

Generalized Parton Distributions at Jefferson Lab

Carlos Muñoz Camacho

Institut de Physique Nucléaire, CNRS/IN2P3 (France)

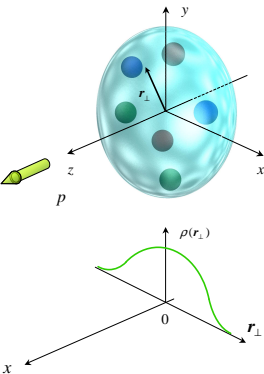
2018 European Nuclear Physics Conference
Bologne (Italy)

Outline

- ① Introduction
- ② Nucleon 3D-imaging & Generalized Parton Distributions (GPDs)
- ③ Deeply Virtual Compton Scattering (DVCS): $ep \rightarrow ep\gamma$
 - Results on both proton and neutron (preliminary)
- ④ Exclusive π^0 electroproduction (DVMP): $eN \rightarrow eN\pi^0$
 - Also: proton + neutron \Rightarrow flavor separation
- ⑤ Plans at 12 GeV
- ⑥ Summary

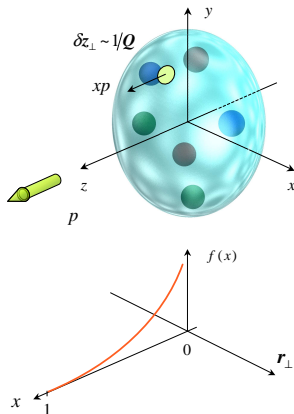
Studying the structure of the nucleon experimentally

Elastic scattering



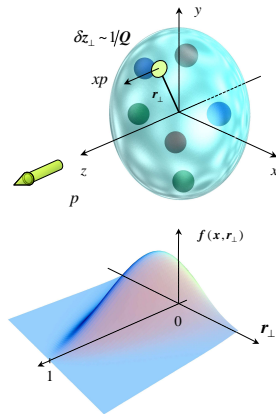
Form factors

Deeply Inelastic Scattering



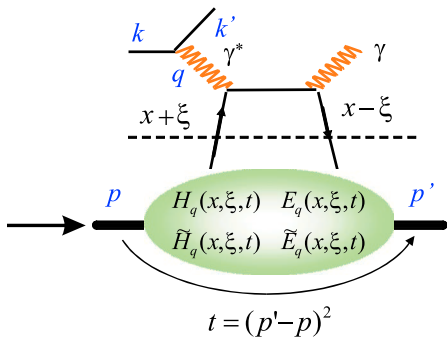
Parton distributions

Hard exclusive processes



Generalized Parton Distributions (GPDs)

Deeply Virtual Compton Scattering (DVCS): $\gamma^* p \rightarrow \gamma p$



High Q^2
Perturbative QCD

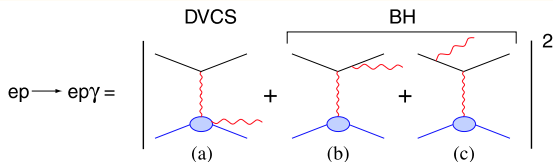
Non-perturbative
GPDs

Handbag diagram

Bjorken limit :

$$Q^2 = \left. \begin{array}{l} -q^2 \rightarrow \infty \\ \nu \rightarrow \infty \end{array} \right\} x_B = \frac{Q^2}{2M\nu} \text{ fixed}$$

DVCS experimentally: interference with Bethe-Heitler



At leading order in $1/Q$ (leading twist) :

$$d^5 \vec{\sigma} - d^5 \overleftarrow{\sigma} = \Im (T^{BH} \cdot T^{DVCS})$$

$$d^5 \vec{\sigma} + d^5 \overleftarrow{\sigma} = |BH|^2 + \Re (T^{BH} \cdot T^{DVCS}) + |DVCS|^2$$

$$\mathcal{T}^{DVCS} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\epsilon} + \dots =$$

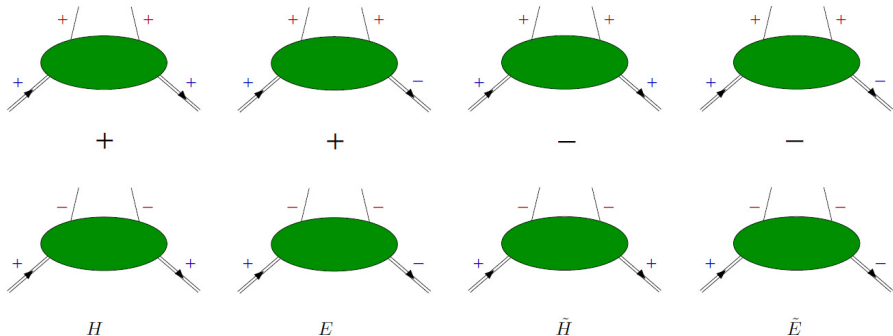
$$\underbrace{\mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi}}_{\text{Access in helicity-independent cross section}} - \underbrace{i\pi H(x = \xi, \xi, t)}_{\text{Access in helicity-dependent cross-section}} + \dots$$

Access in helicity-independent cross section

Access in helicity-dependent cross-section

Leading twist GPDs

8 GPDs related to the different combination of quark/nucleon helicities

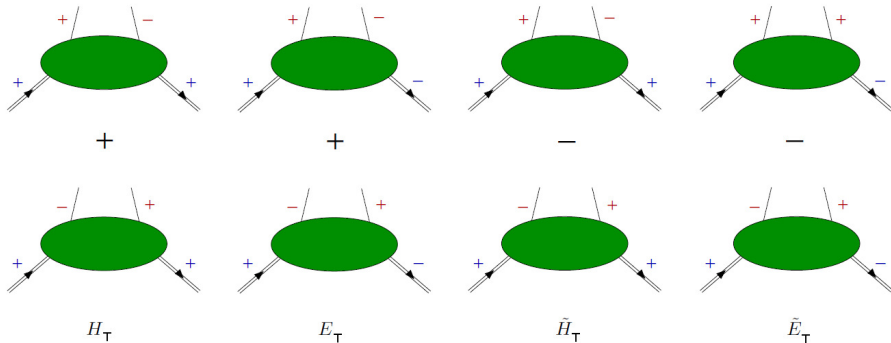


4 chiral-even GPDs: conserve the helicity of the quark

Access through DVCS (and DVMP)

Leading twist GPDs

8 GPDs related to the different combination of quark/nucleon helicities



4 chiral-odd GPDs: flip helicity of the quark

“transversity GPDs”

Experimental access more complicated (π^0 electroproduction?)

