# Beautiful paths to probe physics beyond the standard model of particles 



Jennifer school, Trieste, August $2^{\text {nd }} 2018$

## Program of the 3 lectures

- How to study elementary particles
- direct searches and indirect searches
- experiments through history of particles physics
- Rare B decays
- quest for New Physics (beyond Standard Model)
- two approaches for the same quest (LHCb vs Belle)
- CP Violation
- matter and anti-matter
- fully exploiting our detector ....


## At a B-factory...



How many B candidates can I reconstruct with $1 \mathrm{fb}^{-1}$ ?
$1 \mathrm{fb}^{-1} \rightarrow 1 \times 10^{6}$ B produced
but $\mathrm{BF}\left(\mathrm{B} \rightarrow \mathrm{D}^{0} \pi^{-}\right)=5 \times 10^{-3}$
and $\mathrm{BF}\left(\mathrm{D} \rightarrow \mathrm{K}^{-} \pi^{+}\right)=3.8 \%$ and reconstruction efficiency $\sim 10 \% \ldots$ signal yield $\sim 10$ events !!

## Rediscovering beauty: $\mathbf{B} \rightarrow \mathbf{D}^{(*)} \mathbf{h}+\mathbf{B} \rightarrow \mathbf{J} / \psi \mathbf{K}^{(*)}$

with very limited statistics $\left(<1 \mathrm{fb}^{-1}\right)$, Belle II can rediscover the B meson



Show capacity for charm physics in $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \mathbf{c} \overline{\mathbf{c}}$

- $\mathrm{D}^{0}, \mathrm{D}^{+}, \mathrm{D}^{*}$
- Cabibbo favoured and suppressed modes
... for B-physics
- hadronic modes from $\mathrm{b} \rightarrow \mathrm{C}$
- semileptonic decay modes from $\mathrm{b} \rightarrow \mathrm{c}$
that is for dominant decays.... ... we are looking for rare decays


## Rare B decays

- FCNC are strongly suppressed in the SM: only loops + GIM mechanism
- Any new particle generating new diagrams can change the amplitudes

$\rightarrow$ NP beyond the direct reach of the LHC

Three classes of SM processes

$$
\mathcal{O}_{\mathrm{obs}}=\mathcal{O}_{\mathrm{SM}}+\mathcal{O}_{\mathrm{NP}}
$$



New particles can for example contribute to loop or tree level diagrams by enhancing/suppressing decay rates, introducing new sources of CP violation or modifying the angular distribution of the final-state particles

## indirect search: $K_{L}^{0} \rightarrow \mu \mu$

$\mathrm{K}_{\mathrm{L}} \rightarrow \mu^{+} \mu^{-}$decay can be generated by the box diagram:

in a renormalisable gauge theory, is expected to give a branching ratio of $\mathbf{g}^{4} \sim \alpha^{2} \sim \mathbf{1 0}^{-4}$, with $\alpha$ the fine structure constant.

$\mathrm{K}_{\mathrm{L}}^{0} \rightarrow \mu \mu$ was not observed though expected Now BF is measured to be $(6.84 \pm 0.11) 10^{-9}$


## direct search: $J / \psi \rightarrow$ ee

$\rightarrow$ c quark eventually observed in 1974
[Ting], [Richter] J/ $\psi$



With the measured charm quark mass $\mathrm{m}_{\mathrm{c}} \sim 1.27 \mathrm{GeV}$, the predicted rates are in agreement with observation.

