

Probing Light Hidden Sectors with Pulsar Timing Arrays

7th Workshop on Theory, Phenomenology
and Experiments in Flavour Physics - Capri

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PRISMA

Cluster of Excellence
Precision Physics, Fundamental Interactions
and Structure of Matter

In this talk:

Probing hidden sectors:

Gravitational
waves from
first-order phase
transitions

Cosmological
constraints

Experimental
sensitivity to
GW spectra

Example:
Higgsed dark-
photon model

Talking to the Hidden Sector

The fourth portal

The usual portal suspects:

- Vector: (kinetic mixing, weakly coupled gauge bosons, ...)
- Scalar: (Higgs portals, ...)
- Fermion: (neutrino portal, ...)
- Any combination of the above

Talking to the Hidden Sector

The fourth portal

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Typically we forget about gravity!

The upsides:

- Cannot be switched off
- Gravitational waves may provide brand new insights

Thermalised Hidden Sectors

Phase transition must occur in a thermalised sector to produce observable GW signal

Temperature in hidden sector typically deviates from temperature of the SM thermal bath:

$$\xi_h \equiv \frac{T_h}{T_\gamma} = \left(\frac{g_{*s,\gamma} g_{*s,h}^{\text{dec}}}{g_{*s,\gamma}^{\text{dec}} g_{*s,h}} \right)^{1/3}$$

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Temperature ratio
strongly affects GW
spectrum

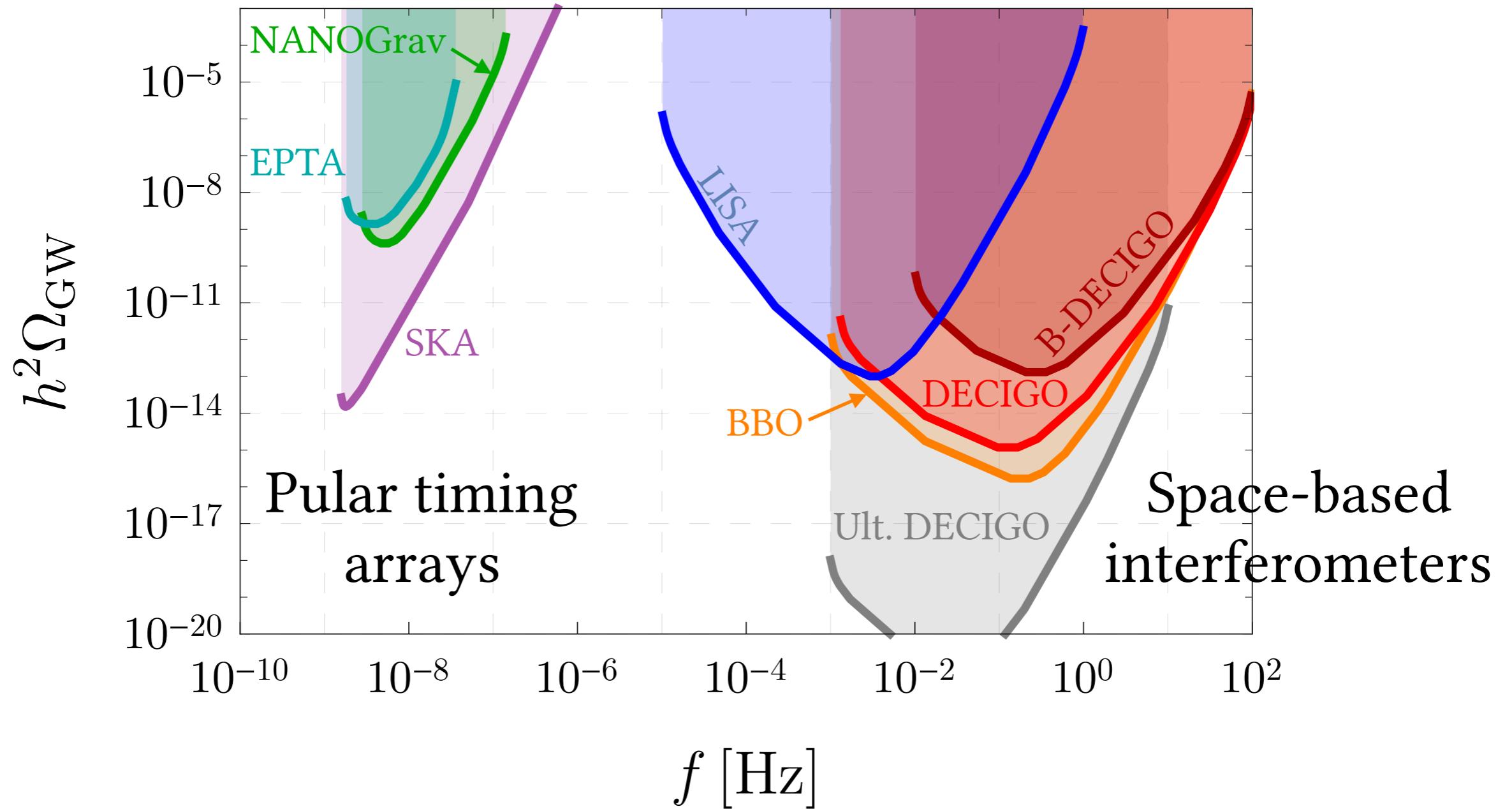
two effects


Stringent constraints
on new thermalised
keV/MeV relativistic
degrees of freedom

Gravitational waves from first-order phase transitions

Gravitational Waves

Sensitivity curves

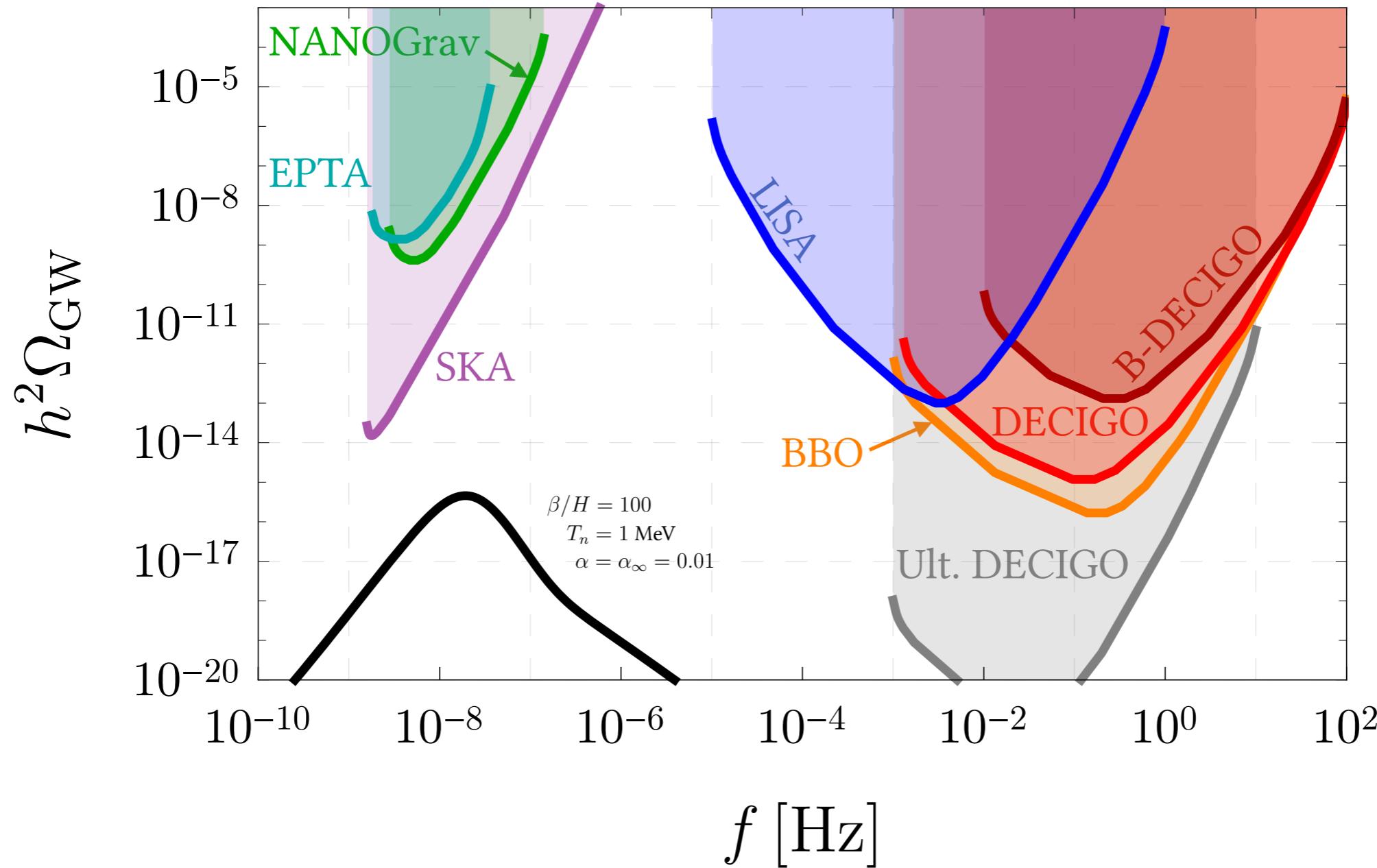


Power-law integrated sensitivity derived from experimental noise curves

[Thrane & Romano - 1310.5300]

Gravitational Waves

Sensitivity curves



Power-law integrated sensitivity derived from experimental noise curves

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GW Stochastic Background

First-order phase transitions

Phase transition controlled by effective potential:

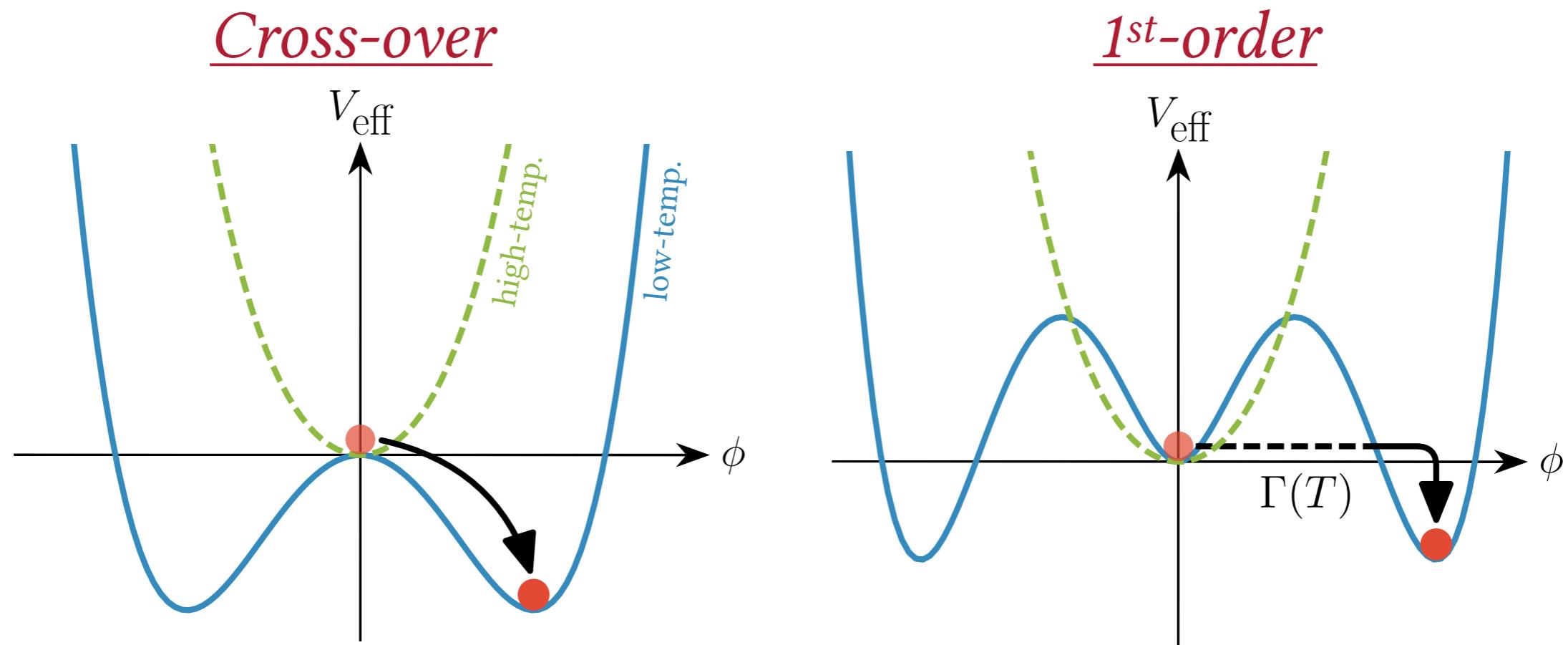
$$V_{\text{eff}}(\phi, T) = V_{\text{tree}}(\phi) + V_{\text{loop}}(\phi) + V_{\text{thermal}}(\phi, T) + V_{\text{daisy-res.}}(\phi, T)$$

