

From Quantum Cellular Automata to Doubly Special Relativity

Alessandro Bisio

Workshop Quantum Foundations

New frontiers in testing quantum mechanics from underground to the space

November 30th 2017

Laboratori Nazionali di Frascati

in collaboration with:

Giacomo Mauro D'Ariano

Paolo Perinotti

Alessandro Tosini

Nicola Mosco

Marco Erba

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JOHN TEMPLETON FOUNDATION

SUPPORTING SCIENCE ~ INVESTING IN THE BIG QUESTIONS

Quantum Theory

Von Neumann, 1932

Each physical system is associated with a Hilbert space

Unit vectors are associated with states of the system

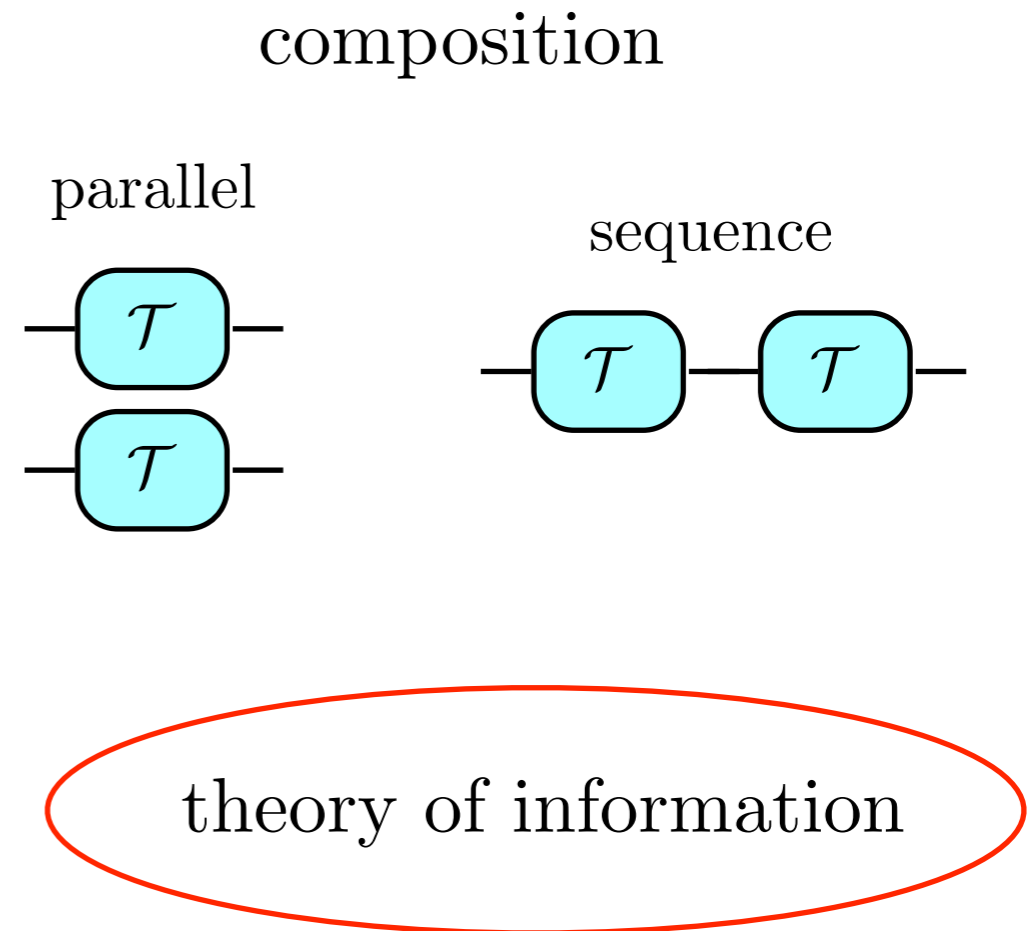
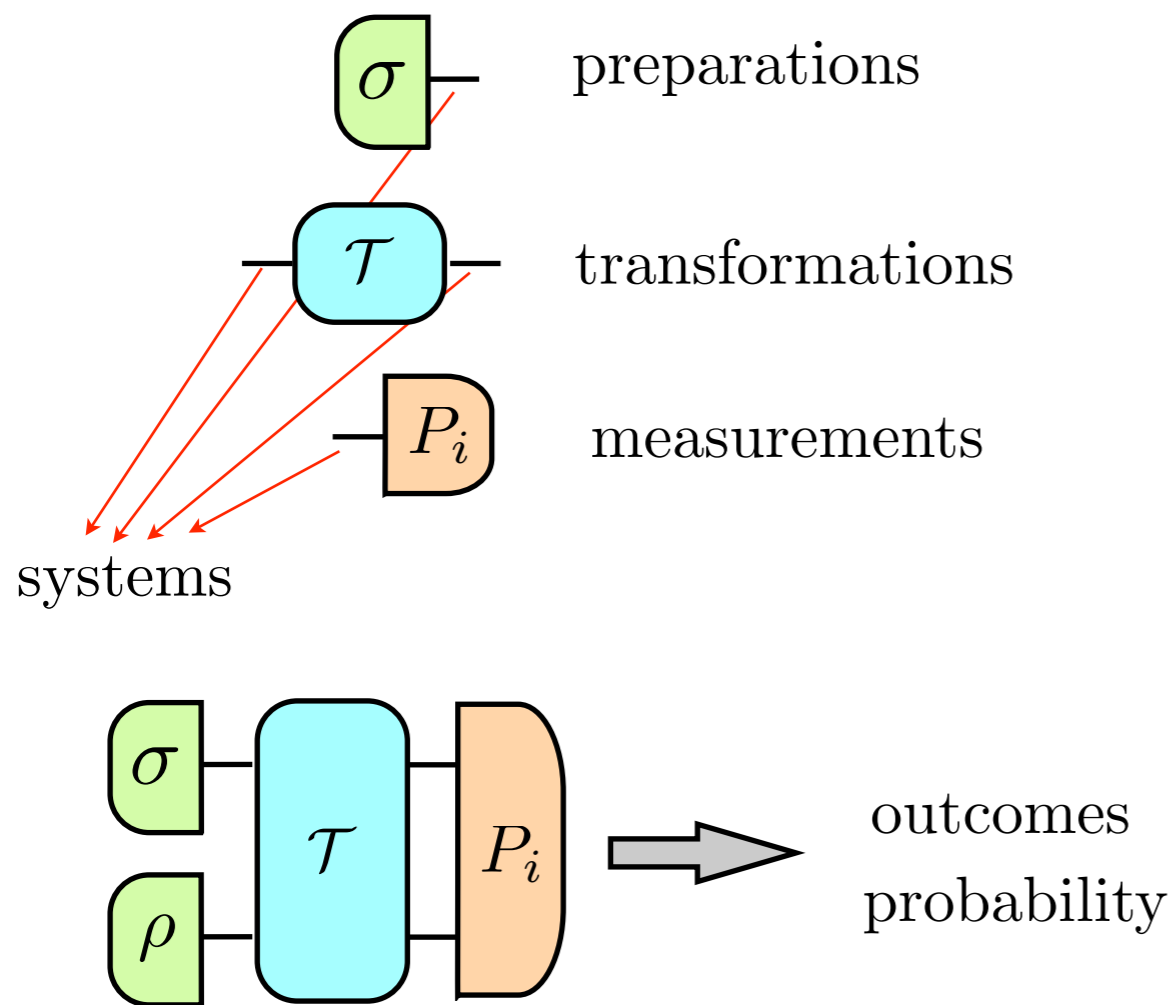
Physical observables are represented by self adjoint operators

The Hilbert space of a composite system is the tensor product of the state spaces associated with the component systems

The probabilities of the outcomes are given by the Born rule

Reconstruction of Quantum Theory

Operational Probabilistic Theory



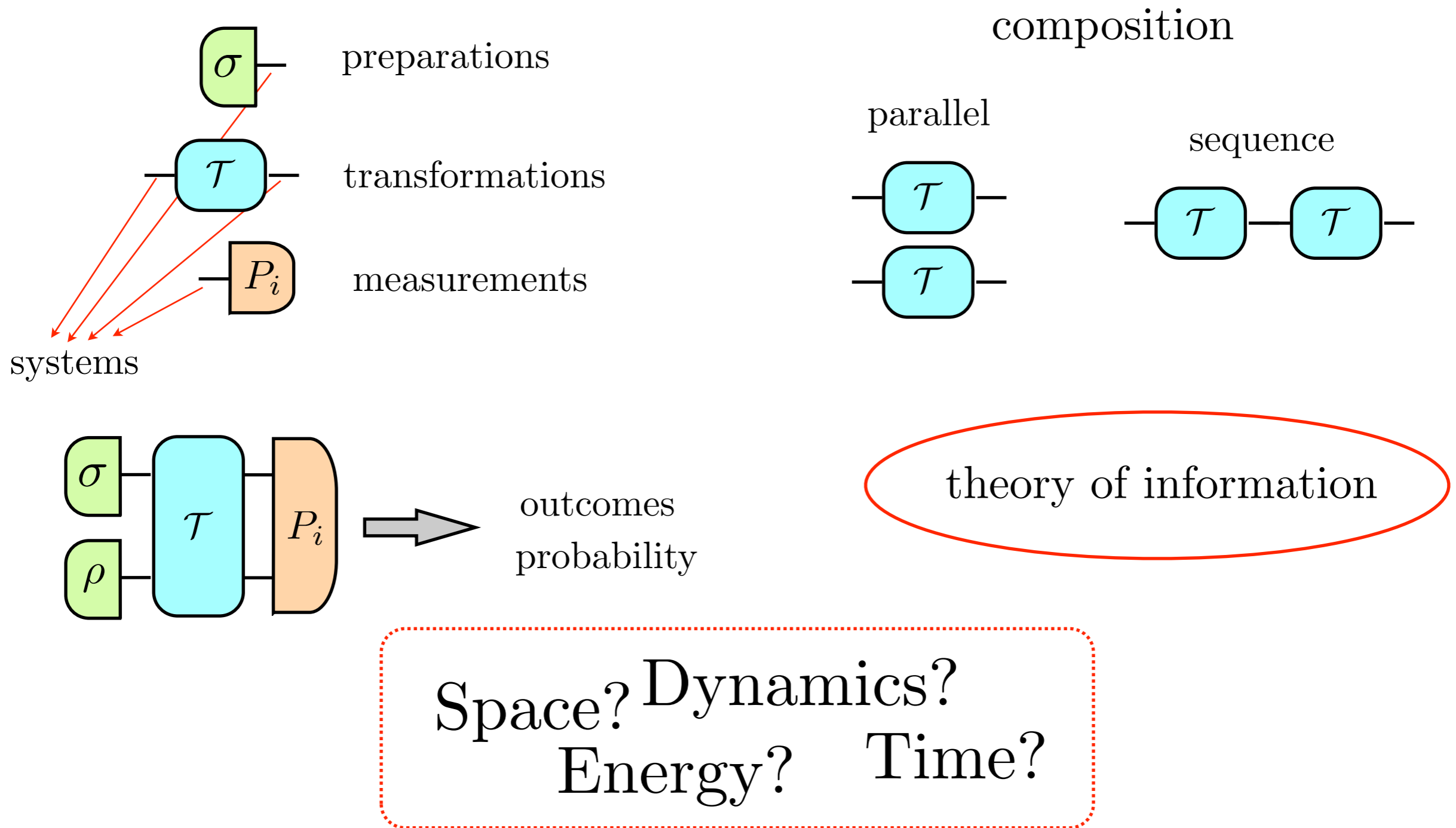
G. Ludwig, *Foundations of Quantum Mechanics* (Springer, New York, 1985).