## Radiative B decays

artist's view ... of the penguin diagram


- 1975: "South Area Experiment" group conceives CLEO

- 1979: First data collected
- 1980: B meson discovered
- 1983: Ds meson discovered
- 1986: CLEO II detector with Csl calorimeter installed
- 1989: b $\rightarrow$ u transitions discovered
- 1993: b $\rightarrow$ s penguin decays discovered
- 1995: CLEO II.V with silicon vertex detector installed
- 1999: CLEO III with RICH installed
- 2003: CLEO-c data collection started
- 2004: hc discovered and D+ meson decay constant measured
- 2008: Running ends on March 3rd
- 2009: 500th paper published

CLEO observation of $\mathrm{B} \rightarrow \mathrm{K}^{*} \gamma$ [1993]

$\mathbf{B} \rightarrow \mathbf{K}^{*} \gamma$ measuremen ${ }^{*}$
CLEO observation of $\mathrm{B} \rightarrow \mathrm{K}^{*} \gamma$ [1993]

(a) $K_{S}^{0} \pi^{0}$


$\begin{array}{ll}\sum_{\sum}^{0} 800 \\ \sum 600 & \text { (c) } K^{+} \pi^{-}\end{array}$

$\mathrm{N}_{\mathrm{s}} \sim 350$ evts Belle, submitted to PRL
$B \rightarrow K^{*} \gamma$ measurements
simultaneous fit of 4 final states $\Rightarrow$ extraction of BFs....

but uncertainty in the hadronization process limits the ability to predict individual exclusive rates from first principles of the theory 9


## $B \rightarrow K^{*} \gamma$ measurements

simultaneous fit of 4 final states $\Rightarrow$ extraction of BFs, $\Delta_{0+}, \mathrm{A}_{\mathrm{CP}}, \Delta \mathrm{A}_{\mathrm{CP}} \ldots$



## $B \rightarrow K^{*} \gamma$ measurements

## simultaneous fit of 4 final states

$\Rightarrow$ extraction of BFs, $\Delta_{0+}, \mathrm{A}_{\mathrm{CP}}, \Delta \mathrm{A}_{\mathrm{CP}}$
isospin asymmetry : $\Delta_{0+}=\frac{\Gamma\left(B^{0} \rightarrow K^{* 0} \gamma\right)-\Gamma\left(B^{+} \rightarrow K^{*+} \gamma\right)}{\Gamma\left(B^{0} \rightarrow K^{* 0} \gamma\right)+\Gamma\left(B^{+} \rightarrow K^{*+} \gamma\right)}$.


evidence of isospin violation in $K^{*} \gamma$ !
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## $\underline{b} \rightarrow \mathbf{s} \gamma$



- Amplitude $\propto V_{\text {ts }}\left|C_{7}\right|$
- First penguin ever observed (93)
- Experiment:

$$
B \simeq 3.10^{-4}
$$

- $\mathrm{SM}: B=(3.36 \pm 0.23) .10^{-4}$
[Misiak et al., hep-ph/0609232]
$\Rightarrow$ [Misiak et al, arXiv: 1503.01789]
- Strong constraint on New Physics



## $\underline{b} \rightarrow \mathbf{s} \gamma \mathbf{S M}$ branching fraction

[Misiak et al, PRL98, 02202, 2007]

- From effective Hamiltonian one gets the BF
- Uncertainties due to $m_{b}$ and $m_{c}$ : normalise to $b \rightarrow c e \nu$ and $b \rightarrow$ ue $\nu$ [Misiak \& Steinhauser, NPB764:62,2007]
- $b \rightarrow \boldsymbol{s} \gamma$ branching fraction calculated at all NNLO orders in 2006

$$
\mathcal{B}\left(B \rightarrow X_{s} \gamma\right)_{E_{\gamma}>1.6 \mathrm{GeV}}=(3.15 \pm 0.23) 10{ }^{4}
$$

BF very stable versus $\mu$


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## How to estimate the branching fraction $b \rightarrow s \gamma$ ? Semi-inclusive (sum-of-exclusive)


