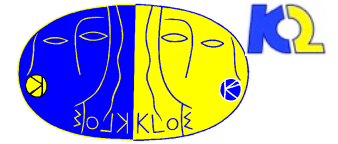

Test of discrete symmetries with neutral kaons at KLOE-2



Antonio Di Domenico
Dipartimento di Fisica, Sapienza Università di Roma
and INFN sezione di Roma, Italy



on behalf of the KLOE-2 collaboration



Workshop on Quantum foundations
LNF - Frascati, Italy 29 November – 1 December 2017

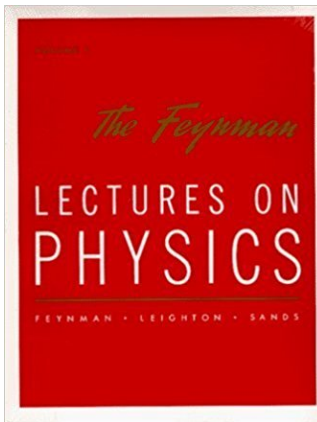
K mesons – more than 70 years history

- 1944 : first indication of a new charged particle with mass $\sim 0.5 \text{ GeV}/c^2$ in cosmic rays (Leprince-Ringuet, Lheritier)
- 1947 : first K^0 observation in cloud chamber - V particle (Rochester, Butler)
- 1955 : introduction of **Strangeness** (Gell-Mann, Nishijima)
 K^0, \bar{K}^0 are two distinct particles (Gell-Mann, Pais) **Strangeness oscillation**
- 1955 prediction of **regeneration** of short-lived particle (Pais, Piccioni)
- 1956 Observation of long lived K_L (BNL Cosmotron)
- 1957 τ - θ puzzle on spin-parity assignment, **P violation** in weak interactions
- 1960: $\Delta m = m_L - m_S$ measured from **regeneration**
- 1964: discovery of **CP violation** (Cronin, Fitch,...)
- 1970 : suppression of FCNC, $K_L \rightarrow \mu\mu$ - GIM mechanism/charm hypothesis
- 1972 : Kobayashi Maskawa six quark model: CP violation explained in SM
- 1992- 2000 : CPLEar: K^0, \bar{K}^0 time evolution and decays, **T, CP, CPT tests**
- 1999-2003 : KTeV and NA48 (prev. E731 and NA31): **direct CP violation** proven : $\varepsilon'/\varepsilon \neq 0$
- 2003-2008 : NA48/2: charged kaon beam, search for direct CP viol.
- 2000-2006 : KLOE at DaΦne: first Φ factory enters in operation, V_{us} and precision tests of the SM, entangled neutral K pairs and **CPT and QM tests**.

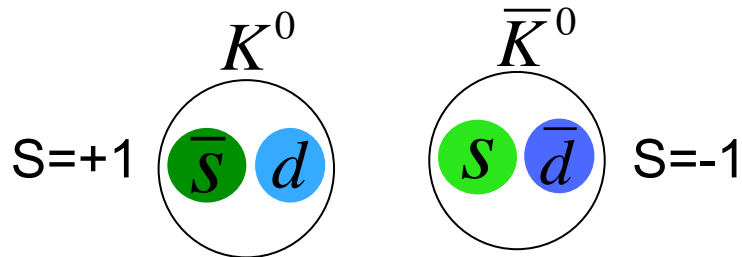
Neutral K meson system: a jewel donated to us by Nature



R. Feynman



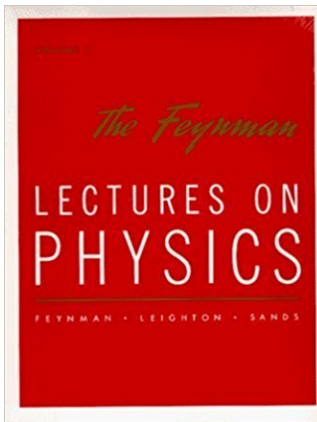
“If there is any place where we have a chance to test the main principles of quantum mechanics in the purest way - does the superposition of amplitudes work or doesn't it? - this is it.”



Neutral K meson system: a jewel donated to us by Nature

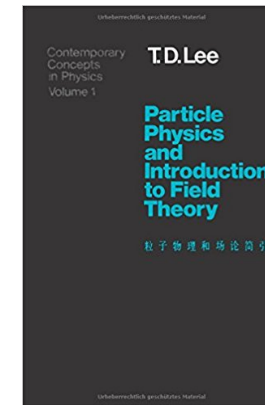
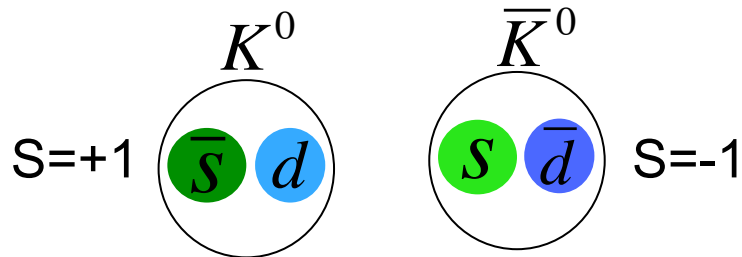


R. Feynman



“If there is any place where we have a chance to test the main principles of quantum mechanics in the purest way - does the superposition of amplitudes work or doesn't it? - this is it.”

“One of the most intriguing physical systems in Nature”

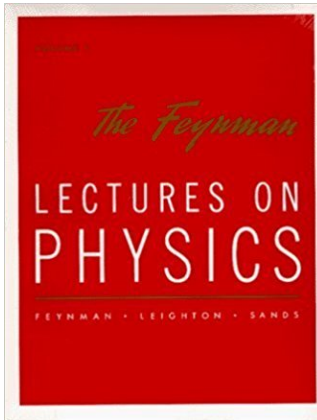


T. D. Lee

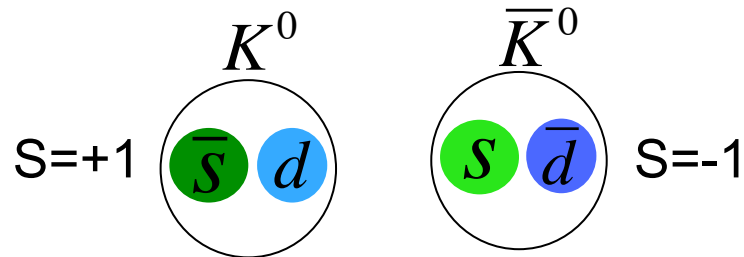
Neutral K meson system: a jewel donated to us by Nature



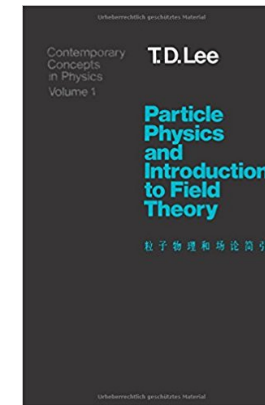
R. Feynman



“If there is any place where we have a chance to test the main principles of quantum mechanics in the purest way - does the superposition of amplitudes work or doesn't it? - this is it.”



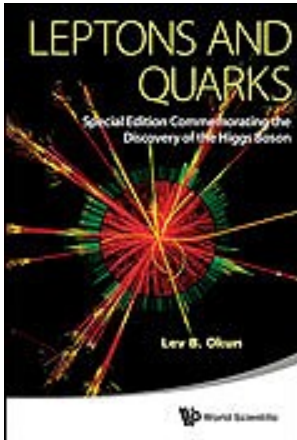
“One of the most intriguing physical systems in Nature”



T. D. Lee



Lev B. Okun

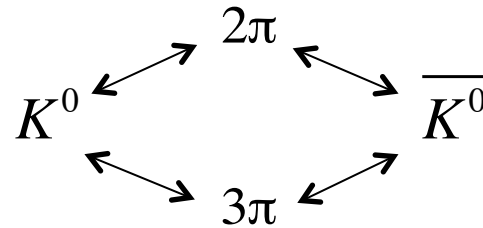


“Neutral K mesons are a unique physical system which appears to be created by nature to demonstrate, in the most impressive manner, a number of spectacular phenomena.”

“If the K mesons did not exist, they should have been invented ‘on purpose’ in order to teach students the principles of quantum mechanics”

The neutral kaon two-level oscillating system in a nutshell

K^0 and \bar{K}^0 can decay to common final states due to weak interactions:
strangeness oscillations



$$|\Psi\rangle = a|K^0\rangle + b|\bar{K}^0\rangle$$

$$i\frac{\partial}{\partial t}\Psi(t) = \mathbf{H}\Psi(t)$$

\mathbf{H} is the effective hamiltonian (non-hermitian), decomposed into a Hermitian part (mass matrix \mathbf{M}) and an anti-Hermitian part ($i/2$ decay matrix Γ):

$$\mathbf{H} = \mathbf{M} - \frac{i}{2}\Gamma = \begin{pmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{pmatrix} - \frac{i}{2} \begin{pmatrix} \Gamma_{11} & \Gamma_{12} \\ \Gamma_{21} & \Gamma_{22} \end{pmatrix}$$

Diagonalizing the effective Hamiltonian:

eigenvalues

$$\lambda_{S,L} = m_{S,L} - \frac{i}{2}\Gamma_{S,L}$$

$$|K_{S,L}(t)\rangle = e^{-i\lambda_{S,L}t}|K_{S,L}(0)\rangle$$

$$\tau_S \sim 90 \text{ ps} \quad \tau_L \sim 51 \text{ ns}$$

$K_L \rightarrow \pi\pi$ violates CP

eigenstates: physical states

$$|K_{S,L}\rangle = \frac{1}{\sqrt{2(1+|\varepsilon_{S,L}|)}} \left[(1 + \varepsilon_{S,L})|K^0\rangle \pm (1 - \varepsilon_{S,L})|\bar{K}^0\rangle \right]$$

$$= \frac{1}{\sqrt{(1+|\varepsilon_{S,L}|)}} \left[|K_{1,2}\rangle + \varepsilon_{S,L}|K_{2,1}\rangle \right]$$

$|K_{1,2}\rangle$ are
CP= ± 1 states

$$\langle K_S | K_L \rangle \cong \varepsilon_S^* + \varepsilon_L \neq 0$$

small CP impurity $\sim 2 \times 10^{-3}$

The neutral kaon two-level oscillating system in a nutshell

$$|K_{S,L}\rangle \propto \left[(1 + \varepsilon_{S,L}) |K^0\rangle \pm (1 - \varepsilon_{S,L}) |\bar{K}^0\rangle \right]$$

CP violation:

$$\varepsilon_{S,L} = \varepsilon \pm \delta$$

T violation:

$$\varepsilon = \frac{H_{12} - H_{21}}{2(\lambda_S - \lambda_L)} = \frac{-i\Im M_{12} - \Im \Gamma_{12}/2}{\Delta m + i\Delta\Gamma/2}$$

CPT violation:

$$\delta = \frac{H_{11} - H_{22}}{2(\lambda_S - \lambda_L)} = \frac{1}{2} \frac{(m_{\bar{K}^0} - m_{K^0}) - (i/2)(\Gamma_{\bar{K}^0} - \Gamma_{K^0})}{\Delta m + i\Delta\Gamma/2}$$

- $\delta \neq 0$ implies CPT violation
- $\varepsilon \neq 0$ implies T violation
- $\varepsilon \neq 0$ or $\delta \neq 0$ implies CP violation

$$\Delta m = m_L - m_S \quad , \quad \Delta\Gamma = \Gamma_S - \Gamma_L$$

$$\Delta m = 3.5 \times 10^{-15} \text{ GeV}$$

$$\text{(with a phase convention } \Im \Gamma_{12} = 0 \text{)} \quad \Delta\Gamma \approx \Gamma_S \approx 2\Delta m = 7 \times 10^{-15} \text{ GeV}$$

The neutral kaon two-level oscillating system in a nutshell

$$|K_{S,L}\rangle \propto \left[(1 + \varepsilon_{S,L}) |K^0\rangle \pm (1 - \varepsilon_{S,L}) |\bar{K}^0\rangle \right]$$

CP violation:

$$\varepsilon_{S,L} = \varepsilon \pm \delta$$

T violation:

$$\varepsilon = \frac{H_{12} - H_{21}}{2(\lambda_S - \lambda_L)} = \frac{-i\Im M_{12} - \Im \Gamma_{12}/2}{\Delta m + i\Delta\Gamma/2}$$

CPT violation:

$$\delta = \frac{H_{11} - H_{22}}{2(\lambda_S - \lambda_L)} = \frac{1}{2} \frac{(m_{\bar{K}^0} - m_{K^0}) - (i/2)(\Gamma_{\bar{K}^0} - \Gamma_{K^0})}{\Delta m + i\Delta\Gamma/2}$$

huge amplification factor!!

