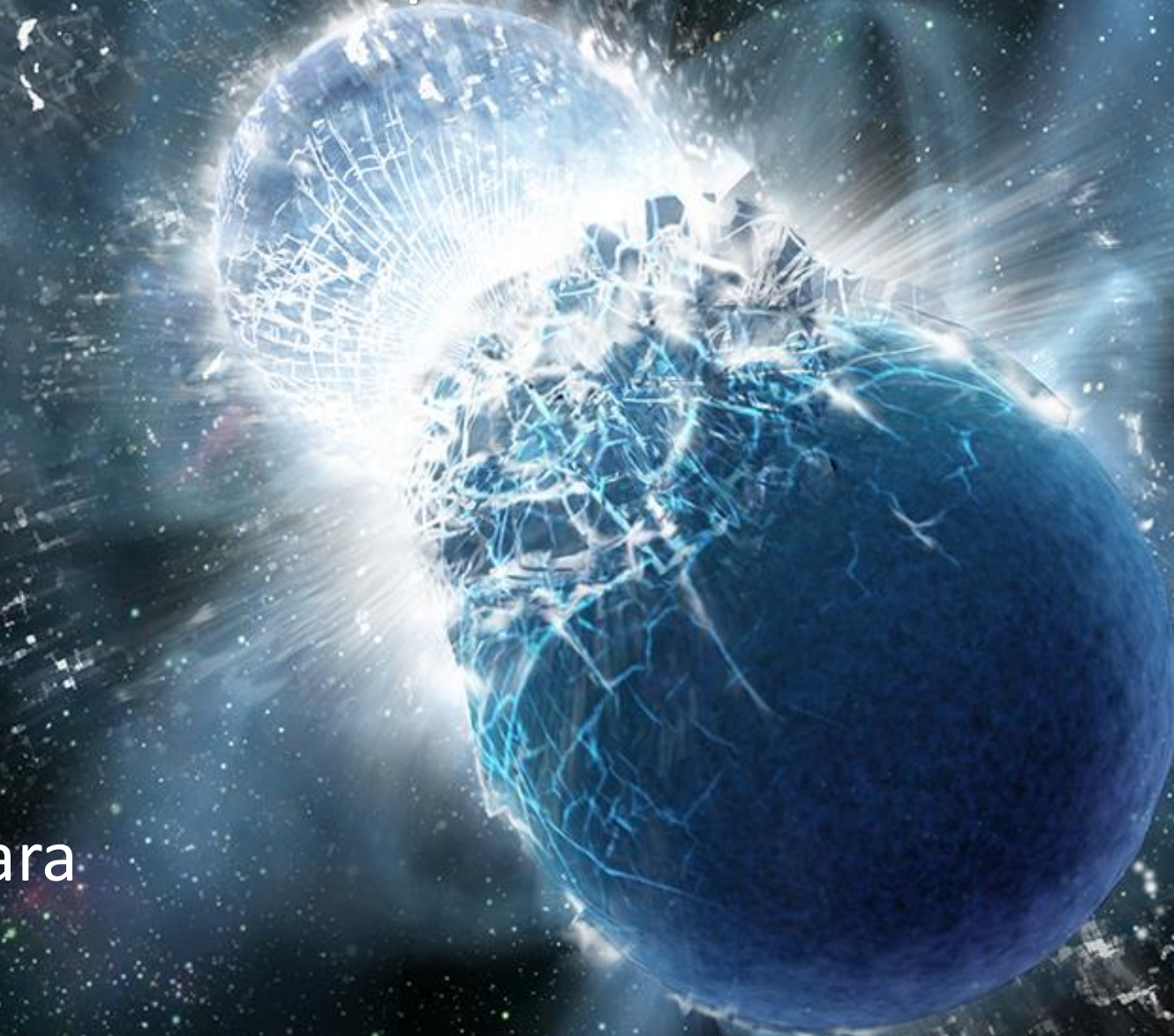


Matter ejection from binary neutron stars merger with two different Equations of State

Traversi Silvia
University of Ferrara

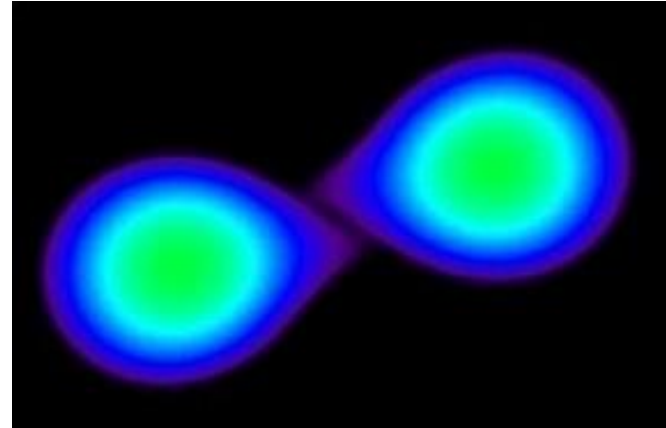


EJECTION MECHANISMS

EJECTION MECHANISMS

➤ **Dynamical ejection:**

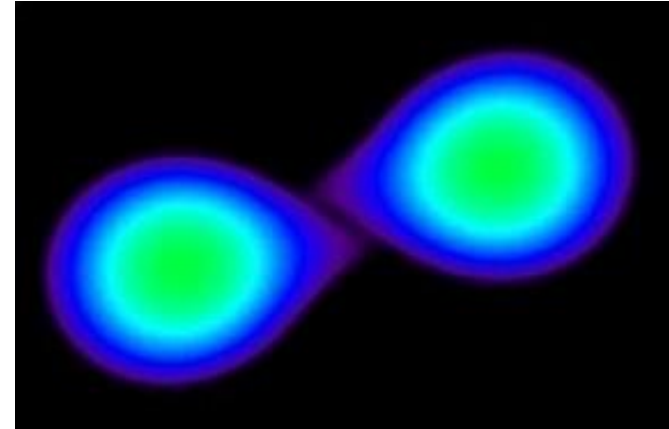
❖ Tidal deformation: equatorial plane



EJECTION MECHANISMS

➤ **Dynamical ejection:**

❖ Tidal deformation: equatorial plane

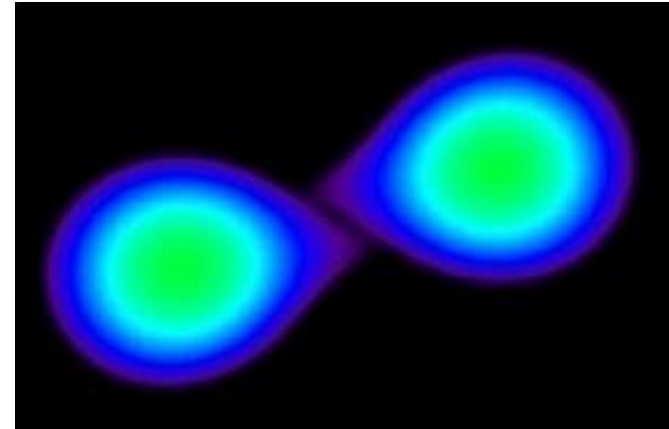


❖ Shock at NSs interface and radial oscillations

EJECTION MECHANISMS

➤ Dynamical ejection:

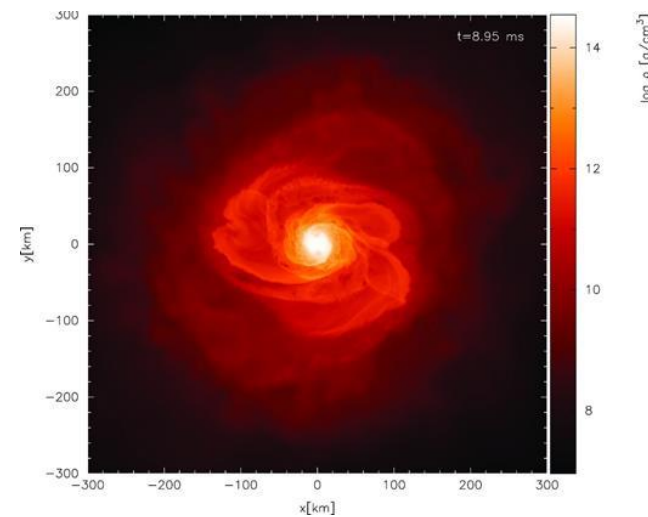
❖ Tidal deformation: equatorial plane

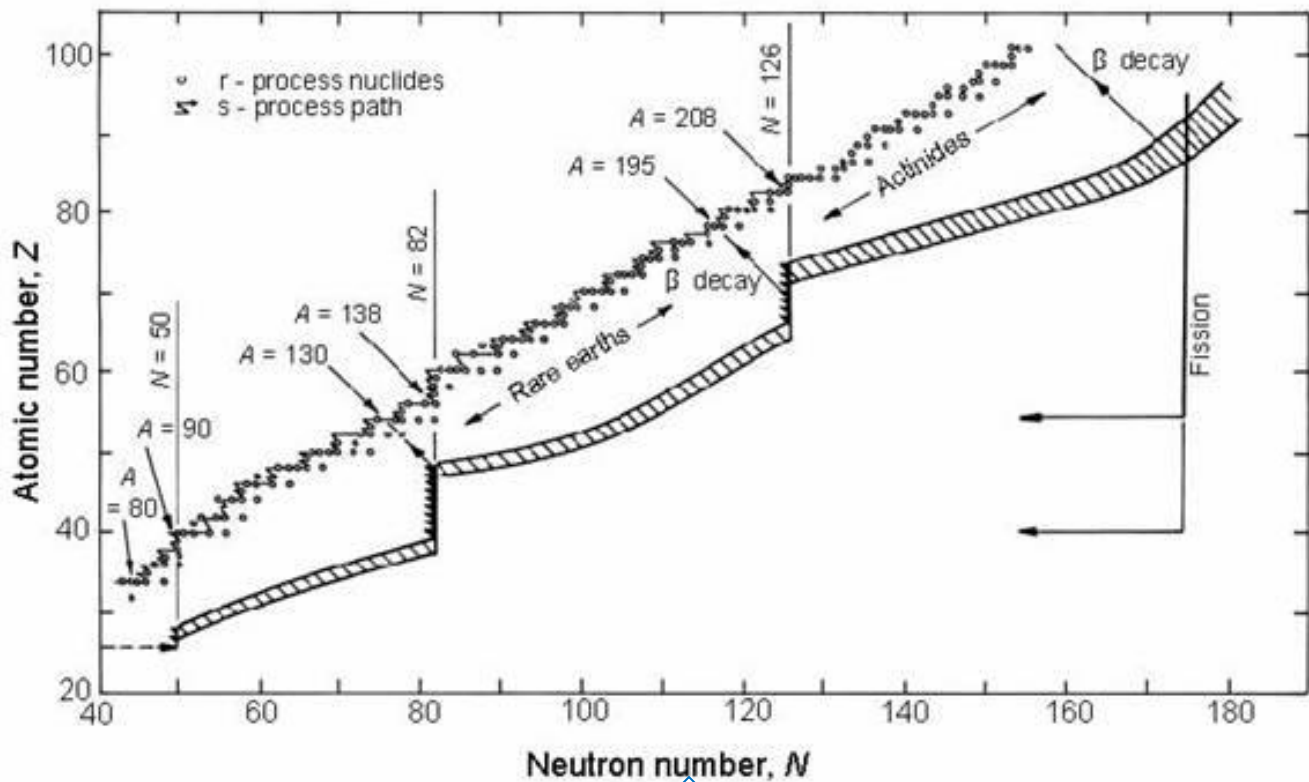


❖ Shock at NSs interface and radial oscillations

➤ Disk: $10^{-3} M_{\odot} < M_{disk} < 0.03 M_{\odot}$

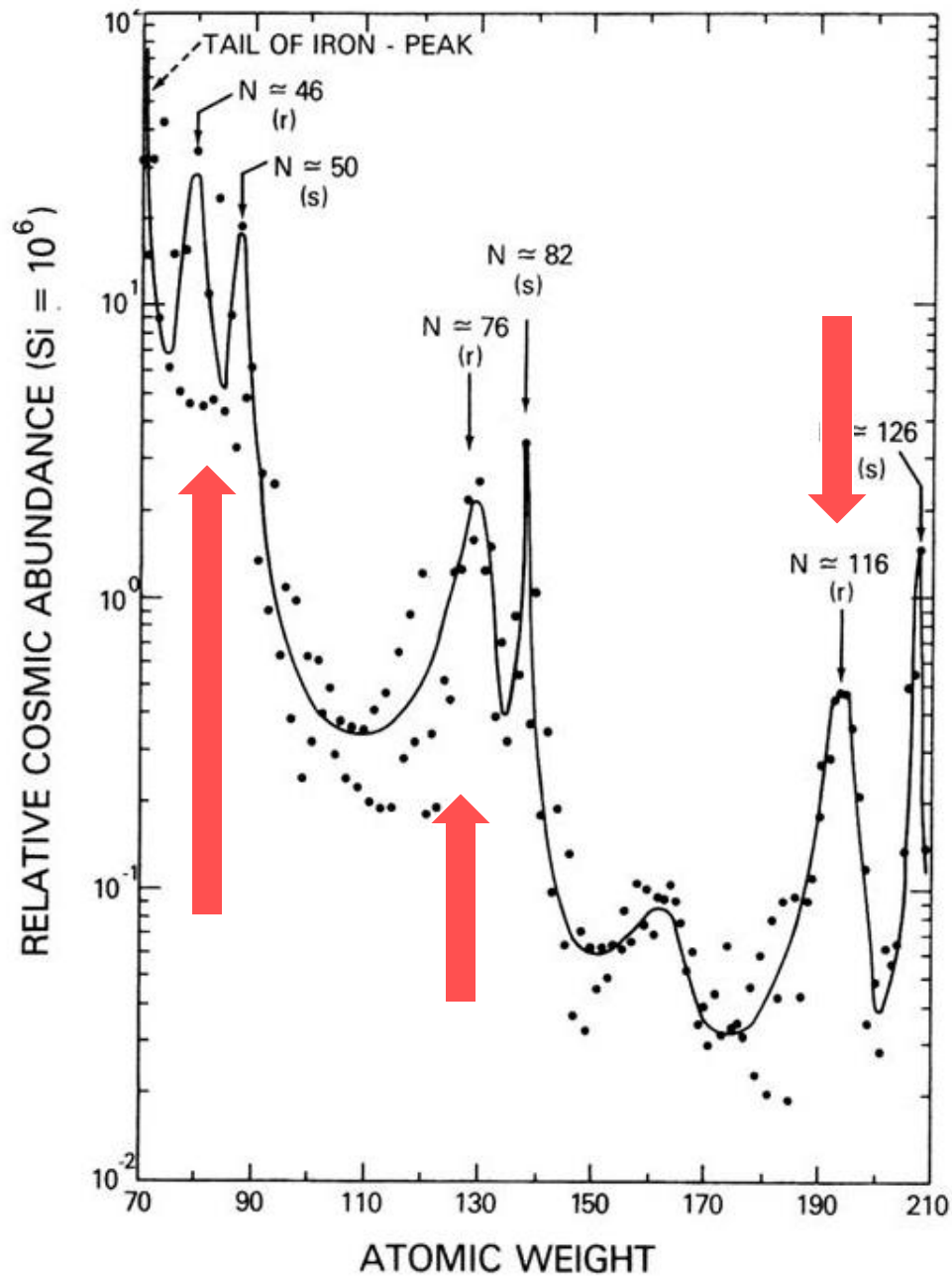
❖ Viscous or neutrino heating





R-process path

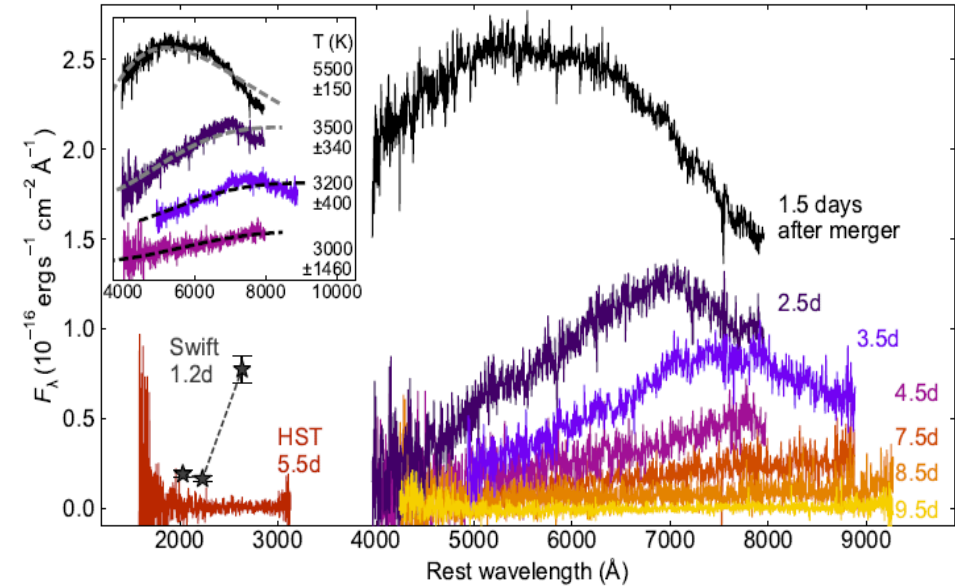
Heavy elements peaks



The kilonova signal

$$t_{peak} \propto \left(\frac{k M_{ej}}{v} \right)^{1/2} \quad L_{peak} \propto \left(\frac{v M_{ej}}{k} \right)^{1/2}$$

$$T_{peak} \propto (v M_{ej})^{-1/8} k^{-3/8}$$

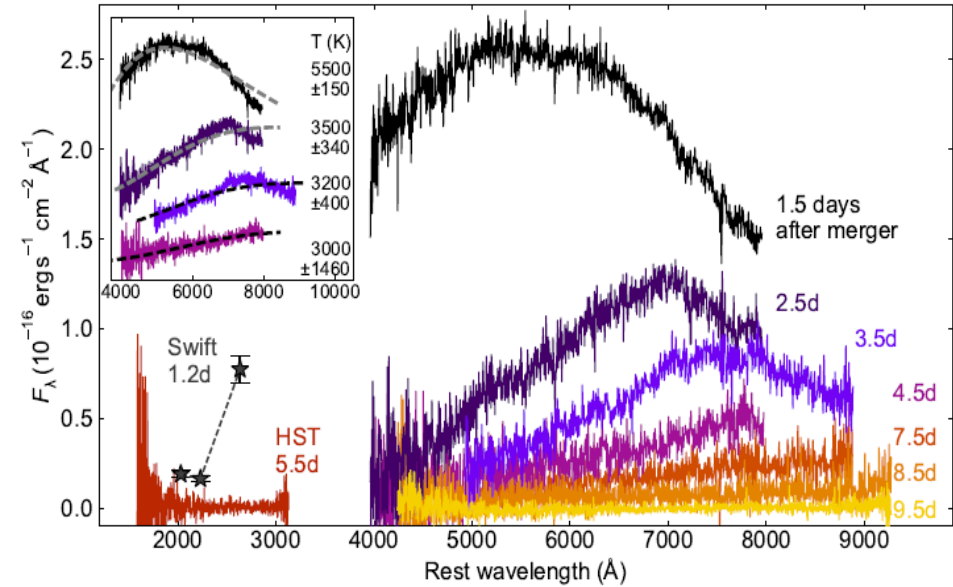


Nicholl et al. [arXiv:astro-ph.HE/1710.05456]

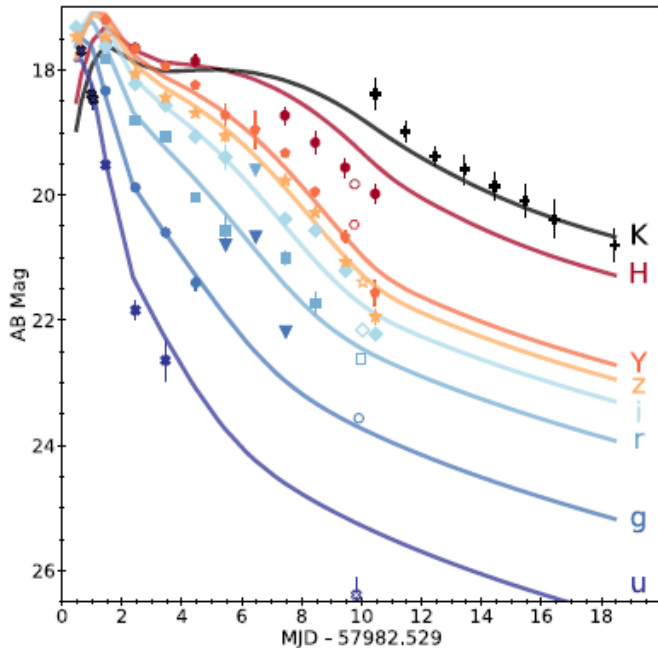
The kilonova signal

$$t_{peak} \propto \left(\frac{k M_{ej}}{v} \right)^{1/2} \quad L_{peak} \propto \left(\frac{v M_{ej}}{k} \right)^{1/2}$$

$$T_{peak} \propto (v M_{ej})^{-1/8} k^{-3/8}$$



Nicholl et al. [arXiv:astro-ph.HE/1710.05456]



Cowperthwaite, P.S.; others.
[arXiv:astro-ph.HE/1710.05840].