[772 MBB]
[arXiv: 1411.7198] 38 modes

$$
\mathrm{M}_{\mathrm{X}_{\mathrm{s}}}<2.8 \mathrm{GeV} / \mathrm{c}^{2}, \mathrm{E}^{*}>1.9 \mathrm{GeV}
$$

| Mode ID | Final State | Mode ID | Final State |
| :---: | :--- | :---: | :--- |
| 1 | $K^{+} \pi^{-}$ | 20 | $K_{S}^{0} \pi^{+} \pi^{0} \pi^{0}$ |
| 2 | $K_{S}^{0} \pi^{+}$ | 21 | $K^{+} \pi^{+} \pi^{-} \pi^{0} \pi^{0}$ |
| 3 | $K^{+} \pi^{0}$ | 22 | $K_{S}^{0} \pi^{+} \pi^{-} \pi^{0} \pi^{0}$ |
| 4 | $K_{S}^{0} \pi^{0}$ | 23 | $K^{+} \eta$ |
| 5 | $K^{+} \pi^{+} \pi^{-}$ | 24 | $K_{S}^{0} \eta$ |
| 6 | $K_{S}^{0} \pi^{+} \pi^{-}$ | 25 | $K^{+} \eta \pi^{-}$ |
| 7 | $K^{+} \pi^{-} \pi^{0}$ | 26 | $K_{S}^{0} \eta \pi^{+}$ |
| 8 | $K_{S}^{0} \pi^{+} \pi^{0}$ | 27 | $K^{+} \eta \pi^{0}$ |
| 9 | $K^{+} \pi^{+} \pi^{-} \pi^{-}$ | 28 | $K_{S}^{0} \eta \pi^{0}$ |
| 10 | $K_{S}^{0} \pi^{+} \pi^{+} \pi^{-}$ | 29 | $K^{+} \eta \pi^{+} \pi^{-}$ |
| 11 | $K^{+} \pi^{+} \pi^{-} \pi^{0}$ | 30 | $K_{S}^{0} \eta \pi^{+} \pi^{-}$ |
| 12 | $K_{S}^{0} \pi^{+} \pi^{-} \pi^{0}$ | 31 | $K^{+} \eta \pi^{-} \pi^{0}$ |
| 13 | $K^{+} \pi^{+} \pi^{+} \pi^{-} \pi^{-}$ | 32 | $K_{S}^{0} \eta \pi^{+} \pi^{0}$ |
| 14 | $K_{S}^{0} \pi^{+} \pi^{+} \pi^{-} \pi^{-}$ | 33 | $K^{+} K^{+} K^{-}$ |
| 15 | $K^{+} \pi^{+} \pi^{-} \pi^{-} \pi^{0}$ | 34 | $K^{+} K^{-} K_{S}^{0}$ |
| 16 | $K_{S}^{0} \pi^{+} \pi^{+} \pi^{-} \pi^{0}$ | 35 | $K^{+} K^{+} K^{-} \pi^{-}$ |
| 17 | $K^{+} \pi^{0} \pi^{0}$ | 36 | $K^{+} K^{-} K_{S}^{0} \pi^{+}$ |
| 18 | $K_{S}^{0} \pi^{0} \pi^{0}$ | 37 | $K^{+} K^{+} K^{-} \pi^{0}$ |
| 19 | $K^{+} \pi^{-} \pi^{0} \pi^{0}$ | 38 | $K^{+} K^{-} K_{S}^{0} \pi^{0}$ |




# Semi-inclusive (sum-of-exclusive) <br> 38 modes 

$\mathrm{M}_{\mathrm{X}_{\mathrm{s}}}<2.8 \mathrm{GeV} / \mathrm{c}^{2}, \mathrm{E}^{*}>1.9 \mathrm{GeV}$
possible but large systematics (difficult to estimate/trust)