Accessing different GDPs

Polarized beam, unpolarized target (BSA)

$$d\sigma_{LU} = \sin \phi \cdot \mathcal{I}m\{F_1 \mathcal{H} + x_B(F_1 + F_2)\tilde{\mathcal{H}} - kF_2\mathcal{E}\}d\phi$$

Unpolarized beam, longitudinal target (ITSA)

$$d\sigma_{UL} = \sin \phi \cdot \mathcal{I}m\{F_1 \tilde{\mathcal{H}} + x_B(F_1 + F_2)(\tilde{\mathcal{H}} + x_B/2\mathcal{E}) - x_B kF_2 \tilde{\mathcal{E}} \dots\}d\phi$$

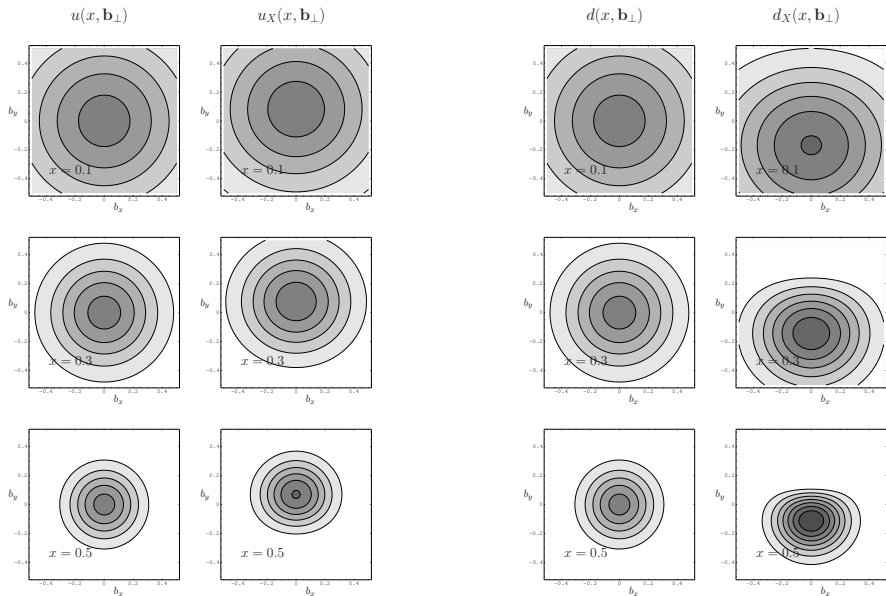
Polarized beam, longitudinal target (BITSA)

$$d\sigma_{LL} = (A + B \cos \phi) \cdot \mathcal{R}e\{F_1 \tilde{\mathcal{H}} + x_B(F_1 + F_2)(\tilde{\mathcal{H}} + x_B/2\mathcal{E}) \dots\}d\phi$$

Unpolarized beam, transverse target (tTSA)

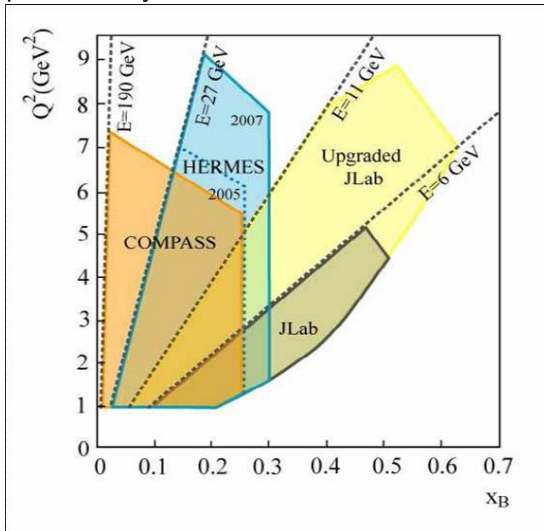
$$d\sigma_{UT} = \cos \phi \cdot \mathcal{I}m\{k(F_2 \mathcal{H} - F_1 \mathcal{E}) + \dots\}d\phi$$

Impact-parameter interpretation of GPDs



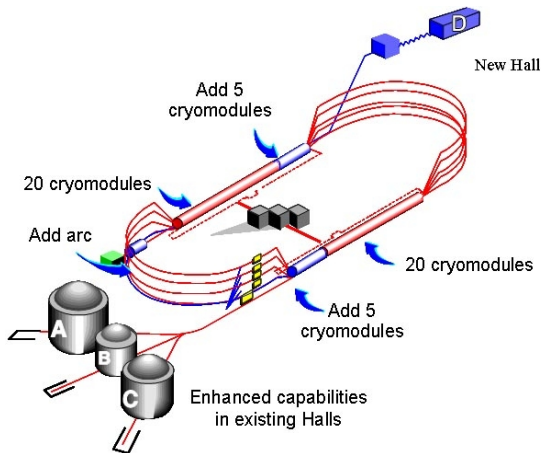
Kinematic coverage

Kinematic complementarity between different facilities:



Jefferson Lab: upgraded to 12 GeV

- 6-12 GeV longitudinally polarized ($>85\%$) continuous electron beam
- High intensity ($>100\ \mu\text{A}$): luminosities $> 10^{38}\ \text{s}^{-1}\ \text{cm}^{-2}$
- 3 experimental Halls (A, B, C) w/ fixed target and dedicated detectors