GW Stochastic Background

First-order phase transitions

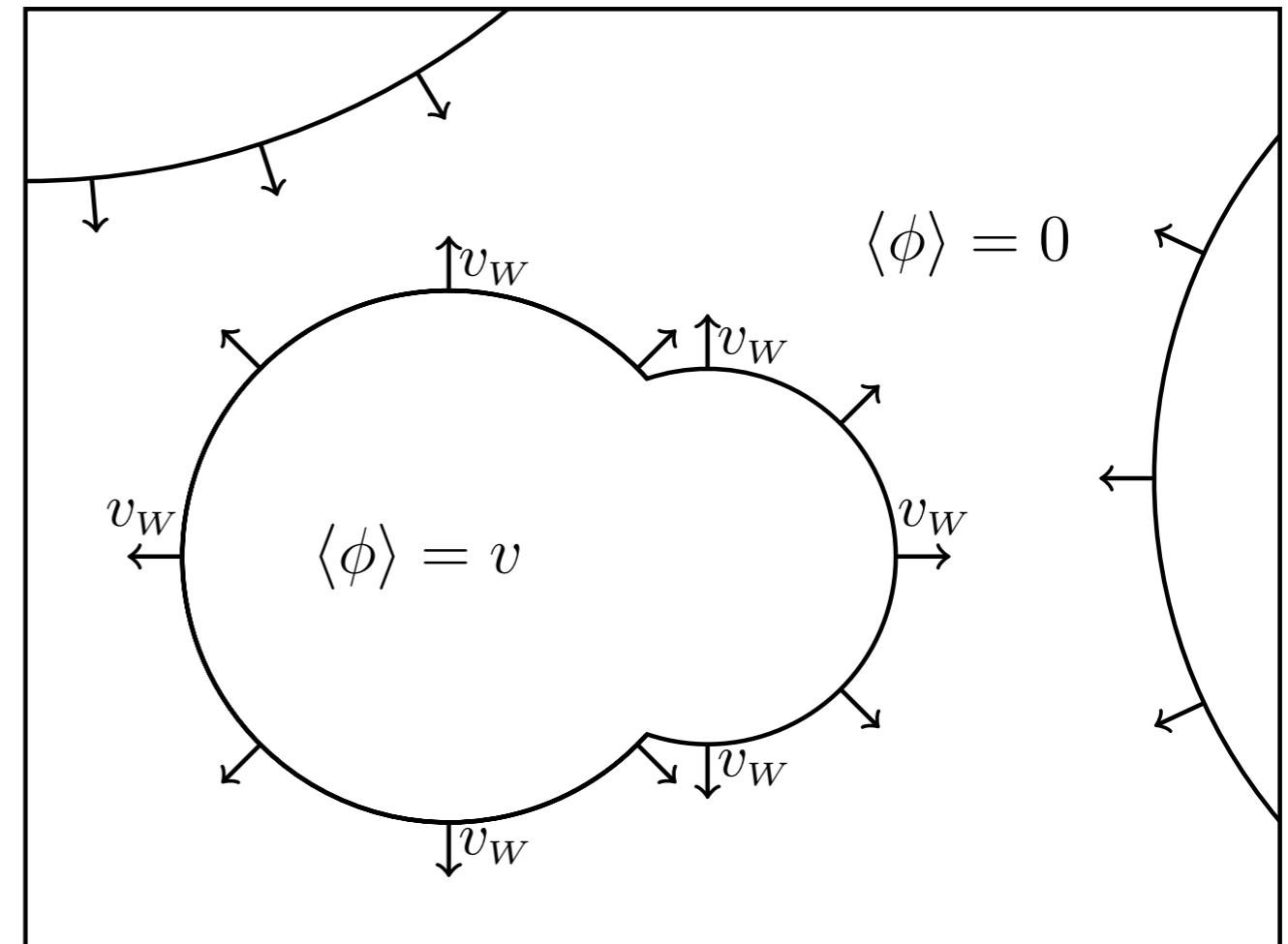
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GW Stochastic Background

Sourcing gravitational waves

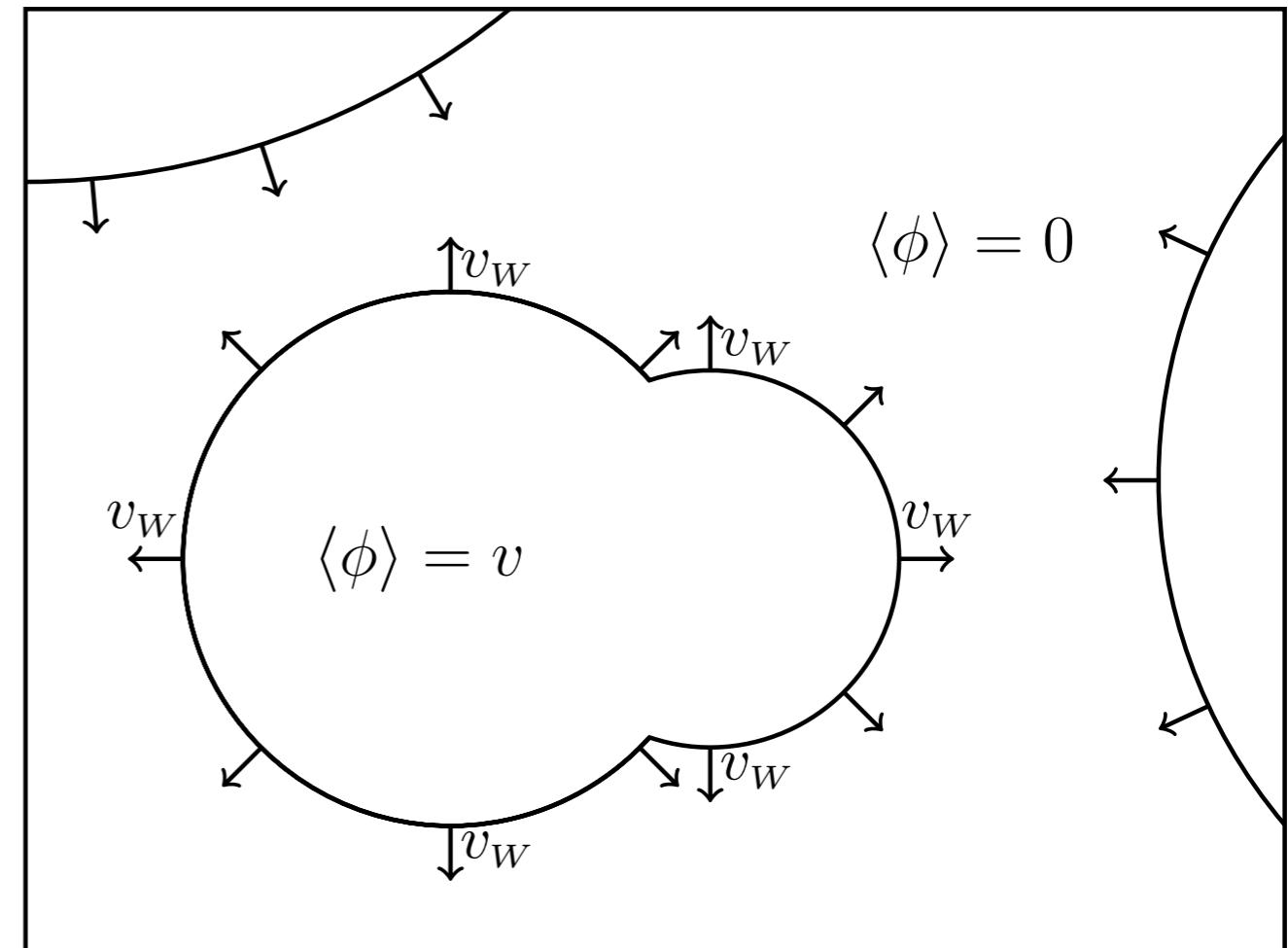


GW Stochastic Background

Sourcing gravitational waves

Nucleation temperature:

$$\Gamma(T_n) \sim H(T_n)$$



GW Stochastic Background

Sourcing gravitational waves

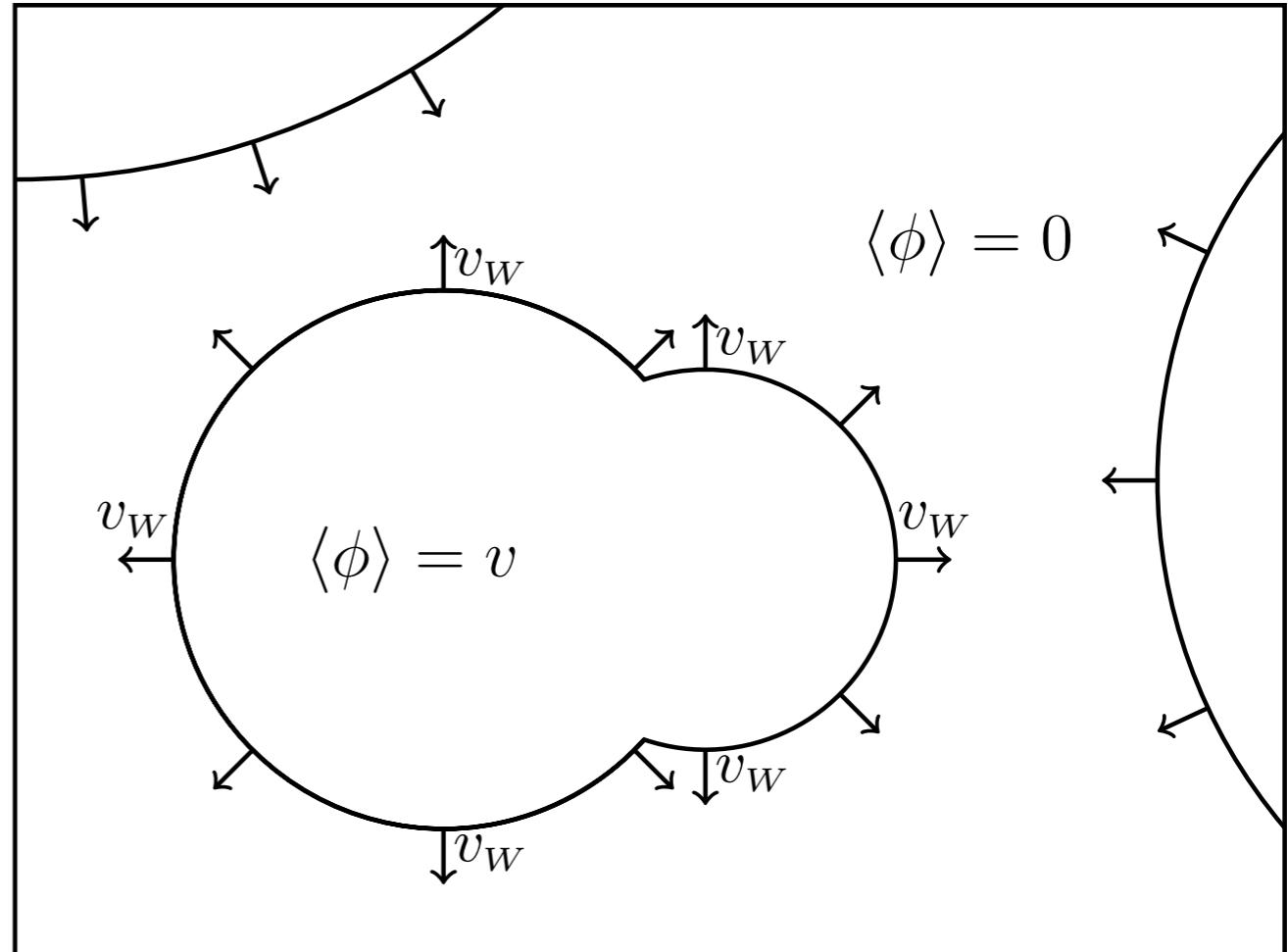
Nucleation temperature:

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Contributions to the spectrum:

[Hindmarsh et al. - 1504.03291]

Ω_ϕ : Collision of bubble walls



GW Stochastic Background

Sourcing gravitational waves

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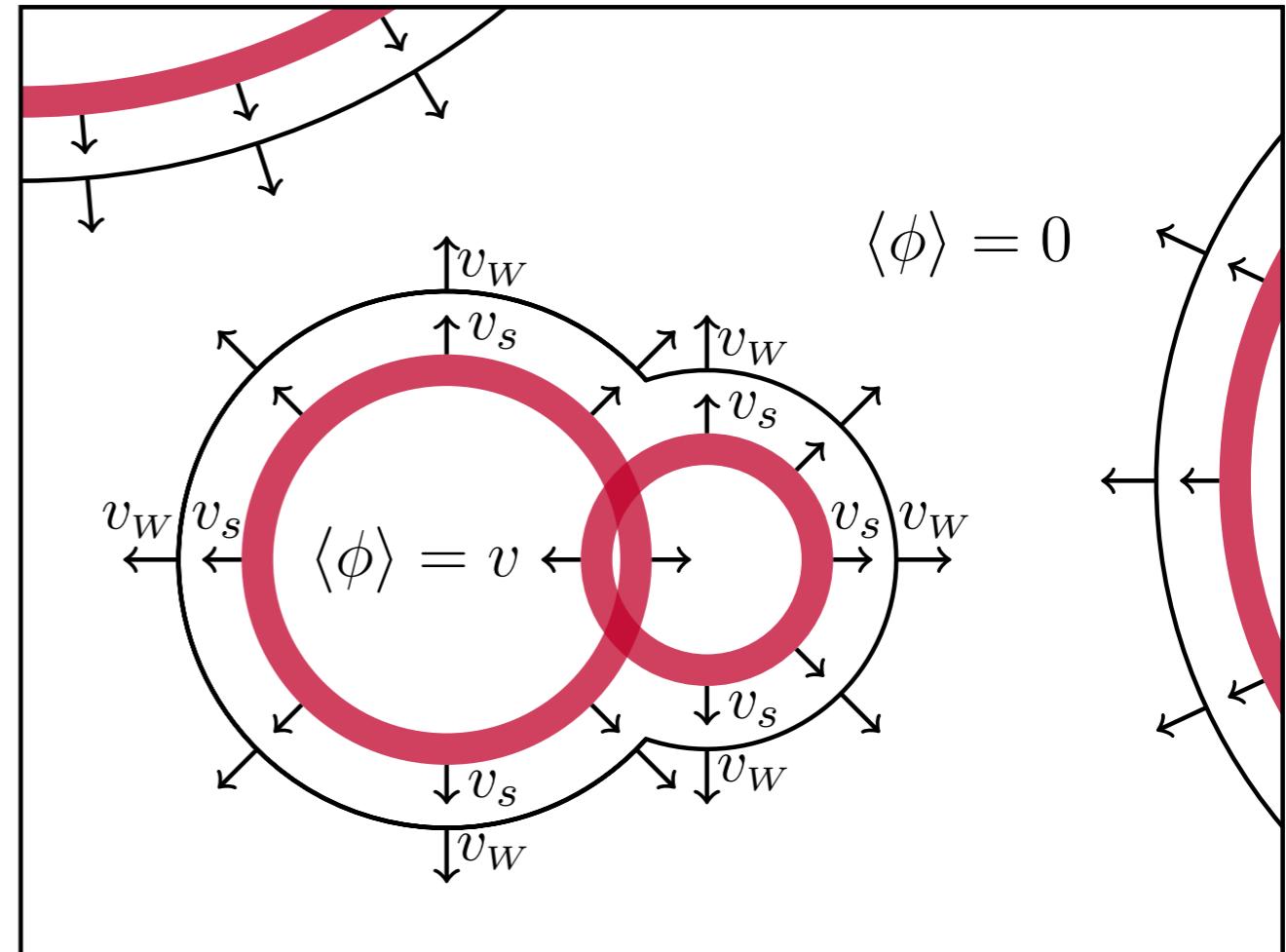
Contributions to the spectrum:

[Hindmarsh et al. - 1504.03291]

Ω_ϕ : Collision of bubble walls

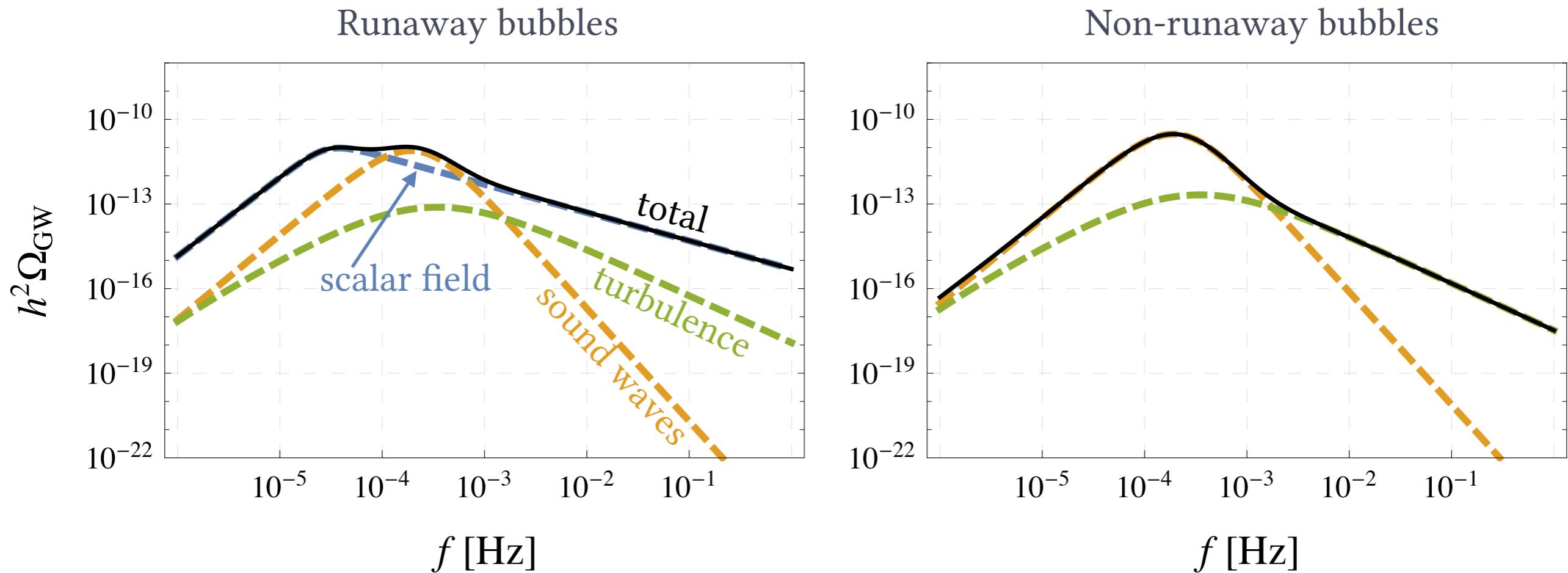
Ω_{sw} : Collision of sound-waves

Ω_{turb} : Plasma turbulence



GW Stochastic Background

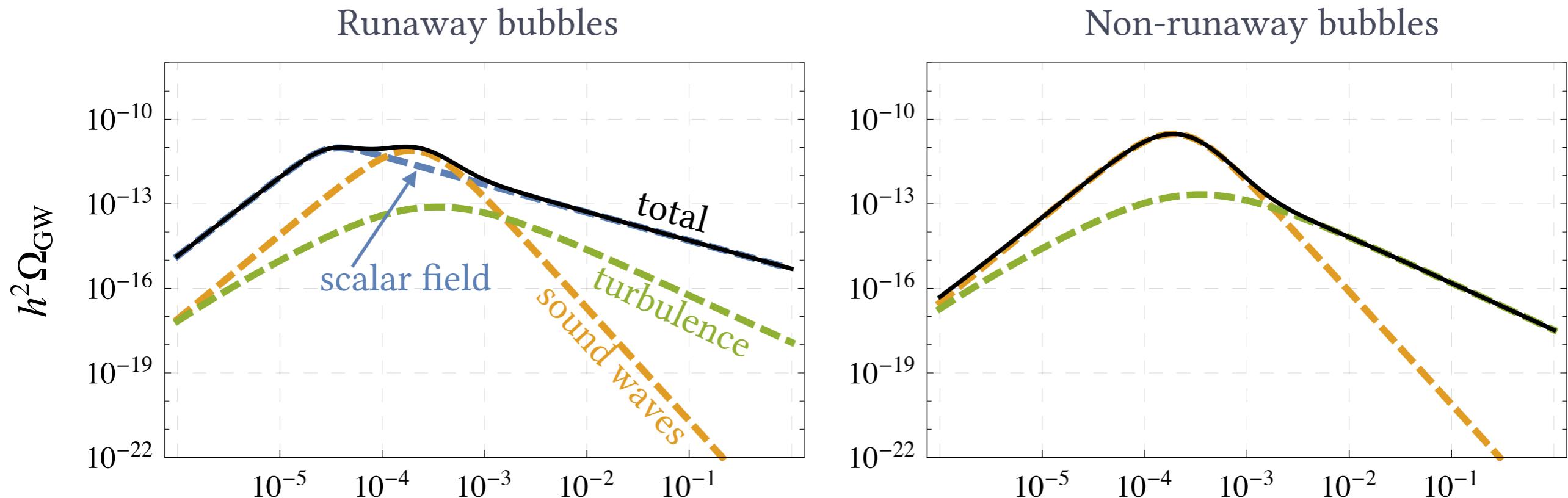
Contributions to the spectrum



1. Ω_ϕ : Initial collision of bubble walls
2. Ω_{sw} : Collision of sound-waves in the plasma
3. Ω_{turb} : Plasma turbulence after sound wave collision

GW Stochastic Background

Contributions to the spectrum



$$f [\text{Hz}]$$
$$\alpha > \alpha_\infty$$

$$f [\text{Hz}]$$
$$\alpha \leq \alpha_\infty$$

Bubble-wall friction:

$$\alpha_\infty \sim \sum_a [m_a(\phi_{\text{broken}}) - m_a(\phi_{\text{unbroken}})]$$

GW Stochastic Background

Parametrising the spectrum

Nucleation temperature:

$$T_n$$

Strength:

$$\alpha \equiv \frac{\epsilon}{\rho_{\text{rel}}}$$

Inverse time-scale:

$$\frac{\beta}{H} \equiv -T_n \left. \frac{dS}{dT} \right|_{T_n}$$

Bubble wall velocity:

$$v_w$$

- Broken-power law spectra $\Omega_{\text{GW}}(f)$ from fitting simulations [Huber et al. - 0806.1828, Hindmarsh et al. - 1504.03291, Caprini et al. - 0909.0622]
- Assuming $v_w \simeq 1$: valid for strong transitions

GW Stochastic Background

Parametrising the spectrum

$$\xi_h \neq 1$$

Nucleation temperature:

$$T_n$$

Strength:

$$\alpha \equiv \frac{\epsilon}{\rho_{\text{rel}}}$$

Inverse time-scale:

$$\frac{\beta}{H} \equiv -T_n \left. \frac{dS}{dT} \right|_{T_n}$$

Bubble wall velocity:

$$v_w$$

$$T_n = \xi_h^{-1}(T_n)_h$$

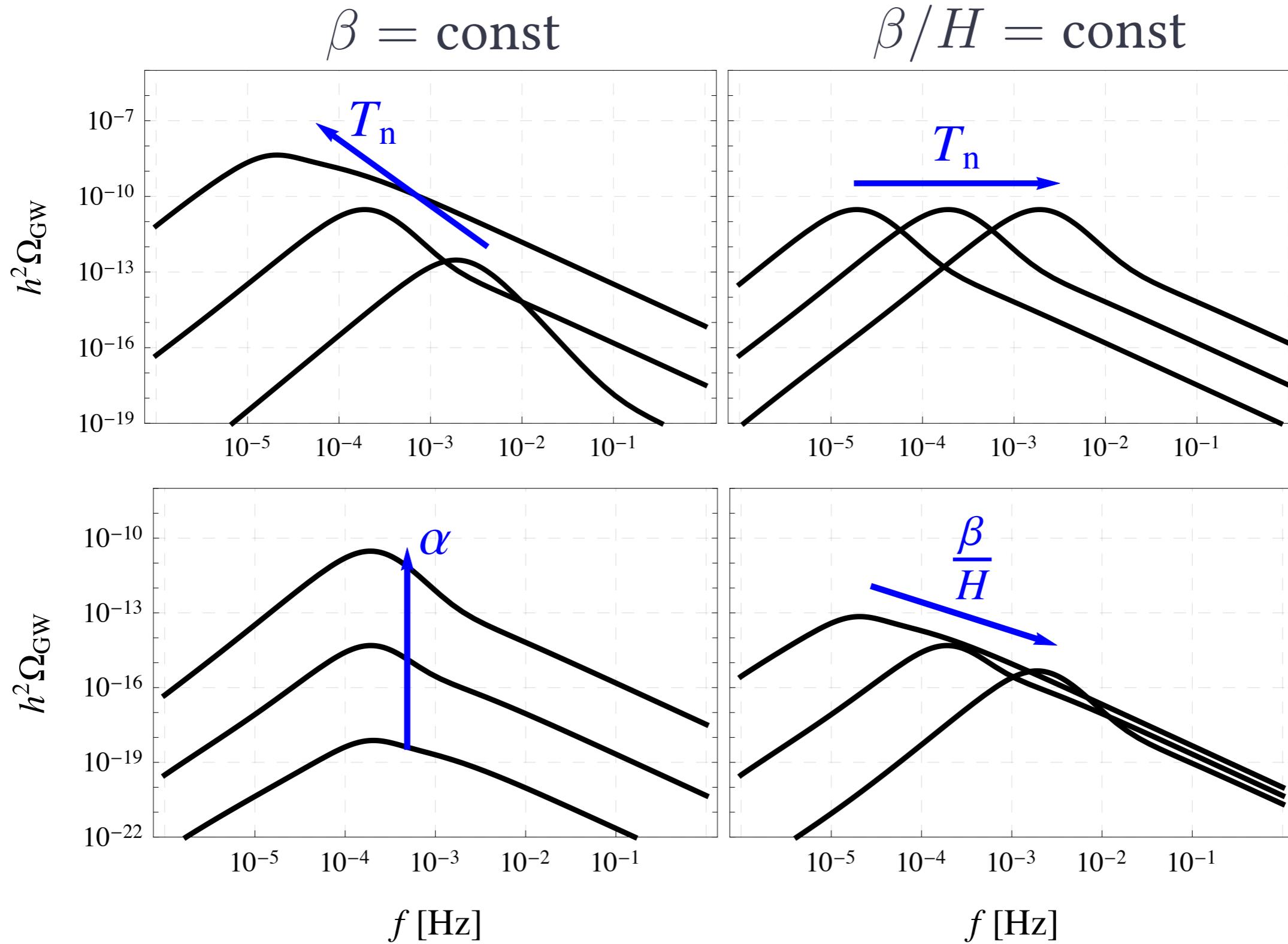
$$\alpha = \xi_h^4 \alpha_h$$

$$\frac{\beta}{H} = \xi_h^2 \left(\frac{\beta}{H} \right)_h$$

- Broken-power law spectra $\Omega_{\text{GW}}(f)$ from fitting simulations [Huber et al. - 0806.1828, Hindmarsh et al. - 1504.03291, Caprini et al. - 0909.0622]
- Assuming $v_w \simeq 1$: valid for strong transitions