| Mode ID | Final State | Mode ID | Final State |
| :---: | :--- | :---: | :--- |
| 1 | $K^{+} \pi^{-}$ | 20 | $K_{S}^{0} \pi^{+} \pi^{0} \pi^{0}$ |
| 2 | $K_{S}^{0} \pi^{+}$ | 21 | $K^{+} \pi^{+} \pi^{-} \pi^{0} \pi^{0}$ |
| 3 | $K^{+} \pi^{0}$ | 22 | $K_{S}^{0} \pi^{+} \pi^{-} \pi^{0} \pi^{0}$ |
| 4 | $K_{S}^{0} \pi^{0}$ | 23 | $K^{+} \eta$ |
| 5 | $K^{+} \pi^{+} \pi^{-}$ | 24 | $K_{S}^{0} \eta$ |
| 6 | $K_{S}^{0} \pi^{+} \pi^{-}$ | 25 | $K^{+} \eta \pi^{-}$ |
| 7 | $K^{+} \pi^{-} \pi^{0}$ | 26 | $K_{S}^{0} \eta \pi^{+}$ |
| 8 | $K_{S}^{0} \pi^{+} \pi^{0}$ | 27 | $K^{+} \eta \pi^{0}$ |
| 9 | $K^{+} \pi^{+} \pi^{-} \pi^{-}$ | 28 | $K_{S}^{0} \eta \pi^{0}$ |
| 10 | $K_{S}^{0} \pi^{+} \pi^{+} \pi^{-}$ | 29 | $K^{+} \eta \pi^{+} \pi^{-}$ |
| 11 | $K^{+} \pi^{+} \pi^{-} \pi^{0}$ | 30 | $K_{S}^{0} \eta \pi^{+} \pi^{-}$ |
| 12 | $K_{S}^{0} \pi^{+} \pi^{-} \pi^{0}$ | 31 | $K^{+} \eta \pi^{-} \pi^{0}$ |
| 13 | $K^{+} \pi^{+} \pi^{+} \pi^{-} \pi^{-}$ | 32 | $K_{S}^{0} \eta \pi^{+} \pi^{0}$ |
| 14 | $K_{S}^{0} \pi^{+} \pi^{+} \pi^{-} \pi^{-}$ | 33 | $K^{+} K^{+} K^{-}$ |
| 15 | $K^{+} \pi^{+} \pi^{-} \pi^{-} \pi^{0}$ | 34 | $K^{+} K^{-} K_{S}^{0}$ |
| 16 | $K_{S}^{0} \pi^{+} \pi^{+} \pi^{-} \pi^{0}$ | 35 | $K^{+} K^{+} K^{-} \pi^{-}$ |
| 17 | $K^{+} \pi^{0} \pi^{0}$ | 36 | $K^{+} K^{-} K_{S}^{0} \pi^{+}$ |
| 18 | $K_{S}^{0} \pi^{0} \pi^{0}$ | 37 | $K^{+} K^{+} K^{-} \pi^{0}$ |
| 19 | $K^{+} \pi^{-} \pi^{0} \pi^{0}$ | 38 | $K^{+} K^{-} K_{S}^{0} \pi^{0}$ |

$\left\{\begin{array}{l}B\left(B \rightarrow X_{s} \gamma\right)=(3.51 \pm 0.17 \pm 0.33) \times 10^{-4} \\ B\left(B \rightarrow X_{s} \gamma\right)=(3.29 \pm 0.19 \pm 0.48) \times 10^{-4}\end{array}\right.$
[syst: cross-feed, peaking BG, $\mathrm{X}_{\mathrm{s}}$ fragmentation]

## $\underline{B \rightarrow} \mathbf{X}_{\mathrm{s}} \gamma$ spectrum

- $\mathrm{b} \rightarrow \mathrm{s} \gamma$ is a 2 -body decay. The energy of the photon in the $b$ quark frame is

$$
\mathrm{E}_{\gamma}=\frac{\mathrm{m}_{\mathrm{b}}}{2}\left(1-\frac{\mathrm{m}_{\mathrm{s}}^{2}}{\mathrm{~m}_{\mathrm{b}}^{2}}\right) \simeq \frac{\mathrm{m}_{\mathrm{b}}}{2}
$$