The GPD experimental program at Jefferson Lab

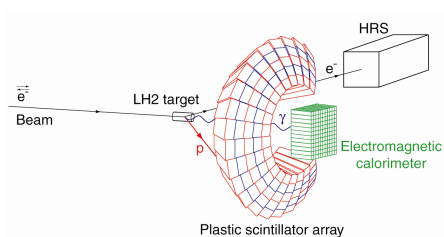
- Hall A: high accuracy, limited kinematic coverage
- Hall B: wide kinematic range, limited precision
- Hall C: high precision program at 11 GeV

Partially overlapping, partially complementary programs
with different experimental setups

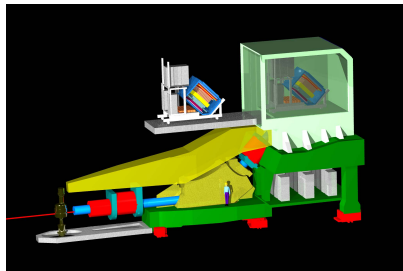
The roadmap:

- Early results (2001) from non-dedicated experiment (CLAS)
- 1st round of dedicated experiments in Halls A/B in 2004/5
- 2nd round on 2008–2010: precision tests + more spin observables
- Compelling DVCS experiments in Halls A+B+C at 11 GeV ($\gtrsim 2016$)

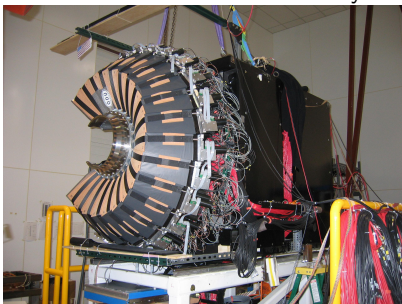
Experimental setup



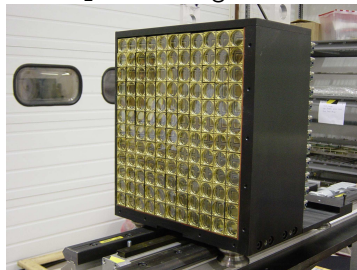
High Resolution Spectrometer



100-channel scintillator array

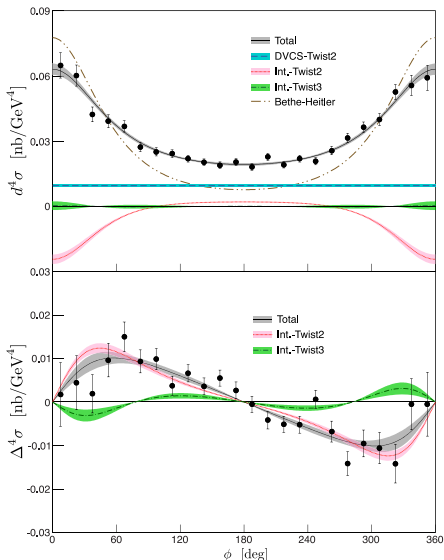


132-block PbF_2 electromagnetic calorimeter



DVCS cross sections: azimuthal analysis

$$Q^2 = 2.36 \text{ GeV}^2, x_B = 0.37, -t = 0.32 \text{ GeV}^2$$



$$ep \rightarrow ep\gamma = \left| \text{DVCS} + \text{BH} \right|^2$$

The diagram illustrates the Feynman diagrams for the process $ep \rightarrow ep\gamma$. It shows the DVCS (Deeply Virtual Compton Scattering) contribution and the BH (Bethe-Heitler) contribution. The DVCS part is represented by a single diagram (a) where the photon is emitted from the quark line. The BH part is represented by two diagrams (b) and (c) where the photon is emitted from the electron line. The diagrams are labeled (a), (b), and (c) respectively.

$$d^4\sigma = \mathcal{T}_{\text{BH}}^2 + \mathcal{T}_{\text{BH}} \text{Re}(\mathcal{T}_{\text{DVCS}}) + \mathcal{T}_{\text{DVCS}}^2$$

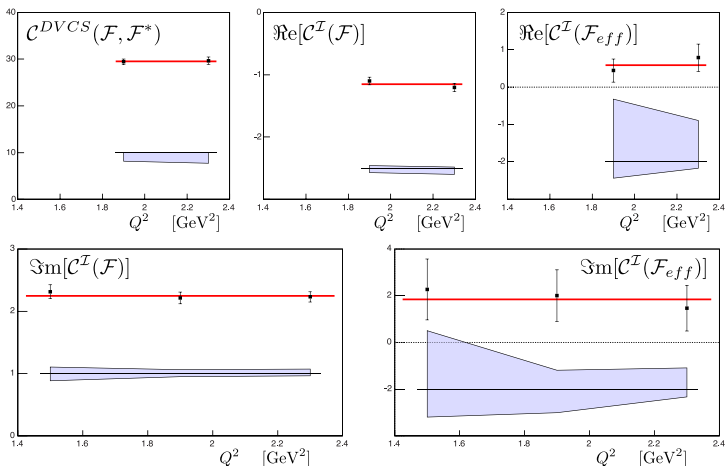
$$\text{Re}(\mathcal{T}_{\text{DVCS}}) \sim c_0^{\mathcal{I}} + c_1^{\mathcal{I}} \cos \phi + c_2^{\mathcal{I}} \cos 2\phi$$

$$\mathcal{T}_{\text{DVCS}}^2 \sim c_0^{\text{DVCS}} + c_1^{\text{DVCS}} \cos \phi$$

$$\Delta^4\sigma = \frac{d^4\vec{\sigma} - d^4\overleftarrow{\sigma}}{2} = \text{Im}(\mathcal{T}_{\text{DVCS}})$$

