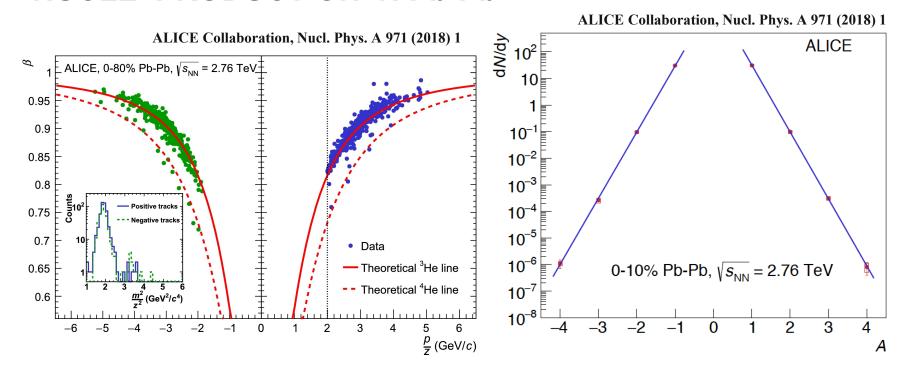


For the full statistics of 2011 (Pb-Pb@2.76 TeV) ALICE identified 10 Anti-Alphas using TPC and TOF

TOF β vs p/z after pre-selection of 3σ in TPC shows clear separation \rightarrow Cut on Alpha's p/z needed to suppress secondary Anti-Alpha re-measured also in the new data sample Pb-Pb@5.02 TeV



NUCLEI PRODUCTION IN Pb-Pb



For the full statistics of 2011 (Pb-Pb@2.76 TeV) ALICE identified 10 Anti-Alphas using TPC and TOF

TOF β vs p/z after pre-selection of 3σ in TPC shows clear separation \rightarrow Cut on Alpha's p/z needed to suppress secondary

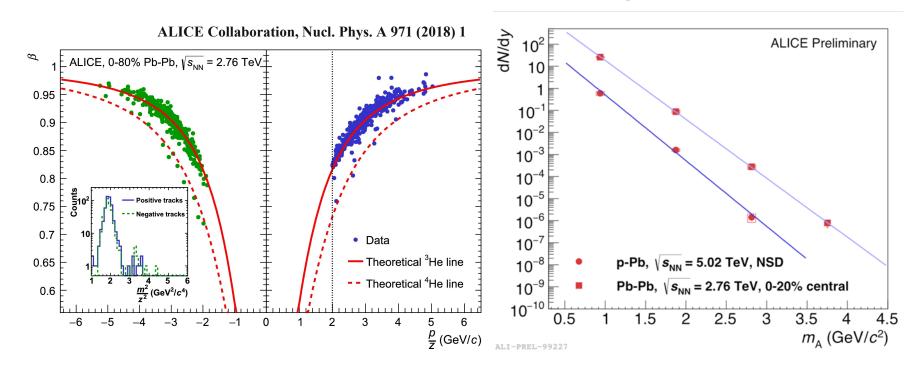
Anti-Alpha re-measured also in the new data sample Pb-Pb@5.02 TeV

Nuclei yields follow an exponential decrease with mass as predicted by the thermal model

In **Pb-Pb** the penalty factor for adding one baryon is ~350 (for particles and antiparticles)



NUCLEI PRODUCTION IN Pb-Pb AND IN p-Pb



For the full statistics of 2011 (Pb-Pb@2.76 TeV) ALICE identified 10 Anti-Alphas using TPC and TOF

TOF β vs p/z after pre-selection of 3σ in TPC shows clear separation \rightarrow Cut on Alpha's p/z needed to suppress secondary Anti-Alpha re-measured also in the new data sample Pb-Pb@5.02 TeV

Nuclei yields follow an exponential decrease with mass as predicted by the thermal model

In Pb-Pb the penalty factor for adding one baryon is ~350 (for particles and antiparticles) and in p-Pb it is ~600



(ANTI-)HYPERTRITON IDENTIFICATION

Decay Channels

$$\frac{{}^{3}_{\Lambda} H \rightarrow {}^{3}_{\Pi} H e + \pi^{-}}{{}^{3}_{\Lambda} \overline{H}} \rightarrow {}^{3}_{\Pi} \overline{H} e + \pi^{+}}$$

$$\frac{{}^{3}_{\Lambda} H \rightarrow {}^{3}_{\Pi} H + \pi^{0}}{{}^{3}_{\Lambda} \overline{H}} \rightarrow {}^{3}_{\Pi} \overline{H} + \pi^{0}}$$

$$\frac{{}^{3}_{\Lambda} H \rightarrow d + p + \pi^{-}}{{}^{3}_{\Lambda} \overline{H}} \rightarrow \overline{d} + \overline{p} + \pi^{+}}$$