L. Hardy, e-print arXiv:quant-ph/0101012.

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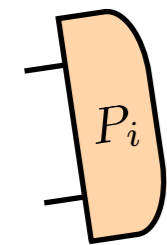
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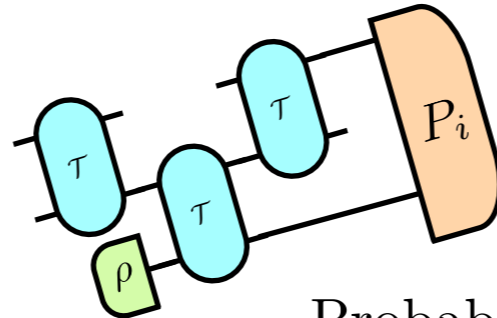
Quantum Theory



$$p(i|j) = \text{Tr}[\rho_i P_j]$$

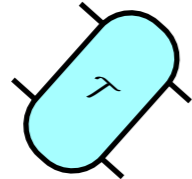


Quantum computation



Probabilistic structure

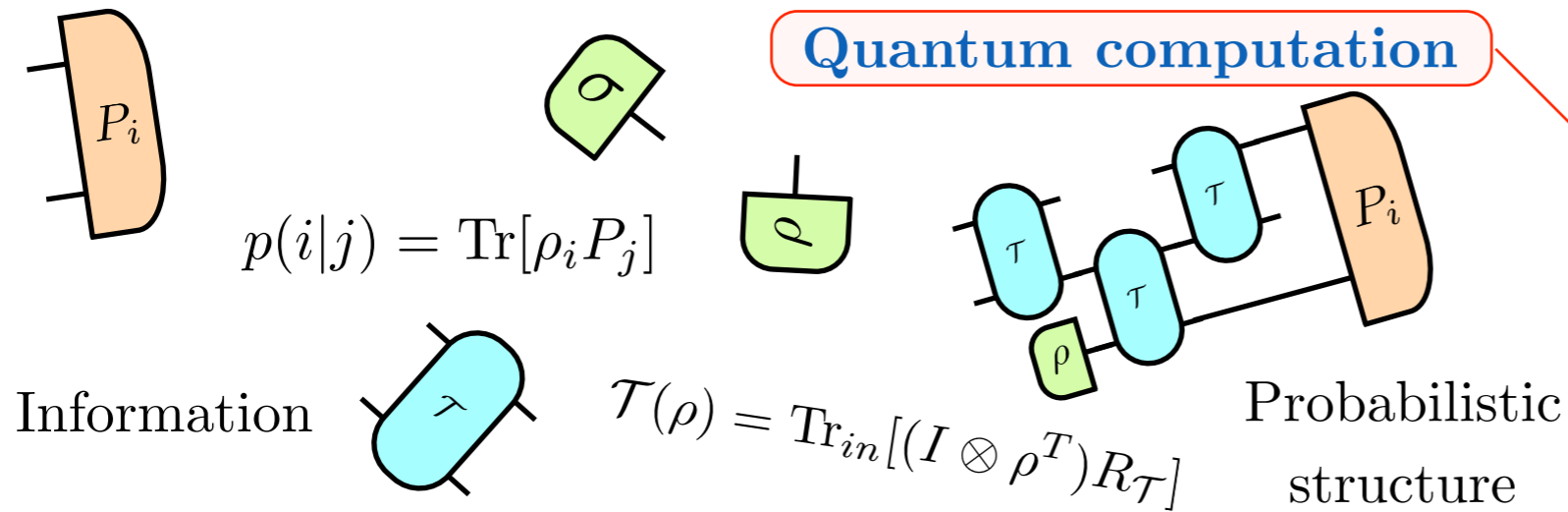
Information



$$\mathcal{T}(\rho) = \text{Tr}_{in}[(I \otimes \rho^T) R_{\mathcal{T}}]$$

Reconstruction of Quantum Field Theory

Quantum Theory



Simulating Physics

with Computers

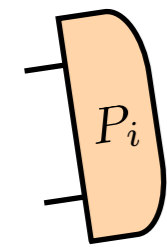
R. P. Feynman, Int. J. Theo. Phys. 21, 467 (1982)

Can a Quantum Computer **exactly** simulate physical systems?

Replace physical laws with quantum algorithms

Reconstruction of Quantum Field Theory

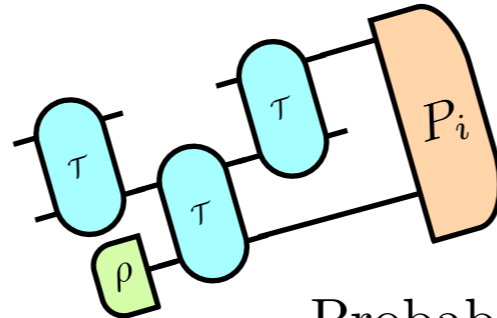
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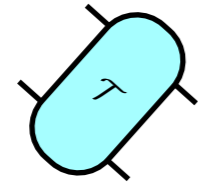


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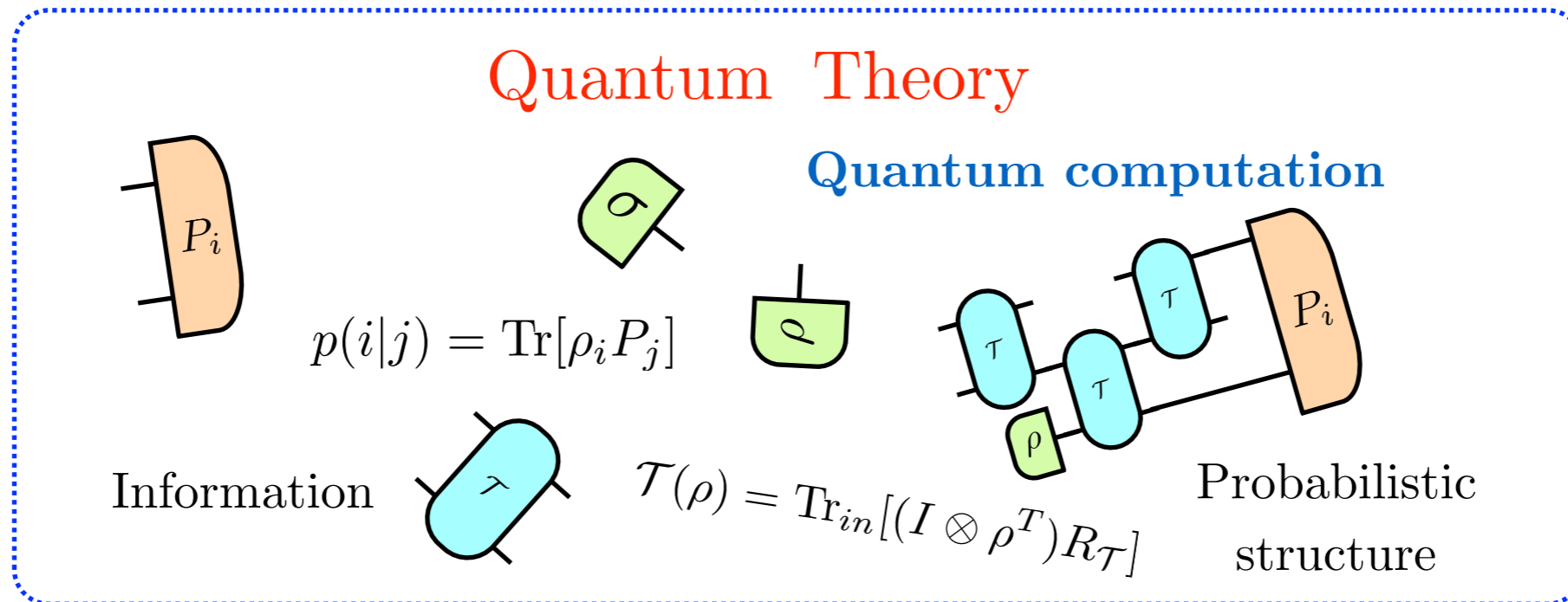
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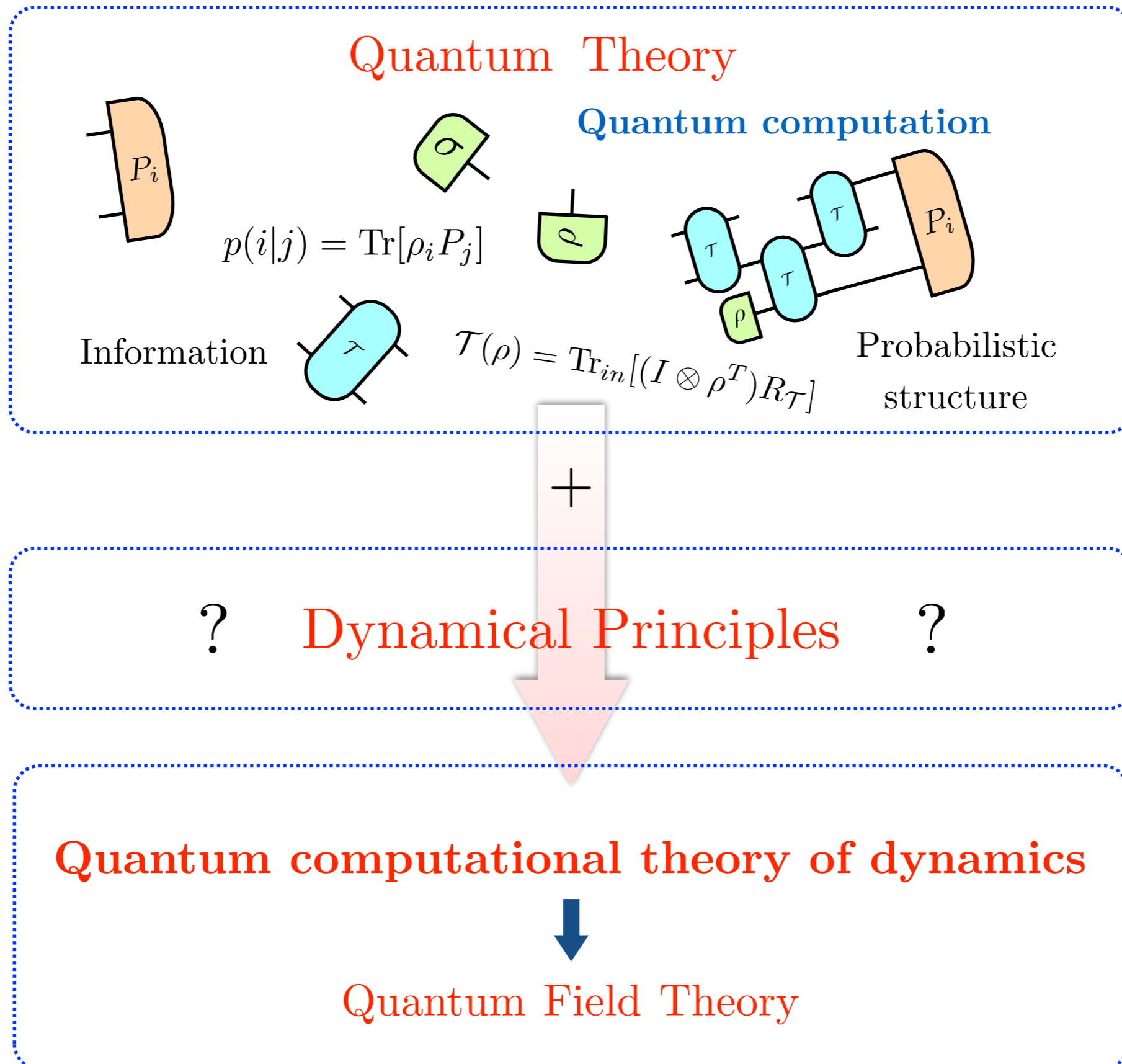
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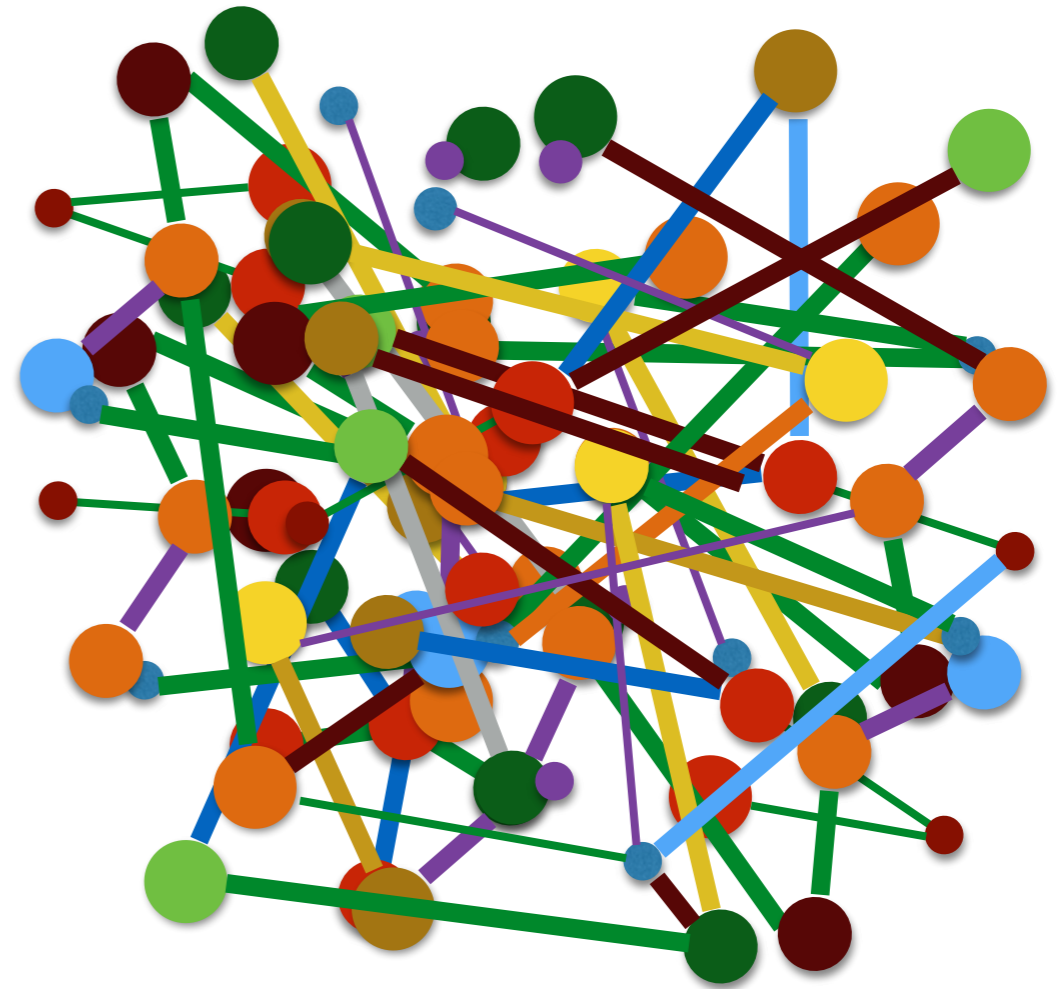


Reconstruction of Quantum Field Theory



What kind of computer?

