

Gauge-Higgs Unification

Theory and Phenomenological Consequences at the LHC

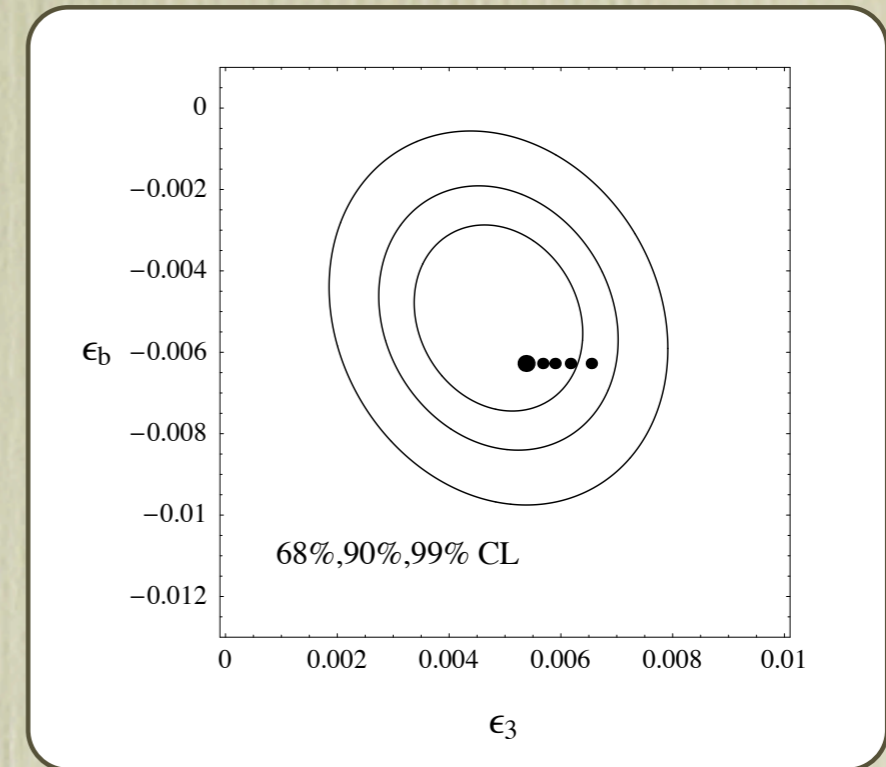
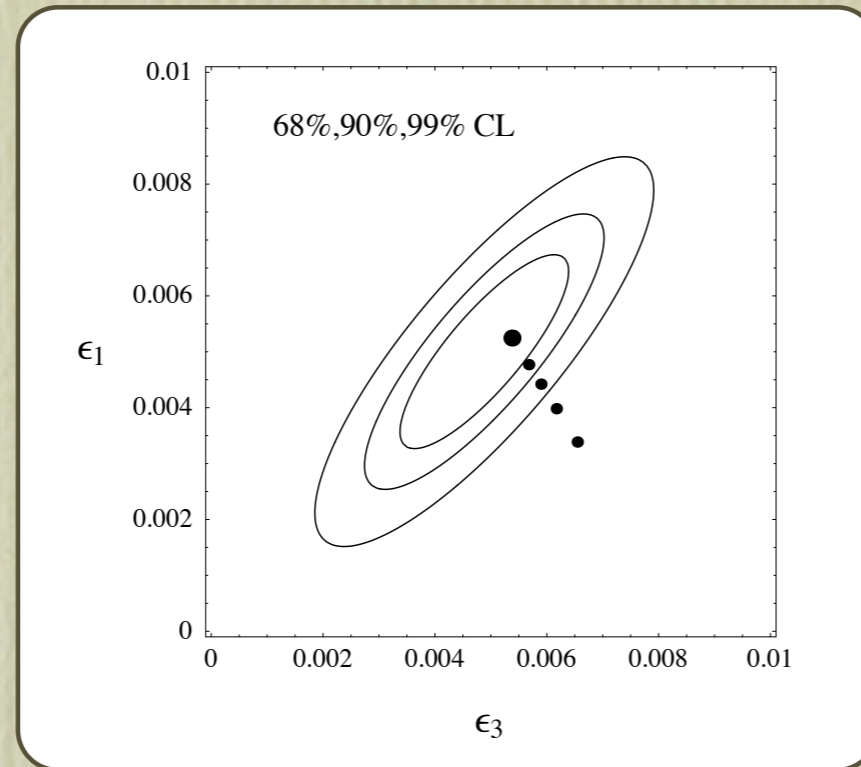
Roberto Contino

Università di Roma La Sapienza



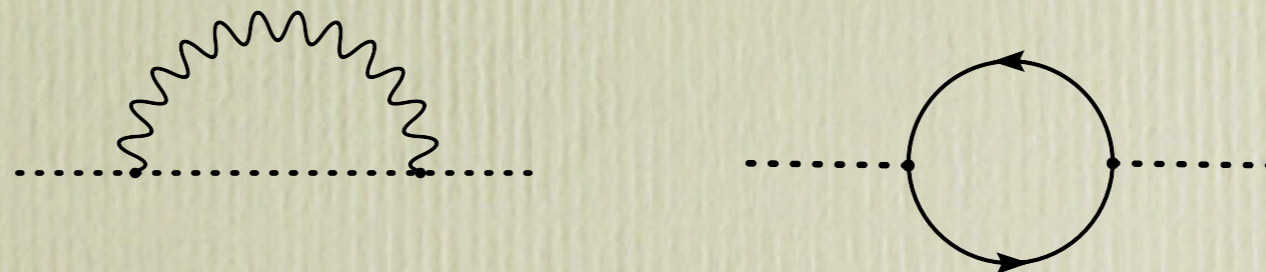
Fact #1 :

LEP and SLD precision data strongly suggest the existence of a **light** Higgs boson, $m_H \sim 100\text{-}300 \text{ GeV}$



Fact #2 :

The instability against radiative correction makes a light (elementary) scalar in the low-energy spectrum highly **unnatural** unless a **symmetry protection** is at work



$$\delta m_h^2 = \left[6 y_t^2 - \frac{3}{4} (3 g_2^2 + g_1^2) - 6 \lambda_4 \right] \frac{\Lambda^2}{8\pi^2}$$

Two examples of symmetry protections



mass of **fermions** and **gauge bosons** are UV-stable:
each protected by a symmetry

Strategy:

relating the Higgs boson to fermions or gauge fields
to acquire their symmetry protection

Two examples of symmetry protections

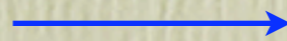


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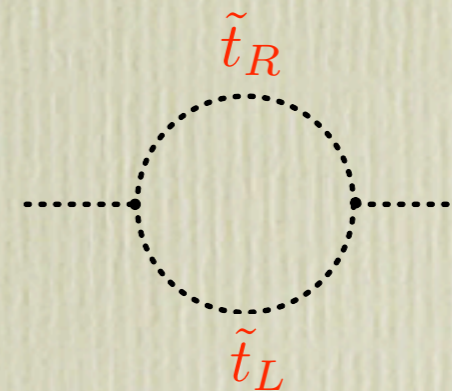
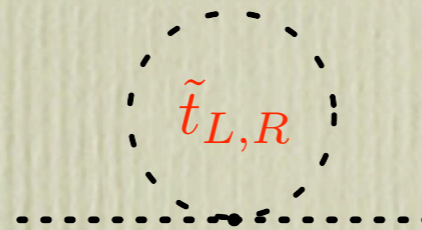
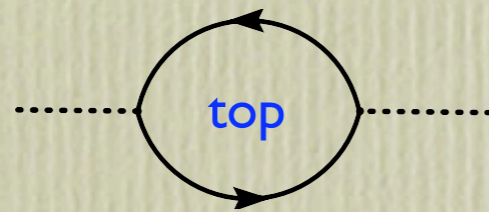
CHIRAL SYMMETRY



SUSY

$$h \subset \begin{pmatrix} \tilde{h} \\ h \end{pmatrix}$$

(fermion protection)



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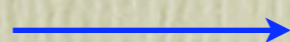


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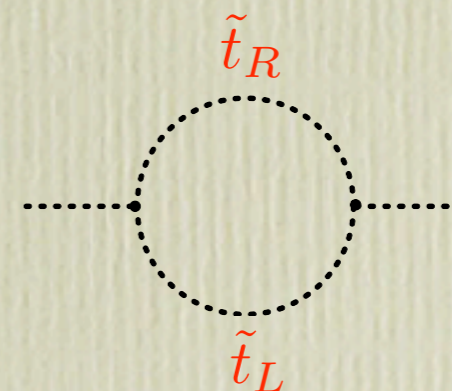
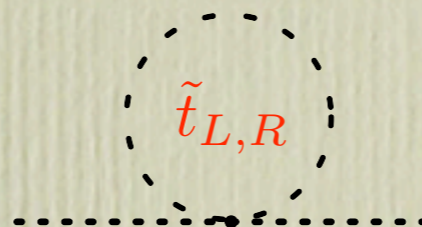
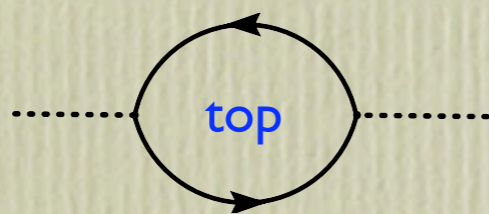
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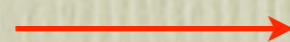
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GAUGE SYMMETRY

(gauge protection)



**GAUGE-HIGGS
UNIFICATION**

$$h = A_5$$

[requires extra dimensions]

Quick introduction to Gauge-Higgs unification

- Suppose a fifth extra spatial dimension exists: $A_M = \{A_\mu, A_5\}$

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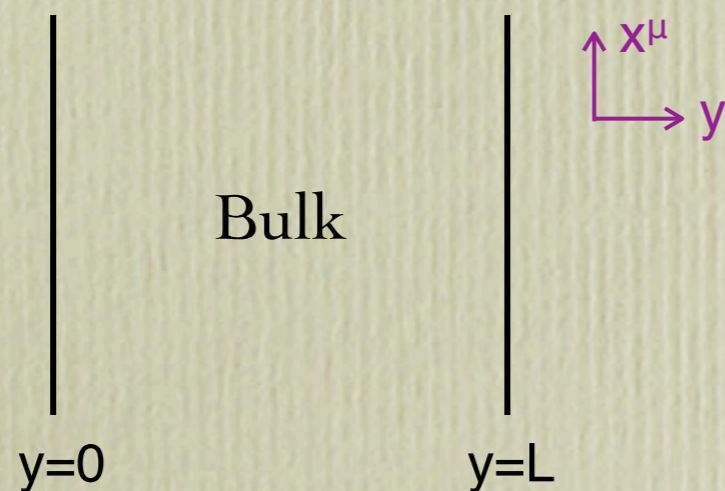
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Consider a **segment** :

Any field propagating into the 5th dimension can be decomposed in **Fourier armonics**



$$\Phi(x, y) = \sum_n \zeta_n(y) \phi^{(n)}(x)$$

and must satisfy definite **boundary conditions** :

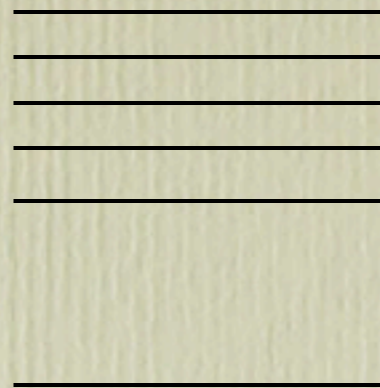
$$\partial_5 \Phi(x, y_i) = 0 \quad \text{Neumann} \quad (+)$$

$$\Phi(x, y_i) = 0 \quad \text{Dirichlet} \quad (-)$$

Each Fourier mode behaves like a 4D field

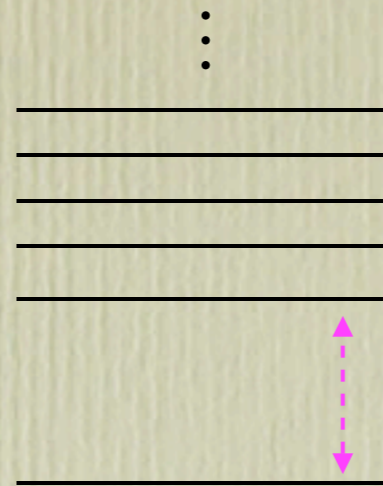
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⋮