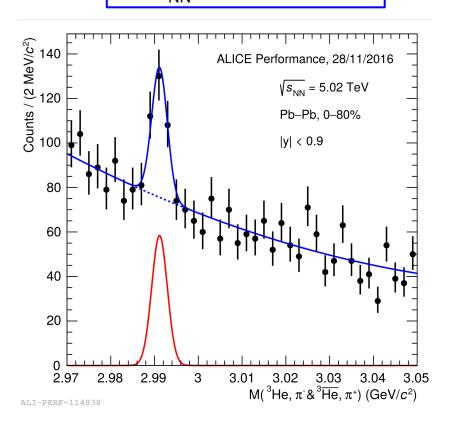
$$\frac{{}^{3}_{\Lambda} H \rightarrow d + n + \pi^{0}}{{}^{3}_{\Lambda} \overline{H}} \rightarrow \overline{d} + \overline{n} + \pi^{0}}$$

- ³_∧H search via two-body decays into charged particles:
- > Two body decay: lower combinatorial background
- ➤ Charged particles: ALICE acceptance and reconstruction efficiency for charged particles higher than for neutrals

Signal extraction:

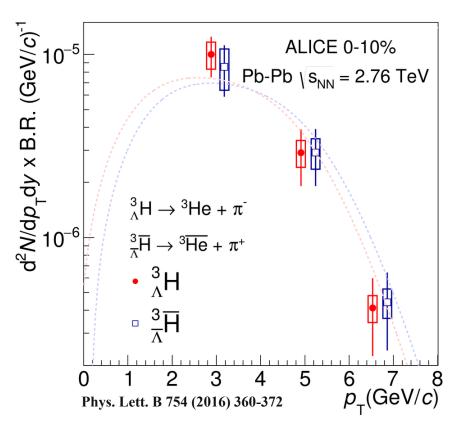
- > Identify ³He and π
- Evaluate (³He,π) invariant mass
- > Apply topological cuts in order to:
 - isolate secondary decay vertex and
 - reduce combinatorial background

New preliminary results at $\sqrt{s_{NN}} = 5.02 \text{ TeV}$





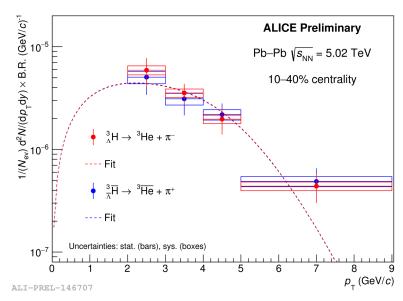
(ANTI-)HYPERTRITON YIELDS



Pb-Pb@2.76Tev:

d*N*/d*y* x B.R. (3 _ΛH \rightarrow 3 He π) yield extracted in three p_T bins for central (0-10%) events

for ${}^3_{\overline{\Lambda}}\overline{H}$ and ${}^3_{\Lambda}H$ separately

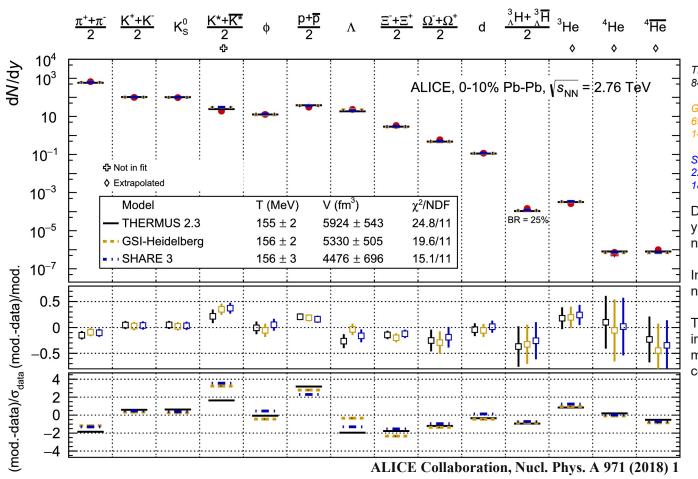


Pb-Pb@5.02Tev:

 $dN/dy \times B.R. (^3_{\Lambda}H \rightarrow ^3He \pi)$ yield extracted in four p_T bins for central (10-40%) events for $^3_{\Lambda}\overline{H}$ and $^3_{\Lambda}H$ separately



THERMAL MODEL FITS



THERMUS: S. Wheaton, et al., CPC 180, 84 (2009)

GSI-Heidelberg: A. Andronic, et al., PLB 697, 203 (2011); PLB 673, 142 (2009) 142

SHARE3: G. Torrieri, et al., CPC 167, 229 (2005); CPC 175, 635 (2006); CPC 185, 2056 (2014)

Different models describe particle yields including light (hyper-) nuclei with $T_{\rm ch}$ of about 156 MeV

