

# Longitudinal structure function measurements from HERA

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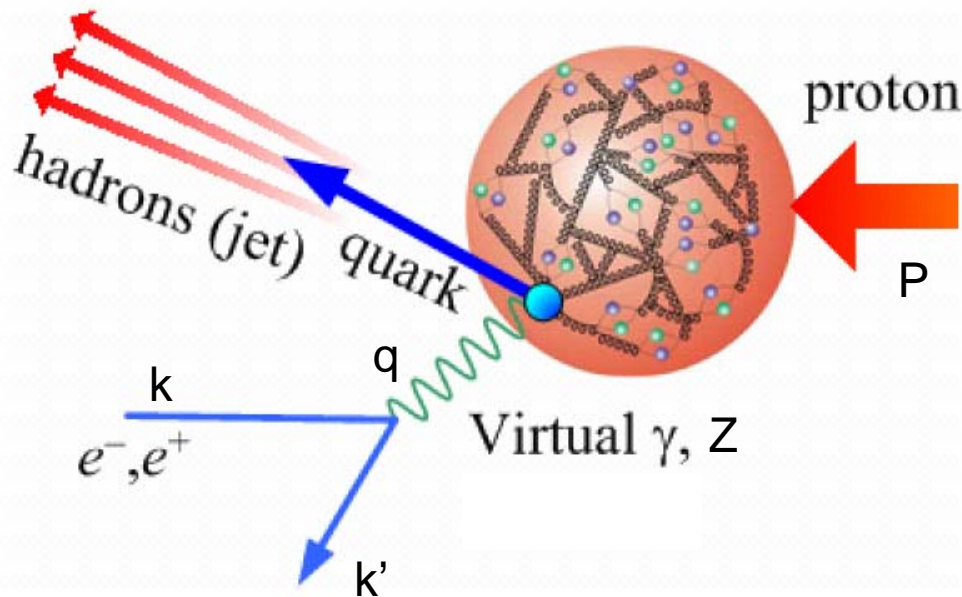
*on behalf of H1 and ZEUS*



- Deep Inelastic Scattering / Structure functions
- Longitudinal structure function  $F_L(x, Q^2)$
- HERA / H1 and ZEUS
- Measurement strategy for  $F_L$
- Experimental details of the  $F_L$  analyses
- $F_L$  results
- Summary

# Deep Inelastic Scattering

*Neutral Current (NC):*  $e^{\pm} p \rightarrow e^{\pm} X$



$Q^2 = -q^2 = -(k-k')^2$  virtuality of  $\gamma^*, Z$

$x = Q^2/2(Pq)$  Bjorken  $x$

$y = (Pq)/(Pk)$  inelasticity

$Q^2 = sxy$   $s=(k+P)^2$

## Factorisation

$$\sigma_{DIS} \sim \hat{\sigma} \otimes pdf(x)$$

$\hat{\sigma}$  – perturbative QCD cross section  
 $pdf$  – universal parton distribution functions

# The Proton Structure Functions

$$\frac{d^2\sigma_{NC}^{e^+p}}{dx dQ^2} = \frac{2\pi\alpha^2 Y_+}{xQ^4} \sigma_r^\pm = \frac{2\pi\alpha^2 Y_+}{xQ^4} \left[ F_2(x, Q^2) - \frac{y^2}{Y_+} F_L(x, Q^2) \mp \frac{Y_-}{Y_+} xF_3(x, Q^2) \right]$$

helicity factors:  $Y_\pm = 1 \pm (1-y)^2$

*dominant contribution:*

$$F_2(x, Q^2) = \sum e_{q_i}^2 x(q_i + \bar{q}_i)$$

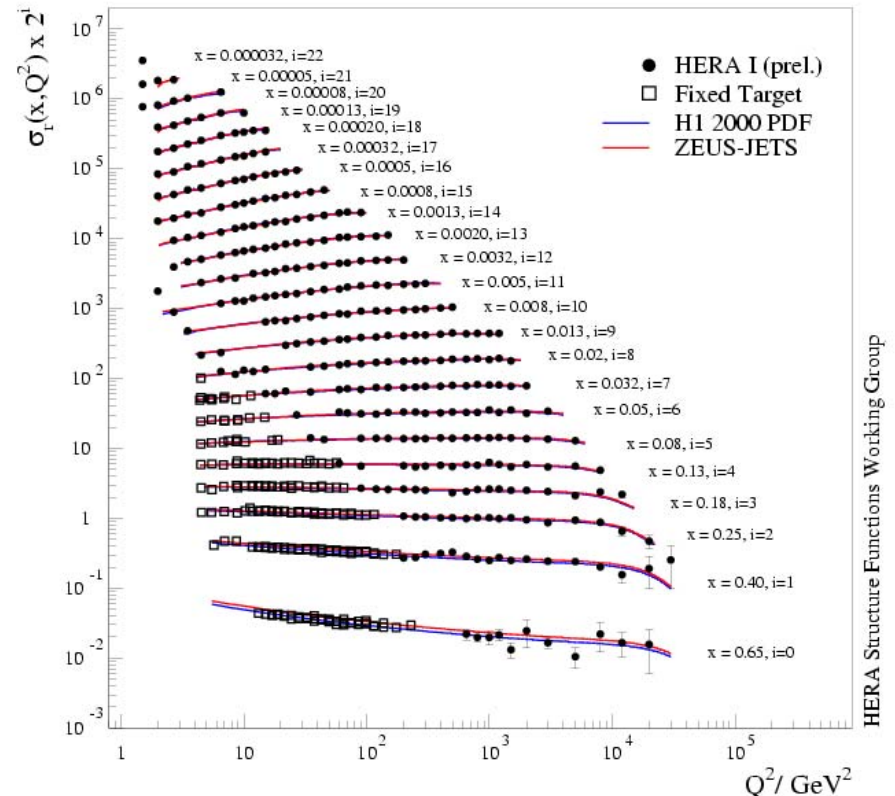
contributes only at high  $Q^2$  ( $\gtrsim M_Z^2$ )

$$xF_3(x, Q^2) = x \sum B_i(q_i - \bar{q}_i)$$

$$F_2 \sim \sigma_L^{\gamma p} + \sigma_T^{\gamma p}, \quad F_L \sim \sigma_L^{\gamma p}$$

$$\rightarrow 0 \leq F_L \leq F_2$$

HERA 1  $e^+p$  Neutral Current Scattering - H1 and ZEUS



# The longitudinal structure function $F_L(x, Q^2)$

- $F_L$  is an independent structure function to be measured at HERA to complete the DIS program
- $F_L$  is a pure QCD effect which allows to make critical tests of the perturbative QCD framework used for pdf determinations
- $F_L$  is directly sensitive to gluon density

in QPM

due to helicity and angular momentum conservation for spin  $\frac{1}{2}$  quarks

