

QCD - experimental part

Jets, Minimum Bias and Underlying Event

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Quinto workshop sulla fisica pp ad LHC



Start with Jets, link with Giulia

Jet reconstruction algorithms at work:

What the two experiments use and how they perform

Jet energy measurement:

brief review of the correction strategies

In situ validation

Analysis with jets:

Inclusive jet cross section

Jet Shapes

Event Shapes

Early measurements: Minimum Bias and Underlying Event

Basic tools

triggers, tracking, particle ID...

Minimum Bias measurement plans

Underlying Event measurement and sensitivity to QCD models

Monte Carlo tuning at LHC

Reconstruction Algorithms



CMS and ATLAS have implemented lots of different algorithm

Cone family (iterative, midpoint, SIS)	ATLAS	CMS
DR	0.4,0.7	0.5,0.7
Seed (if present)	1 GeV	1 GeV
Split and Merging parameter	50%	75%
Input	Towers, Clusters	Towers
Fast K_T	ATLAS	CMS
D parameter	0.4,0.6	0.4,0.6
Input	Towers, Clusters	Towers

The recombination scheme is simply the 4-vector sum of the constituents (E-scheme)

$$E_{jet} = \sum_{const} E_{const}$$
$$P_{jet}^i = \sum_{const} P_{const}^i$$

A careful study of the jet constituent drove their choice

Input to the reconstruction

What is the input for the jet reconstruction algorithm?

The choice is made based on granularity needs, efficiencies, sensitivity to noise, CPU requirements

The simple choice: calorimeter pseudoprojective towers:

The calorimeter granularity drives the size of the tower for ATLAS ($\Delta\eta \times \Delta\Phi = 0.1 \times 0.1$) and CMS (0.087×0.087 in the central region, increasing with the pseudorapidity)

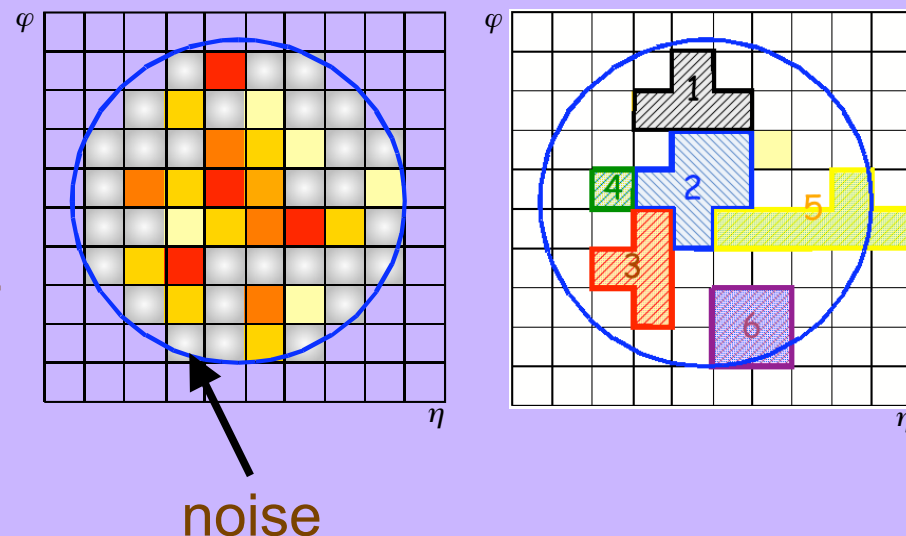
Topological cluster building implemented and widely used in ATLAS

Additional input foreseen for CMS: tracks, Pflow objects (any 4-vector....)

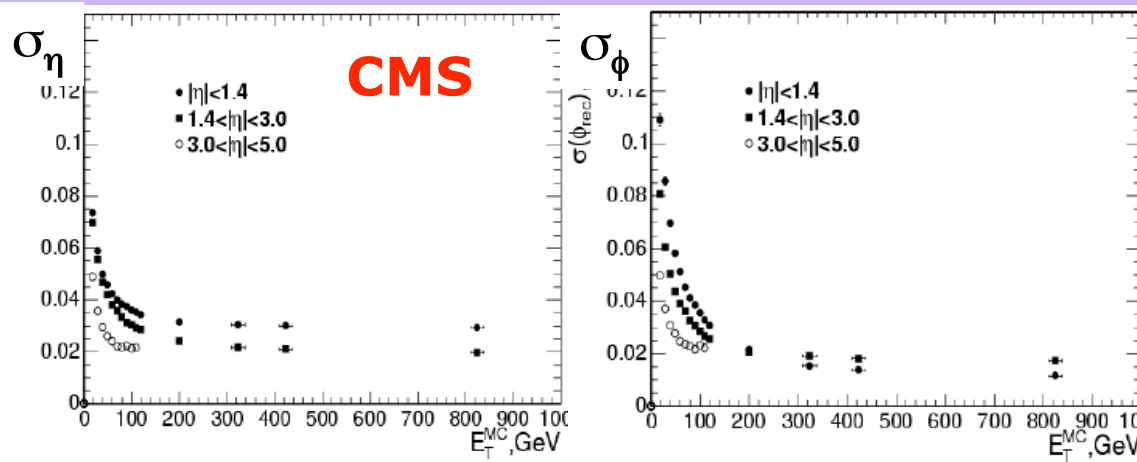
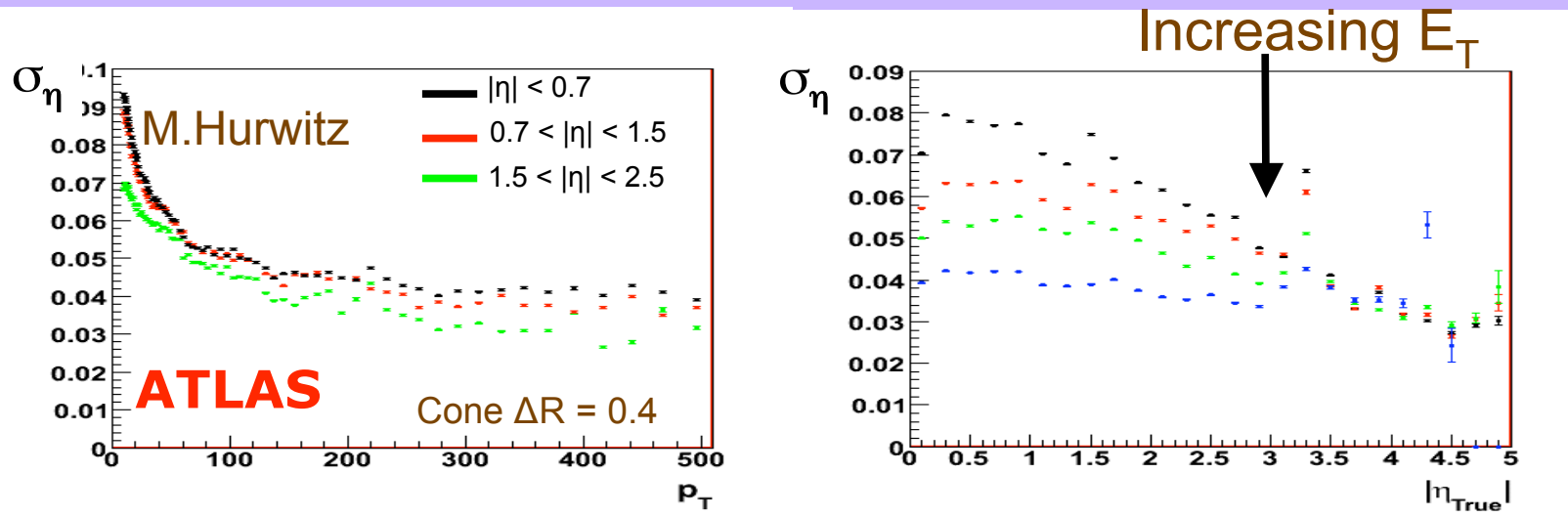
Noise suppression:

implemented at tower level for CMS

ATLAS: not implemented for towers. Topological clusters "based" on noise suppression



Reconstruction performances - angular resolution



Eg:
 $E_T = 100$, $|\eta| < 1$.
 $\sigma \sim [0.04-0.05]$

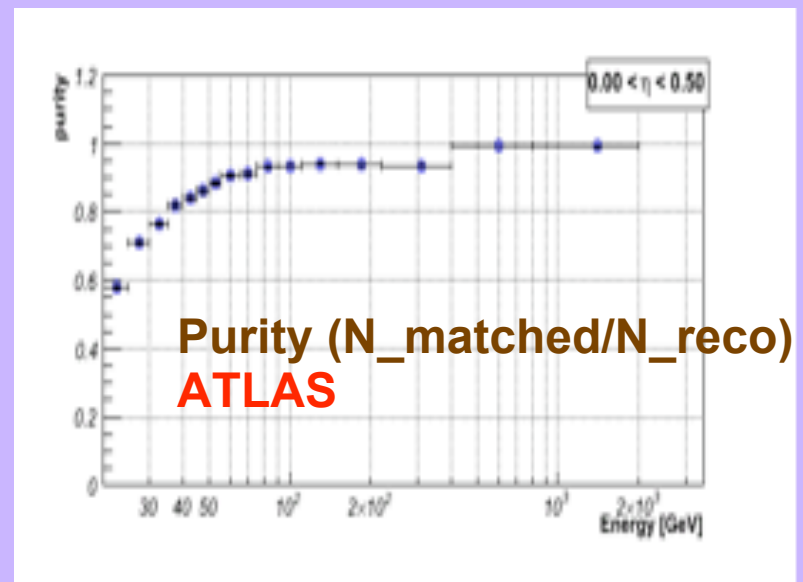
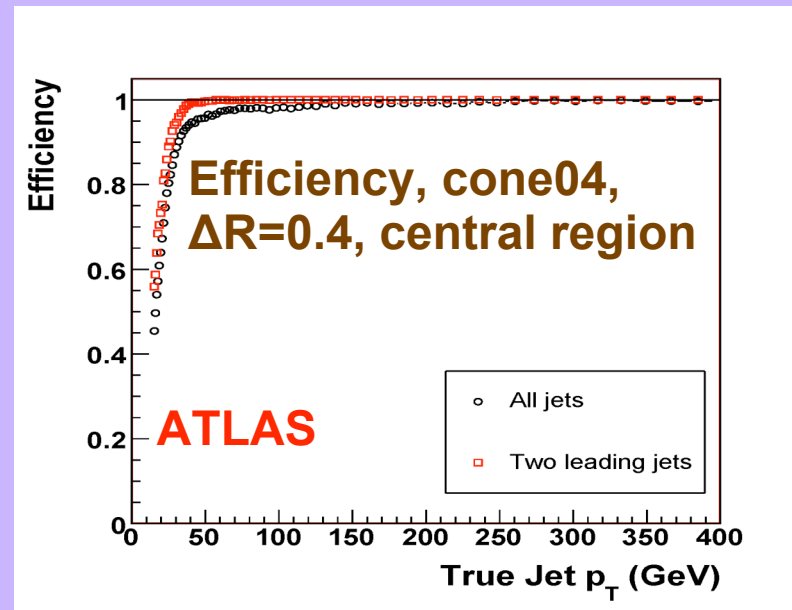
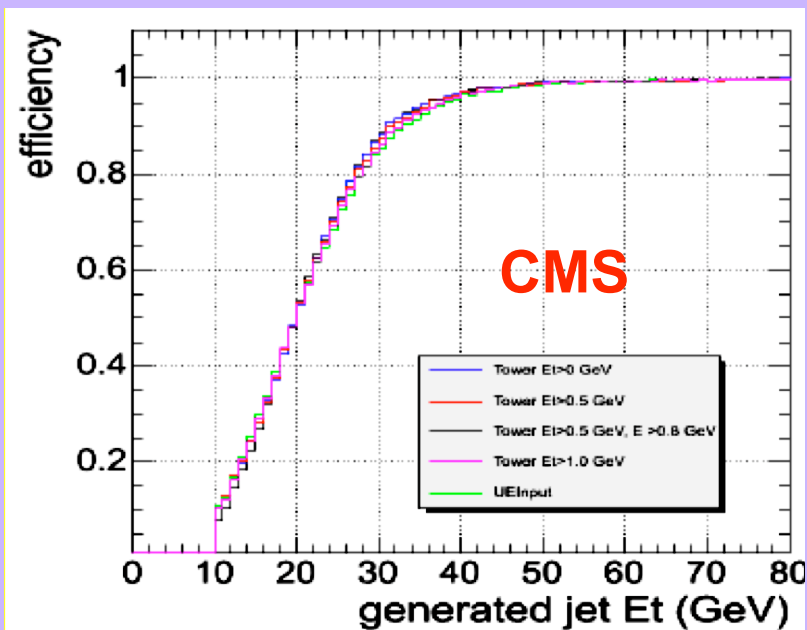
The angular resolution depends on the calorimeter granularity
 It is energy dependent

Reconstruction performances - efficiency

Reconstruction efficiencies:

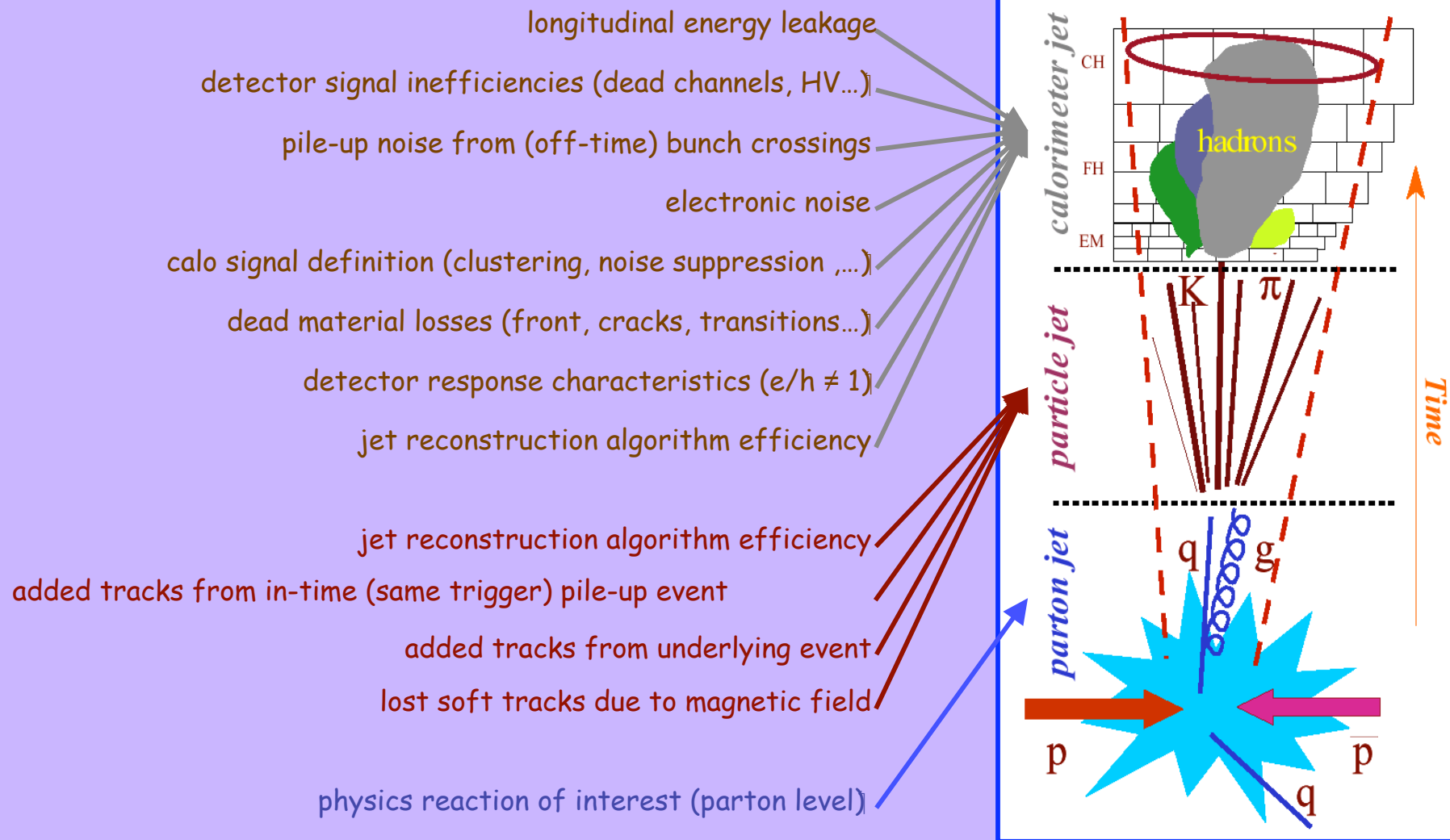
Important ingredient for, e.g.,
the inclusive cross section
measurement

The dependence on the jet
reconstruction algorithm is
found to be small



Reconstruction performances - energy measurement

The jet energy measurement is a complex task:
what enters and where



Reconstruction performances - energy measurement



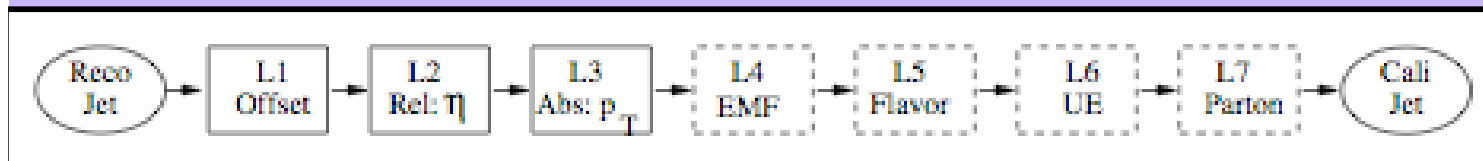
Both ATLAS and CMS are now trying to factorize the problem as much as possible

CMS is discussing an approach a la D0:

first subtract pileup, then make the detector uniform, compute the absolute response using γ +jet, correct for hadronization...

What comes out from this is a parton level jet:

$$E_{jet}^{ptel} = \frac{E_{jet}^{meas} - O}{R_{jet} S_{jet}}$$



ATLAS applies cell level corrections for detector effects only:

After the cell level corrections one gets the particle level jet

Underlying event and hadronization corrections applied (if necessary) to go to the parton level jet

Inclusive Jet Cross Section

Physics motivations:

interesting per-se

("old" and "new" physics)

Background for all the physics channels

Uncertainties from theoretical prediction

Experimental issues (100 pb^{-1}):

Understanding of the **X axis**:

Setting the energy scale and understand its systematic (γ + jets up to $\sim 300 \text{ GeV}$)

Measure the resolution from the data (di-jet balance)

Underlying event (tuning with tracks)

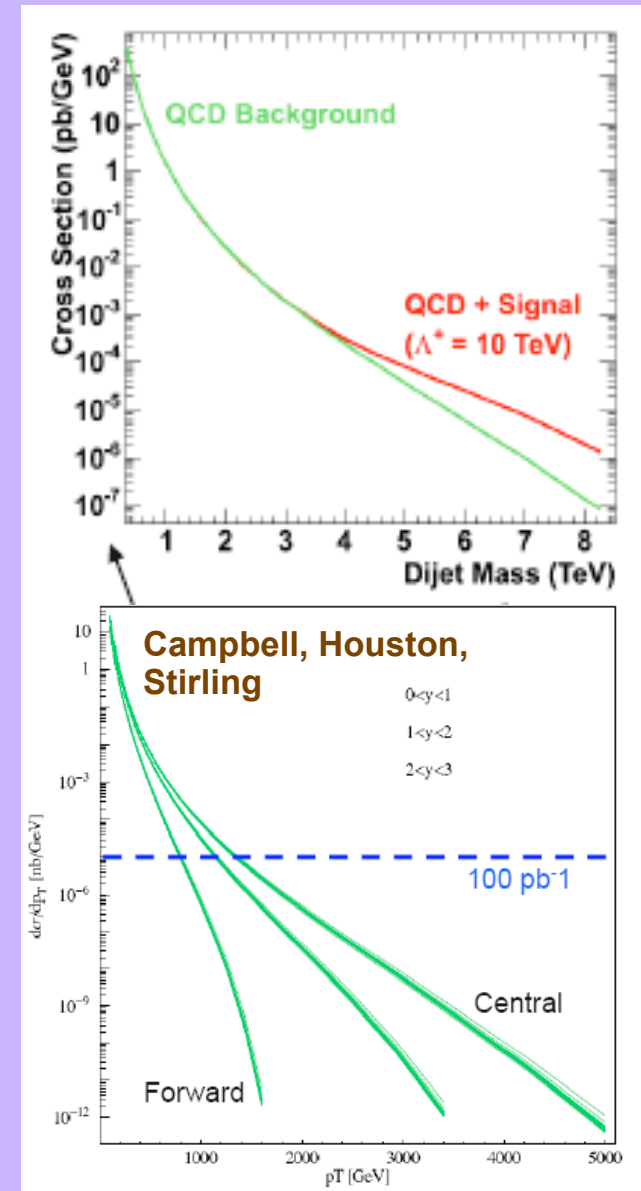
Understand high E_T jets (how do we check the scale at high E_T ?)