Blue Kilonova

$$M_{ej}^B \sim 0.01 M_{\odot}, v_{ej}^B = 0.27c, k^B = 0.5 \text{ cm}^2 \text{ s}^{-1}$$

Purple Kilonova

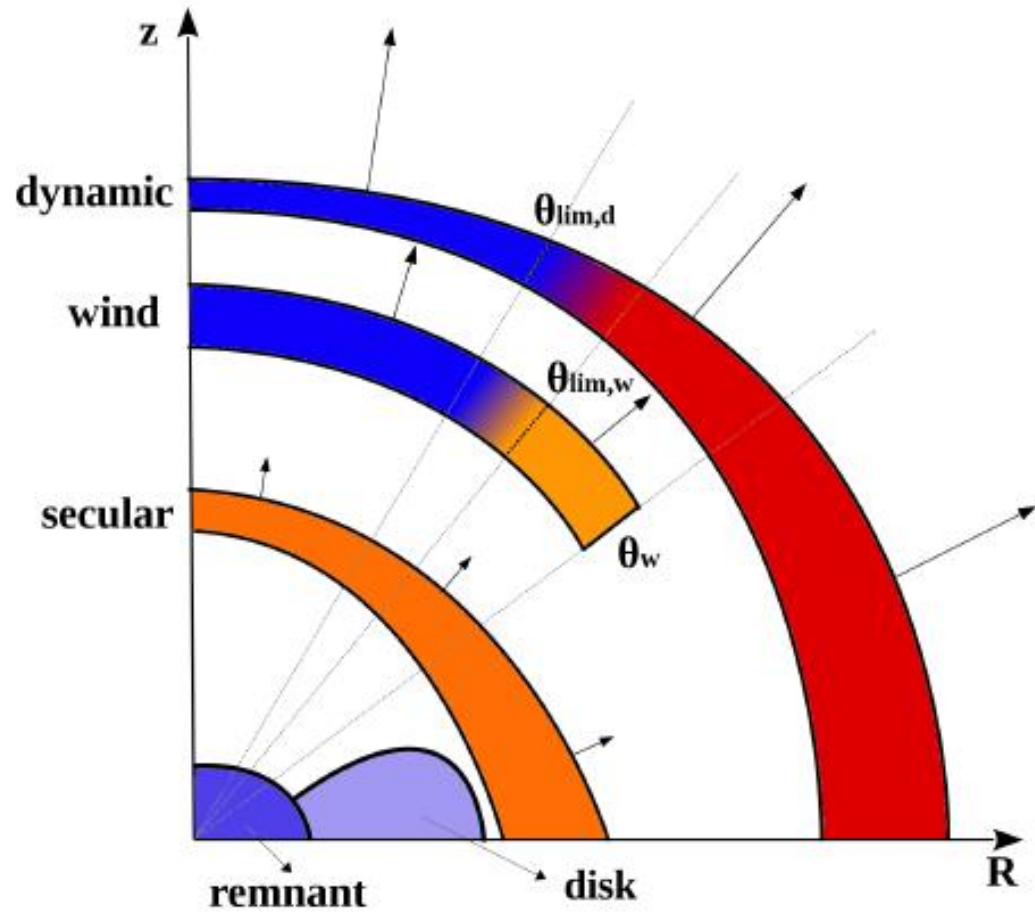
$$M_{ej}^P \sim 0.03 M_{\odot}, v_{ej}^P = 0.11c, k^P = 3 \text{ cm}^2 \text{ s}^{-1}$$

Red Kilonova

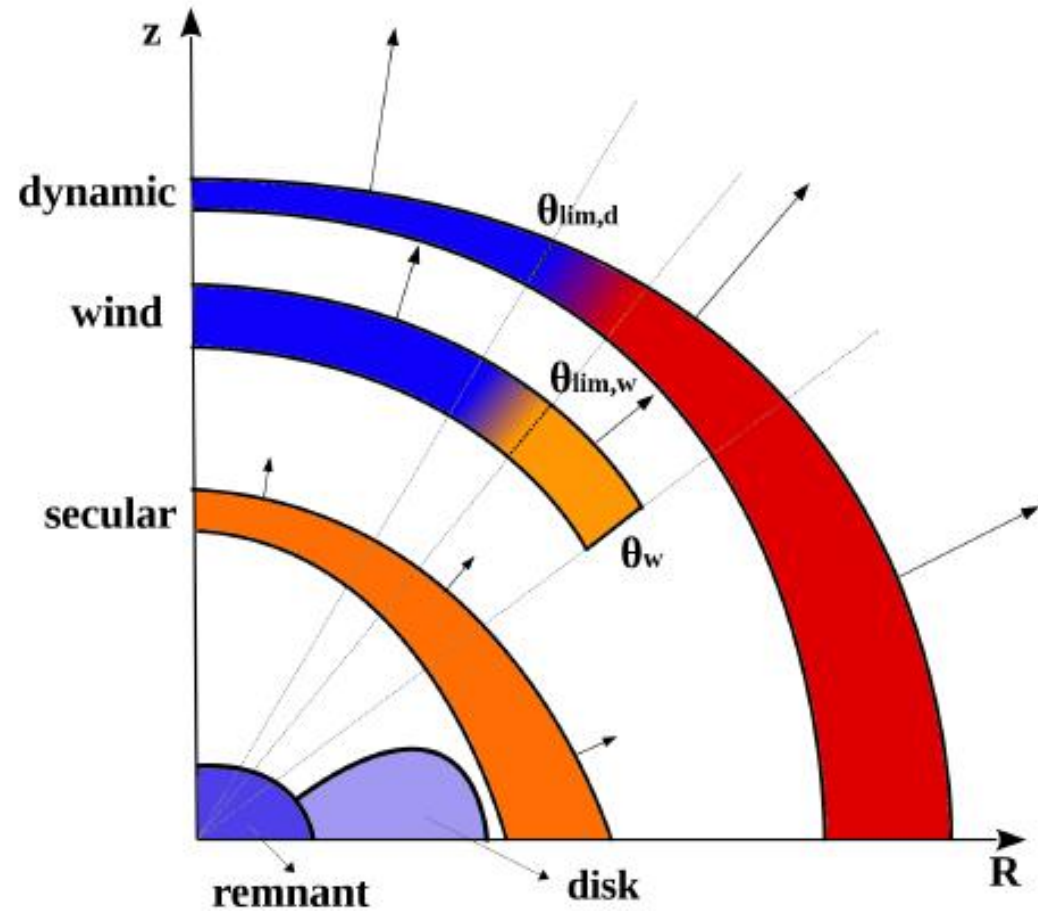
$$M_{ej}^R \sim 0.01 M_{\odot}, v_{ej}^R = 0.16c, k^R = 10 \text{ cm}^2 \text{ s}^{-1}$$