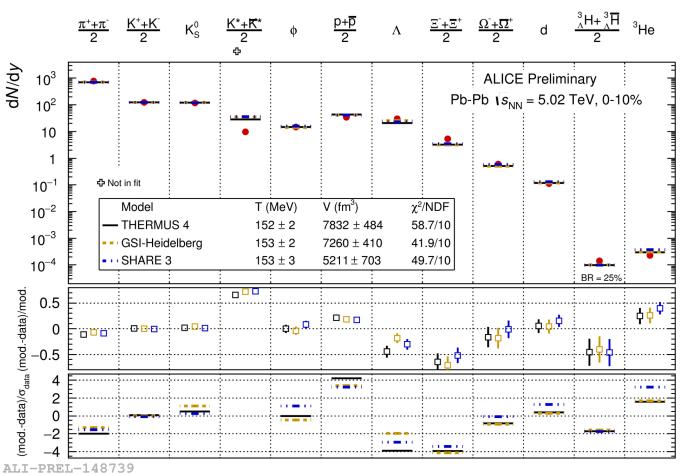
Including nuclei in the fit causes no significant change in T_{ch}

Thermal model is very successful in reproducing the particle yields measured by ALICE in Pb-Pb collisions at $\sqrt{s_{\rm NN}}$ =2.76 TeV

This result, together with the successful Blast-Wave fit to deuteron data suggest that nuclei production happens at the hadronisation, when all the other particles are formed.

ALICE

THERMAL MODEL FITS



THERMUS: S. Wheaton, et al., CPC 180, 84 (2009)

GSI-Heidelberg: A. Andronic, et al., PLB 697, 203 (2011); PLB 673, 142 (2009) 142

SHARE3: G. Torrieri, et al., CPC 167, 229 (2005); CPC 175, 635 (2006); CPC 185, 2056 (2014)

Different models describe particle yields including light (hyper-) nuclei with $T_{\rm ch}$ of about 152 MeV

Including nuclei in the fit causes no significant change in $T_{\rm ch}$

The larger data sample collected in LHC Run2 and improved reconstruction and analysis techniques reduced the uncertainties.

The preliminary data from LHC Run 2 confirm the tensions seen with data from LHC Run 1.

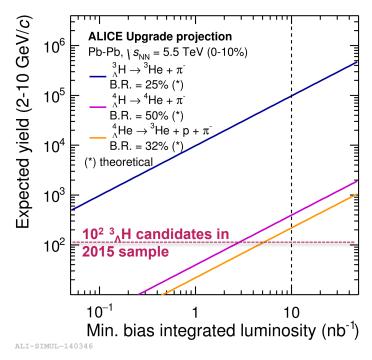


OUTLOOK – ALICE UPGRADE

At the end of Pb-Pb during RUN2 (Nov. 2018) the expected statistics for A=2,3 is >2x

During the Long Shutdown 2 (2019-2020):

- New Inner Tracking System (ITS)
 - √ improved pointing precision
 - √ less material -> thinnest tracker at the LHC
- Upgrade of Time Projection Chamber (TPC):
 - √ new GEM technology for readout chambers
 - ✓ continuous readout
 - √ faster readout electronics
- High Level Trigger (HLT):
 - √ new architecture
 - ✓ on line tracking & data compression
 - √ 50kHz Pb-Pb event rate



At the end of RUN3 (2023): the expected Integrated Luminosity: ~10 nb⁻¹ (~8x10⁹ collisions in the 0-10% centrality class)

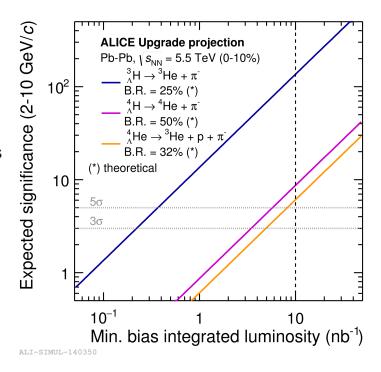


OUTLOOK – ALICE UPGRADE

At the end of Pb-Pb during RUN2 (Nov. 2018) the expected statistics for A=2,3 is >2x

During the Long Shutdown 2 (2019-2020):

- New Inner Tracking System (ITS)
 - √ improved pointing precision
 - √ less material -> thinnest tracker at the LHC
- Upgrade of Time Projection Chamber (TPC):
 - ✓ new GEM technology for readout chambers
 - ✓ continuous readout
 - √ faster readout electronics
- High Level Trigger (HLT):
 - √ new architecture
 - ✓ on line tracking & data compression
 - √ 50kHz Pb-Pb event rate



At the end of RUN3 (2023):

the expected Integrated Lyminosity: ~10 pb-1 (~8x109 colli

the expected Integrated Luminosity: ~10 nb⁻¹ (~8x10⁹ collisions in the 0-10% centrality class)

All the physics which is now done for A = 2 and A = 3 (hyper-)nuclei will be done for A = 4