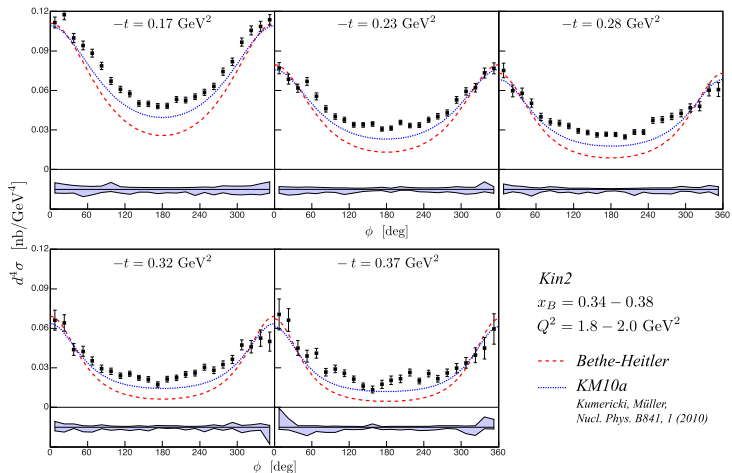
$$\text{Im}(\mathcal{T}_{\text{DVCS}}) \sim s_1^{\mathcal{I}} \sin \phi + s_2^{\mathcal{I}} \sin 2\phi$$

M. Defurne *et al.* Phys. Rev. C 92, 055202

DVCS cross sections: Q^2 -dependance

No Q^2 -dependance within limited range \Rightarrow leading twist dominance

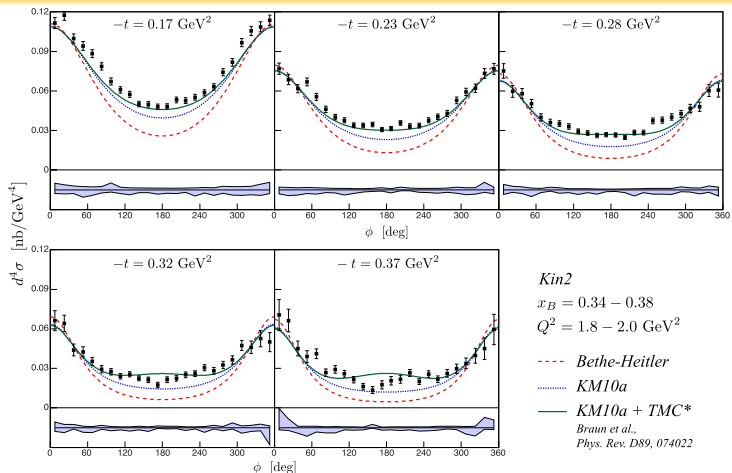
DVCS cross sections: kinematical power corrections



- KM10a: global fit to HERA x-sec & HERMES + CLAS spin asymmetries

Kumericki and Mueller (2010)

DVCS cross sections: kinematical power corrections



- KM10a: global fit to HERA x-sec & HERMES + CLAS spin asymmetries
Kumericki and Mueller (2010)
- Target-mass corrections (TMC): $\sim \mathcal{O}(M^2/Q^2)$ and $\sim \mathcal{O}(t/Q^2)$
Braun, Manashov, Mueller and Pirnay (2014)

Rosenbluth-like separation of the DVCS cross section

$$\sigma(ep \rightarrow ep\gamma) = \underbrace{|BH|^2}_{\text{Known to } \sim 1\%} + \underbrace{\mathcal{I}(BH \cdot DVCS)}_{\text{Linear combination of GPDs}} + \underbrace{|DVCS|^2}_{\text{Bilinear combination of GPDs}}$$

$$\mathcal{I} \propto 1/y^3 = (k/\nu)^3,$$

$$|\mathcal{T}^{DVCS}|^2 \propto 1/y^2 = (k/\nu)^2$$

BKM-2010 – at leading twist \rightarrow 7 independent GPD terms:

$$\{\Re, \Im [\mathcal{C}^{\mathcal{I}}, \mathcal{C}^{\mathcal{I},V}, \mathcal{C}^{\mathcal{I},A}] (\mathcal{F})\}, \quad \text{and} \quad \mathcal{C}^{DVCS}(\mathcal{F}, \mathcal{F}_*).$$

φ -dependence provides 5 independent observables:

$$\sim 1, \sim \cos \varphi, \sim \sin \varphi, \sim \cos(2\varphi), \sim \sin(2\varphi)$$

The measurement of the cross section at **two or more beam energies** for exactly the **same Q^2 , x_B , t kinematics**, provides the additional information in order to extract all leading twist observables independently.

DVCS process: leading twist ambiguity

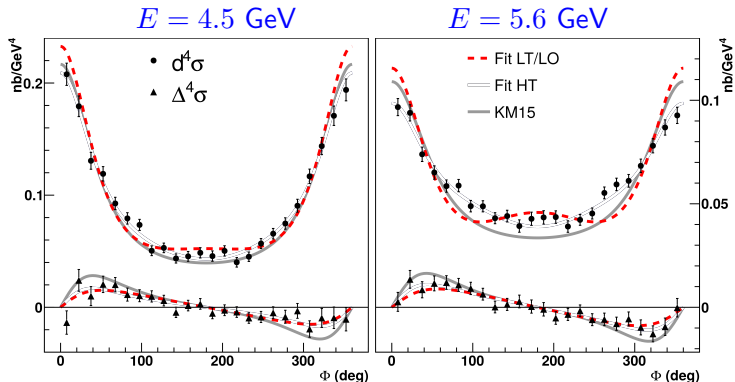
- DVCS defines a preferred axis: light-cone axis
- At finite Q^2 and non-zero t , there is an ambiguity:
 - ① Belitsky et al. (“BKM”, 2002–2010): light-cone axis in plane (q, P)
 - ② Braun et al. (“BMP”, 2014): light-cone axis in plane (q, q')
easier to account for kin. corrections $\sim \mathcal{O}(M^2/Q^2)$, $\sim \mathcal{O}(t/Q^2)$

$$\left. \begin{aligned} \mathcal{F}_{++} &= \mathbb{F}_{++} + \frac{\chi}{2} [\mathbb{F}_{++} + \mathbb{F}_{-+}] - \chi_0 \mathbb{F}_{0+} \\ \mathcal{F}_{-+} &= \mathbb{F}_{-+} + \frac{\chi}{2} [\mathbb{F}_{++} + \mathbb{F}_{-+}] - \chi_0 \mathbb{F}_{0+} \\ \mathcal{F}_{0+} &= -(1 + \chi) \mathbb{F}_{0+} + \chi_0 [\mathbb{F}_{++} + \mathbb{F}_{-+}] \end{aligned} \right\} \xrightarrow[\mathbb{F}_{0+} = 0]{\mathbb{F}_{-+} = 0} \left\{ \begin{aligned} \mathcal{F}_{++} &= (1 + \frac{\chi}{2}) \mathbb{F}_{++} \\ \mathcal{F}_{-+} &= \frac{\chi}{2} \mathbb{F}_{++} \\ \mathcal{F}_{0+} &= \chi_0 \mathbb{F}_{++} \end{aligned} \right.$$