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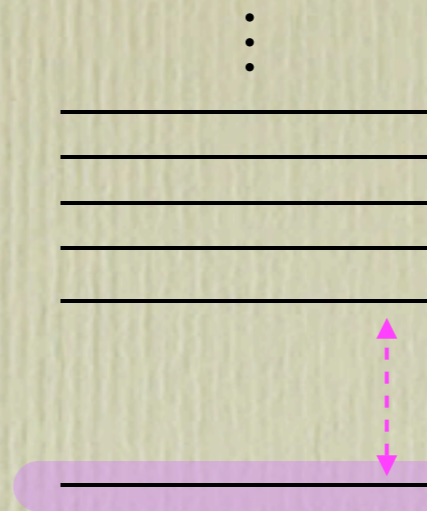
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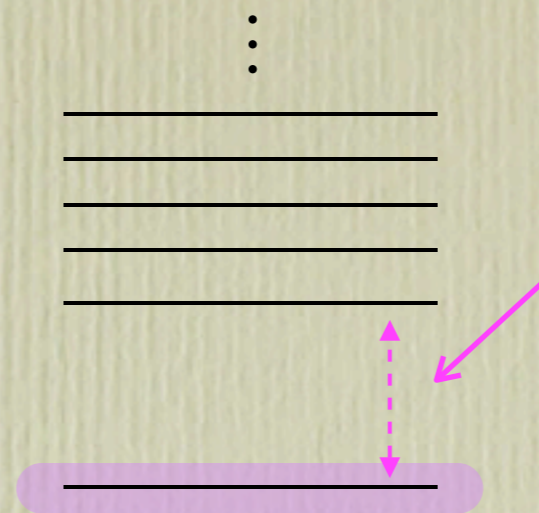
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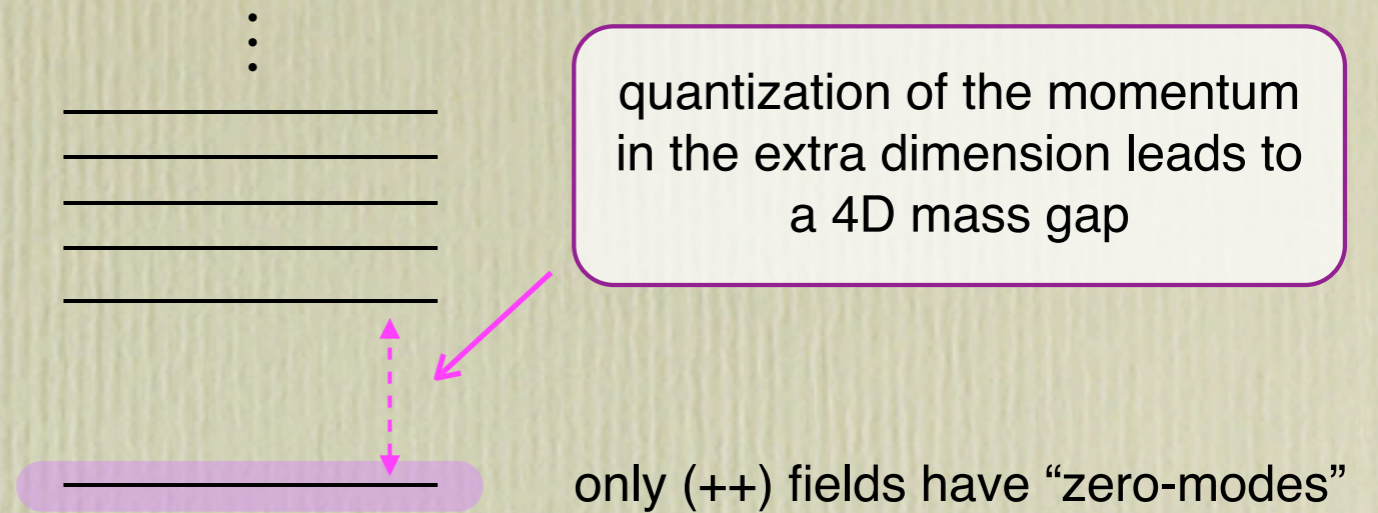
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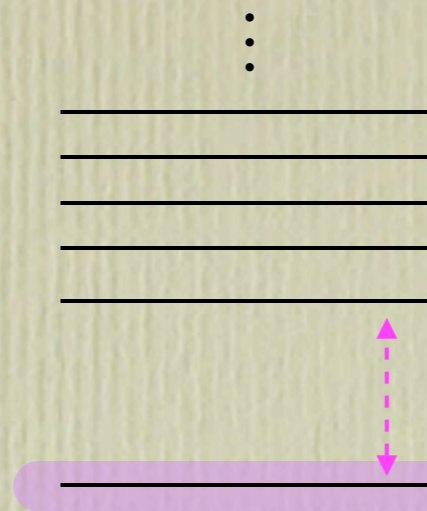
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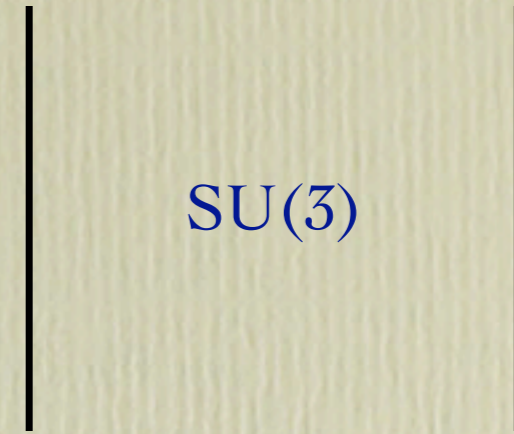
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1. Fermions in 5D are NOT chiral

2. A_5 , being a gauge field, transforms as an adjoint representation \rightarrow not an $SU(2)$ doublet !

Consider for example $SU(3)$ in the bulk
with the following boundary conditions:



$$A_{\mu}^a(++), \quad T^a \in \text{Alg}\{SU(2) \times U(1)\}$$

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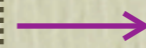


$$D_M = \partial_M + iA_M$$

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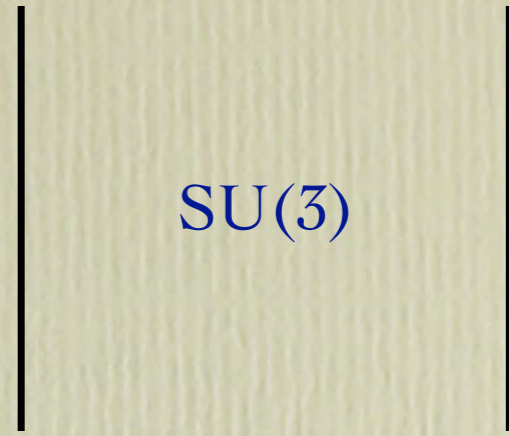
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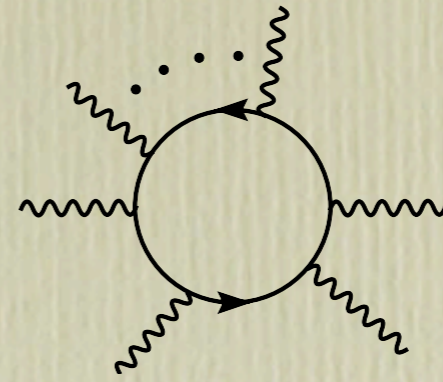
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The Yukawa coupling between doublet and singlet originates from the covariant derivative :

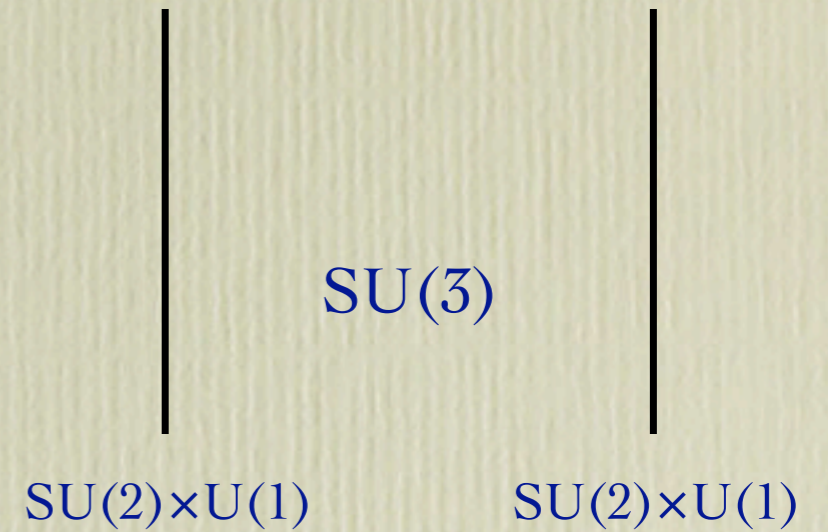
$$\bar{\Psi} i\Gamma^M (\partial_M - iA_M) \Psi \supset \bar{\Psi}_L \gamma^5 T^{\hat{a}} \Psi_R A_5^{\hat{a}} + h.c.$$

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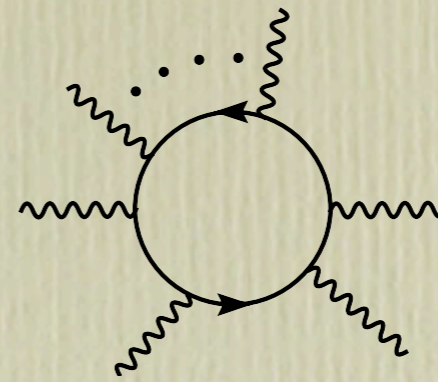
The Higgs potential at 1-loop



- The 5D gauge symmetry forbids a potential for A_5 but it is **globally broken** by the boundary conditions