Understanding the **Y axis** (jet counting):

Luminosity

Jet reconstruction efficiencies

Jet trigger efficiencies (tag and probe proved to work OK)



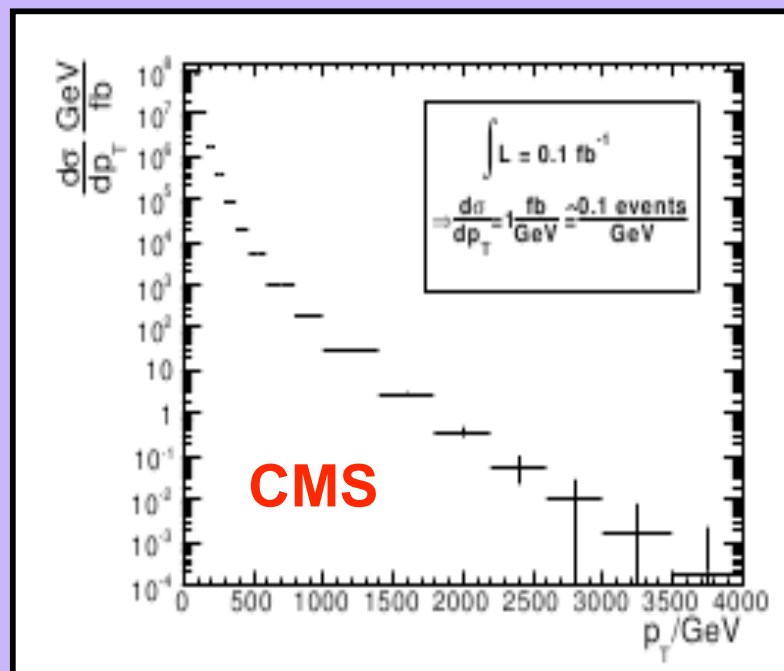
Inclusive Jet Cross Section

Not necessarily an easy measurement:

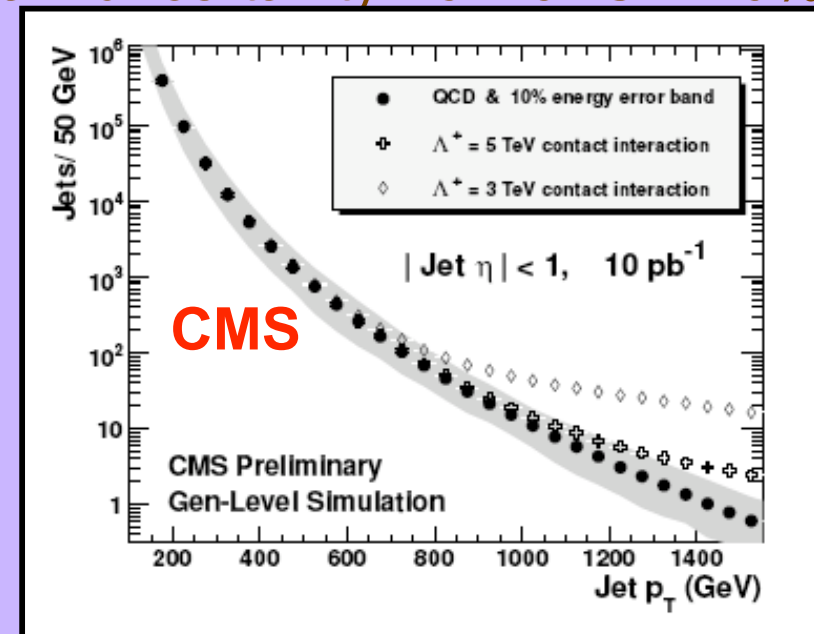
Early data taking characterized by the highest uncertainty on the jet energy scale

Up to what scale can we probe with the first data?

statistical error at 1 TeV is 1.3% with 100 pb^{-1} , 0.4% with 1 fb^{-1}



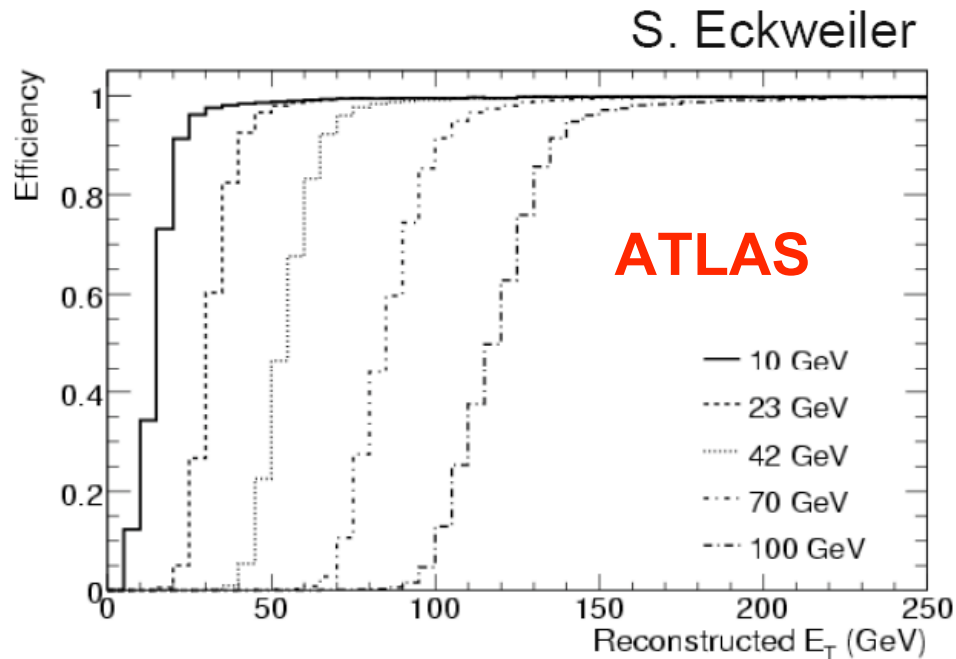
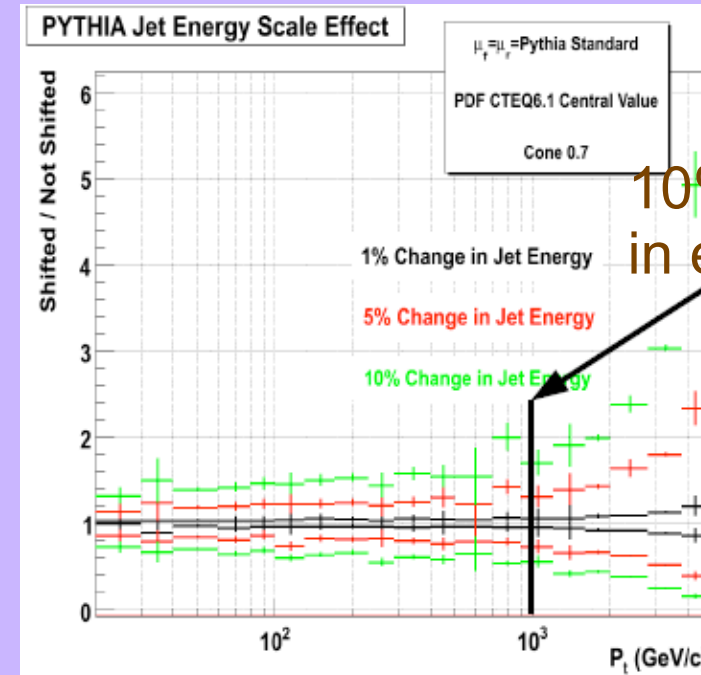
e.g, with 10 pb^{-1} we can probe contact interaction up to $\sim 3 \text{ TeV}$ (main uncertainty from JES $\sim 10\%$)



Inclusive Jet Cross Section

The uncertainty on the jet energy scale is the dominant experimental uncertainty:

10% error on the energy at 1 TeV means **50% error** on the cross section



The jet trigger efficiency can be measured using tag and probe techniques:

The highest E_T jet is biased
Look if the trigger is able to see the other jets

Jet Shape

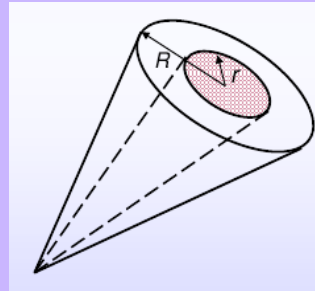
Jet shapes useful to distinguish
gluon and quark originating jets:

Already used at CDF
(hep-ex/0505013v2)

One possibility is to define

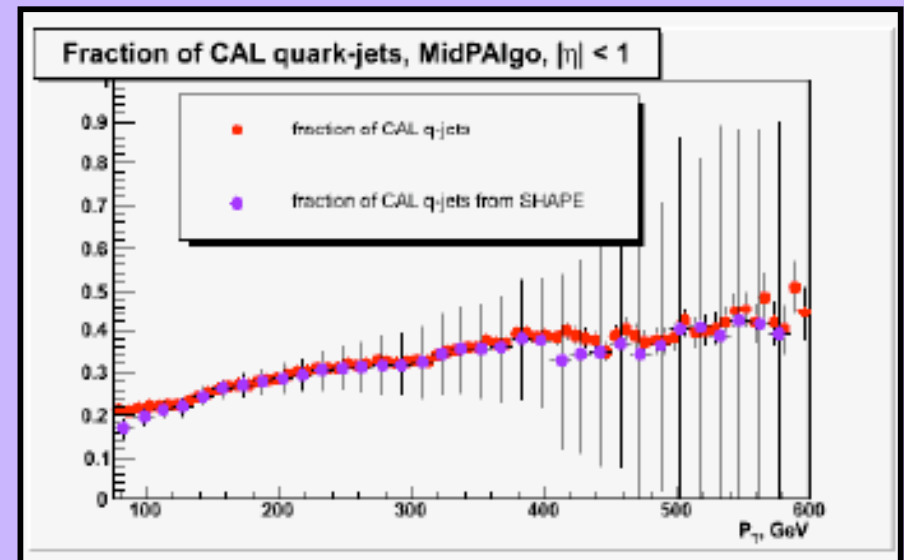
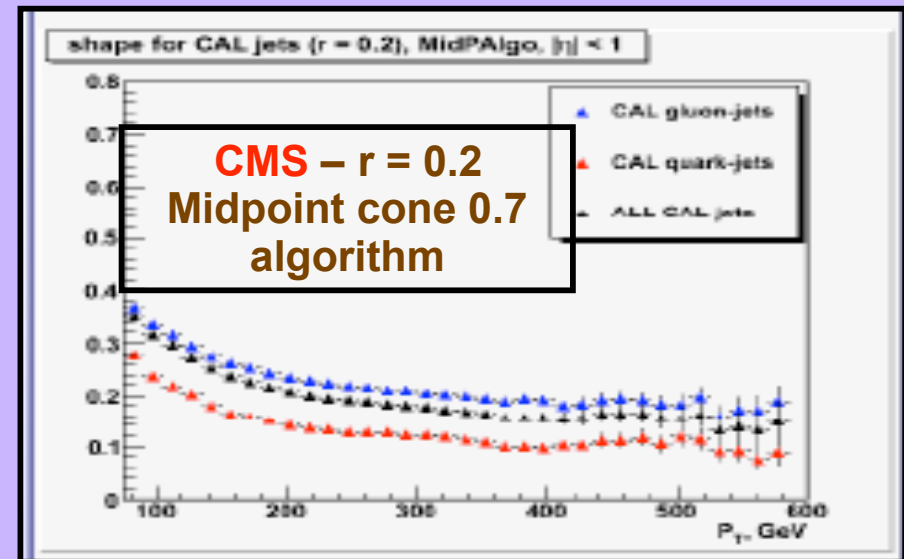
$$\psi = \frac{1}{N_{jets}} \sum_{jets} \frac{P_T(0, r)}{P_T(0, R)}$$

$$SH = 1 - \psi$$



$$F_q^{SH}(P_T) = 1 - \frac{SH_{all-jets}(P_T) - SH_{quark-jets}(P_T)}{SH_{gluon-jets}(P_T) - SH_{quark-jets}(P_T)}$$

Promising to estimate the relative
abundance of gluon and quark
jets



Angular decorrelation

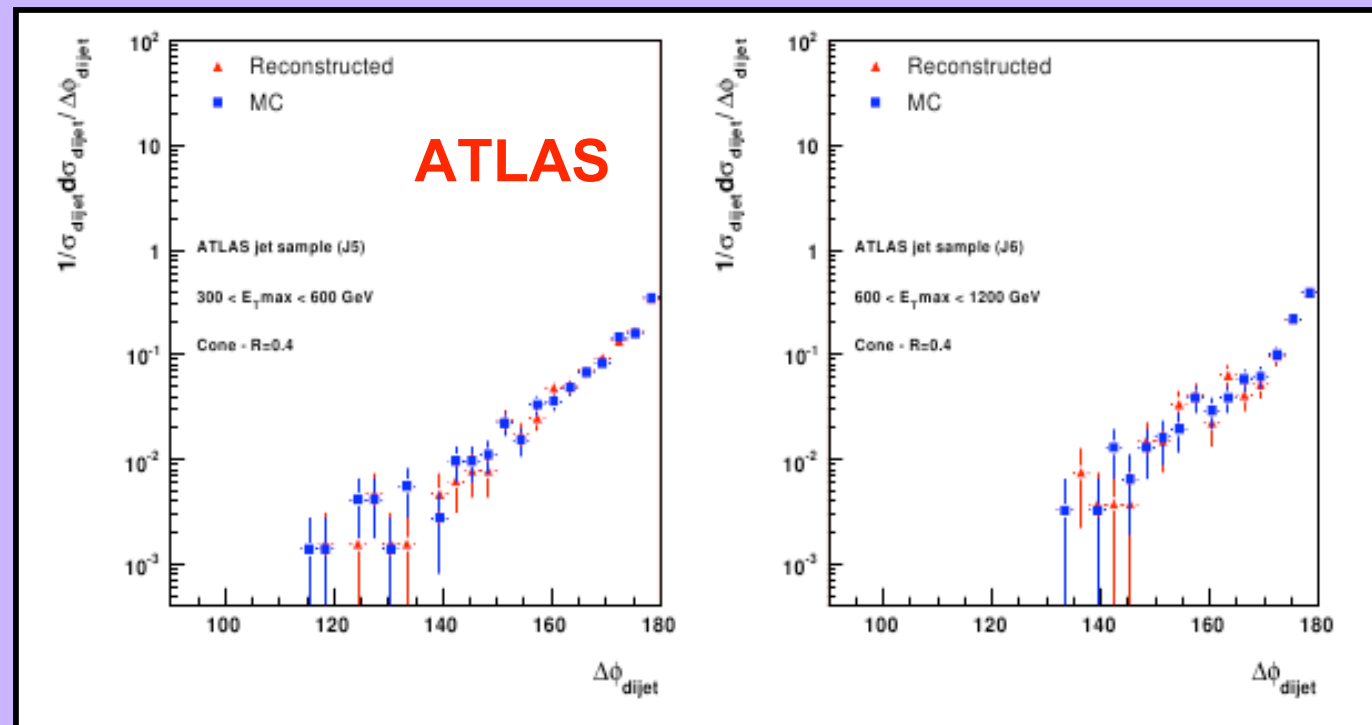
The angular distance ($\Delta\phi$) between the leading and the next to leading jet in QCD “di-jet” events is sensitive in particular to hard radiation :

Already measured by D0 at 1.96 TeV

Not sensitive to “reconstruction details”

It is an important tool to tune the event generators

The quality of the jet reconstruction will allow to repeat the measurement at LHC



Event shape



Something new: a study of the event shapes at CMS

Event shapes are robust against experimental issues

(see Giulia's talk)

Not too sensitive against jet scale and jet resolution issues

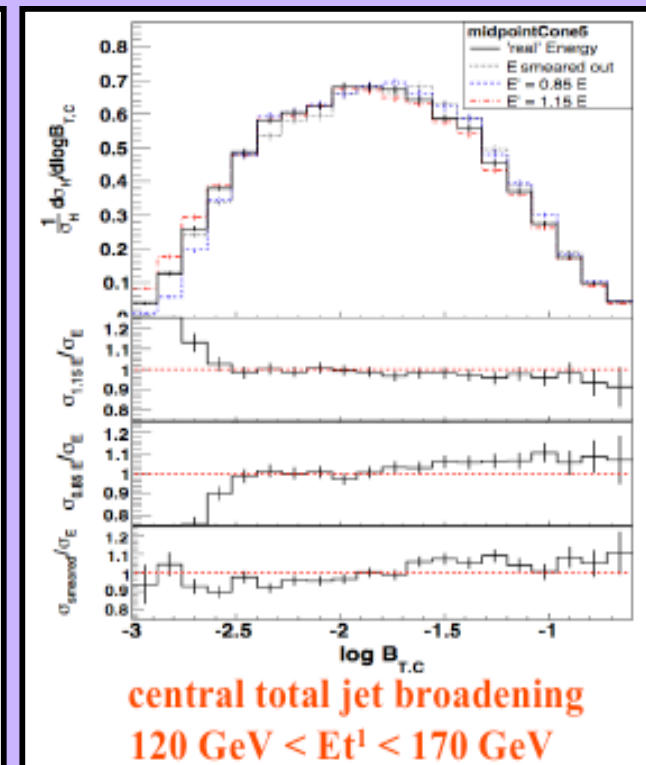
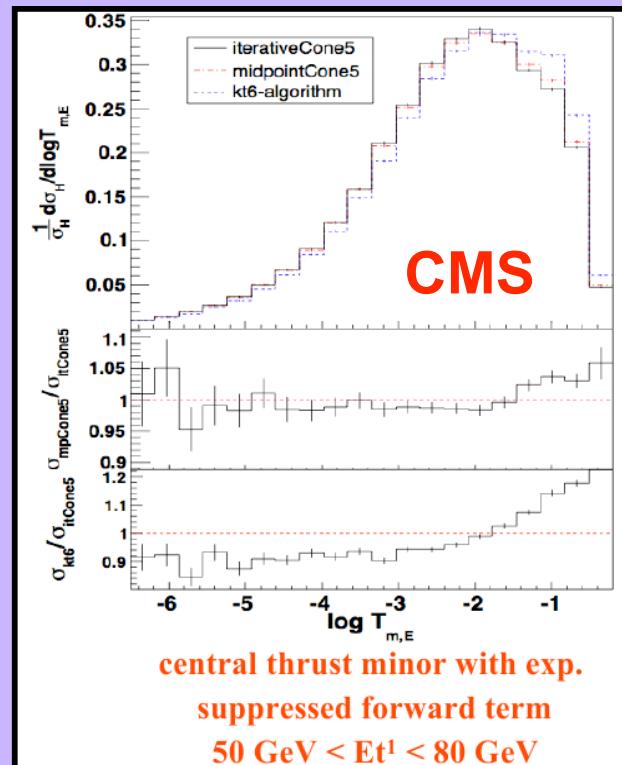
Not too sensitive to the jet reconstruction algorithm

Study on a “toy” MC:

change of $\pm 15\%$ in
the particle jet
scale

Smearing/ No
smearing

Small ($\sim 10\%$)
impact on the
measurement



Event shape

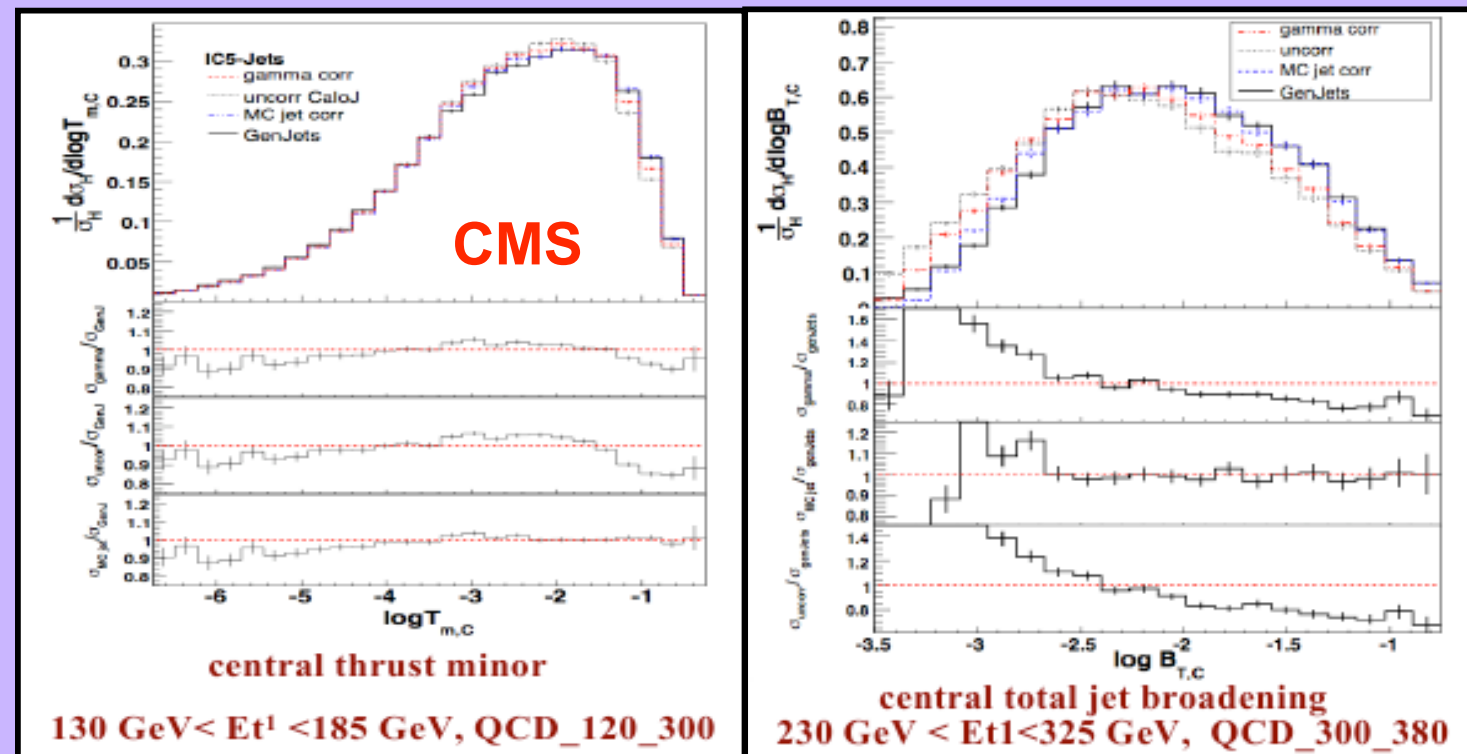
Study carried out also in full simulation:

In general, a good agreement between the particle and reconstructed jets

global thrust shapes well reproduced at reco level

Some of the central shapes have problems with the jet scale and resolution

Low dependency on the jet algorithm confirmed at reco level



QCD with Jet - summary



The most used jet reconstruction algorithms are implemented in ATLAS and CMS

Long and deep studies have been done or are ongoing on their performance

Scale corrections highly factorized both in ATLAS and CMS

Jet scale is an issue for the early data. Reasonable to assume 5%-10% of systematic uncertainty

Many analysis going on both in ATLAS and CMS. The inclusive jet cross section suffers from JES uncertainty

Strategies on how to asses the jet performance with data are in place

In general, the detector reconstruction looks satisfactory with simulated data

The LHC experiments are getting ready for the jet reconstruction and analysis.