(eg. $\chi_0 = 0.25$, $\chi = 0.06$ for $Q^2 = 2 \text{ GeV}^2$, $x_B = 0.36$, $t = -0.24 \text{ GeV}^2$)

E07-007: DVCS beam-energy dependence

- Cross section measured at 2 beam energies and constant Q^2 , x_B , t



- Leading-twist and LO simultaneous fit of both beam energies (dashed line) does not reproduce the data

Light-cone axis in the (q, q') plane (Braun et al.): \mathbb{H}_{++} , $\widetilde{\mathbb{H}}_{++}$, \mathbb{E}_{++} , $\widetilde{\mathbb{E}}_{++}$

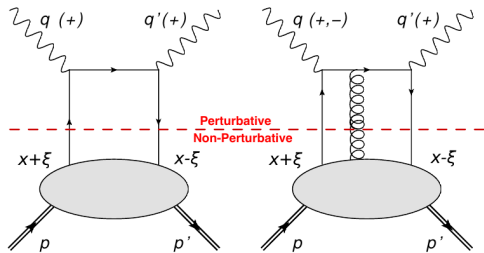
Beyond Leading Order (LO) and Leading Twist (LT)

Two fit-scenarios:

Light-cone axis in
the (q, q') plane (Braun et al.)

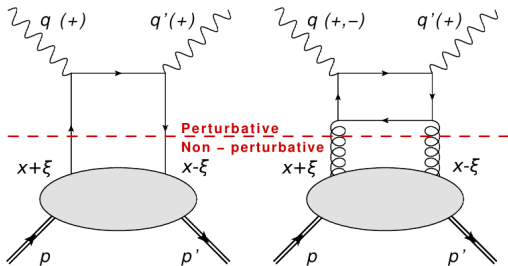
LO/LT + HT

$\mathbb{H}_{++}, \tilde{\mathbb{H}}_{++}, \mathbb{H}_{0+}, \tilde{\mathbb{H}}_{0+}$



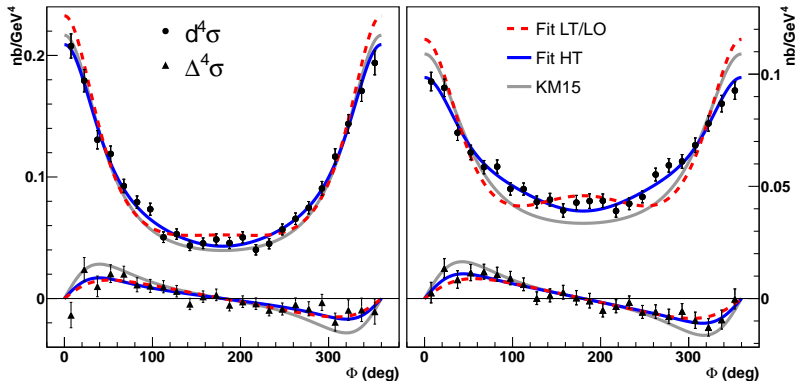
LO/LT + NLO

$\mathbb{H}_{++}, \tilde{\mathbb{H}}_{++}, \mathbb{H}_{-+}, \tilde{\mathbb{H}}_{-+}$



E07-007: DVCS beam-energy dependence

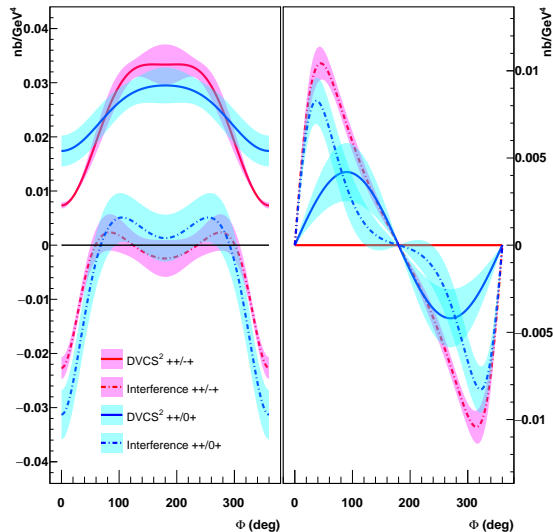
- Cross section measured at 2 beam energies and constant Q^2 , x_B , t



- Leading-twist and LO simultaneous fit of both beam energies (dashed line) does not reproduce the data
- Including either NLO or higher-twist effects (dark solid line) satisfactorily reproduce the angular dependence

DVCS² and \mathcal{I} (DVCS·BH) separation

DVCS² and \mathcal{I} (DVCS·BH) separated in NLO and higher-twist scenarios

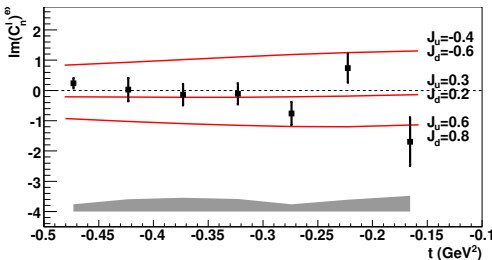


- DVCS² & \mathcal{I} significantly different in each scenario
- Sizeable DVCS² contribution in the higher-twist scenario in the helicity-dependent cross section

Nature Commun. 8, 1408 (2017)