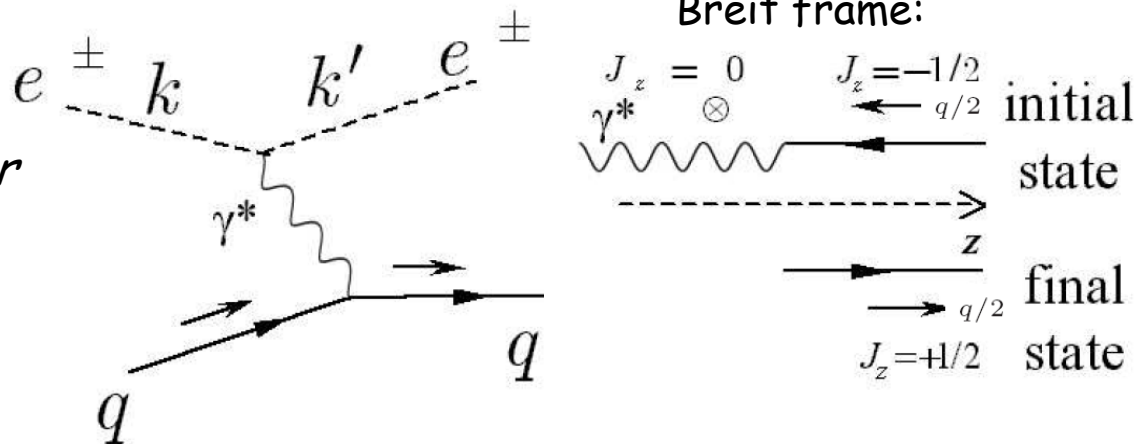
$$F_L \sim \sigma_L^{\gamma p} = 0$$

$$F_L = F_2 - 2xF_1 = 0$$

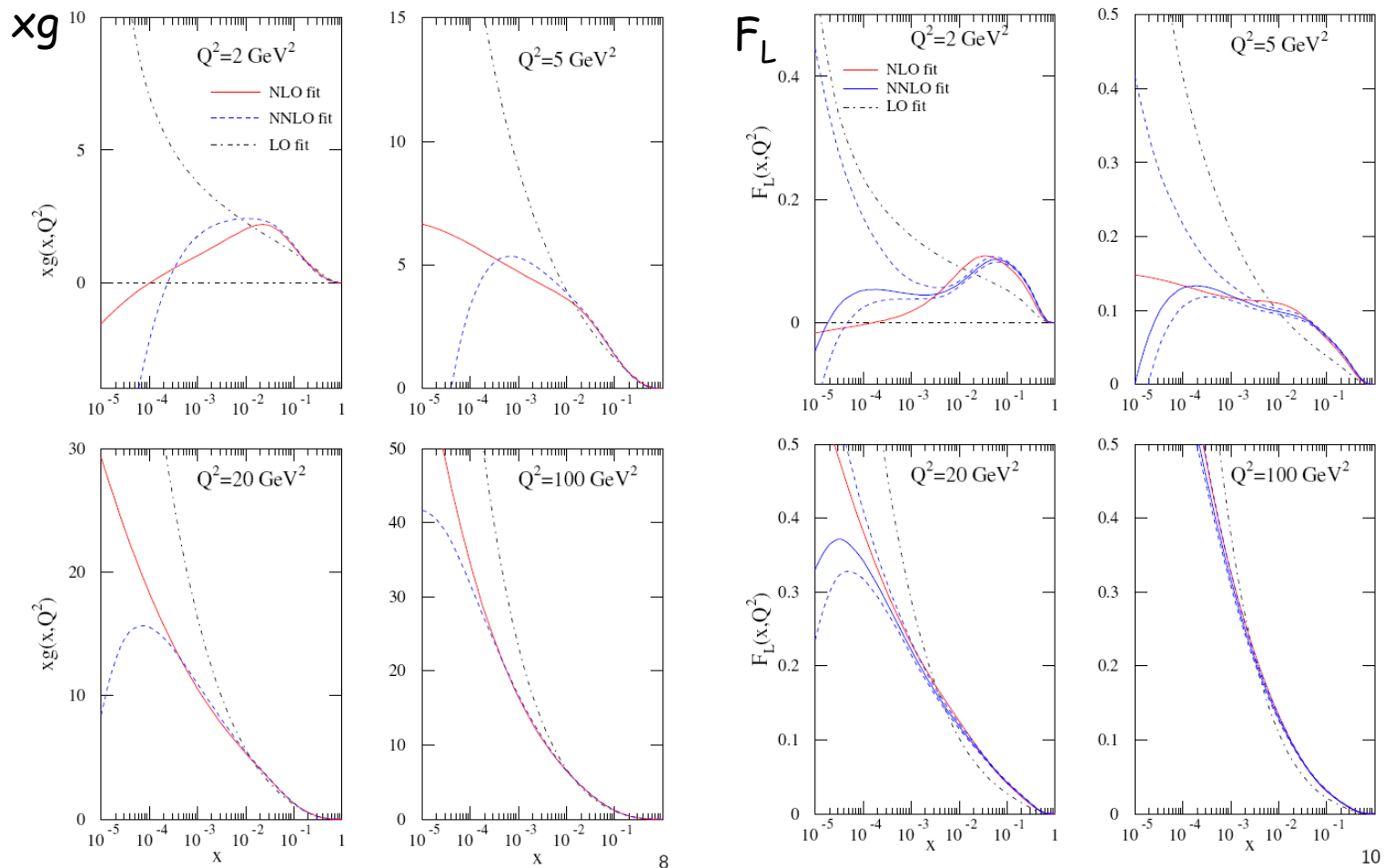
Callan-Gross relation

in QCD:

$$F_L(x, Q^2) = \frac{\alpha_s}{4\pi} x^2 \int_x^1 \frac{dz}{z^3} \left[ \frac{16}{3} F_2 + 8 \sum_q e_q^2 \left(1 - \frac{x}{z}\right) \cdot xg \right]$$