CONCLUSIONS

- Excellent ALICE performance allows for detection of light (anti-)nuclei and (anti-)hypernuclei
- ALICE (anti-)(hyper)nuclei production measurements challenge the production models:
 - ✓ Large production of nuclear clusters measured by ALICE as predicted by the thermal models
 - ✓ Thermal models describe sufficiently well the nuclei production (from proton to ⁴He) in Pb-Pb
 - ✓ Simple coalescence models describe the (anti-)(hyper)nuclei only in small systems (pp, p-Pb)
 - ✓ Deuteron yield and $v_2(p_T)$ in Pb-Pb collisions suggest an early "freeze out", while large effects of re-interactions (favoring late stage coalescence) should be expected
 - ✓ B₂ coalescence parameter evolves smoothly as a function of multiplicity with no discontinuity between different colliding system
- Future LHC runs, RUN2 and RUN3, and ALICE upgrades will allow for precise study of (anti-)(hyper)nuclei production yield (and lifetime)
 - ✓ New and more precise data are expected from the LHC in the next years!



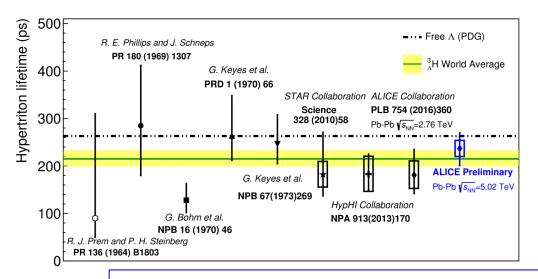


HYPERTRITON LIFETIME DETERMINATION

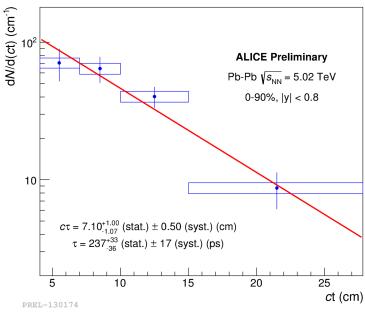
Direct decay time measurement is difficult (~ps), but the excellent determination of primary and decay vertex allows measurement of lifetime via:

$$N(t) = N(0) e^{-\frac{t}{\tau}}$$

where $t = L/(\beta \gamma c)$ and $\beta \gamma c = p/m$ with m the hypertriton mass, p the total momentum and L the decay length



New preliminary results at $\sqrt{s_{NN}}$ = 5.02 TeV



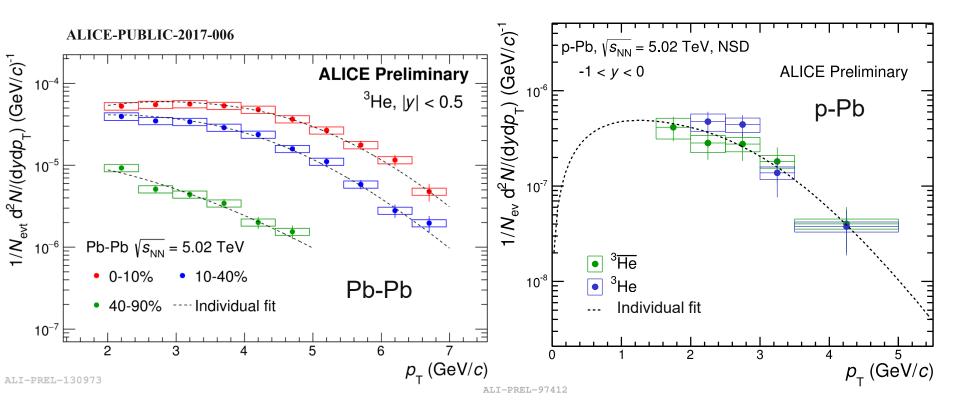
$$c\tau = \left(7.10^{+1.00}_{-1.07}(stat.) \pm 0.50(syst.)\right) cm$$
$$\tau = \left(237^{+33}_{-36}(stat.) \pm 17(syst.)\right) ps$$

ALI-PREL-130195

- Previous heavy-ion experiment results show a trend well below the free ∧ lifetime
- ALICE preliminary result from Pb-Pb at 5.02 TeV is closer to the free Λ lifetime



3-HELIUM

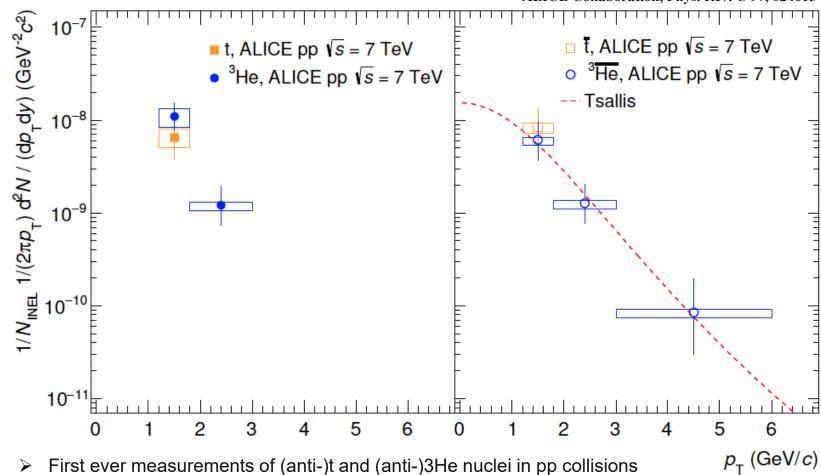


- Dashed curves represent individual Blast-Wave fits
- Spectrum obtained in 3 centrality classes in Pb-Pb and for NSD collisions in p-Pb