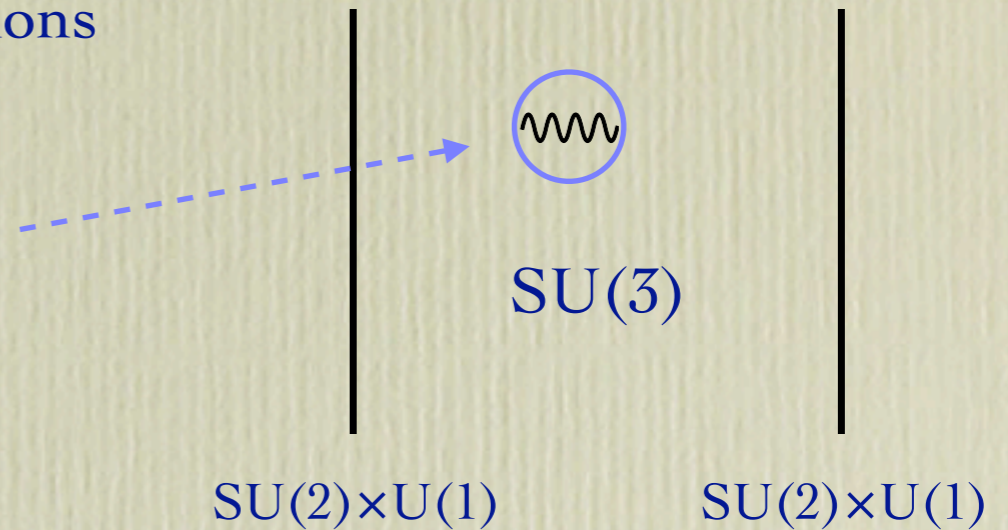



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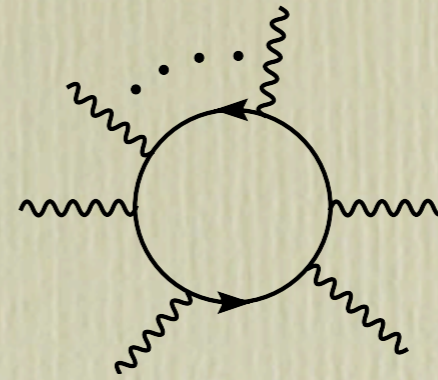
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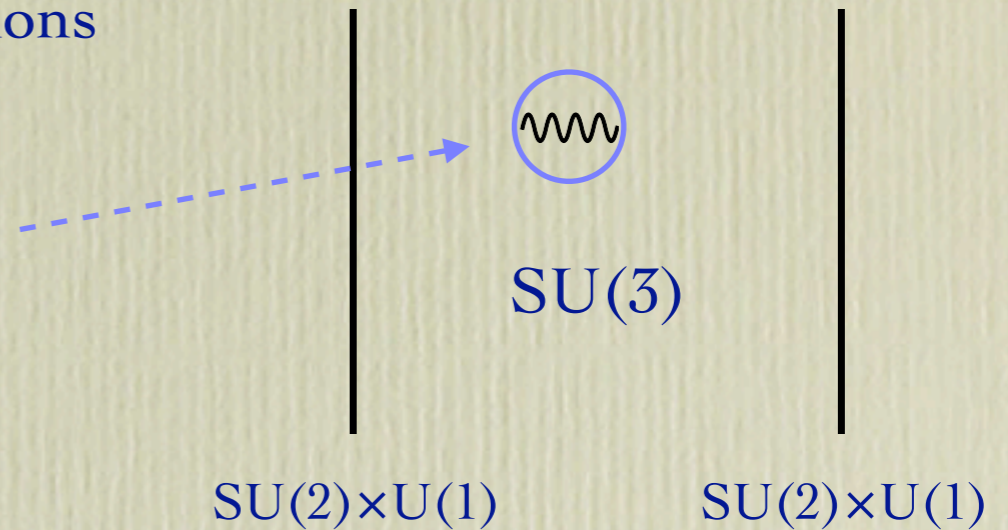
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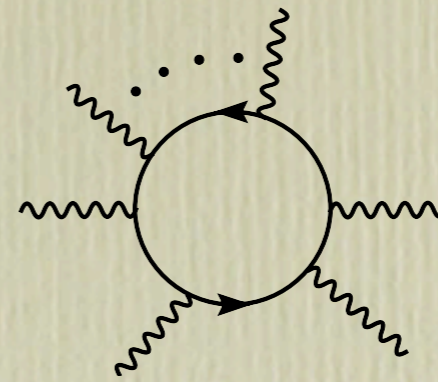


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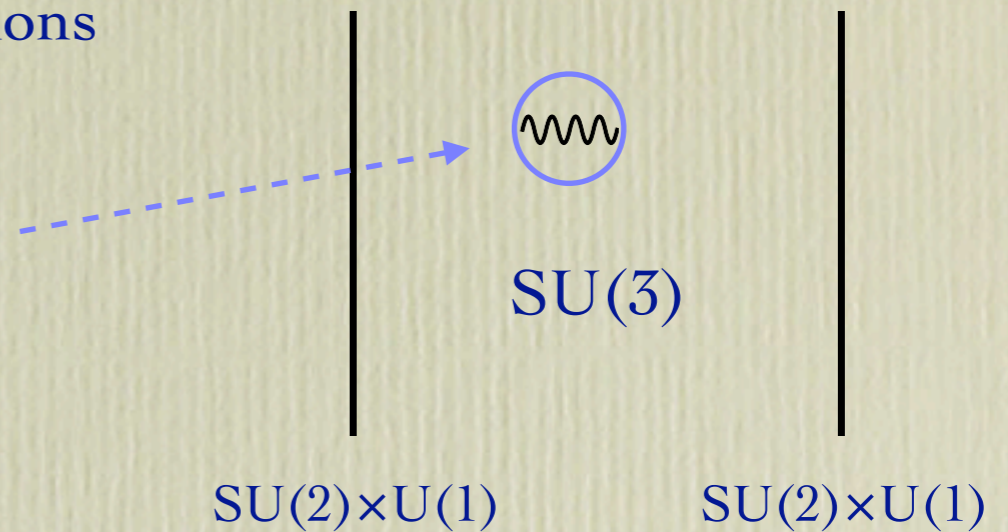
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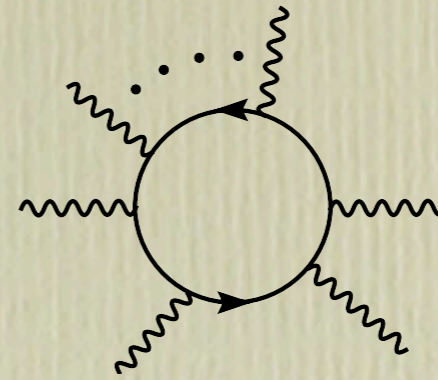
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That is:

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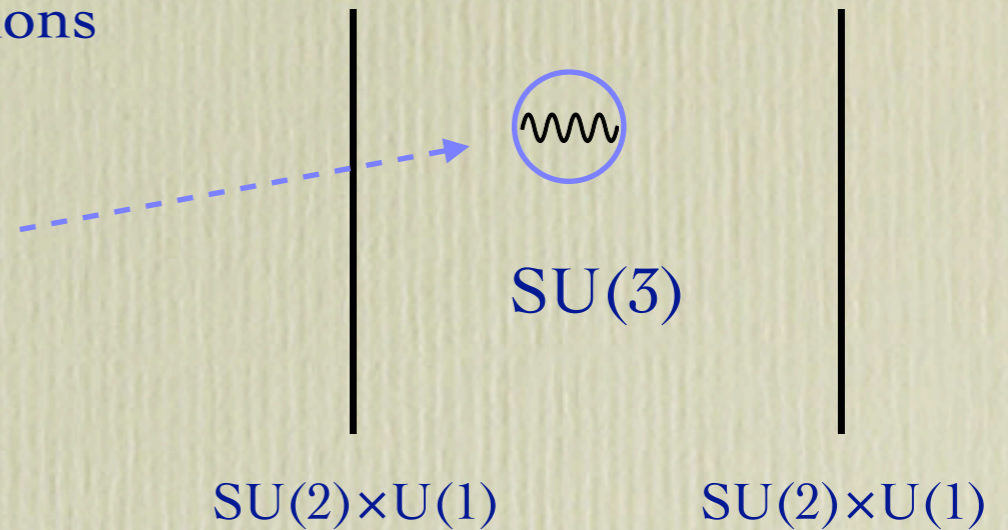
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- From a 4D point of view the quadratic divergence in the top loop is **canceled** by the tower of **Kaluza-Klein modes**

$$\text{Diagram with top loop } t + \sum_n \text{Diagram with top loop } t^{(n)} = \text{finite}$$

The diagram illustrates the cancellation of a quadratic divergence in a top loop. On the left, a Feynman diagram shows a top quark loop (a circle with two arrows) connected to two external wavy lines. This is followed by a plus sign and a summation symbol over n . To the right of the summation is another Feynman diagram, identical in structure but with the top quark loop labeled $t^{(n)}$. This is followed by an equals sign and the word "finite".

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The cancellation happens among states with the same spin !

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Heavy Top partners come in **complete multiplets** of the bulk gauge symmetry

For a symmetry larger than SU(3) there can be fermions with **exotic quantum numbers**

For example: $SO(5) \rightarrow SO(4)$

$$\Psi (= 5 \text{ of } SO(5)) = \left[\begin{array}{l} \mathbf{2}_{7/6} = \begin{pmatrix} T_{5/3} \\ T \end{pmatrix} \\ \mathbf{2}_{1/6} = \begin{pmatrix} t \\ b \end{pmatrix} \\ \mathbf{1}_{2/3} = t \end{array} \right]$$

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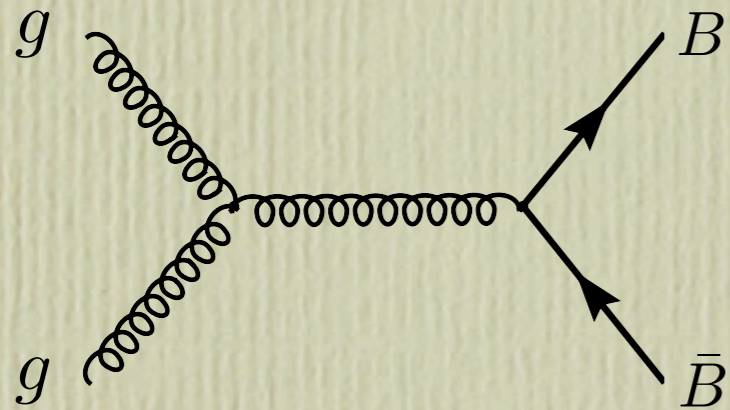
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electric charge $+5/3$

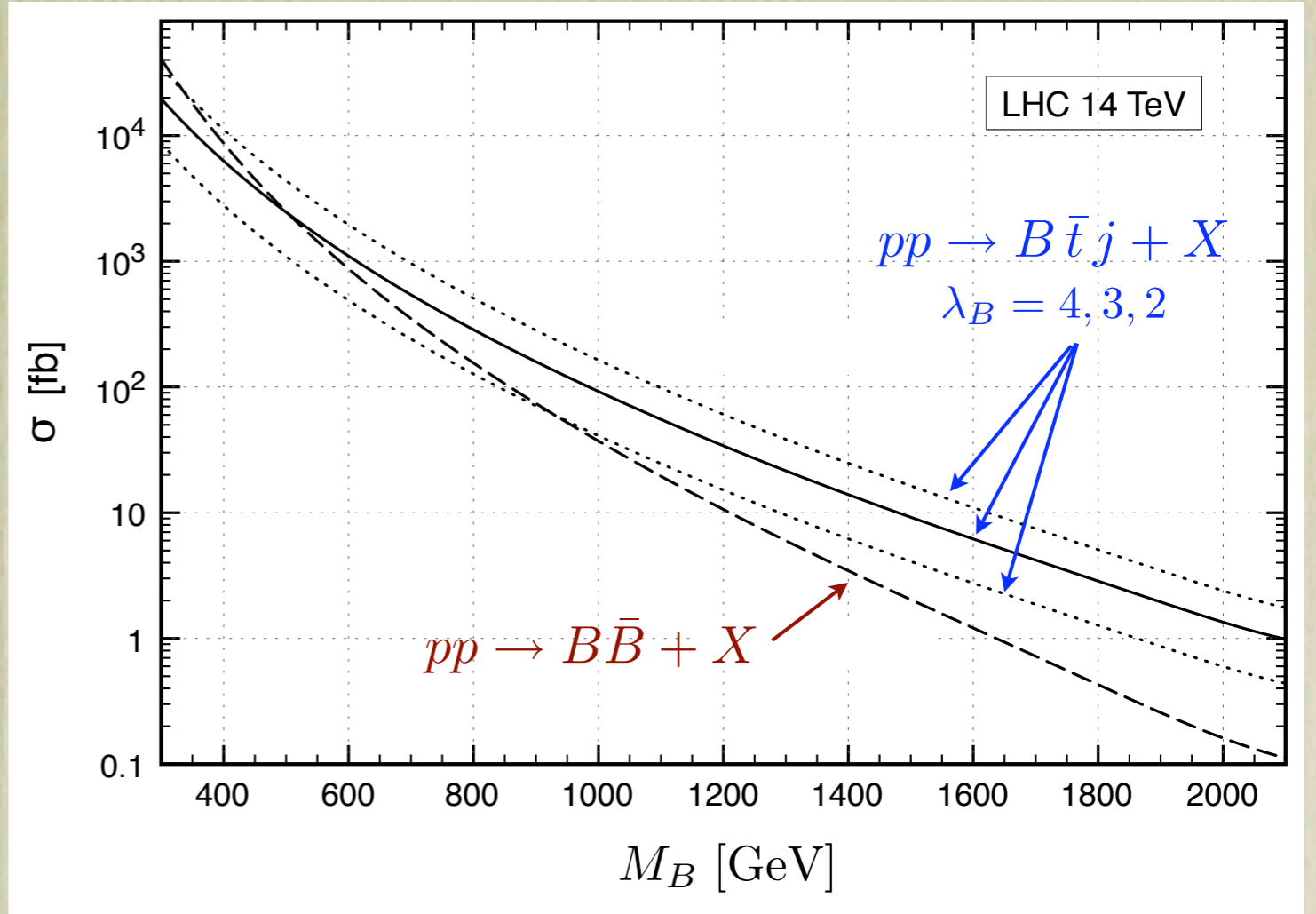
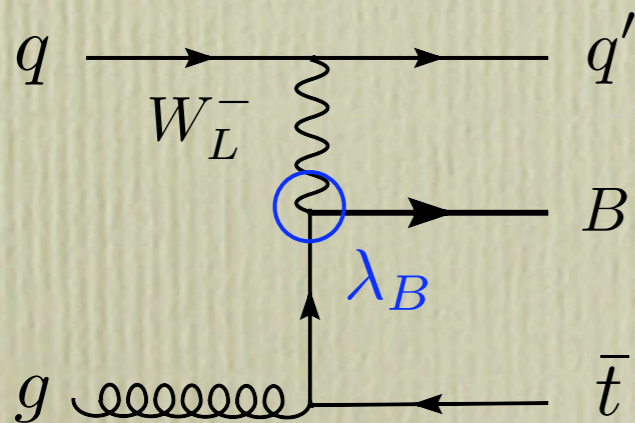
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Discovering the top partners at the LHC