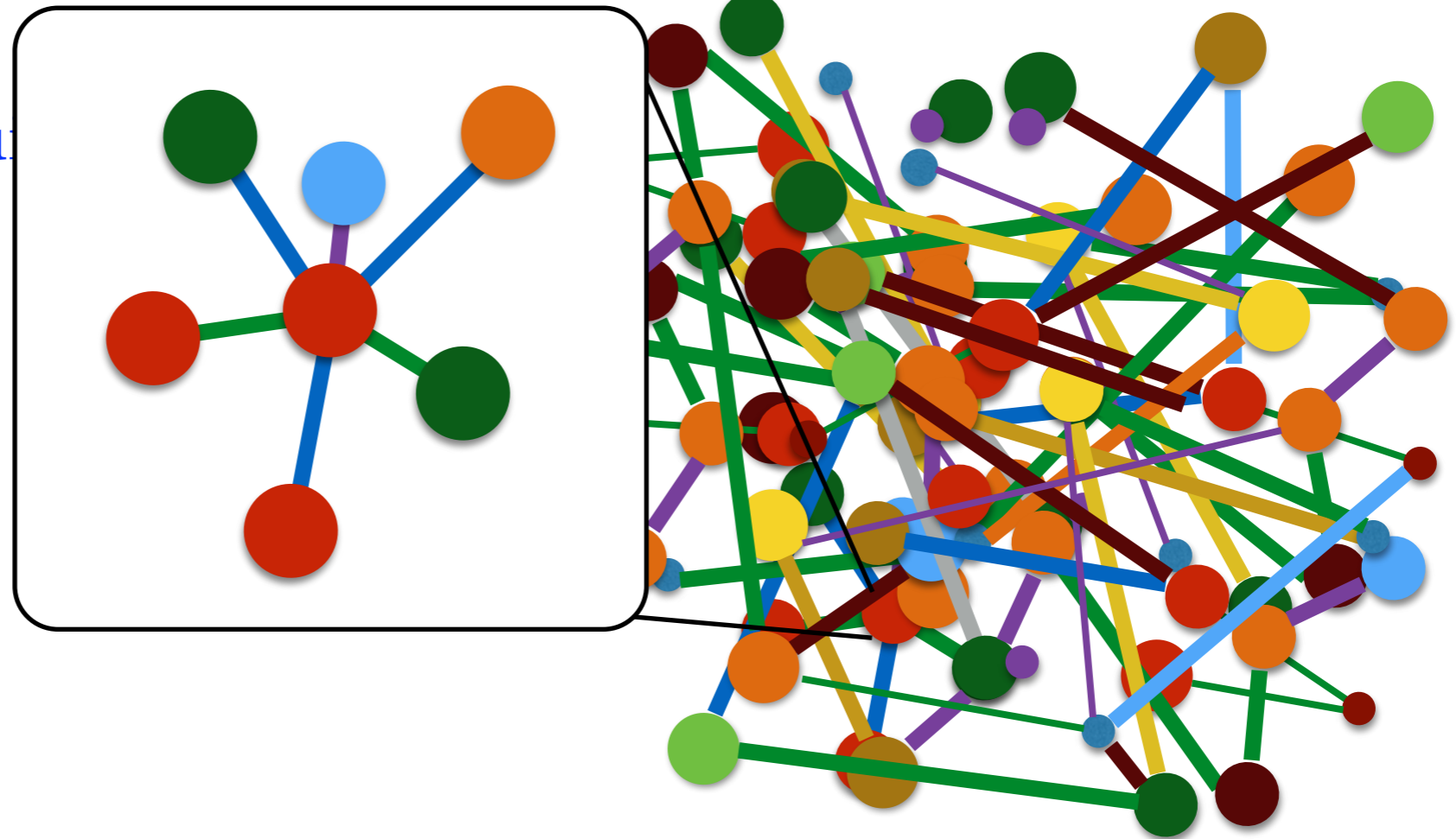
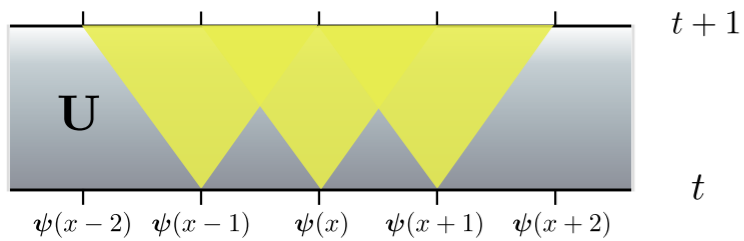
Quantum Circuit



Rules of the game \iff axioms

What kind of computer?

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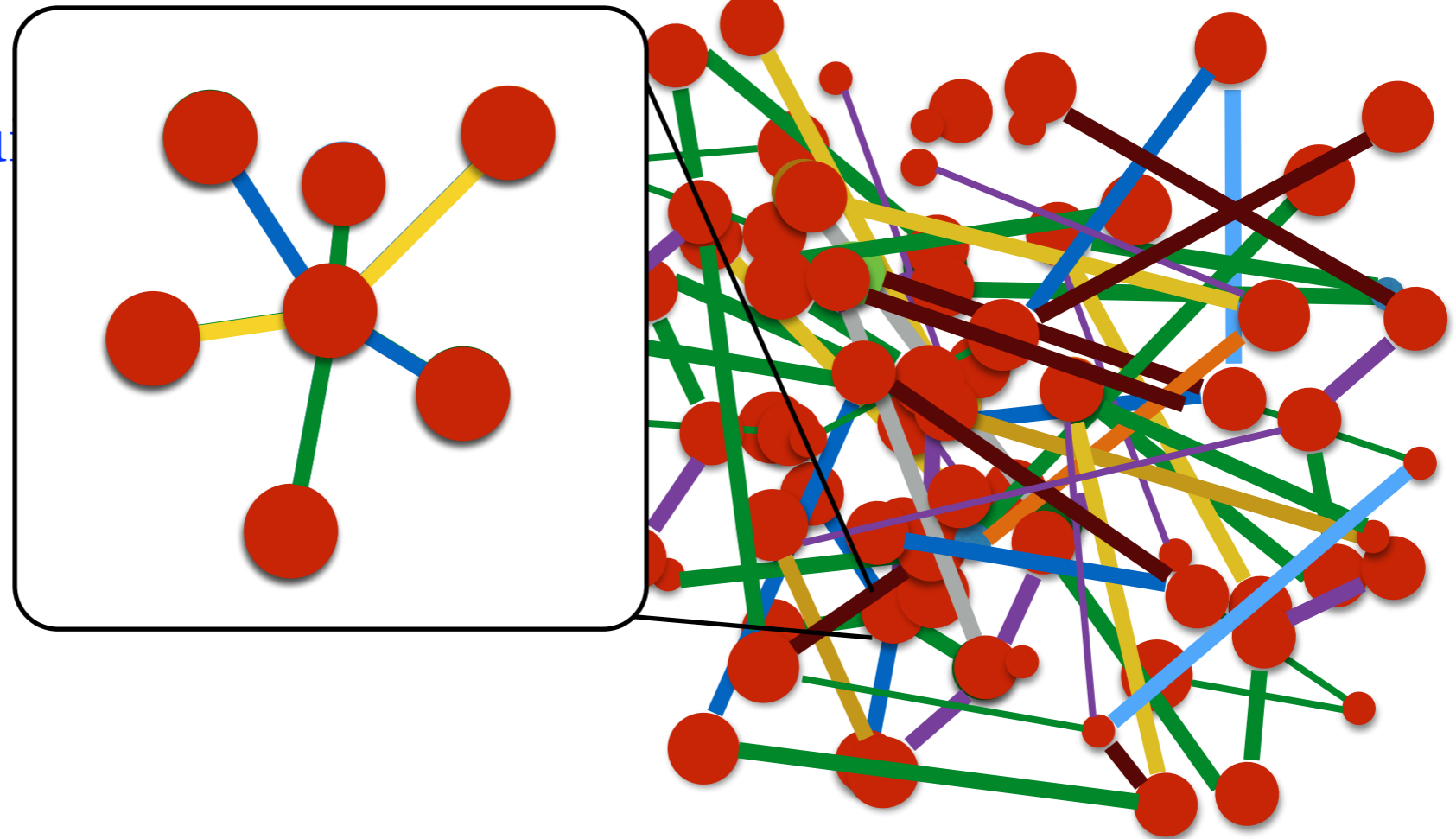
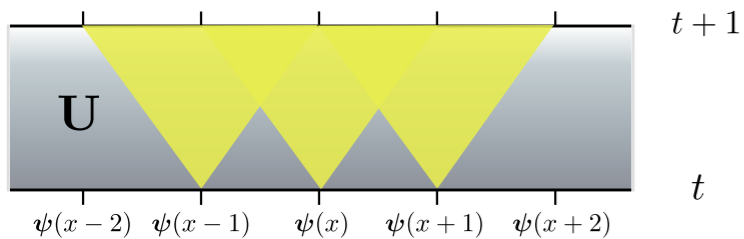
“[...] everything that happens in a finite volume of space and time would have to be exactly analyzable with a finite numbers of logical operations” R. Feynman

Each system interacts with a finite number of neighbors: **locality**

Reversible Quantum Computation: **unitary evolution**

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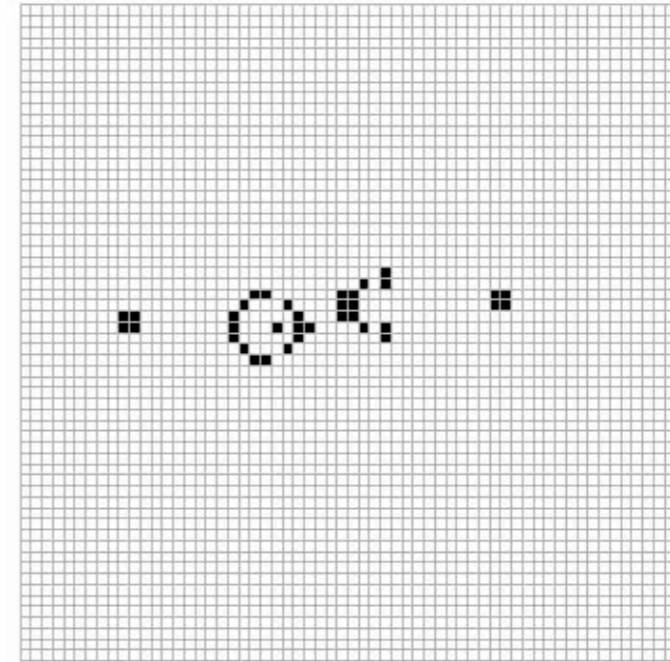
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Quantum Cellular
Automaton

B. Schumacher, R.F. Werner
e-print arXiv:0405174.

like Conway's game of life...



...but **quantum!**

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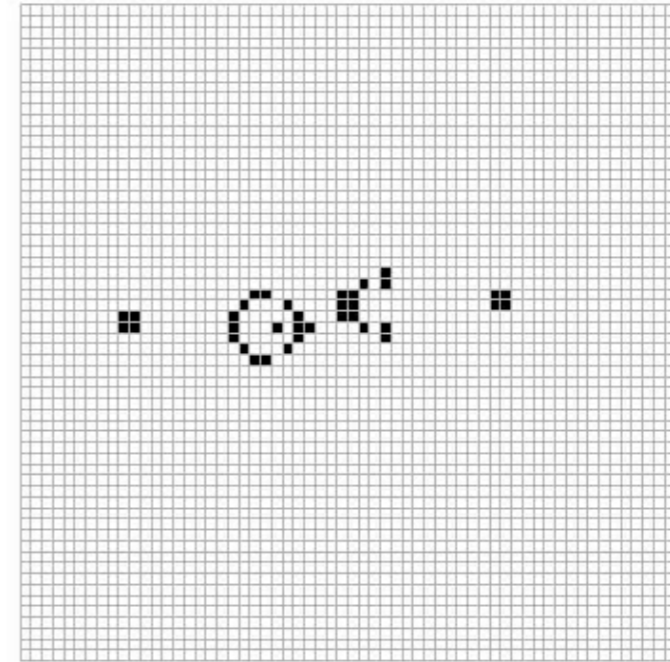
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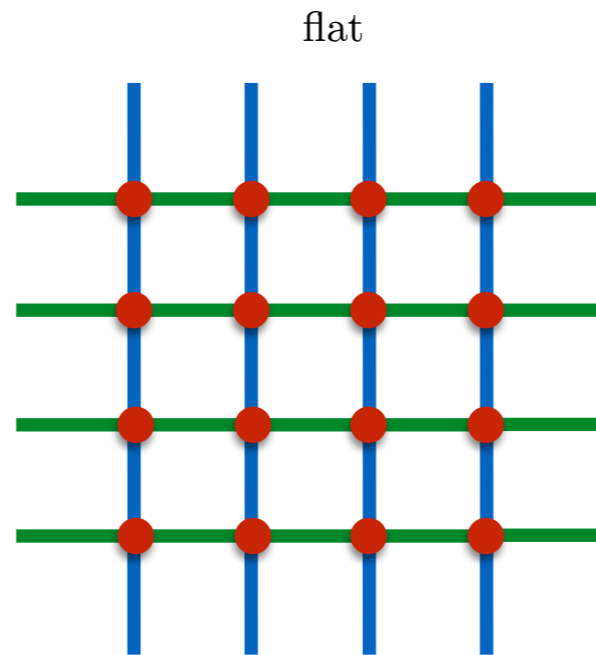
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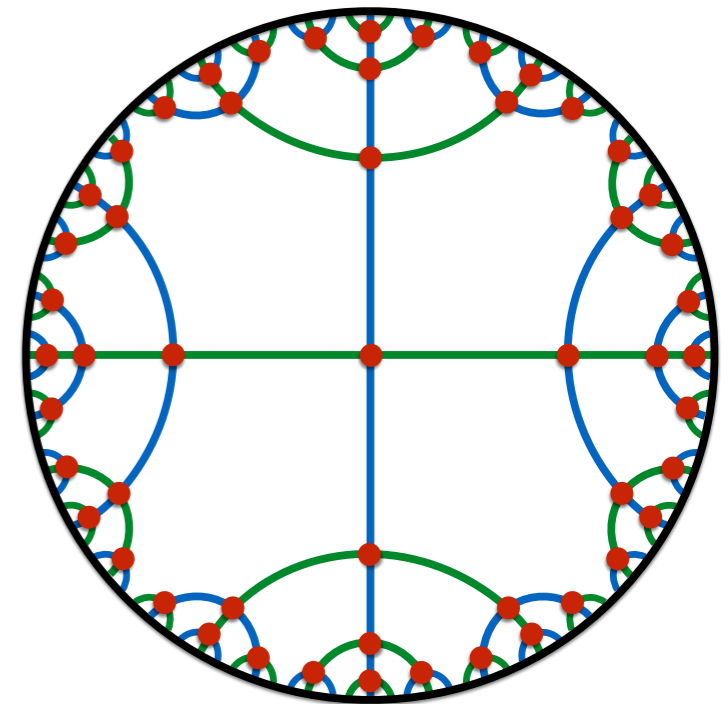