ALICE

3-HELIUM AND TRITON

ALICE Collaboration, Phys. Rev. C 97, 024615

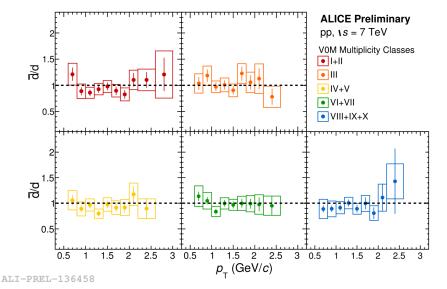


- > t and anti-t measurement difficult
- (anti-)t/(anti-)³He agrees with unity

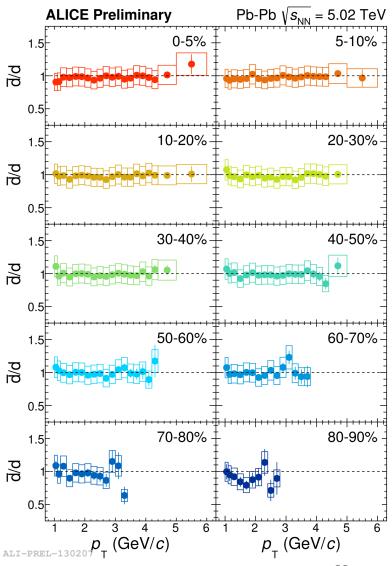


ANTI-NUCLEI PRODUCTION

- Anti-nuclei / nuclei ratios are consistent with unity (similar to other light flavour species)
- ightharpoonup Ratios exhibit constant behavior as a function of p_T and centrality
- Ratios are compatible with unity, in agreement with the coalescence and thermal model expectations
- Also in pp multiplicity intervals, anti-deuterons and deuterons are produced equally



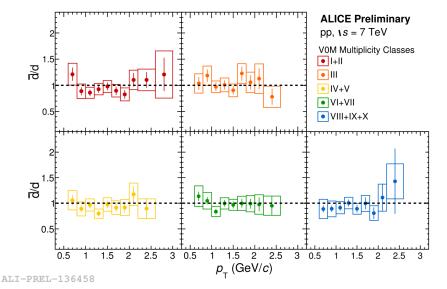
ALICE-PUBLIC-2017-006



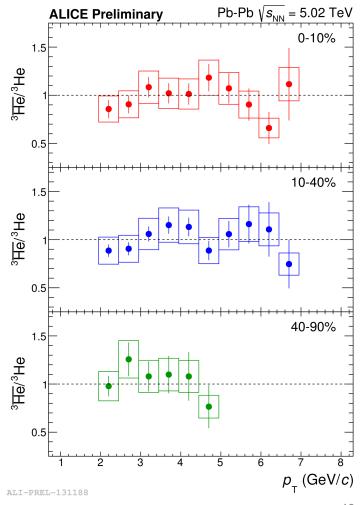


ANTI-NUCLEI PRODUCTION

- Anti-nuclei / nuclei ratios are consistent with unity (similar to other light flavour species)
- ightharpoonup Ratios exhibit constant behavior as a function of p_T and centrality
- Ratios are compatible with unity, in agreement with the coalescence and thermal model expectations
- Also in pp multiplicity intervals, anti-deuterons and deuterons are produced equally



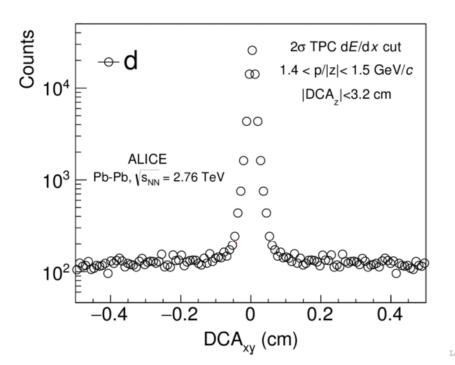
ALICE-PUBLIC-2017-006

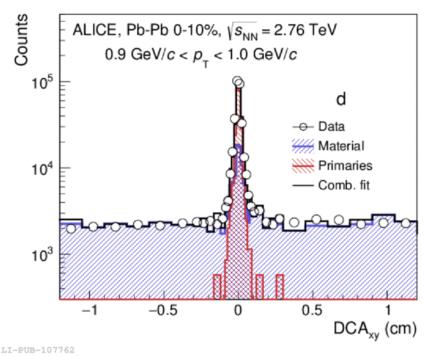




SECONDARY CONTAMINATION

- The measurement of nuclei is strongly affected by background from knock-out from material
- Rejection is possible by fitting the DCA_{XY} distributions with templates