- But we measure $B \rightarrow X_{s} \gamma$ and in the $B$ meson, the b quark is moving which smears the energy spectrum
$\rightarrow$ Mean $\sim \frac{\mathrm{m}_{\mathrm{B}}}{2}$

$\rightarrow$ Width $\sim$ Fermi motion in B meson
- The BF is calculated for some energy cutoff (1.6 GeV). For other cutoffs $\mathrm{E}_{0}$ apply [Misiak et al, (2007)]

$$
\left(\frac{B\left(\mathrm{E}_{\gamma}>\mathrm{E}_{0}\right)}{B\left(\mathrm{E}_{\gamma}>1.6 \mathrm{GeV}\right)}\right) \simeq 1+0.15 \frac{\mathrm{E}_{0}}{1.6 \mathrm{GeV}}-0.14\left(\frac{\mathrm{E}_{0}}{1.6 \mathrm{GeV}}\right)^{2}
$$

## $b \rightarrow$ s $\gamma$ spectrum at Belle



One would like to measure the photon energy spectrum in $b \rightarrow s \gamma$ decays

- Be unbiased: only look at the $\gamma$
- B mesons only decay to $\gamma$ via $b \rightarrow s \gamma$
- But there are indirect $\gamma$ from $\pi^{0}$ and $\eta$ in $B \bar{B}$ events
- ... and a lot more indirect $\pi^{0}$ and $\eta$ in non-B $\bar{B}$ events
$\Rightarrow$ Lots of background at low energy


## $\mathbf{b} \rightarrow \mathbf{s} \gamma$ spectrum at Belle

inclusive $B \rightarrow X_{s} \gamma$ measurement untagged
lepton tag: background suppression, low stat


Example with data sets

- $140 \mathrm{fb}^{-1} \mathrm{ON}$-resonance
- $15 \mathrm{fb}^{-1}$ OFF-resonance

Event selection:

- No kinematic constraints
- Only a high energy photon measured in $\Upsilon(4 S)$ rest frame
- Lower $\mathrm{E}_{\gamma}$ threshold (1.7 GeV)
- Hadronic events with isolated photon(s) in ECL. $\mathrm{E}^{*}>1.5 \mathrm{GeV}$.
- Veto $\gamma$ from $\pi^{0}$ and $\eta$
- Apply event shape cuts to suppress continuum background.


## The spectrum



## The spectrum



Endpoint check:
Photons from $\mathrm{e}^{+} \mathrm{e}^{-}$collisions can have an energy up to 5 GeV

But not if they come from a $B$ decay. The kinematic limit is $\mathrm{E}^{*}=\mathrm{m}_{\mathrm{B}} / 2$.

No significant deviation from 0 observed

## The spectrum


$B \bar{B}$ subtraction :
Using measured $\pi^{0}$ and $\eta$ spectra and some efficiency-corrected MC.

## The spectrum



Raw spectrum after all cuts and background corrections

Signal yield:
$24100 \pm 2200$ events

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## The spectrum



Efficiency corrected spectrum


## $\underline{\mathbf{X}_{\mathrm{s}} \gamma \text { inclusive }}$

Lower $E_{\gamma}$ threshold ( 1.7 GeV ) $\Rightarrow 97 \%$ of the spectrum !


$$
\begin{aligned}
& \boldsymbol{B}\left(\mathbf{B} \rightarrow \mathbf{X}_{\mathbf{s}} \gamma\right)=(\mathbf{3 . 4 5} \pm \mathbf{0 . 1 5} \pm \mathbf{0 . 4 0}) \times \mathbf{1 0}^{-\mathbf{4}}\left(\text { for } \mathrm{E}_{\gamma}^{*}>1.7 \mathrm{GeV}\right) \\
& B\left(\mathrm{~B} \rightarrow \mathrm{X}_{\mathrm{s}} \gamma\right)=(3.21 \pm 0.15 \pm 0.29 \pm 0.08) \times 10^{-4}\left(\text { for } \mathrm{E}_{\gamma}^{*}>1.8 \mathrm{GeV}\right) \\
& B\left(\mathrm{~B} \rightarrow \mathrm{X}_{\mathrm{s}} \gamma\right)=(3.06 \pm 0.41 \pm 0.26) \times 10^{-4}\left(\text { for } \mathrm{E}_{\gamma}^{*}>2.0 \mathrm{GeV}\right)
\end{aligned}
$$