DVCS on the neutron: experiment E03-106 at JLab

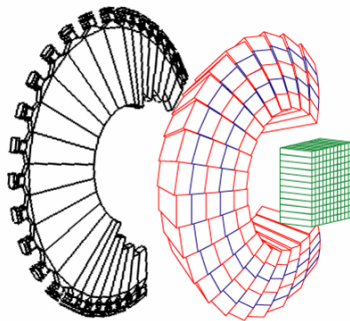
LD₂ target ($F_2^n(t) \gg F_1^n(t)$!)



$$\sigma^{\rightarrow} - \sigma^{\leftarrow} = \Gamma(A \sin \varphi + \dots)$$

$$A = F_1(t)\mathcal{H} + \frac{x_B}{2 - x_B}[F_1(t) + F_2(t)]\tilde{\mathcal{H}} - \underbrace{\frac{t}{4M^2} \cdot F_2(t) \cdot \mathcal{E}}_{\text{Main contribution for neutron}}$$

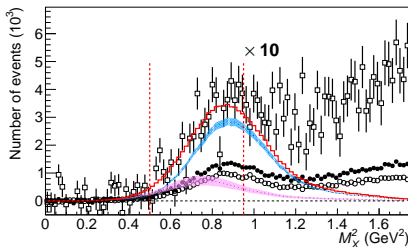
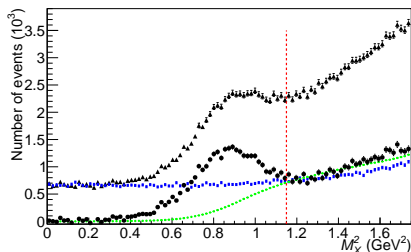
Charged particle veto
in front of scintillator array



E08-025: DVCS off the neutron at different beam energies

- LD₂ as a target $(Q^2 = 1.75 \text{ GeV}^2, x_B = 0.36)$
- Quasi-free p evts subtracted using the (normalized) data from E07-007
- Concurrent running: switching LD₂/LD₂ \rightarrow minimize uncertainties

$$D(e, e\gamma)X - p(e, e\gamma)p = n(e, e\gamma)n + d(e, e\gamma)d$$

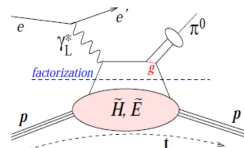


The average momentum transfer to the target is much larger than the np relative momentum, justifying this **impulse approximation**

Preliminary DVCS cross sections off the n & d

DVCS²/Interference separation off the neutron

Potential flavor separation of CFFs combining these data with DVCS off the proton. . .

π^0 electroproduction ($ep \rightarrow ep\pi^0$)

At leading twist:

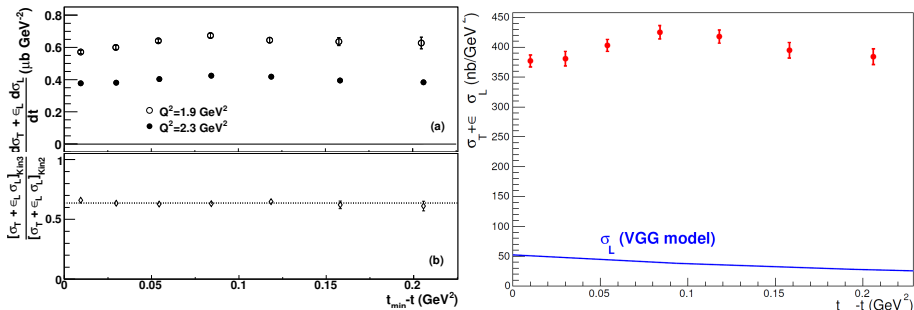
$$\frac{d\sigma_L}{dt} = \frac{1}{2}\Gamma \sum_{h_N, h_{N'}} |\mathcal{M}^L(\lambda_M = 0, h'_N, h_N)|^2 \propto \frac{1}{Q^6} \quad \sigma_T \propto \frac{1}{Q^8}$$

$$\mathcal{M}^L \propto \left[\int_0^1 dz \frac{\phi_\pi(z)}{z} \right] \int_{-1}^1 dx \left[\frac{1}{x - \xi} + \frac{1}{x + \xi} \right] \times \left\{ \Gamma_1 \tilde{H}_{\pi^0} + \Gamma_2 \tilde{E}_{\pi^0} \right\}$$

Different quark weights: flavor separation of GPDs

$$|\pi^0\rangle = \frac{1}{\sqrt{2}}\{|u\bar{u}\rangle - |d\bar{d}\rangle\} \quad \tilde{H}_{\pi^0} = \frac{1}{\sqrt{2}}\left\{\frac{2}{3}\tilde{H}^u + \frac{1}{3}\tilde{H}^d\right\}$$

$$|p\rangle = |uud\rangle \quad H_{DVCS} = \frac{4}{9}H^u + \frac{1}{9}H^d$$

Exclusive π^0 electroproduction cross-sections – Hall A

- $\sigma_T + \epsilon_L \sigma_L \sim Q^{-5}$
(similar to $\sigma_T(ep \rightarrow ep\pi^+)$ measured in Hall C)
- GPDs predict $\sigma_L \sim Q^{-6}$
- σ_T likely to dominate at these Q^2 ,
but L/T separation necessary (\rightarrow new experiment...)