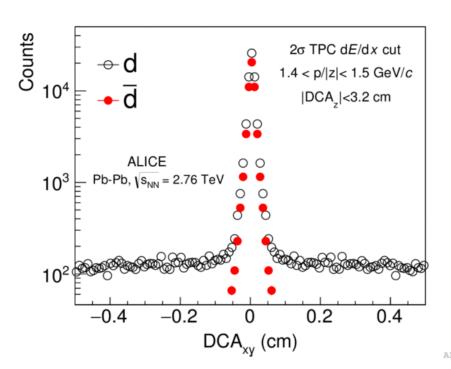


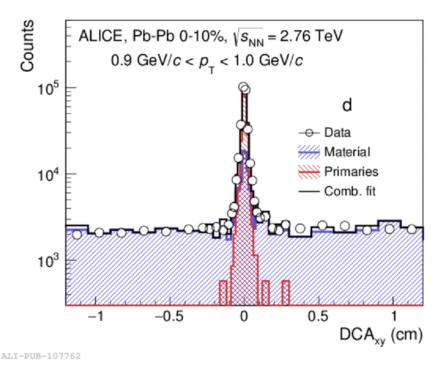




SECONDARY CONTAMINATION

- The measurement of nuclei is strongly affected by background from knock-out from material
- \triangleright Rejection is possible by fitting the DCA_{xy} distributions with templates
- Not relevant for anti-nuclei. However, their measurement suffers from large systematics related to unknown hadronic interaction cross-sections of anti-nuclei in material







An AND H-DIBARYON SEARCH

H-Dibaryon: hypothetical udsuds bound state

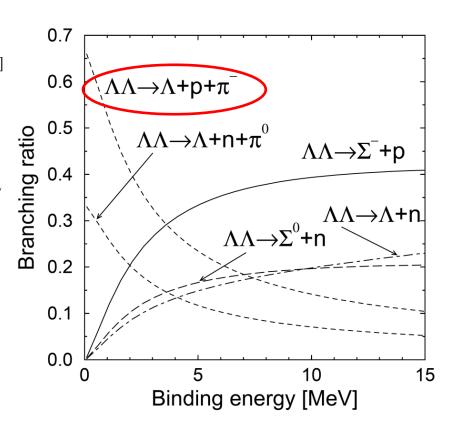
- First predicted by Jaffe [Jaffe, PRL 38, 195617 (1977)]
- Several predictions of bound and also resonant states.
- ➤ Recent Lattice models predict weakly bound states [Inoue et al., PRL 106, 162001 (2011), Beane et al., PRL 106, 162002 (2011)]

If H-Dibaryon is bound: $m_H < \Lambda \Lambda$ threshold

 \succ measurable channel H \rightarrow Λpπ but BR depends on binding energy

Bound state of Λn ?

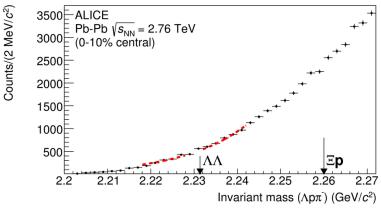
 \triangleright HypHI experiment at GSI sees evidence of a new state: Λn → d + π⁻ [C. Rappold et al. (HypHI collaboration), Phys. Rev. C88, 041001(R) (2013)]

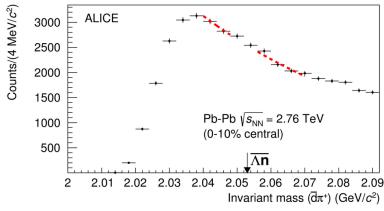


Schaffner-Bielich et al., PRL 84, 4305 (2000)