- $\delta \neq 0$ implies CPT violation
- $\varepsilon \neq 0$ implies T violation
- $\varepsilon \neq 0$ or $\delta \neq 0$ implies CP violation

$$\Delta m = m_L - m_S \quad , \quad \Delta\Gamma = \Gamma_S - \Gamma_L$$

$$\Delta m = 3.5 \times 10^{-15} \text{ GeV}$$

$$(\text{with a phase convention } \Im \Gamma_{12} = 0) \quad \Delta\Gamma \approx \Gamma_S \approx 2\Delta m = 7 \times 10^{-15} \text{ GeV}$$

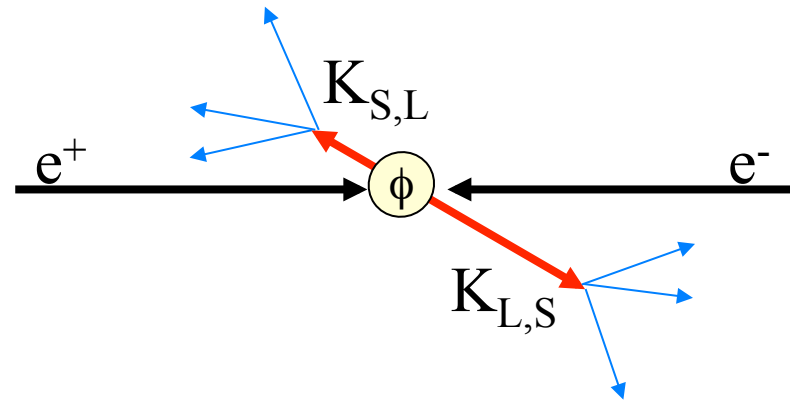
Neutral kaons at a ϕ -factory

Production of the vector meson ϕ in e^+e^- annihilations:

- $e^+e^- \rightarrow \phi$ $\sigma_\phi \sim 3 \mu\text{b}$
 $W = m_\phi = 1019.4 \text{ MeV}$
- $\text{BR}(\phi \rightarrow K^0\bar{K}^0) \sim 34\%$
- $\sim 10^6$ neutral kaon pairs per pb^{-1} produced in an antisymmetric quantum state with $J^{PC} = 1^{--}$:

$$\mathbf{p}_K = 110 \text{ MeV}/c$$

$$\lambda_S = 6 \text{ mm} \quad \lambda_L = 3.5 \text{ m}$$



$$|i\rangle = \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right]$$

$$= \frac{N}{\sqrt{2}} \left[|K_S(\vec{p})\rangle |K_L(-\vec{p})\rangle - |K_L(\vec{p})\rangle |K_S(-\vec{p})\rangle \right]$$

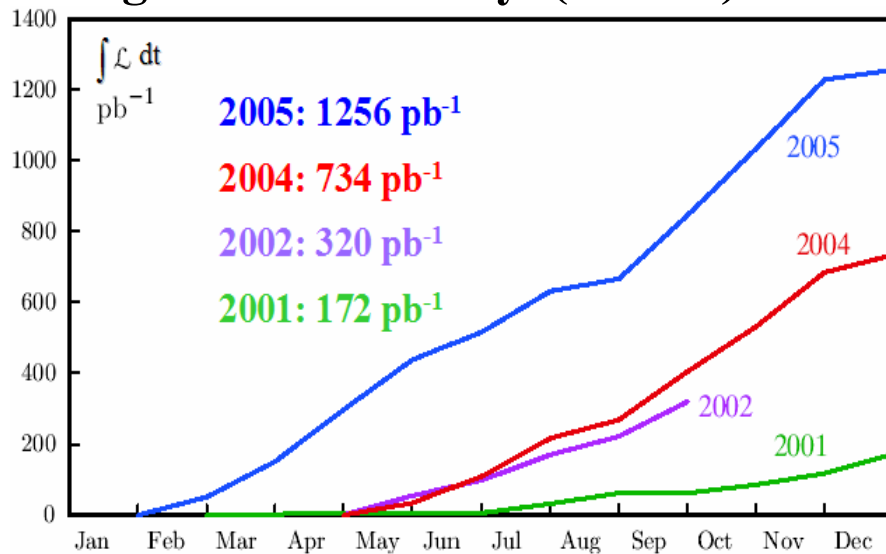
$$N = \sqrt{(1 + |\varepsilon_S|^2)(1 + |\varepsilon_L|^2)} / (1 - \varepsilon_S \varepsilon_L) \cong 1$$

The KLOE detector at the Frascati ϕ -factory DAFNE

DAFNE
collider

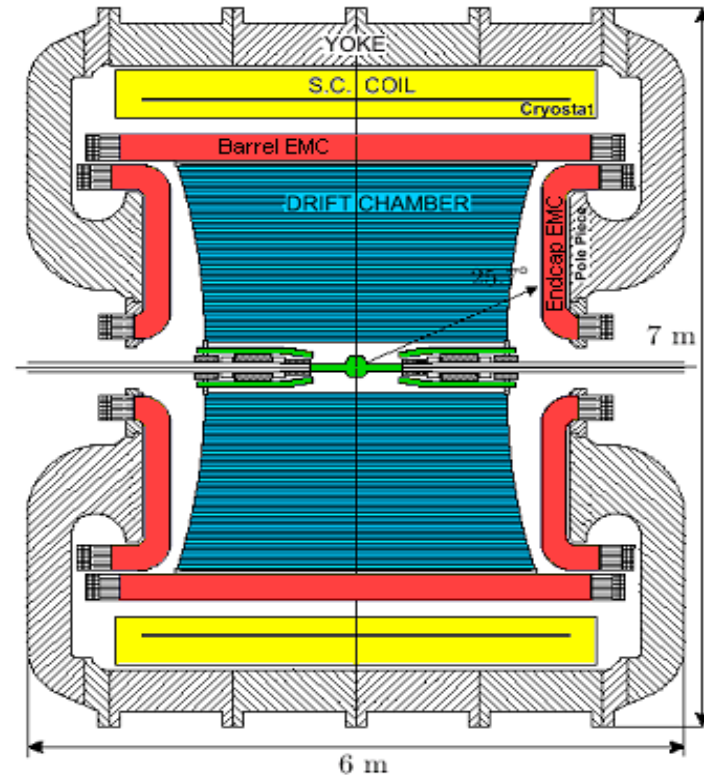


Integrated luminosity (KLOE)



Total KLOE $\int \mathcal{L} dt \sim 2.5 \text{ fb}^{-1}$
 (2001 - 05) $\rightarrow \sim 2.5 \times 10^9$ $K_S K_L$ pairs

KLOE detector



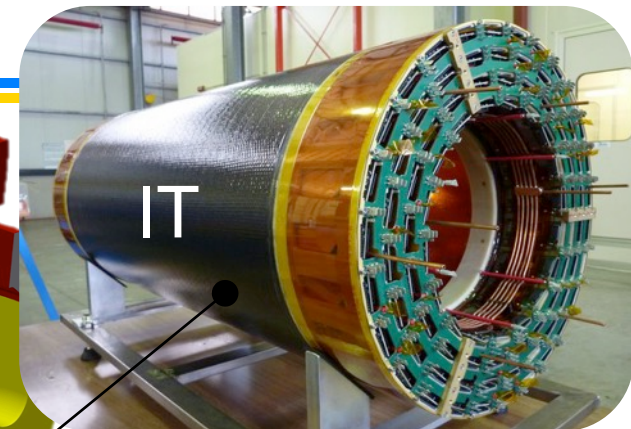
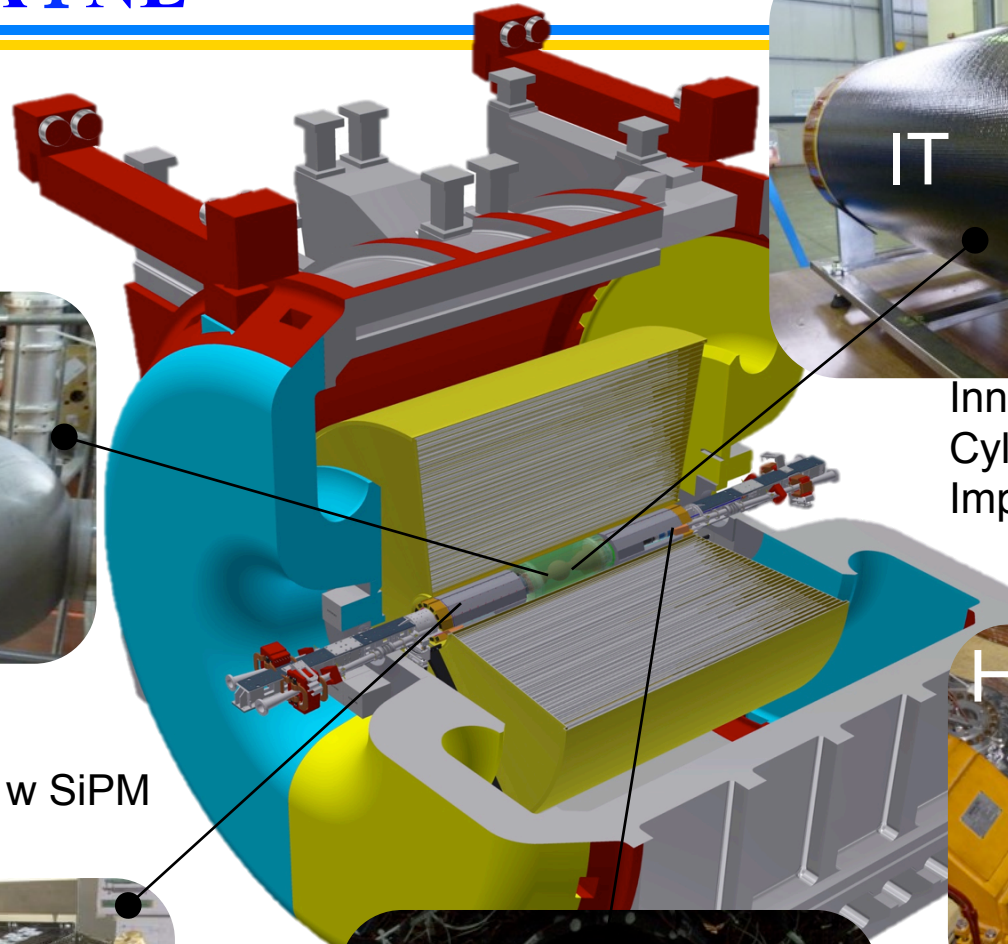
Lead/scintillating fiber calorimeter
 drift chamber
 4 m diameter \times 3.3 m length
 helium based gas mixture

KLOE-2 at DAΦNE

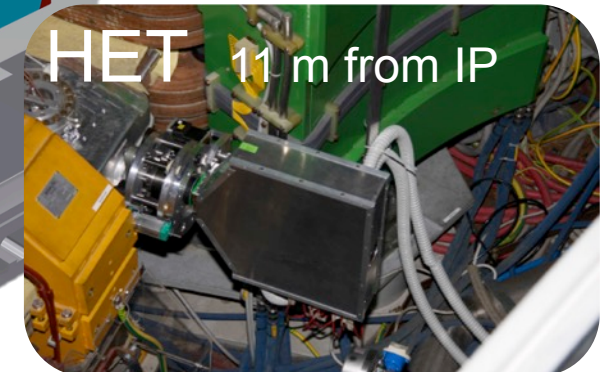
LYSO Crystal w SiPM
Low polar angle



Tungsten / Scintillating Tiles w SiPM
Quadrupole Instrumentation



Inner Tracker – 4 layers of
Cylindrical GEM detectors
Improve track and vtx reconstr.
First CGEM in HEP expt.

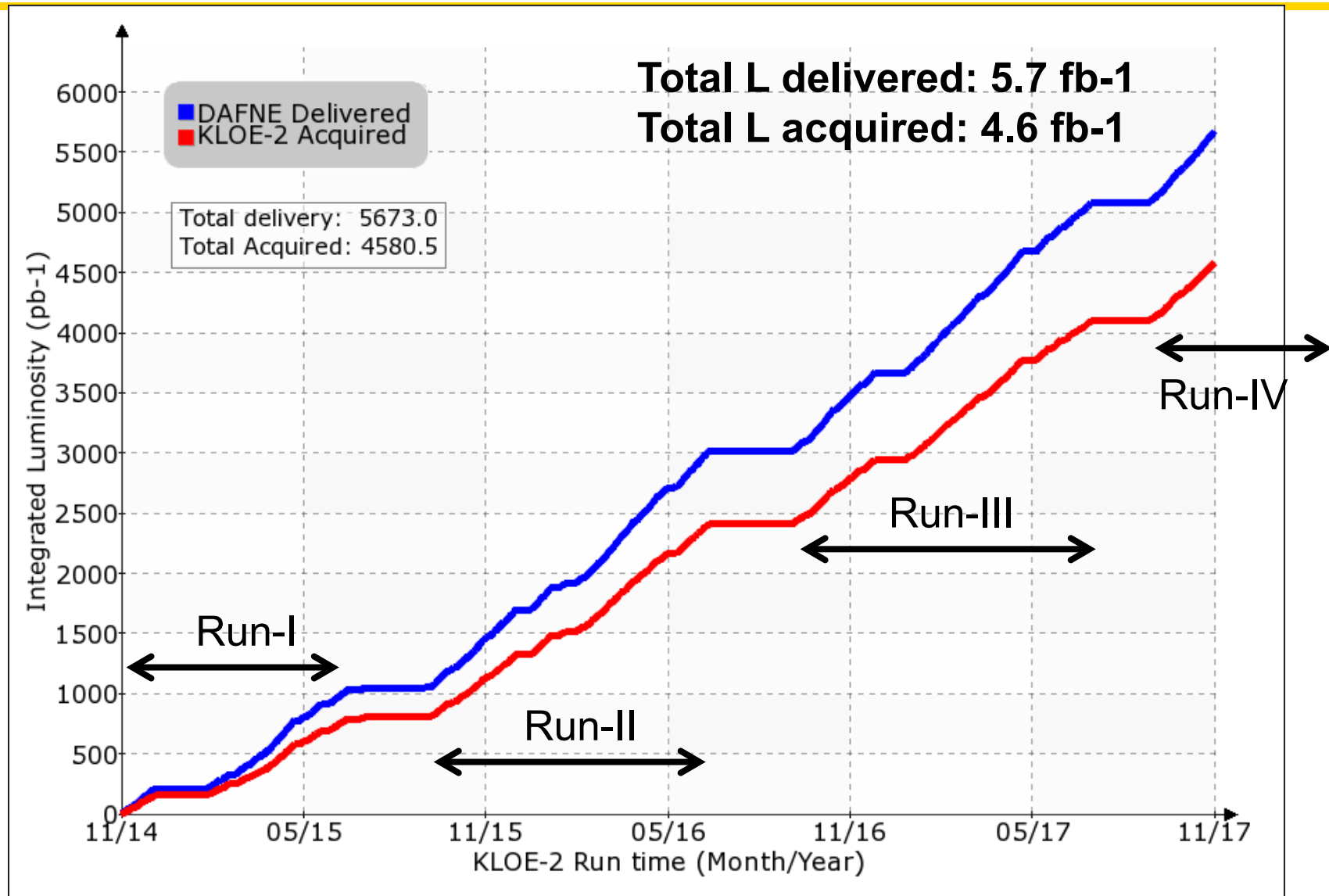


Scintillator hodoscope +PMTs



calorimeters LYSO+SiPMs
at ~ 1 m from IP

KLOE-2 Data Taking



KLOE-2 goal: L acquired > 5 fb⁻¹ => L delivered > ~ 6.2 fb⁻¹ by 31 March 2018

KLOE-2 Physics

KAON Physics:

- CPT and QM tests with kaon interferometry
- Direct T and CPT tests using entanglement
- CP violation and CPT test:
 $K_S \rightarrow 3\pi^0$
direct measurement of $\text{Im}(\varepsilon'/\varepsilon)$ (lattice calc. improved)
- CKM V_{us} :
 K_S semileptonic decays and A_S (also CP and CPT test)
 $K_{\mu 3}$ form factors, $Kl3$ radiative corrections
- χpT : $K_S \rightarrow \gamma\gamma$
- Search for rare K_S decays

Hadronic cross section

- Measurement of a_{μ}^{HLO} in the space-like region using Bhabha process
- ISR studies with 3π , 4π final states
- F_{π} with increased statistics

EPJC (2010) 68, 619 + procs LNF WS 2016 (in publication)

Dark forces:

- Improve limits on:
 $U\gamma$ associate production
 $e^+e^- \rightarrow U\gamma \rightarrow \pi\pi\gamma, \mu\mu\gamma$
- Higgstrahlung
 $e^+e^- \rightarrow Uh' \rightarrow \mu^+\mu^- + \text{miss. energy}$
- Leptophobic B boson search
 $\phi \rightarrow \eta B, B \rightarrow \pi^0\gamma, \eta \rightarrow \gamma\gamma$
 $\eta \rightarrow B\gamma, B \rightarrow \pi^0\gamma, \eta \rightarrow \pi^0\gamma\gamma$
- Search for U invisible decays

Light meson Physics:

- η decays, ω decays, TFF $\phi \rightarrow \eta e^+e^-$
- C,P,CP violation:
improve limits on $\eta \rightarrow \gamma\gamma, \pi^+\pi^-, \pi^0\pi^0, \pi^0\pi^0\gamma$
- improve $\eta \rightarrow \pi^+\pi^-e^+e^-$
- χpT : $\eta \rightarrow \pi^0\gamma\gamma$
- Light scalar mesons: $\phi \rightarrow K_S K_S \gamma$
- $\gamma\gamma$ Physics: $\gamma\gamma \rightarrow \pi^0$ and π^0 TFF
- light-by-light scattering
- axion-like particles

List of KLOE CP/CPT/QM tests with neutral kaons

Mode	Test	Param.	KLOE measurement
$K_L \rightarrow \pi^+\pi^-$	CP	BR	$(1.963 \pm 0.012 \pm 0.017) \times 10^{-3}$
$K_S \rightarrow 3\pi^0$	CP	BR	$< 2.6 \times 10^{-8}$
$K_S \rightarrow \pi e \nu$	CP	A_S	$(1.5 \pm 10) \times 10^{-3}$
$K_S \rightarrow \pi e \nu$	CPT	$\text{Re}(x_-)$	$(-0.8 \pm 2.5) \times 10^{-3}$
$K_S \rightarrow \pi e \nu$	CPT	$\text{Re}(y)$	$(0.4 \pm 2.5) \times 10^{-3}$
All $K_{S,L}$ BRs, η 's etc... (unitarity)	CP CPT	$\text{Re}(\epsilon)$ $\text{Im}(\delta)$	$(159.6 \pm 1.3) \times 10^{-5}$ $(0.4 \pm 2.1) \times 10^{-5}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	QM	ζ_{00}	$(0.1 \pm 1.0) \times 10^{-6}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	QM	ζ_{SL}	$(0.3 \pm 1.9) \times 10^{-2}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	CPT & QM	α	$(-10 \pm 37) \times 10^{-17} \text{ GeV}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	CPT & QM	β	$(1.8 \pm 3.6) \times 10^{-19} \text{ GeV}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	CPT & QM	γ	$(0.4 \pm 4.6) \times 10^{-21} \text{ GeV}$ compl. pos. hyp. $(0.7 \pm 1.2) \times 10^{-21} \text{ GeV}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	CPT & QM	$\text{Re}(\omega)$	$(-1.6 \pm 2.6) \times 10^{-4}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	CPT & QM	$\text{Im}(\omega)$	$(-1.7 \pm 3.4) \times 10^{-4}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	CPT & Lorentz	Δa_0	$(-6.2 \pm 8.8) \times 10^{-18} \text{ GeV}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	CPT & Lorentz	Δa_Z	$(-0.7 \pm 1.0) \times 10^{-18} \text{ GeV}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	CPT & Lorentz	Δa_X	$(3.3 \pm 2.2) \times 10^{-18} \text{ GeV}$
$K_S K_L \rightarrow \pi^+\pi^-, \pi^+\pi^-$	CPT & Lorentz	Δa_Y	$(-0.7 \pm 2.0) \times 10^{-18} \text{ GeV}$

K_S semileptonic charge asymmetry CP and CPT test

K_S semileptonic charge asymmetry

K_S and K_L semileptonic charge asymmetry

$$A_{S,L} = \frac{\Gamma(K_{S,L} \rightarrow \pi^- e^+ \nu) - \Gamma(K_{S,L} \rightarrow \pi^+ e^- \bar{\nu})}{\Gamma(K_{S,L} \rightarrow \pi^- e^+ \nu) + \Gamma(K_{S,L} \rightarrow \pi^+ e^- \bar{\nu})} = 2\Re \varepsilon \pm 2\Re \delta - 2\Re y \pm 2\Re x_-$$

T CPT viol. in mixing
 \downarrow \downarrow
 CPTV in $\Delta S = \Delta Q$ $\Delta S \neq \Delta Q$ decays

$A_{S,L} \neq 0$ signals CP violation

$A_S \neq A_L$ signals CPT violation

$$A_L = (3.322 \pm 0.058 \pm 0.047) \times 10^{-3}$$

KTEV PRL88,181601(2002)

$$A_S = (1.5 \pm 9.6 \pm 2.9) \times 10^{-3}$$

KLOE PLB 636(2006) 173

Data sample: L=410 pb⁻¹

$$A_S - A_L = 4(\Re \delta + \Re x_-)$$

$$\Re x_- = (-0.8 \pm 2.5) \times 10^{-3}$$

CPT & $\Delta S = \Delta Q$ viol.

$$A_S + A_L = 4(\Re \varepsilon - \Re y)$$

$$\Re y = (0.4 \pm 2.5) \times 10^{-3}$$

CPT viol.

input from other experiments

KLOE PLB 636(2006) 173

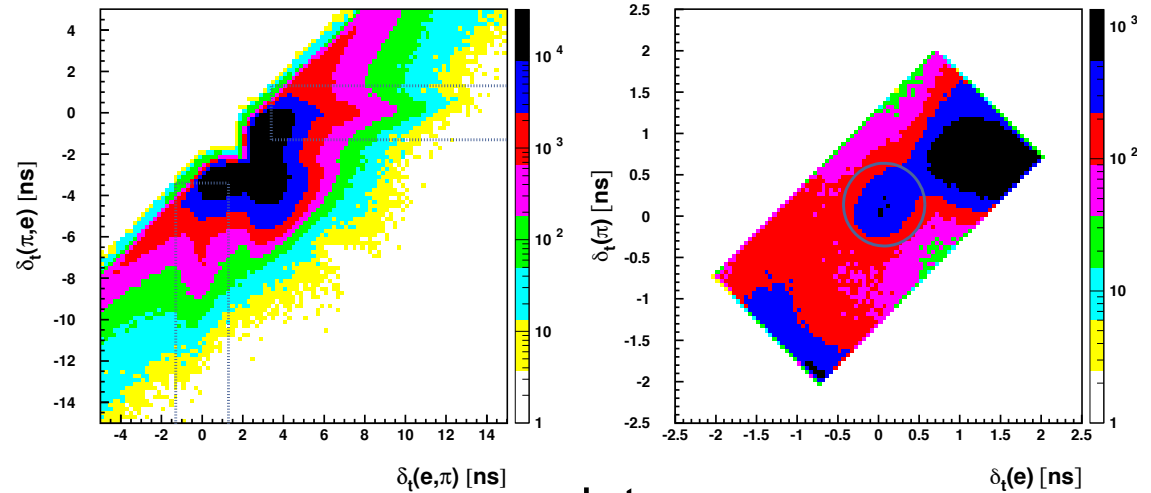
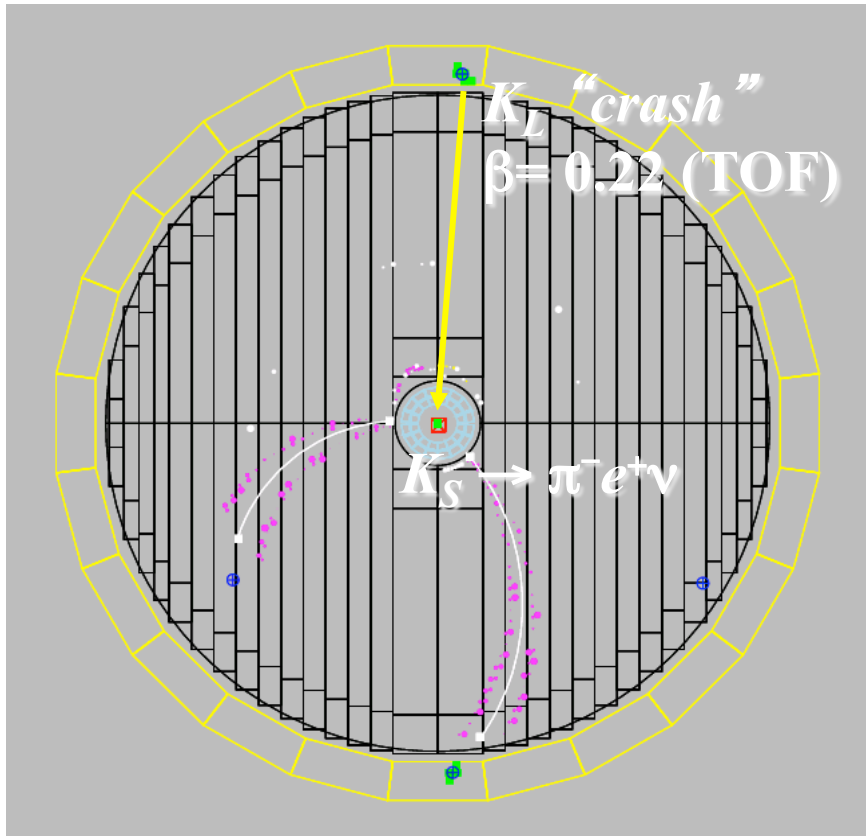
K_S semileptonic charge asymmetry