Minimum Bias and Underlying Event

Motivations



Exploring Fundamental aspects of hadron-hadron collisions

Describe QCD@LHC in the best way

Not enough to rescale conclusions from Tevatron to 14 TeV

[different Q^2 , x range and energy dependence of the cutoffs]

Structure of hadrons

Factorization of interactions

spin offs on other relevant physics

Calibration of major physics tools

Low, medium and high-PT QCD affected by “surrounding” processes which affect:

Pile up understanding, jet energy, isolation performances,

vertexing, detector response, High-PT background...

Tuning of Monte Carlo Models

Both **not-perturbative** and **perturbative** aspects

Remnants, I-FSR radiation, **MPI...**

(UE activity, minijet, hard scattering)

Understanding the detector

occupancies

background

...

Introduction



A proton-proton collision is a combination of “soft” and occasionally “hard” process

$$\frac{d\sigma}{dt} \sim A + B\alpha_s + C\alpha_s^2 + D\alpha_s^3 + \dots \quad \dots \text{pQCD well describe hard process}$$

Modeling the collision:

+pQCD is applied to describe high-PT parton-parton scattering
 +attempt to extend high-PT treatment to low-PT region
 (using the cutoff P_{t0} to regularize the x-section)

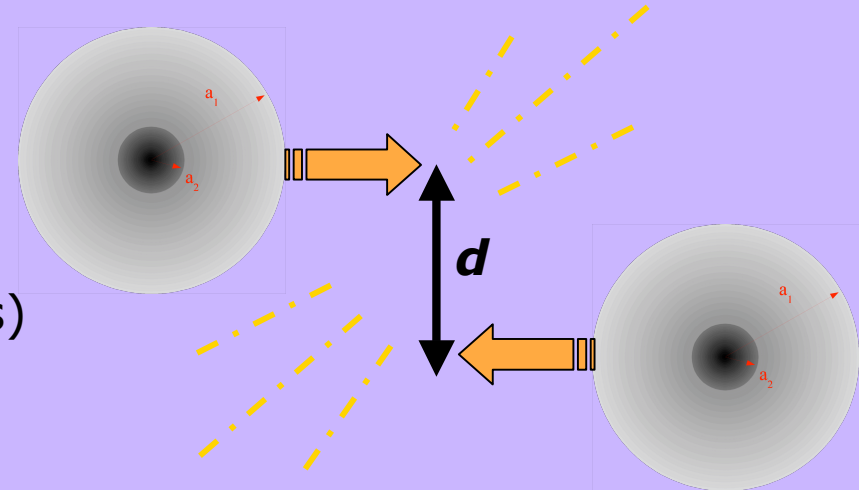
+MPI can occur

$$\sigma_{\text{int}} = \int_{P_{T\text{min}}}^{s/4} \frac{d\sigma}{dp_t^2} dp_t^2 \quad P_{T\text{min}}^{(s)} = P_{T\text{min}}^{(s')} \left(\frac{s}{s'} \right)^\epsilon$$

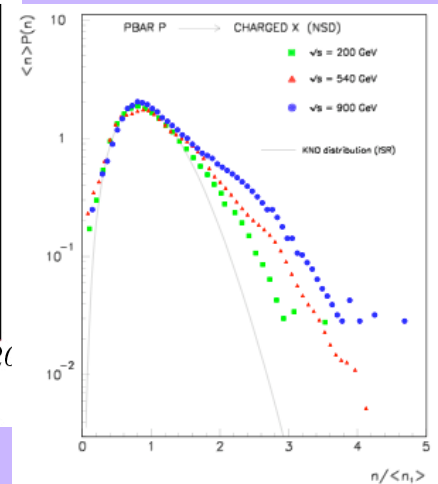
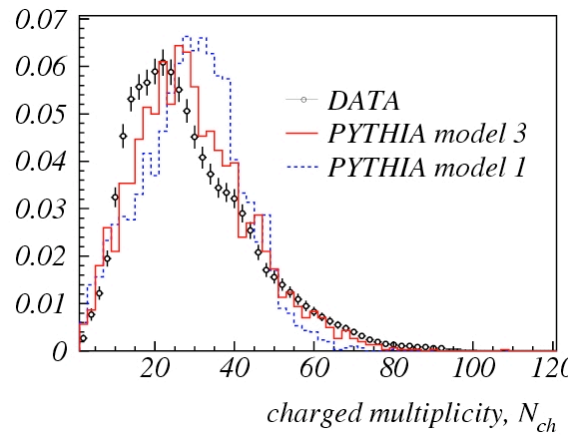
$$n \sim \sigma_{\text{int}}$$

Lower $P_{T\text{min}}$ -> higher n (# interactions)

Lower d -> higher probability to have an hard-scattering



Past observation and MPI



ISR, FSR, SPECTATORS are not enough to account from the observed multiplicities, P_T spectra, KNO scaling violation (CDF, AFS, UA1...)

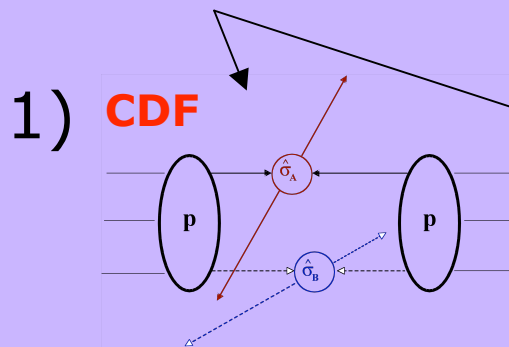
The Pythia solution:

[T. Sjöstrand et al. PRD 36 (1987) 2019]

Multiple Parton Interactions (MPI)

Main Parameter: P_T cut-off P_{T0}

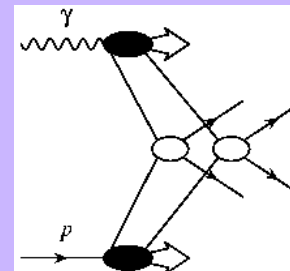
Inspired by observations of double high P_T scatterings



$$p\bar{p} \rightarrow \gamma / \pi^0 + 3 jets + X$$

Phys. Rev. Lett. **79** (4) 584 (1997);
Phys. Rev. D **56** (7) 3811 (1997).

2)



Photoproduction at HERA

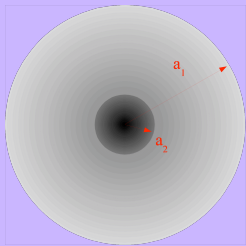
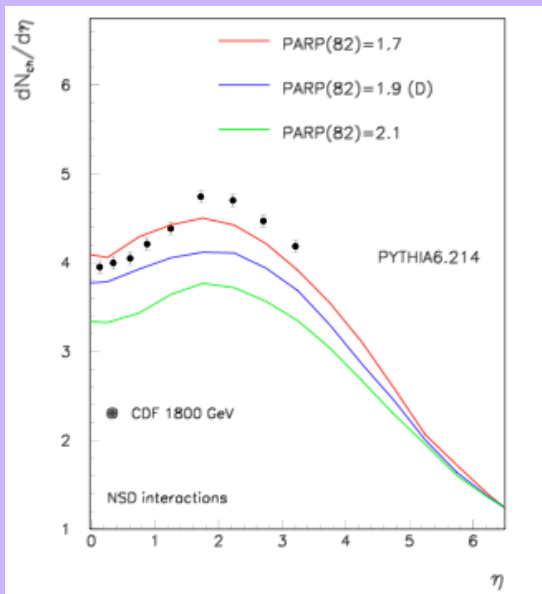
Z. Phys. C**72** (1996) 637.

$$\sigma(\widehat{P}_T) \rightarrow \sigma(\widehat{P}_T) \cdot \frac{(\widehat{P}_T)^4}{((\widehat{P}_{T0})^2 + (\widehat{P}_T)^2)^2}$$

Modeling the MPI (pythia)

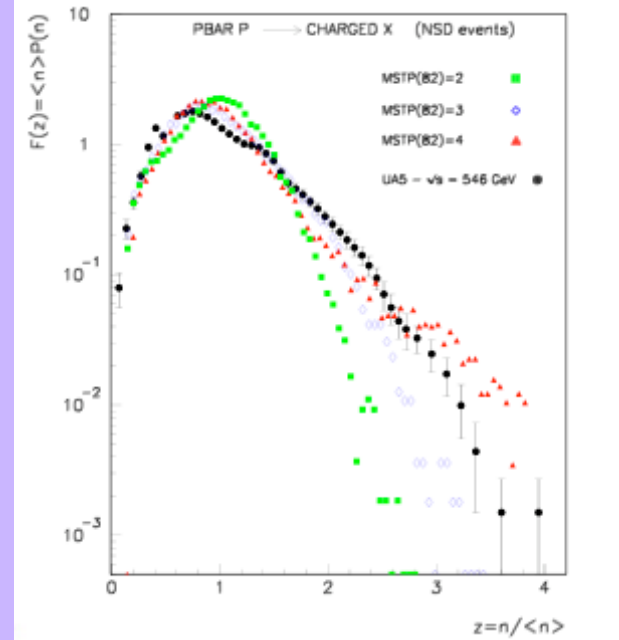
There are two options for multiple parton scattering [MSTP(82)] :

- "simple" scenario abrupt cut at $p_{t_{min}}$
- "complex" scenario smooth transition



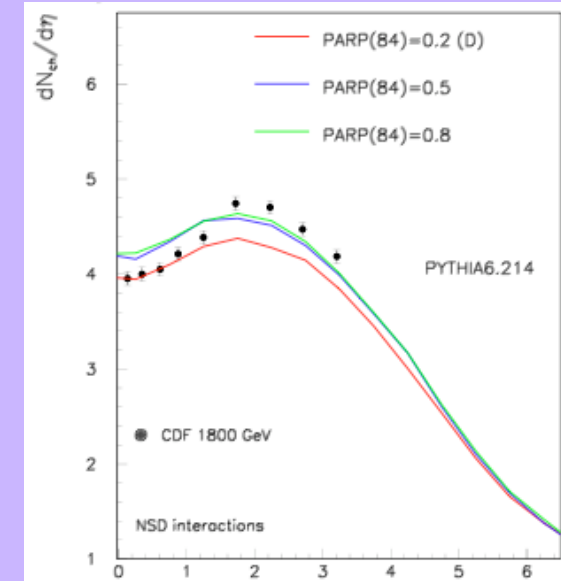
Proton matter distribution [PARP(84)]

$$\rho(r) \propto \frac{1-\beta}{a_1^3} \exp\left\{-\frac{r^2}{a_1^2}\right\} + \frac{\beta}{a_2^3} \exp\left\{-\frac{r^2}{a_2^2}\right\}$$



PT cut [PARP(82), PARP(90)]

$$P_{T_{min}}^{(s)} = P_{T_{min}}^{(s')} \left(\frac{s}{s'} \right)^{\epsilon}$$



Definitions

Minimum Bias (main component of the pileup):

Generic p-p interaction (hard, soft, SD, DD...)

σ_{tot}	=	σ_{Elastic}	+	σ_{SD}	+	σ_{DD}	+	σ_{HardCore}
(14 TeV)		~20 mb		~15 mb		~10 mb		~55 mb = ~100 mb

Underlying Event:

All the activity from a single p-p interaction superimposed to the hard scattering process

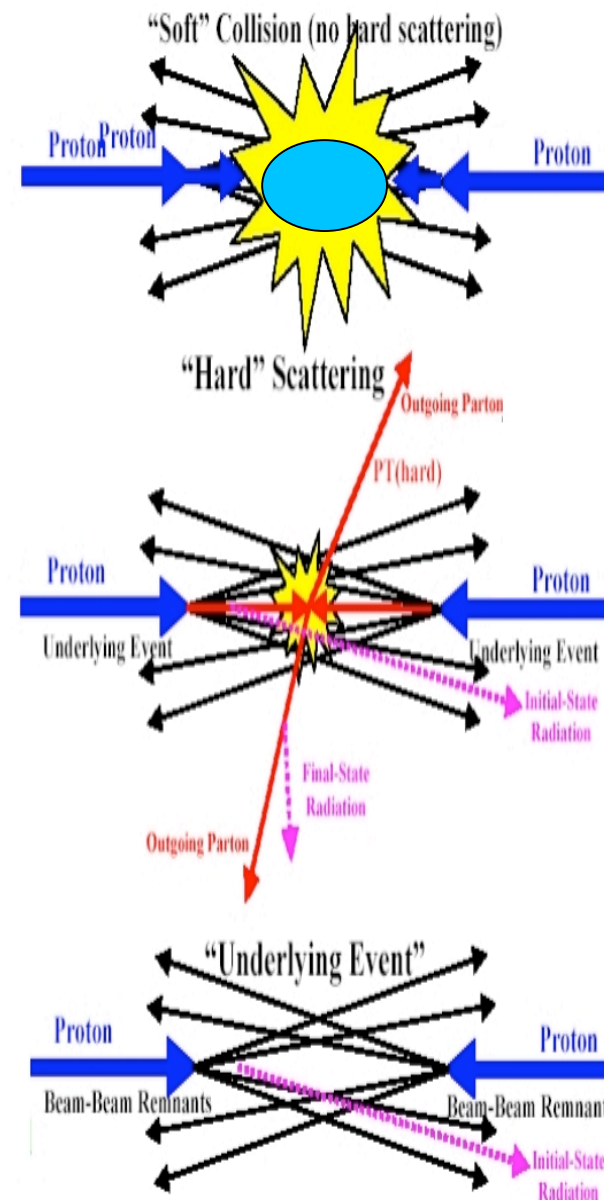
- + Initial and final state radiation
- + Spectators
- + beam-beam remnant
- + **Multiple Parton Interactions**

The UE is related to the hard scattering

- + primary vertex sharing
- + “pedestal effect”
(events with high pt jet in the final state show an higher UE activity)
- + color and flavor connected

UE! = Minimum Bias

but phenomenological aspects are similar



pQCD models

Generators setup used
(details in backup slides)

+ **Pythia Tune DW**

(from TuneA)

OLD MPI model,
IP CORRELATIONS

+ **Pythia Tune DWT**

DW and
default PT-cut-off evolution

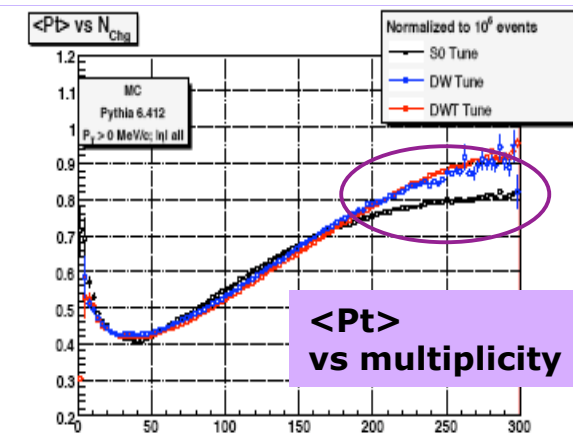
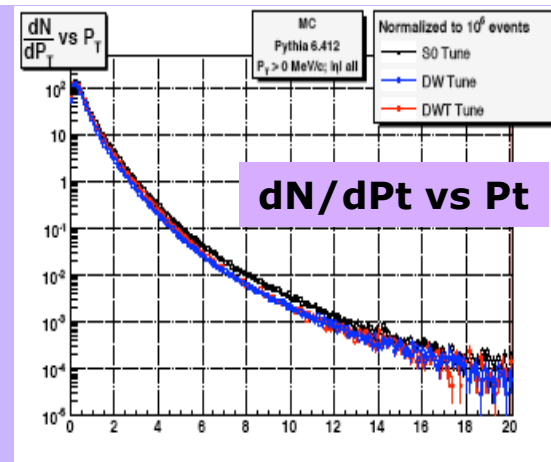
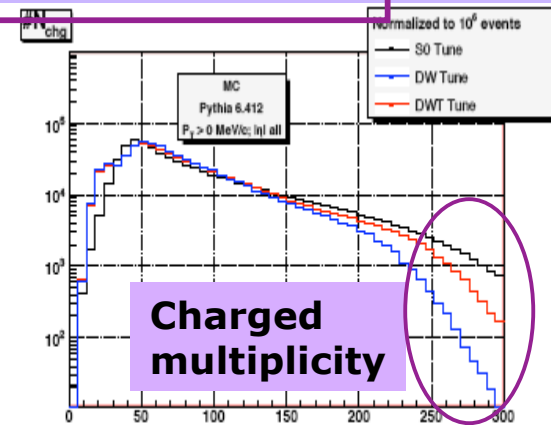
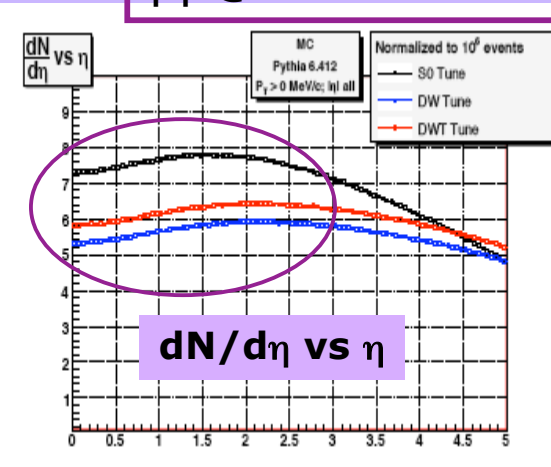
+ **Pythia Tune S0**

New MPI, more correlations

+ HERWIG

NO MPI, reference

Eg: effect on base MB observables
pp@14 TeV



All these Pythia Tunes describe the UE@Tevatron, but show several **differences extrapolating to LHC energy**

Not enough to re-scale conclusions to 14 TeV

[different Q^2 , x range and energy dependence of the cutoffs]

MB Triggers



- + Using forward detectors
- + Random trigger (inefficient for $N_{int} \ll 1$)
- + Pile up from other streams
(jet or lepton triggers as example)

Best solution, at the moment, seems to be the one based on forward detectors

The use of pile up from other streams could introduce biases (under study)

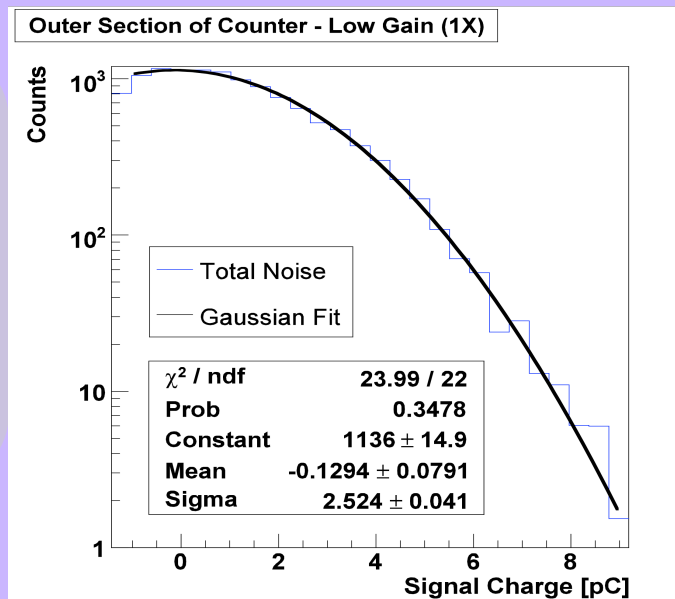
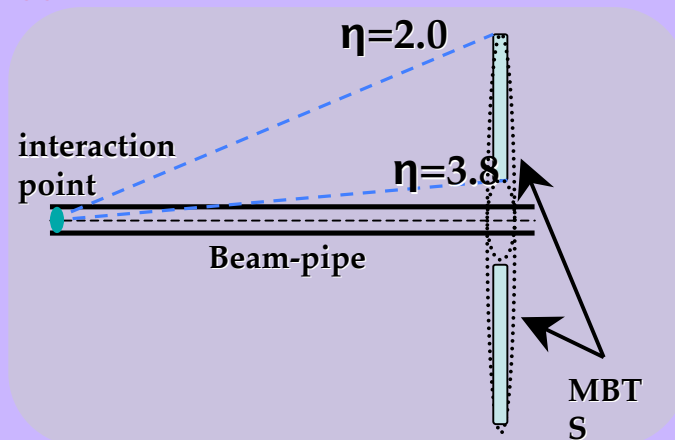
MB Triggers - **ATLAS**



MBTS

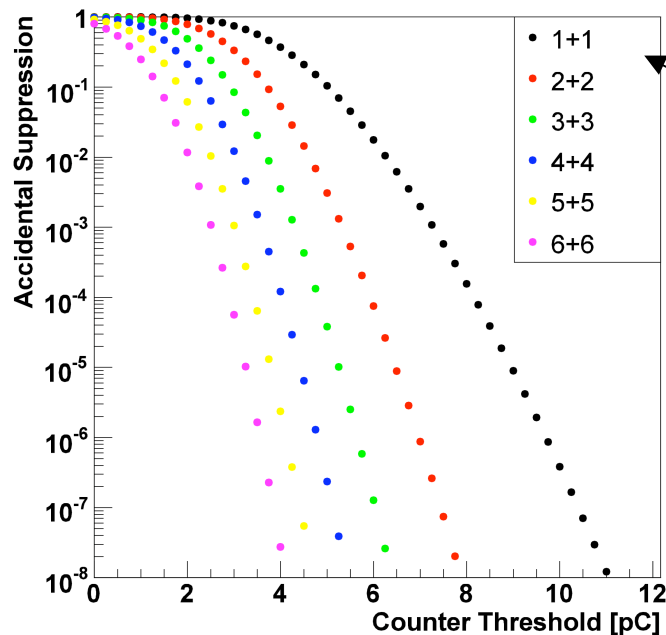
(minimum bias
trigger scintillator)

Trigger scintillation counters mounted on end of LAr calorimeter covering same radii as the inner detector