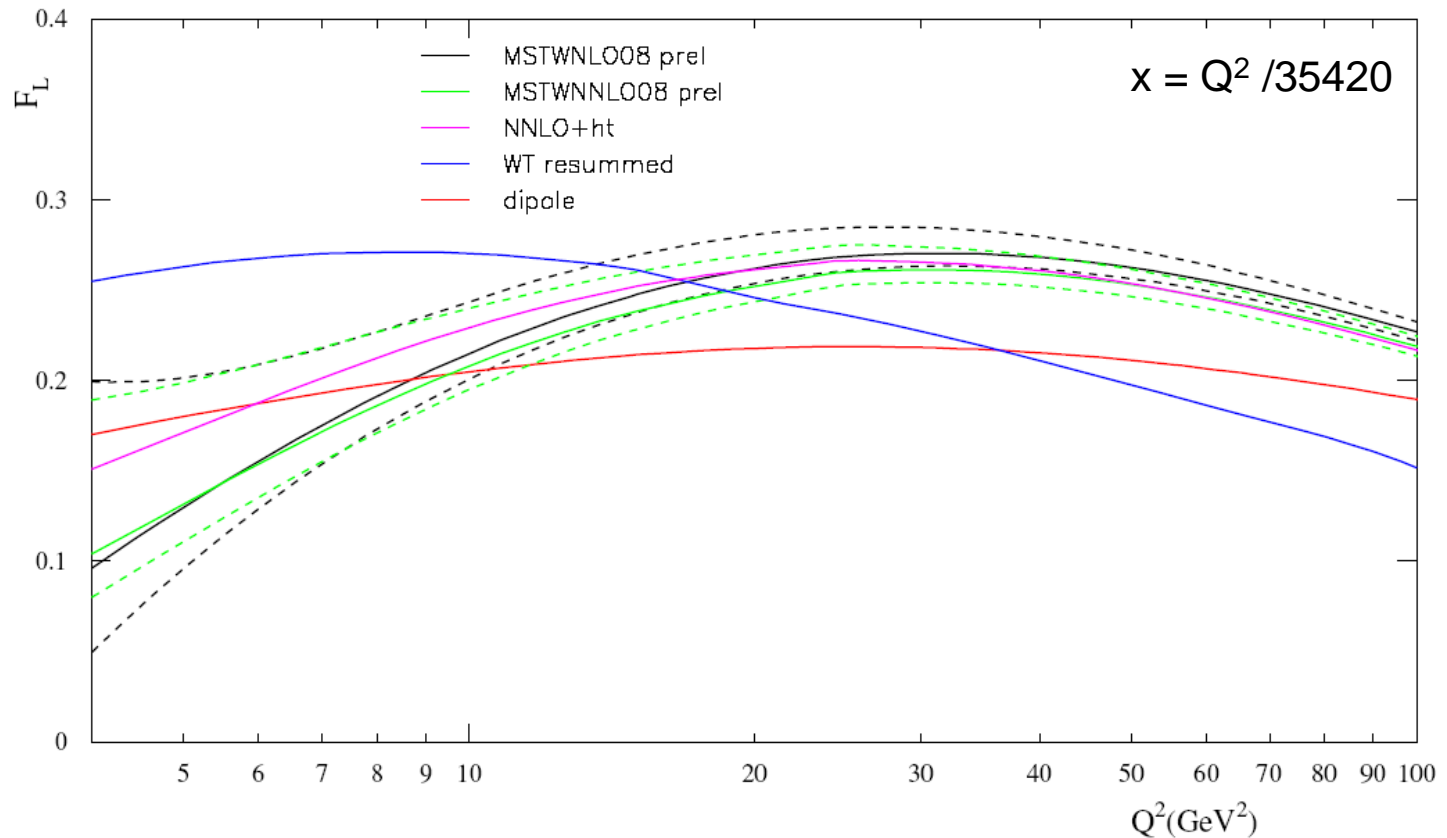


# Gluon and $F_L$ in LO – NLO – NNLO (MSTW)



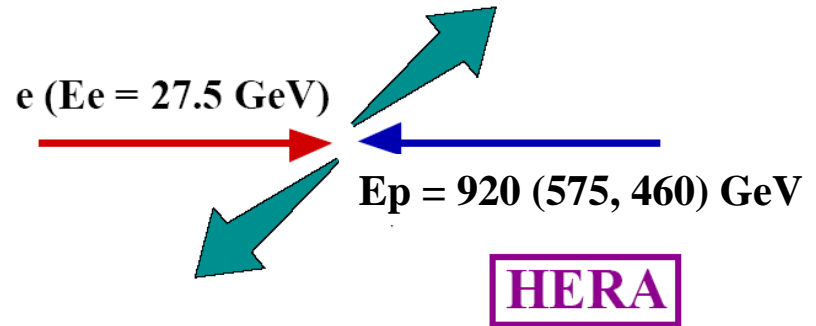
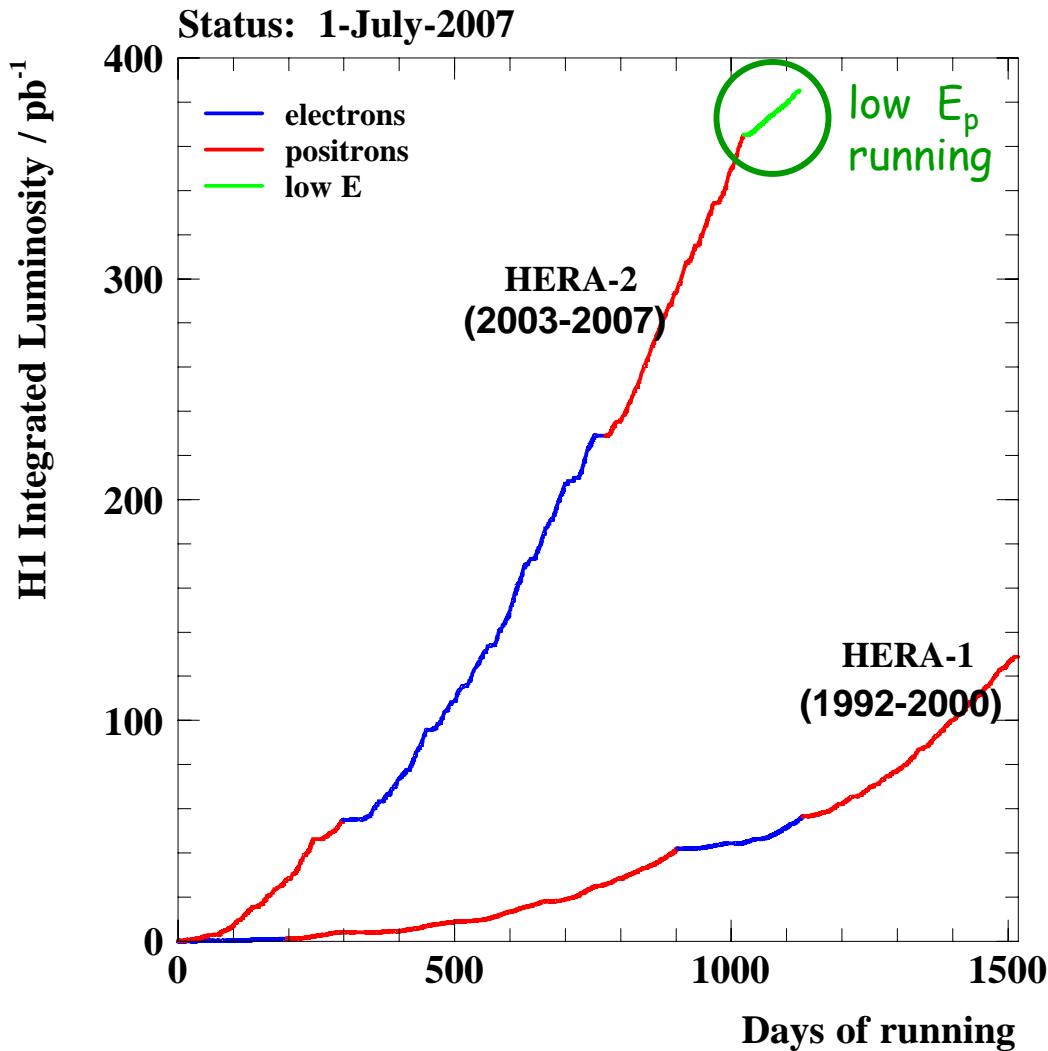
→ poor stability for gluon at small  $x$  → similarly for  $F_L$  but less prominent

# Theory predictions for $F_L$ in the HERA domain



- firm NLO/NNLO QCD predictions for  $Q^2 > 10 \text{ GeV}^2$
- spread of predictions at  $Q^2$  below  $10 \text{ GeV}^2$

# HERA (1992–2007)



peak luminosity  $5 \cdot 10^{31} \text{ cm}^{-2} \text{ sec}^{-1}$   
 $Q^2_{\text{max}} = 10^5 \text{ GeV}^2$   
 $\lambda_{\text{max}} \sim 1/1000 r_{\text{proton}}$   
 longitudinal  $e$ -beam polarisation

H1+ZEUS in total  $\sim 1 \text{ fb}^{-1}$

about equally shared between

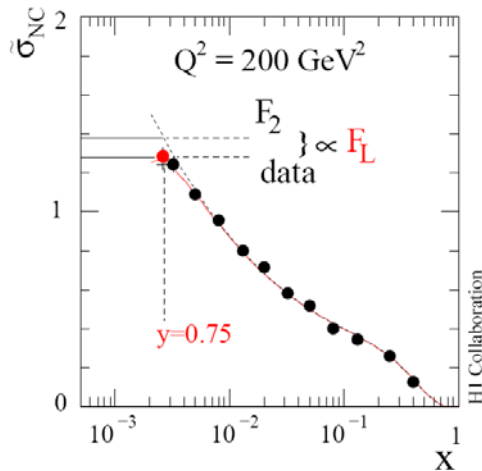
- experiments (H1, ZEUS)
- $e^+$  and  $e^-$ ,
- positive and negative  $P_e$

→ low proton energy run for direct  $F_L$  measurements  
 $13 \text{ pb}^{-1} E_p = 460 \text{ GeV}$   
 $7 \text{ pb}^{-1} E_p = 575 \text{ GeV}$

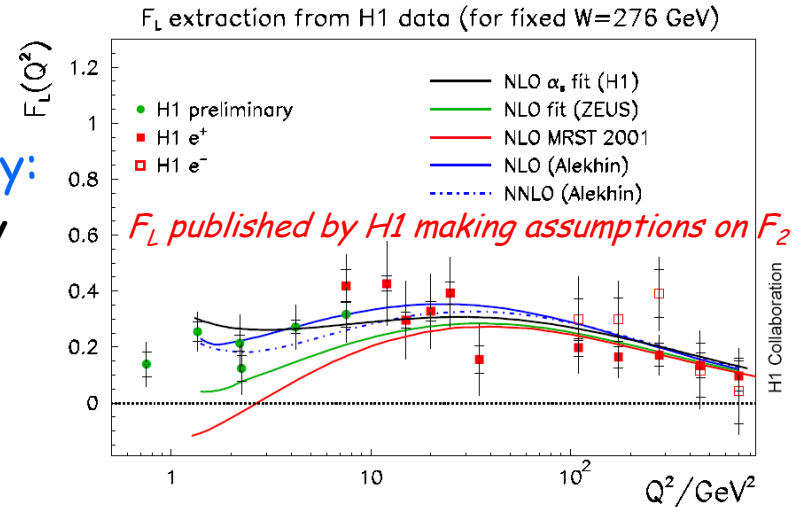
# Measurement strategy for $F_L$

$$\tilde{\sigma}_{NC} = \frac{d^2\sigma_{NC}^{ep}}{dx dQ^2} / \left( \frac{2\pi\alpha^2}{xQ^4} Y_+ \right) = F_2 - \frac{y^2}{1 + (1-y)^2} F_L$$

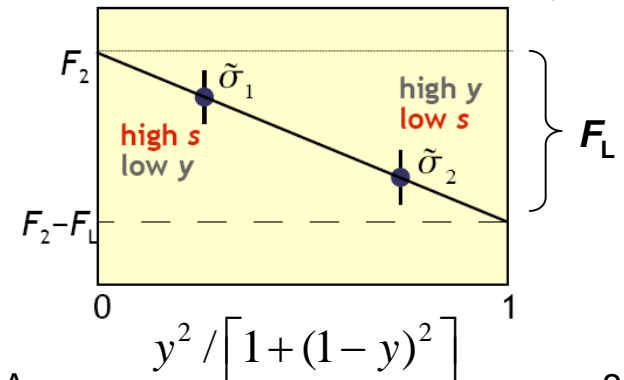
sensitivity to  $F_L$   
only at high  $y$



→ one possible way:  
measure  $\sigma$  at high  $y$   
and assume  $F_2$



→ free from theoretical assumption:  
measure  $\sigma$  at the same  $x$  &  $Q^2$  and different  $y$   
by changing the proton beam energy ( $y = Q^2/sx$ )