$$|i\rangle \propto [|K_S\rangle |K_L\rangle - |K_L\rangle |K_S\rangle]$$

- KLOE 1.7 fb⁻¹; ~ 4 × statistics w.r.t. previous measurement

- Pre-selection
- PID with time of flight technique

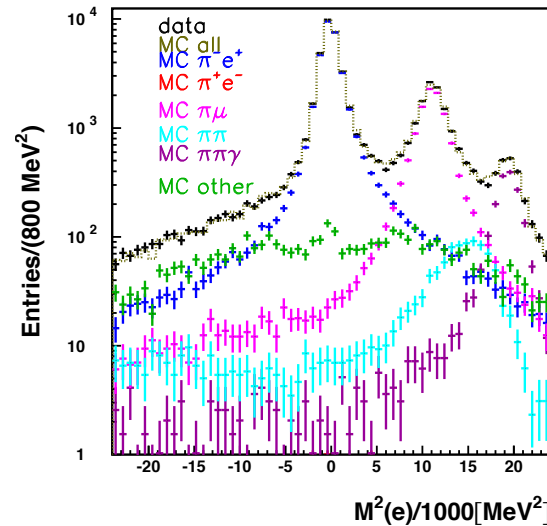
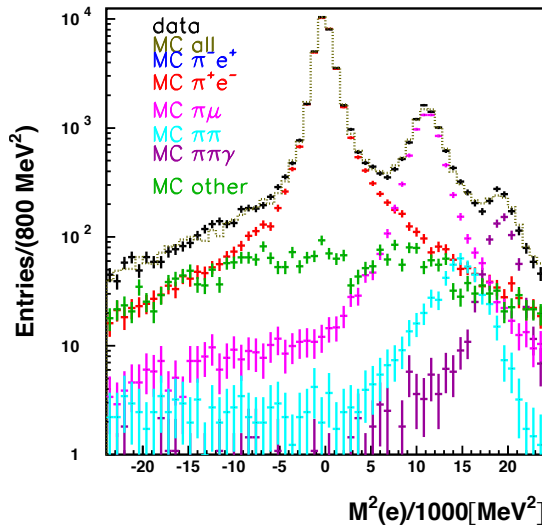
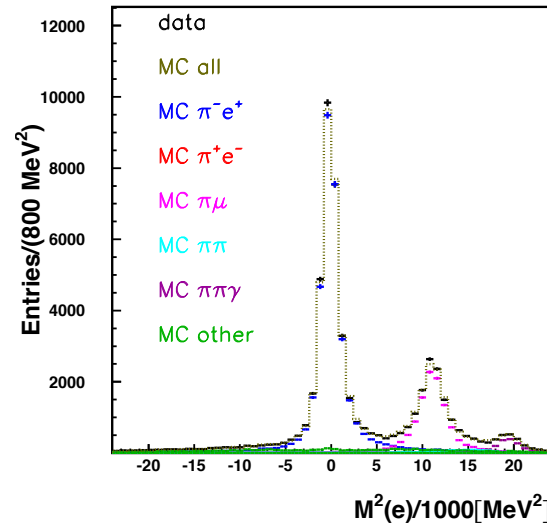
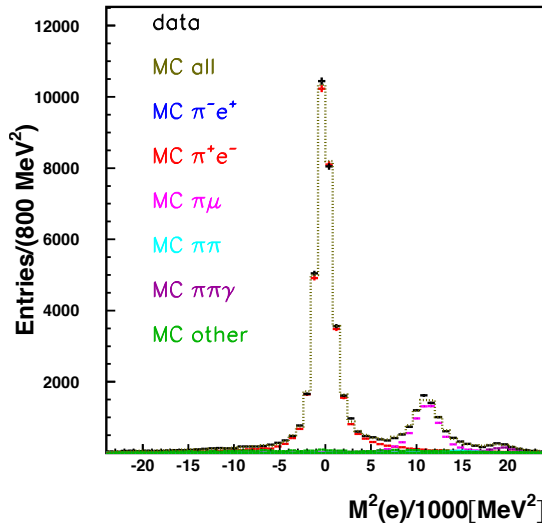
$$\delta t(m_X) = t_{cl} - \frac{L}{c\beta(m_X)} \quad \delta_{t,ab} = \delta t(m_a) - \delta t(m_b)$$



data

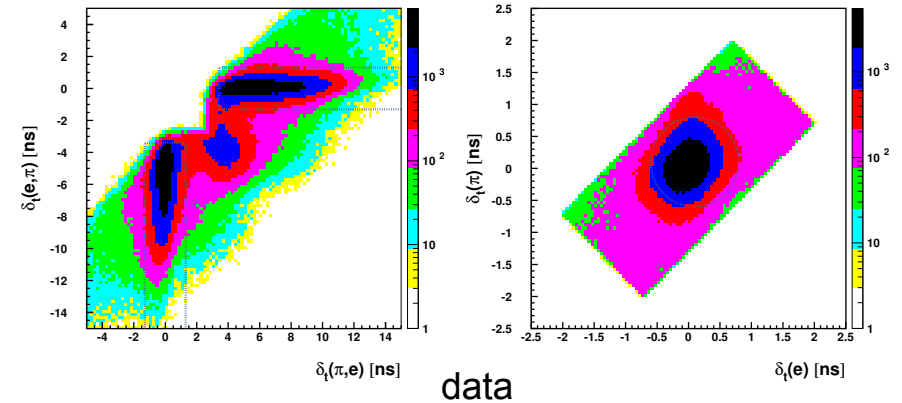
K_S tagged by K_L interaction in EmC
Efficiency ~ 30% (largely geometrical)

K_S semileptonic charge asymmetry

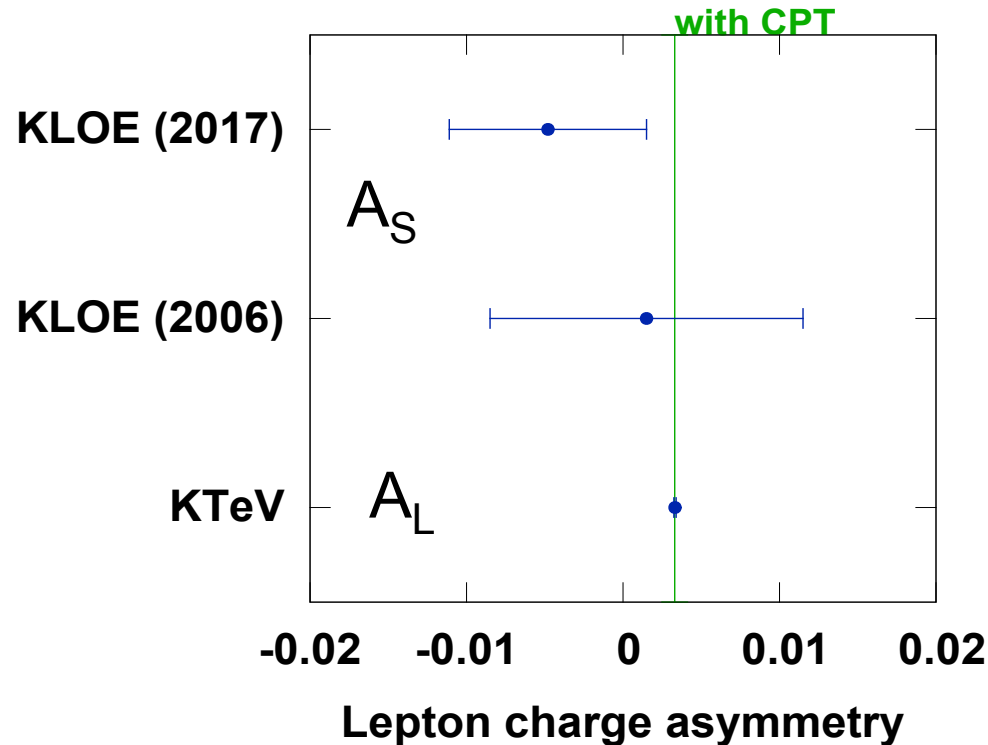


- Fit of $M^2(e)$ distribution varying MC normalizations of signal and bkg contributions
- Control sample: $K_L \rightarrow \pi e \nu$ close to IP tagged by $K_S \rightarrow \pi^0 \pi^0$
- track to EMC cluster and TOF efficiency correction from data c.s.

$K_L \rightarrow \pi e \nu$



K_S semileptonic charge asymmetry at KLOE



Data sample: $L=1.7 \text{ fb}^{-1}$

FINAL

KLOE (2017)

$$A_S = (-4.8 \pm 5.7 \pm 2.6) \times 10^{-3}$$

It will improve the CPT test ($\text{Im}\delta$) using Bell-Steinberger relationship

$$\text{with KLOE-2 data: } \delta A_S(\text{stat}) \rightarrow \sim 3 \times 10^{-3}$$

$$A_S - A_L = 4(\Re\delta + \Re x_-) \longrightarrow \Re x_- = (-2.3 \pm 1.6) \times 10^{-3} \quad \text{CPT \& } \Delta S = \Delta Q \text{ viol.}$$

$$A_S + A_L = 4(\Re\epsilon - \Re y) \longrightarrow \Re y = (2.0 \pm 1.6) \times 10^{-3} \quad \text{CPT viol.}$$

input from other experiments

Direct test of CPT and T in neutral kaon transitions

Testing CPT

The three discrete symmetries of QM, C (charge conjugation: $q \rightarrow -q$), P (parity: $x \rightarrow -x$), and T (time reversal: $t \rightarrow -t$) are known to be violated in nature both singly and in pairs. Only CPT appears to be an exact symmetry of nature.

CPT theorem holds for any QFT formulated on flat space-time which assumes: (1) Lorentz invariance (2) Locality (3) Unitarity (i.e. conservation of probability).

Extension of CPT theorem to a theory of quantum gravity far from obvious.
(e.g. CPT violation appears in several QG models)

huge effort in the last decades to study and shed light on QG phenomenology
 \Rightarrow Phenomenological CPTV parameters to be constrained by experiments

Consequences of CPT symmetry: equality of masses, lifetimes, $|q|$ and $|\mu|$ of a particle and its anti-particle.

Neutral meson systems offer unique possibilities to test CPT invariance; e.g. taking as figure of merit the fractional difference between the masses of a particle and its anti-particle:

neutral K system	neutral B system	proton- anti-proton
$\left m_{K^0} - m_{\bar{K}^0} \right / m_K < 10^{-18}$	$\left m_{B^0} - m_{\bar{B}^0} \right / m_B < 10^{-14}$	$\left m_p - m_{\bar{p}} \right / m_p < 10^{-8}$

Other interesting CPT tests: e.g. the study of anti-hydrogen atoms, etc..

Testing CPT

The three discrete symmetries of QM, C (charge conjugation: $q \rightarrow -q$), P (parity: $x \rightarrow -x$), and T (time reversal: $t \rightarrow -t$) are known to be violated in nature both singly and in pairs. Only CPT appears to be an exact symmetry of nature.

CPT theorem holds for any QFT formulated on flat space-time which assumes:
(1) Lorentz invariance (2) Locality (3) Unitarity (i.e. conservation of probability).

Extension of CPT theorem to a theory of quantum gravity far from obvious.
(e.g. CPT violation appears in several QG models)

huge effort in the last decades to study and shed light on QG phenomenology
 \Rightarrow Phenomenological CPTV parameters to be constrained by experiments

Consequences of CPT symmetry: equality of masses, lifetimes, $|q|$ and $|\mu|$ of a particle and its anti-particle.

Neutral meson systems offer unique possibilities to test CPT invariance; e.g. taking as figure of merit the fractional difference between the masses of a particle and its anti-particle:

neutral K system

$$\left| \frac{m_{K^0} - m_{\bar{K}^0}}{m_K} \right| < 10^{-18}$$

neutral B system

$$\left| \frac{m_{B^0} - m_{\bar{B}^0}}{m_B} \right| < 10^{-14}$$

proton- anti-proton

$$\left| \frac{m_p - m_{\bar{p}}}{m_p} \right| < 10^{-8}$$

Other interesting CPT tests: e.g. the study of anti-hydrogen atoms, etc..

Direct CPT test in transitions

- Is it possible to test the CPT symmetry directly in transition processes between kaon states, rather than comparing masses, lifetimes, or other intrinsic properties of particle and anti-particle states?
- CPT violating effects may not appear at first order in diagonal mass terms (K^0 vs \bar{K}^0 survival probabilities) while they can manifest at first order in transitions (non-diagonal terms).
- In standard WWA the test is related to $\text{Re}\delta$, a genuine CPT violating effect independent of $\Delta\Gamma$, i.e. not requiring the decay as an essential ingredient.
- Clean formulation required. Possible spurious effects induced by CP violation in the decay and/or a violation of the $\Delta S = \Delta Q$ rule have to be well under control.

Probing CPT: J. Bernabeu, A.D.D., P. Villanueva, JHEP 10 (2015) 139

Time-reversal violation: J. Bernabeu, A.D.D., P. Villanueva, NPB 868 (2013) 102

Definition of states

Let us also consider the states $|K_+\rangle$, $|K_-\rangle$ defined as follows: $|K_+\rangle$ is the state filtered by the decay into $\pi\pi$ ($\pi^+\pi^+$ or $\pi^0\pi^0$), a pure CP = +1 state; analogously $|K_-\rangle$ is the state filtered by the decay into $3\pi^0$, a pure CP = -1 state. Their orthogonal states correspond to the states which cannot decay into $\pi\pi$ or $3\pi^0$, defined, respectively, as

$$\begin{aligned} |\tilde{K}_-\rangle &\equiv \tilde{N}_- [|K_L\rangle - \eta_{\pi\pi}|K_S\rangle] \\ |\tilde{K}_+\rangle &\equiv \tilde{N}_+ [|K_S\rangle - \eta_{3\pi^0}|K_L\rangle] \end{aligned}$$

$$\begin{aligned} \eta_{\pi\pi} &= \frac{\langle \pi\pi | T | K_L \rangle}{\langle \pi\pi | T | K_S \rangle} \\ \eta_{3\pi^0} &= \frac{\langle 3\pi^0 | T | K_S \rangle}{\langle 3\pi^0 | T | K_L \rangle} \end{aligned}$$

Orthogonal bases: $\{K_+, \tilde{K}_-\}$ $\{\tilde{K}_+, K_-\}$

Even though the decay products are orthogonal, the filtered $|K_+\rangle$ and $|K_-\rangle$ states can still be non-orthogonal.