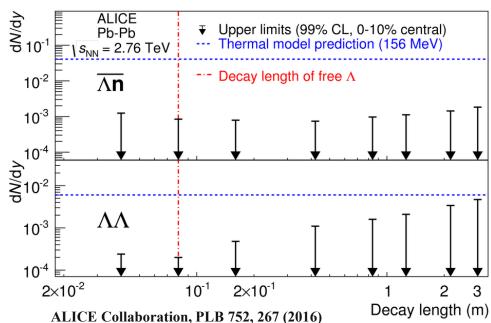


An AND H-DIBARYON SEARCH





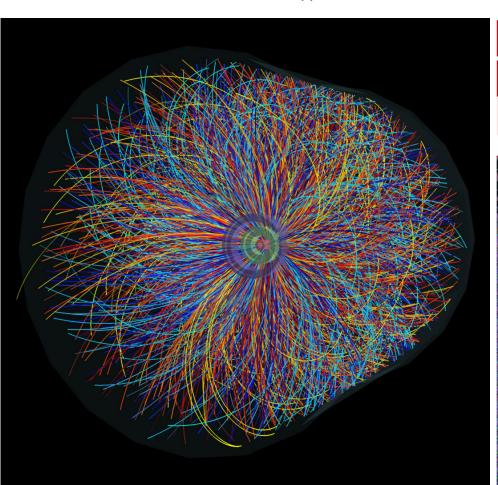
- ➤ No signal visible
- ➤ The upper limits for exotica are lower than the thermal model expectation by a factor 20
- ➤ Thermal models with the same temperature describe precisely the production yield of deuterons, ³He and ³_ΛH
- ➤ The existence of such states with the assumed B.R., mass and lifetime is questionable



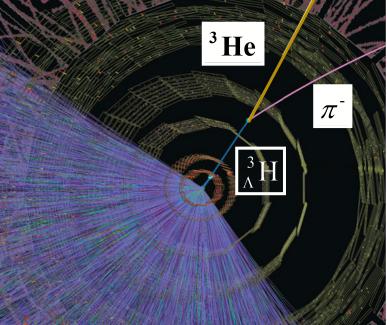


THE EXPERIMENTAL CHALLENGE

The challenge: extract the ³_AH signal from an overwhelming background



At √s _{NN} =	2.76 TeV	5.02 TeV
Centrality	$dN_{ch}/d\eta \; (\eta < 0.5)$	
0-5 %	1601 ± 60	1943 ± 54

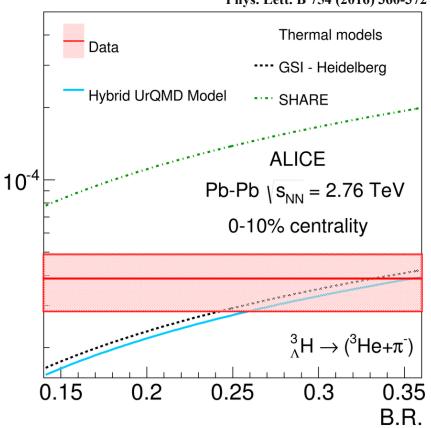


K. Aamodt et al. (ALICE Collaboration) Phys. Rev. Lett. 106, 032301 (2011); J. Adam et al (ALICE Collaboration) Phys. Rev. Lett. 116, 222302 (2016)



COMPARISON WITH THEORETICAL PREDICTIONS

Phys. Lett. B 754 (2016) 360-372

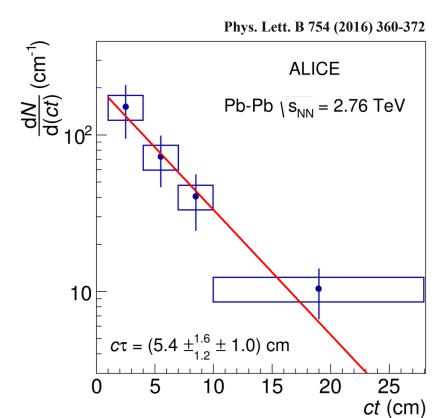


Three different theoretical predictions drawn as a function of BR($^{3}_{\Lambda}H \rightarrow {}^{3}He+\pi^{-}$) after being multiplied by BR:

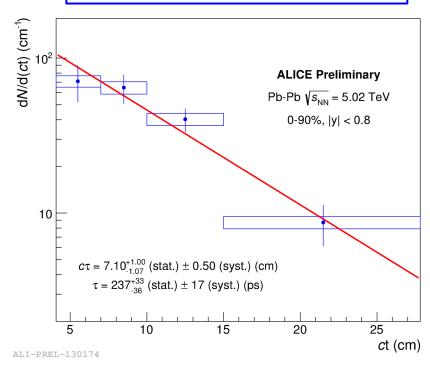
- ✓ Hybrid UrQMD: combines the hadronic transport
 approach with an initial hydrodynamical stage for the
 hot and dense medium (J. Steinheimer et al., Phys.
 Lett. B 714, 85 (2012))
- ✓ **GSI-Heidelberg**: equilibrium statistical model with T_{chem} =156 MeV (**A. Andronic et al., Phys. Lett. B** 697, 203 (2011))
- SHARE: non-equilibrium thermal model with T_{chem} =138.3 MeV (M. Petráň et al., Phys. Rev. C 88, 034907 (2013))
- Great sensitivity to theoretical models parameters
- Non–equilibrium statistical thermal model (Petráň-Rafelski SHARE) provides better global fitting ($\chi^2 \sim 1$) to lower mass hadrons but **misses** $^3_{\Lambda}$ **H** and light nuclei
- \triangleright Experimental data closest to equilibrium thermal model with $T_{\rm chem}$ = 156 MeV and to Hybrid UrQMD