The electron fraction

- ❖ Shock: $Y_e > 0.25 - 0.3$
Lanthanide poor



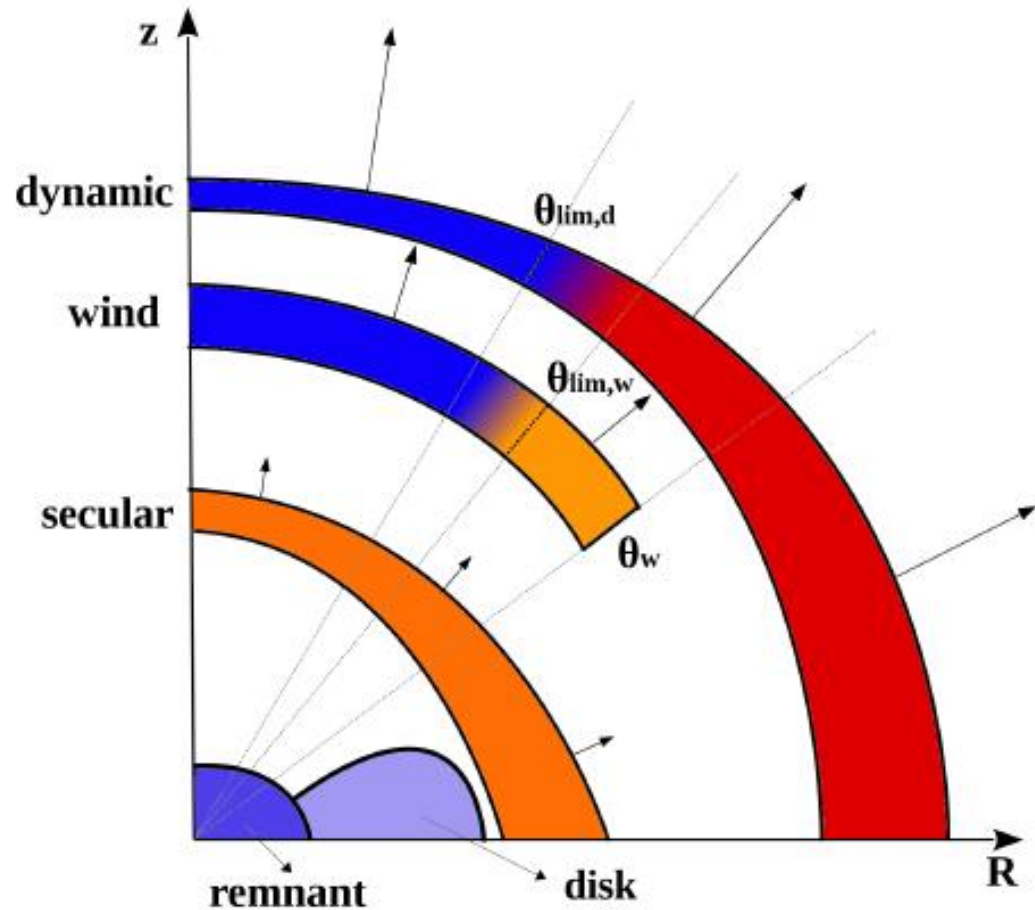
The electron fraction



❖ Shock: $Y_e > 0.25 - 0.3$
Lanthanide poor

❖ Tidal ejection: $Y_e < 0.1$
Lanthanide rich

The electron fraction



❖ Shock: $Y_e > 0.25 - 0.3$
Lanthanide poor

❖ Tidal ejection: $Y_e < 0.1$
Lanthanide rich

❖ Wind in polar direction: $Y_e = 0.3 - 0.4$
Lanthanide poor

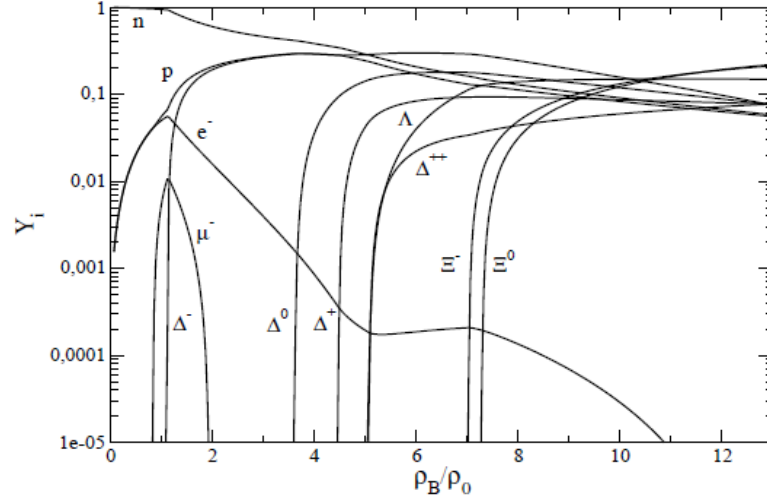
❖ Wind and secular: $Y_e \sim 0.2$
Depends on the lifetime of the HMNS

The two families scenario

1. Hadronic stars: $M_{max} \sim 1.6M_{\odot}$, $R \sim 10km$

2. Quark stars: $M_{max} \sim 2M_{\odot}$, $R > 11km$

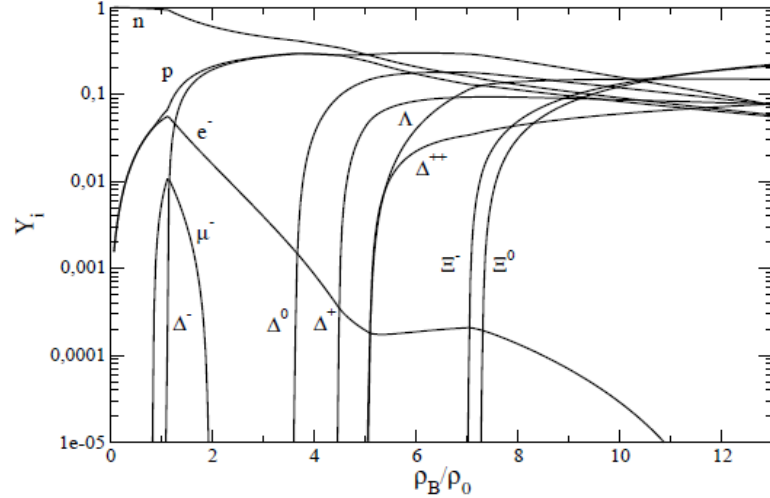
The two families scenario



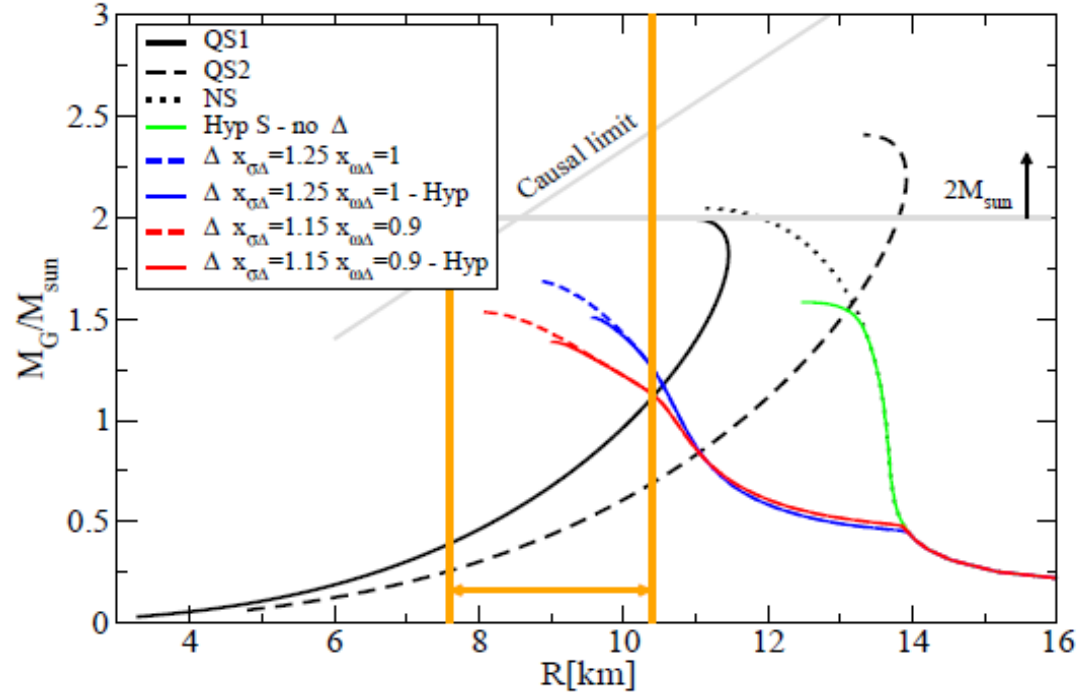
Appearance of hyperons and Δ -resonances
Softening of the EOS

1. Hadronic stars: $M_{max} \sim 1.6M_{\odot}$, $R \sim 10km$
2. Quark stars: $M_{max} \sim 2M_{\odot}$, $R > 11km$

The two families scenario

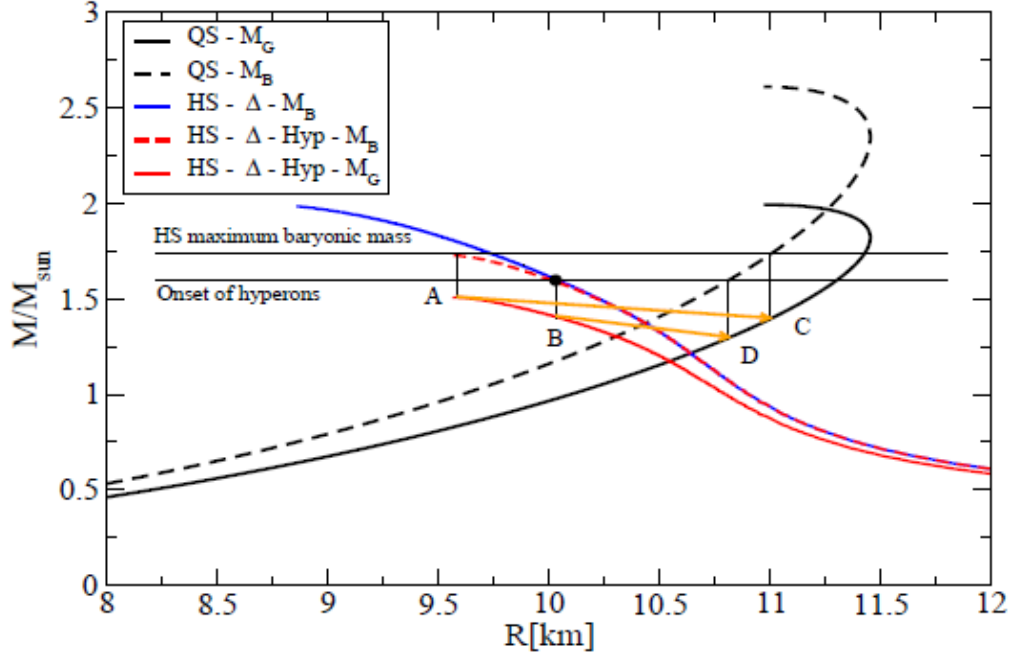
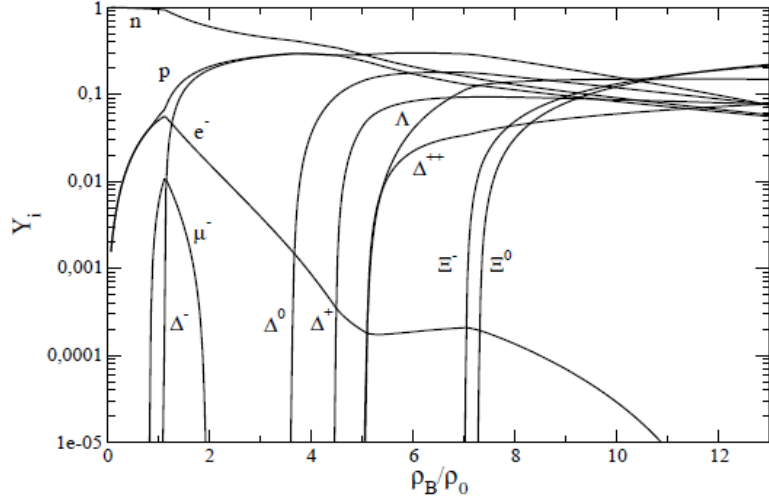