GW Stochastic Background

Scaling behaviour

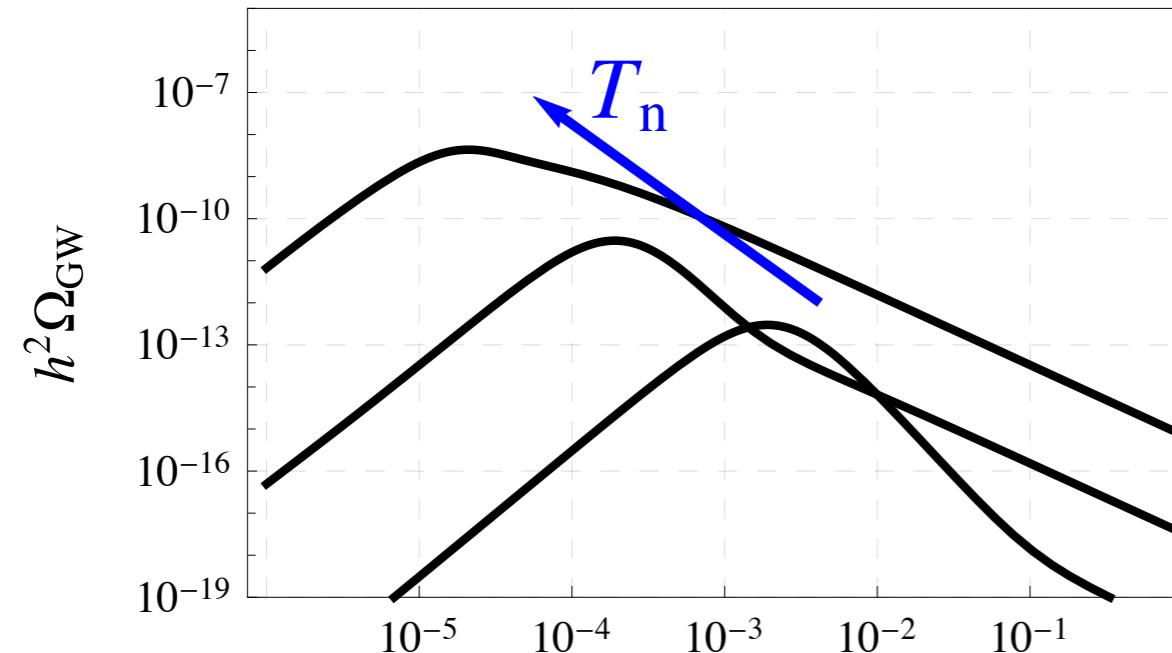


GW Stochastic Background

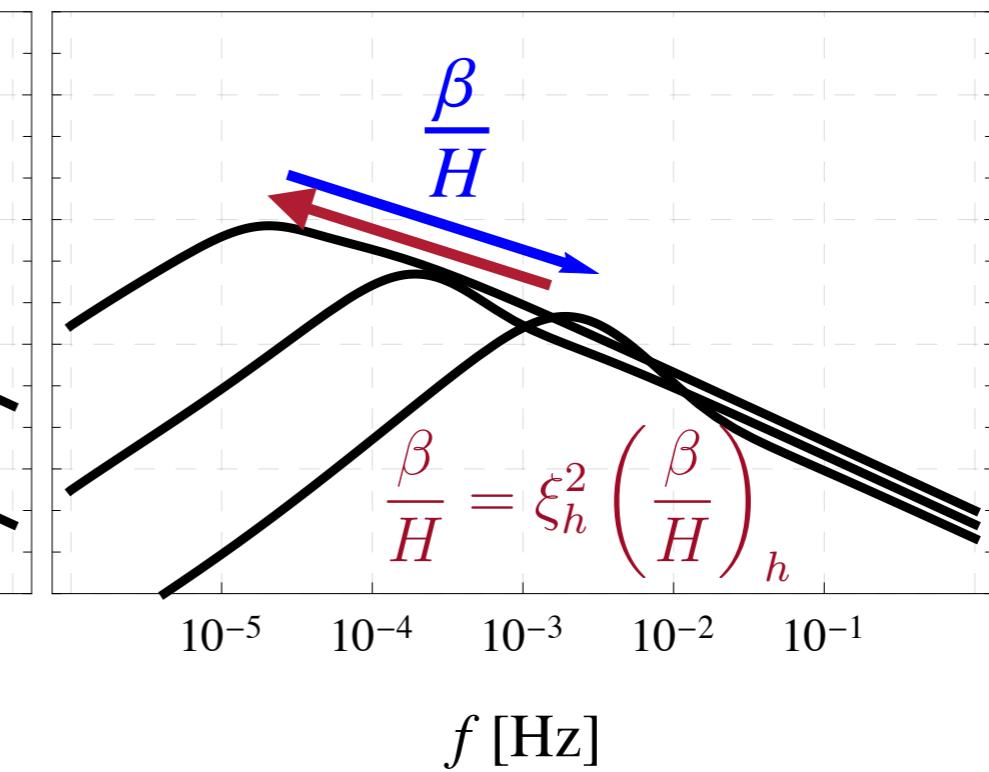
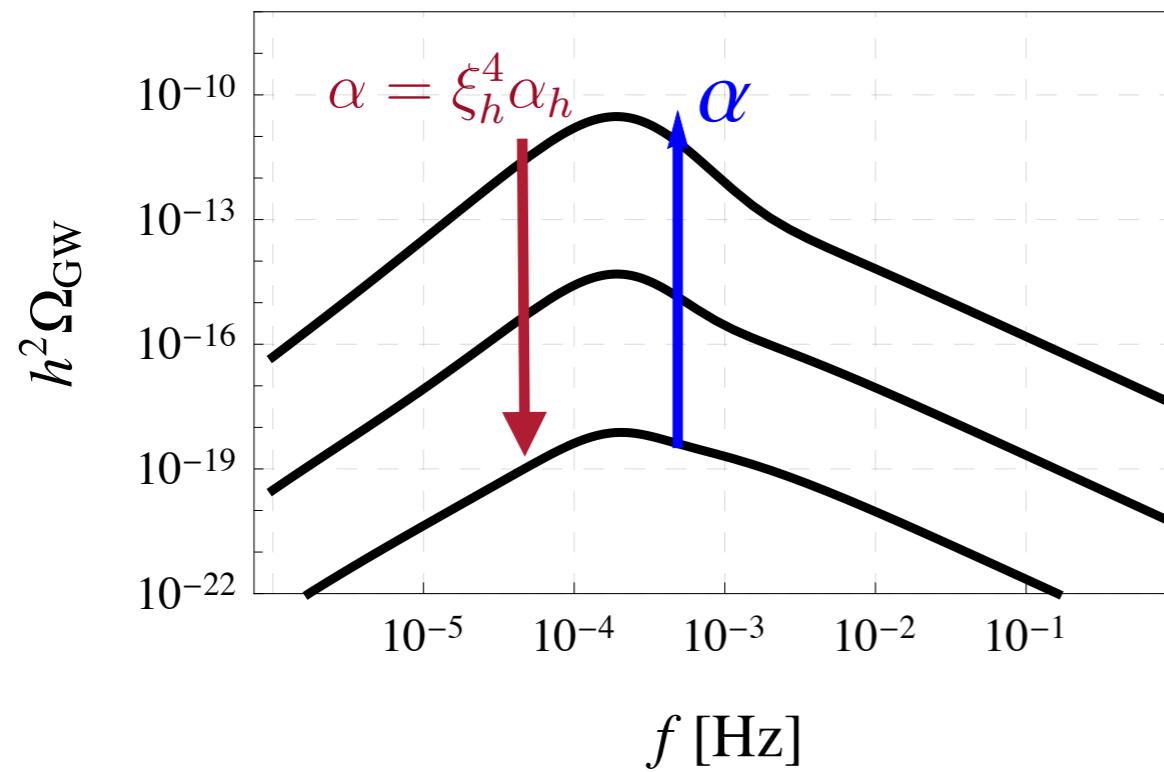
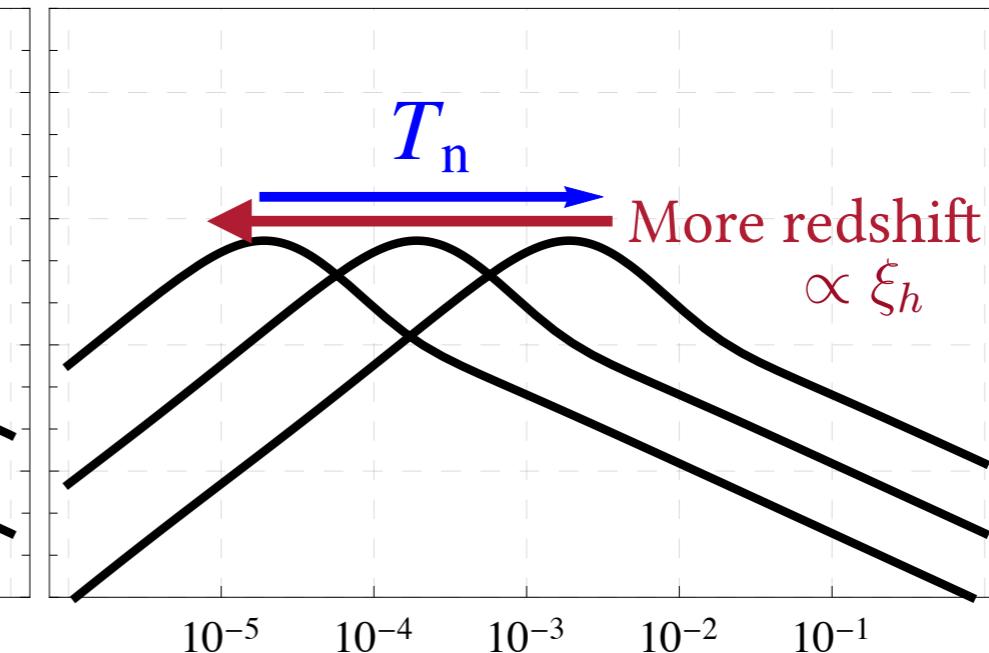
ξ_h decreasing

Scaling behaviour

$\beta = \text{const}$



$\beta/H = \text{const}$



Cosmological Constraints

Cosmological Constraints

Thermal history

Non-standard thermal history essential ingredient for MeV/keV scale phase transitions

$$N_{\text{eff}} = \frac{8}{7} \left(\frac{\rho_R - \rho_\gamma}{\rho_\gamma} \right) \left(\frac{T_\gamma^{\text{SM}}}{T_\nu^{\text{SM}}} \right)^4$$

Extra relativistic DOFs until after phase transition

Three possibilities:

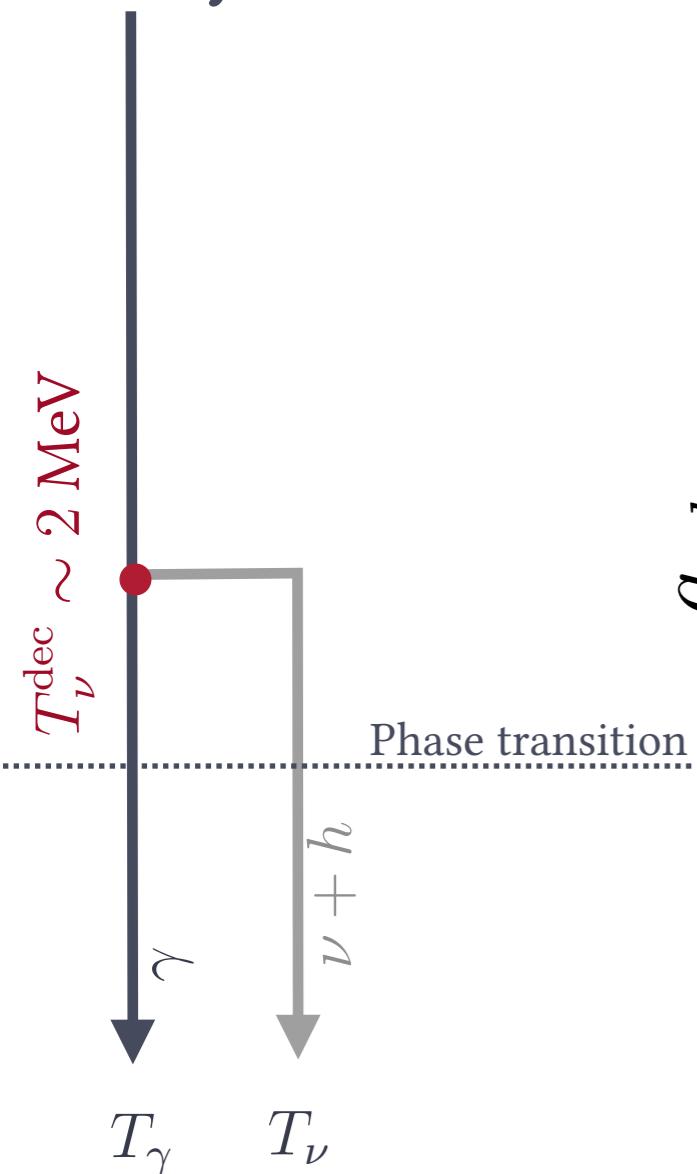
- Hidden sector remains coupled to photons or neutrinos
- Hidden sector decouples at early times
- Hidden sector decouples at early times but re-equilibrates with neutrinos below T_ν^{dec}