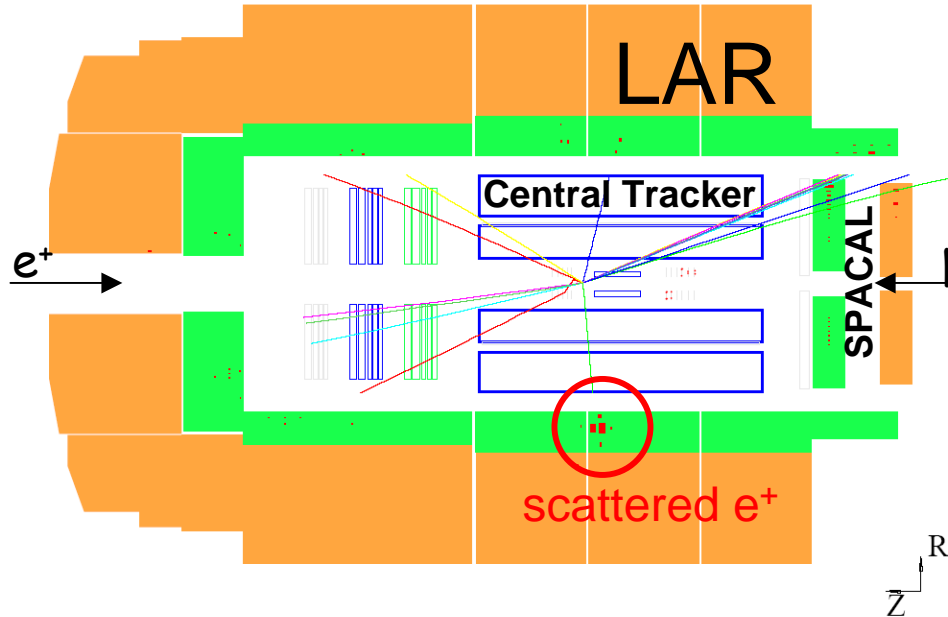


# H1 and ZEUS

$E'_e > 3 \text{ GeV}$  ( $y \approx 0.90$ )

$$y = 1 - (E'_e/E_e) \sin^2(\vartheta_e/2)$$

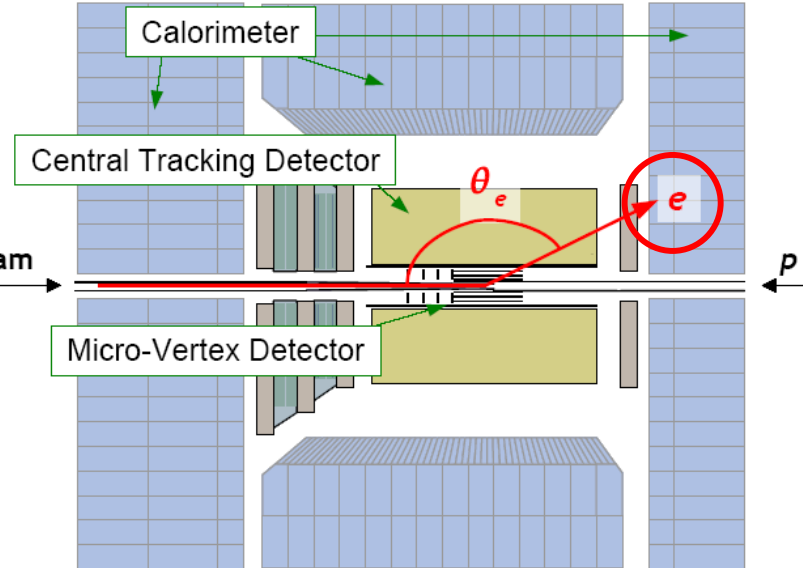
$E'_e > 6 \text{ GeV}$  ( $y \approx 0.76$ )



$F_L$  measurements in H1

$Q^2$ range ( $\text{GeV}^2$ )			
medium $Q^2$	12-90	Spacal+CT	DESY-08-053
high $Q^2$	35-800	LAr+CT	H1 prel.
low $Q^2$	5-15	Spac+BST	to come

V.Chekelian, 28.06.2008  
PIC 2008



$F_L$  measurements in ZEUS

$$\vartheta_e < 168^\circ$$

$$24 \leq Q^2 \leq 110 \text{ GeV}^2$$

more to come

FL measurements from HERA

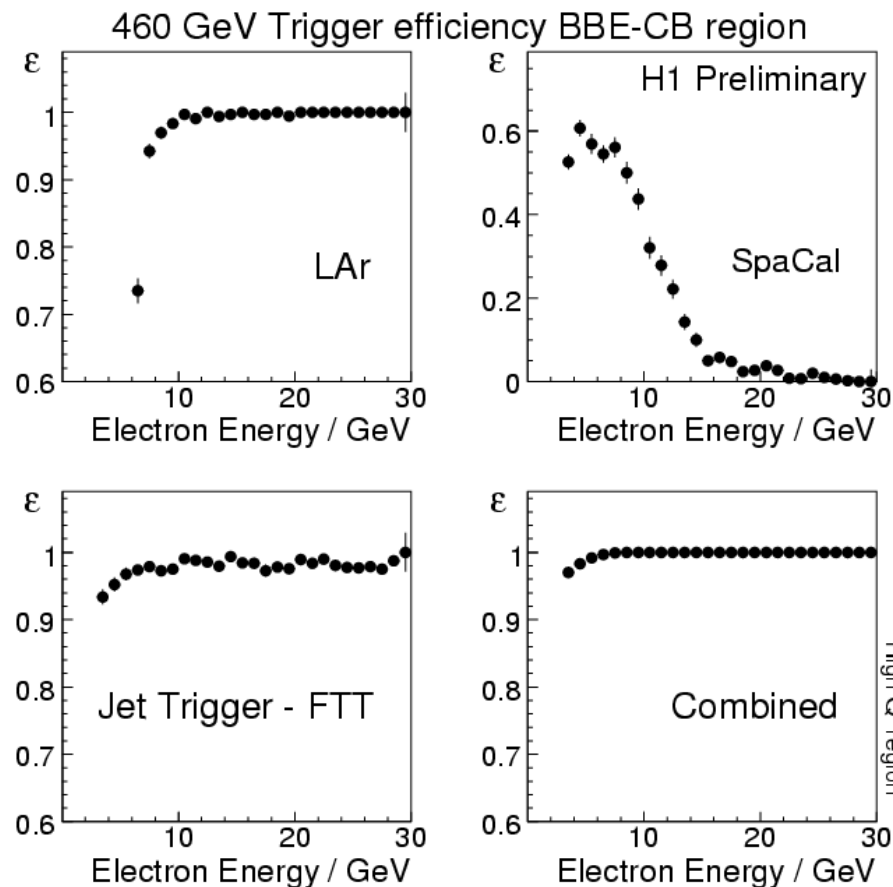
# Hardware & software improvements

H1: new trigger hardware since fall 2006:

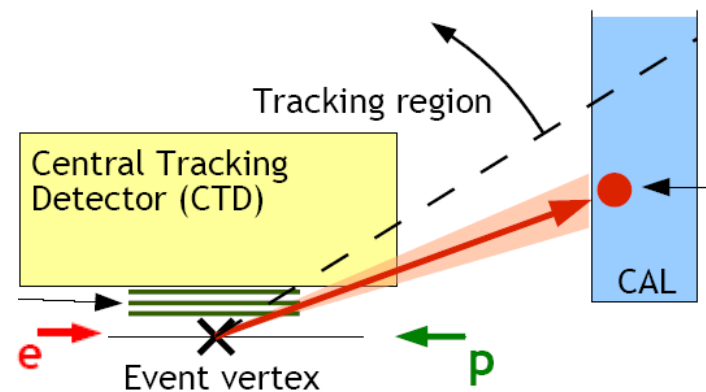
- Jet Trigger (real time clustering in LAr)
- Fast Track Trigger (FTT)

ZEUS: new tool is developed to extend the tracking region:

- acceptance of the track reconstruction is limited to  $\vartheta < 154^\circ$



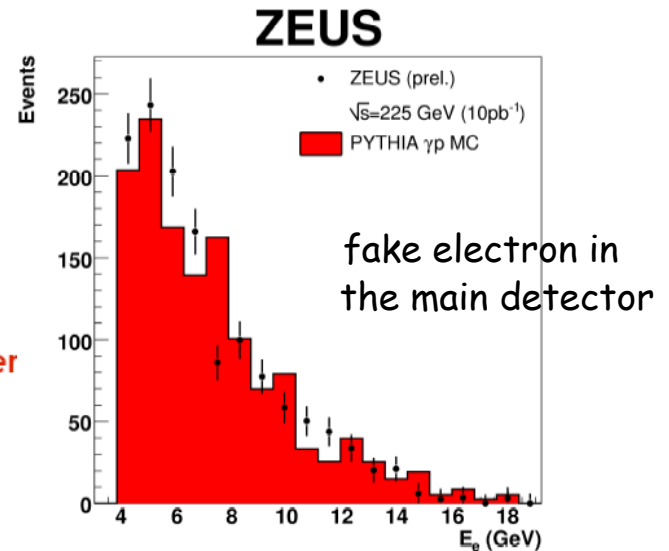
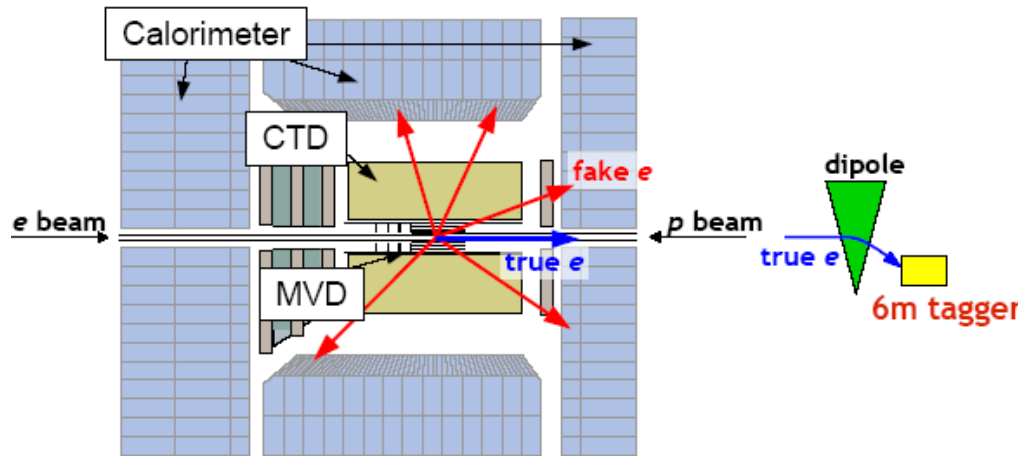
→ combined trigger eff.  $\approx 100\%$  for  $E_e > 3$  GeV



- use single hits in the tracking detector along a road from primary vertex to el. candidate in CAL taking into account the charge of the scattered electron

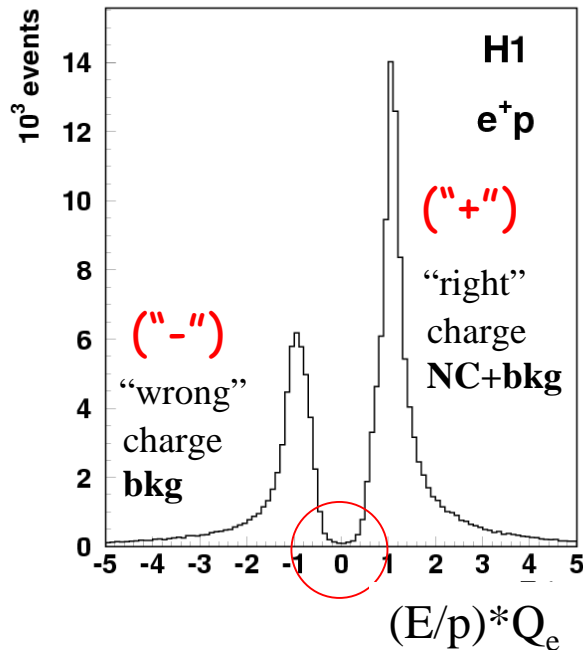
→ reject neutral particles up to  $\vartheta \approx 168^\circ$

# Photoproduction background estimation using 6m electron tagger (ZEUS)



- in photoproduction ( $Q^2 \approx 0$ ) quasi-real photon interacts with the proton
- electron with reduced energy goes along the e beam direction, bends in the dipole magnet and hits the electron tagger located at 6 m
- fraction of  $\gamma p$  events is measured in 6m tagger and used to normalize PYTHIA  $\gamma p$  MC for each  $E_p$  period
- H1 uses similar technique for  $E_p=920 \text{ GeV}$  at  $\gamma < 0.56$

# $\gamma p$ bkg identification up to $y=0.90$ (H1)



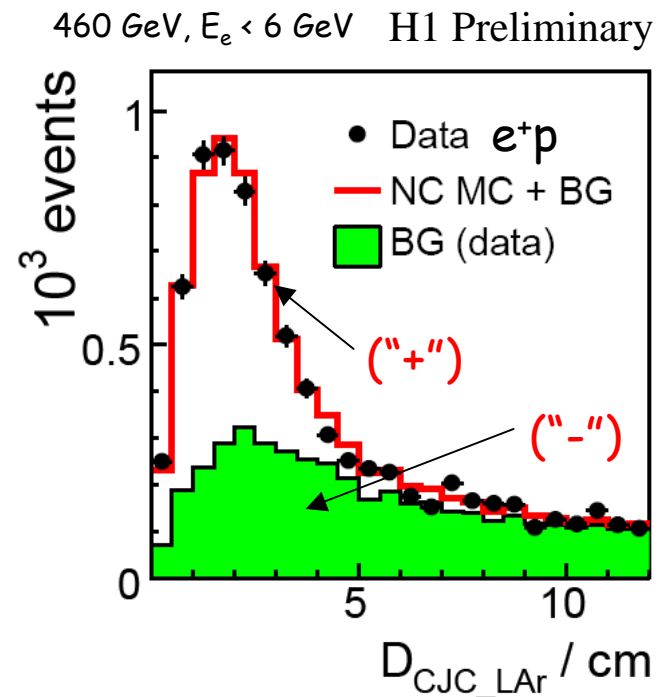
**Electric charge** of the scattered electron using track from the primary interaction, pointing to the electron cluster:

- good charge measurement resolution
- wrong assignment of the charge  $< 1\%$

1. identify and exclude half of  $\gamma p$  bkg require the “right” charge for el.
2. estimate and subtract remaining  $\gamma p$  bkg using “wrong” charge el.

taken into account in statistical subtraction:

- **charge asymmetry in  $\gamma p$  data** due to antiprotons determined using “wrong charge” el. candidates in the  $e^\pm p$  HERA II data and in  $\gamma p$  events identified by the 6 m electron tagger

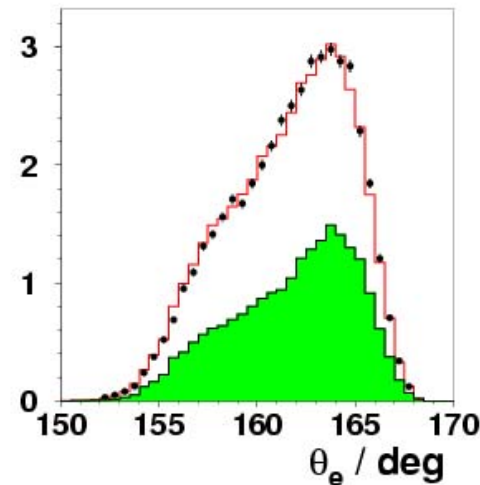
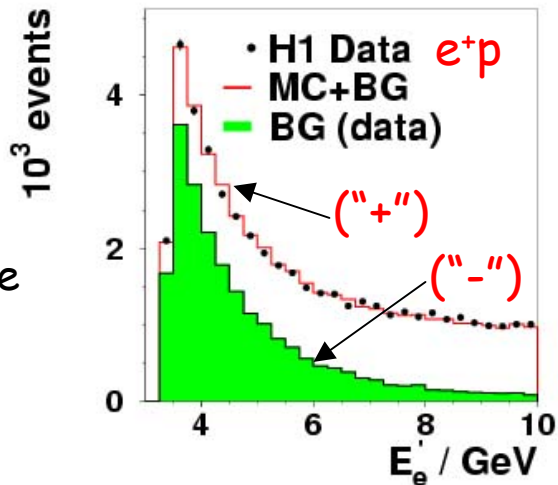