Quantum Cellular
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on a
Cayley Graph

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curved (e.g. hyperbolic)



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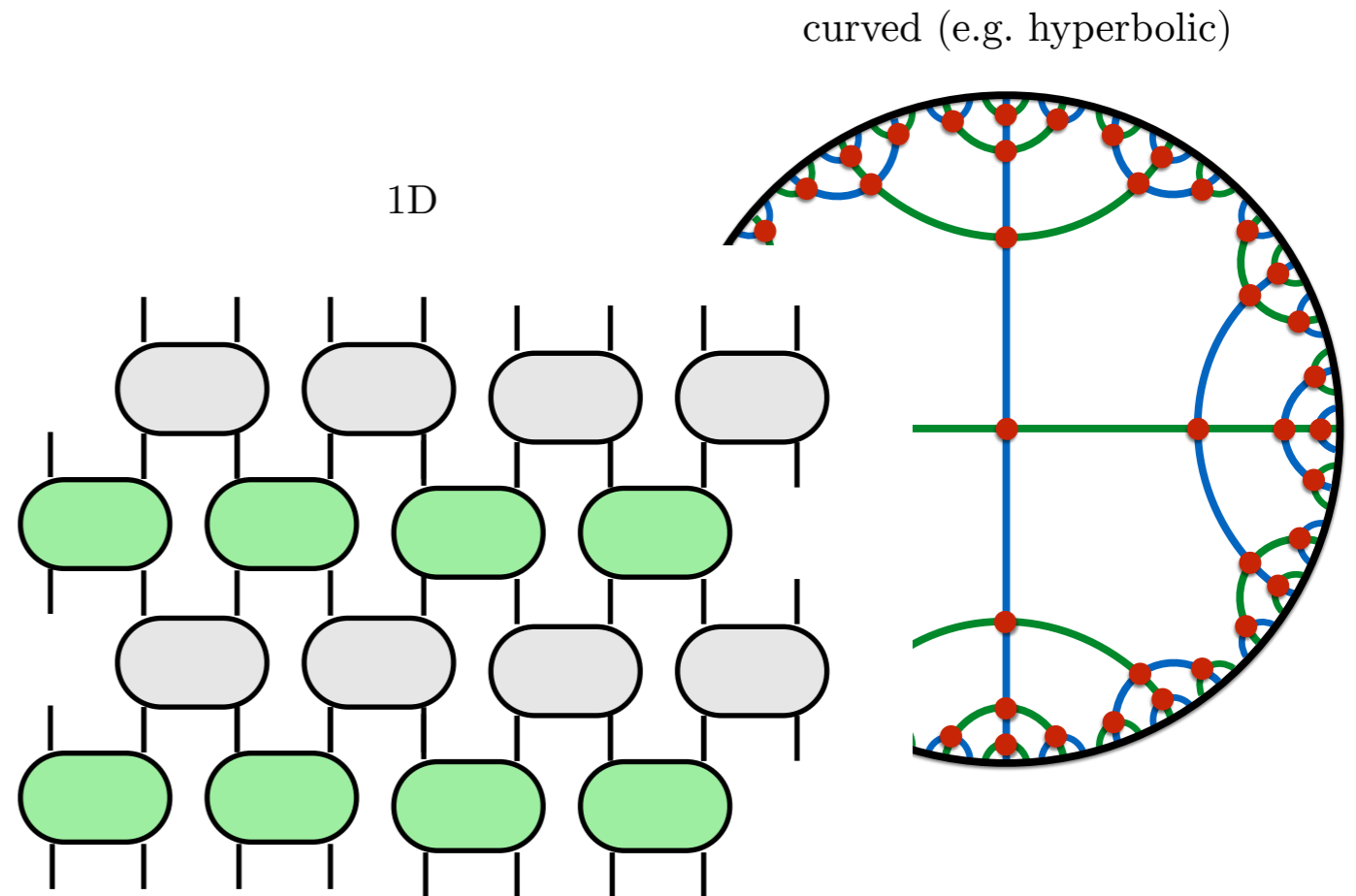


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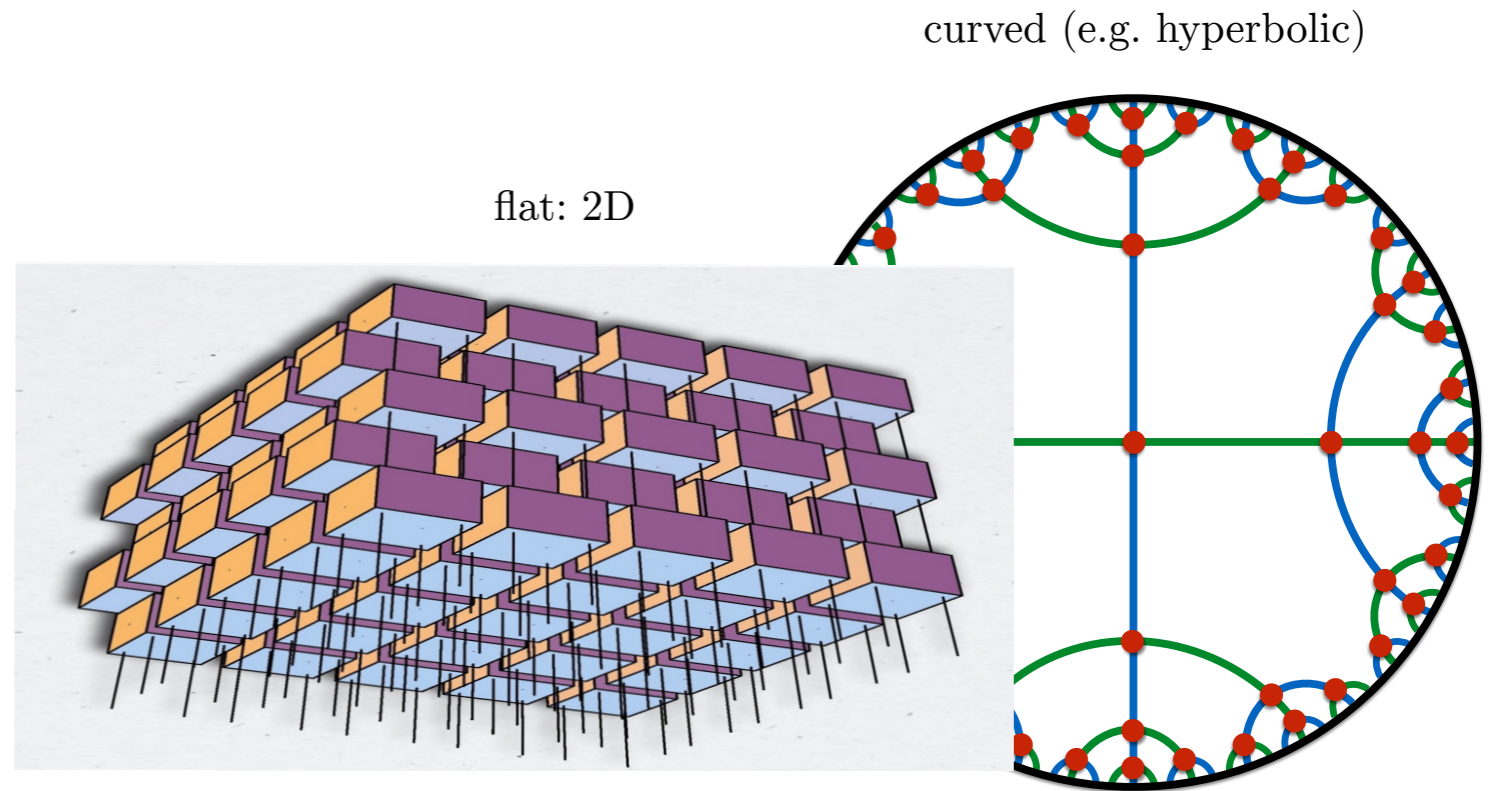


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Towards free Quantum Field Theory

Cayley graph **quasi-isometrically** embeddable in flat space



The group G must be **virtually abelian** \implies We restrict to the abelian case \mathbb{Z}^3

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Linear evolution $U \psi_i^s U^\dagger = \mathbf{U}_{i,j}^{s,r} \psi_j^r \iff$ **Quantum Walk**

field operators \leftarrow ψ_i^s
site label \leftarrow i
internal degree of freedom (e.g. spinorial index) \leftarrow s
 $\mathbf{U}_{i,j}^{s,r}$ \leftarrow unitary matrix

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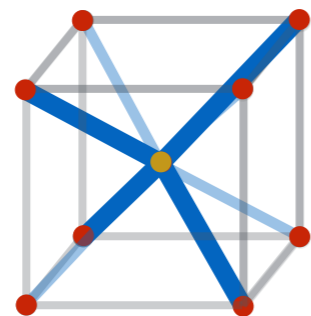
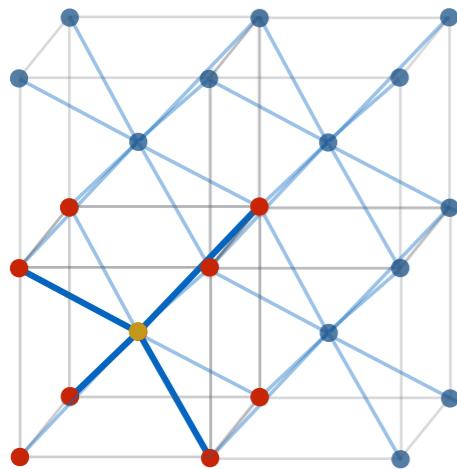
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Weyl Quantum Walk