- Most precise measurement of $B\left(\mathrm{~B} \rightarrow \mathrm{X}_{\mathrm{s}} \gamma\right)$ (lowest $\mathrm{E}_{\gamma}^{*}$ threshold)
- Crucial input for global fit to extract $\left|\mathrm{V}_{\mathrm{ub}}\right|$ and $\mathrm{B} \rightarrow \mathrm{X}_{\mathrm{s}} \gamma$ decay rate
$\circ B$ is given for $\mathrm{E}_{\gamma}$ thresholds: 1.7, 1.8, 1.9, 2.0 GeV
- Systematic error is dominated b24ff - resonance subtraction !


## $\mathbf{B} \rightarrow \mathbf{X}_{\mathrm{s}} \gamma$ as an illustration



W

Sensitive to NP


NNLO SM calculation :
$B_{S M}\left(\mathrm{~B} \rightarrow \mathrm{X}_{\mathrm{s}} \gamma\right)=(3.36 \pm 0.23) \times 10^{-4}$ $\left(\right.$ for $\left.E_{\gamma}>1.6 \mathrm{GeV}\right) \quad$ M.Misiak et al. [arXiv: 1503.01789] (central value increased by $6.4 \%$ compared to 2007 value ) PRL 98, 022002 (2007)

The lower $\gamma$ energy threshold, the smaller the model uncertainties in SM, but the larger background in measurement

Charged Higgs (2 HDM Type II) bound
(up- and down-type quarks couple to separate doublets)


## $\underline{B \rightarrow} \mathbf{X}_{s} \boldsymbol{\gamma}$

$$
\begin{aligned}
& \text { WA: } B\left(B \rightarrow X_{s} \gamma\right)=(3.49 \pm 0.20) \times 10^{-4}\left(\text { for } E_{\gamma}>1.6 \mathrm{GeV}\right) \\
& \mathrm{Vm}: B\left(B \rightarrow X_{\mathrm{s}} \gamma\right)=(3.36 \pm 0.23) \times 10^{-4}\left(\text { for } E_{\gamma}>1.6 \mathrm{GeV}\right)
\end{aligned}
$$

[Misiak et al, arXiv:1503.01789]

## Charged Higss bound (2 HDM TypeII): $M_{H^{+}}>400 \mathrm{GeV} @ 95 \%$ C.L.

[arXiv:1706.07414]
THDM Type II - Flavour constraints


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## Rare B decays atLHCb

## LHCb is

- 1075 members, from 68 institutes in 17 countries (September 2014)
- Dedicated experiment for precision measurements of CP violation and rare decays
- Beautiful, charming, strange physics program

- $p p$ collisions at $\sqrt{s}=8(13) \mathrm{TeV}$ in RunI (RunII)
- $\quad b \bar{b}$ quark pairs produced correlated in the forward region
- Luminosity of $4 \times 10^{32} \mathrm{~cm}^{-2} \mathrm{~s}^{-1}$


## LHCb



## LHCb

## Tracking system

Measure displaced vertices and momentum of particles


Vertex and IP resolution $\sigma(\mathrm{IP}) \sim 24 \mu \mathrm{~m}$ at $\mathrm{P}_{\mathrm{T}}=2 \mathrm{GeV} / \mathrm{c}$ $\sigma_{\mathrm{BV}} \sim 16 \mu \mathrm{~m}$ in $\mathrm{x}, \mathrm{y}$