# High $y$ region at medium $Q^2$ (H1)

$E_p = 460 \text{ GeV}$

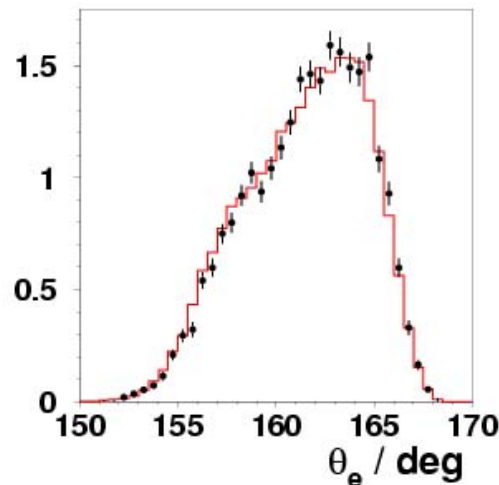
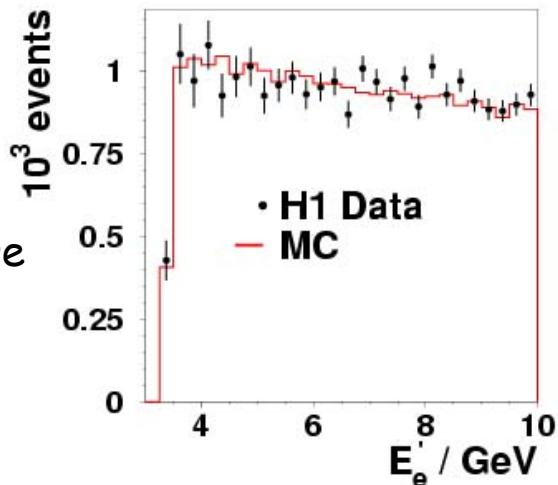
$E'_e < 10 \text{ GeV}$

before  
"wrong" charge  
subtraction



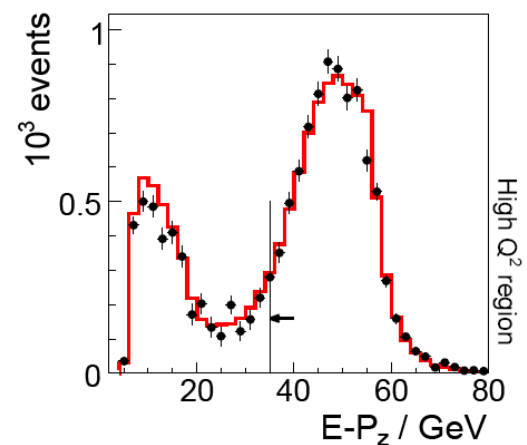
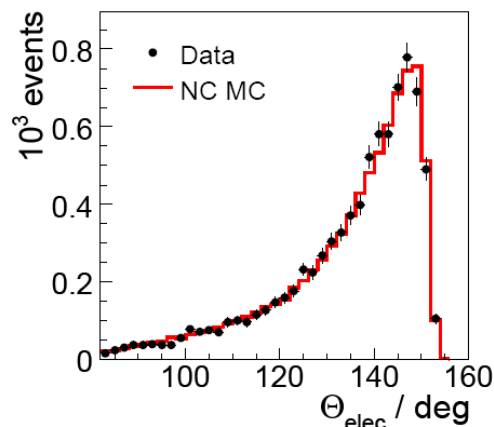
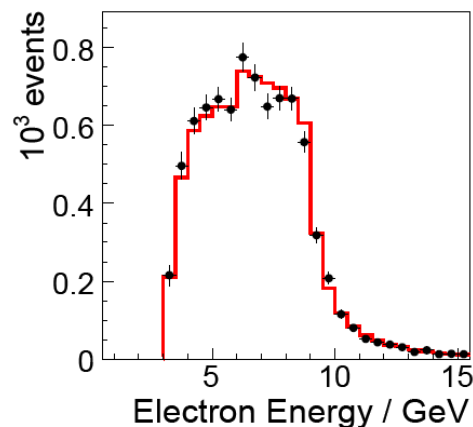
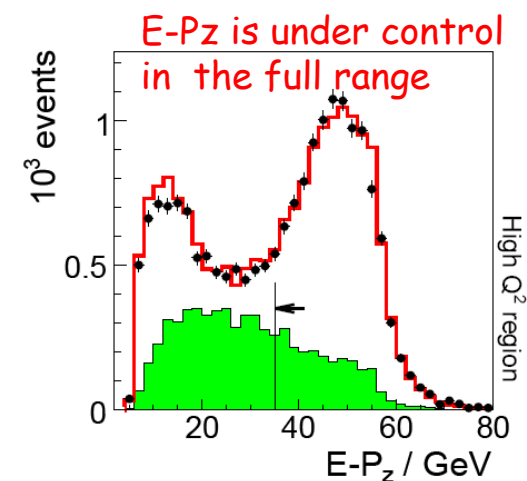
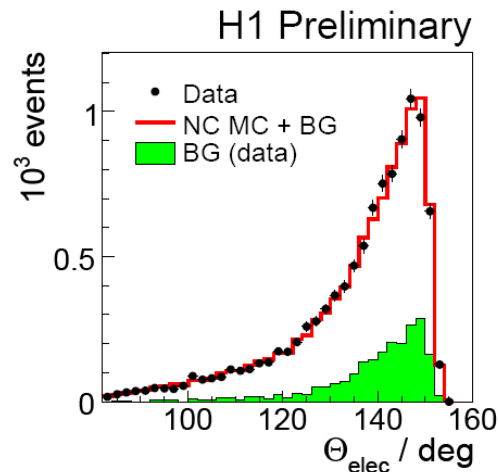
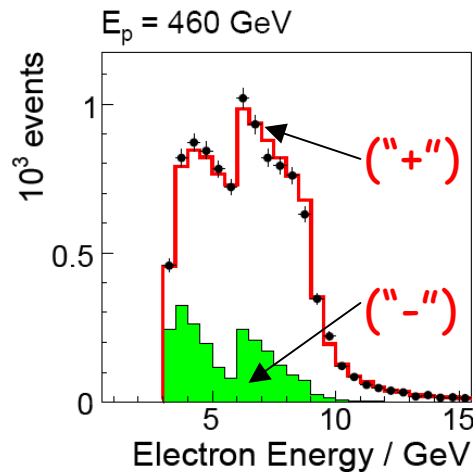
$\gamma p$  background (green)  
concentrates at low  $E'_e$

after  
"wrong" charge  
subtraction



the data are well  
understood in terms  
of MC

# High $y$ region ( $0.70 < y < 0.90$ ) at high $Q^2$ (H1)



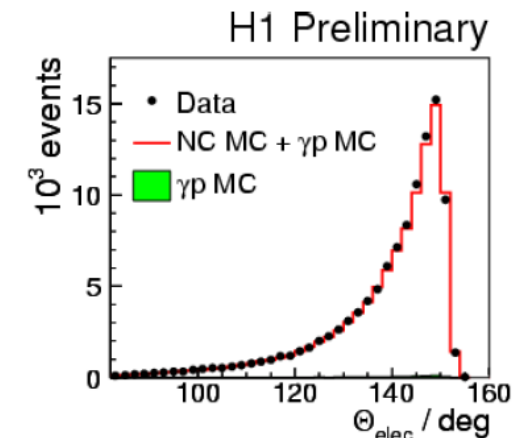
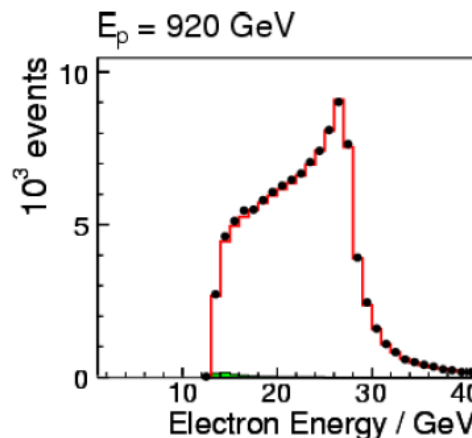
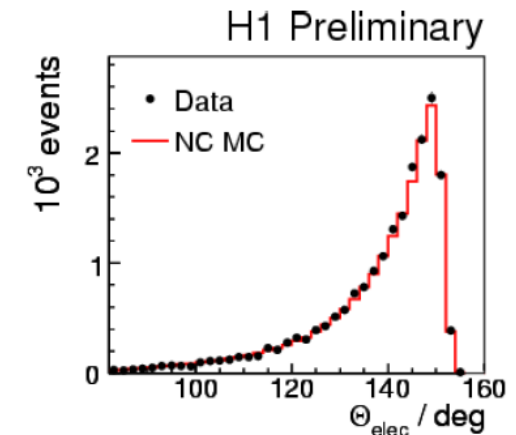
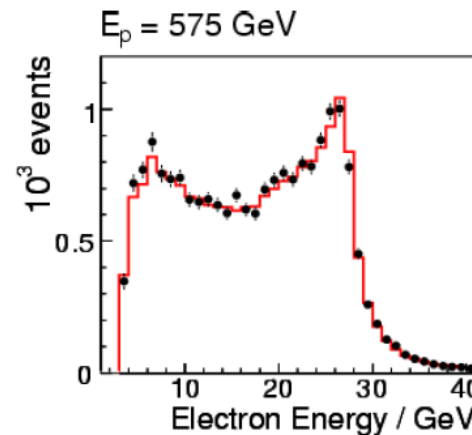
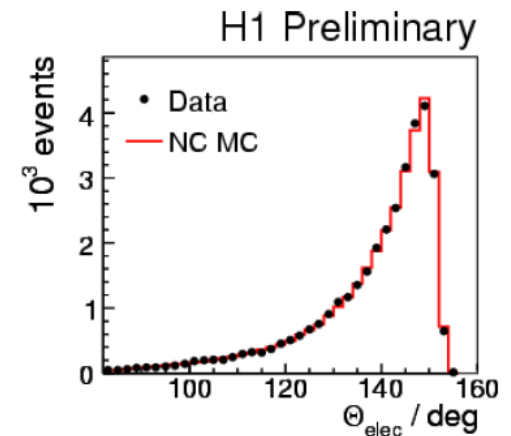
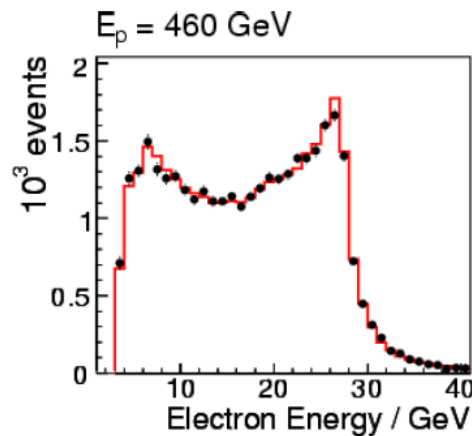
→ step at  $E_e = 6$  GeV is due to selection requirements