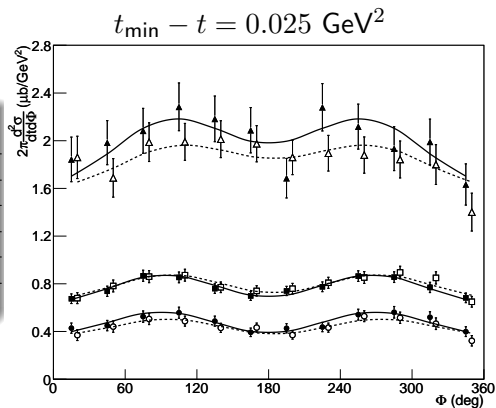
E. Fuchey et al., Phys. Rev. C83 (2011), 025125

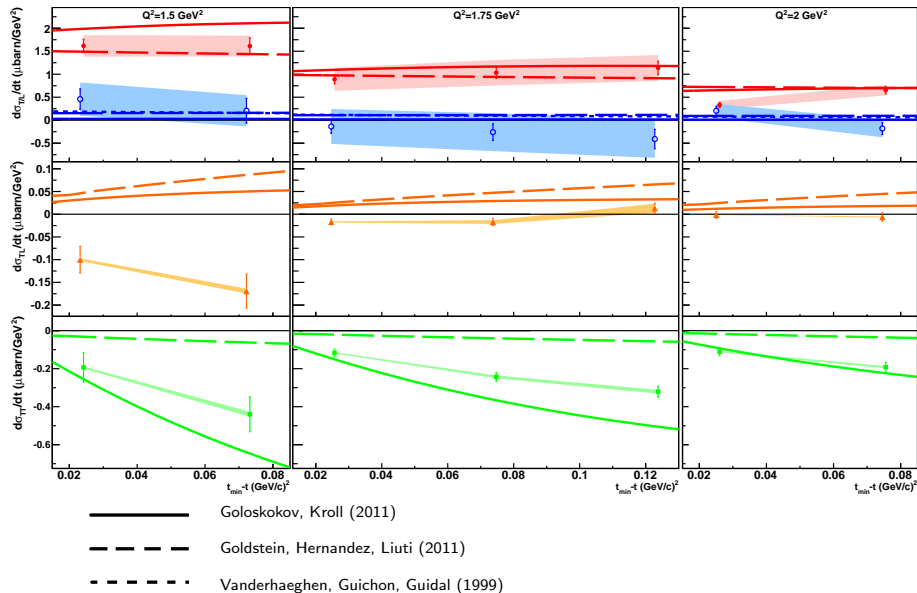
Rosenbluth separation

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} = \frac{1}{2\pi} \Gamma(Q^2, x_B, E) \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{TL}}{dt} \cos\phi + \epsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi \right]$$

Kinematics

Setting	Q^2 (GeV ²)	x_B	E^{beam} (GeV)	ϵ
Kin1	1.50	0.36	3.355	0.52
			5.55	0.84
Kin2	1.75	0.36	4.455	0.65
			5.55	0.79
Kin3	2.00	0.36	4.455	0.53
			5.55	0.72

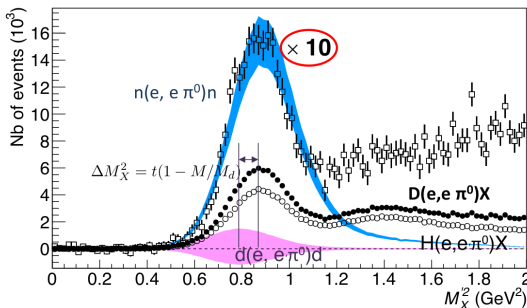


π^0 separated response functions

E08-025: DVCS and π^0 off quasi-free neutrons

- LD₂ as a target
- Quasi-free p evts subtracted using the (normalized) data from E07-007
- Concurrent running: switching LD2/LD2 \rightarrow minimize uncertainties

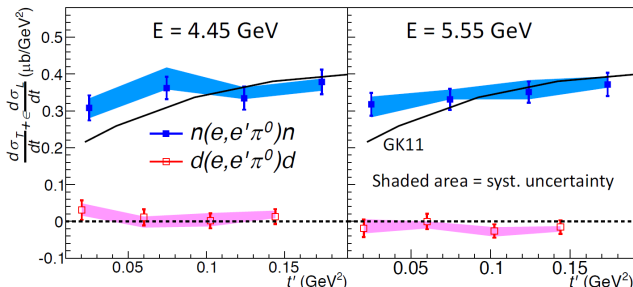
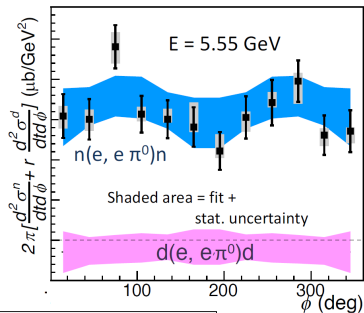
$$D(e, e\pi^0)X - p(e, e\pi^0)p = n(e, e\pi^0)n + d(e, e\pi^0)d$$



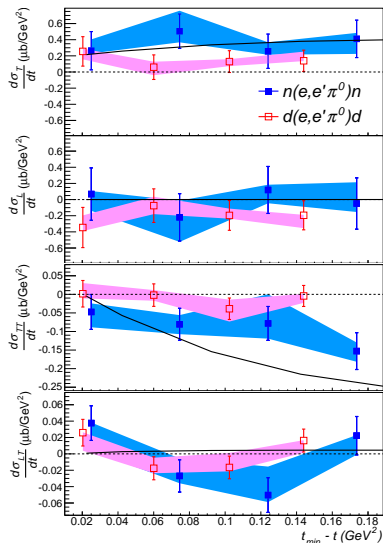
The average momentum transfer to the target is much larger than the np relative momentum, justifying this **impulse approximation**

π^0 electroproduction cross section off the neutron

- Cross section off coherent d found negligible within uncertainties
- Very low E_{beam} dependence of the n cross section \rightarrow dominance of σ_T



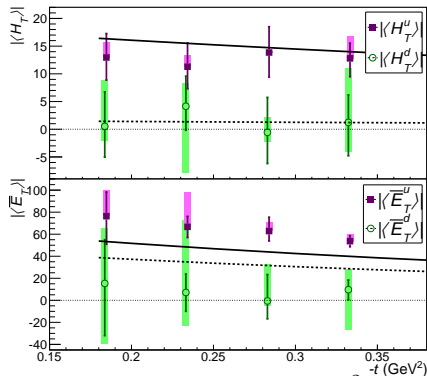
Separated π^0 cross section off the neutron



M. Mazouz et al, Phys.Rev.Lett. 118 (2017)

In the modified factorization approach (KG):