Condition of orthogonality:

$$\eta_{\pi\pi} + \eta_{3\pi^0}^* = \epsilon_L + \epsilon_S^* \longrightarrow \begin{aligned} |K_+\rangle &\equiv |\tilde{K}_+\rangle \\ |K_-\rangle &\equiv |\tilde{K}_-\rangle \end{aligned}$$

Neglect direct CP violation. Similarly any $\Delta S = \Delta Q$ rule violation for $|K^0\rangle$ and $|\bar{K}^0\rangle$

Direct test of CPT symmetry in neutral kaon transitions

CPT symmetry test

Reference		\mathcal{CPT} -conjugate	
Transition	Decay products	Transition	Decay products
$K^0 \rightarrow K_+$	$(\ell^-, \pi\pi)$	$K_+ \rightarrow \bar{K}^0$	$(3\pi^0, \ell^-)$
$K^0 \rightarrow K_-$	$(\ell^-, 3\pi^0)$	$K_- \rightarrow \bar{K}^0$	$(\pi\pi, \ell^-)$
$\bar{K}^0 \rightarrow K_+$	$(\ell^+, \pi\pi)$	$K_+ \rightarrow K^0$	$(3\pi^0, \ell^+)$
$\bar{K}^0 \rightarrow K_-$	$(\ell^+, 3\pi^0)$	$K_- \rightarrow K^0$	$(\pi\pi, \ell^+)$

One can define the following ratios of probabilities:

$$R_{1,\mathcal{CPT}}(\Delta t) = P [K_+(0) \rightarrow \bar{K}^0(\Delta t)] / P [K^0(0) \rightarrow K_+(\Delta t)]$$

$$R_{2,\mathcal{CPT}}(\Delta t) = P [K^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow \bar{K}^0(\Delta t)]$$

$$R_{3,\mathcal{CPT}}(\Delta t) = P [K_+(0) \rightarrow K^0(\Delta t)] / P [\bar{K}^0(0) \rightarrow K_+(\Delta t)]$$

$$R_{4,\mathcal{CPT}}(\Delta t) = P [\bar{K}^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow K^0(\Delta t)]$$

Any deviation from $R_{i,\mathcal{CPT}}=1$ constitutes a violation of CPT-symmetry

J. Bernabeu, A.D.D., P. Villanueva, JHEP 10 (2015) 139

Direct test of symmetries with neutral kaons

Reference	T -conjugate	CP -conjugate	CPT -conjugate
$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$
$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$
$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$
$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$
$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$
$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$
$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$
$K_+ \rightarrow K_-$	$K_- \rightarrow K_+$	$K_+ \rightarrow K_-$	$K_- \rightarrow K_+$
$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$
$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$
$K_- \rightarrow K_+$	$K_+ \rightarrow K_-$	$K_- \rightarrow K_+$	$K_+ \rightarrow K_-$
$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$

Direct test of symmetries with neutral kaons

Conjugate=
reference

Reference	T -conjugate	CP -conjugate	CPT -conjugate
$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$
$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$
$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$
$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$
$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$
$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$
$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$
$K_+ \rightarrow K_-$	$K_- \rightarrow K_+$	$K_- \rightarrow K_-$	$K_- \rightarrow K_+$
$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$
$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$
$K_- \rightarrow K_+$	$K_+ \rightarrow K_-$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_-$
$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$

Direct test of symmetries with neutral kaons

Conjugate=
reference

already in the
table with
conjugate as
reference

Reference	<i>T</i> -conjugate	<i>CP</i> -conjugate	<i>CPT</i> -conjugate
$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$
$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$
$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$
$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$
$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$
$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$
$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$
$K_+ \rightarrow K_-$	$K_- \rightarrow K_+$	$K_- \rightarrow K_-$	$K_- \rightarrow K_+$
$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$
$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$
$K_- \rightarrow K_+$	$K_+ \rightarrow K_-$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_-$
$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$

Direct test of symmetries with neutral kaons

Conjugate=
reference



already in the
table with
conjugate as
reference



Two identical
conjugates
for one reference



Reference	<i>T</i> -conjugate	<i>CP</i> -conjugate	<i>CPT</i> -conjugate
$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$
$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$
$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$
$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$
$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$
$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$
$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$
$K_+ \rightarrow K_-$	$K_- \rightarrow K_+$	$K_+ \rightarrow K_-$	$K_- \rightarrow K_+$
$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$
$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$
$K_- \rightarrow K_+$	$K_+ \rightarrow K_-$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_-$
$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$

Direct test of symmetries with neutral kaons

Conjugate=
reference

Reference	<i>T</i> -conjugate	<i>CP</i> -conjugate	<i>CPT</i> -conjugate
$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$	$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$
$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$
$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$
$\bar{K}^0 \rightarrow K^0$	$K^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow \bar{K}^0$	$K^0 \rightarrow K^0$	$K^0 \rightarrow K^0$
$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow K^0$
$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow K^0$
$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$	$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$
$K_+ \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_+$	$K_+ \rightarrow K^0$	$K^0 \rightarrow K_+$
$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_+$
$K_+ \rightarrow K_-$	$K_- \rightarrow K_+$	$K_+ \rightarrow K_-$	$K_- \rightarrow K_+$
$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$	$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$
$K_- \rightarrow \bar{K}^0$	$\bar{K}^0 \rightarrow K_-$	$K_- \rightarrow K^0$	$K^0 \rightarrow K_-$
$K_- \rightarrow K_+$	$K_+ \rightarrow K_-$	$K_+ \rightarrow K_+$	$K_+ \rightarrow K_-$
$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$	$K_- \rightarrow K_-$

already in the
table with
conjugate as
reference

4 distinct tests
of *T* symmetry

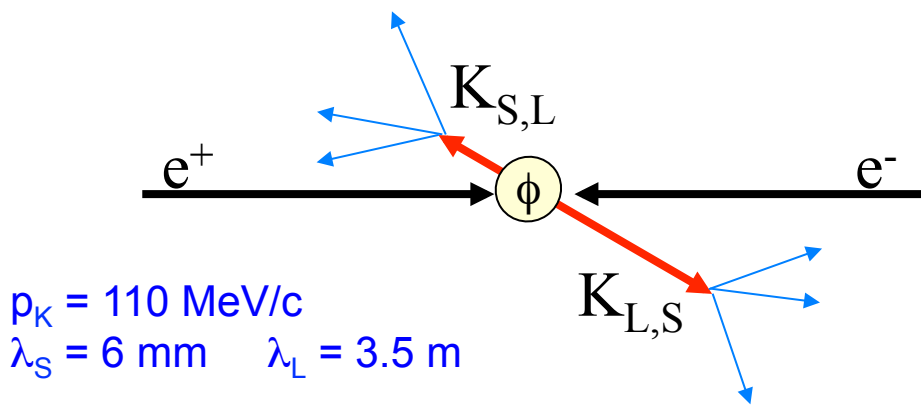
4 distinct tests
of *CP* symmetry

4 distinct tests
of *CPT* symmetry

Two identical
conjugates
for one reference

Quantum entanglement as a tool

- The in \leftrightarrow out states inversion required in a DIRECT test of CPT (or T) can be performed exploiting the properties of the quantum entanglement.
- In maximally entangled systems the complete knowledge of the system as a whole is encoded in the state, no information on single subsystems is available.
- Once a measurement is performed on one subsystem, then the information is immediately transferred to its partner, which is prepared in the orthogonal state
- $\sigma(e^+e^- \rightarrow \phi) \sim 3 \text{ mb}$; $W = m_\phi = 1019.4 \text{ MeV}$ $BR(\phi \rightarrow K^0\bar{K}^0) \sim 34\%$
 $\sim 10^6/\text{pb}^{-1}$ KK pairs produced in an antisymmetric quantum state with $J^{PC} = 1^{--}$:



$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{N}{\sqrt{2}} \left[|K_S(\vec{p})\rangle |K_L(-\vec{p})\rangle - |K_L(\vec{p})\rangle |K_S(-\vec{p})\rangle \right]
 \end{aligned}$$

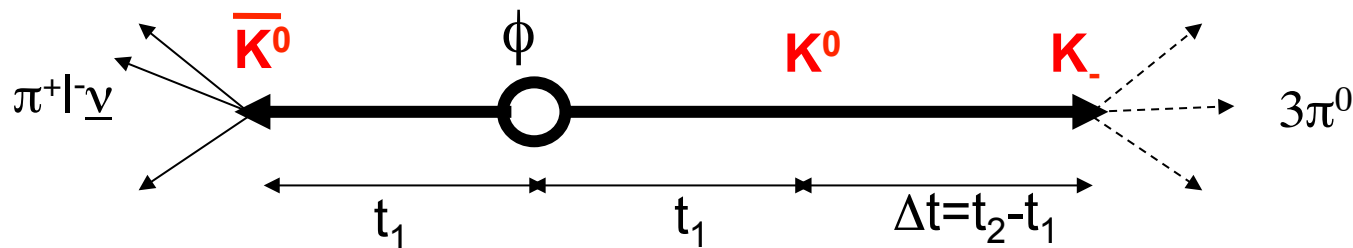
$$N = \sqrt{\frac{(1 + |\epsilon_S|^2)(1 + |\epsilon_L|^2)}{(1 - \epsilon_S \epsilon_L)}} \cong 1$$

Entanglement in neutral kaon pairs

- EPR correlations at a ϕ -factory can be exploited to study transitions involving orthogonal “CP states” K_+ and K_-

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[|K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(-\vec{p})\rangle |K_+(\vec{p})\rangle \right]
 \end{aligned}$$

- decay as filtering measurement
- entanglement \rightarrow preparation of state

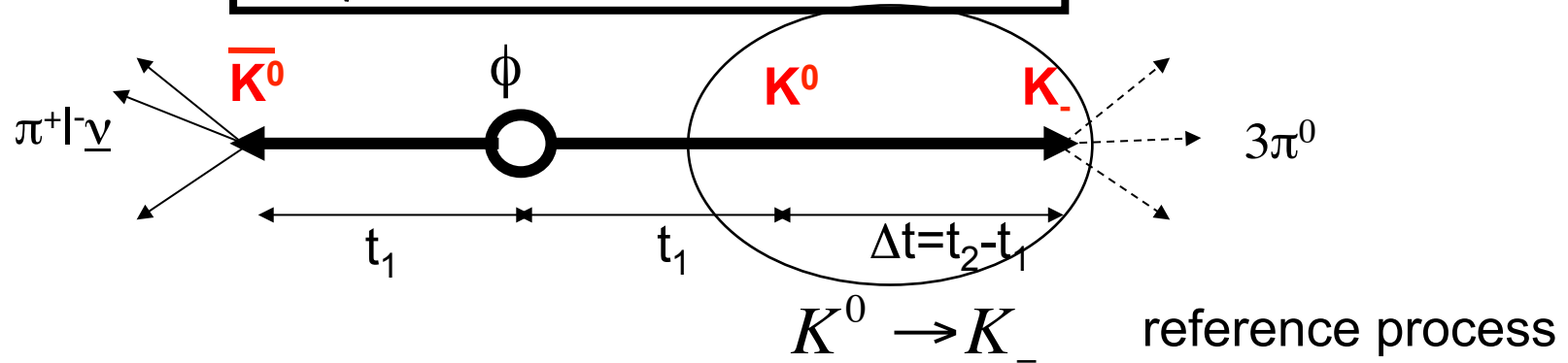


Entanglement in neutral kaon pairs

- EPR correlations at a ϕ -factory can be exploited to study transitions involving orthogonal “CP states” K_+ and K_-

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[|K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(-\vec{p})\rangle |K_+(\vec{p})\rangle \right]
 \end{aligned}$$

- decay as filtering measurement
- entanglement \rightarrow preparation of state

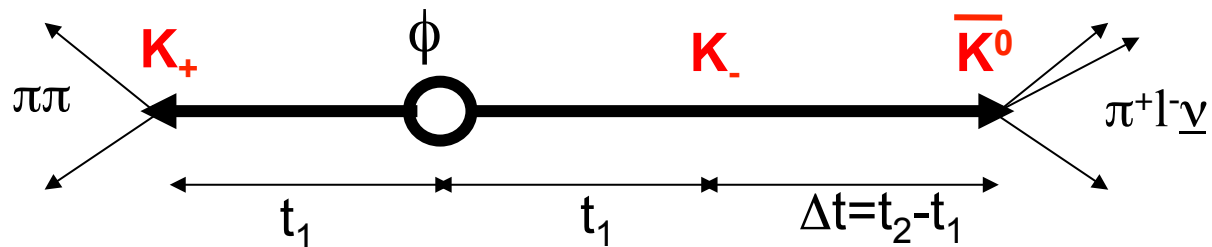
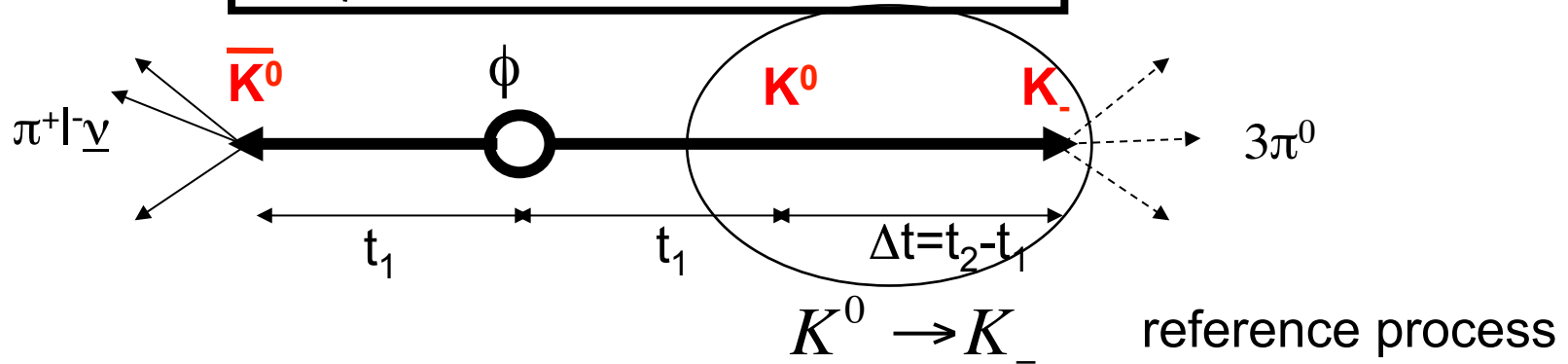