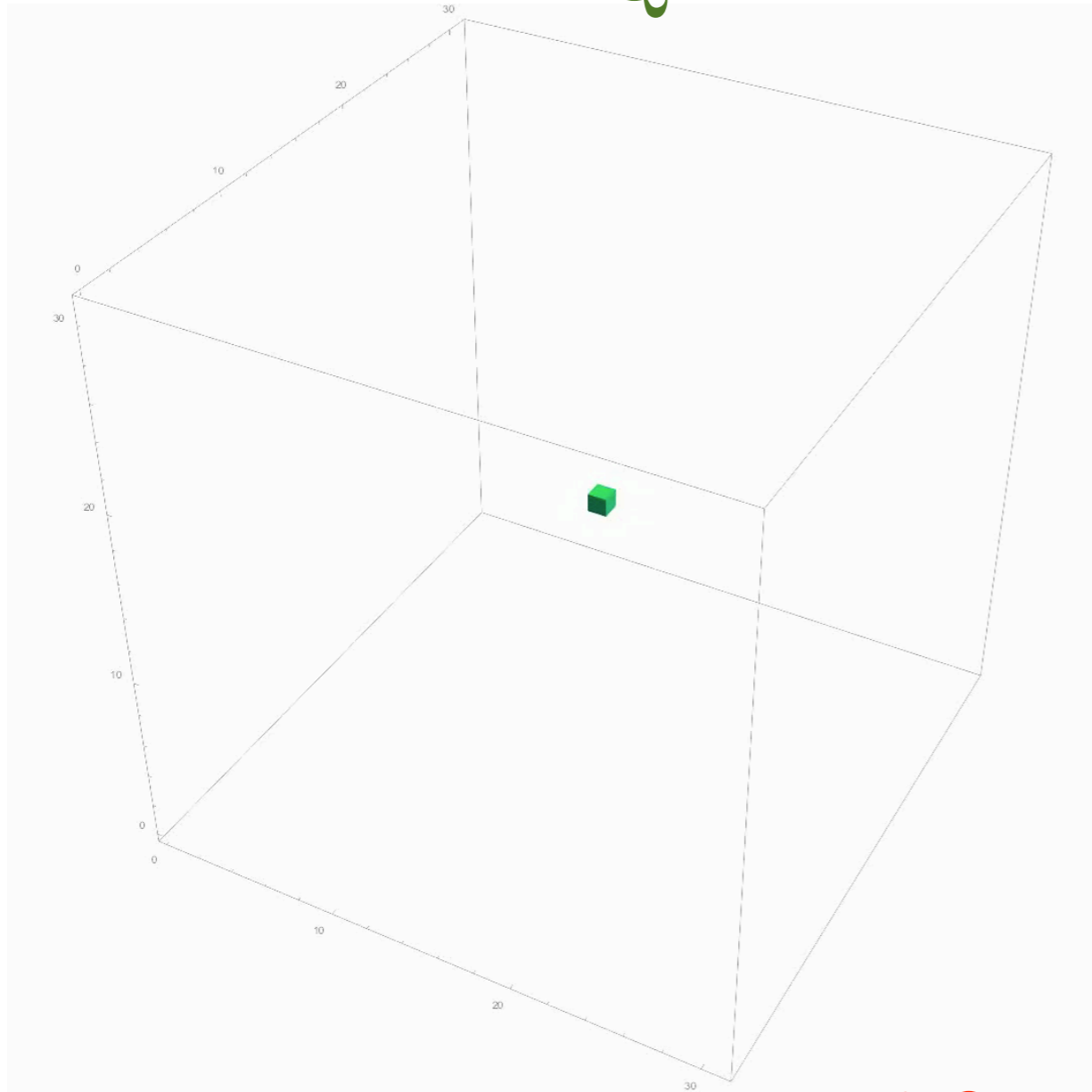


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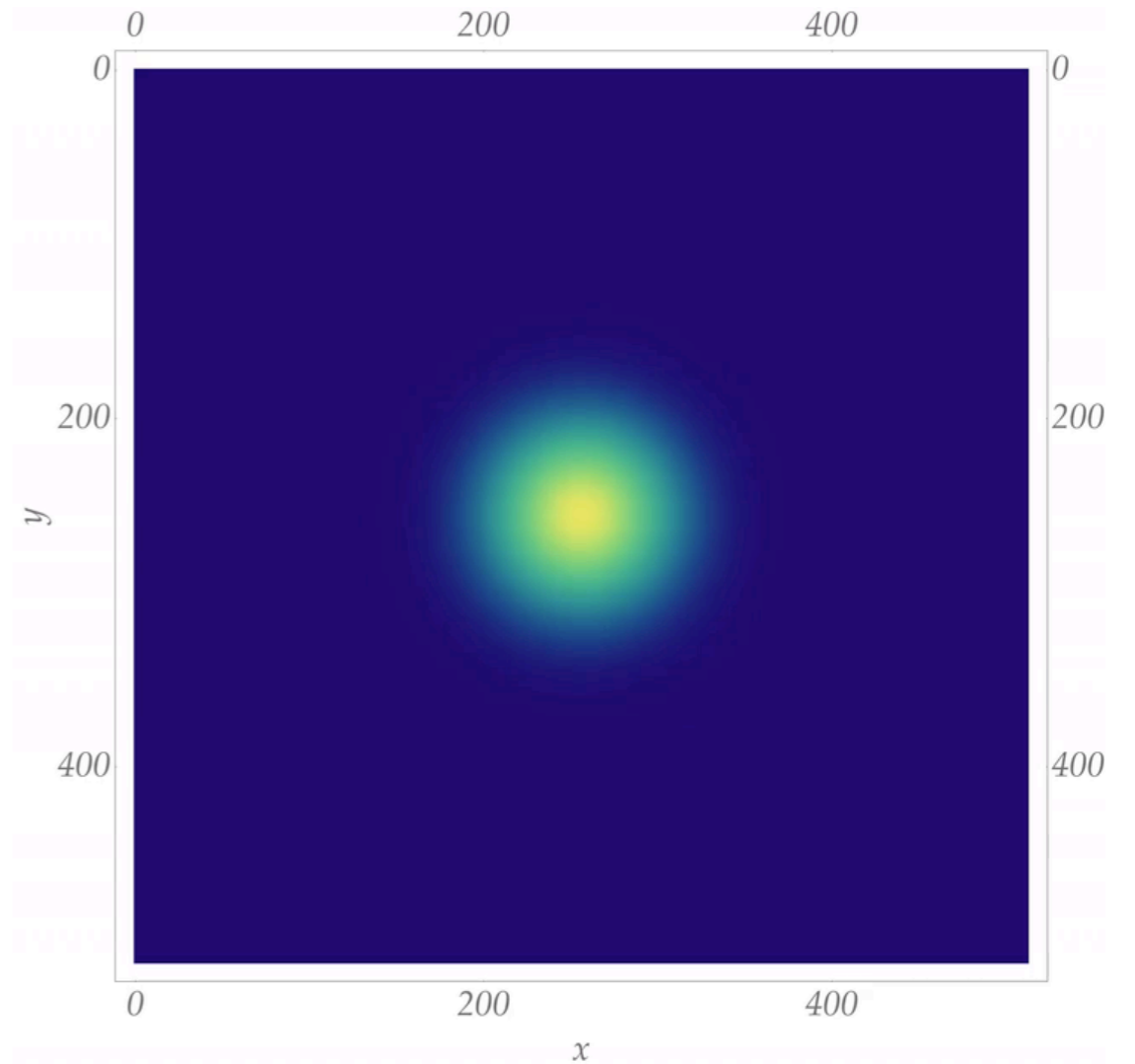
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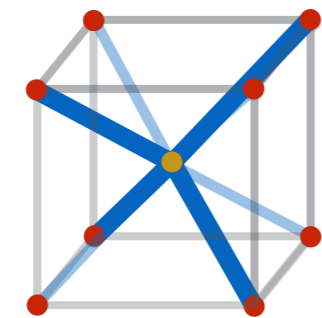
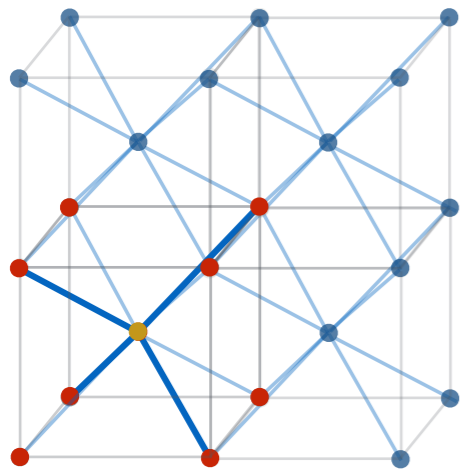
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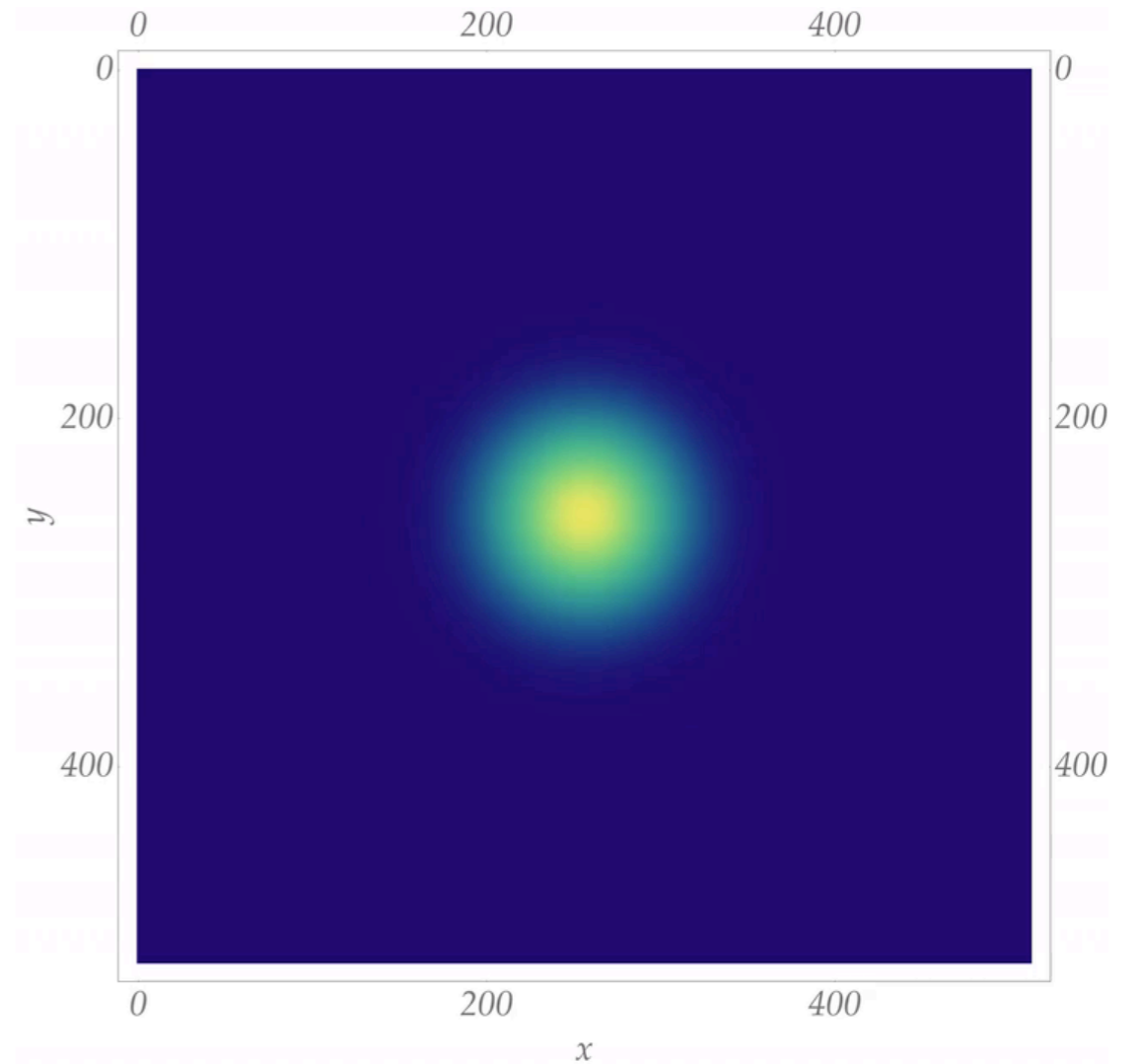
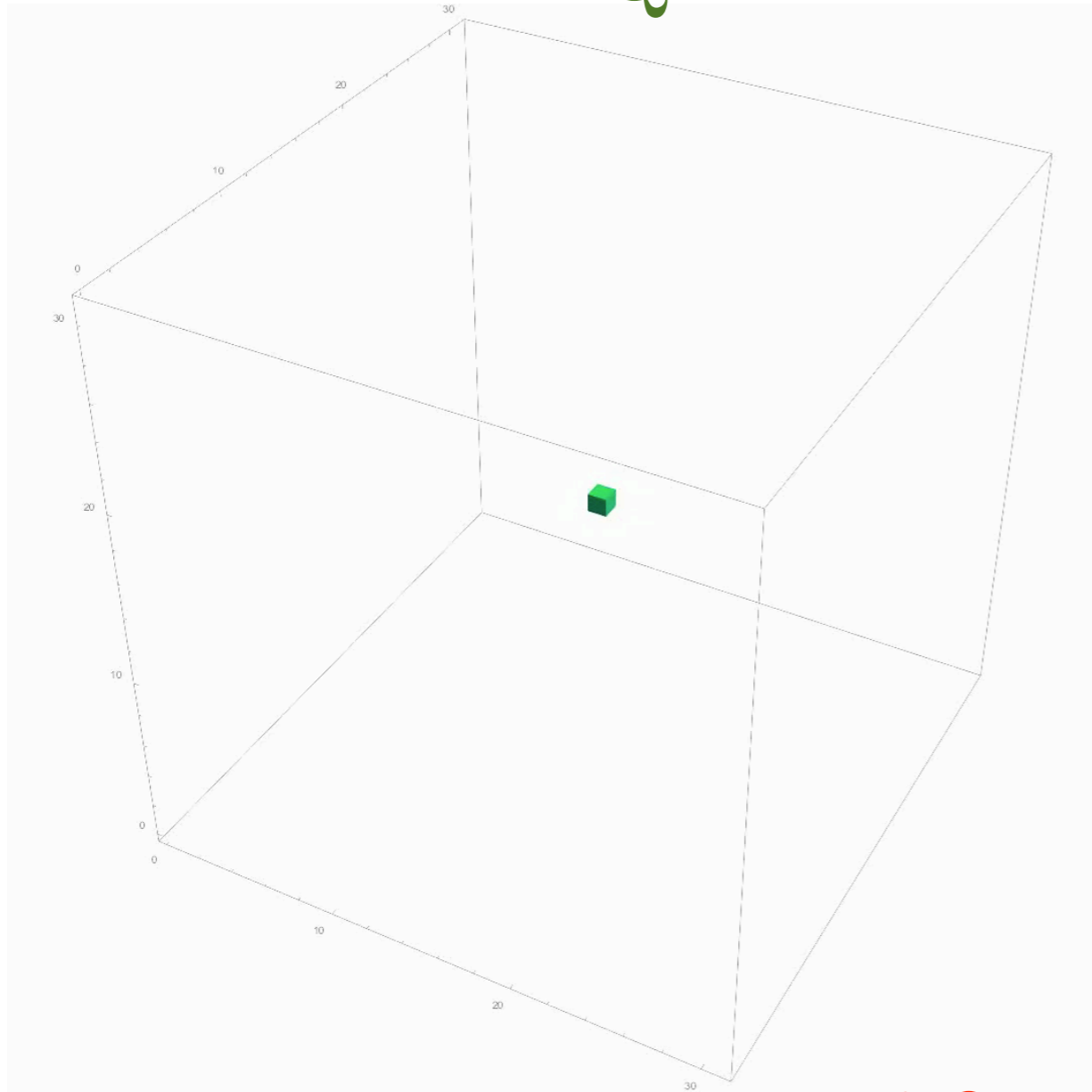
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Fourier \longrightarrow