Appearance of hyperons and Δ -resonances
Softening of the EOS

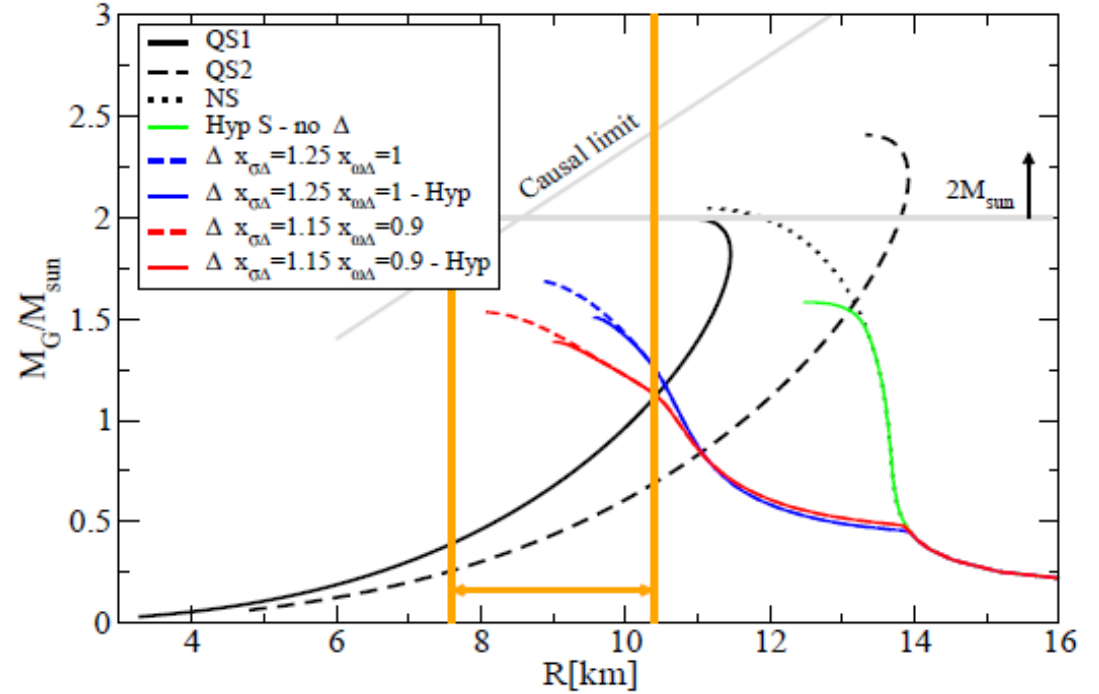


1. Hadronic stars: $M_{max} \sim 1.6M_{\odot}$, $R \sim 10km$
2. Quark stars: $M_{max} \sim 2M_{\odot}$, $R > 11km$

The two families scenario



Appearance of hyperons and Δ -resonances
Softening of the EOS



1. Hadronic stars: $M_{\text{max}} \sim 1.6M_{\odot}$, $R \sim 10\text{km}$

2. Quark stars: $M_{\text{max}} \sim 2M_{\odot}$, $R > 11\text{km}$

The framework of the simulations

Two Equations of state

SFHo

N, e^\pm , γ , nuclei

β - equilibrium condition

Hadronic 2 families EOS

= SFHo for low densities

Hyperons and Δ in the nucleus

The framework of the simulations

Two Equations of state

SFHo

N, e^\pm , γ , nuclei

β - equilibrium condition

Hadronic 2 families EOS

= SFHo for low densities

Hyperons and Δ in the nucleus

- ❖ Masses of the symmetric binaries from $1.15 - 1.15 M_\odot$ to $1.3 - 1.3 M_\odot$
- ❖ Einstein Toolkit code: full General Relativistic framework

The framework of the simulations

Two Equations of state

SFHo

$N, e^{\pm}, \gamma, \text{nuclei}$

β - equilibrium condition

Hadronic 2 families EOS

= SFHo for low densities

Hyperons and Δ in the nucleus

- ❖ Masses of the symmetric binaries from $1.15 - 1.15 M_{\odot}$ to $1.3 - 1.3 M_{\odot}$
- ❖ Einstein Toolkit code: full General Relativistic framework

The threshold mass

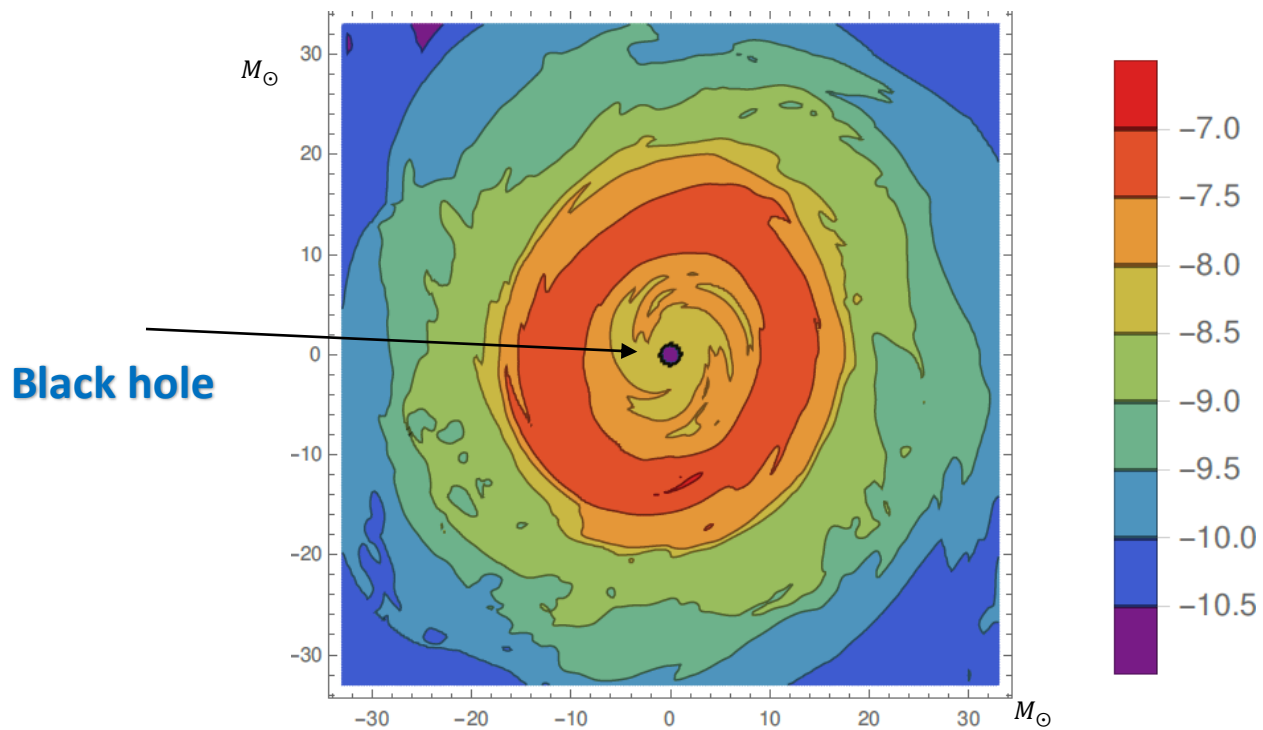
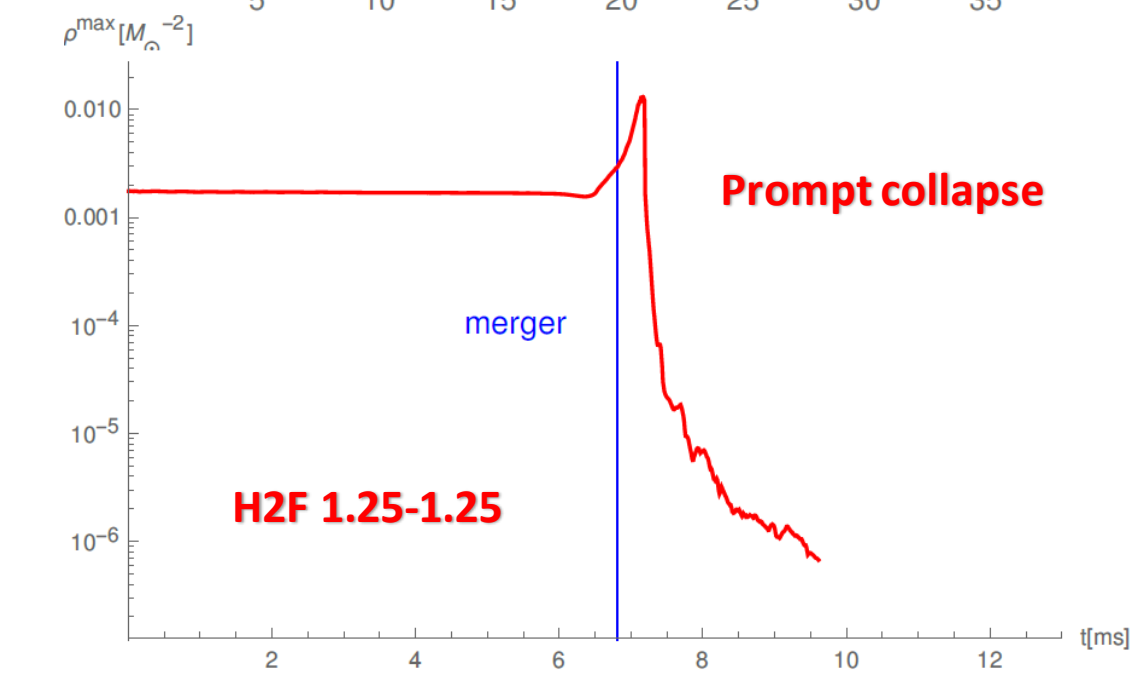
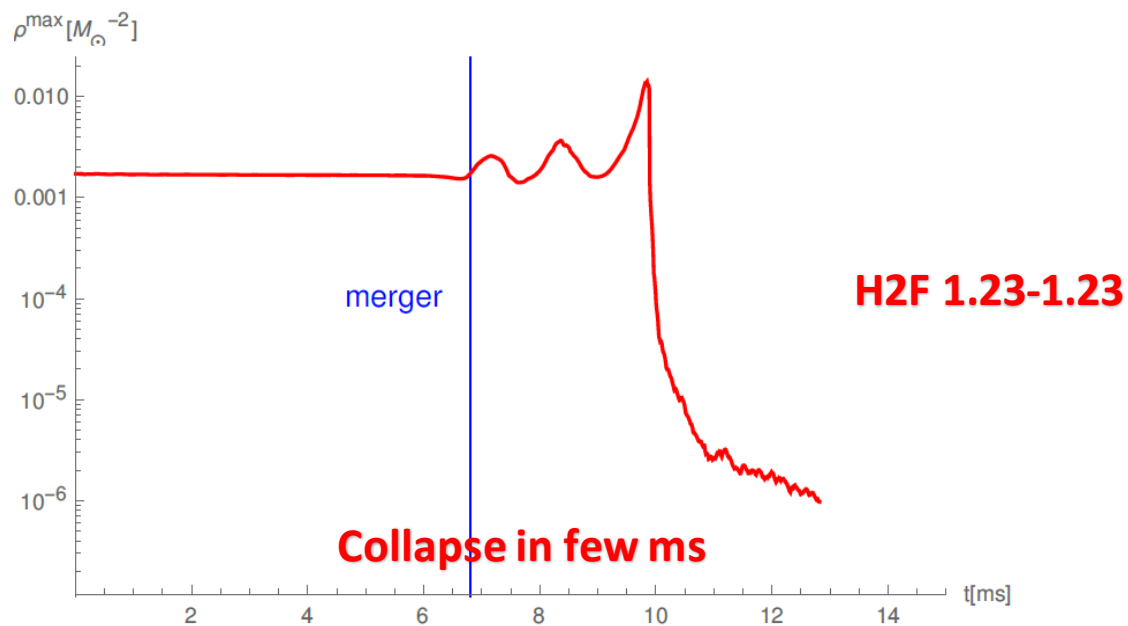
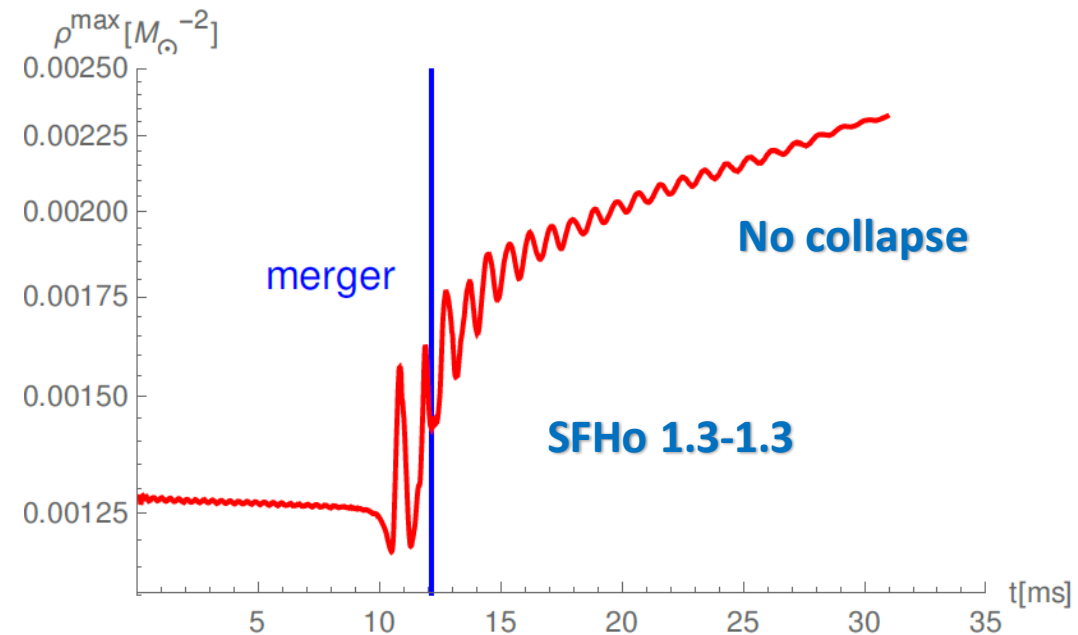
$$M_{thresh} = (1.54 \pm 0.05)M_{TOV}$$