Entanglement in neutral kaon pairs

- EPR correlations at a ϕ -factory can be exploited to study transitions involving orthogonal “CP states” K_+ and K_-

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[|K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(\vec{p})\rangle |K_+(-\vec{p})\rangle \right]
 \end{aligned}$$

- decay as filtering measurement
- entanglement \rightarrow preparation of state

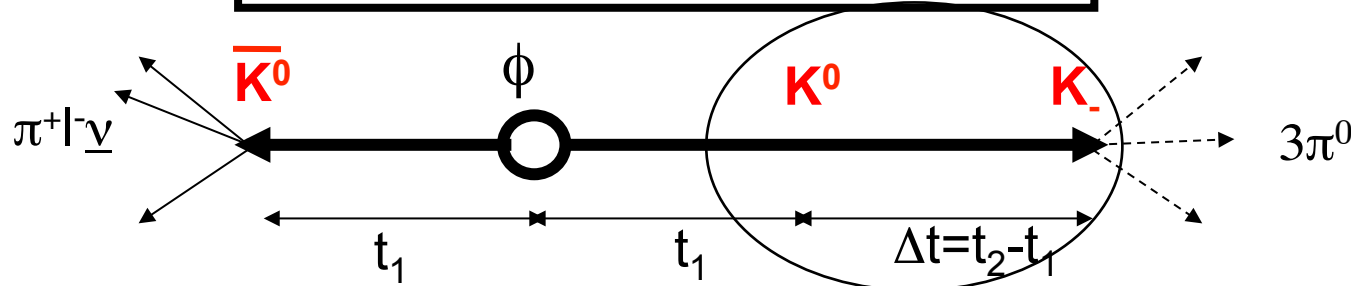


Entanglement in neutral kaon pairs

- EPR correlations at a ϕ -factory can be exploited to study transitions involving orthogonal “CP states” K_+ and K_-

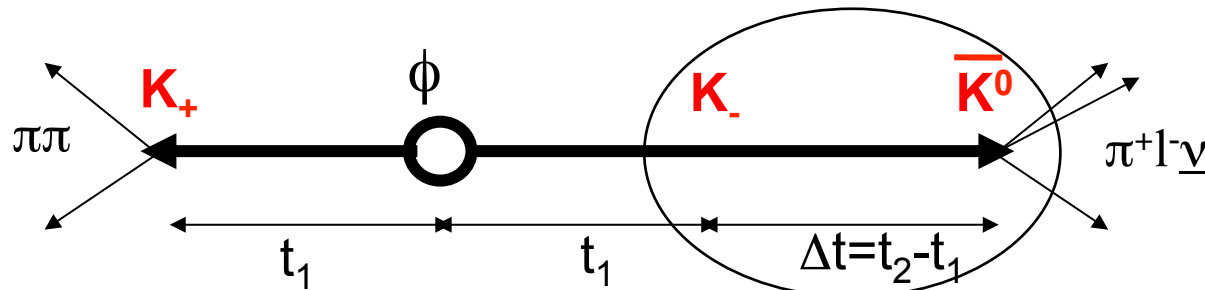
$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[|K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(-\vec{p})\rangle |K_+(\vec{p})\rangle \right]
 \end{aligned}$$

- decay as filtering measurement
- entanglement \rightarrow preparation of state



$K^0 \rightarrow K_-$ reference process

$K_- \rightarrow \bar{K}^0$ CPT-conjugated process

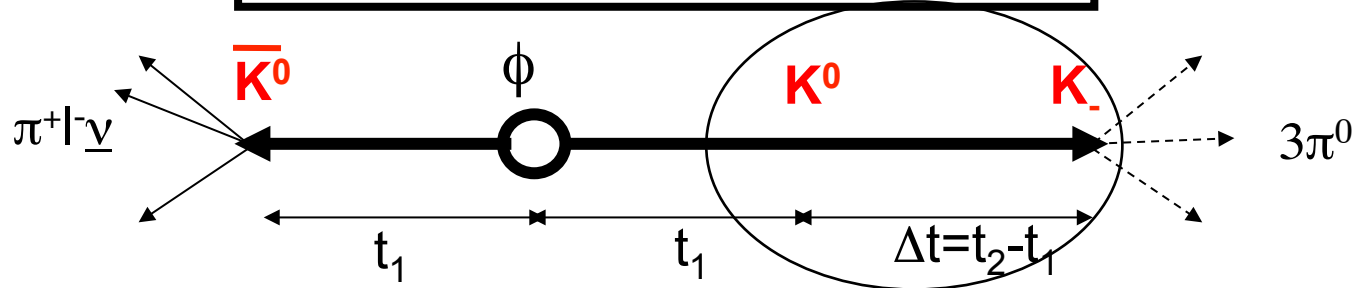


Entanglement in neutral kaon pairs

- EPR correlations at a ϕ -factory can be exploited to study transitions involving orthogonal “CP states” K_+ and K_-

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[|K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(-\vec{p})\rangle |K_+(\vec{p})\rangle \right]
 \end{aligned}$$

- decay as filtering measurement
- entanglement \rightarrow preparation of state

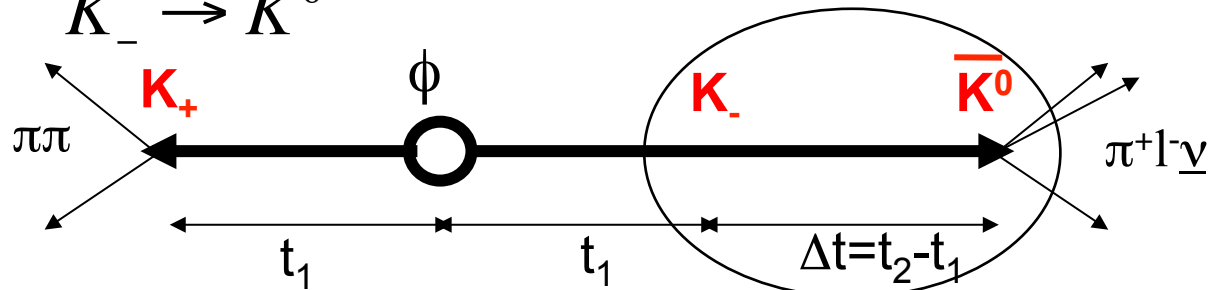


$K^0 \rightarrow K_-$ reference process

Note: CP and T conjugated process

$$\bar{K}^0 \rightarrow K_- \quad K_- \rightarrow K^0$$

$K_- \rightarrow \bar{K}^0$ CPT-conjugated process



Direct test of CPT symmetry in neutral kaon transitions

Two observable ratios of double decay intensities

$$R_{2,\text{CPT}}^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^-, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^-; \Delta t)}$$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^+, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^+; \Delta t)}$$

for $\Delta t > 0$ $R_{2,\text{CPT}}^{\text{exp}}(\Delta t) = R_{2,\text{CPT}}(\Delta t) \times D_{\text{CPT}}$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) = R_{4,\text{CPT}}(\Delta t) \times D_{\text{CPT}}$$

for $\Delta t < 0$ $R_{2,\text{CPT}}^{\text{exp}}(\Delta t) = R_{1,\text{CPT}}(|\Delta t|) \times D_{\text{CPT}}$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) = R_{3,\text{CPT}}(|\Delta t|) \times D_{\text{CPT}}$$

with D_{CPT} constant $D_{\text{CPT}} = \frac{\text{BR}(K_L \rightarrow 3\pi^0) \Gamma_L}{\text{BR}(K_S \rightarrow \pi\pi) \Gamma_S}$

Direct test of CPT symmetry in neutral kaon transitions

Explicitly in standard Wigner Weisskopf approach for $\Delta t > 0$:

$$R_{2,\text{CPT}}^{\text{exp}}(\Delta t) = \frac{P[\text{K}^0(0) \rightarrow \text{K}_-(\Delta t)]}{P[\text{K}_-(0) \rightarrow \bar{\text{K}}^0(\Delta t)]} \times D_{\text{CPT}}$$

$$\simeq |1 - 2\delta|^2 \left| 1 + 2\delta e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\text{CPT}}$$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) = \frac{P[\bar{\text{K}}^0(0) \rightarrow \text{K}_-(\Delta t)]}{P[\text{K}_-(0) \rightarrow \text{K}^0(\Delta t)]} \times D_{\text{CPT}}$$

$$\simeq |1 + 2\delta|^2 \left| 1 - 2\delta e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\text{CPT}}$$

For comparison the ratio of survival probabilities: Vanishes for $\Delta\Gamma \rightarrow 0$

$$\frac{I(\ell^-, \ell^+; \Delta t)}{I(\ell^+, \ell^-; \Delta t)} = \frac{P[\text{K}^0(0) \rightarrow \text{K}^0(\Delta t)]}{P[\bar{\text{K}}^0(0) \rightarrow \bar{\text{K}}^0(\Delta t)]}$$

$$\simeq |1 - 4\delta|^2 \left| 1 + \frac{8\delta}{1 + e^{+i(\lambda_S - \lambda_L)\Delta t}} \right|^2$$

As an illustration of the different sensitivity: it vanishes up to second order in CPTV and decoherence parameters α, β, γ (Ellis, Mavromatos et al. PRD1996)

Impact of the approximations

In general K_+ and K_-
(and K_0 and \tilde{K}_0)
can be non-orthogonal

Direct CP (CPT) violation

$$\eta_{\pi\pi} = \epsilon_L + \epsilon'_{\pi\pi}$$

$$\eta_{3\pi^0} = \epsilon_S + \epsilon'_{3\pi^0}$$

CPT cons. and CPT viol.

$\Delta S = \Delta Q$ violation

$$x_+, x_-$$

Orthogonal
bases

$$\{K_+, \tilde{K}_-\} \quad \{\tilde{K}_+, K_-\}$$

$$\{\tilde{K}_0, K_0\} \text{ and } \{\tilde{K}_0, K_0\}$$

Explicitly for $\Delta t > 0$:

$$\begin{aligned} R_{2,\text{CPT}}^{\text{exp}}(\Delta t) &= \frac{P[\tilde{K}_0(0) \rightarrow K_-(\Delta t)]}{P[\tilde{K}_-(0) \rightarrow K_0(\Delta t)]} \times D_{\text{CPT}} \\ &= |1 - 2\delta + 2x_+^* - 2x_-^*|^2 \left| 1 + (2\delta + \epsilon'_{3\pi^0} - \epsilon'_{\pi\pi}) e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\text{CPT}} \end{aligned}$$

$$\begin{aligned} R_{4,\text{CPT}}^{\text{exp}}(\Delta t) &= \frac{P[\tilde{K}_0(0) \rightarrow K_-(\Delta t)]}{P[\tilde{K}_-(0) \rightarrow K_0(\Delta t)]} \times D_{\text{CPT}} \\ &= |1 + 2\delta + 2x_+ + 2x_-|^2 \left| 1 - (2\delta + \epsilon'_{3\pi^0} - \epsilon'_{\pi\pi}) e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\text{CPT}} \end{aligned}$$

Impact of the approximations

$$\begin{aligned} \frac{R_{2,\text{CPT}}^{\text{exp}}(\Delta t)}{R_{4,\text{CPT}}^{\text{exp}}(\Delta t)} &\simeq (1 - 8\Re\delta - 8\Re x_-) \left| 1 + 2(\eta_{3\pi^0} - \eta_{\pi\pi}) e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \\ &= (1 - 8\Re\delta - 8\Re x_-) \left| 1 + 2(2\delta + \epsilon'_{3\pi^0} - \epsilon'_{\pi\pi}) e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \end{aligned}$$

The double ratio constitutes one of the most robust observables for the proposed CPT test. In the limit $\Delta t \gg \tau_S$ it exhibits a pure and genuine CPT violating effect, even without assuming negligible contaminations from direct CP violation and/or $\Delta S = \Delta Q$ rule violation.

