

Elliptic non-abelian DT invariants of \mathbb{C}^3

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- Free field realization of Z_N .

Brane Construction

- Setup (type IIB string):

	x_1	x_2	x_3	x_4	x_5	x_6	x_7	x_8	x_9	x_{10}
D1	—	—	•	•	•	•	•	•	•	•
D7	—	—	—	—	—	—	—	—	•	•

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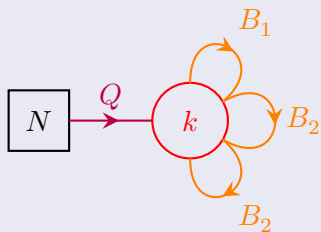
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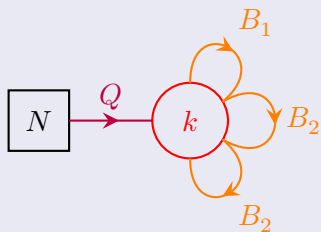
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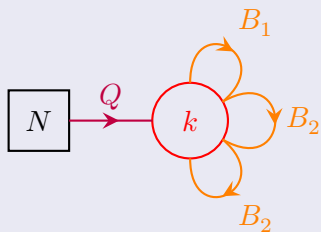
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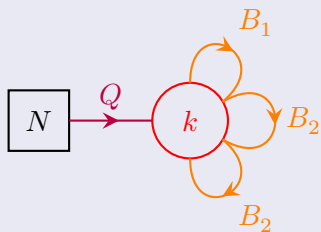
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Interactions

$$W = \text{tr } B_1[B_2, B_3]$$

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Study the Elliptic Genus of this moduli space

Elliptic Genus

The quantity we want to compute is:

$$Z_{N,k}(\vec{\xi}, \vec{\epsilon}) = \text{tr}_{\text{RR}}(-1)^F e^{-2\pi(\tau_2 H + \tau_1 P)} e^{2\pi i \xi_\alpha F_\alpha} e^{2\pi i \epsilon_i R_i} ,$$

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Effect in Path Integral computation

Turn on a flat background gauge field $A^{(F)}$

$$\xi_\alpha = \int_t A_\alpha^{(F)} - \tau \int_s A_\alpha^{(F)}$$

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Path Integral \Rightarrow Finite Dimensional Integral

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Given the quantum numbers of various multiplets

	Q	B_1	B_2	B_3	fugacity
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there is a recipe [Benini-Eager-Hori-Tachikawa] to get

$$Z_{N,k}(\vec{\xi}, \vec{\epsilon}) = \frac{1}{k!} \int_{\text{JK-contour}} du_1 \dots du_k \text{Rational function of Jacobi } \theta_1(\tau | \text{fugacities})$$

Anomalies

- $\theta_1(\tau|u + a + b\tau) = (-1)^{a+b} e^{-2\pi i b u} e^{-i\pi b^2 \tau} \theta_1(\tau|u)$ for $a, b \in \mathbb{Z}$;

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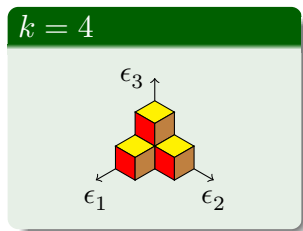
Pole of the integrand JK allowed $\Leftrightarrow N$ Colored Plane Partition of k