$$M_{thresh} = (2.43 - 3.38C_{max})M_{TOV}$$

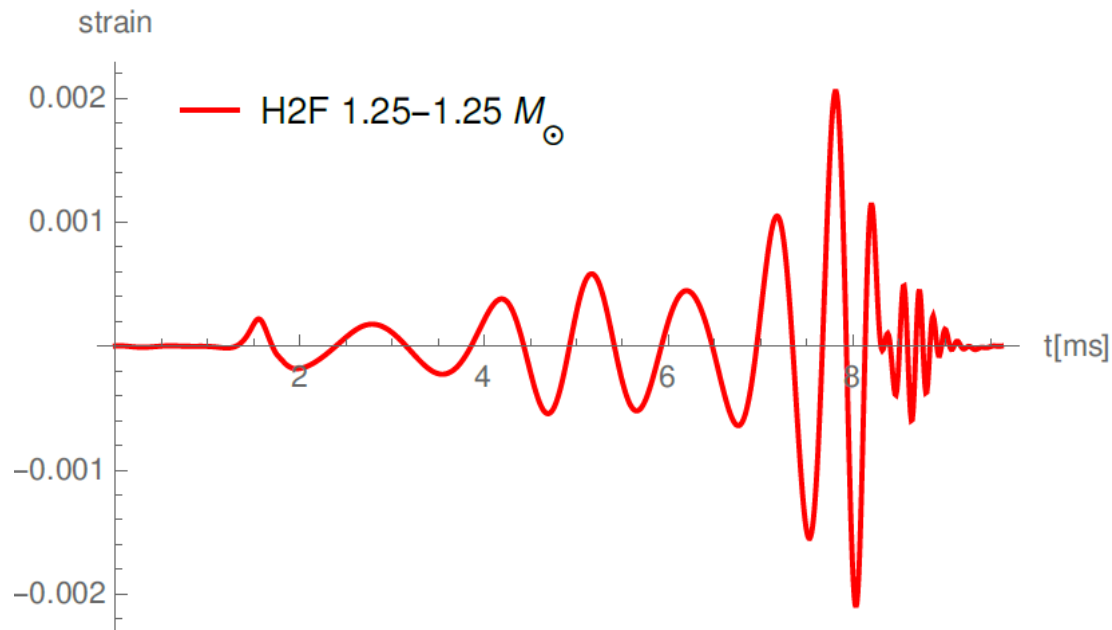
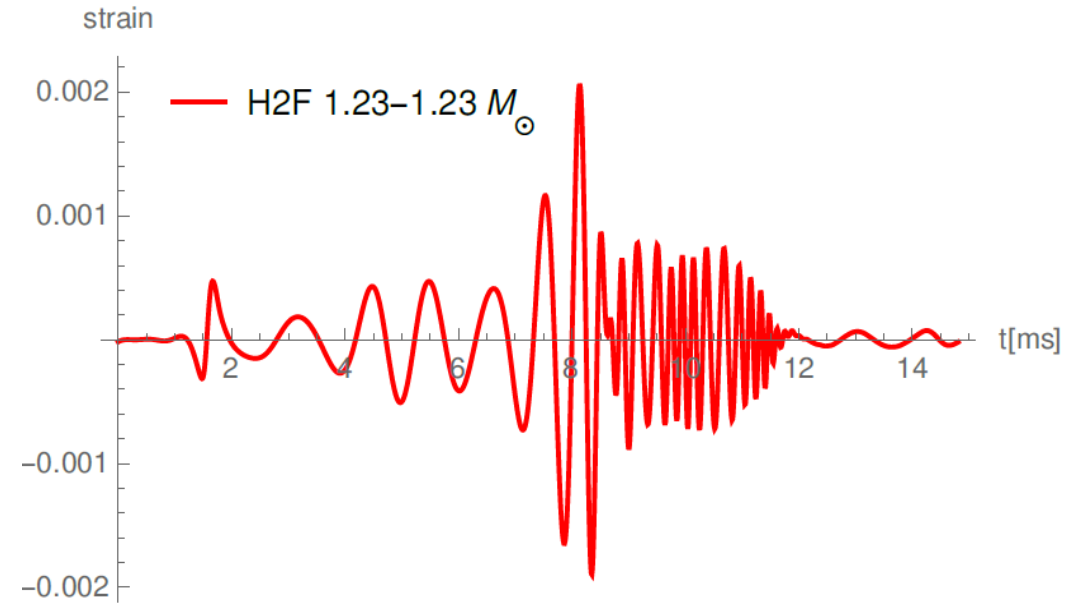
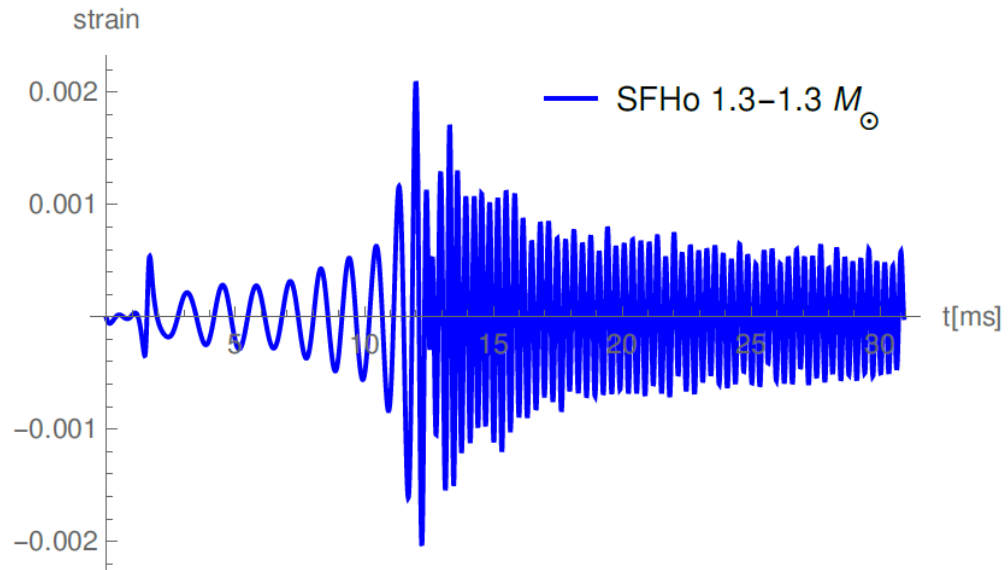
$$M_{thresh}^{SFHo} > 2.8M_{\odot}$$