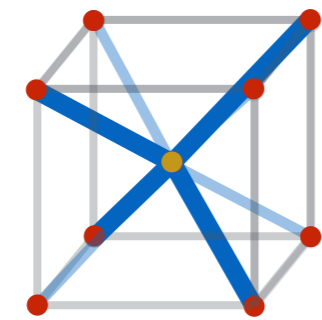
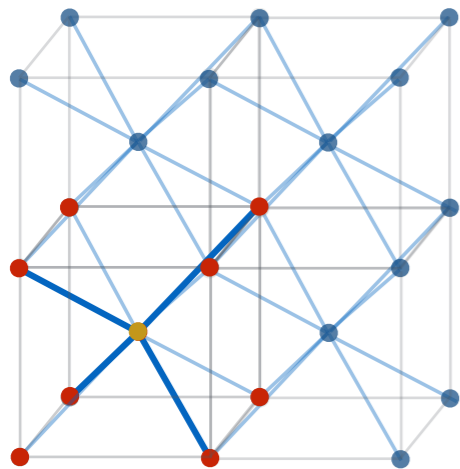
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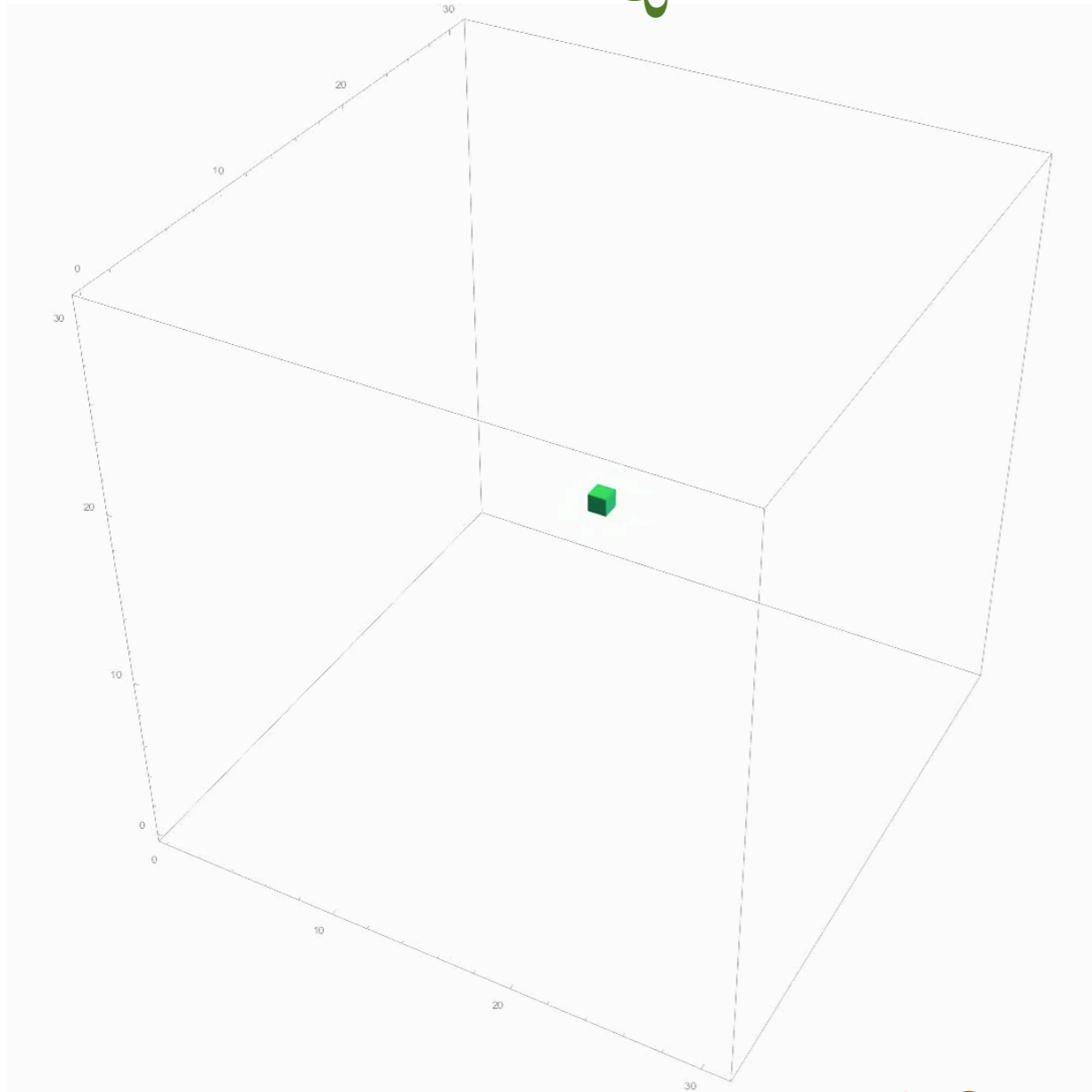
$h \in S$ → generators
 T_h → translations

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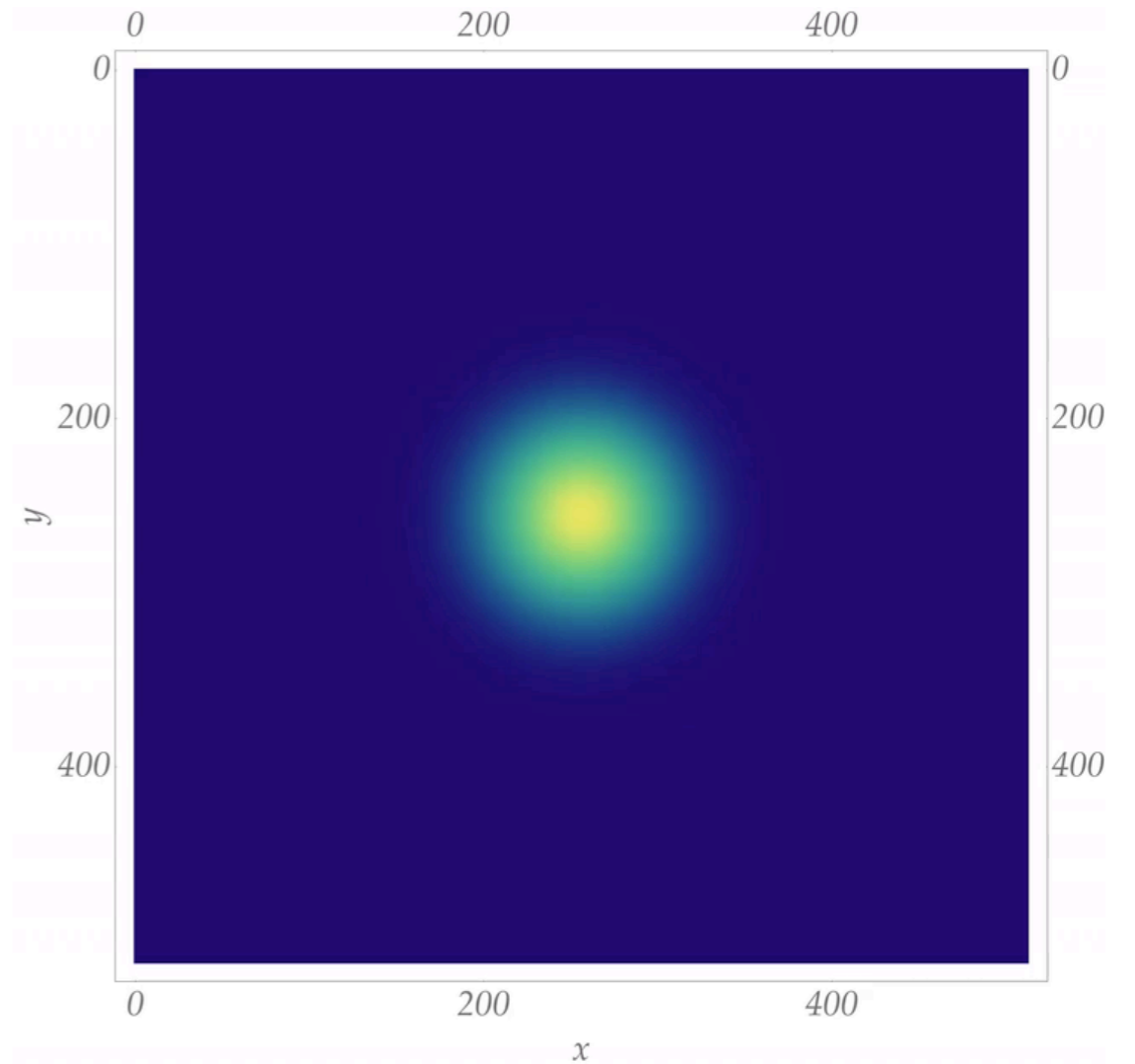
\mathbf{B} → Brillouin zone

Momentum cutoff

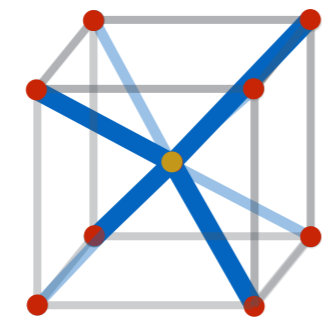
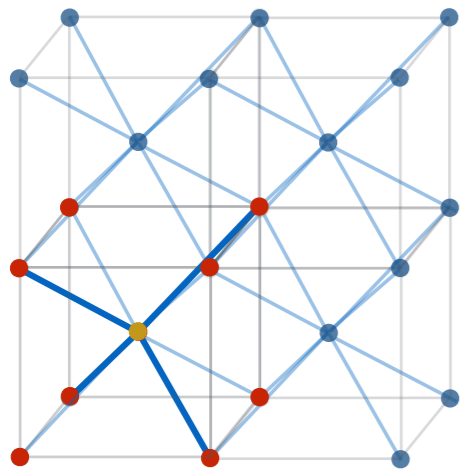
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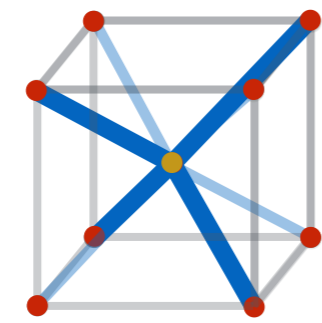
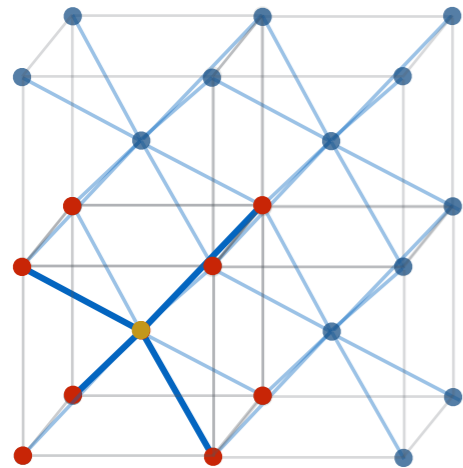
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Weyl Quantum Walk



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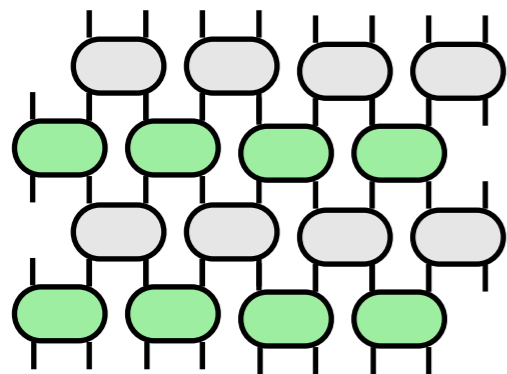
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(1+1)-dimensional Dirac Quantum Walk



$$\mathbf{U} = \begin{pmatrix} nS & -im \\ -im & nS^\dagger \end{pmatrix} \quad S\psi(x) = \psi(x+1) \quad n^2 + m^2 = 1, \quad 0 \leq m \leq 1 \quad \begin{matrix} \nearrow \\ \text{bounded rest} \\ \text{mass} \end{matrix}$$

$$\xrightarrow{m, k \rightarrow 0} i\partial_t \psi(k, t) = \begin{pmatrix} -k & m \\ m & k \end{pmatrix} \psi(k, t) \quad \text{Dirac equation}$$

$$\cos^2(\omega_A) = (1 - m^2) \cos^2(k) \quad \xrightarrow{m, k \rightarrow 0} \omega_A^2 - k^2 = m^2$$

Relativistic dispersion relation

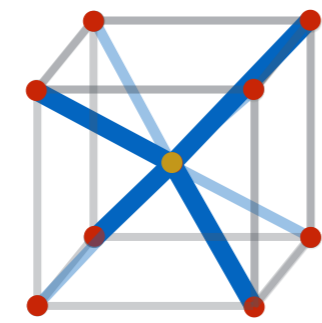
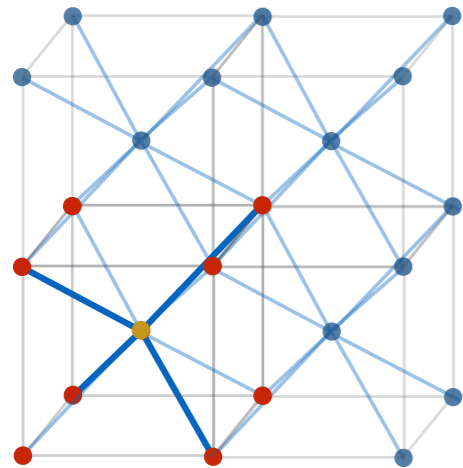
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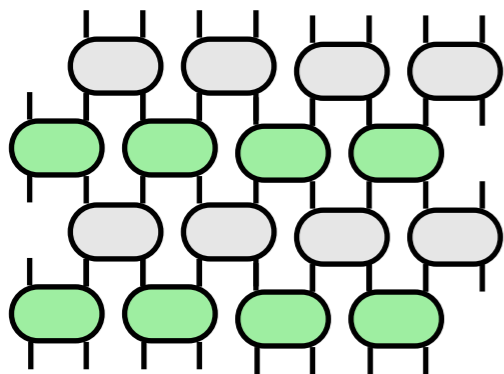
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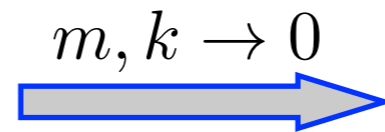
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QCA and Lorentz transformations

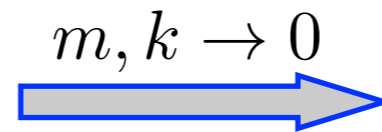
Quantum Cellular Automata



Free Quantum Field Theory
Lorentz invariant equations

QCA and Lorentz transformations

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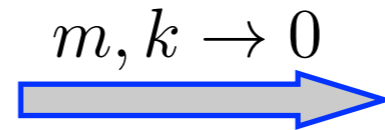
**The observer is the same!
Boosted observer?**



Relativity

QCA and Lorentz transformations

Quantum Cellular Automata



Free Quantum Field Theory
Lorentz invariant equations

**The observer is the same!
Boosted observer?**



Relativity

1D Dirac QW dispersion relation

$$\cos^2(\omega) = (1 - m^2) \cos^2(k)$$

classical kinematics
emergent from the automaton

non Lorentz invariant

Lorentz transformation

$$\begin{pmatrix} \omega' \\ k' \end{pmatrix} = \gamma \begin{pmatrix} 1 & -\beta \\ -\beta & 1 \end{pmatrix} \begin{pmatrix} \omega \\ k \end{pmatrix}$$

$$\gamma := \frac{1}{\sqrt{1 - \beta^2}}$$

**Violations of Lorentz invariance
at ultra-relativistic scales**

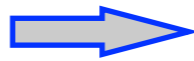
different
transformation

or

privileged
reference frame

Deformed relativity

A simple speculation
from Quantum Gravity



In whose reference frame is the Planck energy the threshold for new phenomena?