**the requirement  $E\text{-}P_z > 35$  GeV :**

- rejects  $\gamma p$  background
- rejects initial state radiation (ISR)

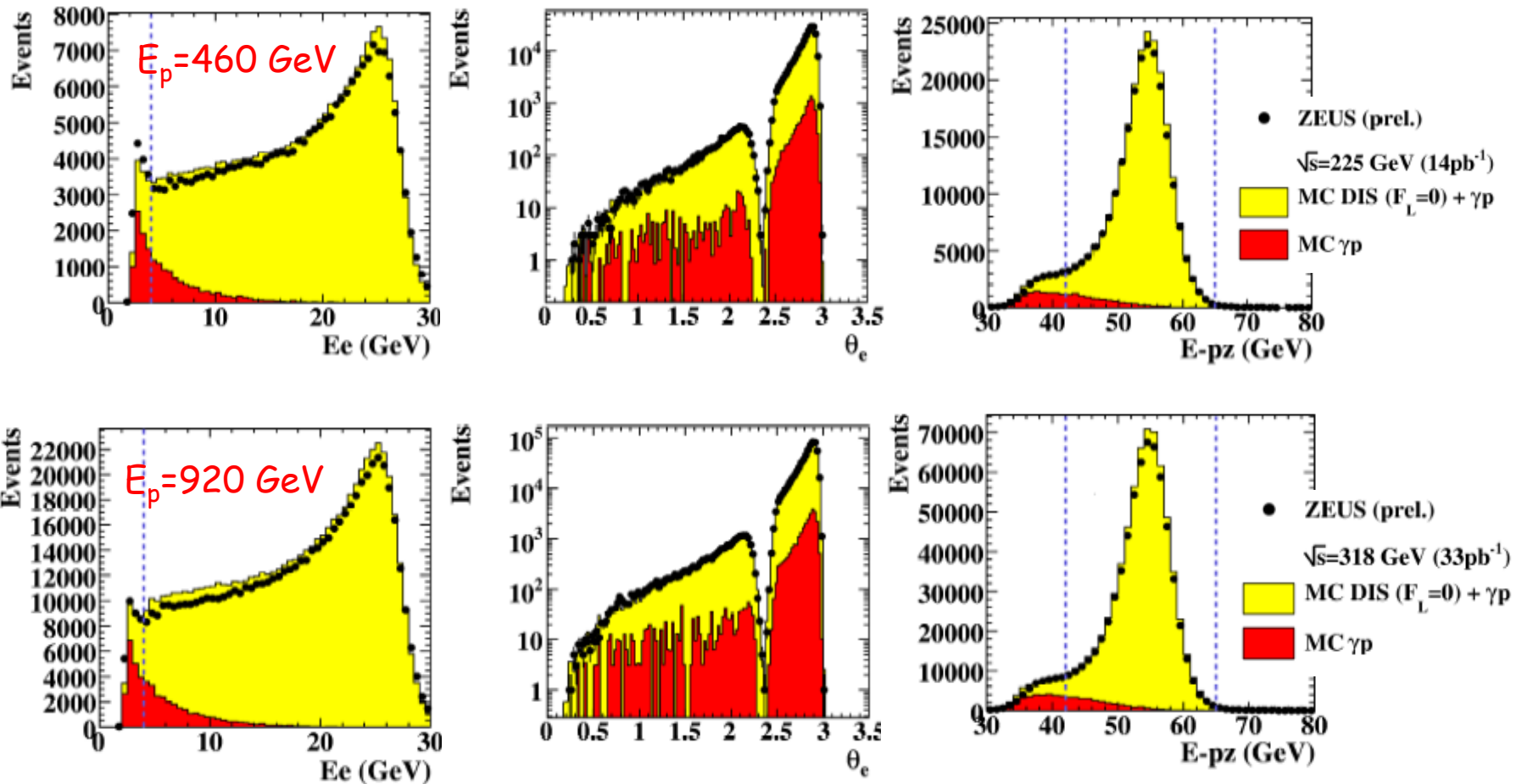
# Full $y$ range at high $Q^2$ after $\gamma p$ background subtraction (H1)

for  $E_p = 920 \text{ GeV}$  ( $y < 0.56$ )  
 $\gamma p$  bkg is taken from PYTHIA MC  
checked using 6m electron tagger