10m (p) $/ \mathrm{p}=0.4 \%-0.6 \%$ resolution
$\sigma(\mathrm{p}) / \mathrm{p}=0.4 \%-0.6 \%$ for $\mathrm{p} \in[0,100] \mathrm{GeV} / \mathrm{c}$ $\sigma\left(\mathrm{m}_{\mathrm{B}}\right) \sim 24 \mathrm{MeV}$ for two body decays

## LHCb

## Particle identification

Distinguish between pions, kaons, protons, electrons and muons


Kaon identification
$\epsilon_{\mathrm{K}} \sim 95 \%, \epsilon_{\pi \rightarrow \mathrm{K}}$ few $\%$

Muon identification
$\epsilon_{\mu}=98 \%, \epsilon_{\pi \rightarrow \mu}=0.6 \%$

## LHCb

## Trigger system <br> Write out 5000 events/sec



## Belle(II), LHCb side by side

## Belle (II)

$$
\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \mathrm{Y}(4 \mathrm{~S}) \rightarrow \mathrm{b} \overline{\mathrm{~b}}
$$

at $Y(4 S)$ : 2 B's $^{\prime}\left(B^{0}\right.$ or $\left.B^{+}\right)$and nothing else $\Rightarrow$ clean events
$\sigma_{\mathrm{b} \overline{\mathrm{b}}} \sim 1 \mathrm{nb} \Rightarrow 1 \mathrm{fb}^{-1}$ produces $10^{6} \mathrm{~B} \overline{\mathrm{~B}}$
$\sigma_{\mathrm{bb}} / \sigma_{\text {total }} \sim 1 / 4$
(in the context of B anomalies) LHCb
$p p \rightarrow b \bar{b} X$
production of $\mathrm{B}^{+}, \mathrm{B}^{0}, \mathrm{~B}_{\mathrm{s}}, \mathrm{B}_{\mathrm{c}}, \Lambda_{\mathrm{b}} \ldots$
but also a lot of other particles in the event
$\Rightarrow$ lower reconstruction efficiencies
$\sigma_{\text {bб }}$ much higher than at the $\mathrm{Y}(4 \mathrm{~S})$

|  | $\sqrt{\mathbf{s}}[\mathbf{G e V}]$ | $\boldsymbol{\sigma}_{\mathbf{b 5}}[\mathbf{n b}]$ | $\boldsymbol{\sigma}_{\mathbf{b 5}} / \boldsymbol{\sigma}_{\mathbf{t a t}}$ |
| :---: | :---: | :---: | :---: |
| HERA pA | 42 GeV | $\sim 30$ | $\sim 10^{-6}$ |
| Tevatron | 2 TeV | 5000 | $\sim 10^{-3}$ |
| LHC | 8 TeV | $\sim 3 \times 10^{5}$ | $\sim 5 \times 10^{-3}$ |
|  | 14 TeV | $-6 \times 10^{5}$ | $\sim 10^{-2}$ |

$\mathbf{b} \overline{\mathbf{b}}$ production cross-section $\sim 5 \times$ Tevatron $, \sim 500,000 \times$ BaBar/Belle ! !
$\sigma_{\mathrm{b}} / \sigma_{\text {total }}$ much lower than at the $\mathrm{Y}(4 \mathrm{~S})$
$\Rightarrow$ lower trigger efficiencies

B mesons live relativey long
mean decay length $\beta \gamma c \tau \sim 200 \mu \mathrm{~m}$
mean decay length $\beta \gamma c \tau \sim 7 \mathrm{~mm}$
data taling period(s)

$$
[1999-2010]=1 \mathrm{ab}^{-1}
$$

[run I: 2010-2012] $=3 \mathrm{fb}^{-1}$,
[run II: 2015-2018] $=2 \mathrm{fb}^{-1} \rightarrow 8 \mathrm{fb}^{-1}$ ?
(near) future
[Belle II from 2018] $\rightarrow 50 \mathrm{ab}^{-1}$
[LHCb upgrade from 2020]