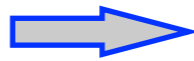
Preserve relativity principle
Lorentz group

AND

invariant energy scale

Deformed relativity

A simple speculation
from Quantum Gravity



In whose reference frame is the Planck energy the threshold for new phenomena?

Preserve relativity principle
Lorentz group

AND

invariant energy scale



Modify the action of Lorentz group

non-linear action in momentum space

$$L_{\beta}^D := \mathcal{D}^{-1} \circ L_{\beta} \circ \mathcal{D},$$

$$L_{\beta} = \gamma \begin{pmatrix} 1 & -\beta \\ -\beta & 1 \end{pmatrix}$$

momentum space
is more fundamental

\mathcal{D} is a non-linear map

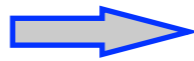
- $J_{\mathcal{D}}(0, 0) = I$

- singular point \iff invariant energy

- invertible

Deformed relativity

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which \mathcal{D} ?

G. Amelino-Camelia, Physics Letters B 510, 255 (2001).

J. Magueijo, L. Smolin, Phys. Rev. Lett. 88, 190403 (2002).

Deformed relativity and QW

Dirac QW dispersion relation

$$\cos^2(\omega) = (1 - m^2) \cos^2(k) \quad \Rightarrow \quad \frac{\sin^2(\omega)}{\cos^2(k)} - \tan^2(k) = m^2$$

$$\tilde{\omega}^2 - \tilde{k}^2 = m^2$$

A. Bibeau-Delisle, AB, G. M. D'Ariano, P. Perinotti, A. Tosini, EPL 109, 50003 (2015).

AB, G. M. D'Ariano, P. Perinotti, Phil. Trans. R. Soc. A 374 20150232 (2016).

AB, G. M. D'Ariano, P. Perinotti, Phys. Rev. A 94, 042120 (2016).

Deformed relativity and QW

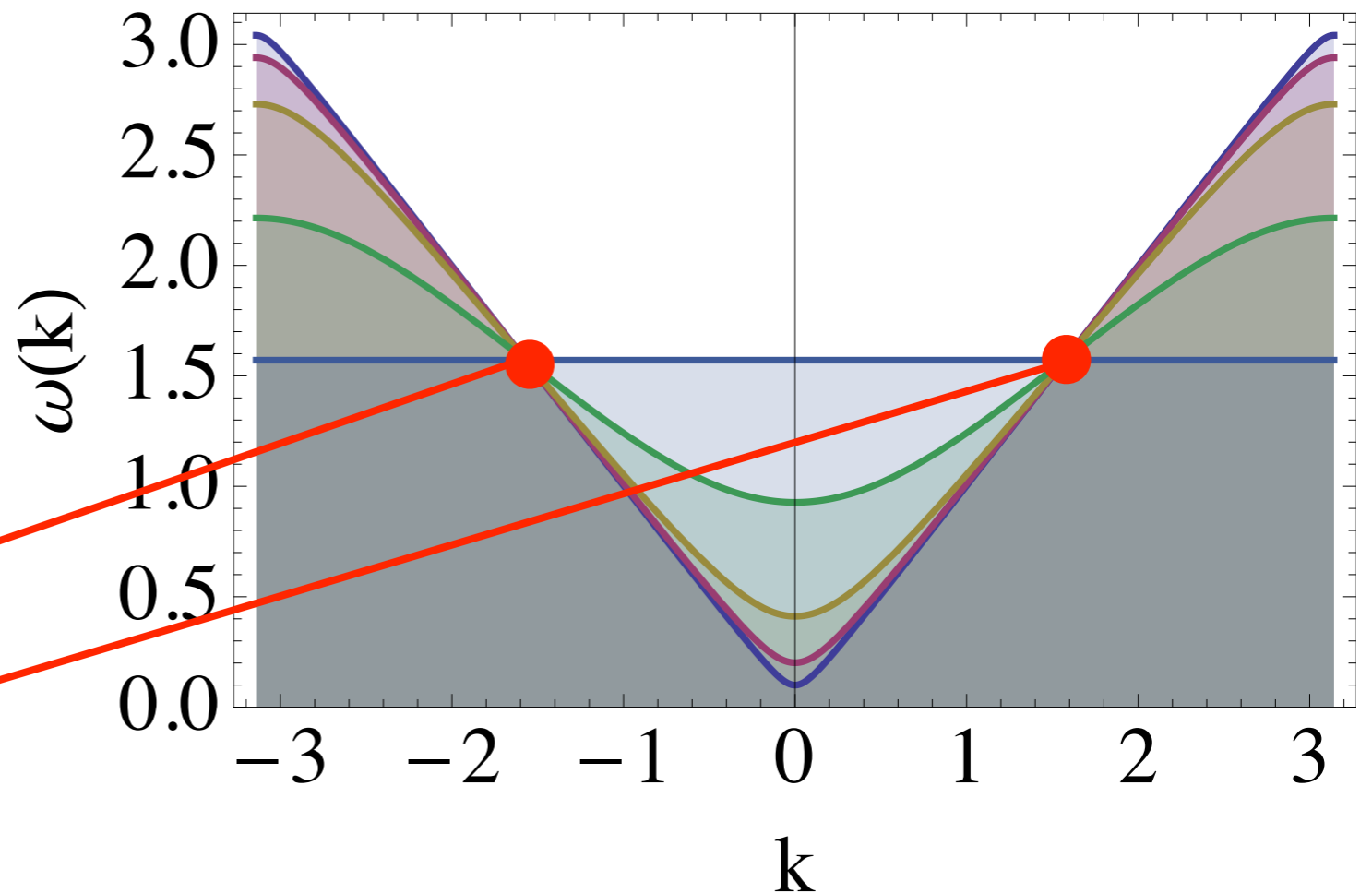
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$$\mathcal{D} \begin{pmatrix} \omega \\ k \end{pmatrix} = \begin{pmatrix} \frac{\sin(\omega)}{\cos(k)} \\ \tan(k) \end{pmatrix}$$

$$-\frac{\pi}{2} \leq k \leq \frac{\pi}{2}$$



$$\omega_{inv} = \frac{\pi}{2}$$

A. Bibeau-Delisle, AB, G. M. D'Ariano, P. Perinotti, A. Tosini, EPL 109, 50003 (2015).

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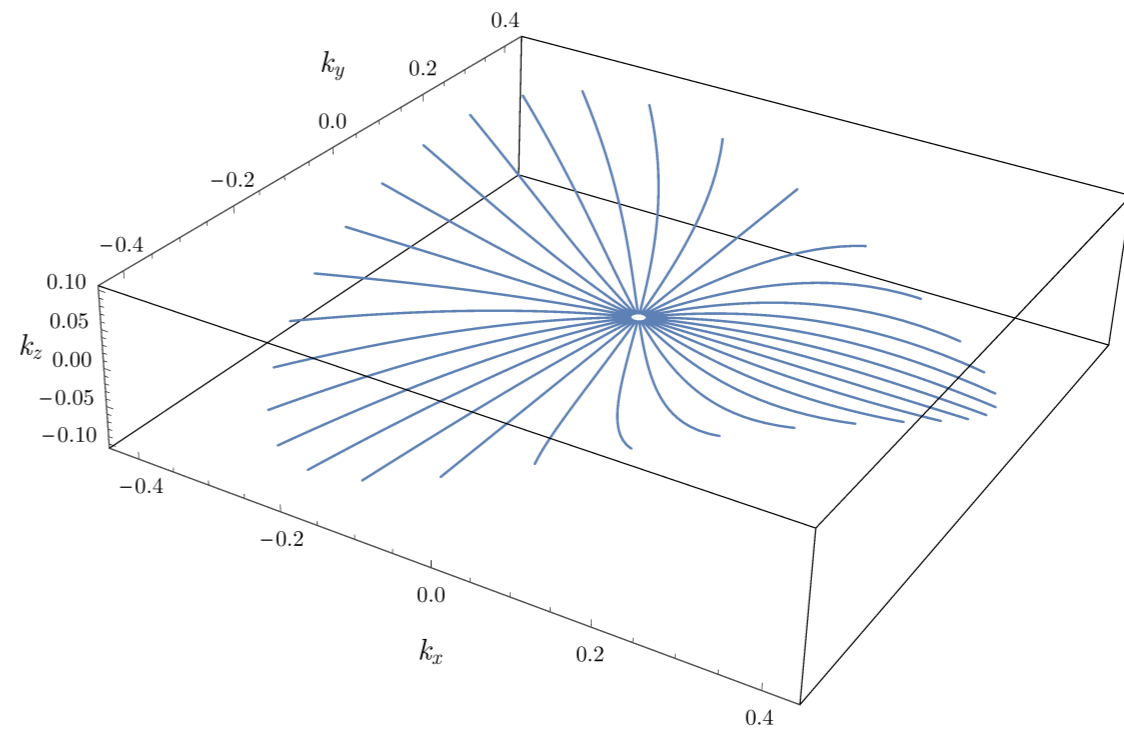
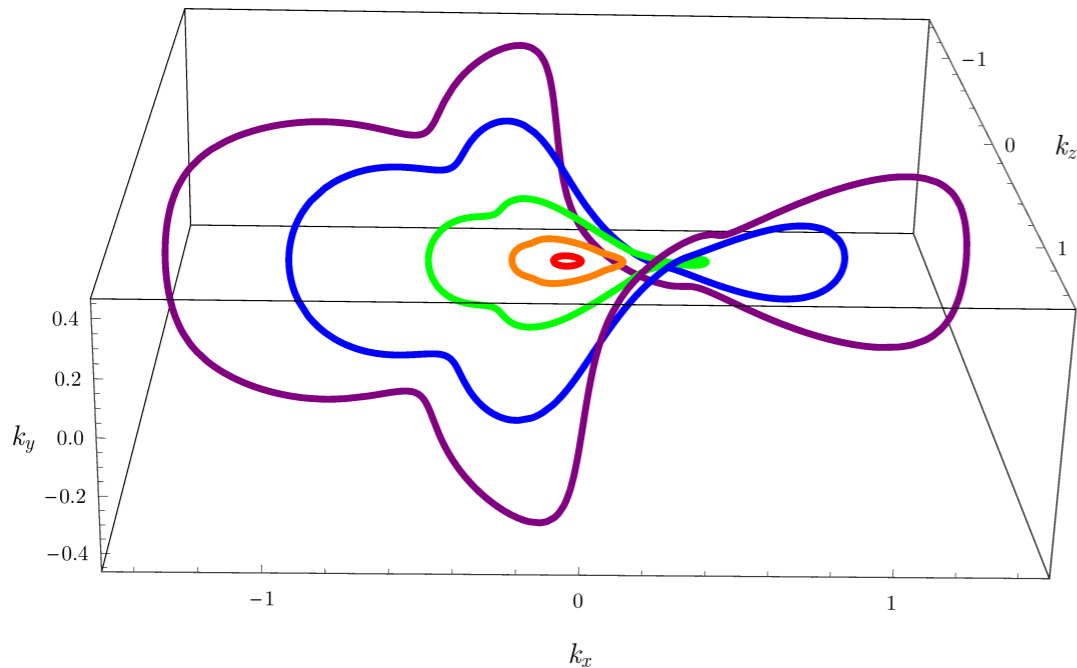
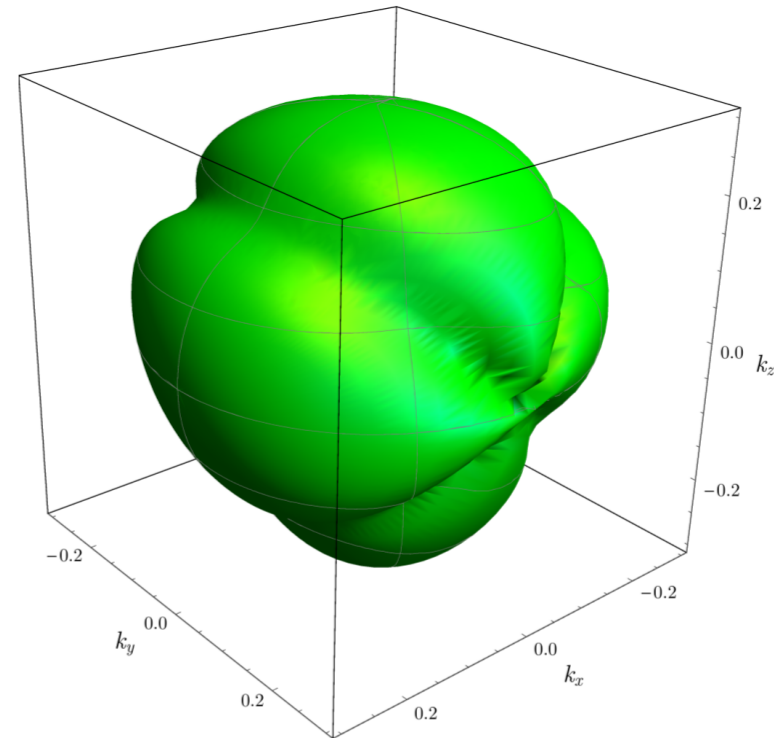
AB, G. M. D'Ariano, P. Perinotti, Phys. Rev. A 94, 042120 (2016).