Noise spectrum
fits well to a
Gaussian with
 $\sigma=2.52\text{pC}$

Beyond 9pC
(almost 4σ),
non-Gaussian
behavior is
possible

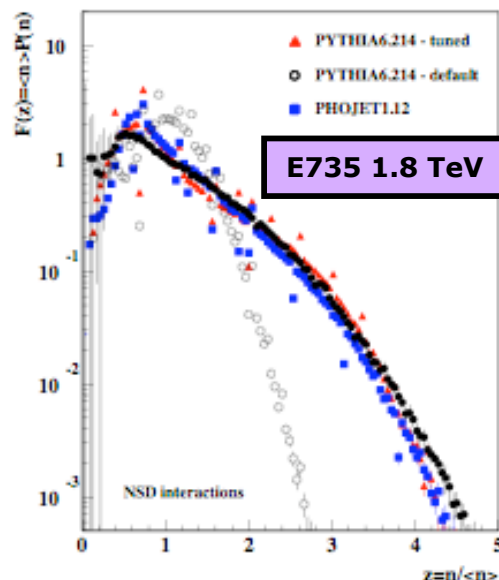
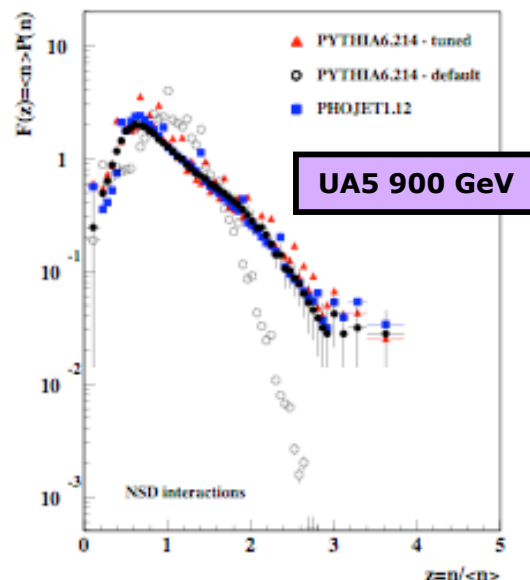


of backward-forward coincidences

+ Accidental rate from noise must be
suppressed to $\sim\text{Hz}$, limited by EF
output-rate of 100Hz

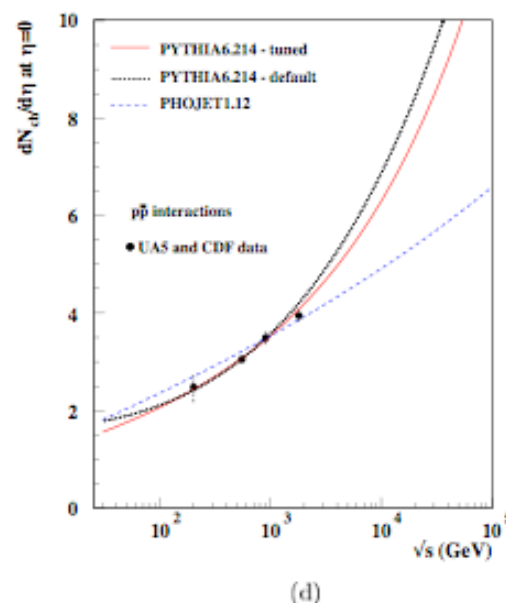
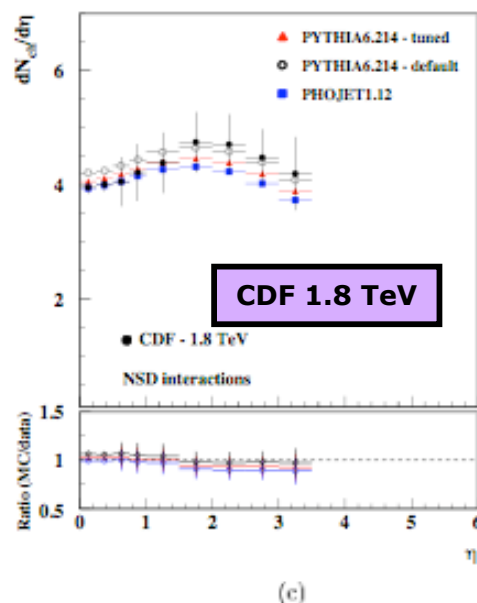
+ Suppression required of
(10^{-6}) 10^{-7} at (900 GeV) 14 TeV

MB measurement - **ATLAS** predictions



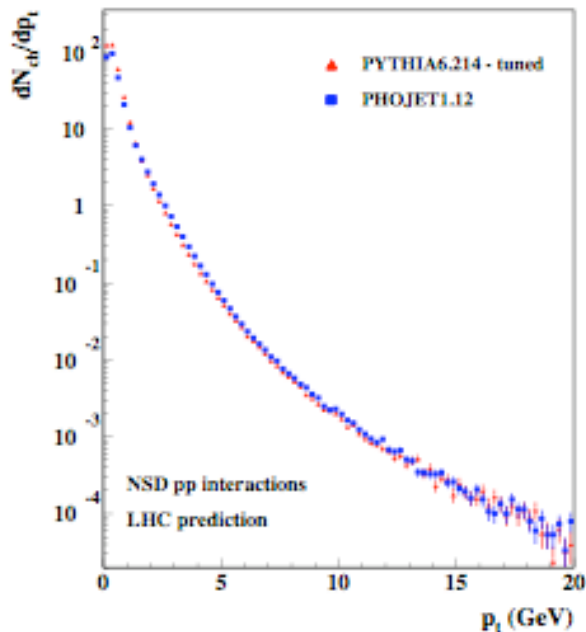
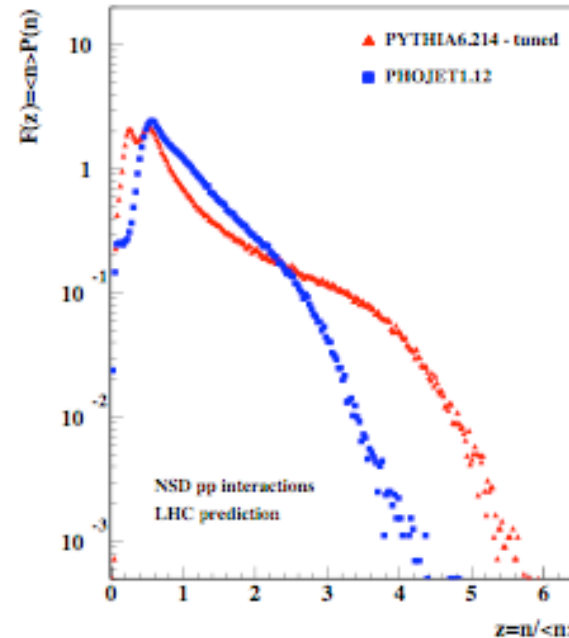
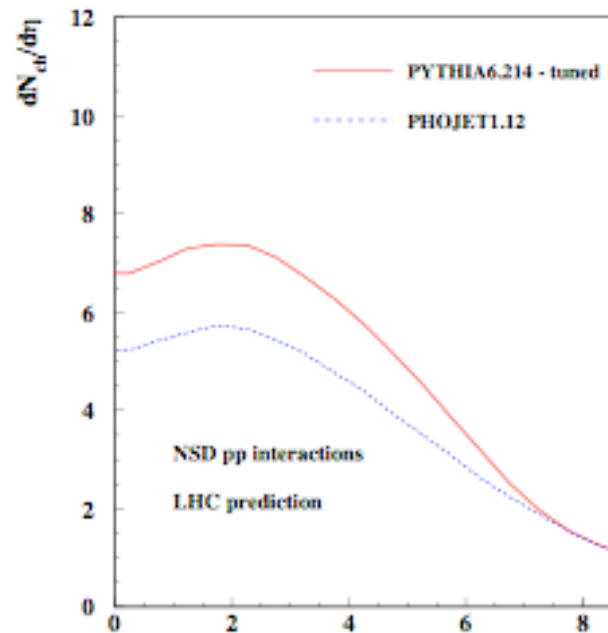
Starting from KNO scaling
on past experiment:

- + Observables are defined
- + Models are exploited
- + **Tune based on MB and UE observables** are chosen



Models	$\chi^2_{\text{min-bias}}$	χ^2_{UE}	χ^2_{global}
PHOJET1.12	8.15	7.35	7.95
PYTHIA6.214 - tuned	11.62	2.07	9.27
CDF tuning	26.66	1.31	20.43
ATLAS - TDR	38.98	15.03	33.09
PYTHIA6.214 - default	68.01	22.68	56.87

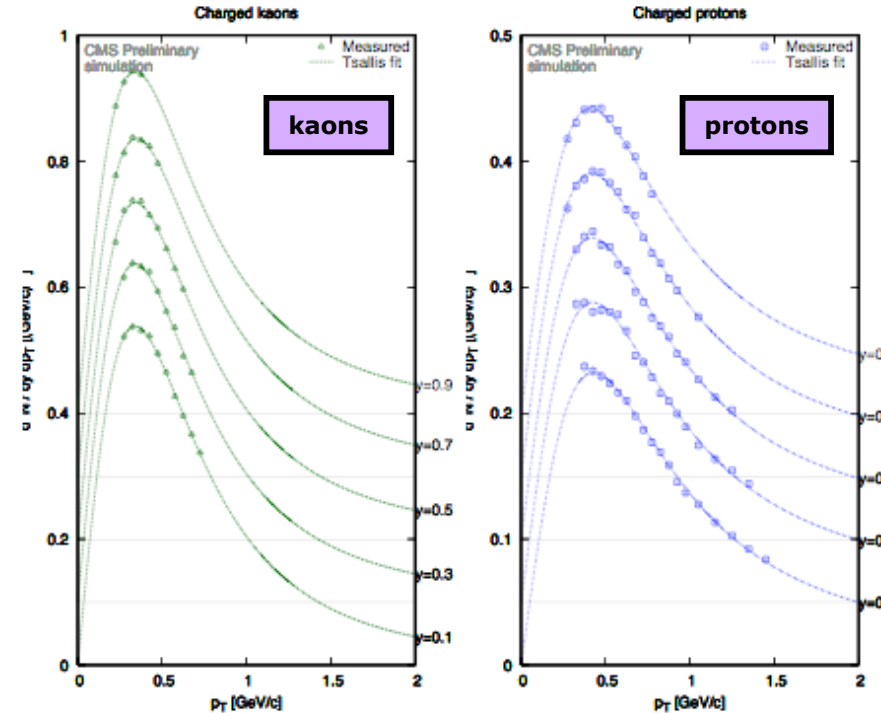
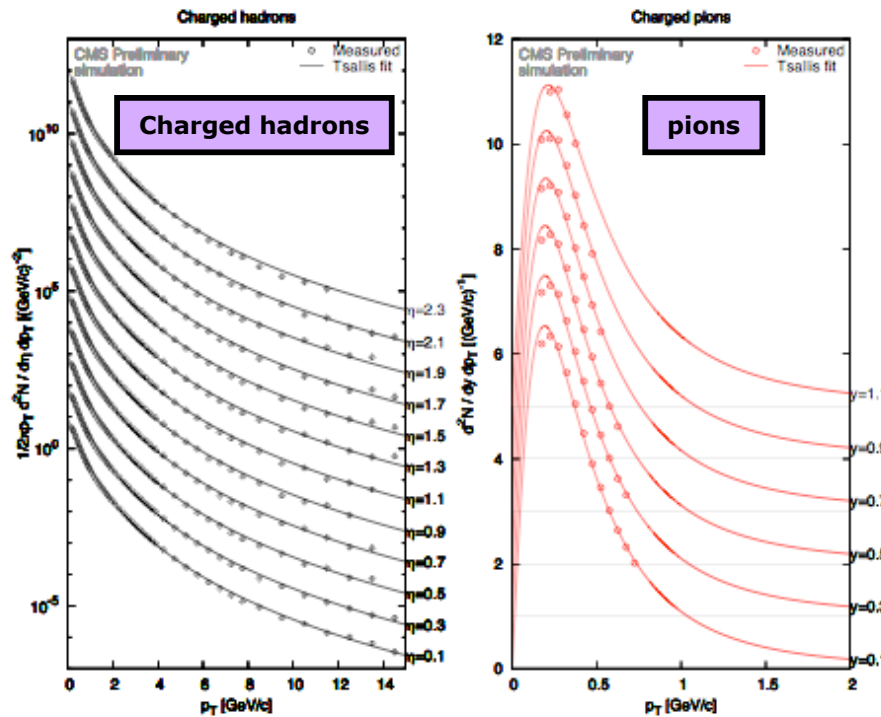
MB measurement plan - **ATLAS**



Foreseen measurements:

- + NSD events
- + charged spectra (p_T , η)
- + fragmentation

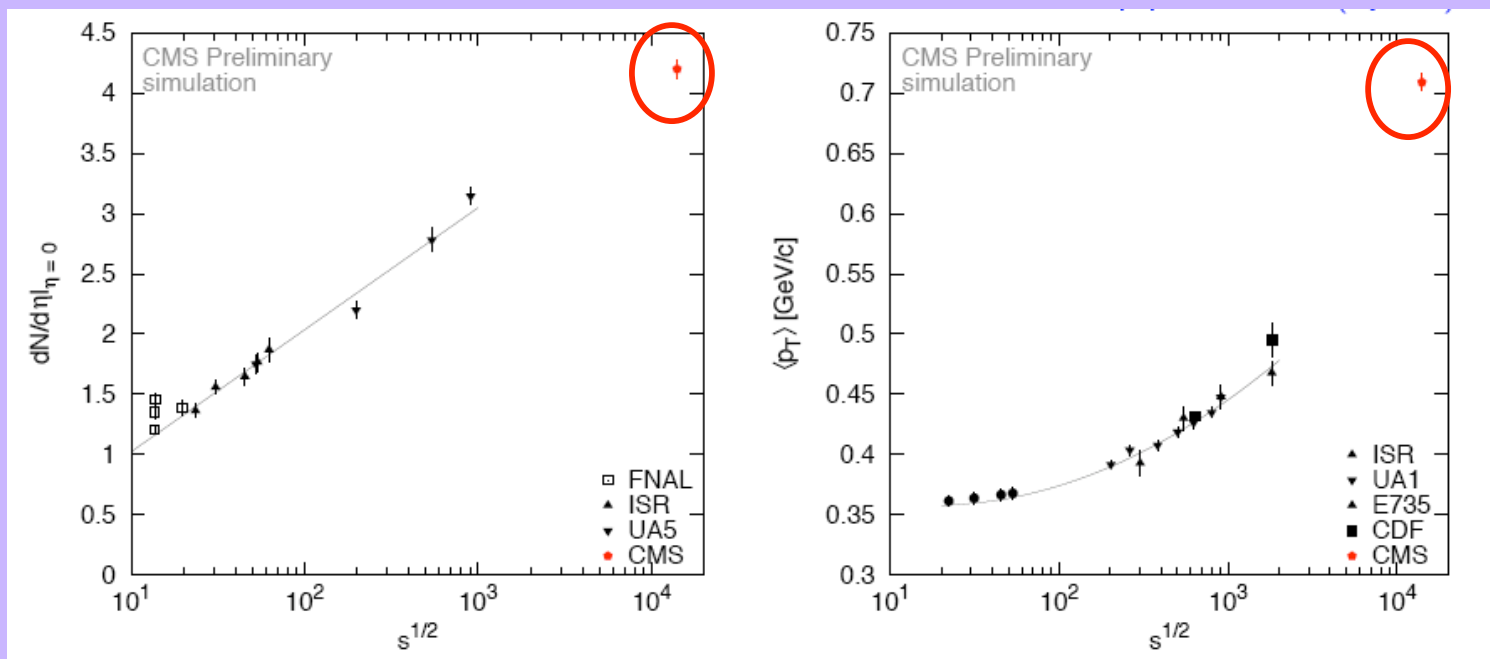
MB measurement - CMS



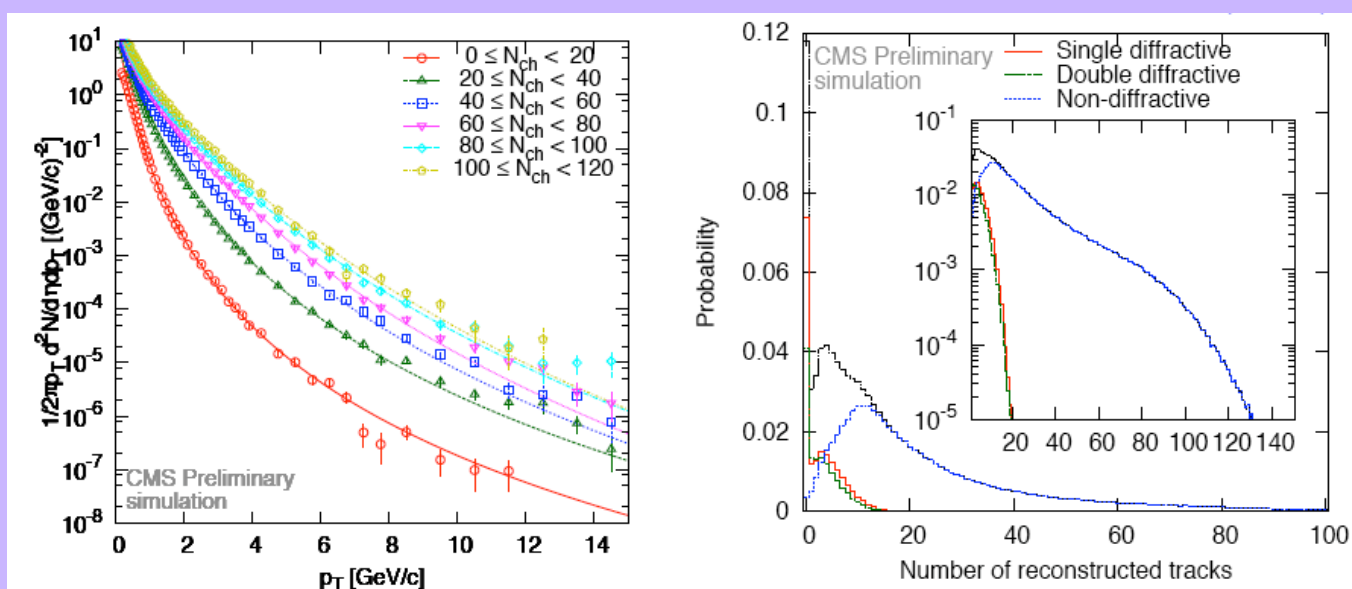
$$E \frac{d^3N}{dp^3} = \frac{dN}{dy} \frac{(n-1)(n-2)}{2\pi nT [nT + (n-2)m]} \left[1 + \frac{E_T(p_T)}{nT} \right]^{-n}$$

From Tsallis fit and expected performances on tracking, particle corrections are calculated and applied (PID performed with dE/dx)

MB measurement - CMS



LHC
expected
multiplicity
and
average
particle p_T



p_T
distribution
and
particle
multiplicity

UE observables definition

From charged jet

Topological structure of p-p collision from charged tracks

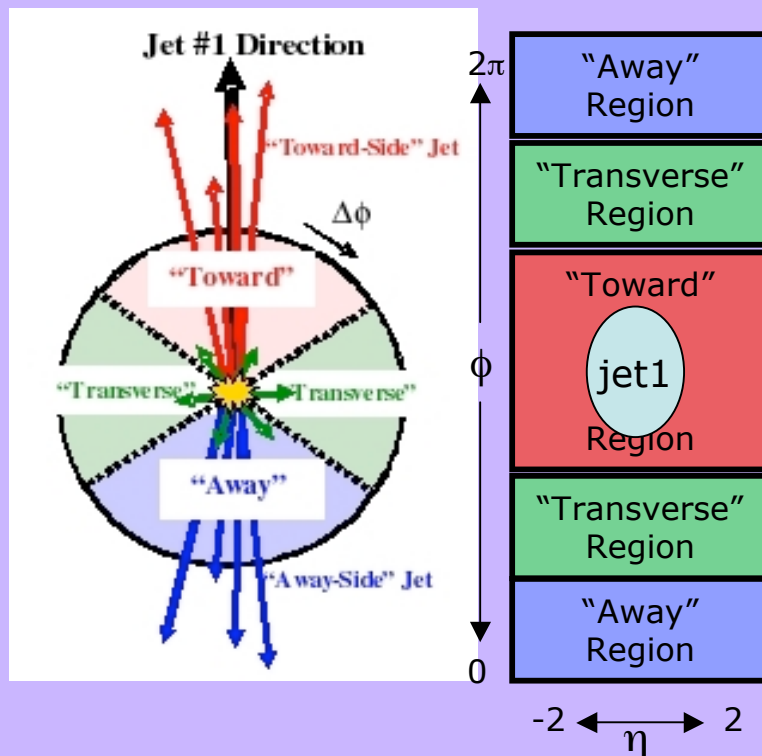
Charged jet definition -> ICA algorithm with massless tracks as input

The leading Ch_jet1 defines:

+ a **direction in the ϕ plane**

+ the PT is used as **reference for the energy scale** of the interaction

The transverse region is particularly sensitive to the UE



Main observables:

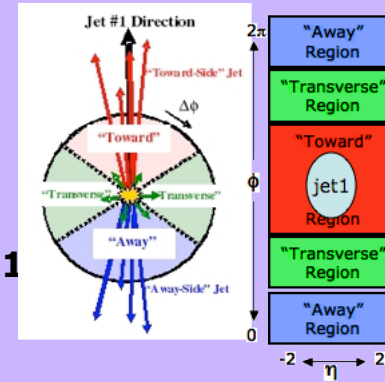
- ❖ $dN/d\eta d\phi$, charged density
- ❖ $d(PT_{\text{sum}})/d\eta d\phi$, energy density