Cosmological Constraints

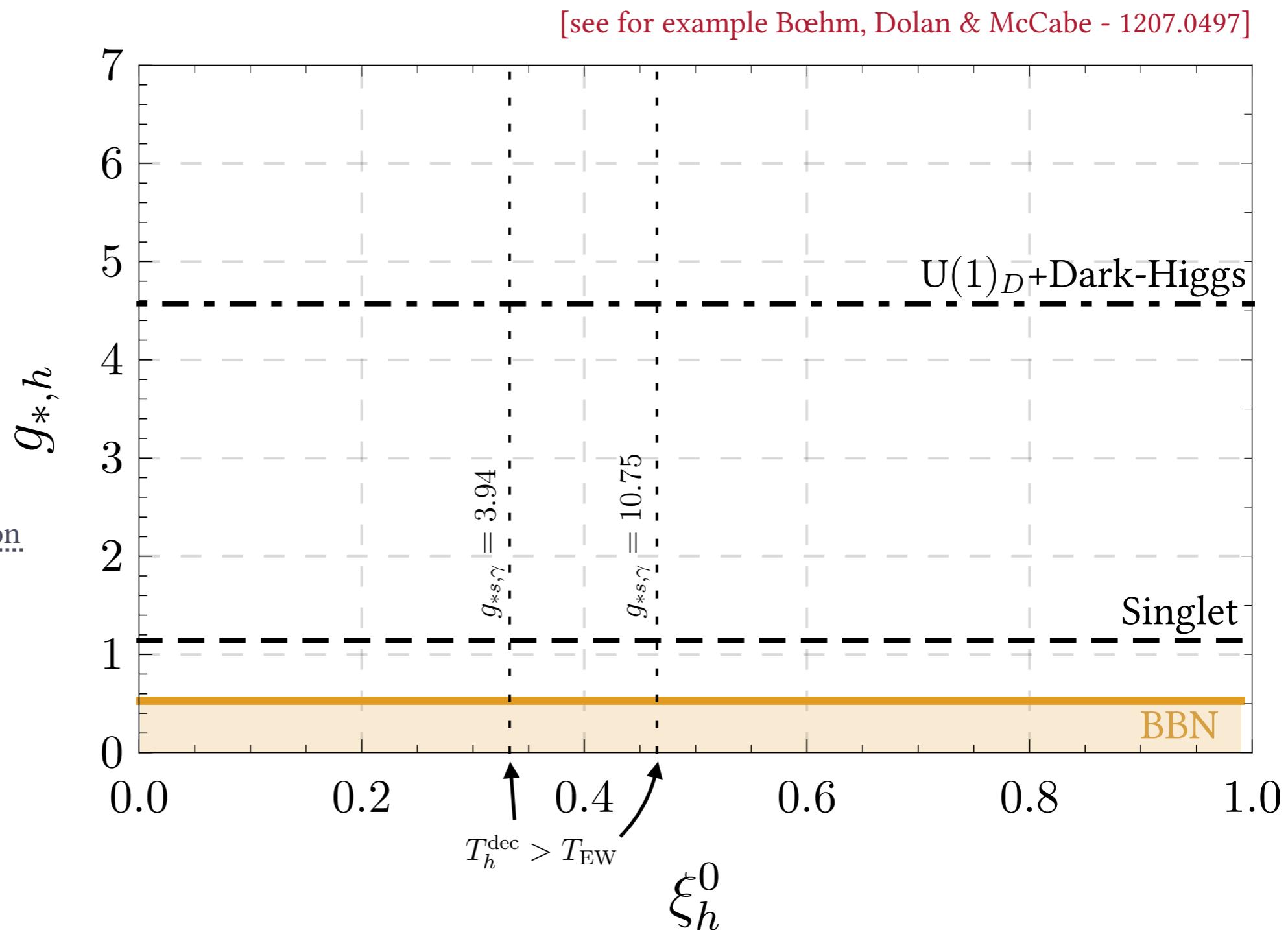
N_{eff} constraints

Hidden sector thermalised with neutrinos:

Early times



Late times

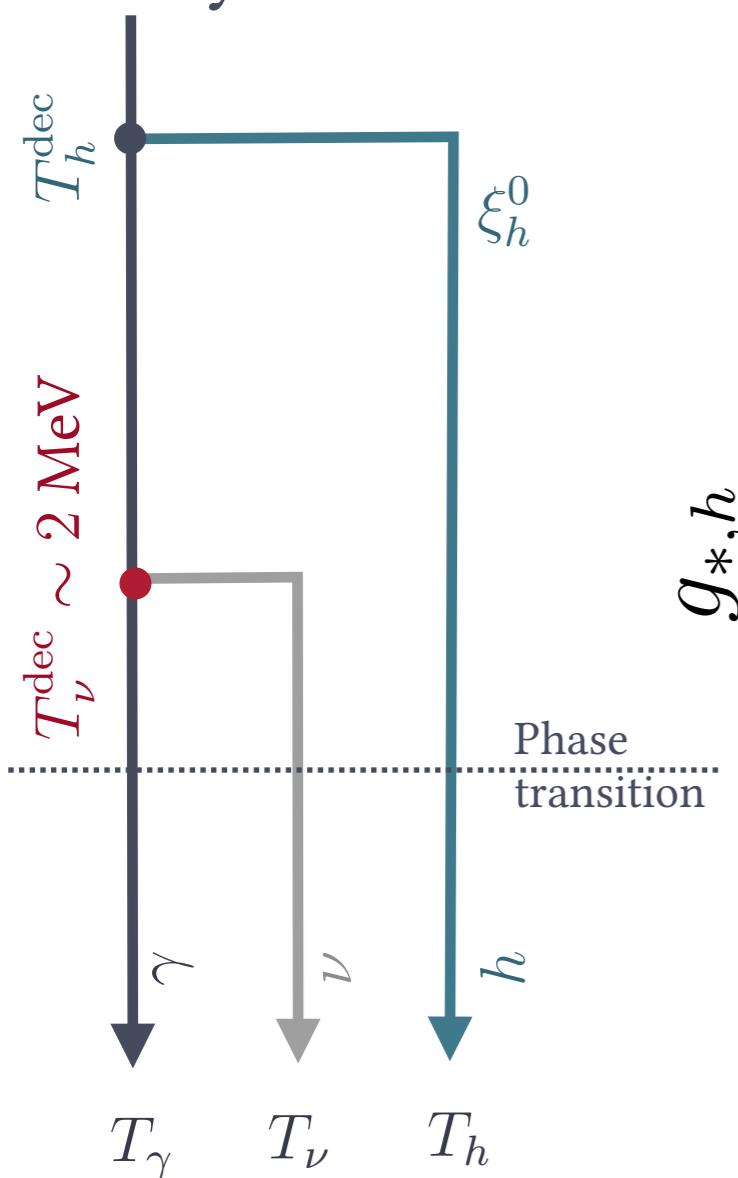


Cosmological Constraints

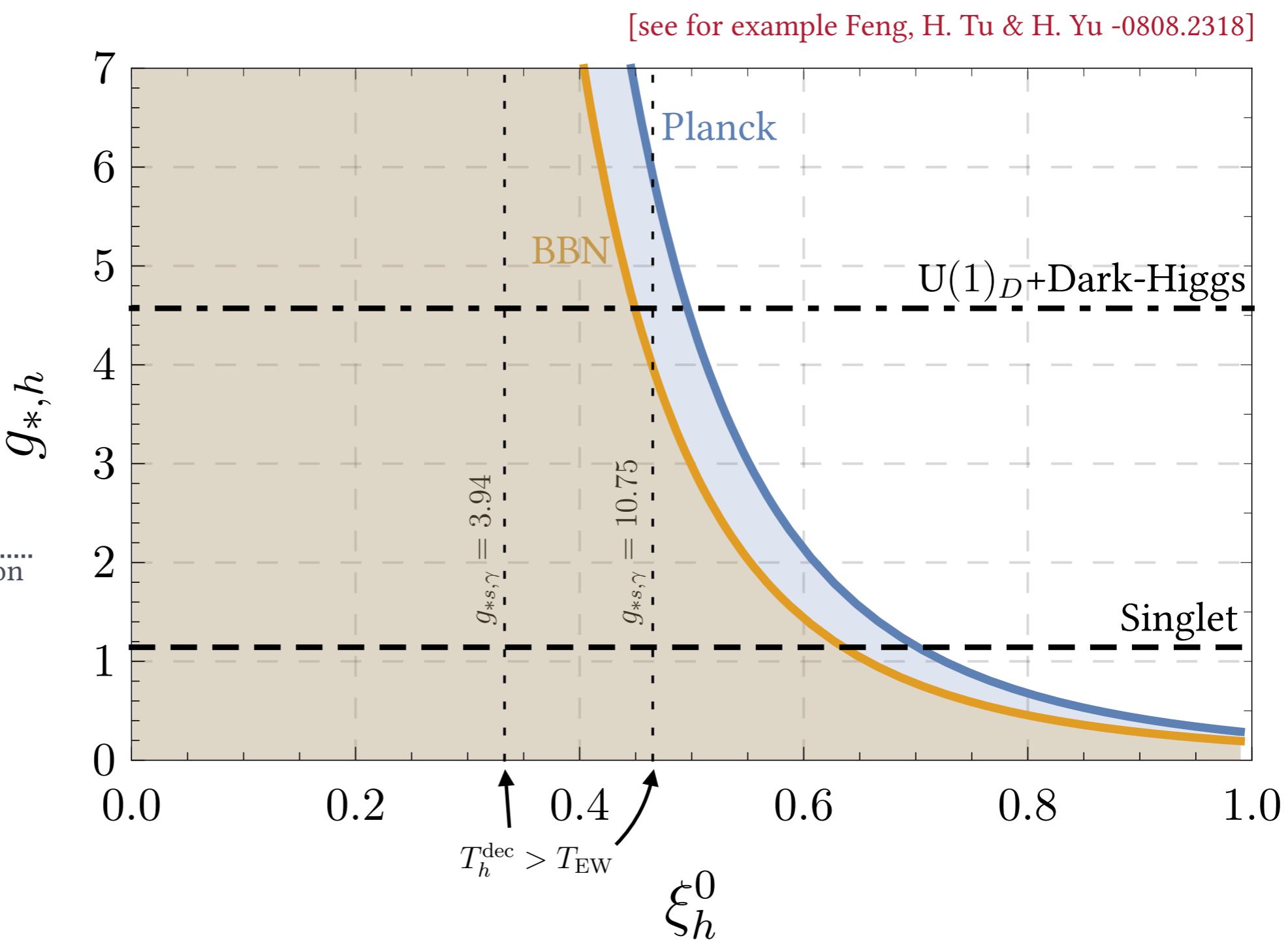
N_{eff} constraints

Decoupled hidden sector:

Early times



Late times

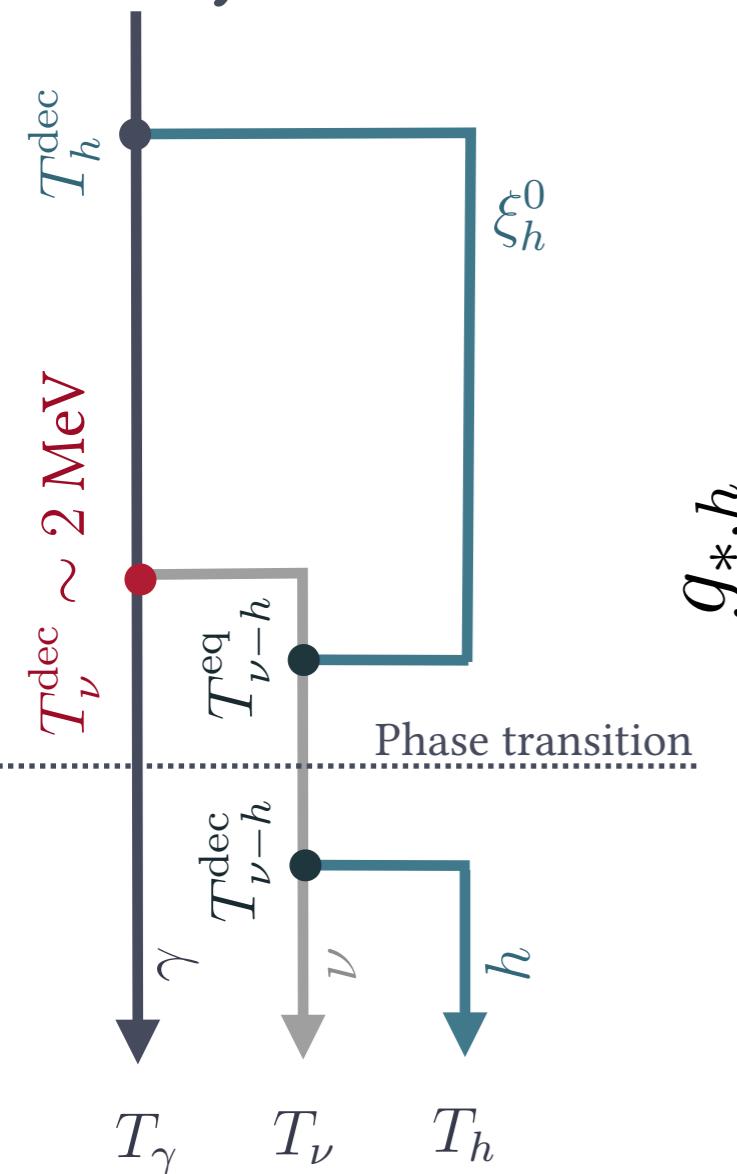


Cosmological Constraints

N_{eff} constraints

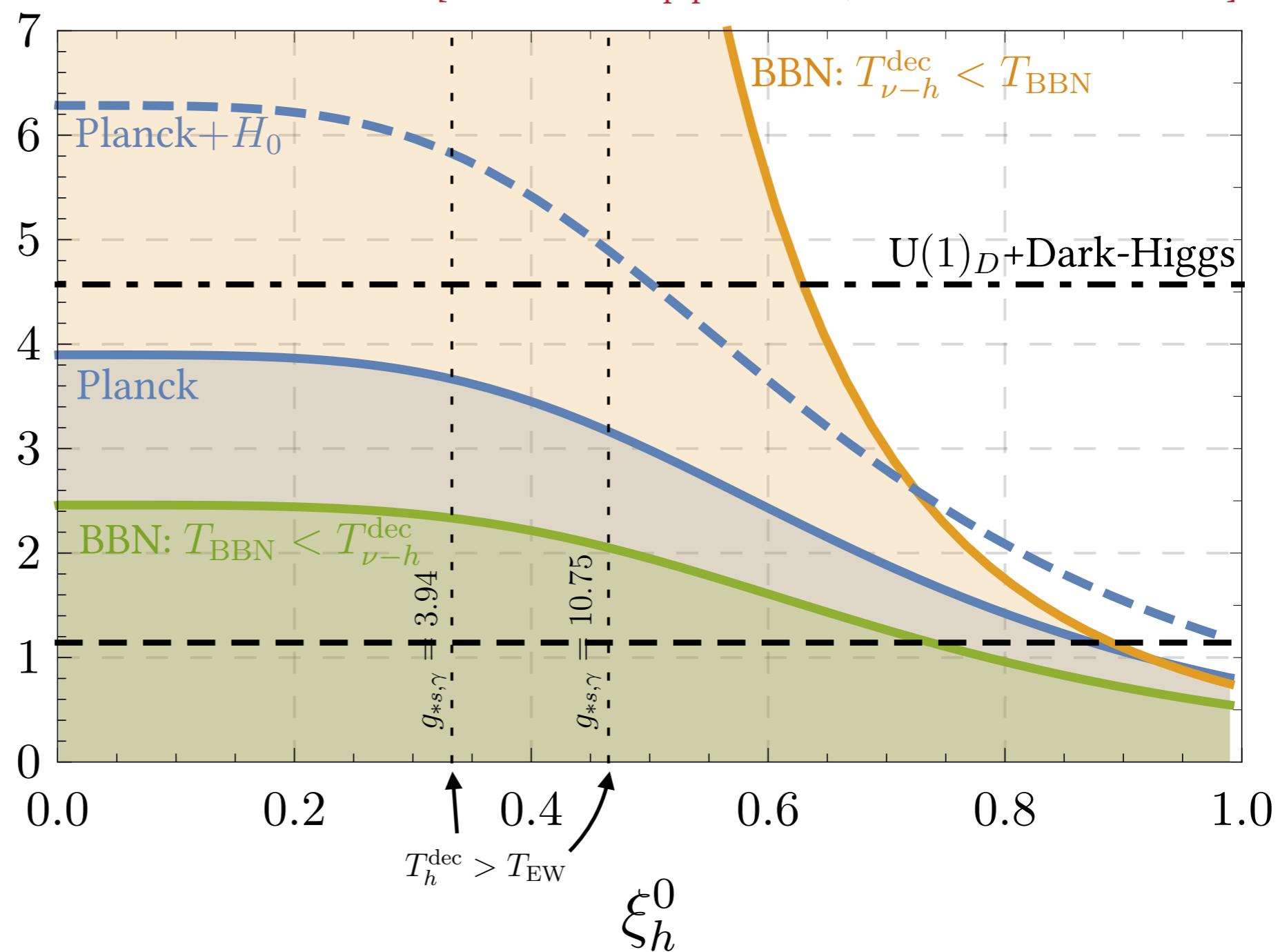
Re-equilibration with neutrinos:

Early times



Late times

[Chacko et al. hep-ph/0312267; Berlin & Blinov 1706.07046]

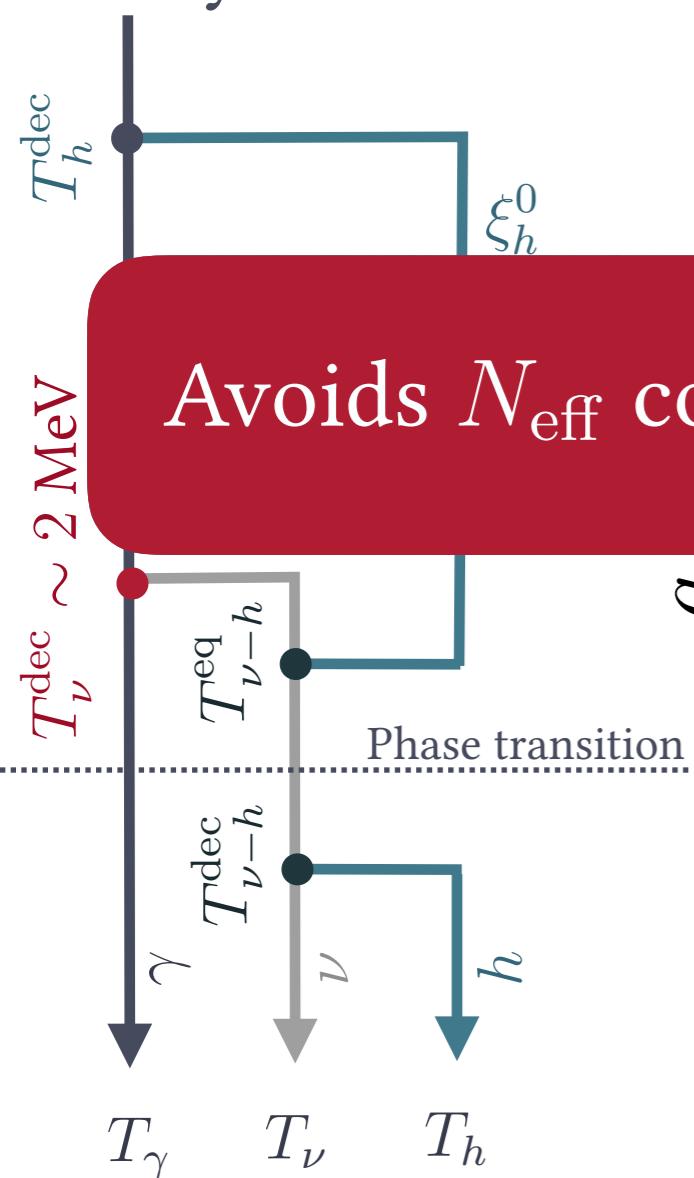


Cosmological Constraints

N_{eff} constraints

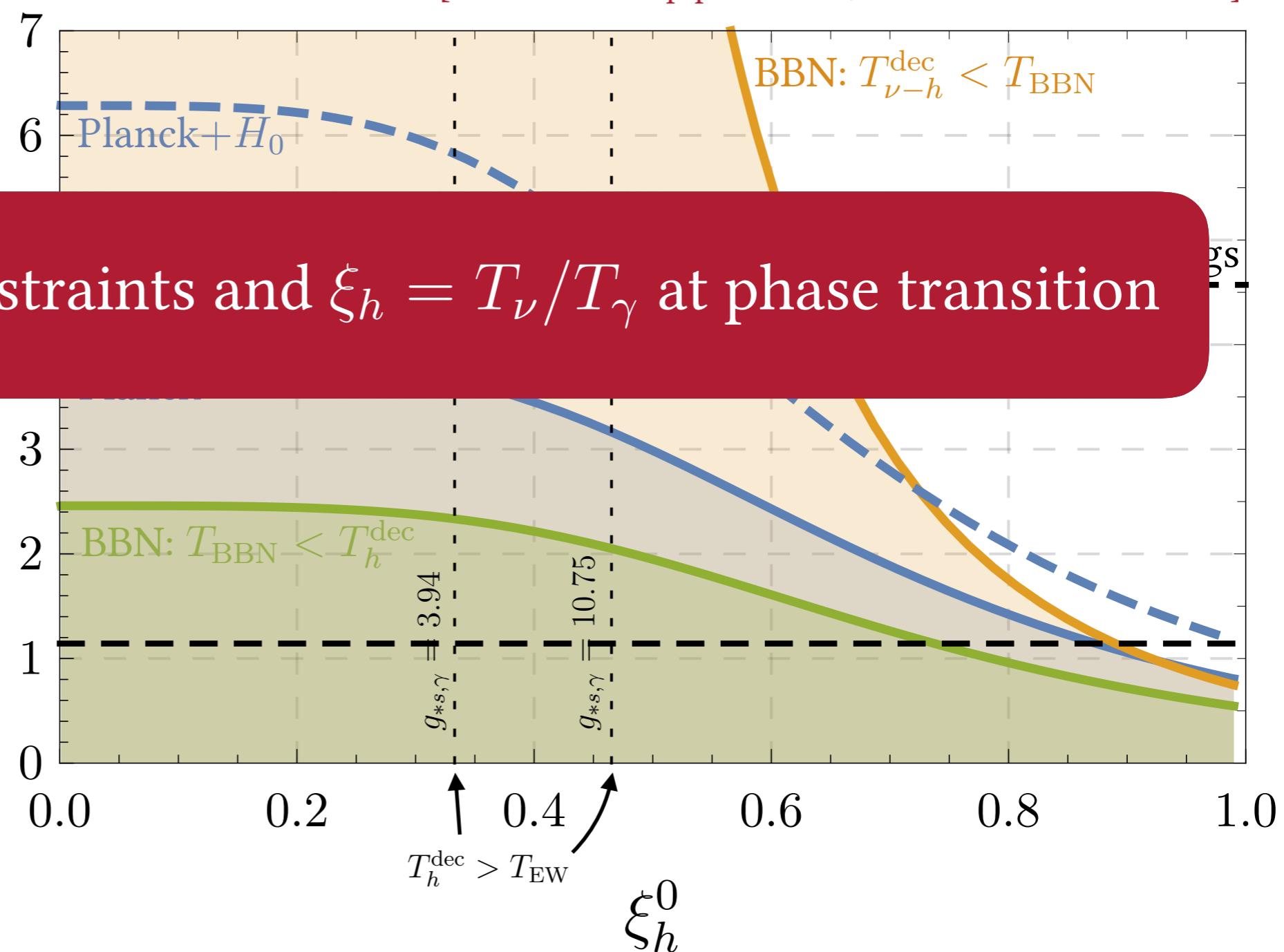
Re-equilibration with neutrinos:

Early times



Avoids N_{eff} constraints and $\xi_h = T_\nu/T_\gamma$ at phase transition

[Chacko et al. hep-ph/0312267; Berlin & Blinov 1706.07046]