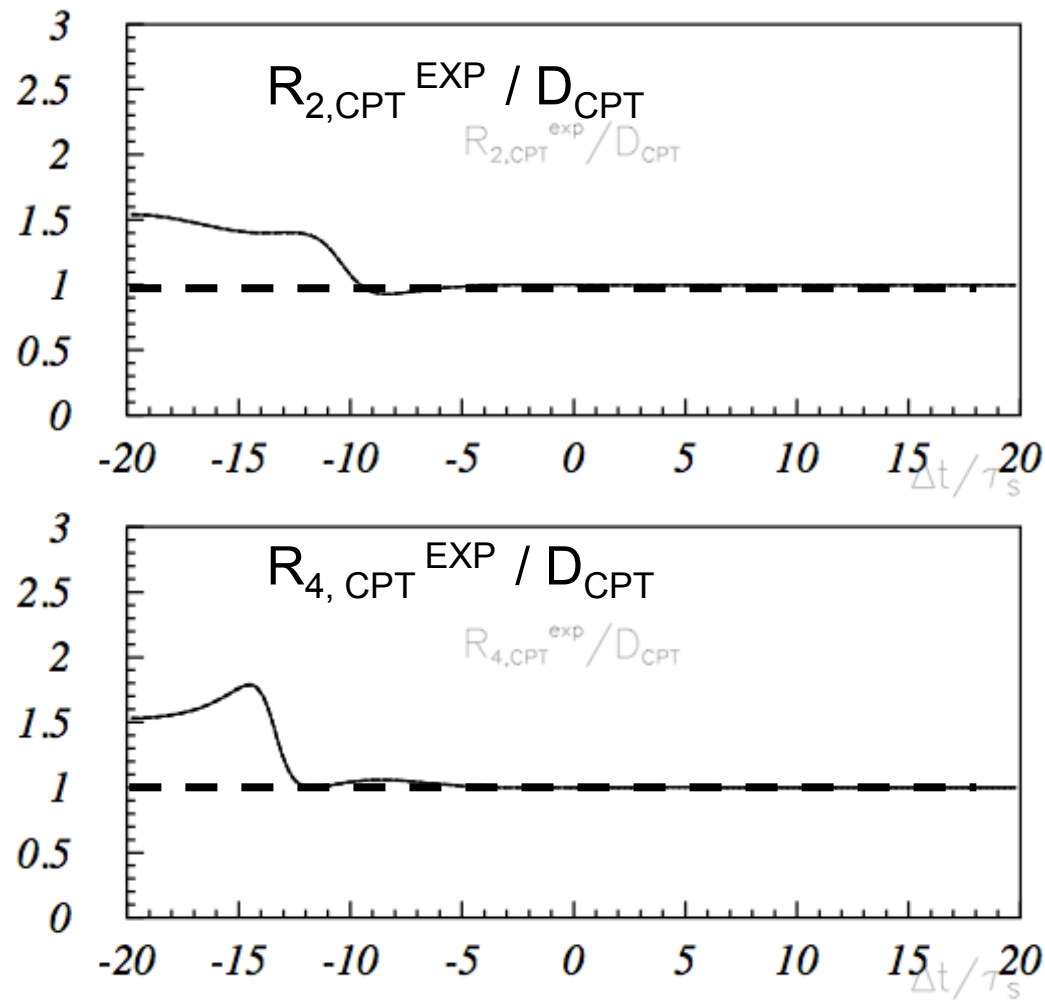
$$\text{DR}_{\text{CPT}} = \frac{R_{2,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)}{R_{4,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)} = 1 - 8\Re\delta - 8\Re x_-$$

There exists a connection with charge semileptonic asymmetries of K_S and K_L

$$\text{DR}_{\text{CPT}} = \frac{R_{2,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)}{R_{4,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)} = \frac{1 + A_L}{1 - A_L} \times \frac{1 - A_S}{1 + A_S} \simeq 1 + 2(A_L - A_S)$$

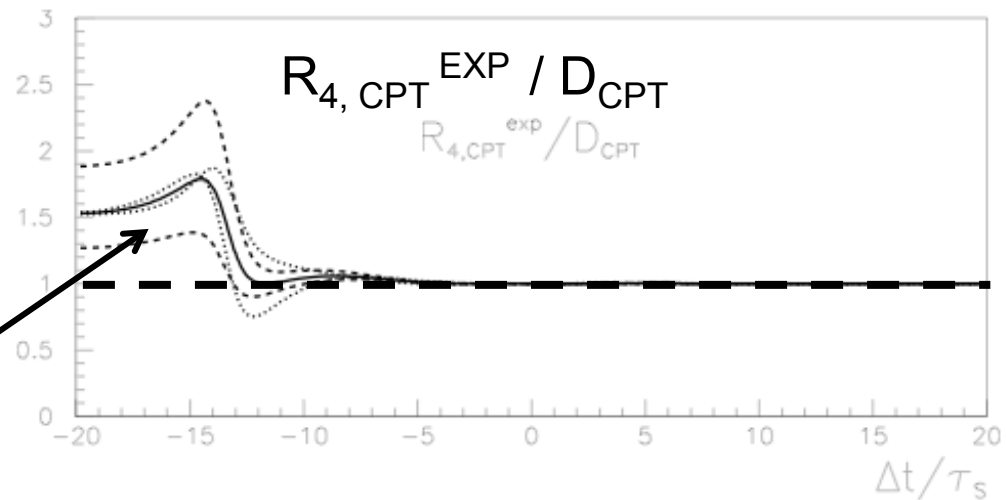
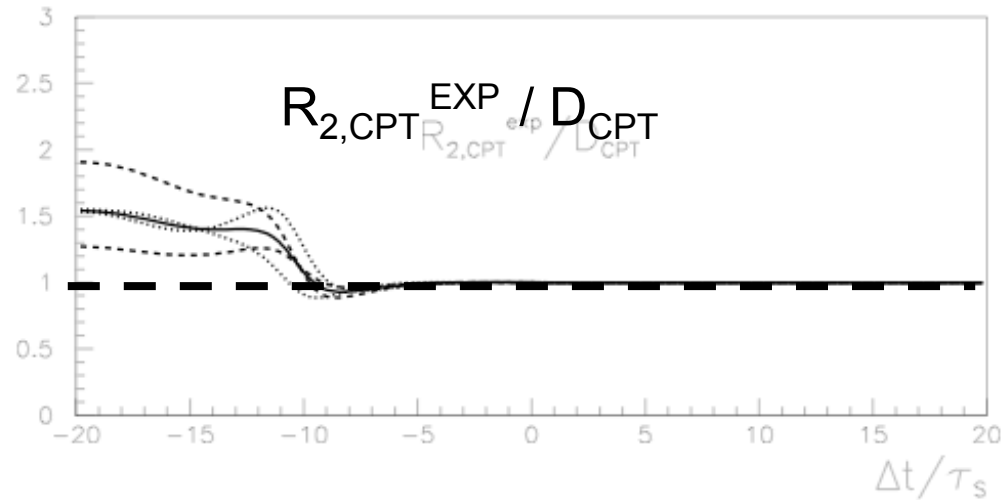
Direct test of CPT in transitions with neutral kaons

for visualization purposes, plots with
 $\text{Re}(\delta)=3.3 \cdot 10^{-4}$ $\text{Im}(\delta)=1.6 \cdot 10^{-5}$



Direct test of CPT in transitions with neutral kaons

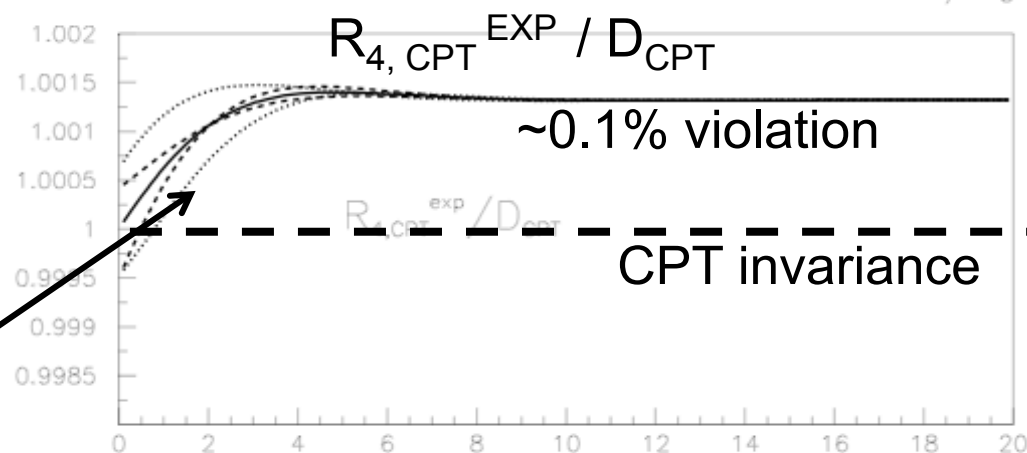
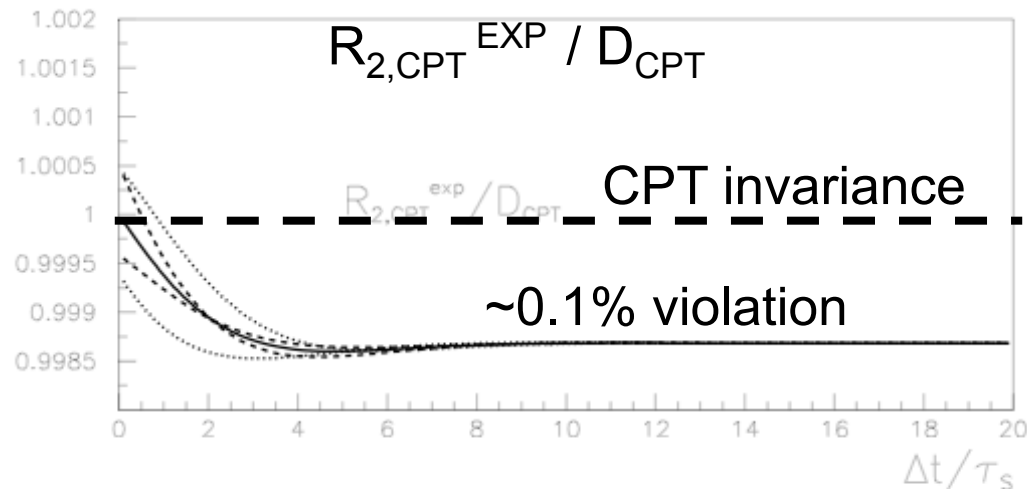
for visualization purposes, plots with
 $\text{Re}(\delta)=3.3 \cdot 10^{-4}$ $\text{Im}(\delta)=1.6 \cdot 10^{-5}$



Modifications due to direct CP violation effects (unrealistically amplified $\sim x100$)

Direct test of CPT in transitions with neutral kaons

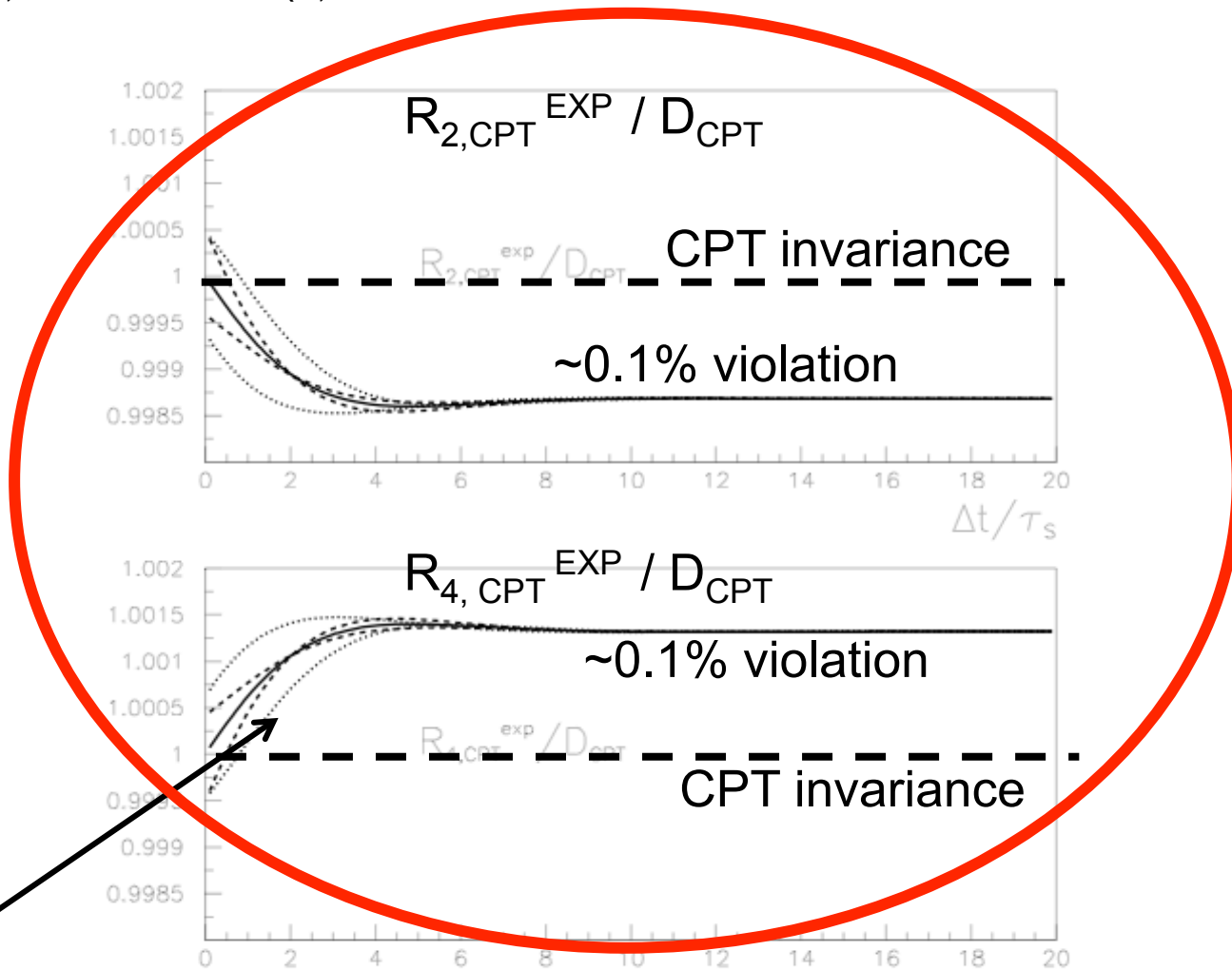
for visualization purposes, plots with
 $\text{Re}(\delta)=3.3 \cdot 10^{-4}$ $\text{Im}(\delta)=1.6 \cdot 10^{-5}$



Modifications due to direct CP violation effects (unrealistically amplified $\sim x100$)

Direct test of CPT in transitions with neutral kaons

for visualization purposes, plots with
 $\text{Re}(\delta)=3.3 \cdot 10^{-4}$ $\text{Im}(\delta)=1.6 \cdot 10^{-5}$



measurable
at KLOE/KLOE-2

Modifications due to direct CP violation effects (unrealistically amplified $\sim x100$)

Direct test of Time Reversal symmetry with neutral kaons

Two observable ratios of double decay intensities

$$R_{2,\mathcal{T}}^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^-, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^+; \Delta t)}$$

$$R_{4,\mathcal{T}}^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^+, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^-; \Delta t)}$$