Feasibility@LHC

UE Observables VS

$\Delta\phi$ Leading Charged Jet - Track

Uncorrected distributions from **10pb⁻¹**



Tracks:
PT>0.9
|η|<2

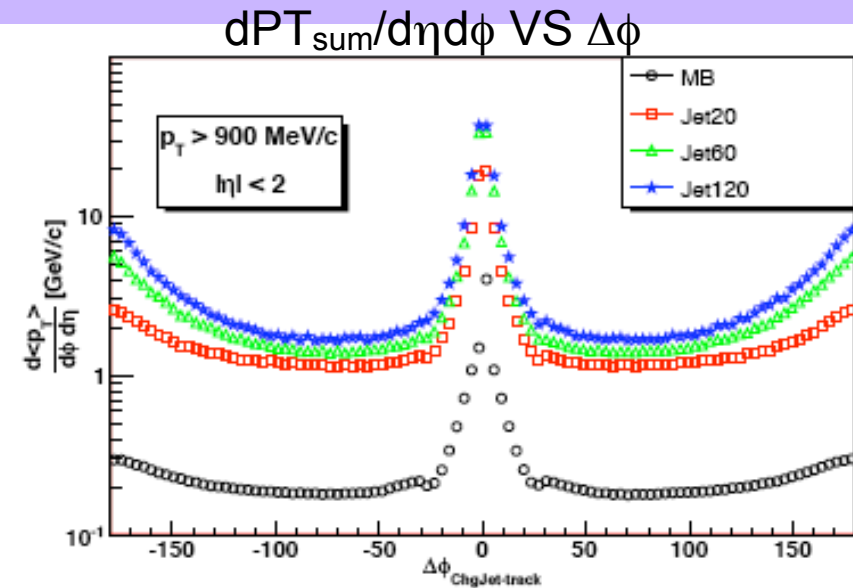
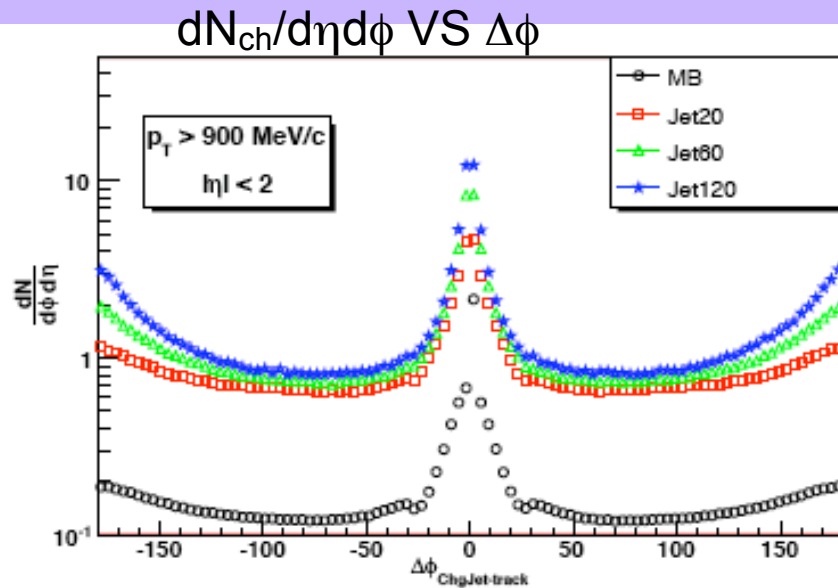
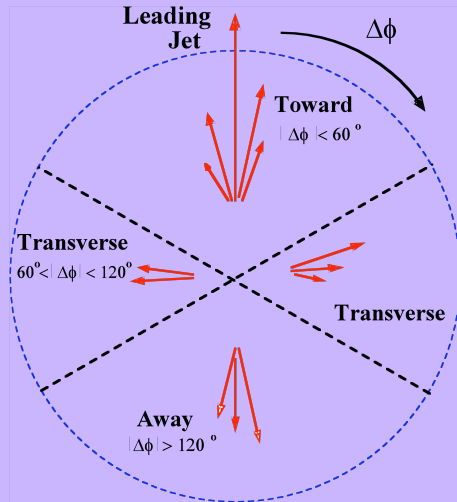


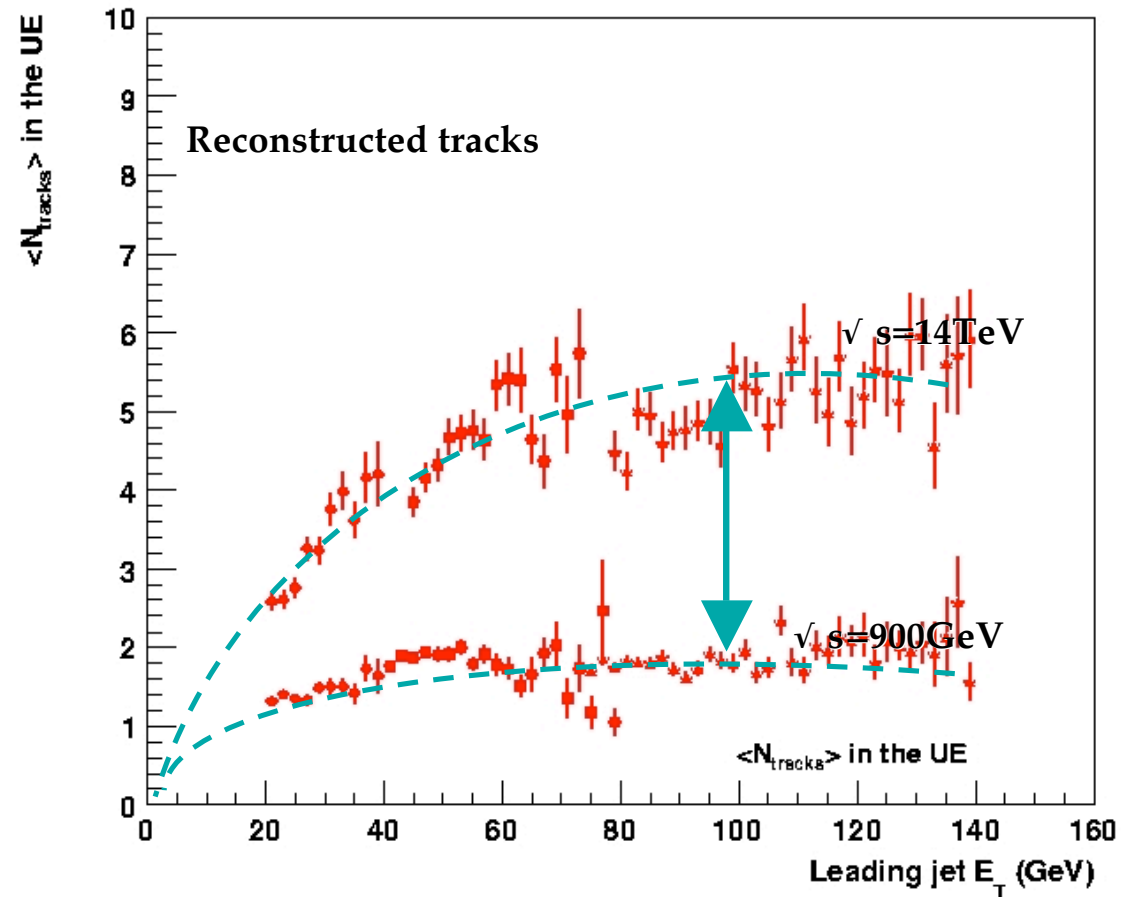
Figure 2: Density of charged particles, $dN/d\eta d\phi$ (left), and average charged $\sum p_t$ density (right) versus the azimuthal distance between charged tracks and leading charged jet. Reconstructed data from different triggers are superimposed.

ATLAS - jet events

Results in the "transverse" region

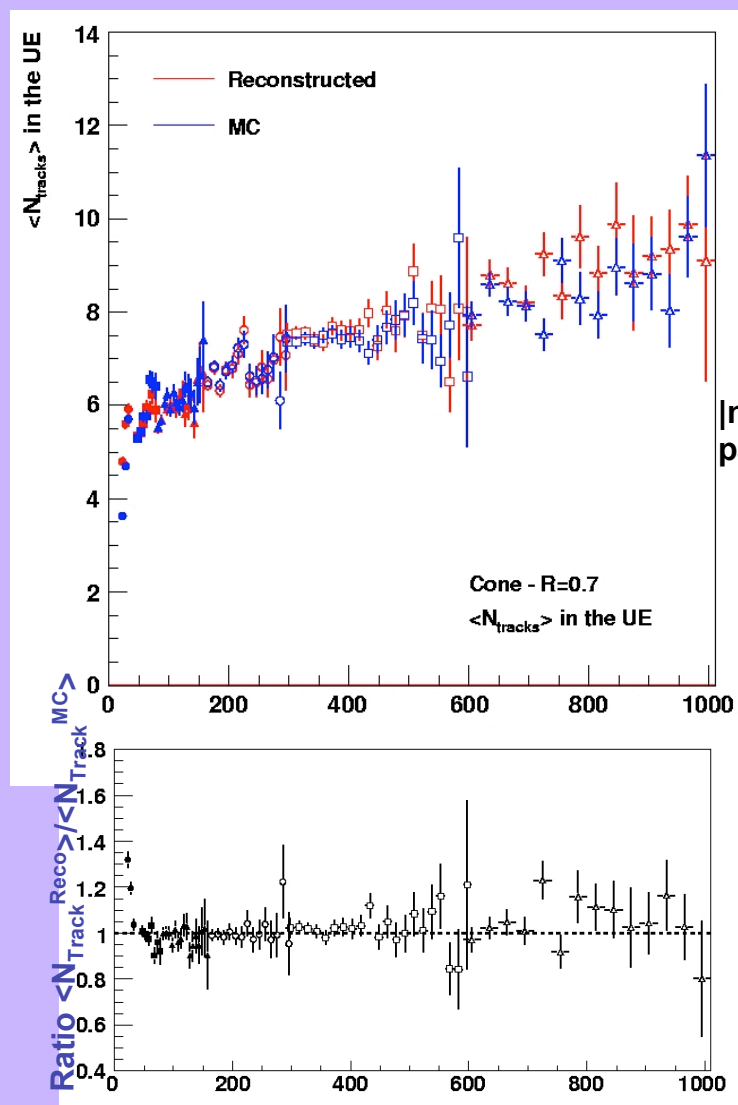


Multiplicity of charged particles with $p_T > 0.5$ GeV and $|\eta| < 1$ in region transverse to leading jet



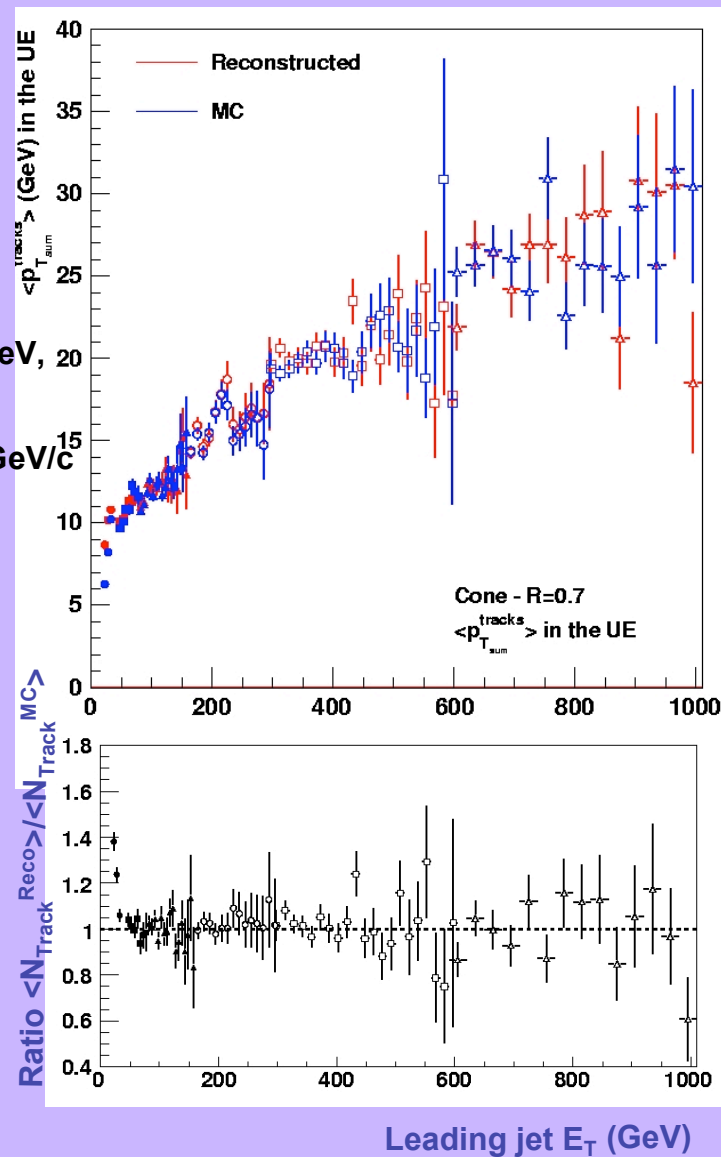
~ 15 days of data taking enough to cover up to $p_T(\text{leading jet}) \sim 40$ GeV

ATLAS - jet events up to TeV



$N_{\text{jets}} > 1,$
 $|\eta_{\text{jet}}| < 2.5,$
 $E_{\text{T}}^{\text{jet}} > 10 \text{ GeV},$
 $|\eta_{\text{track}}| < 2.5,$
 $p_{\text{T}}^{\text{track}} > 1.0 \text{ GeV}/c$

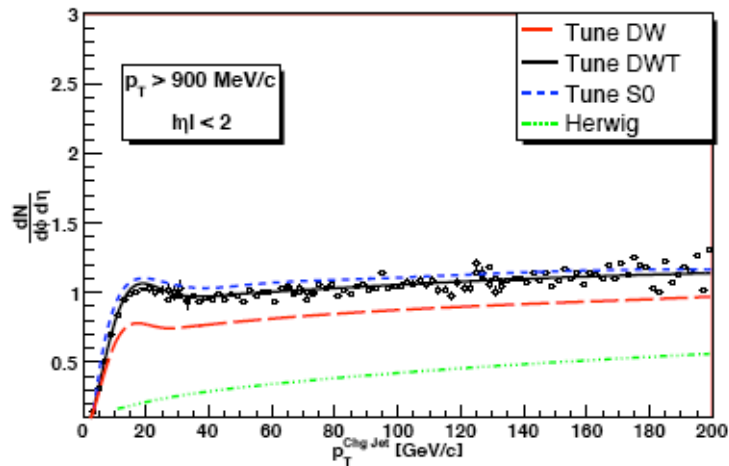
Cone - R=0.7
 $\langle N_{\text{tracks}} \rangle$ in the UE



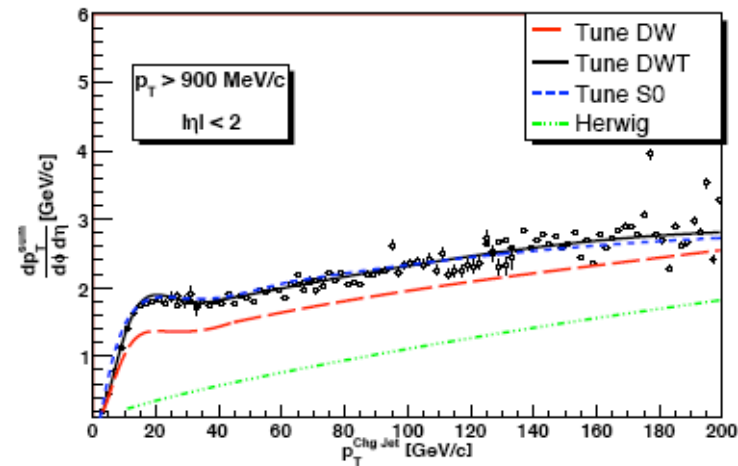
Leading jet E_T (GeV)

Very good RECO/MC agreement

CMS - jet events - 100 pb-1



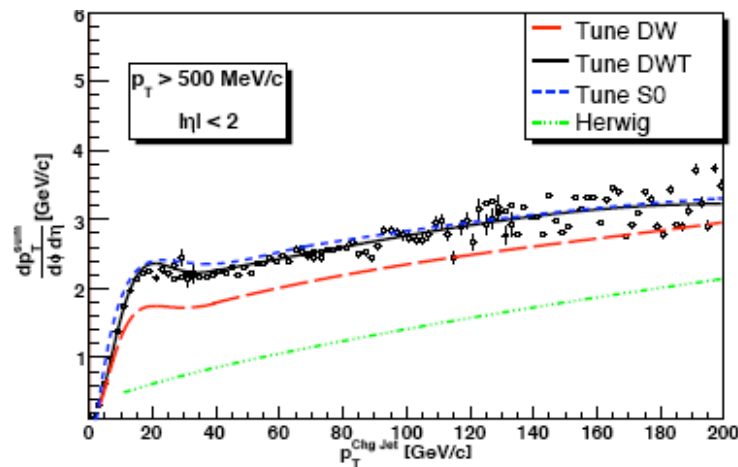
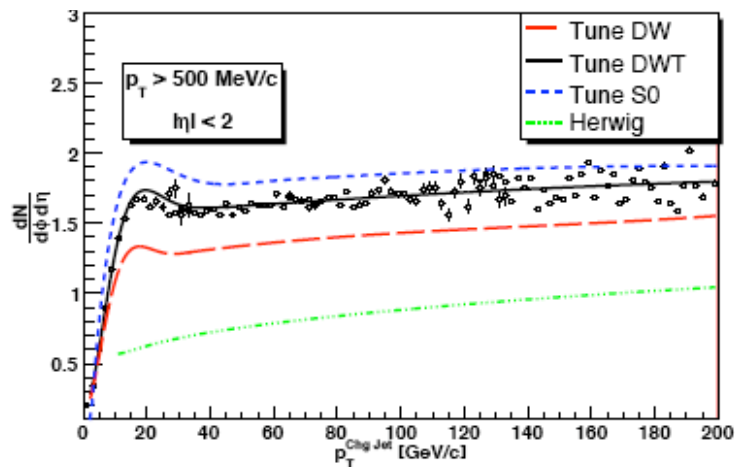
$dN/d\eta d\phi$ vs PT_chg



$d(P_{T_sum})/d\eta d\phi$ vs PT_chg

Figure 6: Same as figure 4 but for corrected distribution.

$p_T > 900$ MeV/c



$p_T > 500$ MeV/c

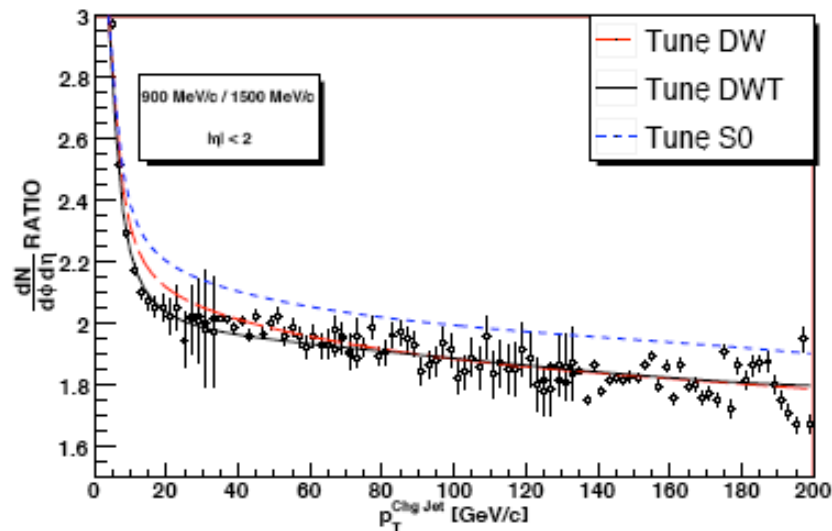
Effect from correction
Good discrimination power

- > get back the DWT
- > DW/DWT (from 900 MeV)
- S0/DWT using tracking from 500 MeV

CMS - jet events - 100 pb⁻¹



$dN/d\eta d\phi$ ratio vs PT_{chg}



$d(PT_{\text{sum}})/d\eta d\phi$ ratio vs PT_{chg}

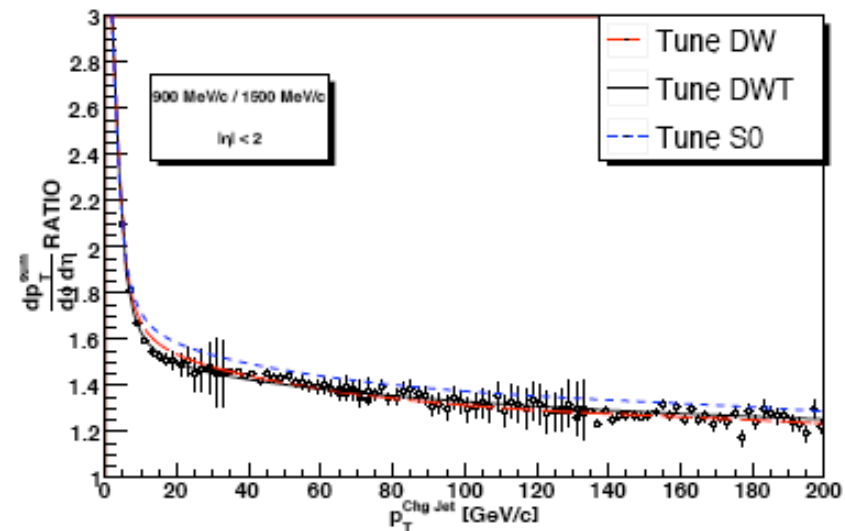


Figure 9: Ratio plots using track with $p_T > 1.5 \text{ GeV}/c$ and $p_T > 0.9 \text{ GeV}/c$ for density of charged particles $dN/d\eta d\phi$ (left) and average charged $\sum p_t$ density (right) versus the transverse momentum of the leading charged jet.

Uncorrected data.

Ratios of observables using minimum p_T of 900 MeV/c and 1.5 GeV/c

- +No need to apply corrections, absorbed in the ratio
- +Additional discrimination between tunes

summary



Minimum Bias measurement plans exist for both experiments

Measuring charged hadron spectra will allow to calibrate and understand soon the detectors and establishing a solid basis for exclusive physics:

Monte Carlo predictions, based on Tevatron data, greatly differ if extrapolated to LHC energy (MPI component)

Underlying Event activity is studied in the transverse region of charged jet events (studies exist in CMS on DY events - not presented here)

Measuring UE will allow us to

tune the energy dependence models (largely related to the MPI) ->
improve the QCD understanding in pp collisions ->

fundamental for all the LHC measurements ("old" and "new" physics)

Strategy proposed by CMS:

1 pb⁻¹ -> tools calibration (tracking, triggers, correction and response function...)