Pair production

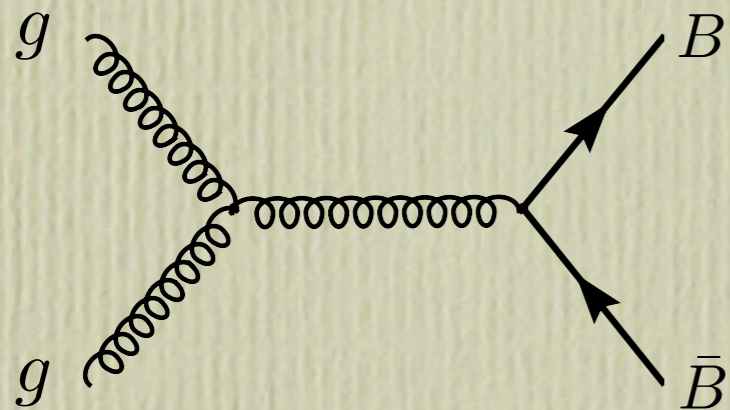


Single production

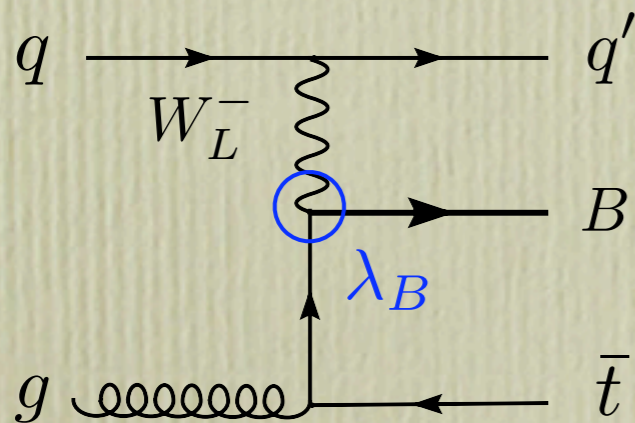


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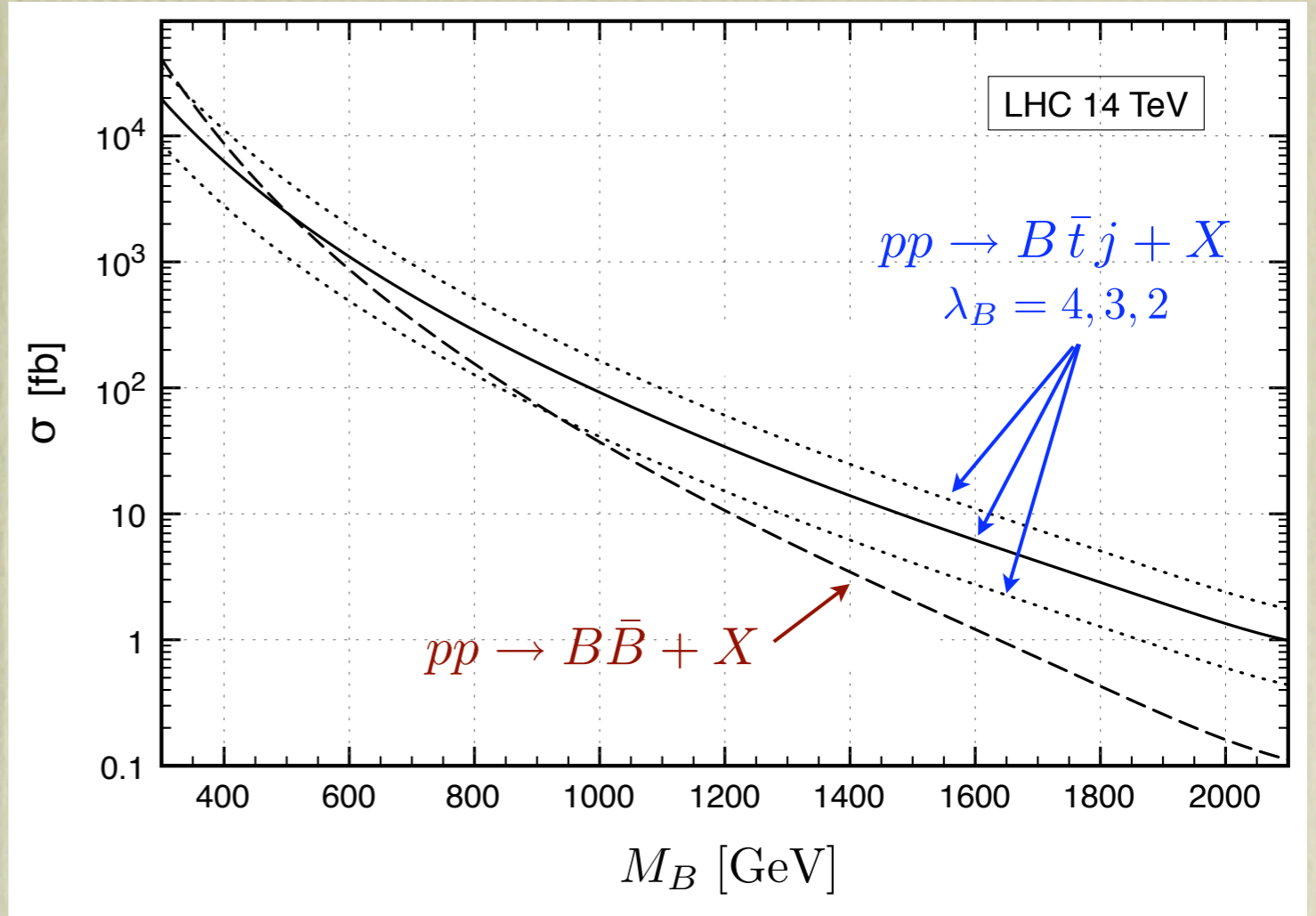
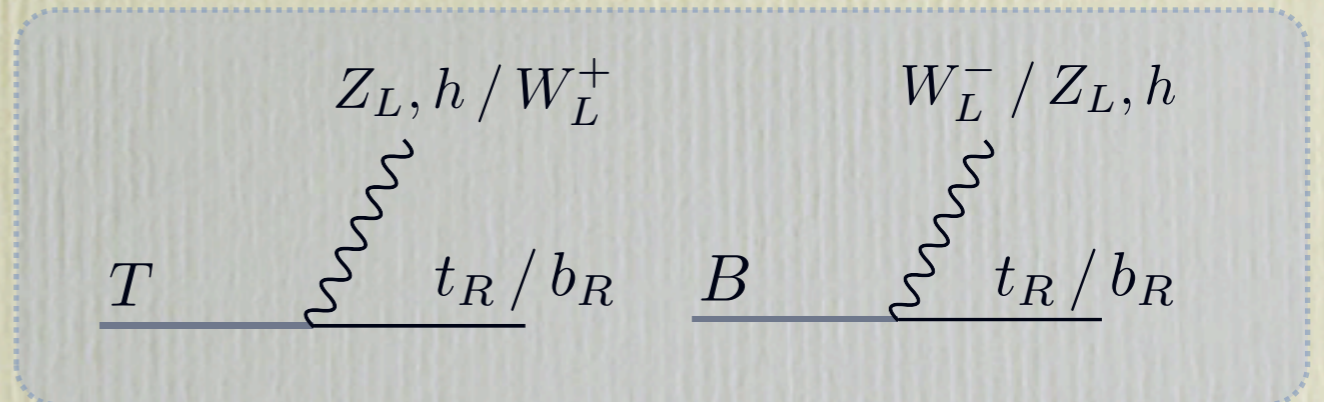
Pair production



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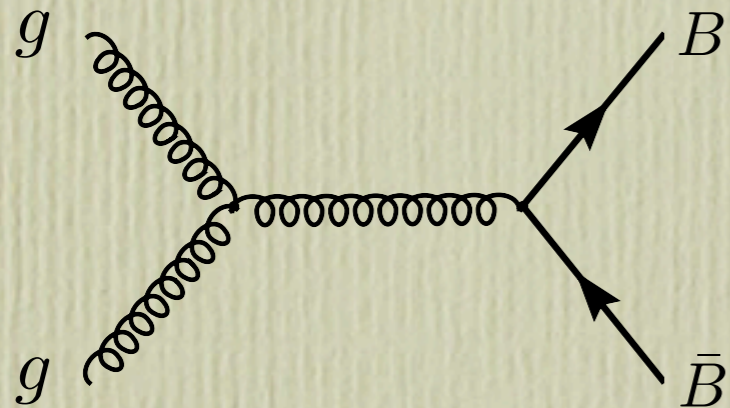