$$M_{thresh}^{H2F} \sim 2.5 - 2.6M_{\odot}$$

Maximum density




Gravitational waves




$$M_{thresh}^{SFHo} > 2.6 M_{\odot}$$

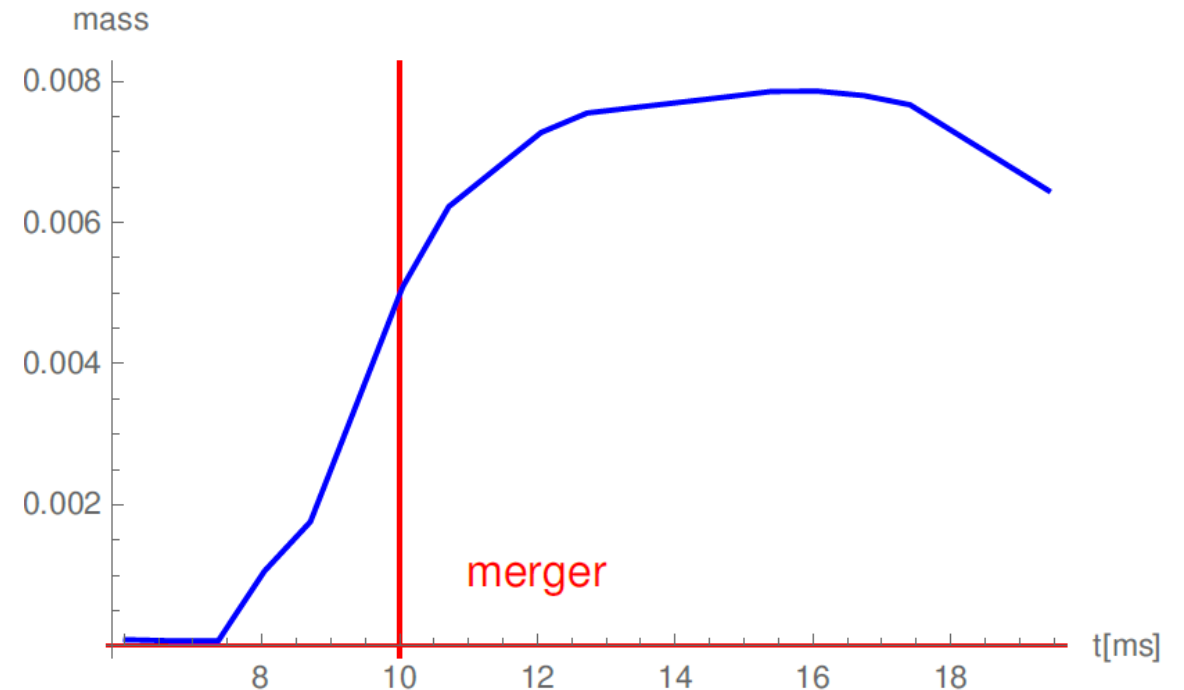
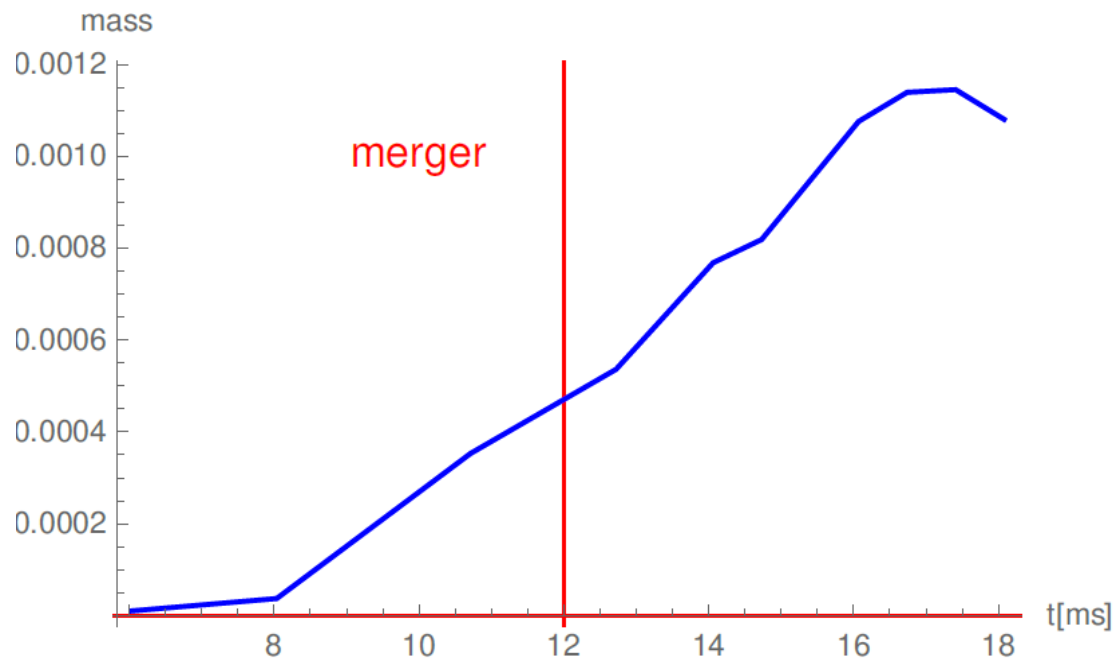
$$M_{thresh}^{H2F} \sim 2.48 M_{\odot}$$

Dynamical ejecta

- Calculate the **Lorentz factor γ** the **time component of the 4-velocity u_t** for fixed time and all cells
- Identify the cells containing unbound fluid  **Geodesic criterion**
 $-u_t > 1$
- Obtain **the mass** integrating on the cells
- Select the time with **maximum** unbound mass

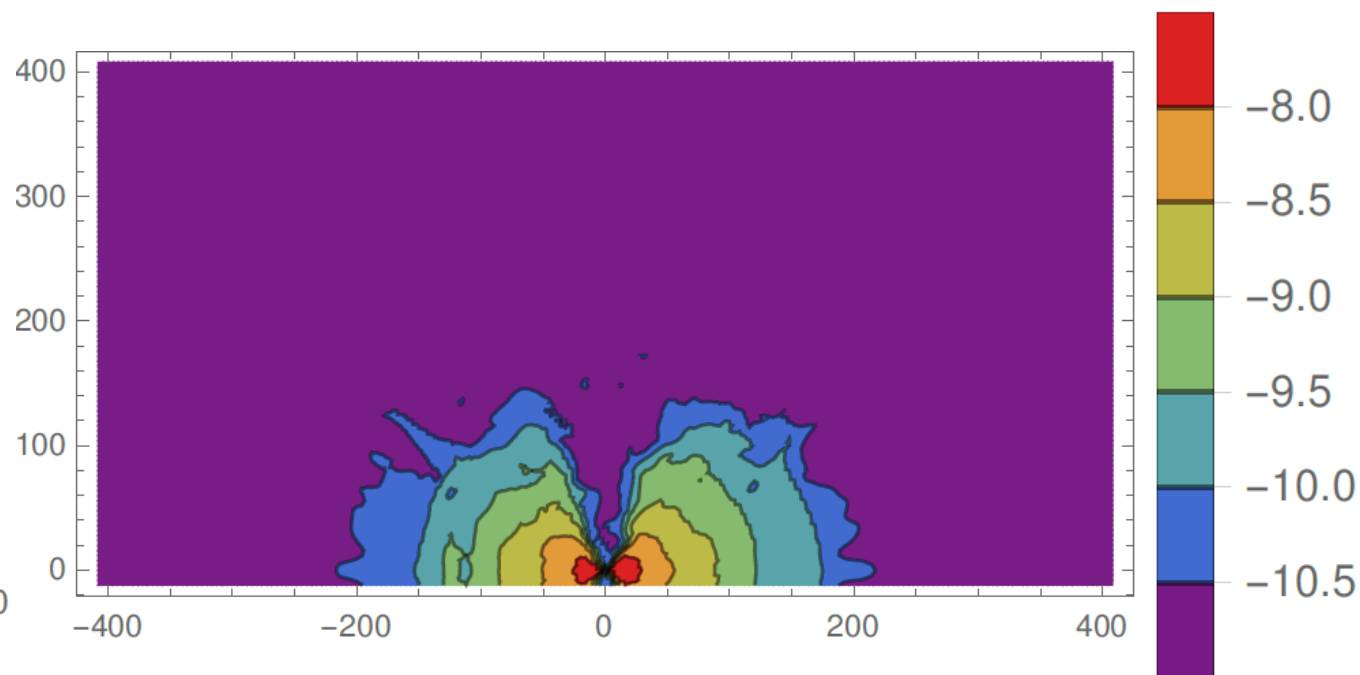
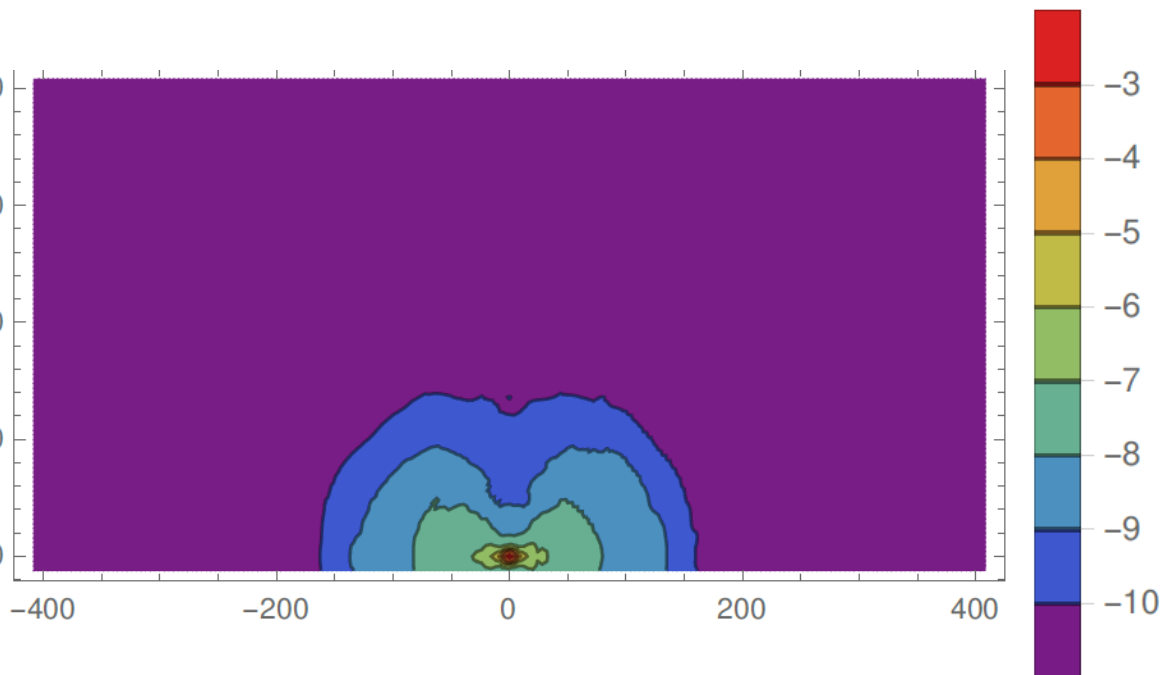
Dynamical ejecta