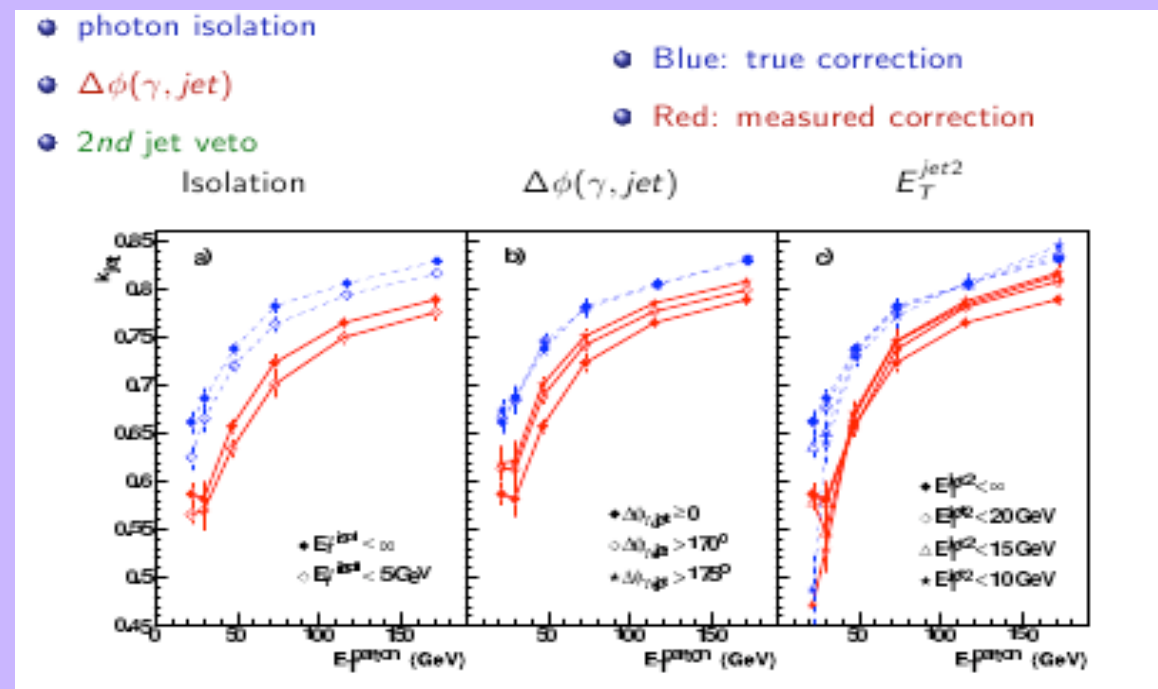
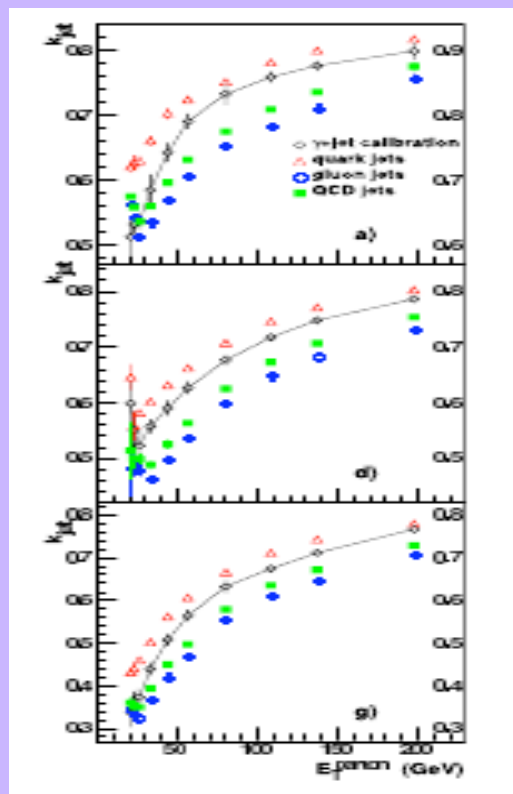
10 pb⁻¹ -> start control of systematics and discriminating between models

100 pb⁻¹ -> deeper discrimination, enhanced considering ratio distribution

Back up - Jets

Jet Energy Scale at CMS

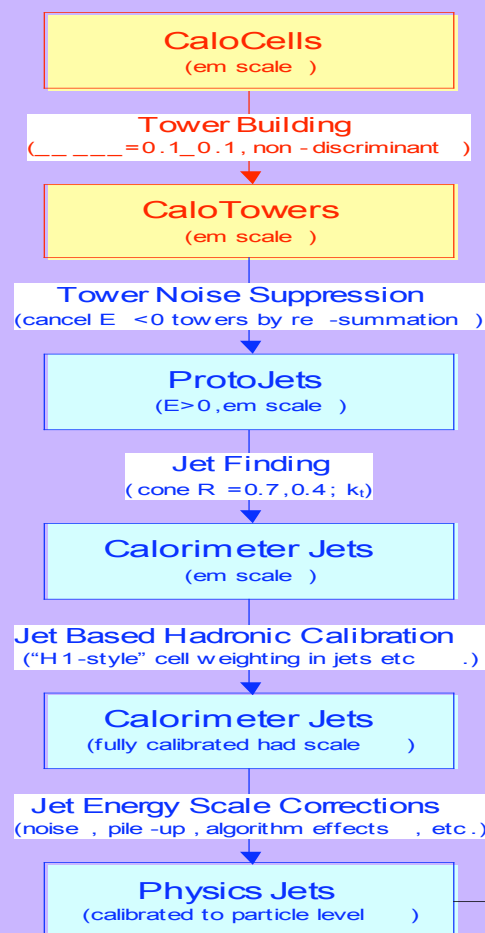
- Investigating several approaches to compute the absolute P_T corrections:
 - Derivation from the data through γ +jets or missing E_T projection method: careful study of the introduced bias
 - Bias studied also as a function of the jet type



Jet energy scale in ATLAS



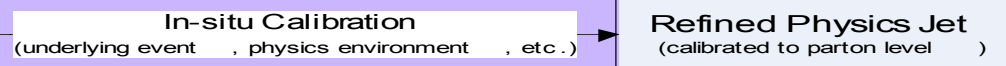
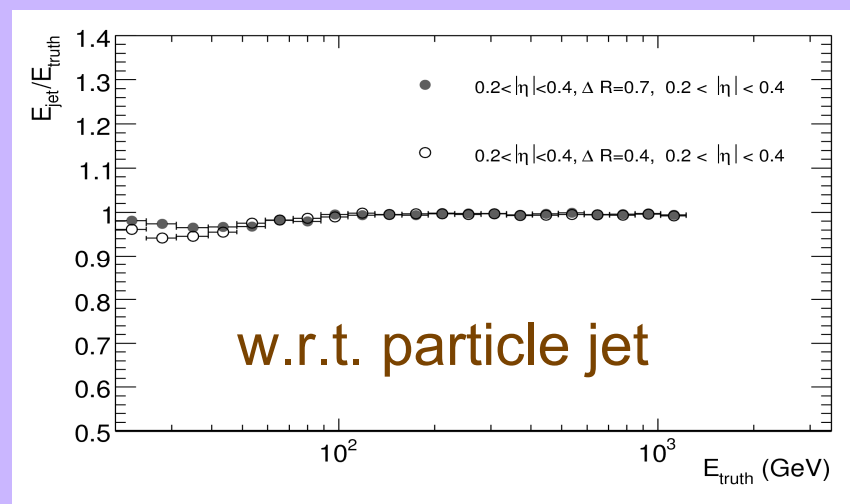
- The **real** strategy will come with the data.
- The **current** strategy:



$$E_{rec}^{jet} = \sum_{i=cells} w_i \frac{E_i}{\sqrt{V_i}} \sqrt{E_i}$$

$$w_i \frac{E_i}{\sqrt{V_i}} = \sum_{j=0}^3 a_j \log \frac{E_i}{V_i}$$

**Fully MC based corrections:
the importance of the test
beam**



Current strategy: a summary



- Topological clustering as a noise suppression tool
- The main step is the correction at particle jet:
 - The corrections will tent to improve the jet resolution. The corrections will be heavily MC based.
- Real data (in situ) will be used for:
 - Check the uniformity of the jet scale with the di-jet balance
 - Check the MC predictions in “easy” events (γ +jet, Z+jet)
 - Check the scale of high P_T jets with bootstrap techniques
- UE subtraction: MC based after tuning
- Pile-up: topological clustering is supposed to account for that:
 - Anyhow, strategies for subtraction based on the number of vertices are in place

H1 jet corrections

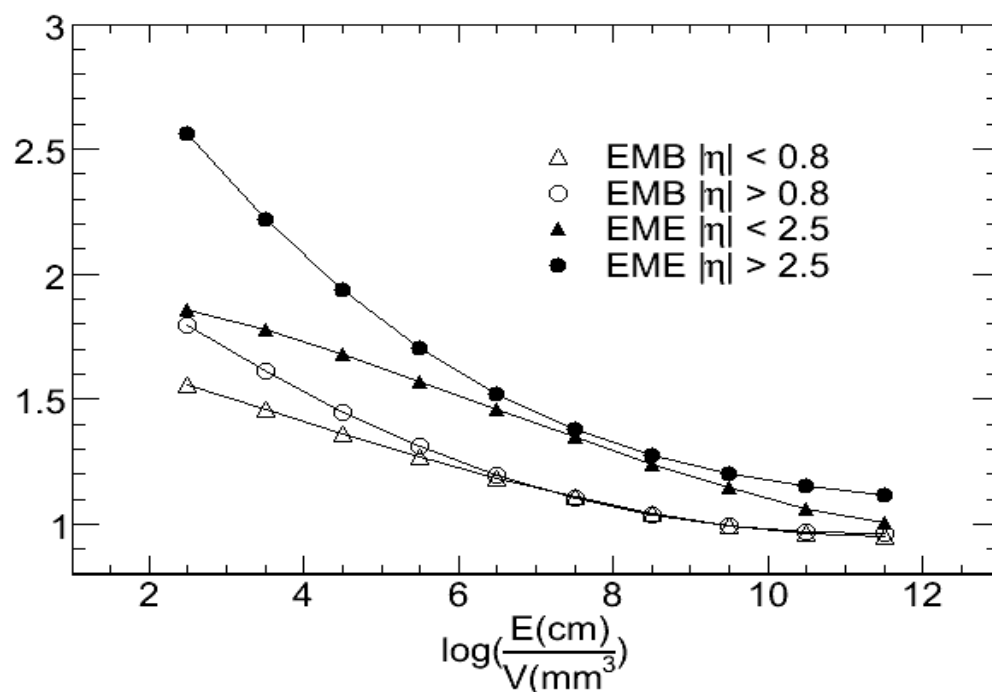
- The total energy is computed as follows:
- The weights depend on the **cell energy density** (and pseudorapidity, through a rough binning)

$$E_{rec}^{jet} = \sum_{i=cells} w_i \frac{E_i}{V_i} E_i$$

$$w_i \frac{E_i}{V_i} = \sum_{j=0}^3 a_j \log \frac{E_i}{V_i}$$

The weights enhance the response of low energy density cells (Had energy).

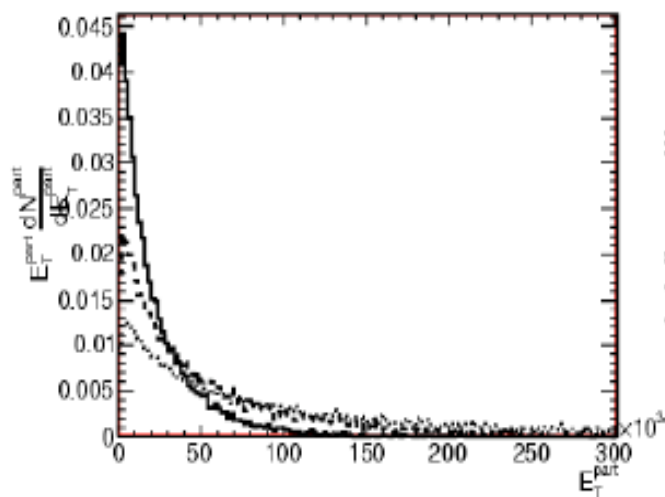
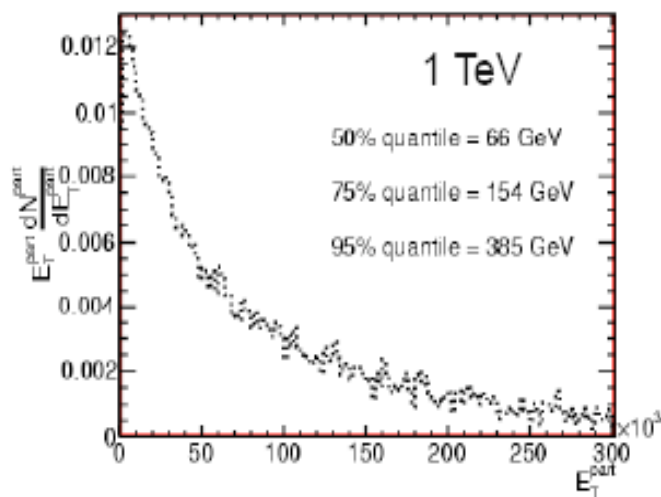
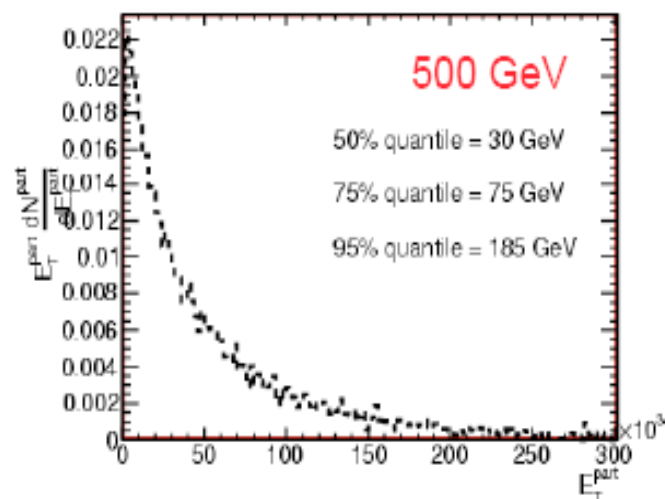
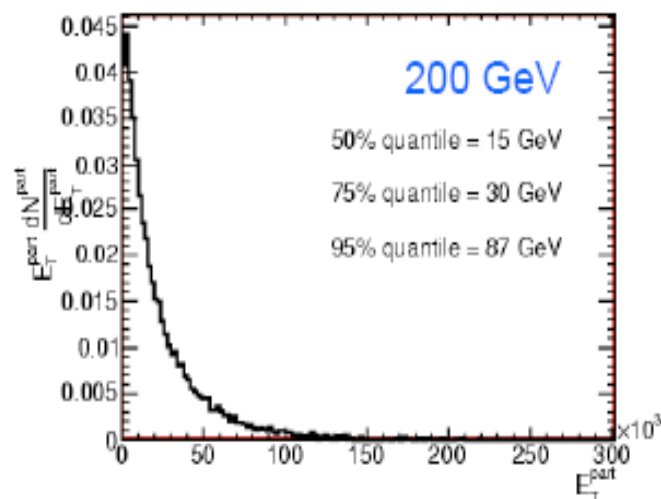
High density is associated with EM deposits: the weights go to 1



H1 jet corrections

- The weights w_i are **the same for all the reco algos**, all inputs.
 - They are the same still hard coded for EtMiss reconstruction
- Crack and Gap regions not taken into account in the weight computing. For each reco algo and input clusters, a **scale factor** is computed:
 - Intended to correct for non-linearities **caused by the gap and the crack**
 - It corrects also for **residual reconstruction algorithm dependent effects**

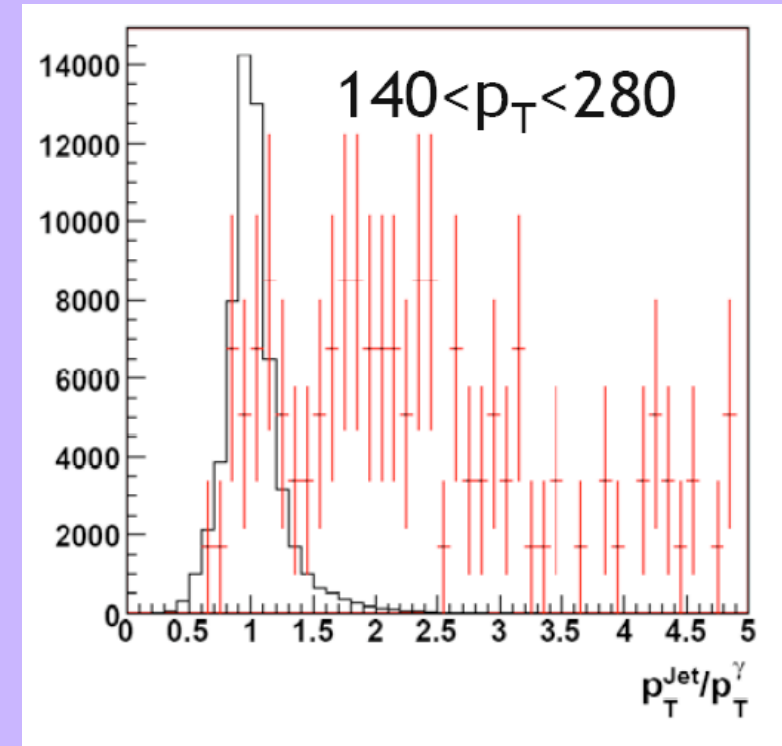
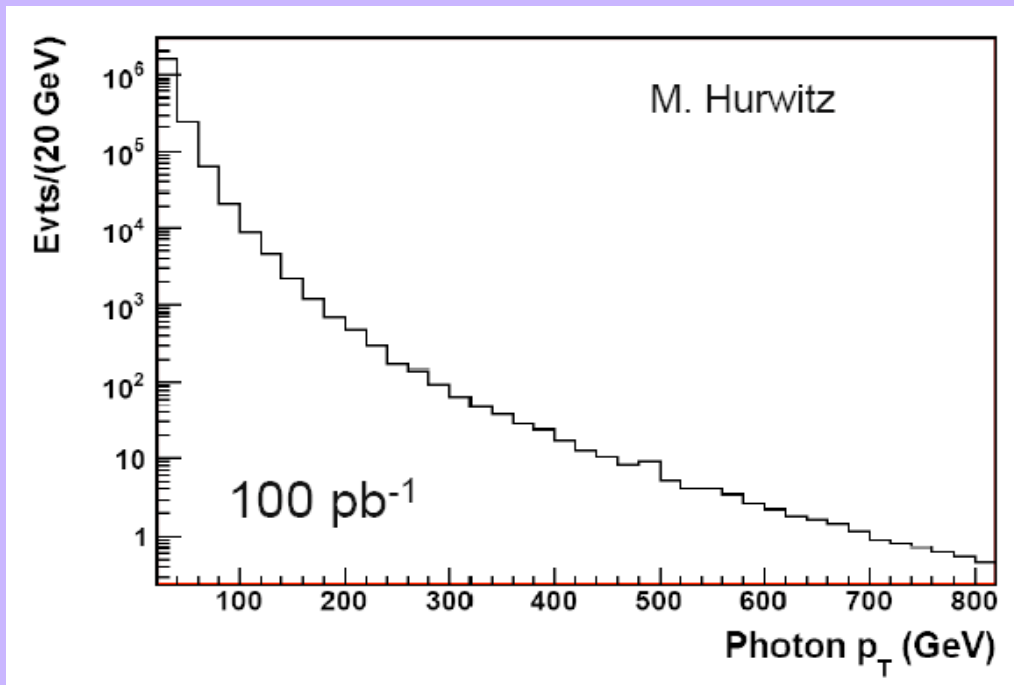
Jet Constituents



I. Viverelli

Inclusive Jet Cross Section (4)

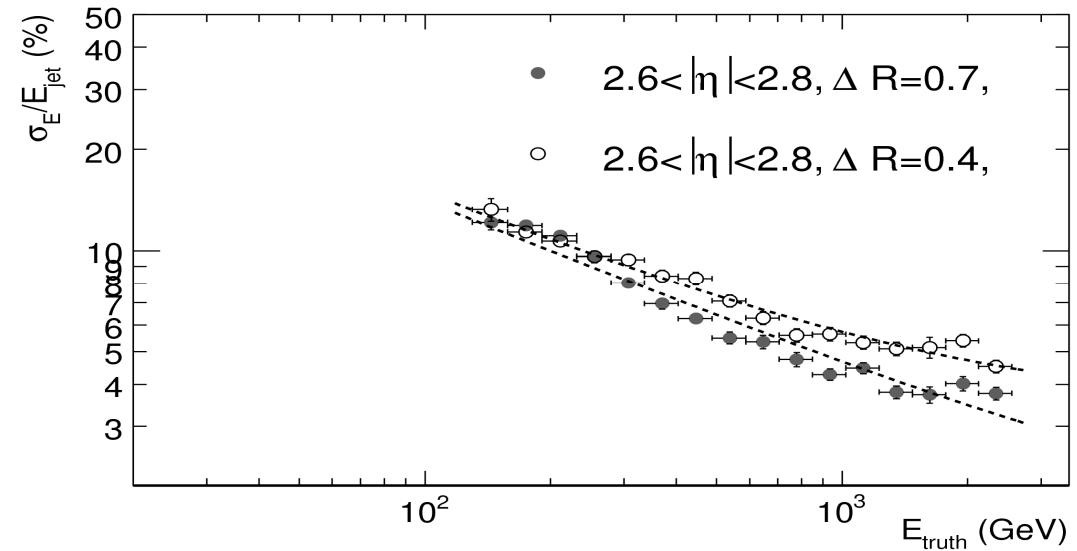
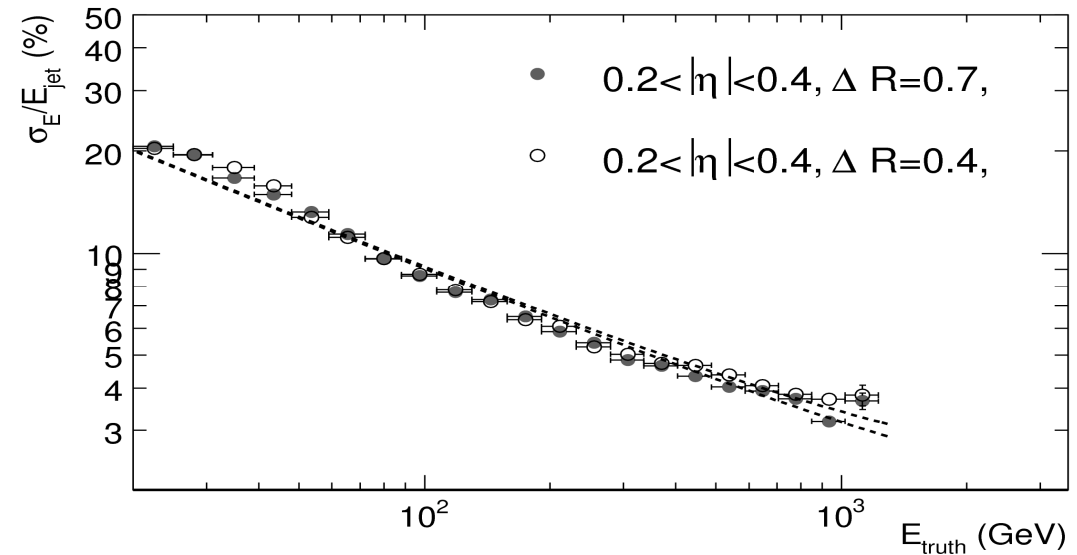
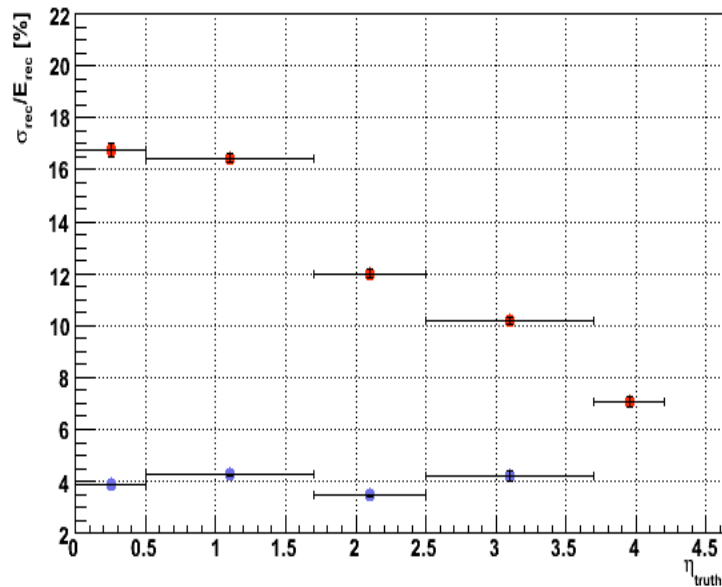
- Check the energy scale with γ +jet
- The balance looks good, but:
 - Large background from DiJet
 - The balance **does not work** for photon-like jets (in red in the picture)