Decay modes

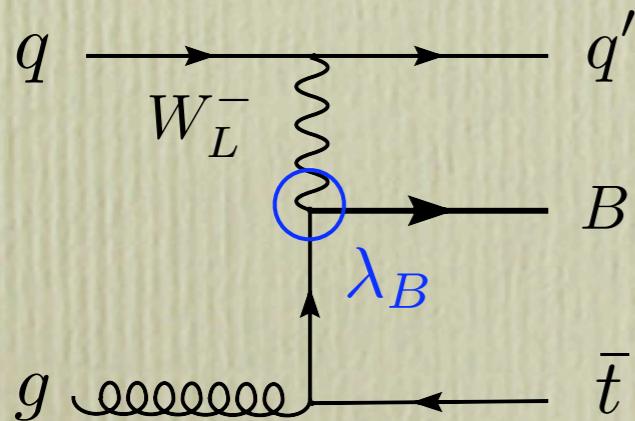


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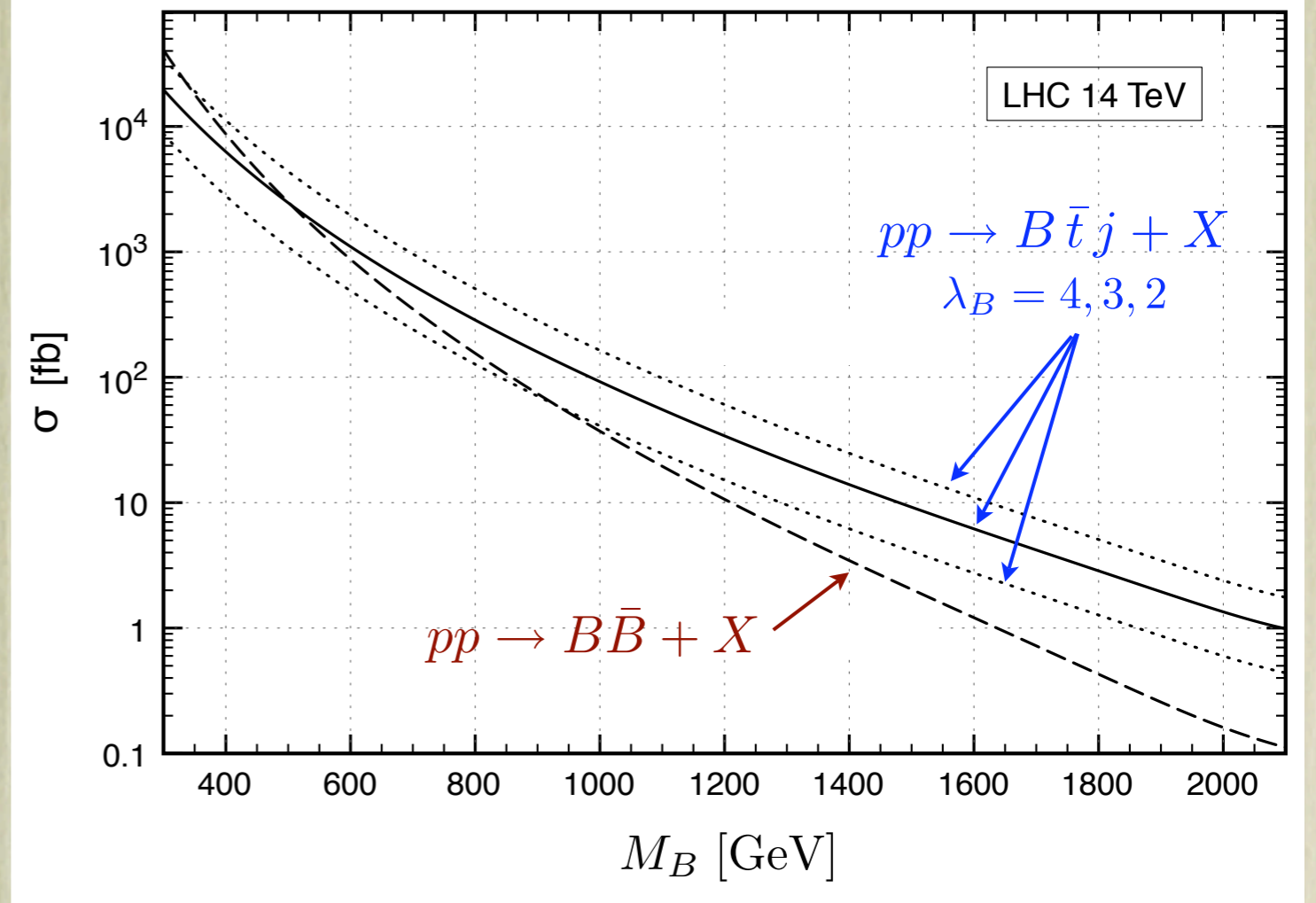
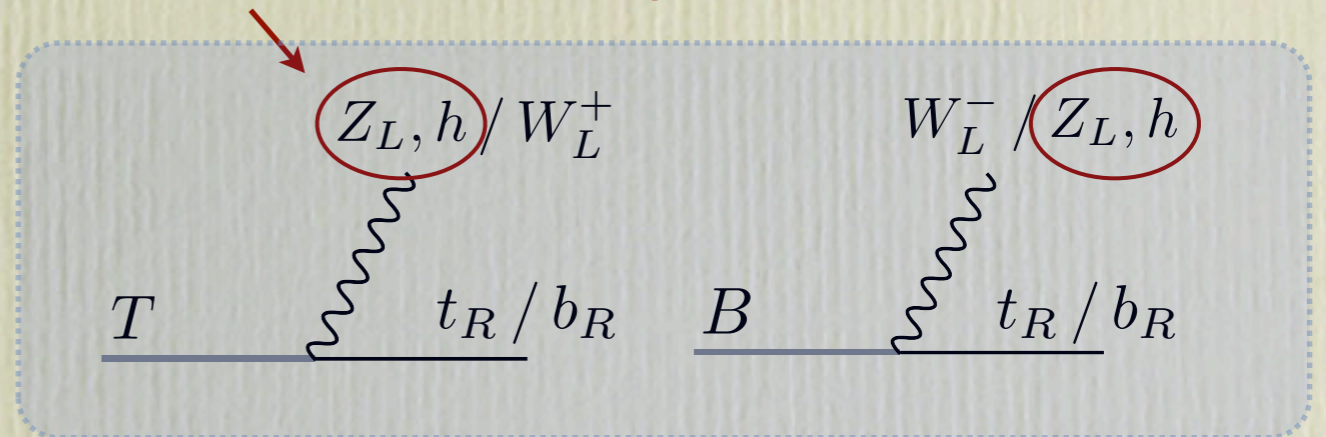


Single production



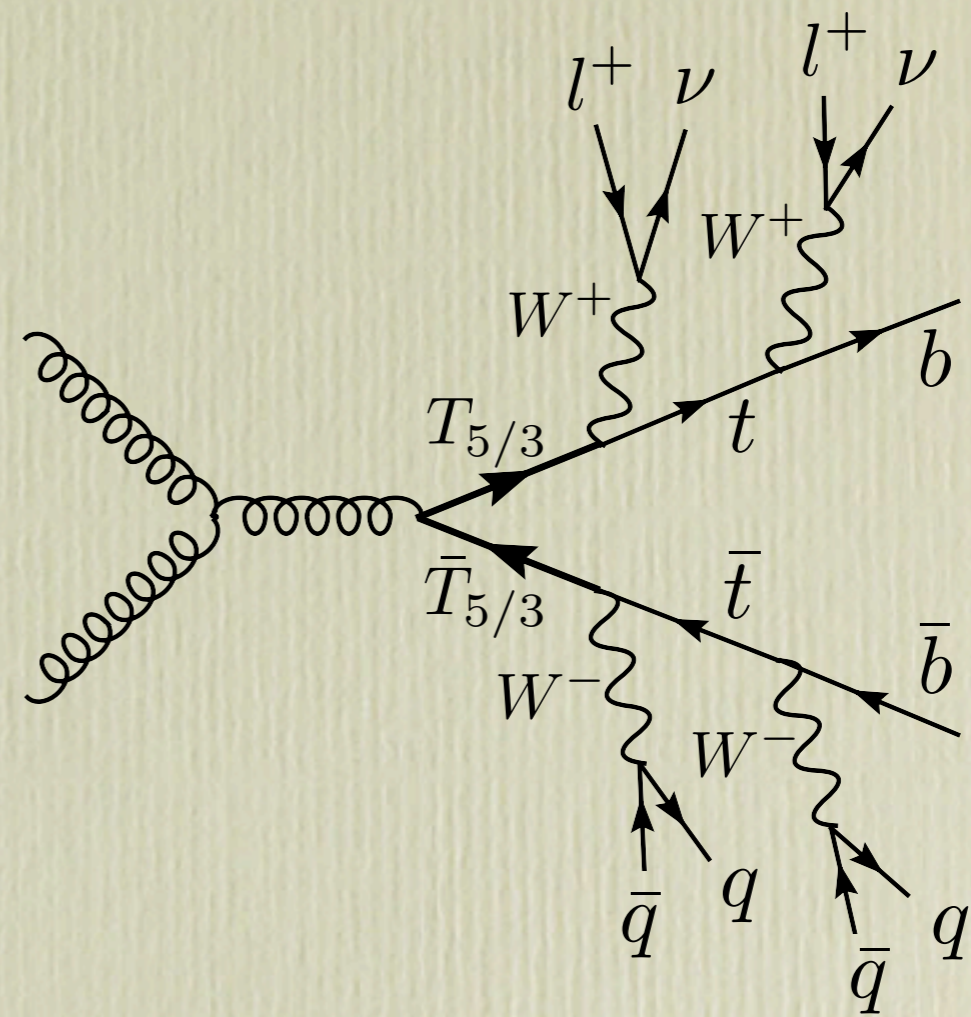
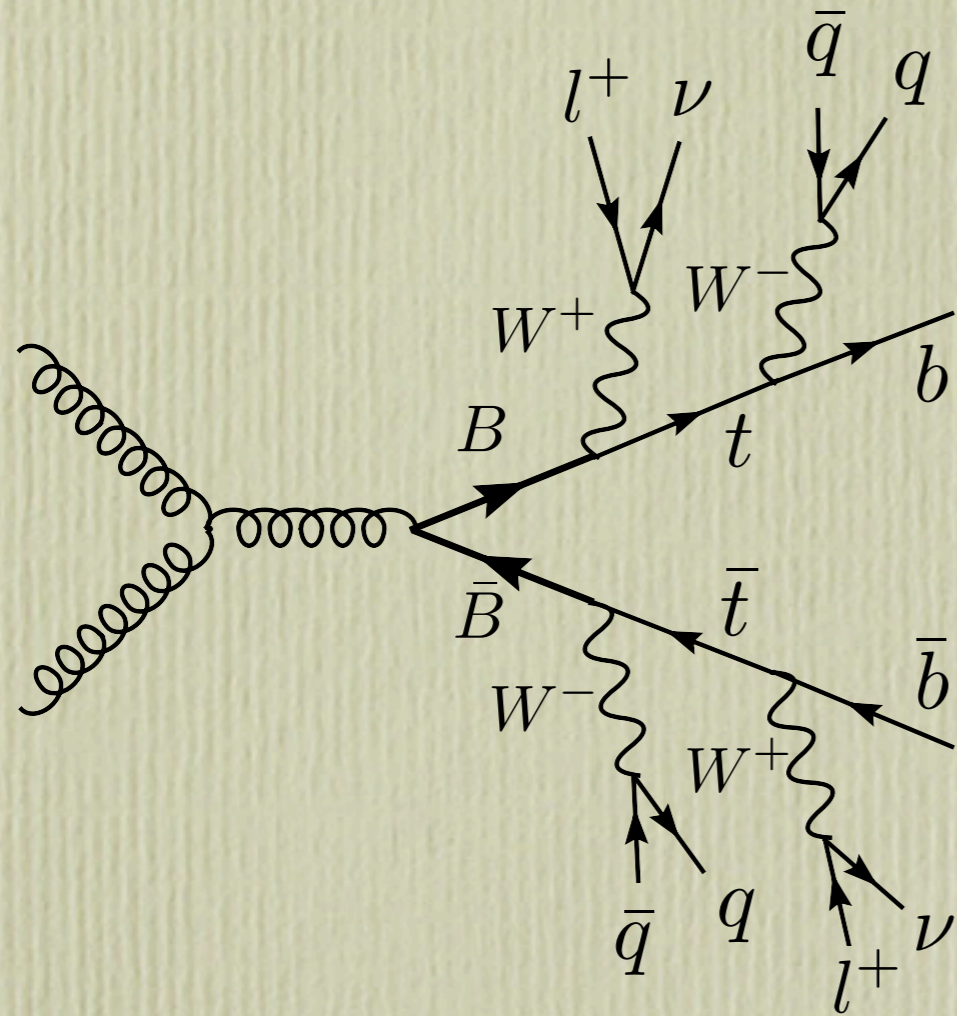
Decay modes

FCNC : absent for a 4th generation !