Deformed relativity and QW: 3 spatial dimensions

Weyl QW dispersion relation

$$n_{\mu}(k)\sigma^{\mu}\psi(k) = 0$$

$$n(k) := \begin{pmatrix} \sin \omega \\ s_x c_y c_z + c_x s_y s_z \\ c_x s_y c_z - s_x c_y s_z \\ c_x c_y s_z + s_x s_y c_z \end{pmatrix} \quad \begin{aligned} c_i &= \cos\left(\frac{k_i}{\sqrt{3}}\right) \\ s_i &= \sin\left(\frac{k_i}{\sqrt{3}}\right) \end{aligned}$$



AB, G. M. D'Ariano, P. Perinotti, Foundations of Physics, 47, 8, 1065(2017).

AB, G. M. D'Ariano, P. Perinotti, Phil. Trans. R. Soc. A 374 20150232 (2016).

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Deformed relativity in position space

The model is defined in
the momentum space



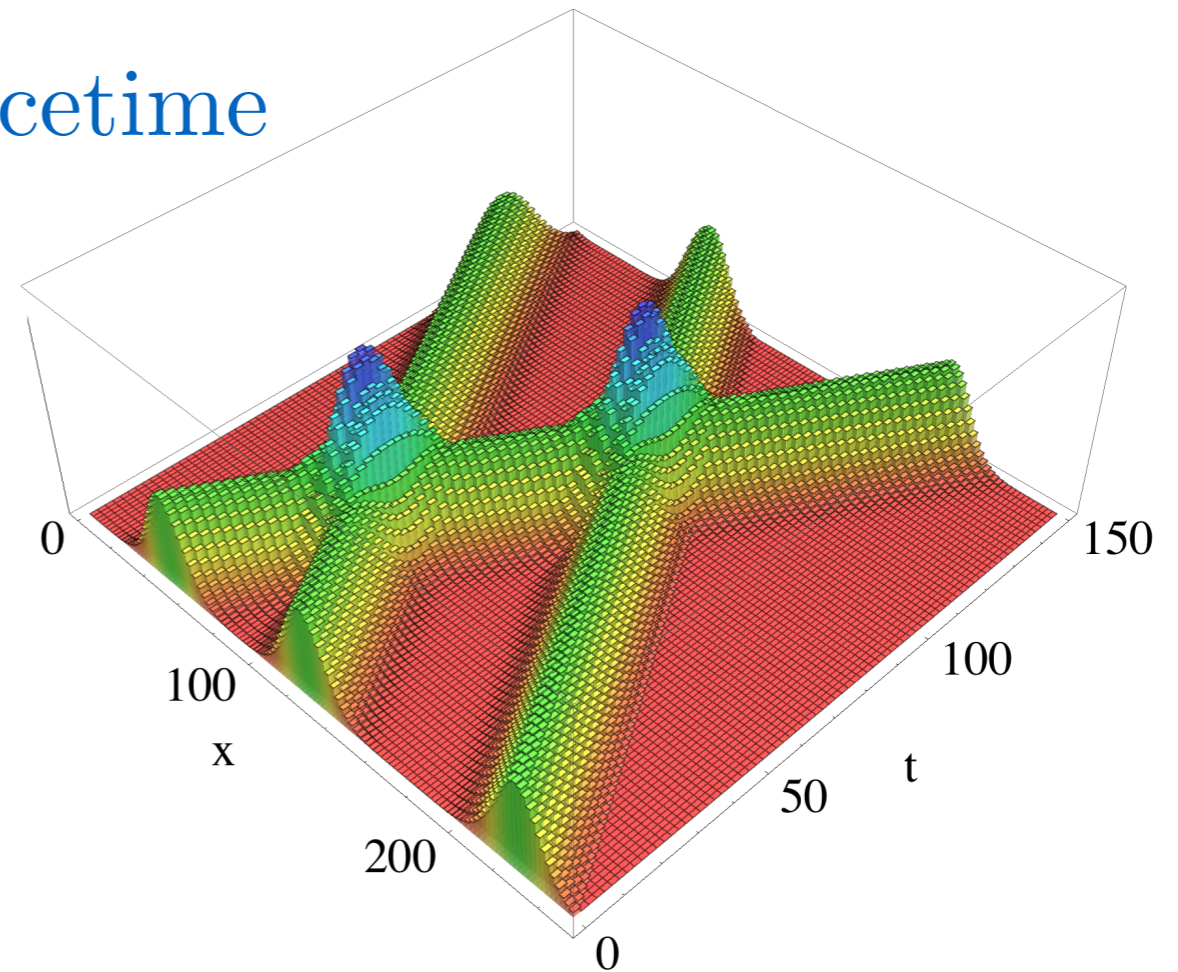
Transformations in the position space?

Operational toy model of spacetime

coincidence of wavepackets



point in spacetime



A. Bibeau-Delisle, AB, G. M. D'Ariano, P. Perinotti, A. Tosini, EPL 109, 50003 (2015).

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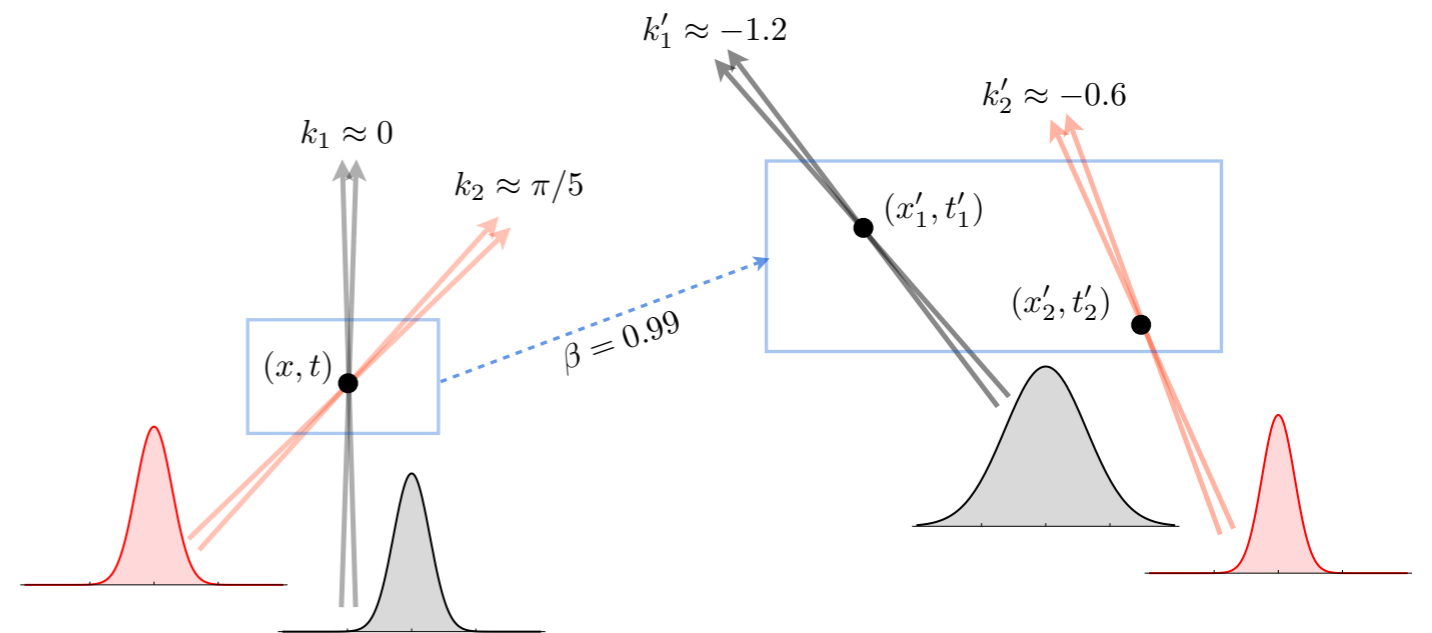
Transformations in the position space?

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transformation of the coincidence points:
$$\begin{pmatrix} t' \\ x' \end{pmatrix} \approx \begin{pmatrix} -\partial_{\omega'} k & \partial_{k'} k \\ \partial_{\omega'} \omega & -\partial_{k'} \omega \end{pmatrix} \begin{pmatrix} t \\ x \end{pmatrix}$$
 $k' = k'_0$

momentum-dependent spacetime

Relative Locality

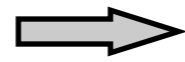
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AB, G. M. D'Ariano, P. Perinotti, Phys. Rev. A 94, 042120 (2016).

Deformed relativity in position space

Relative locality



Observer-dependent spacetime

“before”

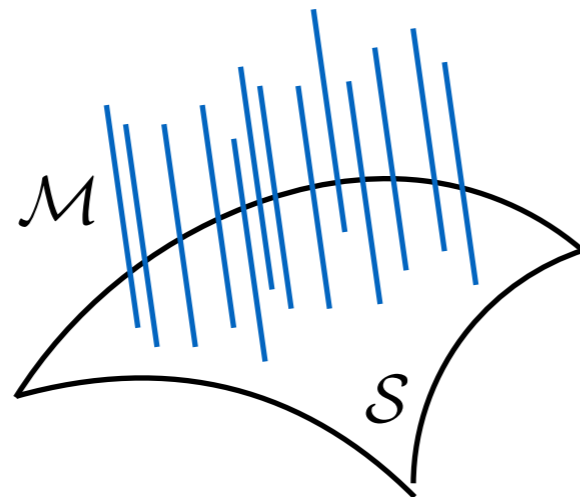
spacetime \mathcal{S}

“objective arena”

flat momentum space \mathcal{M}

phase space

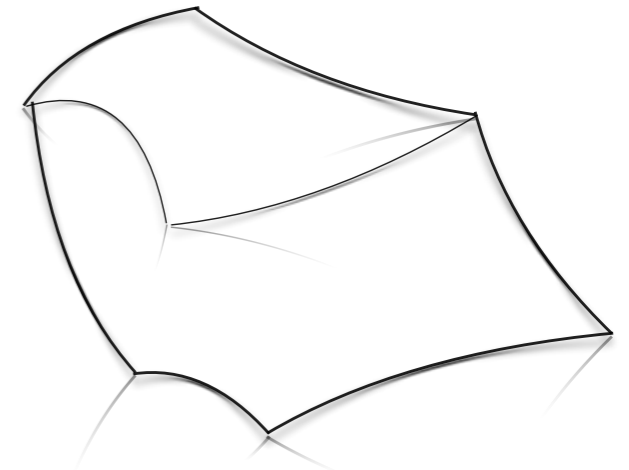
$$\mathcal{P} = \mathcal{T}^* \mathcal{S}$$



“after”

phase space

$$\mathcal{P} \neq \mathcal{T}^* \mathcal{S}$$



no canonical projection that
gives a description of
processes in spacetime

R. Schutzhold, W. G. Unruh, JETP Lett. 78, 431 (2003).

G. Amelino-Camelia, L. Freidel, J. Kowalski-Glikman, L. Smolin, Phys. Rev. D 84, 084010 (2011).

A. Bibeau-Delisle, AB, G. M. D’Ariano, P. Perinotti, A. Tosini, EPL 109, 50003 (2015).

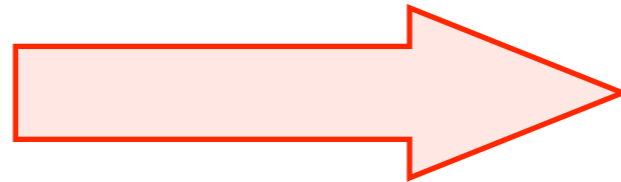
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A final overlook

Main idea

Quantum
theory



“Quantum computational
field theory”

Quantum “ab initio” theory of dynamics

A final overlook

Main idea

Quantum
theory

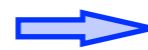


“Quantum computational
field theory”

Quantum “ab initio” theory of dynamics

QCA model

Free Quantum Field Theory



Interactions

Energy?
Momentum?

A final overlook

Main idea

Quantum theory



“Quantum computational field theory”

Quantum “ab initio” theory of dynamics

QCA model

Free Quantum Field Theory



Interactions

Energy?
Momentum?

Boost?

Deformed relativity

momentum space DSR



operational toy-model
of spacetime



emergent
spacetime