- $d\sigma_T \propto \left[(1 - \xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8M^2} |\langle \bar{E}_T \rangle|^2 \right]$
- $d\sigma_{TT} \propto \frac{t'}{8M^2} |\langle \bar{E}_T \rangle|^2$



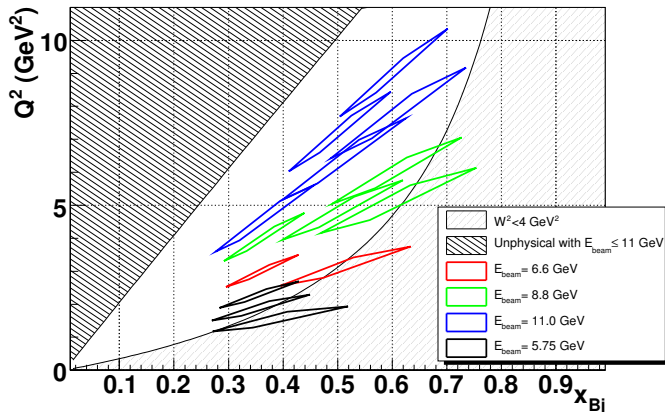
$$|\langle H_T^{p,n} \rangle|^2 = \frac{1}{2} \left| \frac{2}{3} \langle H_T^{u,d} \rangle + \frac{1}{3} \langle H_T^{d,u} \rangle \right|^2$$

E12-06-114: JLab Hall A at 11 GeV

JLab12 with 3, 4, 5 pass beam

(6.6, 8.8, 11.0 GeV beam energy)

DVCS measurements in Hall A/JLab

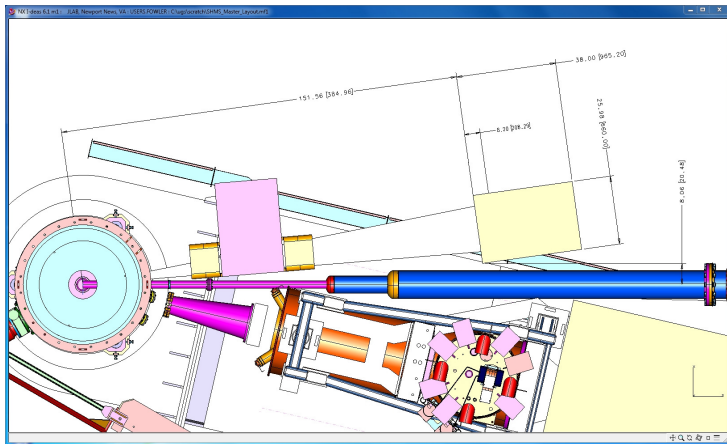


1 year of operations in JLab/Hall A

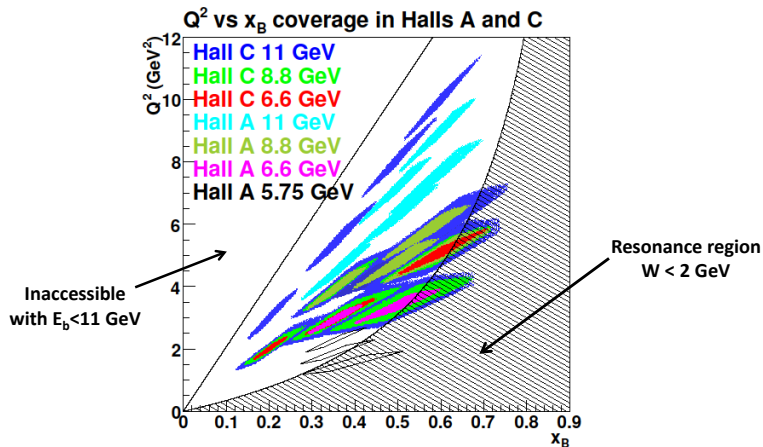
88 days
250k events/setting

E12-13-010: DVCS in Hall C

- HMS ($p < 7.3\text{GeV}$): scattered electron
- PbWO_4 calorimeter: γ/π^0 detection
- Sweeping magnet

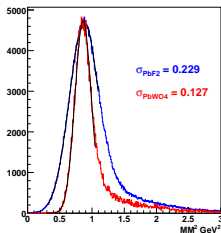
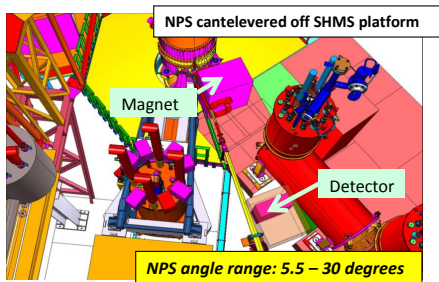


E12-13-010: beam energy separation in Hall C

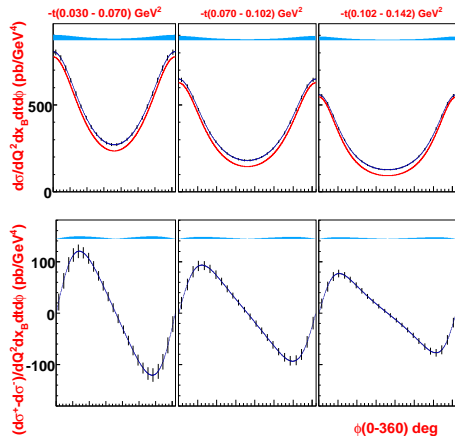


Approved by the PAC, possible running in $\gtrsim 2021$

Projections



- PbF₂ → PbWO₄
- Improved E resolution wrt Hall A



Summary

- Recent high precision DVCS cross sections from Hall A at JLab
- Need of higher twist and/or NLO contributions to fully describe the data (eg. in global GPD fits)
- First separation of DVCS² and BH-DVCS interference in the $eN \rightarrow e\gamma N$ cross section, off the proton and neutron
- L/T separation of π^0 electroproduction cross section off neutron: dominance of σ_T measured
- Flavor separation of transversity GPD convolutions within the modified factorization approach
- Approved program of experiments in Hall A and C to continue these high precision DVCS measurements at 12 GeV

Back-up