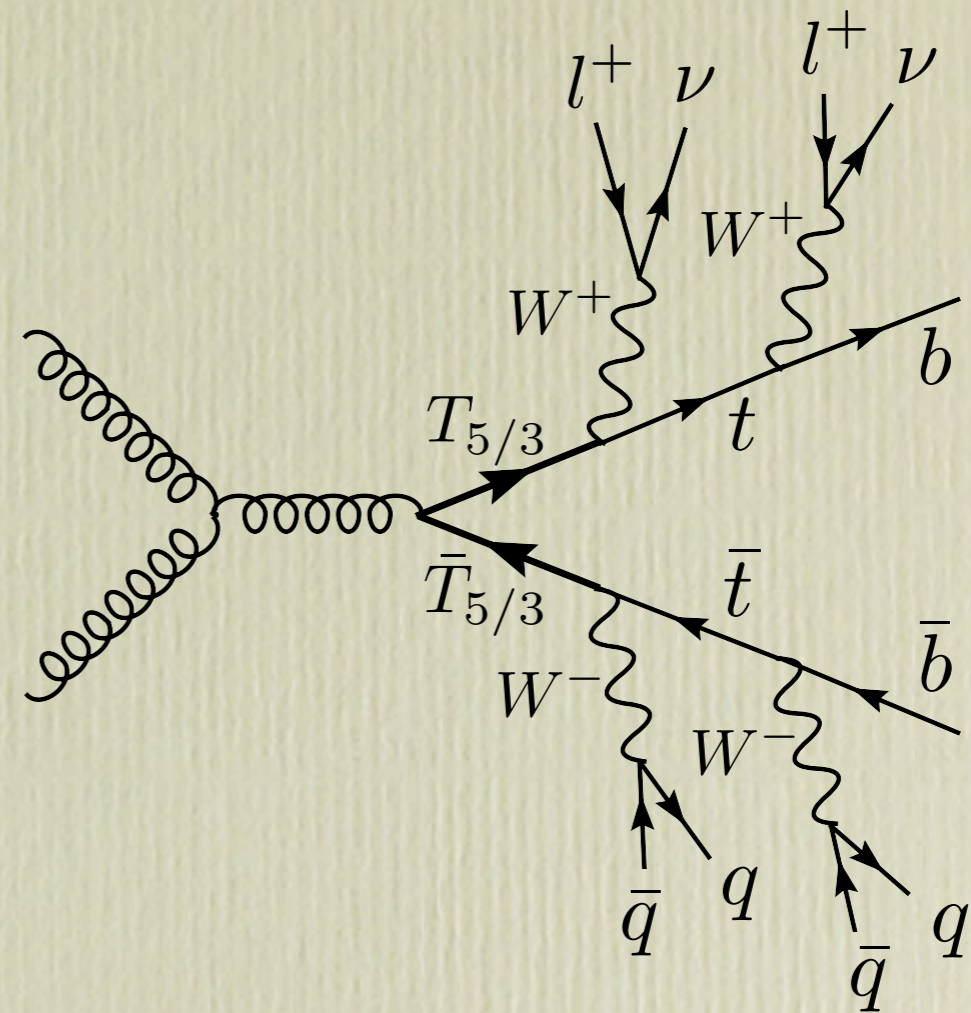
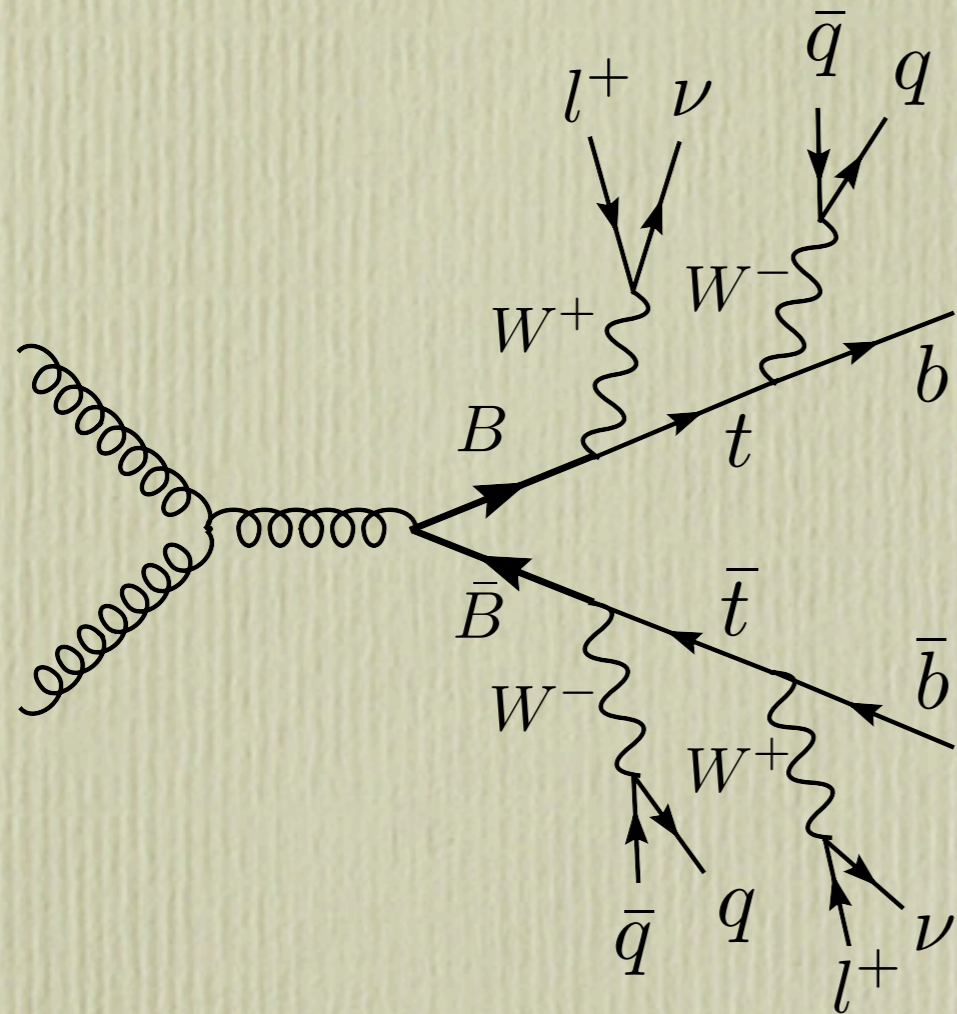
👉 Example: Look for $B\bar{B}$ and $T_{5/3}\bar{T}_{5/3}$ in same-sign di-lepton final states

[R.C. and G.Servant JHEP 0806:026 (2008)]



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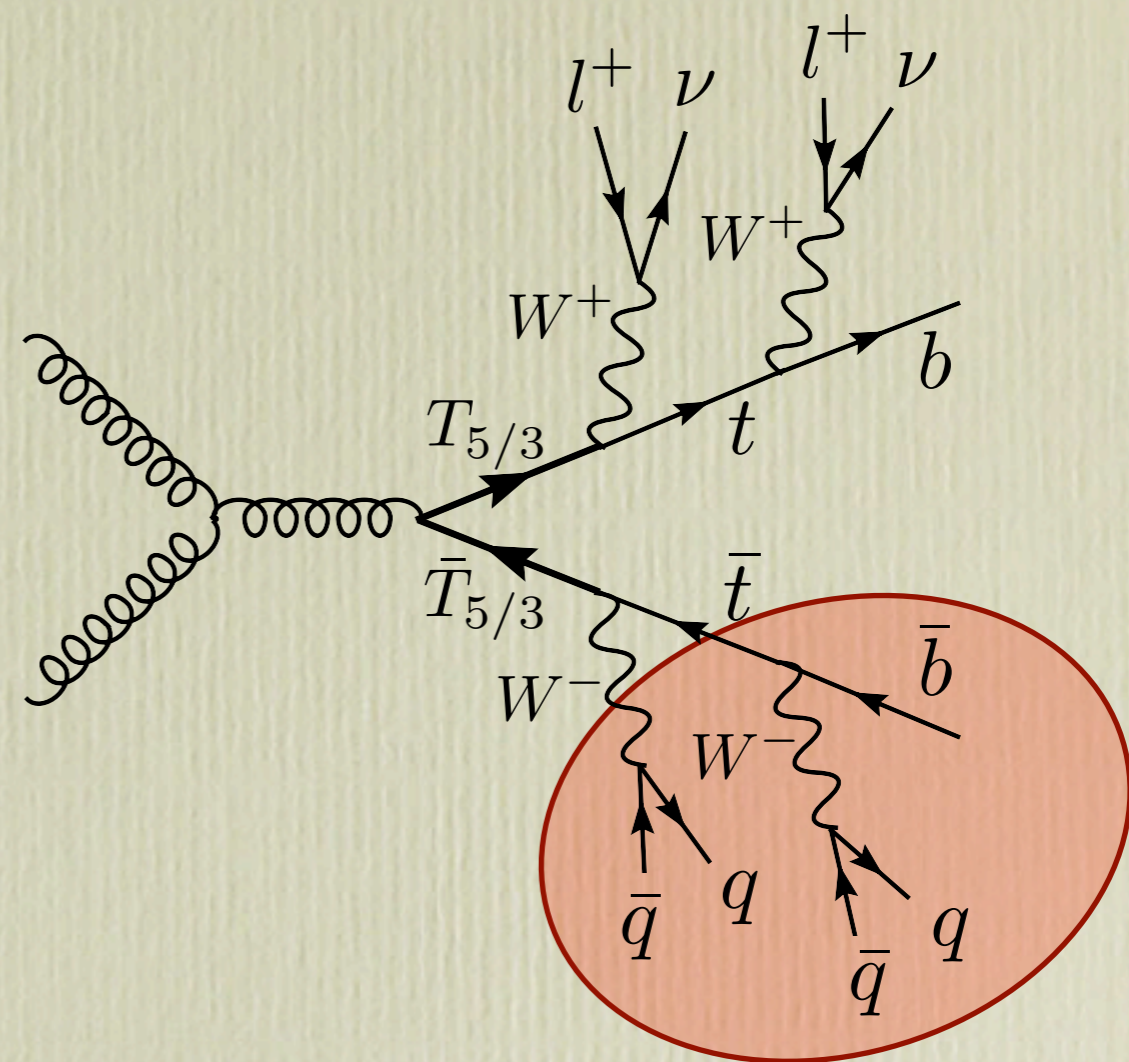
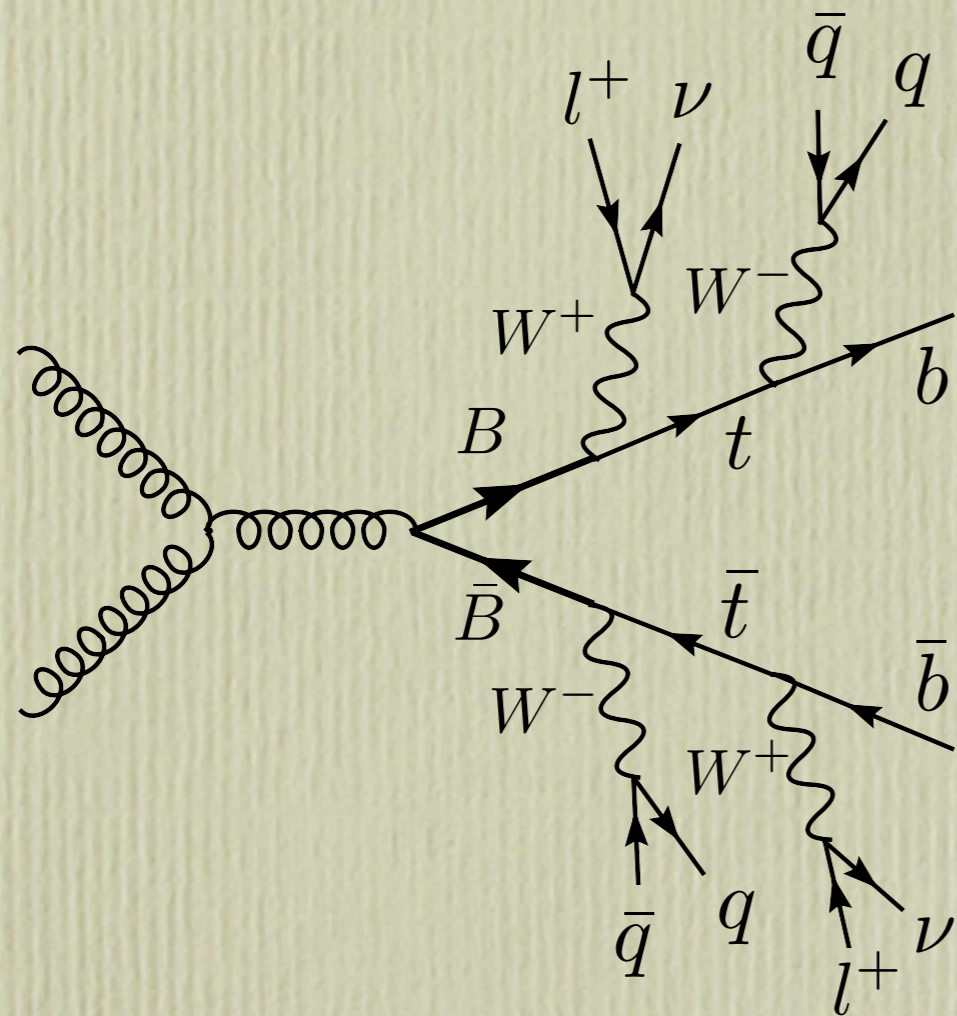
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✓ $t\bar{t} + jets$ is not a background [except for charge mis-ID]

✓ For the $T_{5/3}$ case one can reconstruct the resonant (tW) invariant mass

Signal and Background Simulation

Signal and SM background have been simulated using:

- ❖ MadGraph/MadEvent [MatrixElement] + Pythia [Showering - no hadronization or und.event]
- ❖ Quark/Jet matching a la MLM
- ❖ Jets reconstructed with a cone algorithm (GetJet) with $\Delta R = 0.4$, $E_T^{min} = 30$ GeV
- ❖ Jet energy and momentum smeared by $100\%/\sqrt{E}$ to simulate the detector resolution