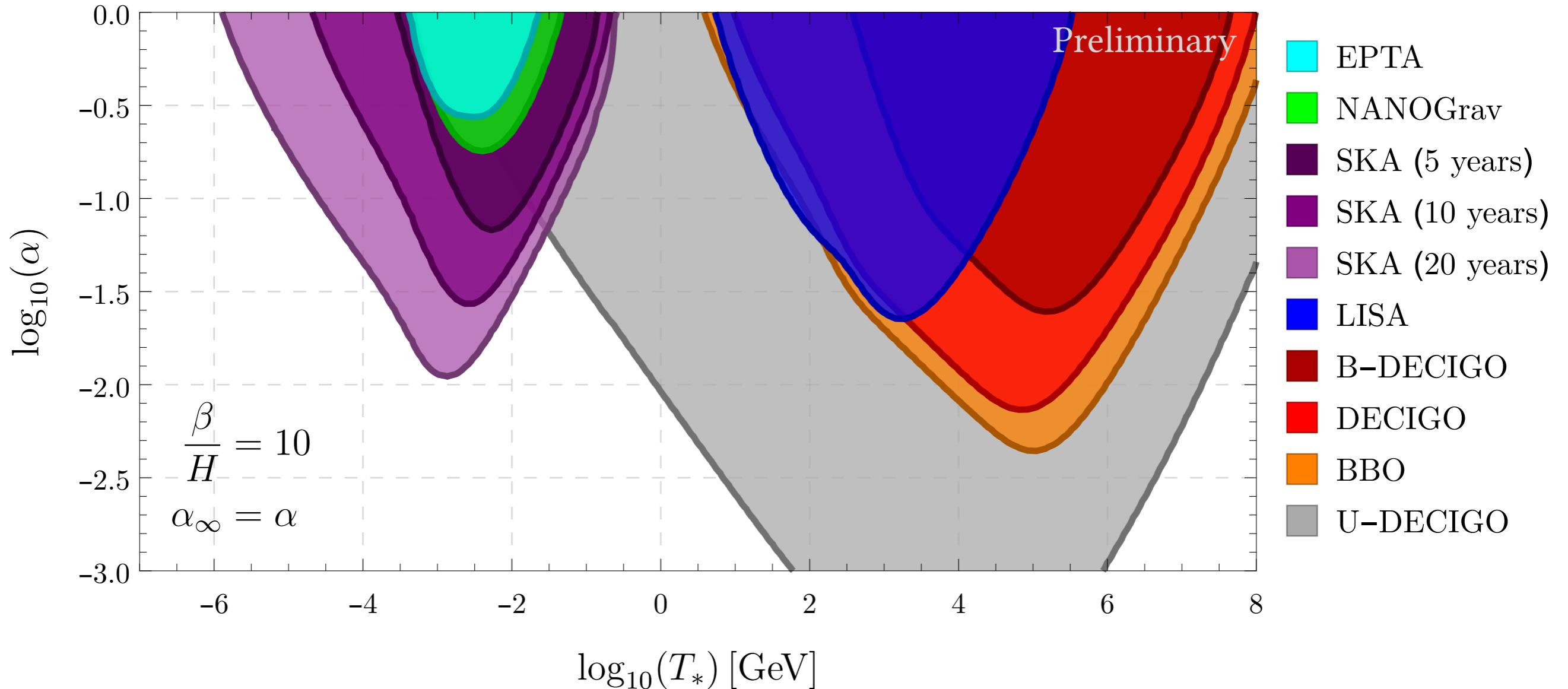


Late times

Experimental Sensitivity to gravitational wave spectra

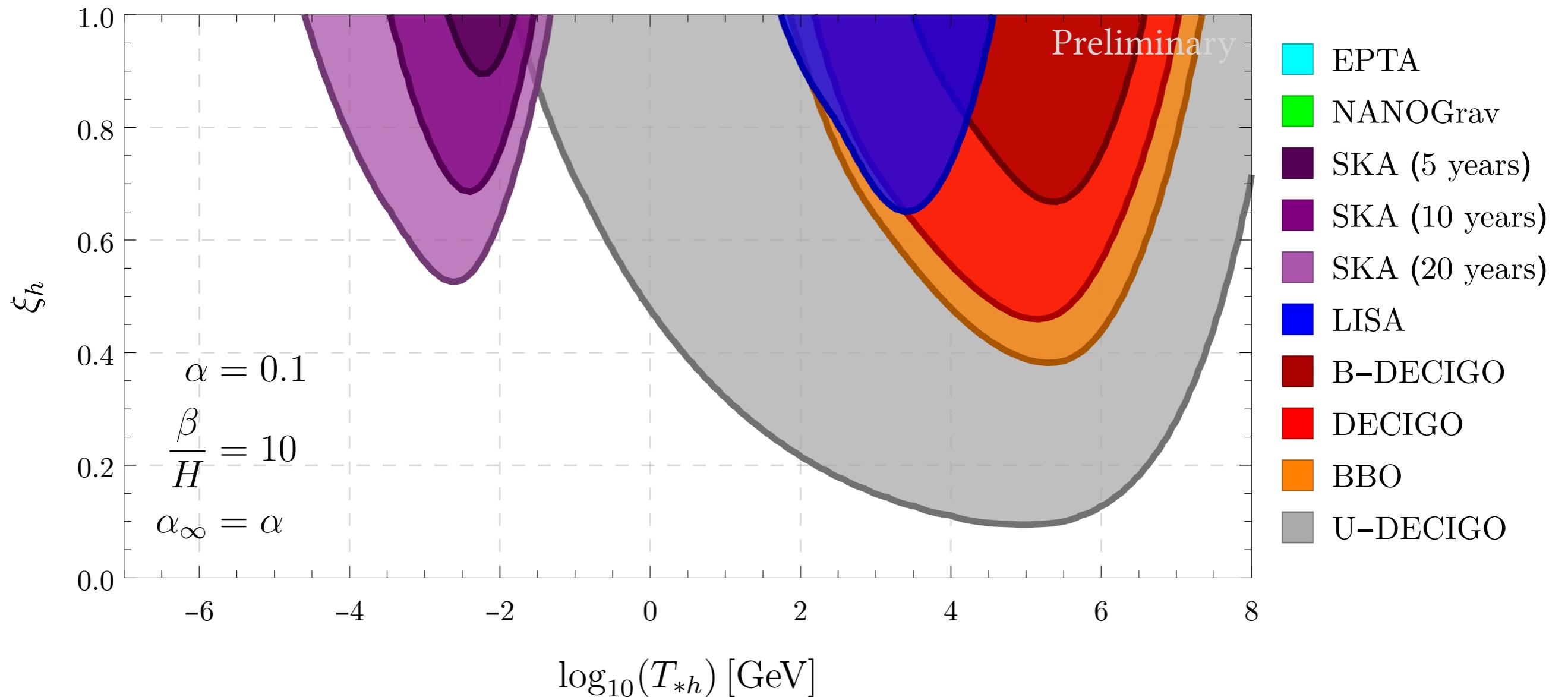
Detectability

Non-runaway bubbles



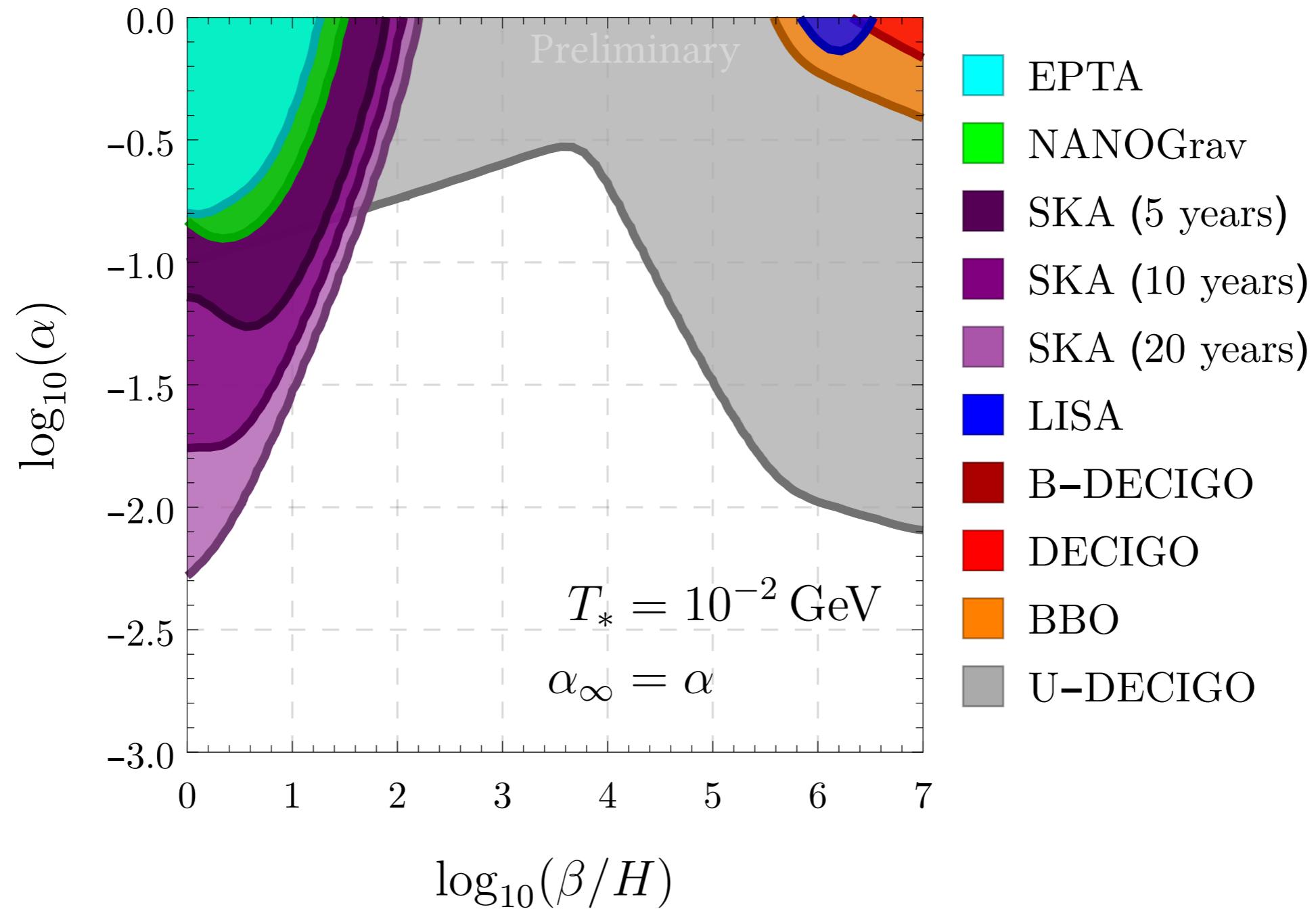
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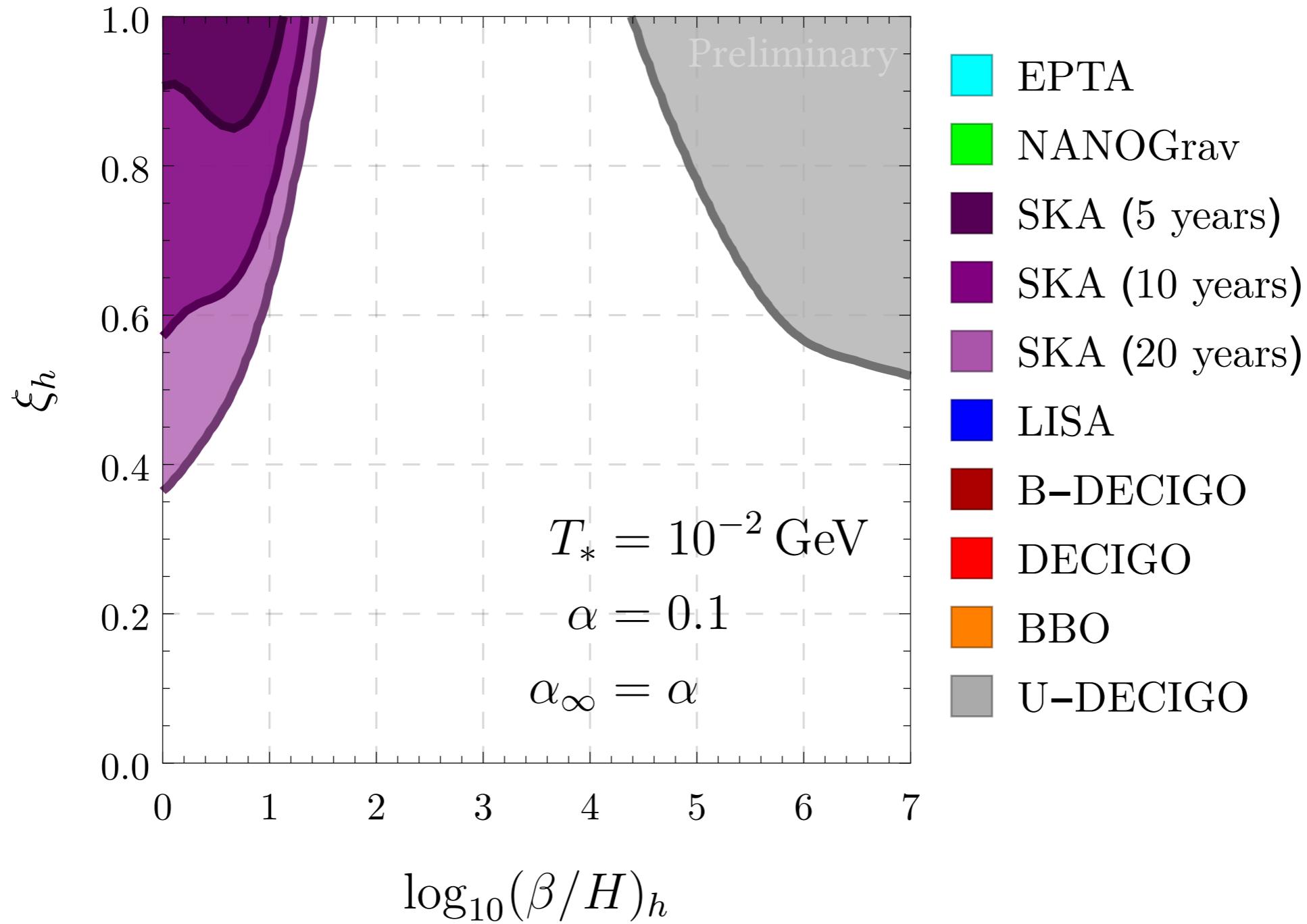
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Detectability