(Colored) Plane Partition

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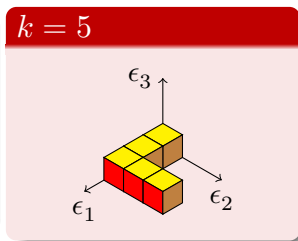
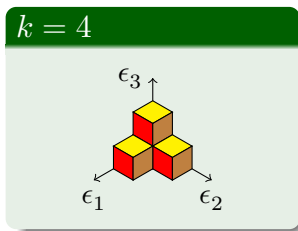
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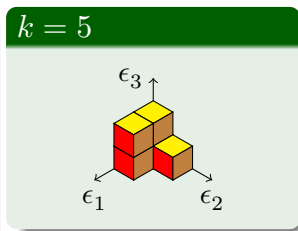
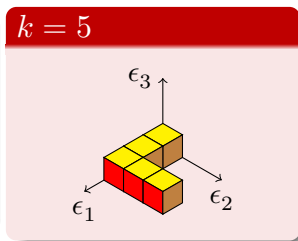
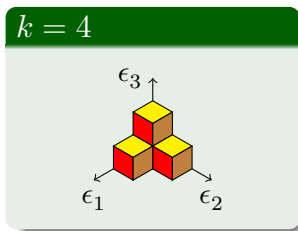
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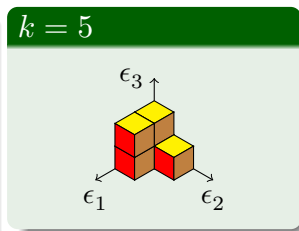
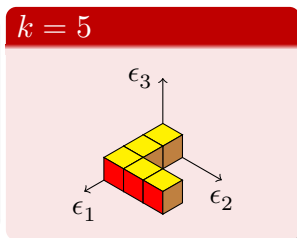
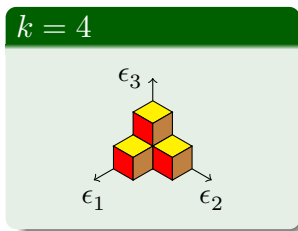
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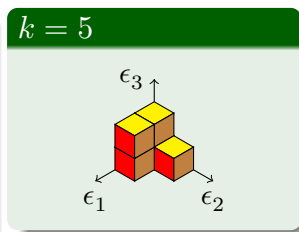
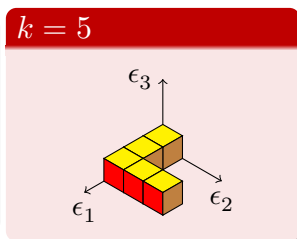
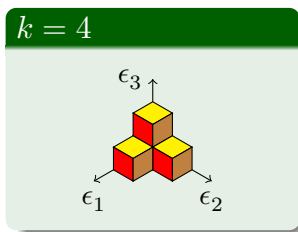


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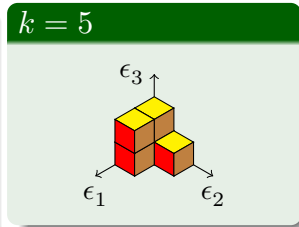
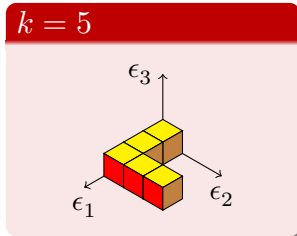
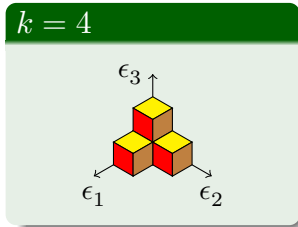
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- The generating function of CPP is $\Phi(v)^N$.

Result for $Z_{N,k}$

We have for $\epsilon = \frac{n}{N}$ with $n \in \mathbb{Z}$ that

$$\begin{aligned} Z_{N,k}(\vec{\xi}, \vec{\epsilon}) &= \sum \text{Residues at pole corresponding to } N \text{ CPP of } k \\ &= \begin{cases} (-1)^{nk} \Phi_{\frac{k}{N} \gcd(n,N)}^{\gcd(n,N)} & \text{if } \frac{N}{\gcd(n,N)} \mid k \\ 0 & \text{otherwise} \end{cases} \end{aligned}$$

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- Interpretation on terms of F theory.

Dimensional reduction to D0/D6

Study the refined Witten index $\tilde{Z}_{N,k}$ of $\mathcal{N} = 4$ GQM on S^1

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- Same computation with replacement $\theta_1 \mapsto \sin$;
- The result can be written in a compact form resumming it as before $\tilde{Z}_N(\vec{\xi}, \vec{\epsilon}; v) = \sum_{k=0}^{\infty} \tilde{Z}_{N,k}(\vec{\xi}, \vec{\epsilon}) v^k$, it is

$$\tilde{Z}_N(\vec{\epsilon}; v) = \text{PE}_{v, \vec{p}} \left[- \frac{(1-p_1 p_2)(1-p_1 p_3)(1-p_2 p_3)}{(1-p_1)(1-p_2)(1-p_3)} \times \right. \\ \left. \times p^{-\frac{N}{2}} \frac{1-p^N}{1-p} \frac{v}{(1-vp^{-\frac{N}{2}})(1-vp^{\frac{N}{2}})} \right]$$

$$(p_i \equiv e^{2\pi i \epsilon_i} \text{ and } p \equiv e^{2\pi i \epsilon}).$$

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- Non trivial dependence on N ;
- Lift to M-theory: Taub-NT $_N \times \mathbb{C}^3$ fibered on S^1 ;
- Computation of 11d-SUGRA index in the Ω -background gives exactly the argument of PE.