- 100 pb^{-1} : γ +jet can probe up to few hundreds GeV:
 - We need a way to get up to the TeV scale

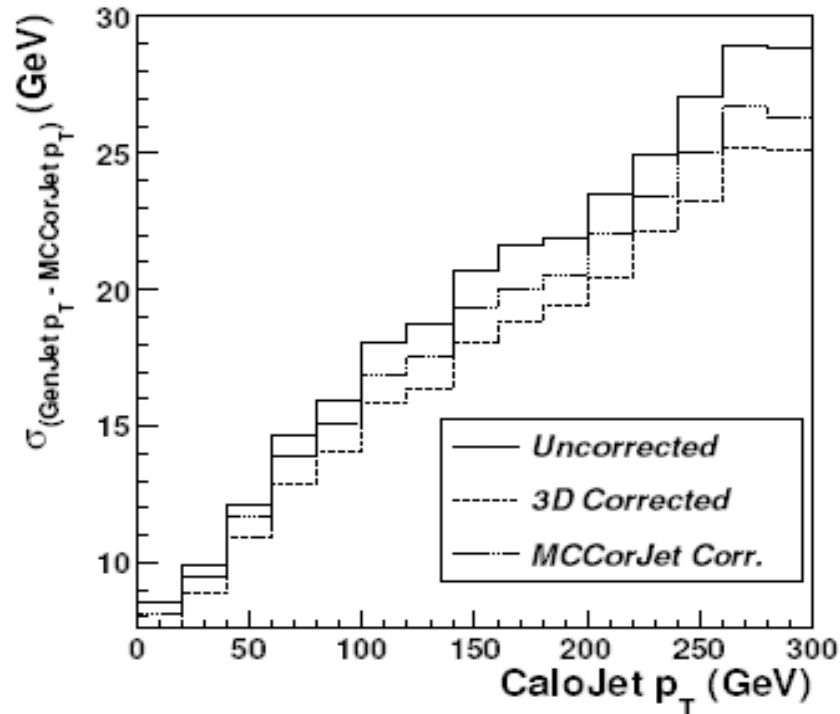
Jet Resolution (ATLAS)

- Resolution:
 - Central region: $\sim 12\%$ at 50 GeV, $\sim 8\%$ at 100 GeV

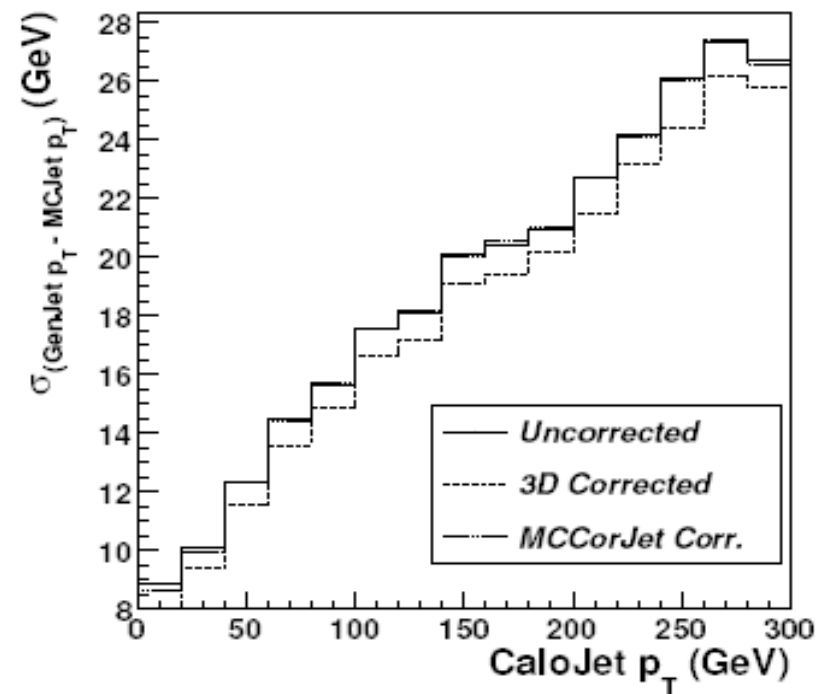


Jet Resolution (CMS)

Jet Resolution for $\eta < 2.5$

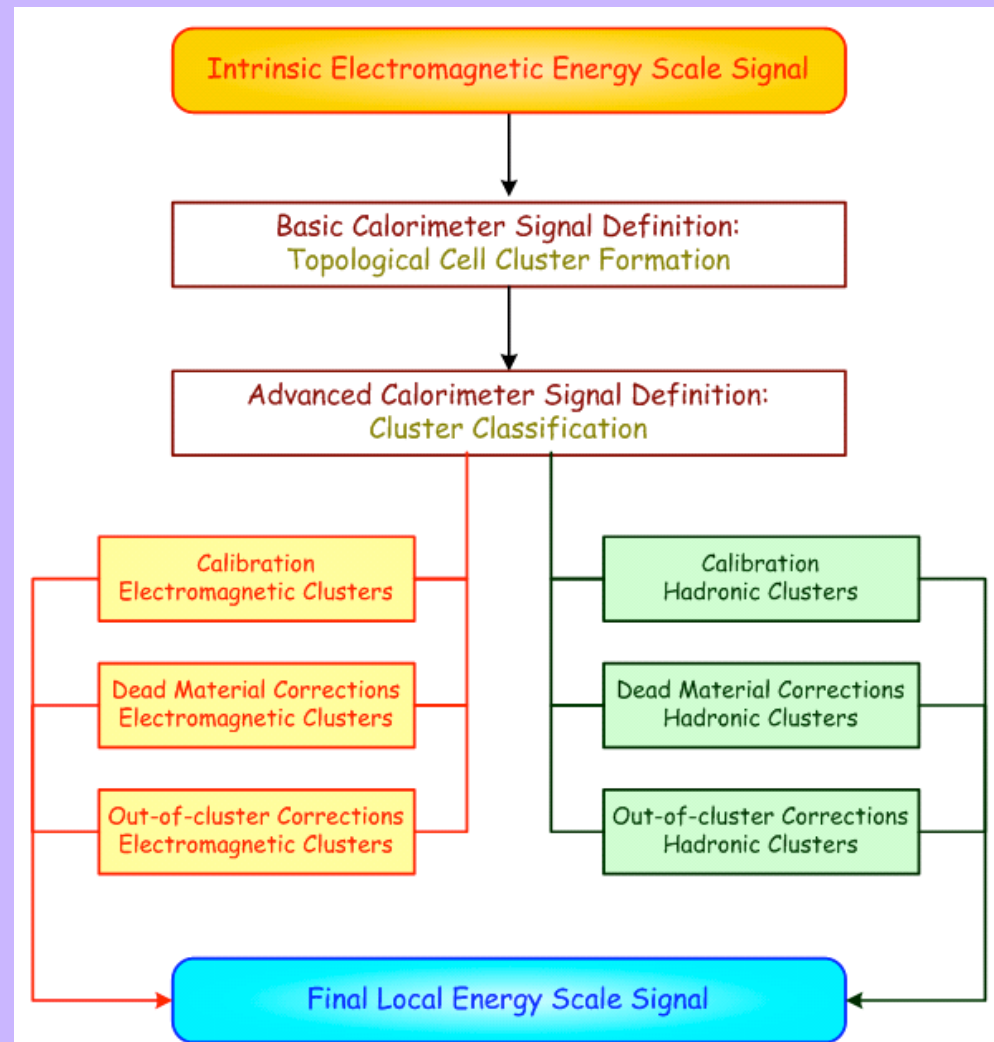


Jet Resolution for $\eta < 1.0$



ATLAS: toward a Local Hadron Calibration

- The local hadron calibration tries to exploit the high ATLAS calorimeter granularity
- Classification of the clusters based on cell energy density and shape variables
- Corrections applied based on the classification



Back up - MBUE

Parameter (PYTHIA v.6412+)	A	ATLAS	DW	DWT	S0
UE model MSTP(81)	1	1	1	1	21
UE infrared regularisation scale PARP(82)	2.0	1.8	1.9	1.9409	1.85
UE scaling power with \sqrt{s} PARP(90)	0.25	0.16	0.25	0.16	0.16
UE hadron transverse mass distribution MSTP(82)	4	4	4	4	5
UE parameter 1 PARP(83)	0.5	0.5	0.5	0.5	1.6
UE parameter 2 PARP(84)	0.4	0.5	0.4	0.4	n/a
UE total gg fraction PARP(86)	0.95	0.66	1.0	1.0	n/a
ISR infrared cutoff PARP(62)	1.0	1.0	1.25	1.25	(= PARP(82))
ISR renormalisation scale prefactor PARP(64)	1.0	1.0	0.2	0.2	1.0
ISR Q_{max}^2 factor PARP(67)	4.0	1.0	2.5	2.5	n/a
ISR infrared regularisation scheme MSTP(70)	n/a	n/a	n/a	n/a	2
ISR FSR off ISR scheme MSTP(72)	n/a	n/a	n/a	n/a	0
FSR model MSTJ(41)	2	2	2	2	(<i>p_T - ordered</i>)
FSR Λ_{QCD} PARJ(81)	0.29	0.29	0.29	0.29	0.14
BR colour scheme MSTP(89)	n/a	n/a	n/a	n/a	1
BR composite z enhancement factor PARP(79)	n/a	n/a	n/a	n/a	2
BR primordial k_T width $< k_T >$ PARP(91)	1.0	1.0	2.1	2.1	n/a
BR primordial k_T UV cutoff PARP(93)	5.0	5.0	15.0	15.0	5.0
CR model MSTP(95)	n/a	n/a	n/a	n/a	6
CR strength ξ_R PARP(78)	n/a	n/a	n/a	n/a	0.2
CR gg fraction (old model) PARP(85)	0.9	0.33	1.0	1.0	n/a

Table 3.1: PYTHIA parameters, divided into main categories: UE (underlying event), ISR (initial state radiation), FSR (final state radiation), BR (beam remnants), and CR (colour reconnections). The UE reference energy for all models is PARP(89)=1800GeV, and all dimensionful parameters are given in units of GeV.

MB Triggers - ATLAS - efficiencies and bias



Shown below for each coincidence logic:

Required counter threshold to suppress accidental rate to 1Hz

Corresponding trigger efficiency for each type of event

Coincidence Logic	Counter Threshold [pC]	Trigger Efficiency			
		900 GeV		14 TeV	
		Non-Diff	Single/ Double Diff.	Non-Diff	Single/ Double Diff.
1+1	10.75	98%	35%	99%	42%
2+2	7.75	97%	27%	98%	36%
3+3	6.25	94%	20%	97%	30%
4+4	5.25	90%	15%	95%	24%
5+5	4.50	85%	10%	92%	19%
6+6	4.00	79%	7%	89%	14%

Data sets and triggers



Samples	N_RECO	σ [pb]
MB	1500000	5,52E+010
QCD 20-30	97000	6,32E+008
QCD 30-50	100000	1,63E+008
QCD 50-80	120000	2,16E+007
QCD 80-120	40000	3,08E+006
QCD 120-170	28000	4,94E+005
QCD 170-230	30000	1.01E+005
QCD 230-300	54000	2,45E+004

MB Trigger definition:
10+10 cells with $E > 1\text{GeV}$ in HF
[CMS AN 2007_017]

Reminder:
Spring07 GEN samples have
Pythia 6.227 with DWT.

Triggers “used”:

MB: Bandwidth 1 Hz

Jet20: Bandwidth 2.5 Hz

Jet60: Bandwidth 2.8 Hz

Jet120: Bandwidth 2.4 Hz

Often MC event weight $\gg 1$
[for 100pb^{-1} Feasibility studies]

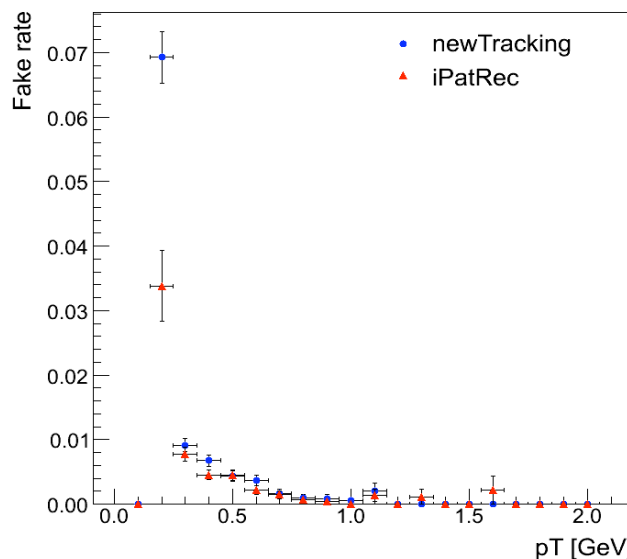
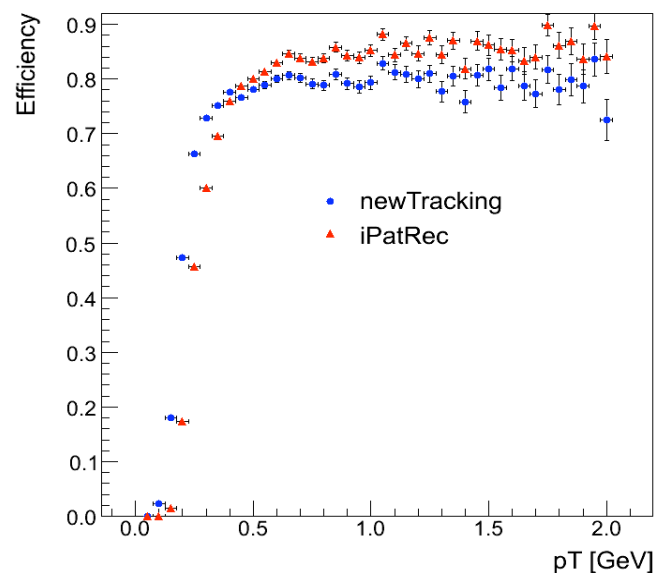
3 luminosity points considered:

1, 10 and 100pb^{-1}

ATLAS - central UE activity - Tools definition

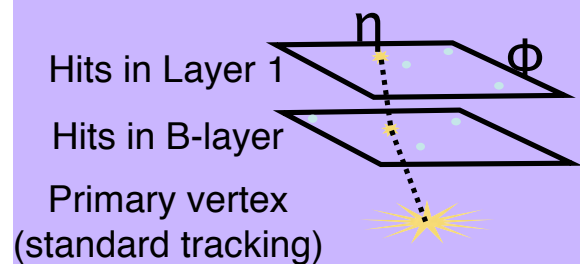


Standard tracking extended to low PT

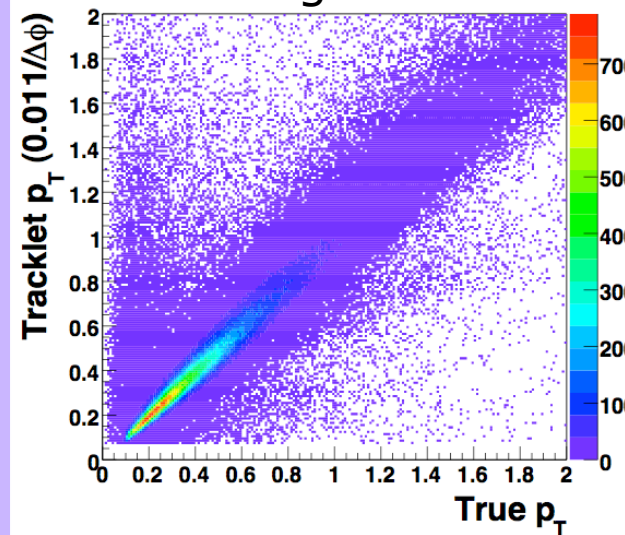
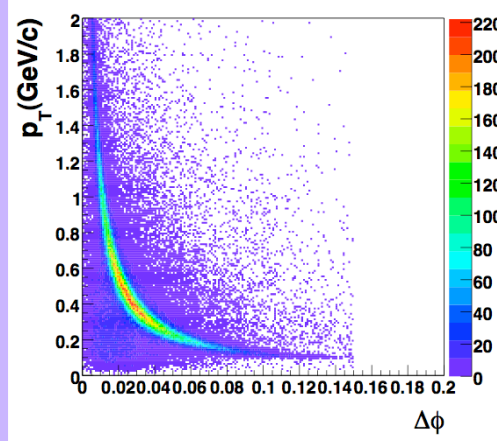


Tracklet method

Exploring methods with sensitivity at lower p_T with even fewer pixel-hits



Useful for 0-2 GeV PT range

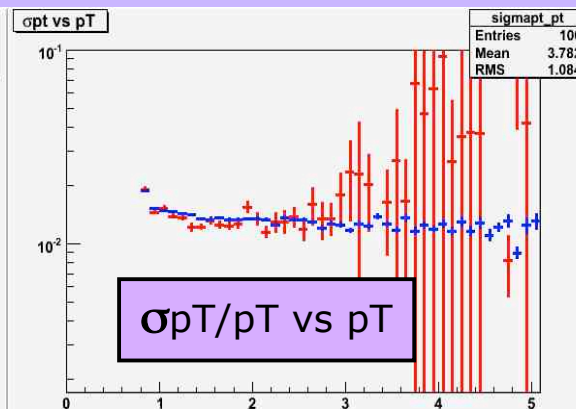
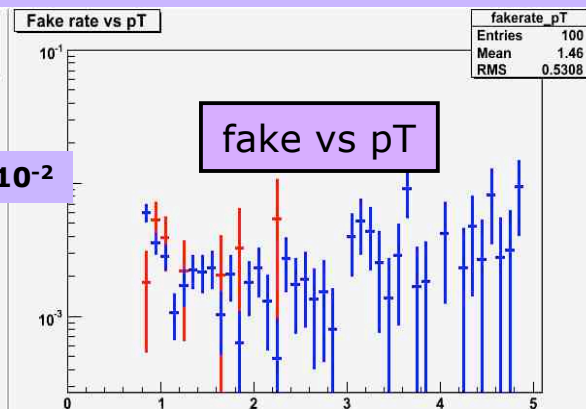
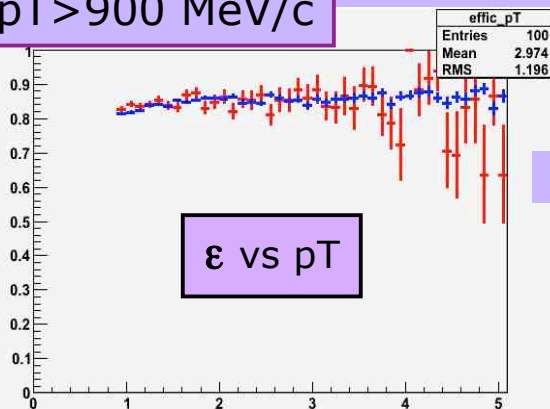