## $\underline{B}_{(s)} \rightarrow \mu \mu:$ ultra rare processes...

loop diagram + suppressed in SM + theoretically clean = an excellent place to look for new physics


## Leptonic decays

$$
\begin{aligned}
& B_{(s)}^{0} \rightarrow \ell^{+} \ell^{-} \\
& B R\left(B_{(q)}^{0} \rightarrow \ell^{+} \ell^{-}\right)=\frac{\tau_{B} G_{F}^{4} M_{W}^{2} s i n^{4} \theta_{W}}{8 \pi^{\varphi}}\left|C_{10} V_{t b} V_{t q}^{*}\right| F_{B}^{2} m_{B} m_{\ell}^{2} \times \left\lvert\, \sqrt{1-\frac{4 m_{\ell}^{2}}{m_{B}^{2}}}\right.
\end{aligned}
$$

Branching ratio proportional to the lepton mass squared

$$
\frac{B R\left(B_{(q)}^{0} \rightarrow \tau^{+} \tau-\right)}{B R\left(B_{(q)}^{0} \rightarrow \mu^{+} \mu^{-}\right)} \sim \frac{m_{\tau}^{2}}{m_{\mu}^{2}} \quad \frac{B R\left(B_{(q)}^{0} \rightarrow \mu^{+} \mu^{-}\right)}{B R\left(B_{(q)}^{0} \rightarrow e^{+} e^{-}\right)} \sim \frac{m_{\mu}^{2}}{m_{e}^{2}}
$$

Helicity suppression, same reason why the pion decays into muon instead of electron $\Rightarrow$ true only in SM

All parameters either measurable or calculable with high precision valid only in Minimal Flavour Violating Models (where the flavour structure is described only by CKM)

In a ''general'' NP scenarios, the branching ratio of B leptonic decay is given by

$$
B R\left(B_{s}^{0} \rightarrow \mu^{+} \mu^{-}\right) \propto\left(1-\frac{4 m_{\ell}^{2}}{m_{B}^{2}}\right)\left|C_{S}-C_{S}^{\prime}\right|^{2}+\left|\left(C_{P}-C_{P}^{\prime}\right)^{2}+2 \frac{m_{\ell}^{2}}{m_{B}^{2}}\left(C_{10}-C_{10}^{\prime}\right)\right|^{2}
$$

## $\underline{B}_{(\mathrm{s})} \rightarrow \mu \mu:$ ultra rare processes...



## $\mathbf{B}_{\mathrm{s}} \rightarrow \mu^{+} \mu^{-}$results

$\mathrm{B}\left(\mathrm{B}_{\mathrm{s}}^{0} \rightarrow \mu^{+} \mu^{-}\right)=\left(2.8_{-0.6}^{+0.7}\right) \times 10^{-9}$
first observation : 6.2 $\sigma$ significance $\mathrm{B}\left(\mathrm{B}^{0} \rightarrow \mu^{+} \mu^{-}\right)=\left(3.9_{-1.4}^{+1.6}\right) \times 10^{-10}$
first evidence: $3.0 \sigma$ significance

## [arXiv:1703.05747]

SM: heavy state decays to $\mu^{+} \mu^{-}$ first lifetime measurement:

$$
\tau\left(\mathrm{B}_{\mathrm{s}} \rightarrow \mu^{+} \mu^{ \pm}\right)=2.04 \pm 0.44 \pm 0.05 \mathrm{ps}
$$


$\mathrm{B}\left(\mathrm{B}_{\mathrm{s}}^{0} \rightarrow \mu^{+} \mu^{-}\right)=\left(3.0 \pm 0.6_{-0.2}^{+0.3}\right) \times \mathbf{1 0}^{-9}(7.8 \sigma$ significance $)$
$B\left(B^{0} \rightarrow \mu^{+} \mu^{-}\right)<3.4 \times 10^{-10} @ 90 \%$ CL

## Constraints on NP models

From D. Straub, arXiv:1205.6094