Further reduction: D(-1)/D5

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$$\bar{Z}_N(\vec{\epsilon}; v) = \sum_{k=0}^{\infty} \bar{Z}_{N,k}(\vec{\xi}, \vec{\epsilon}) = [\Phi(v)]^{-N \frac{(\epsilon_1 + \epsilon_2)(\epsilon_1 + \epsilon_3)(\epsilon_2 + \epsilon_3)}{\epsilon_1 \epsilon_2 \epsilon_3}} ;$$

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- This is an example of *trivial factorization* [Nekrasov].

Free field realization of Z_N

The propagator of free a free boson on torus is

$$\langle \phi(u, \bar{u}) \phi(w, \bar{w}) \rangle_{T^2} = \log G(u, \bar{u}; w, \bar{w}) ,$$

where

$$G(u, \bar{u}; w, \bar{w}) = e^{-\frac{2\pi}{\tau_2} [\Im(u-w)]^2} \left| \frac{\theta_1(\tau|u-w)}{2\pi\eta^3(\tau)} \right|^2 .$$

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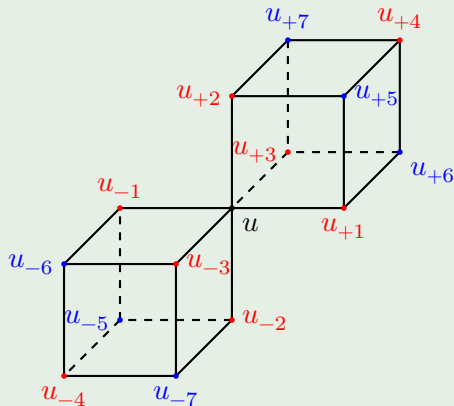
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- It contains Jacobi θ_1 function;
- Differently from the planar case, it is *not* a sum of a holomorphic and anti-holomorphic part because of the \mathfrak{I}^2 .

Multi-local Vertex Operator

Consider the following vertex operator

$$\mathcal{V}_{\vec{\epsilon}}(u) = \prod_{i=1}^7 :e^{\lambda_i \phi_i(u_{+i})} :: e^{-\lambda_i \phi_i(u_{-i})} : ,$$



with

$$\begin{aligned} u_{\pm 1} &= u \pm \frac{\epsilon_1}{2}, & u_{\pm 2} &= u \pm \frac{\epsilon_2}{2}, \\ u_{\pm 3} &= u \pm \frac{\epsilon_3}{2}, & u_{\pm 4} &= u \pm \frac{\epsilon}{2}, \\ u_{\pm 5} &= u \pm \frac{\epsilon_1 + \epsilon_2}{2}, & u_{\pm 6} &= u \pm \frac{\epsilon_1 + \epsilon_3}{2}, \\ u_{\pm 7} &= u \pm \frac{\epsilon_2 + \epsilon_3}{2}; \end{aligned}$$

and

$$\vec{\lambda} = (1, 1, 1, 1, i, i, i).$$

Source (“Hamiltonian”) operator

Let us consider the following operator:

$$H = \frac{1}{2\pi i} \oint_{\Gamma} \partial \phi_4(u) \omega(w) dw ,$$

where the contour Γ is around $u = 0$ and encloses all $u_{\pm i}$, and

$$\omega(w) = - \sum_{\alpha=1}^N \log \theta_1 \left(\tau \left| w + \xi_{\alpha} - \frac{\epsilon}{2} \right. \right)$$

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It can be a starting point for a generalization of *Hirota Bilinear Equations*.

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What we did

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- Following the same program for threefold with more complex topology (i.e. the conifold);
- Explore the D0/D8 case [Nekrasov: "Magnificent four"];
- Understand the possible link between the free field realization point of view and integrable hierarchies.

Thank you for your attention!