- Calculate the **Lorentz factor Γ** the **time component of the 4-velocity u_t** for fixed time and all cells
- Identify the cells containing unbound fluid  **Geodesic criterion**
 $-u_t > 1$
- Obtain the **mass** integrating on the cells
- Select the time with **maximum** unbound mass



Mass of the disk

- Consider the last time available (20 ms after merger): quasi stationary state
- High density cut at $10^{13} \text{ g cm}^{-3} \sim 10^{-5} M_{\odot}^{-2}$
- Low density cut at $5 * 10^6 \text{ g cm}^{-3} \sim 10^{-11} M_{\odot}^{-2}$ or $5 * 10^7 \text{ g cm}^{-3} \sim 10^{-10} M_{\odot}^{-2}$



Preliminary results

Simulation	collapse	dynamical ejecta		disk
	tc - tmdens	Mej	Vav a tmax	Mdisk (cut 10^{-11})
SFHo 1,2	no	0.00127	0.18	0.1006
H2F 1,2	4,39	0.0056	0.13	0.025
SFHo 1,2,3	no	0.001144	0.18	0.092
H2F 1,2,3	3,6	0.00786	0.15	0,0214
SFHo 1,3	no	0.0019	0.16	0.1230
H2F 1,3	0	10^{-5}	0.35	

Preliminary results

Simulation	collapse	dynamical ejecta		disk
	tc - tmdens	Mej	Vav a tmax	Mdisk (cut 10 ⁽⁻¹¹⁾)
SFHo 1,2	no	0.00127	0.18	0.1006
H2F 1,2	4,39	0.0056	0.13	0.025
SFHo 1,2,3	no	0.001144	0.18	0.092
H2F 1,2,3	3,6	0.00786	0.15	0,0214
SFHo 1,3	no	0.0019	0.16	0.1230
H2F 1,3	0	10 ⁽⁻⁵⁾	0.35	

❖ Dynamical ejecta: SFHo $\sim 10^{-3} M_{\odot}$, H2F $\sim 5 * 10^{-3} - 10^{-2} M_{\odot}$

- For softer EOS the **shock component** is more relevant
- For configurations near the threshold mass, the **radial oscillations** can eject matter

Preliminary results

Simulation	collapse	dynamical ejecta		disk
	tc - tmdens	Mej	Vav a tmax	Mdisk (cut 10^{-11})
SFHo 1,2	no	0.00127	0.18	0.1006
H2F 1,2	4,39	0.0056	0.13	0.025
SFHo 1,2,3	no	0.001144	0.18	0.092
H2F 1,2,3	3,6	0.00786	0.15	0,0214
SFHo 1,3	no	0.0019	0.16	0.1230
H2F 1,3	0	10^{-5}	0.35	

❖ Dynamical ejecta: SFHo $\sim 10^{-3} M_{\odot}$, H2F $\sim 5 * 10^{-3} - 10^{-2} M_{\odot}$

- For softer EOS the **shock component** is more relevant
- For configurations near the threshold mass, the **radial oscillations** can eject matter

❖ Disk: SFHo $\sim 10^{-1} M_{\odot}$, H2F $\sim 10^{-2} M_{\odot}$

- For stiff EOS the **tidal deformability** is bigger

Interpretation of GW170817

Total mass of about $2.72M_{\odot}$: in the two families scenario, the event cannot be a NS-NS merger because it would lead to prompt collapse

Interpretation of GW170817

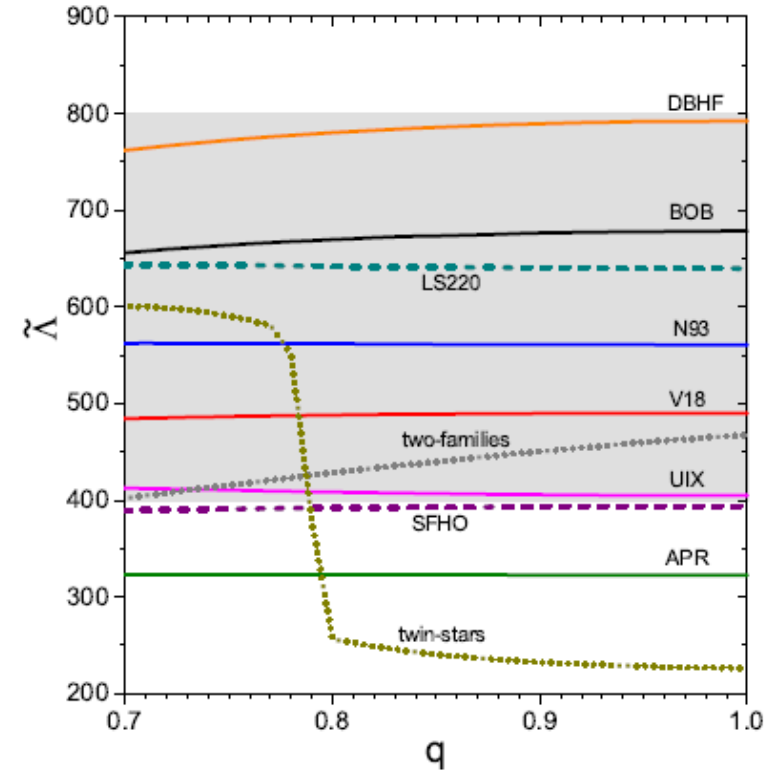
Total mass of about $2.72M_{\odot}$: in the two families scenario, the event cannot be a NS-NS merger because it would lead to prompt collapse



Hadronic star – Quark star merger?

The effective tidal deformability is inside the lower and upper limits imposed by GW170817

For the future: we would like to perform simulations with the quarks EOS



Interpretation of GW170817

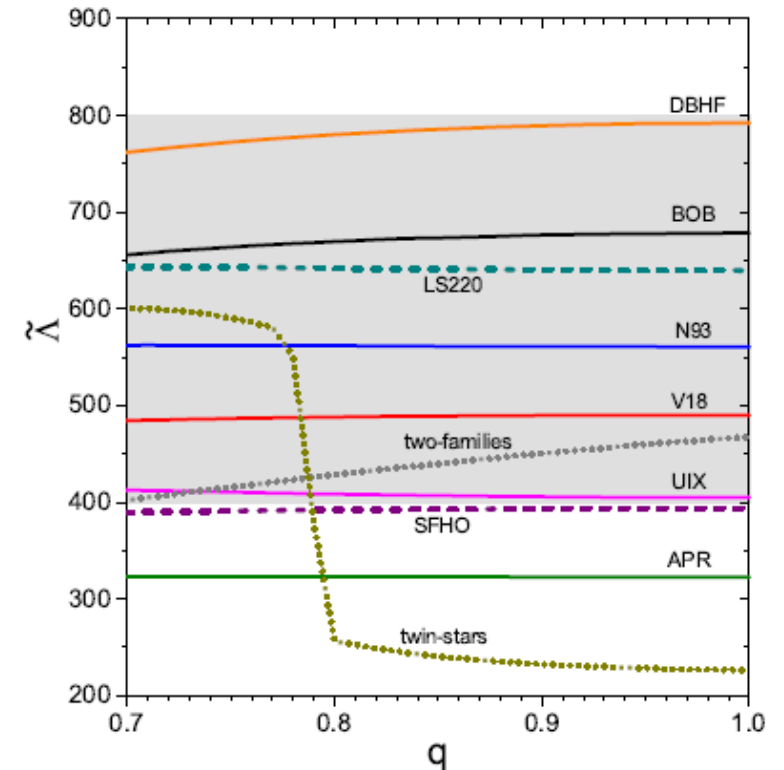
Total mass of about $2.72M_{\odot}$: in the two families scenario, the event cannot be a NS-NS merger because it would lead to prompt collapse



Hadronic star – Quark star merger?

The effective tidal deformability is inside the lower and upper limits imposed by GW170817

For the future: we would like to perform simulations with the quarks EOS



Thanks for the attention