HYPERTRITON LIFETIME DETERMINATION



New preliminary results at $\sqrt{s_{NN}}$ = 5.02 TeV



$$c\tau = \left(5.4^{+1.6}_{-1.2}(stat.) \pm 1.00(syst.)\right)cm$$
$$\tau = \left(181^{+54}_{-39}(stat.) \pm 33(syst.)\right)ps$$

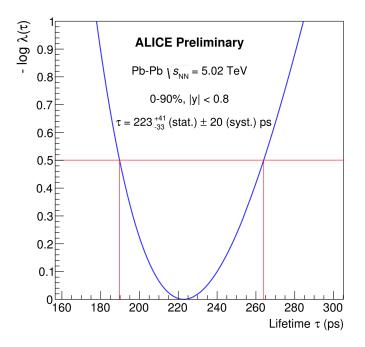
$$c\tau = \left(7.10^{+1.00}_{-1.07}(stat.) \pm 0.50(syst.)\right)cm$$
$$\tau = \left(237^{+33}_{-36}(stat.) \pm 17(syst.)\right)ps$$



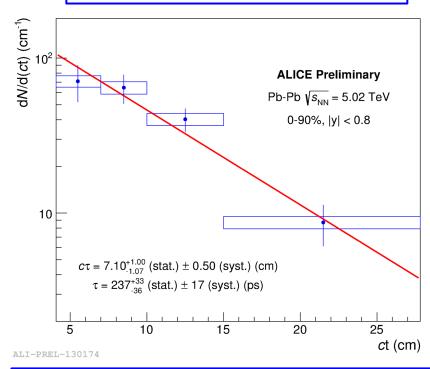
HYPERTRITON LIFETIME DETERMINATION

Two methods for estimation:

- ct spectra fit (exponential fit to the differential yield in different ct bins)
- ct unbinned fit as crosscheck method



New preliminary results at $\sqrt{s_{NN}}$ = 5.02 TeV

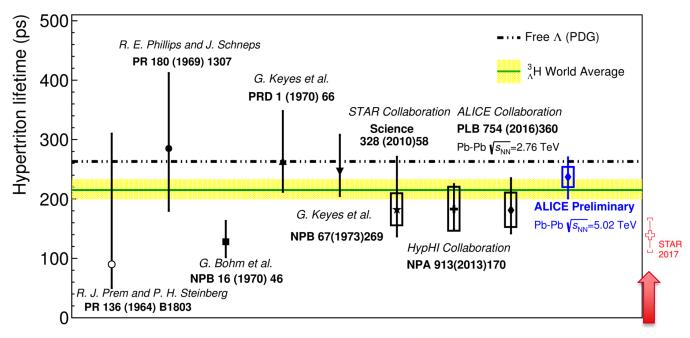


$$c\tau = \left(7.10^{+1.00}_{-1.07}(stat.) \pm 0.50(syst.)\right)cm$$
$$\tau = \left(237^{+33}_{-36}(stat.) \pm 17(syst.)\right)ps$$

ALI-PREL-130191



HYPERTRITON LIFETIME WORLD AVERAGE



ALI-PREL-130195

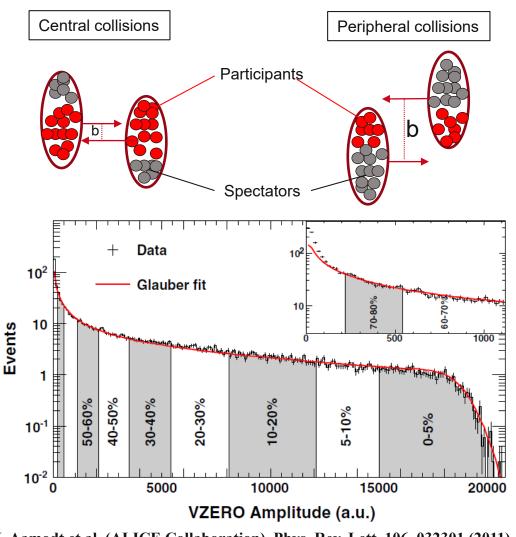
STAR Collaboration, arXiv:1710.00436v1 [nucl-ex]

$$\tau = \left(142^{+24}_{-21}(stat.) \pm 31(syst.)\right) ps$$

- Previous heavy-ion experiment results show a trend well below the free Λ lifetime
- ALICE preliminary result from Pb-Pb at 5.02 TeV is closer to the free Λ lifetime.
- STAR result from Au-Au collision is about 50% shorter than the free Λ lifetime
- The puzzle of the ³ H lifetime is still open

ALICE

COLLISION GEOMETRY



K. Aamodt et al. (ALICE Collaboration), Phys. Rev. Lett. 106, 032301 (2011)

- Nuclei are extended objects
- Geometry not directly measurable
- Centrality (percentage of the total cross section of the nuclear collision) connected to observables via Glauber model
- Data classified into centrality percentiles for which the average impact parameter, number of participants, and number of binary collisions can be determined