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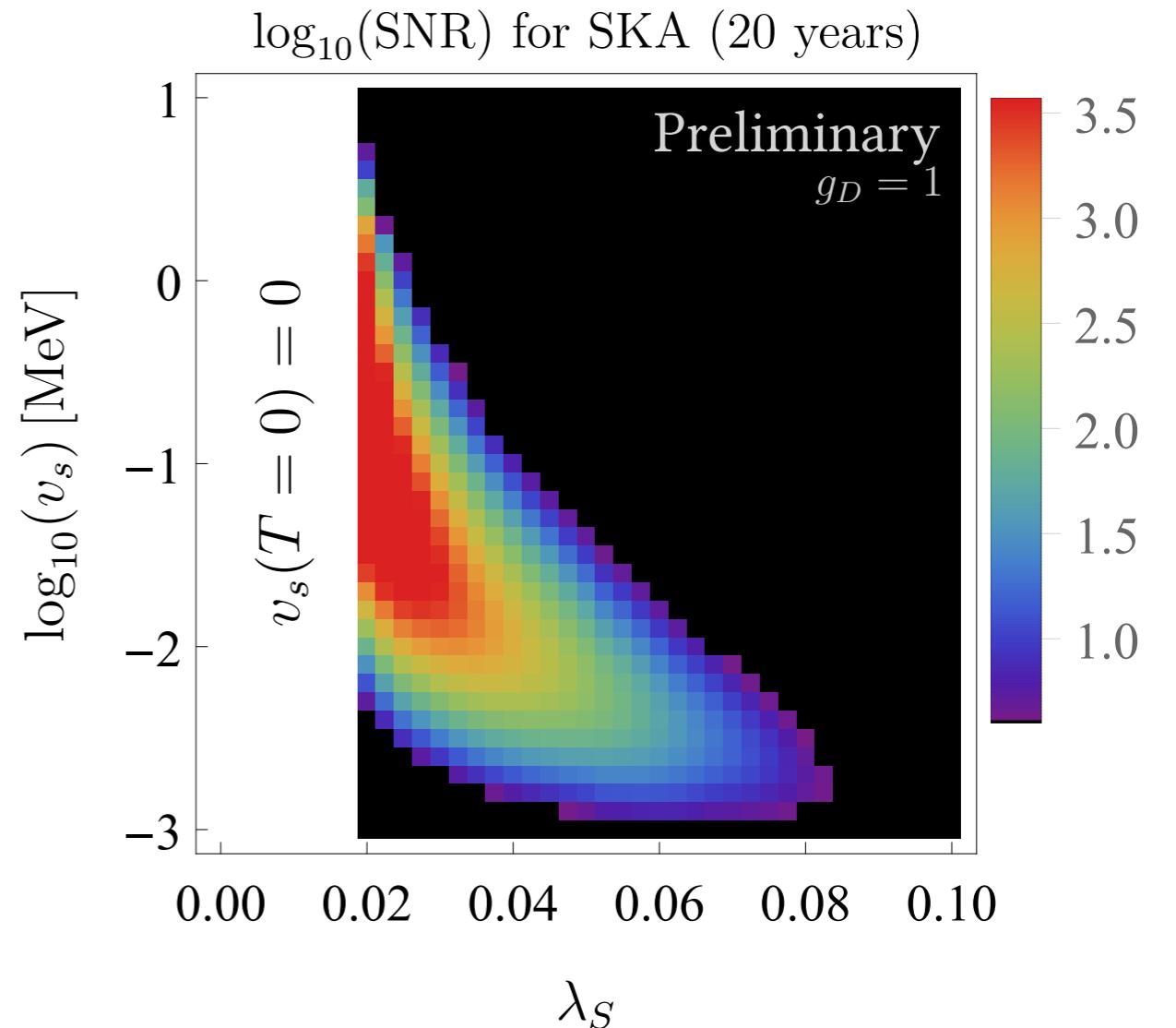
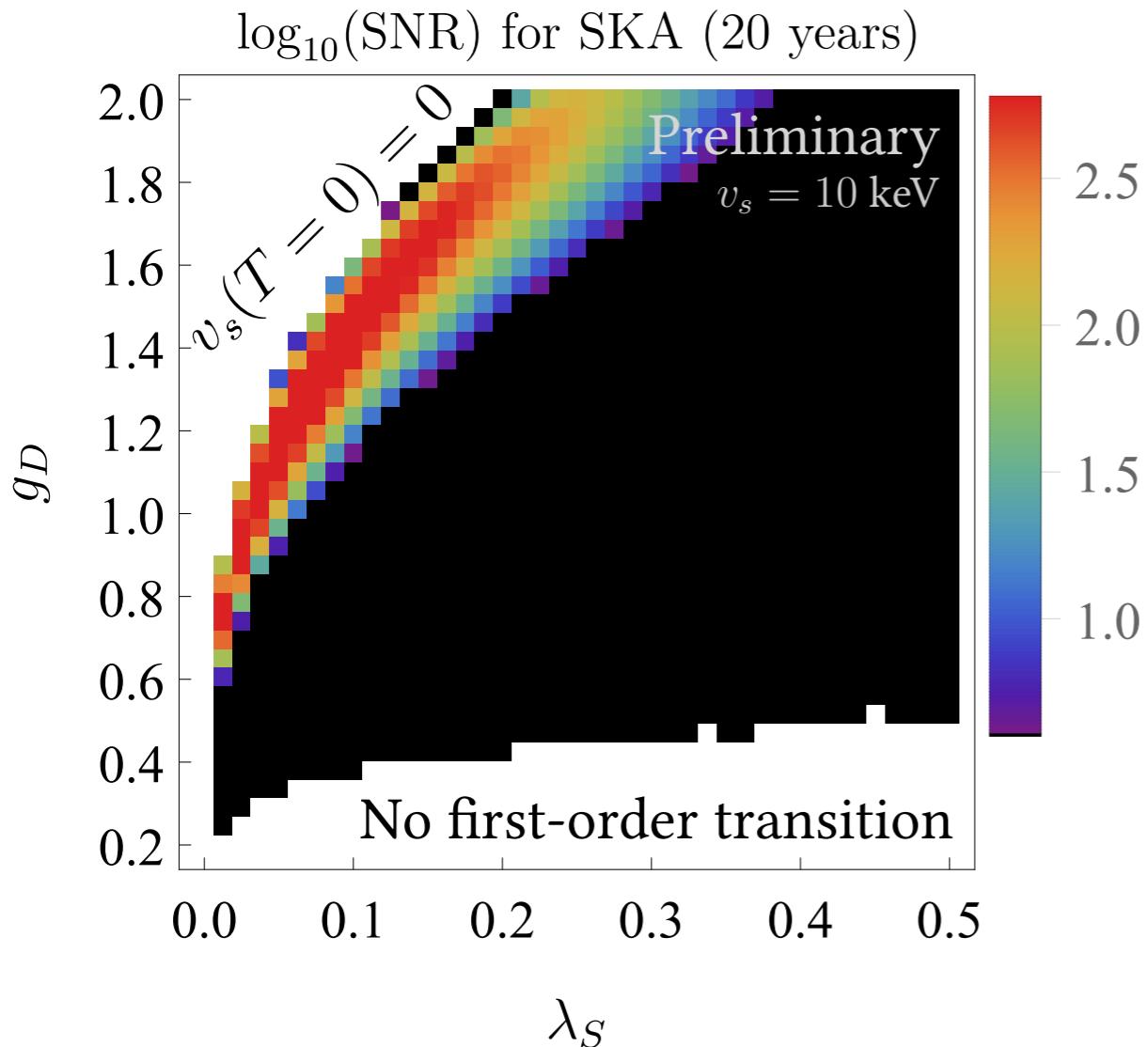


Higgsed Dark-Photon Model

Dark-Photon

First-order phase transition induced by $U(1)_D$ gauge boson

[for space-based experimental sensitivity see Addazi & Marcianò - 1703.03248, Hashino et. al - 1802.02947]

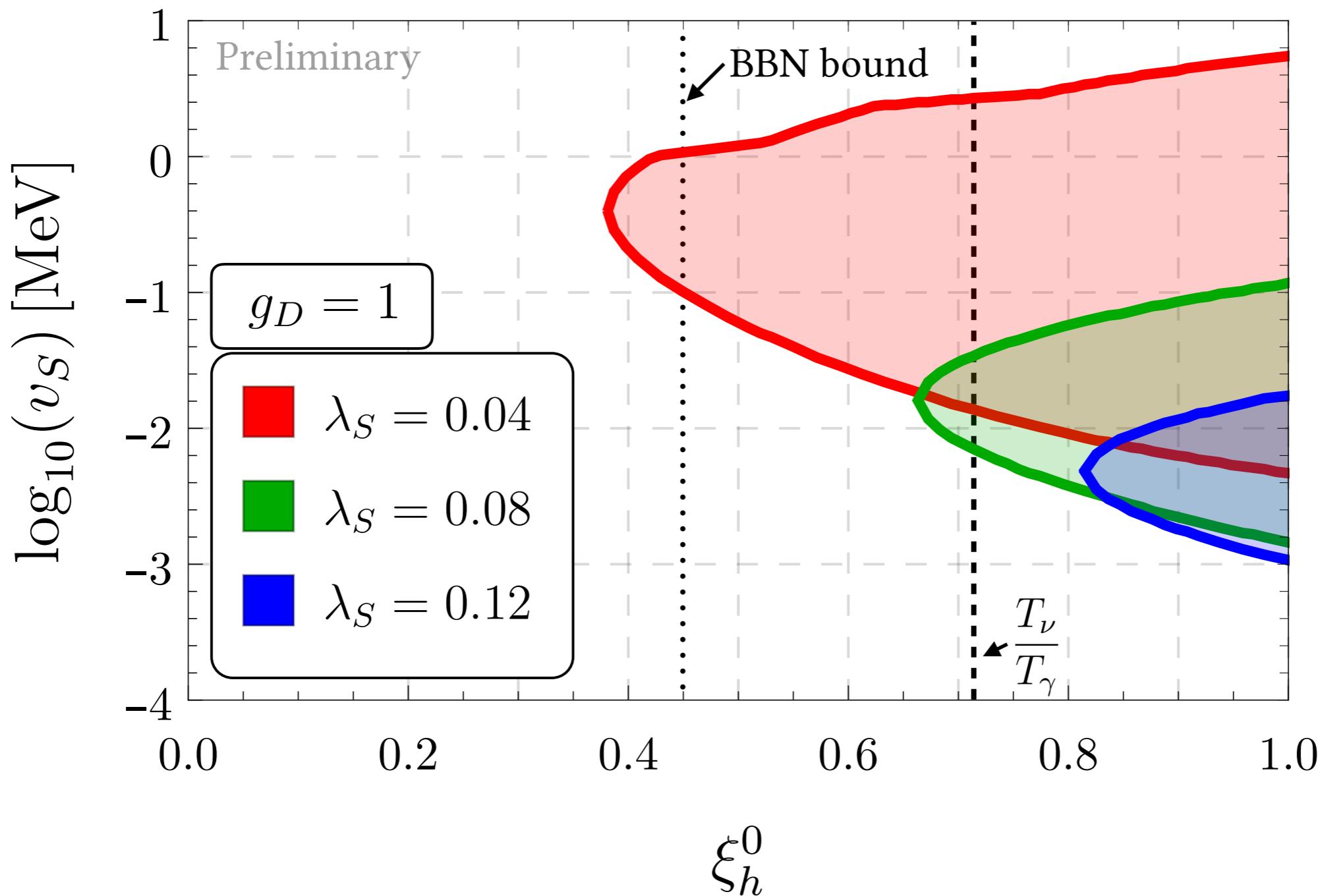


Sensitivity to scenarios where: $m_{A'} > m_{h_D}$

Dark-Photon

Sensitivity

SKA (20 years)



Outlook

Work in progress

This talk

Temperature ratio and thermal history play important roles determining both cosmology constraints and GW signal

Sensitivity to models that are strongly constrained from cosmology

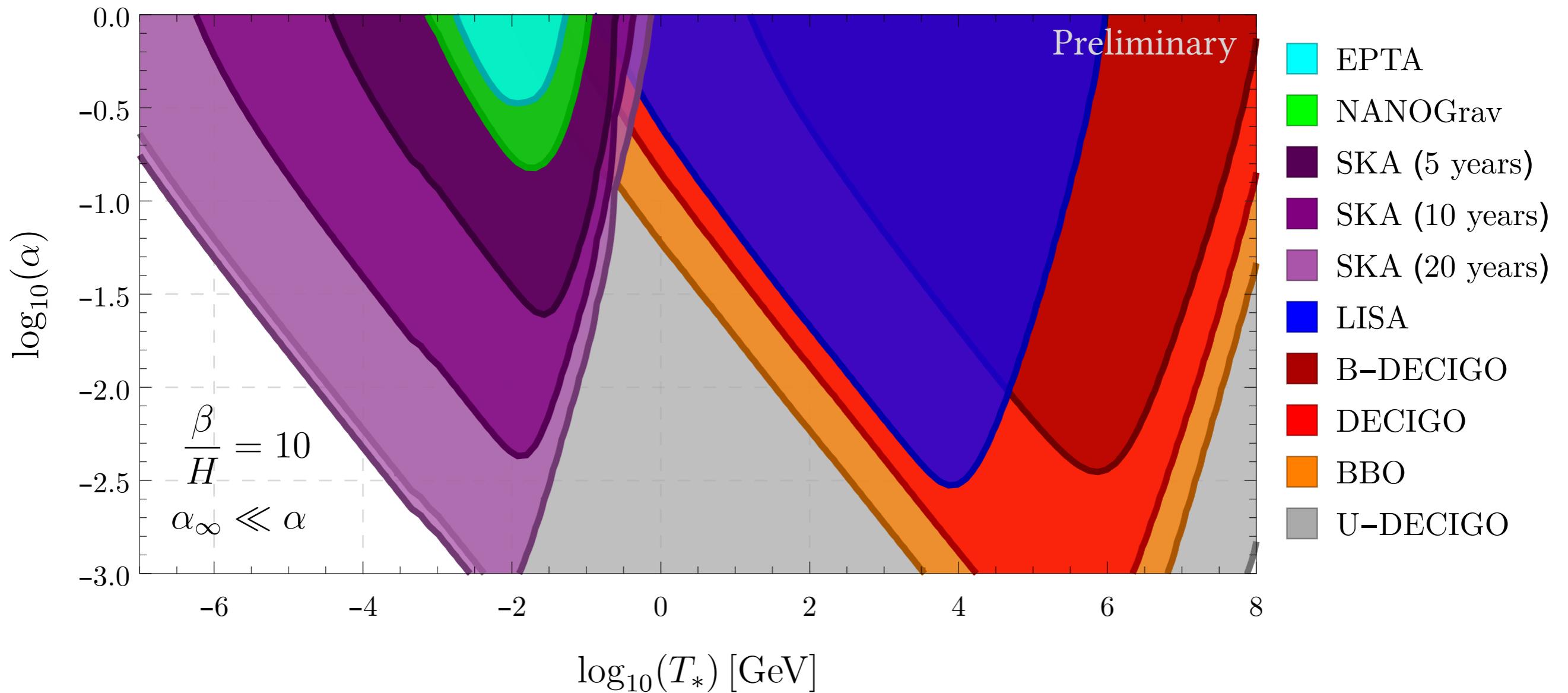
Additional models with fewer DOFs

Automation of GW spectrum calculation through extensions of existing tools

Additional Material

Detectability

Runaway bubbles



Detectability

Runaway bubbles