Tracking performances on **MB** and **QCD**

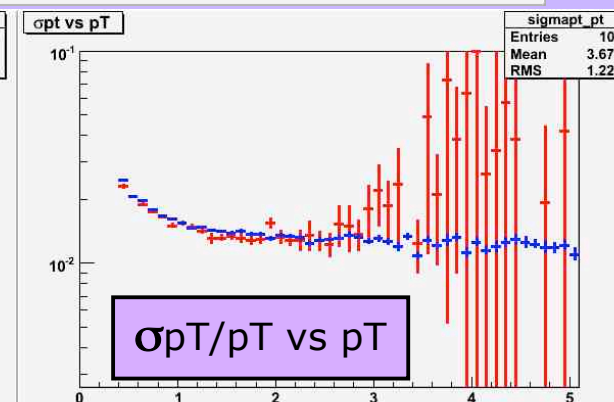
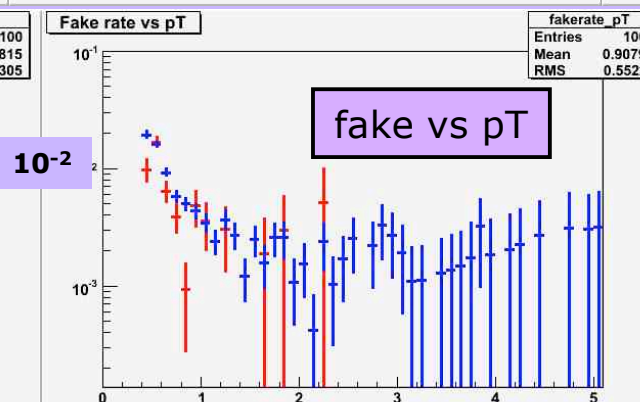
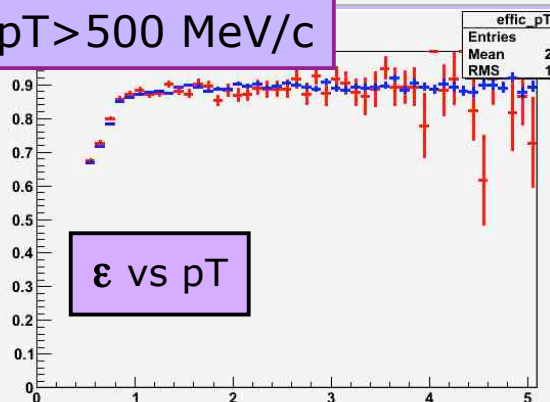


Association criteria based on 50% of shared hits

pT > 900 MeV/c



pT > 500 MeV/c



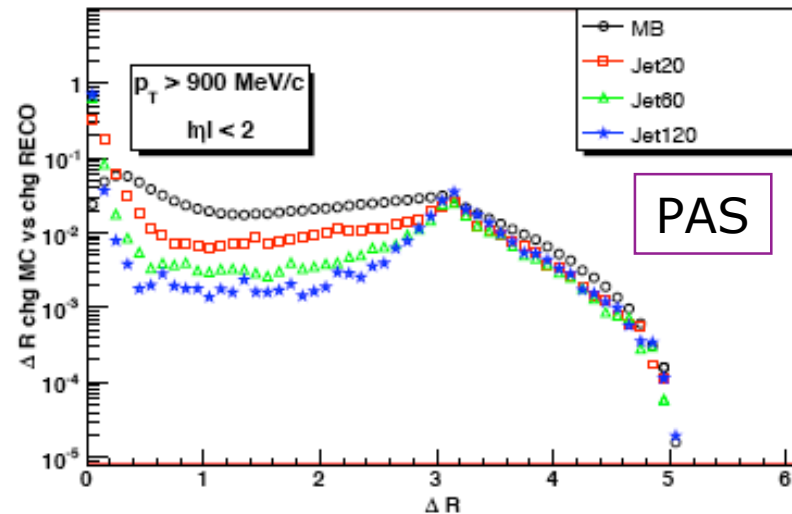
Tracking conf:

min $P_T \geq 0.5$
 $\chi^2/\text{ndof} \leq 5$
 $d_0 \leq 3.5$
 $z_0 \leq 30$
 $5 \leq \text{Hit} \leq 7$ & Lost = 0
 $\text{Hit} \geq 8$ & Lost = 1

- + 900 MeV is the standard CMSSW tracking
- + seeding and tracking from 500 MeV is possible with sufficient high efficiency (from 70% to 90%) and under control fakes (~2% from 500 MeV)
- + Tracking from 500 MeV is used to enhance discriminative power of the observables in the transverse region (see previous)

Used Tools

ΔR reco/calor



PT ratio reco/calor

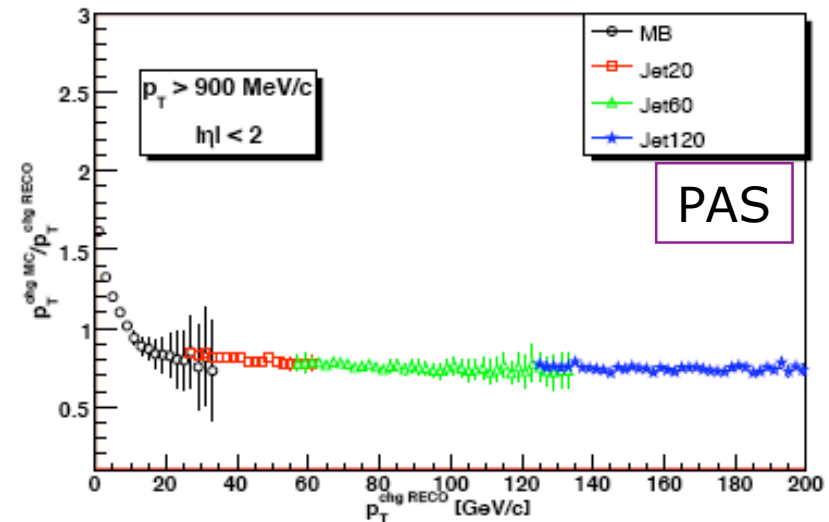
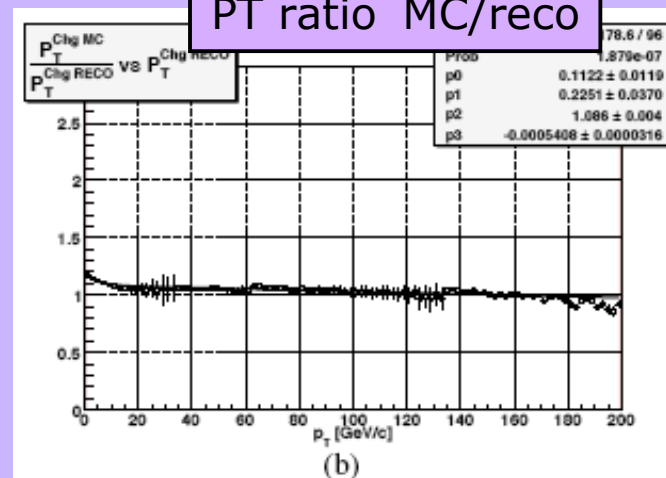


Figure 1: Charged Jet performances. *Left*: $\Delta R (= \sqrt{\Delta\phi + \Delta\eta})$ between the charged and the calorimetric leading jets. *Right*: charged jet response function ($P_T^{\text{charged}}/P_T^{\text{calorimetric}}$).

Charged jet instead of calorimetric:

- + access to low PT region
- + intrinsically free from pile up
- + better control of systematic effects at startup

PT ratio MC/reco



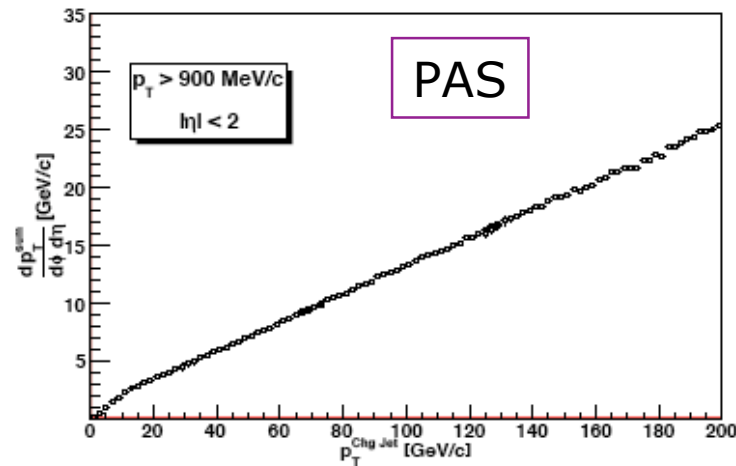
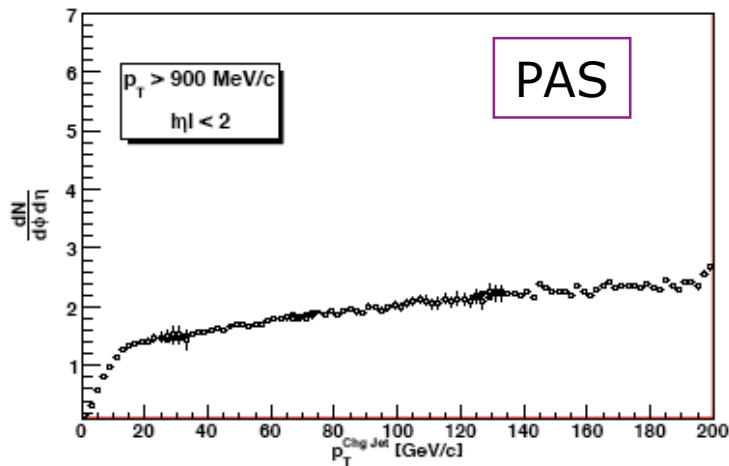
Absolute energy calibration of the leading charged jet reported with the fit adopted in the correction procedure

Feasibility@LHC - the **toward** and **transverse** region - 10 pb⁻¹



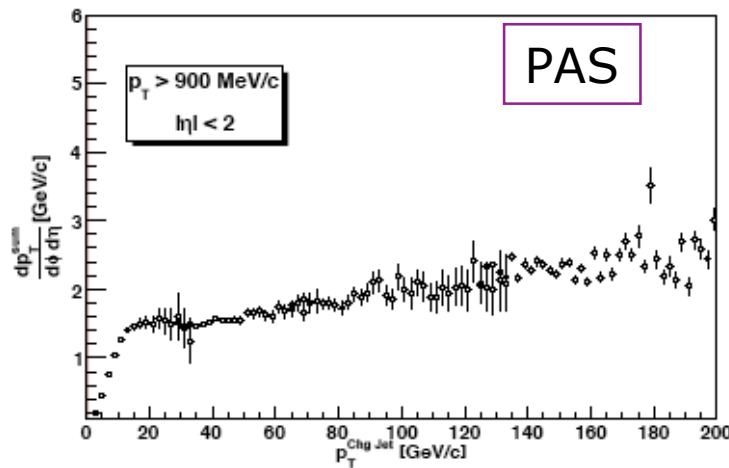
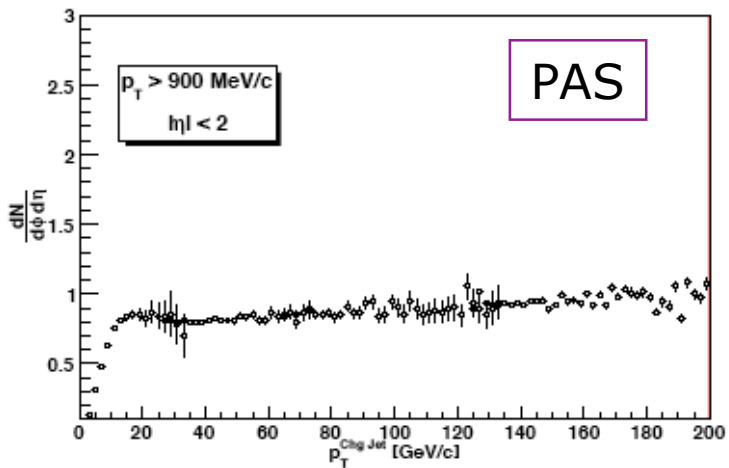
$dN/d\eta d\phi$ vs PT_{chg}

$d(PT_{sum})/d\eta d\phi$ vs PT_{chg}



toward

Figure 3: Toward region. Uncorrected distribution. Density of charged particles $dN/d\eta d\phi$ (left) and average charged $\sum p_t$ density (right) versus the transverse momentum of the leading charged particle jet.



transverse

Figure 4: Same as figure 3 but in the transverse region.

Feasibility@LHC - correction functions - 10 pb-1

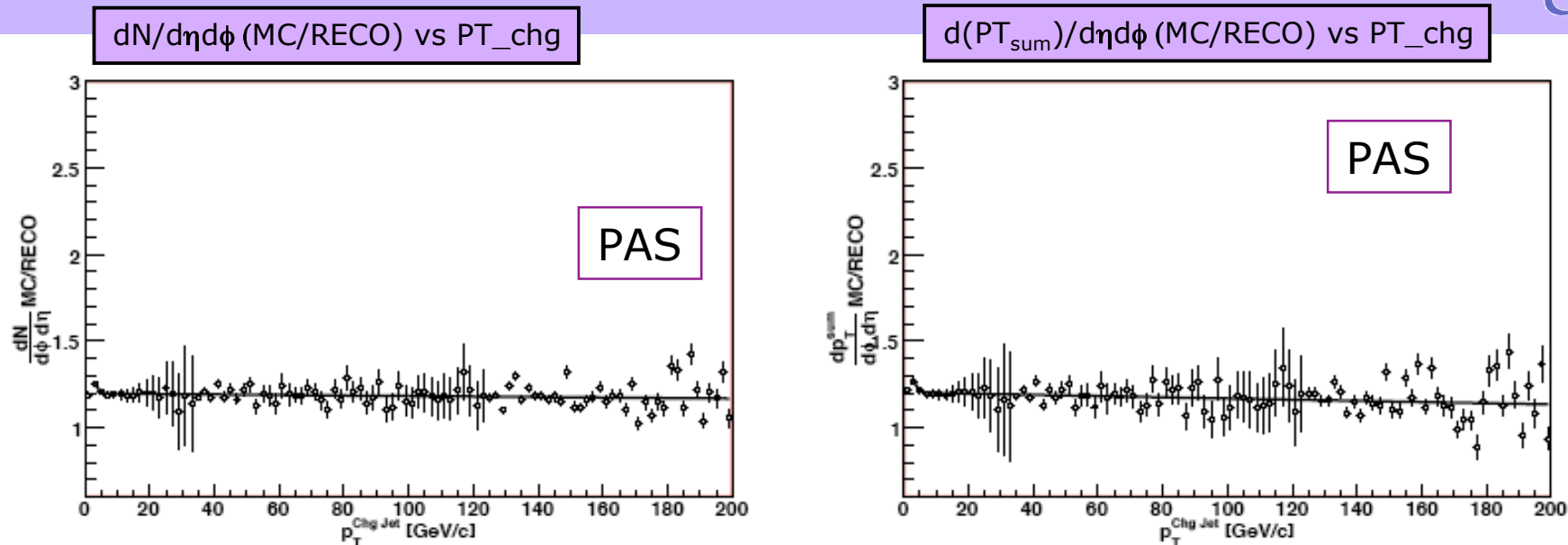


Figure 5: Ratio between generator level DWT expectations and reconstructed data for $dN/d\eta d\phi$ (left) and average charged $\sum p_t$ density (right). The corrections as a function of the charged jet transverse momentum are then obtained from the fit.

Correction/calibration procedure a la CDF [PRD 65 (2003) 092002]:

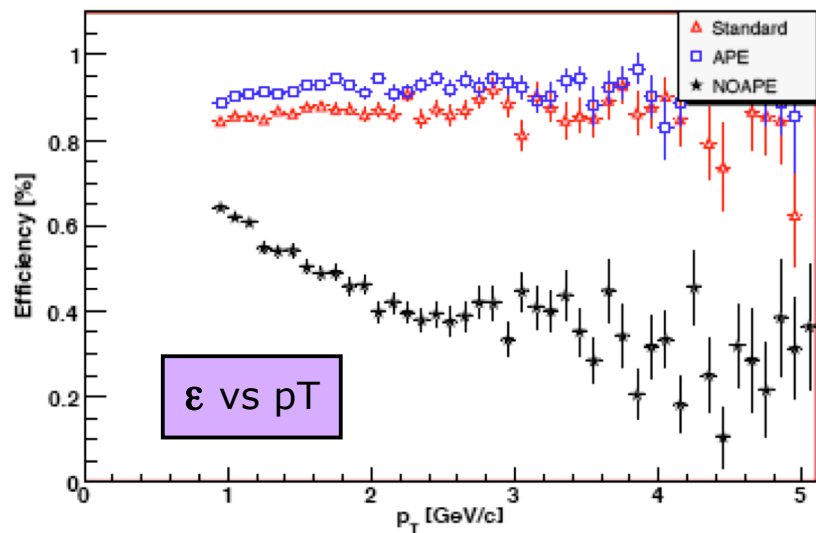
- + Use the PT_chg calibration function vs PT_chg (RECO)
- + Response function vs PT_jet (MC) applied on a event by event basis

[At CDF it provides the same performances of the particle level correction]

Feasibility@LHC - startup scenario - 1 pb-1

Tracking at startup

Ideal aligned detector
Misaligned
Misaligned + APE



Survey + LAS

- Pixel Barrel
 - $dX = dY = 225 \mu\text{m}$
 - $dZ = 337 \mu\text{m}$
 - $\alpha = \beta = \gamma = 10 \mu\text{rad}$
- TIB Layers
 - $dX = dY = dZ = 750 \mu\text{m}$
 - $\alpha = \beta = \gamma = 488 \mu\text{rad}$
- TIB Rods
 - $du = dv = dw = 450 \mu\text{m}$
 - $\alpha = \beta = \gamma = 293 \mu\text{rad}$
- TOB Rods
 - $du = dv = dw = 100 \mu\text{m}$
 - $\alpha = \beta = \gamma = 40 \mu\text{rad}$

APE = Alignment Position Error

Performances on MB sample

- + Efficiency and fake performances are recovered using APE (additional error to the hit taking into account alignment precision)
- + transverse momentum resolution partially recovered
- + Higher efficiency with APE is due to the MS effect recovered by a larger search window

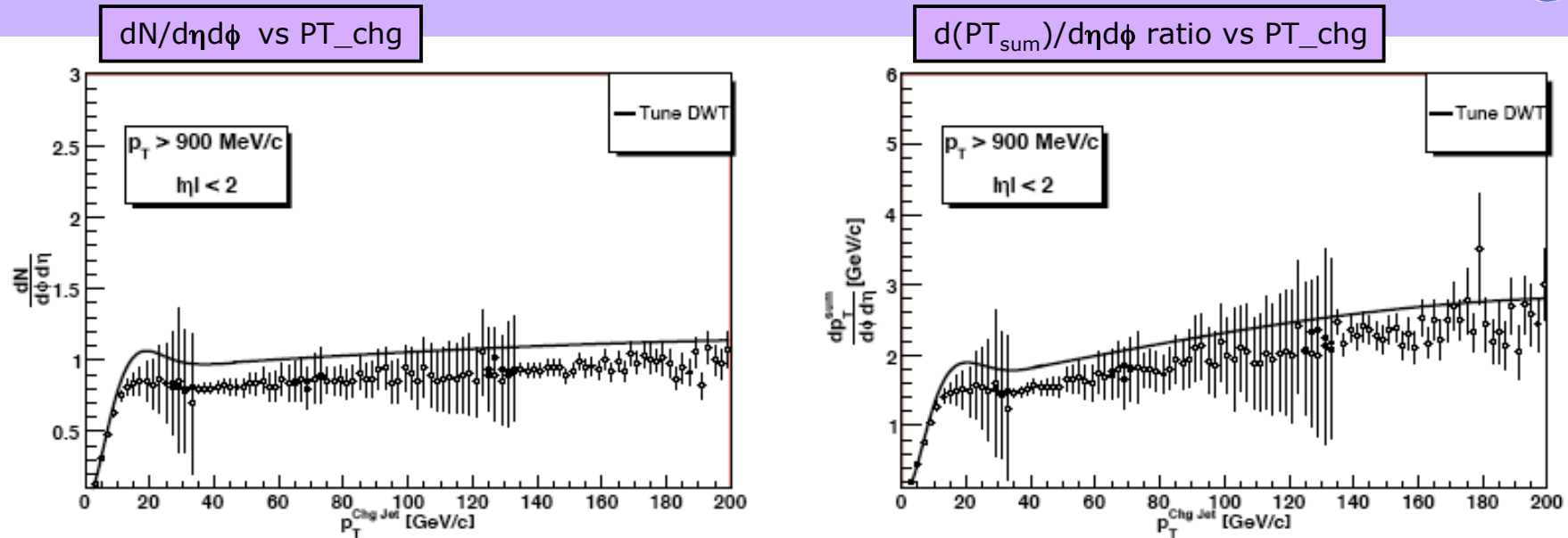


Figure 11: Transverse region. Uncorrected distribution at startup using APE. Density of charged particles $dN/d\eta d\phi$ (left) and the average charged $dpt/d\eta d\phi$ density (right) versus the transverse momentum of the leading charged jet.

Uncorrected data.

Recovering the tracking at startup assure the possibility to build UE observables from first days of data taking and start building correction functions

Feasibility@LHC - systematics



Systematics from adopted triggers and detector related inefficiency (bad channels, dishomogeneities...) are not considered, second order effect

- $\sigma_{\epsilon} = \sigma_{fake} = 1\% (2\%)$
- $\sigma_{Pt} = 2\% (4\%)$

← considered uncertainties from tracking after 100 pb⁻¹
(x2) -> "safety" factor

$$\sigma_{dN} = \sigma_{\epsilon}$$

$$\sigma_{dPT}^2 = dN^2 * \sigma_{Pt}^2 + dPT^2 * \sigma_{\epsilon}^2$$

$$\sigma_{dN_{corr}}^2 = dN^2 * (\sigma_{\epsilon}^2 + \sigma_{fake}^2) + (1 - \epsilon + fake)^2 * \sigma_{dN}^2$$

$$\sigma_{dPT_{corr}}^2 = dN_{corr}^2 * \sigma_{Pt}^2 + dPT_{corr}^2 * (\sigma_{dN_{corr}} / dN_{corr})^2$$

Eg: fixing a point: PT_chg = 100 GeV/c, 100 pb⁻¹

$$dN/d\eta d\phi = 1.07 \pm 0.02 \text{ (statistical)} \pm 0.02 \text{ (systematic)}$$

pT > 900 MeV/c

$$PT_{sum} = 2.25 \pm 0.06 \text{ (statistical)} \pm 0.07 \text{ (systematic)}$$

$$dN/d\eta d\phi = 1.75 \pm 0.03 \text{ (statistical)} \pm 0.03 \text{ (systematic)}$$

pT > 500 MeV/c

$$PT_{sum} = 2.85 \pm 0.08 \text{ (statistical)} \pm 0.12 \text{ (systematic)}$$