Direct test of Time Reversal symmetry with neutral kaons

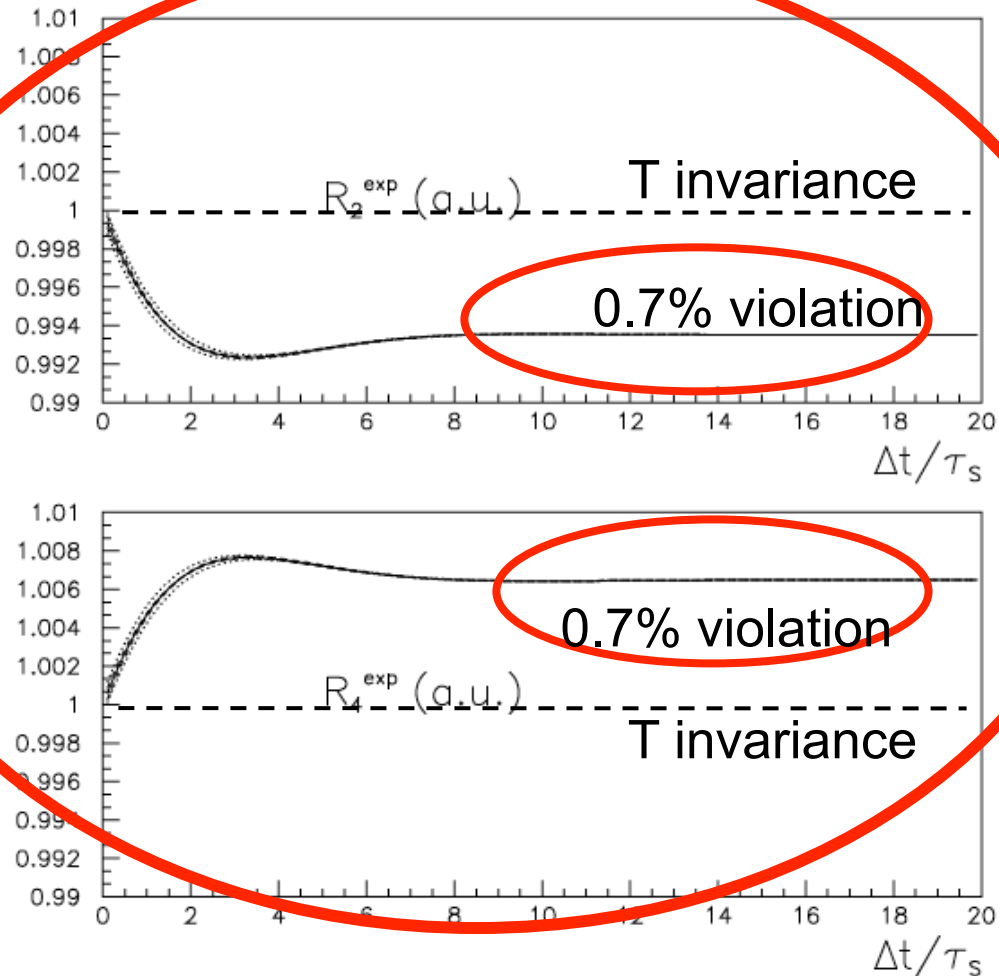
Explicitly in standard
Wigner Weisskopf
approach
for $\Delta t > 0$:

$$\begin{aligned} R_{2,\mathcal{T}}^{\text{exp}}(\Delta t) &= \frac{P[\text{K}^0(0) \rightarrow \text{K}_-(\Delta t)]}{P[\text{K}_-(0) \rightarrow \text{K}^0(\Delta t)]} \times D_{\mathcal{T},2} \\ &= (1 - 4\Re\epsilon) \left| 1 + 2\epsilon e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\mathcal{CPT}} \end{aligned}$$

$$\begin{aligned} R_{4,\mathcal{T}}^{\text{exp}}(\Delta t) &= \frac{P[\bar{\text{K}}^0(0) \rightarrow \text{K}_-(\Delta t)]}{P[\text{K}_-(0) \rightarrow \bar{\text{K}}^0(\Delta t)]} \times D_{\mathcal{T},4} \\ &= (1 + 4\Re\epsilon) \left| 1 - 2\epsilon e^{-i(\lambda_S - \lambda_L)\Delta t} \right|^2 \times D_{\mathcal{CPT}} \end{aligned}$$

Direct test of Time Reversal symmetry with neutral kaons

plots with CPV
 $\text{Re}\epsilon$ and $\text{Im}\epsilon$
 values



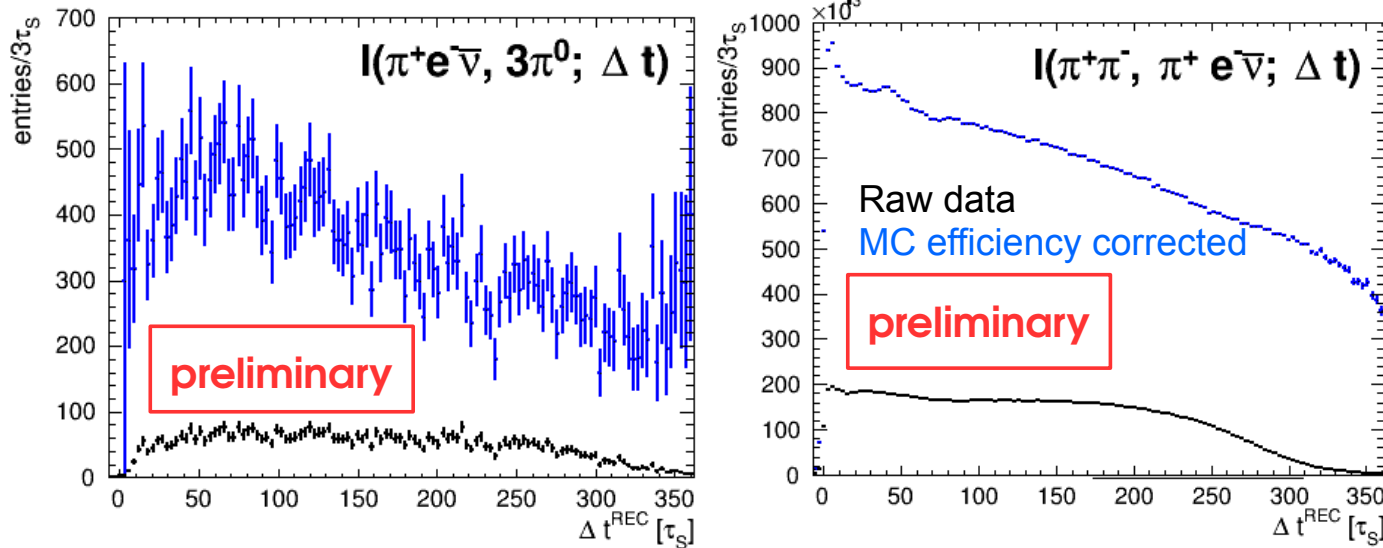
measurable
 at KLOE-2

$$R_2(\Delta t \gg \tau_S) = 1 - 4\text{Re}(\epsilon) \sim 0.993$$

$$R_4(\Delta t \gg \tau_S) = 1 + 4\text{Re}(\epsilon) \sim 1.007$$

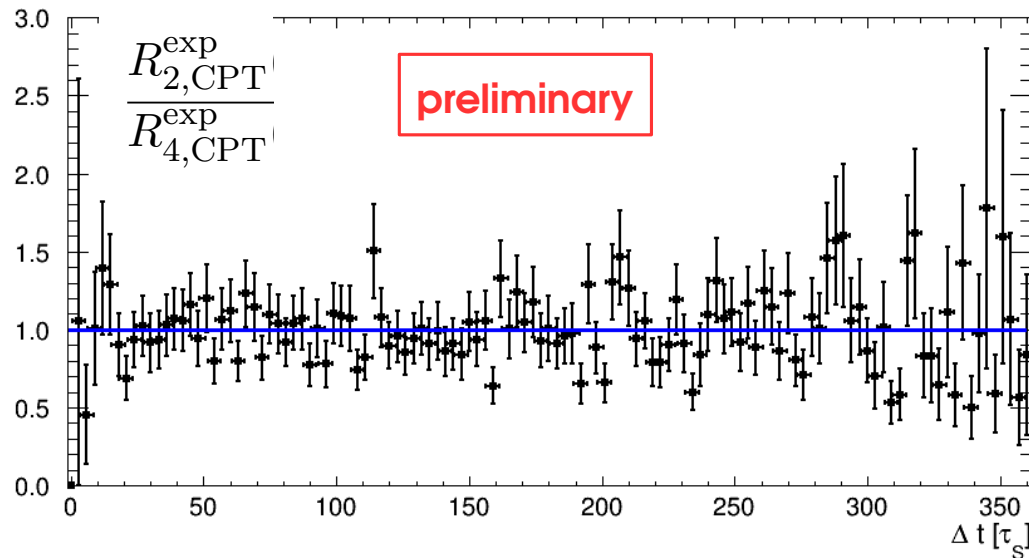
Direct test of CPT in transitions with neutral kaons at KLOE

KLOE data sample: $L=1.7 \text{ fb}^{-1}$



$$R_{2,\text{CPT}}^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^-, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^-; \Delta t)}$$

$$R_{4,\text{CPT}}^{\text{exp}}(\Delta t) \equiv \frac{I(\ell^+, 3\pi^0; \Delta t)}{I(\pi\pi, \ell^+; \Delta t)}$$



CPT test with the double ratio DR_{CPT} :

$$\frac{R_{2,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)}{R_{4,\text{CPT}}^{\text{exp}}(\Delta t \gg \tau_S)} = 1 - 8\Re\delta - 8\Re x_-$$

← $R_2/R_4=1$

- $K_L \rightarrow 3\pi^0$ vtx reconstr. with GPS-like technique
- Analysis in progress: efficiency correction from data control samples
- KLOE-2 can reach a precision $O(10^{-3})$ on R_2/R_4

List of other KLOE CP/CPT/QM tests with neutral kaons

Mode	Test	Param.	KLOE measurement
$K_L \rightarrow \pi^+ \pi^-$	CP	BR	$(1.963 \pm 0.012 \pm 0.017) \times 10^{-3}$
$K_S \rightarrow 3\pi^0$	CP	BR	$< 2.6 \times 10^{-8}$
$K_S \rightarrow \pi e \nu$	CP	A_S	$(1.5 \pm 10) \times 10^{-3}$
$K_S \rightarrow \pi e \nu$	CPT	$\text{Re}(x)$	$(-0.8 \pm 2.5) \times 10^{-3}$
$K_S \rightarrow \pi e \nu$	CPT	$\text{Re}(y)$	$(0.4 \pm 2.5) \times 10^{-3}$
All $K_{S,L}$ BRs, η 's etc... (unitarity)	CP CPT	$\text{Re}(\epsilon)$ $\text{Im}(\delta)$	$(159.6 \pm 1.3) \times 10^{-5}$ $(0.4 \pm 2.1) \times 10^{-5}$
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	QM	ζ_{00}	$(0.1 \pm 1.0) \times 10^{-6}$
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	QM	ζ_{SL}	$(0.3 \pm 1.9) \times 10^{-2}$
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	CPT & QM	α	$(-10 \pm 37) \times 10^{-17}$ GeV
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	CPT & QM	β	$(1.8 \pm 3.6) \times 10^{-19}$ GeV
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	CPT & QM	γ	$(0.4 \pm 4.6) \times 10^{-21}$ GeV compl. pos. hyp. $(0.7 \pm 1.2) \times 10^{-21}$ GeV
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	CPT & QM	$\text{Re}(\omega)$	$(-1.6 \pm 2.6) \times 10^{-4}$
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	CPT & QM	$\text{Im}(\omega)$	$(-1.7 \pm 3.4) \times 10^{-4}$
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	CPT & Lorentz	Δa_0	$(-6.2 \pm 8.8) \times 10^{-18}$ GeV
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	CPT & Lorentz	Δa_Z	$(-0.7 \pm 1.0) \times 10^{-18}$ GeV
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	CPT & Lorentz	Δa_X	$(3.3 \pm 2.2) \times 10^{-18}$ GeV
$K_S K_L \rightarrow \pi^+ \pi^-, \pi^+ \pi^-$	CPT & Lorentz	Δa_Y	$(-0.7 \pm 2.0) \times 10^{-18}$ GeV

Updated results for end 2017

Expected updated results for 2018

Conclusions

- The entangled neutral kaon system at a ϕ -factory is an excellent laboratory for the study of discrete symmetries and fundamental principles of QM.
- The analysis of the full KLOE data set is being completed:
 - a new measurement of the KS semileptonic charge asymmetry
 - the analysis for first test of T and CPT in neutral kaon transitions processes is ongoing.
- It is possible to directly test CPT in transition processes for the first time between neutral kaon states. The proposed CPT test is model independent and fully robust.
- **VERY CLEAN CPT TEST**. Possible spurious effects are well under control, e.g. direct CP violation, $\Delta S = \Delta Q$ rule violation, decoherence effects.
- Several CPTV and/or decoherence parameters have been measured at KLOE, in some cases with a precision reaching the interesting [Planck's scale region](#);
- All results are consistent with no CPT symmetry violation and no decoherence;
- The KLOE-2 experiment at the upgraded DAFNE is currently taking data with the plan to collect $L > 5 \text{ fb}^{-1}$ by end of March 2018;
- All these tests are going to be improved at KLOE-2; a statistical sensitivity of $O(10^{-3})$ could be reached on the newly proposed observables.