	σ [fb]	$\sigma \times BR(l^\pm l^\pm)$ [fb]
$T_{5/3}\bar{T}_{5/3}/B\bar{B} + jets$ ($M = 500$ GeV)	2.5×10^3	104
$T_{5/3}\bar{T}_{5/3}/B\bar{B} + jets$ ($M = 1$ TeV)	37	1.6
$t\bar{t}W^+W^- + jets$ ($\supset t\bar{t}h + jets$)	121	5.1
$t\bar{t}W^\pm + jets$	595	18.4
$W^+W^-W^\pm + jets$ ($\supset hW^\pm + jets$)	603	18.7
$W^\pm W^\pm + jets$	340	15.5

SM bckg
[$m_h = 180$ GeV]

Strategy and cuts

★ We demand at least 5 hard jets ($p_T \geq 30$ GeV) :

$$l^\pm l^\pm + n \text{ jets} + \cancel{E}_T \quad (n \geq 5)$$

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$$\begin{array}{l} \underline{\text{jets}} : \left\{ \begin{array}{l} p_T(1\text{st}) \geq 100 \text{ GeV} \\ p_T(2\text{nd}) \geq 80 \text{ GeV} \\ n_{jet} \geq 5, \quad |\eta_j| \leq 5 \end{array} \right. \quad \underline{\text{leptons}} : \left\{ \begin{array}{l} p_T(1\text{st}) \geq 50 \text{ GeV} \\ p_T(2\text{nd}) \geq 25 \text{ GeV} \\ |\eta_l| \leq 2.4, \quad \Delta R_{lj} \geq 0.4 \end{array} \right. \quad \cancel{E}_T \geq 20 \text{ GeV} \end{array}$$

	signal ($M = 500$ GeV)	signal ($M = 1$ TeV)	$t\bar{t}W$	$t\bar{t}WW$	WWW	$W^\pm W^\pm$
Efficiencies (ϵ_{main})	0.42	0.43	0.074	0.12	0.008	0.01
σ [fb] $\times BR \times \epsilon_{main}$	44.2	0.67	1.4	0.62	0.15	0.16

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Extra Cuts
for $M=1$ TeV:

$$p_T(1\text{st jet}) \geq 200 \text{ GeV}$$

$$\sum_i |\vec{p}_T(l_i)| \geq 300 \text{ GeV}$$

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Efficiencies (ϵ_{disc})	0.65	0.091	0.032	0.16	0.18
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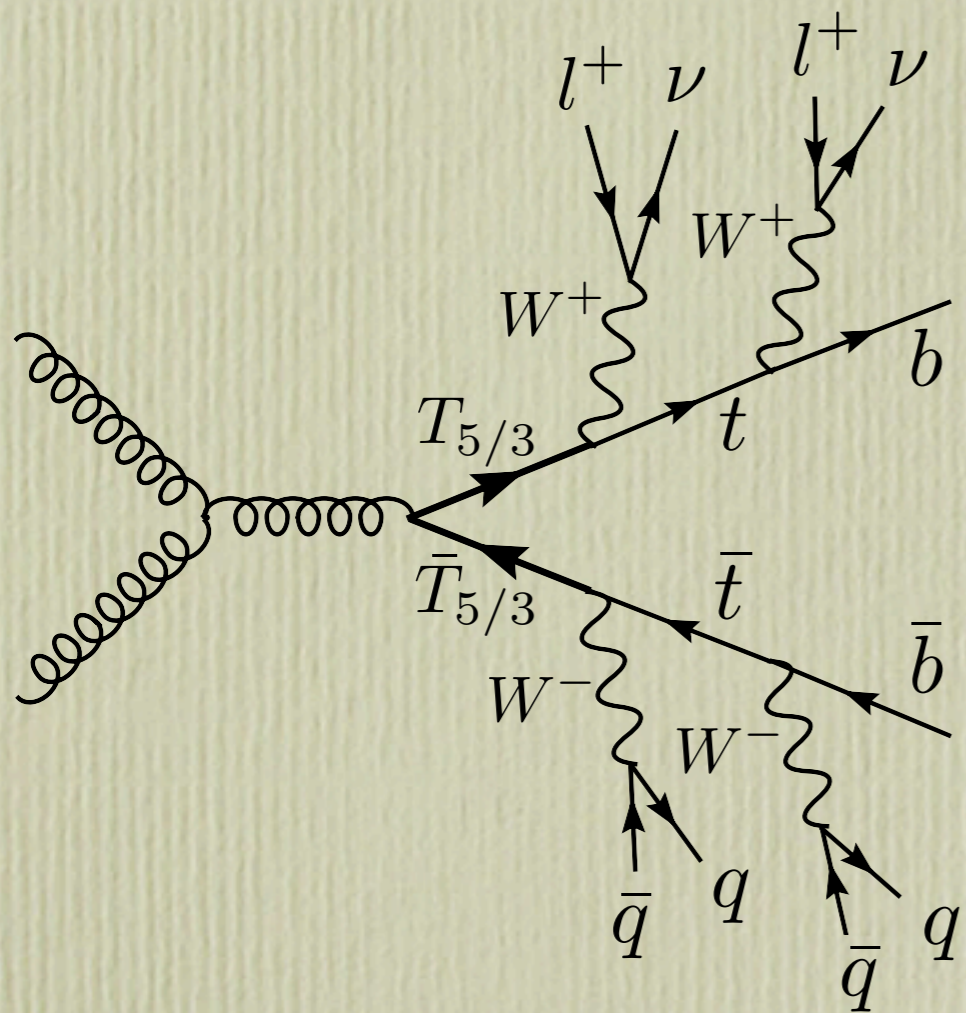
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Discovery Potential:

		L_{disc}		L_{disc}
$M = 500$ GeV	$T_{5/3} + B$	56 pb^{-1}	$M = 1$ TeV	$T_{5/3} + B$
	B only	147 pb^{-1}		B only
				48 fb^{-1}

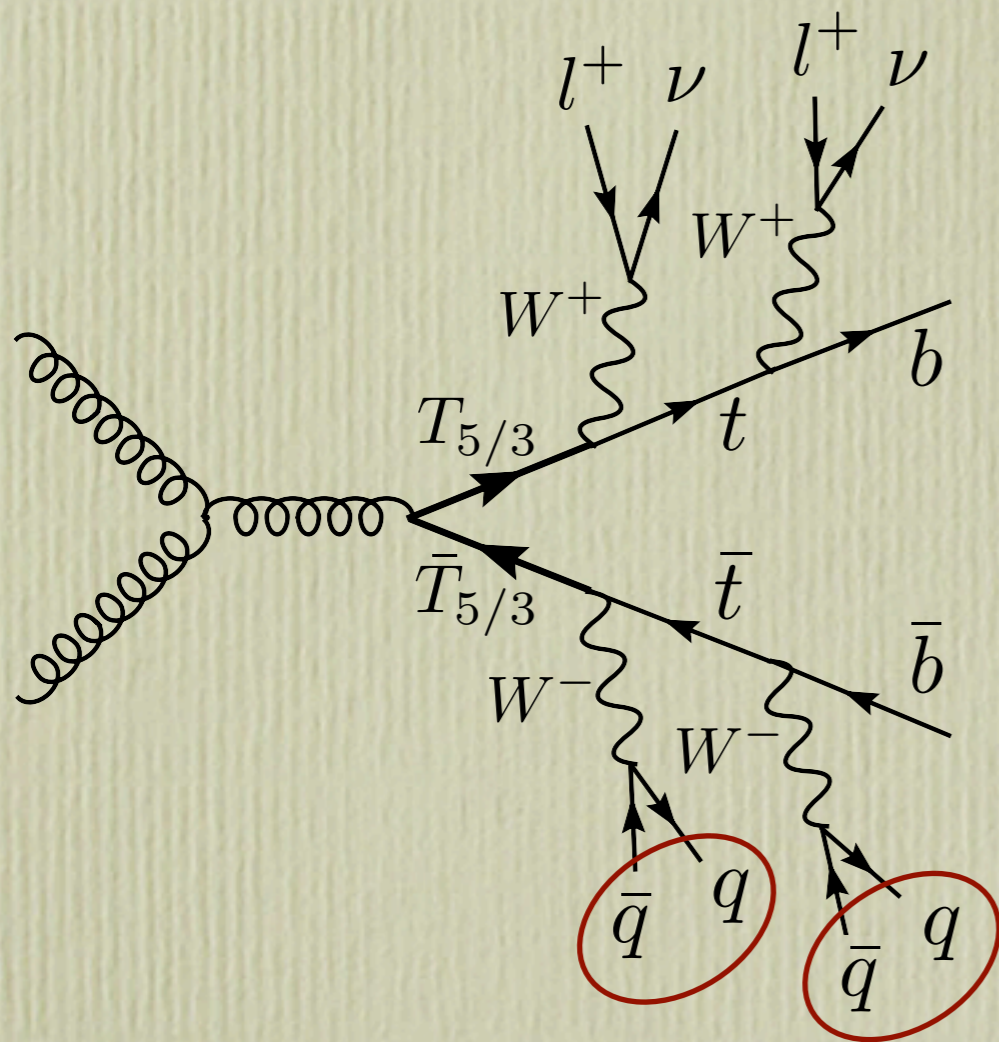
Mass Reconstruction

$M=500$ GeV



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1. Reconstruct 2 W 's

$$|M(jj) - m_W| \leq 20 \text{ GeV}$$

$$\Delta R_{jj}(\text{1st pair}) \leq 1.5$$

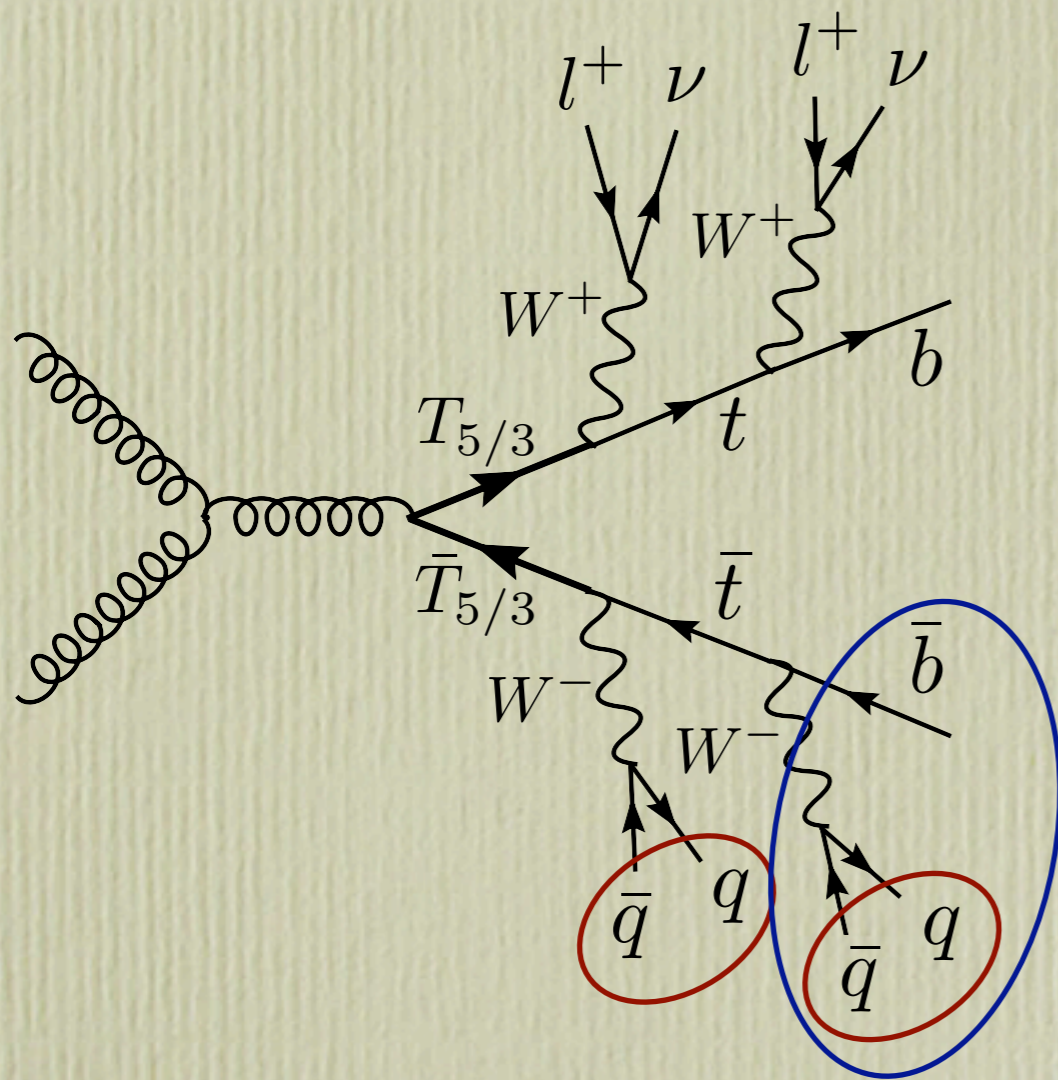
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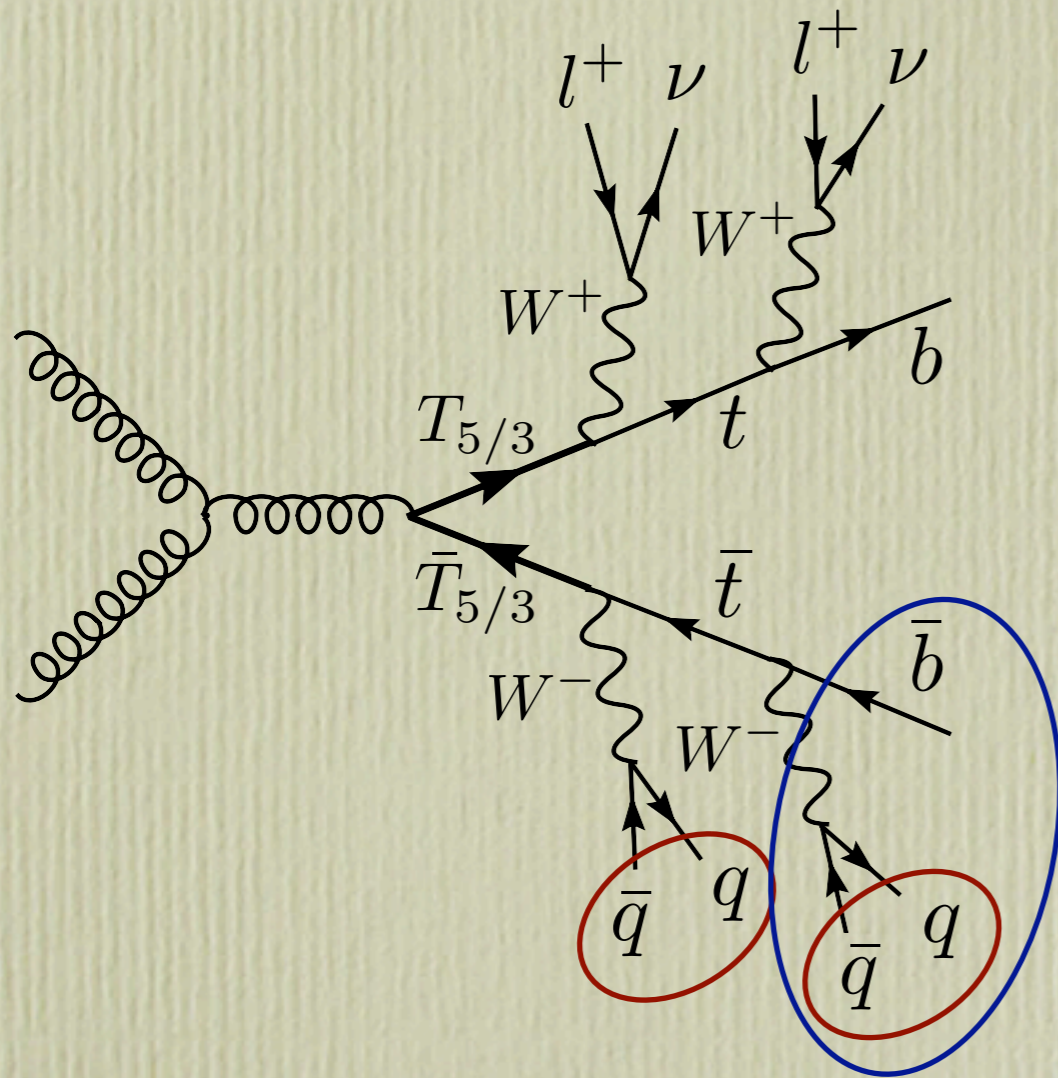
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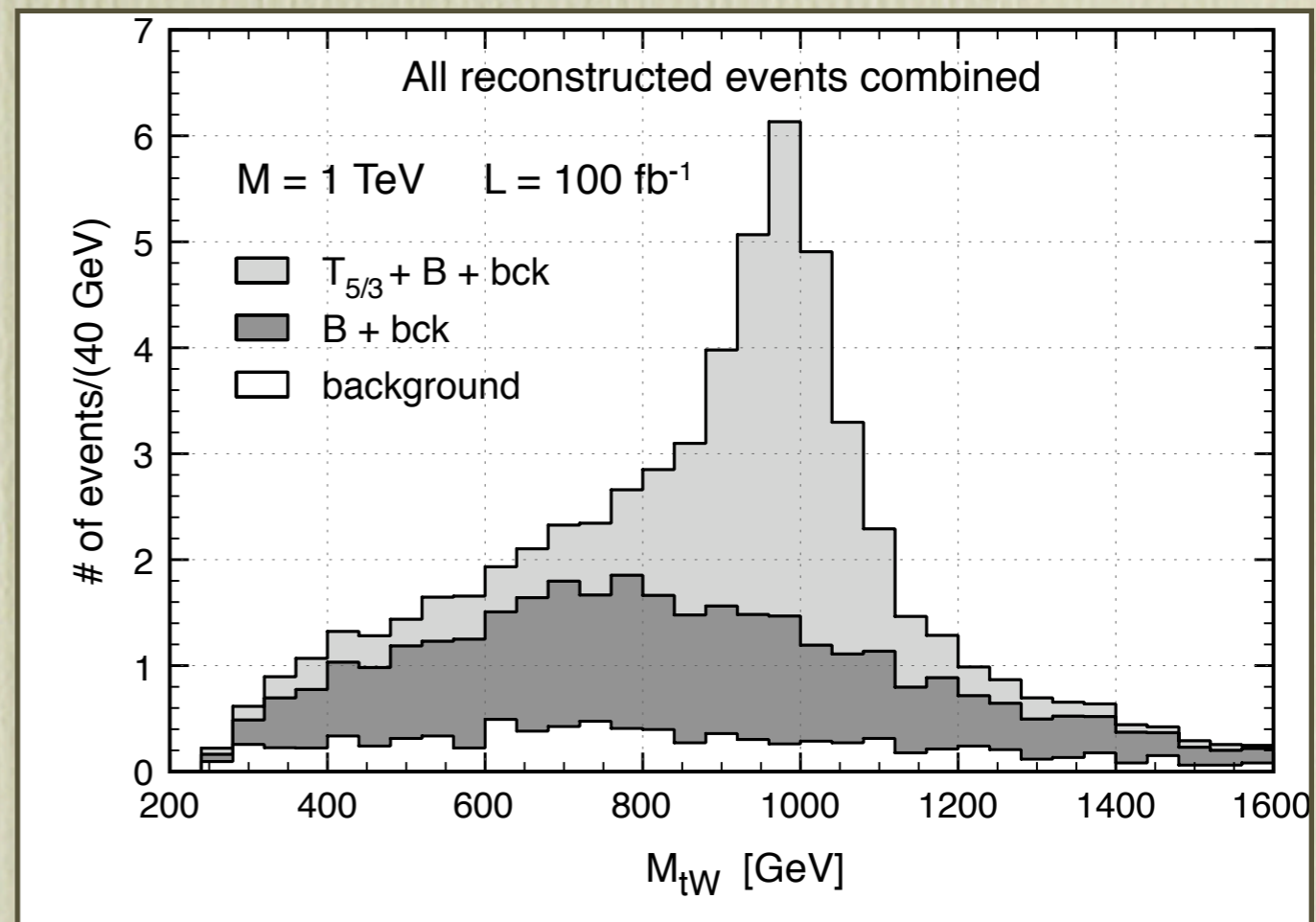
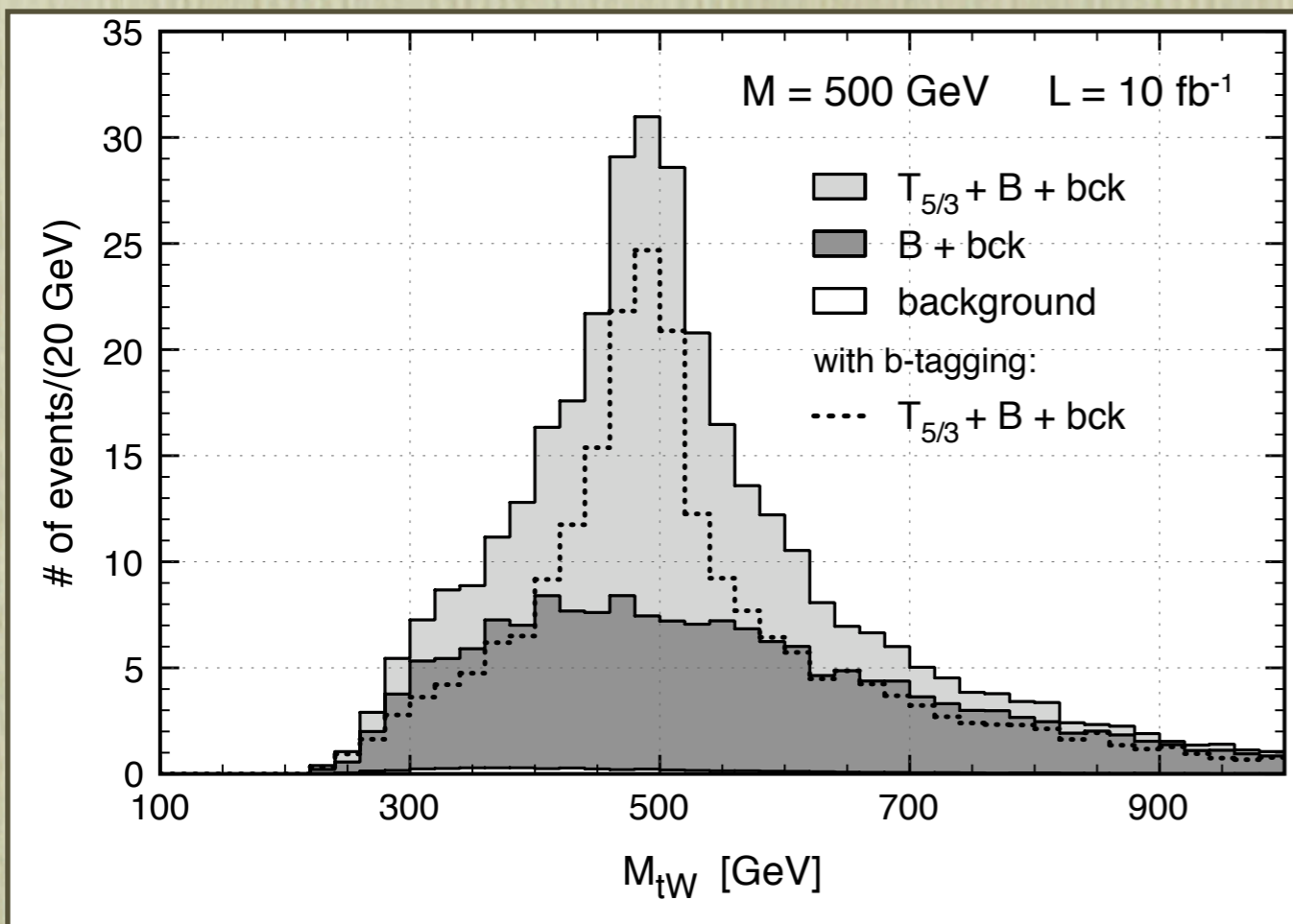
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	signal ($M = 500$ GeV)	$t\bar{t}W$	$t\bar{t}WW$	WWW	WW
ϵ_{2W}	0.62	0.36	0.49	0.29	0.15
ϵ_{top}	0.65	0.56	0.64	0.35	0.35



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- Prediction: Heavy tops and bottoms and possibly other exotic fermions
- Same-sign di-lepton channels promising for discovering B and $T_{5/3}$

Extra Slides