# ZEUS: control plots ( $E_p = 460, 920$ GeV)

MC is shown without  $F_L$  contribution

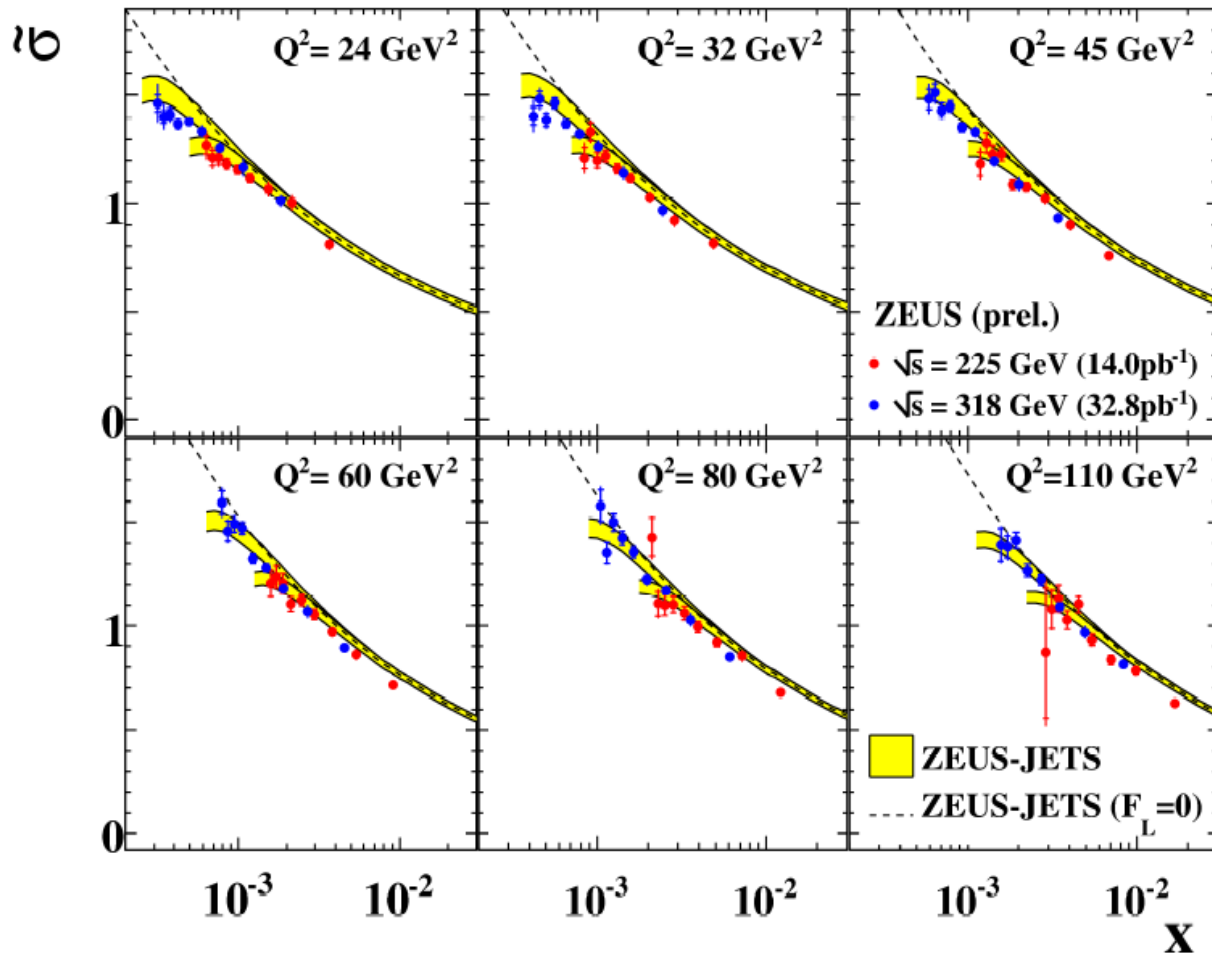


cuts indicated by lines:  $E'_e > 6$  GeV  
 $42 < E-p_z < 65$  GeV



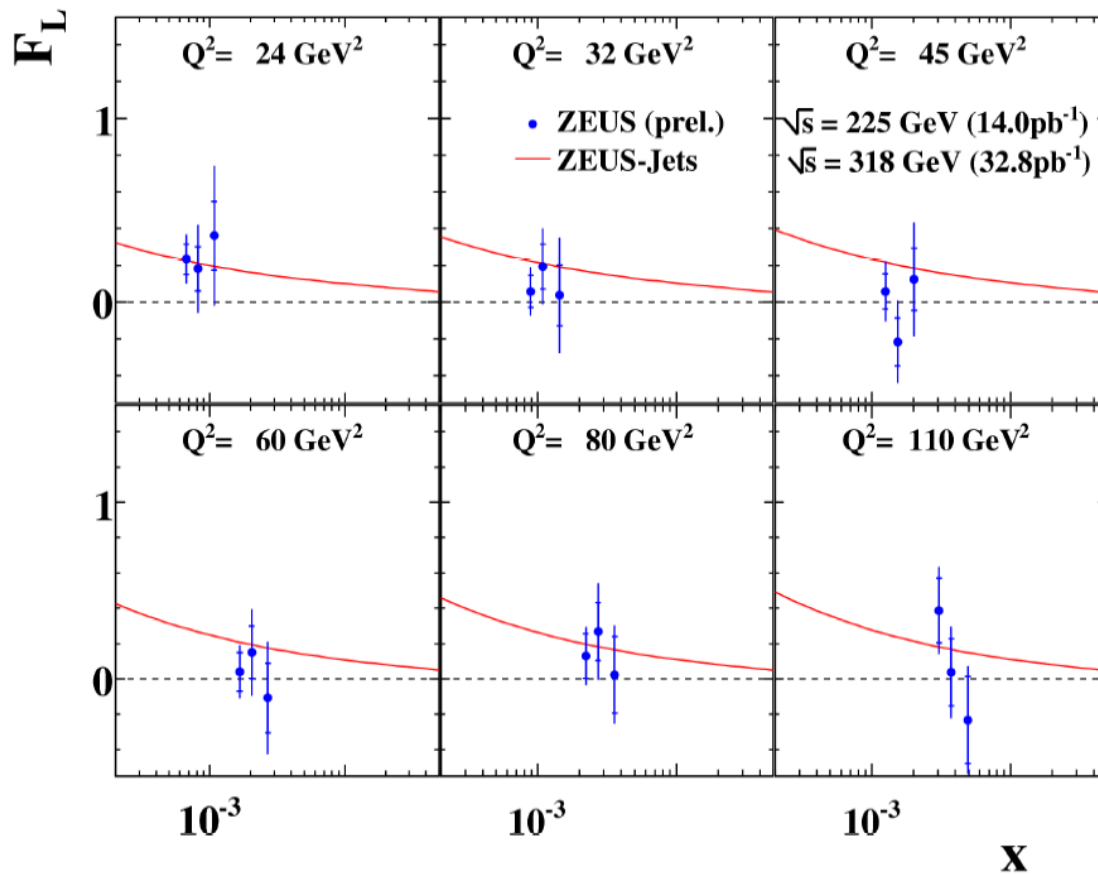
# NC cross sections for $E_p = 460, 920$ GeV

## ZEUS



# $F_L(x, Q^2)$ from ZEUS

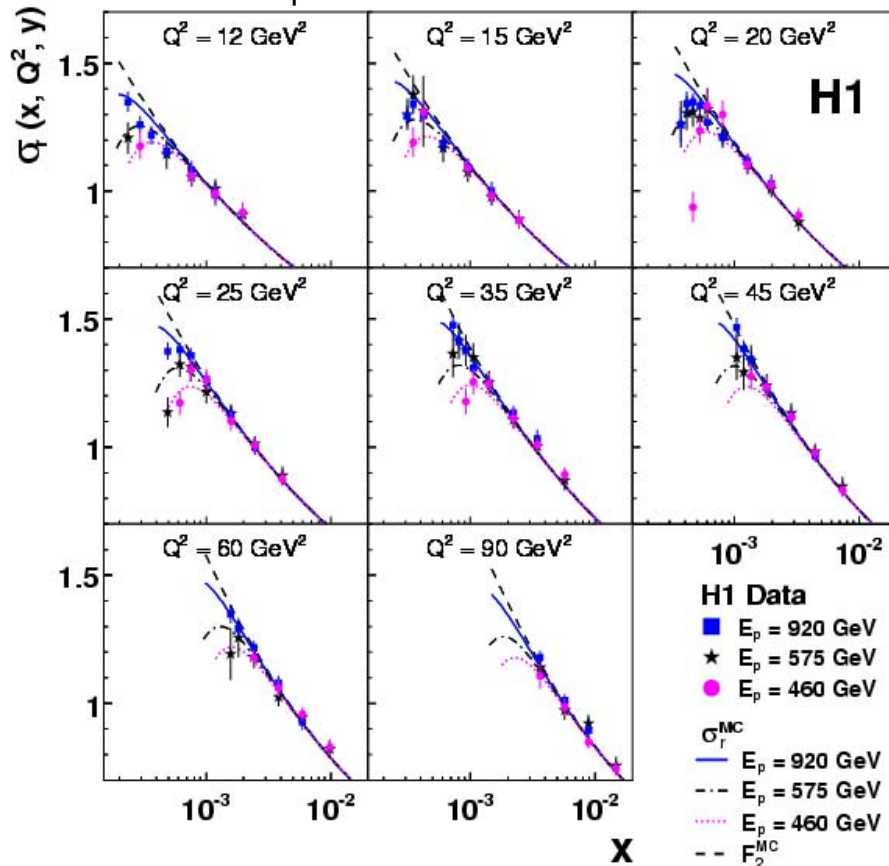
**ZEUS**



$F_L$  measurements are consistent within errors with QCD calculations and with  $F_L=0$

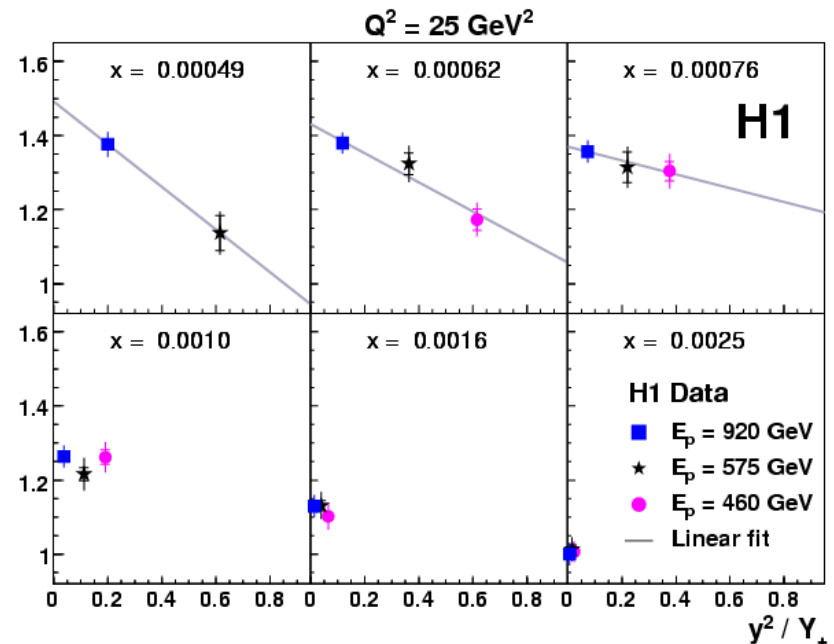
# NC cross sections at medium $Q^2$ (H1)

$E_p = 460, 575, 920$  GeV



$$\tilde{\sigma}_{NC} = F_2 - \frac{y^2}{1 + (1 - y)^2} F_L$$

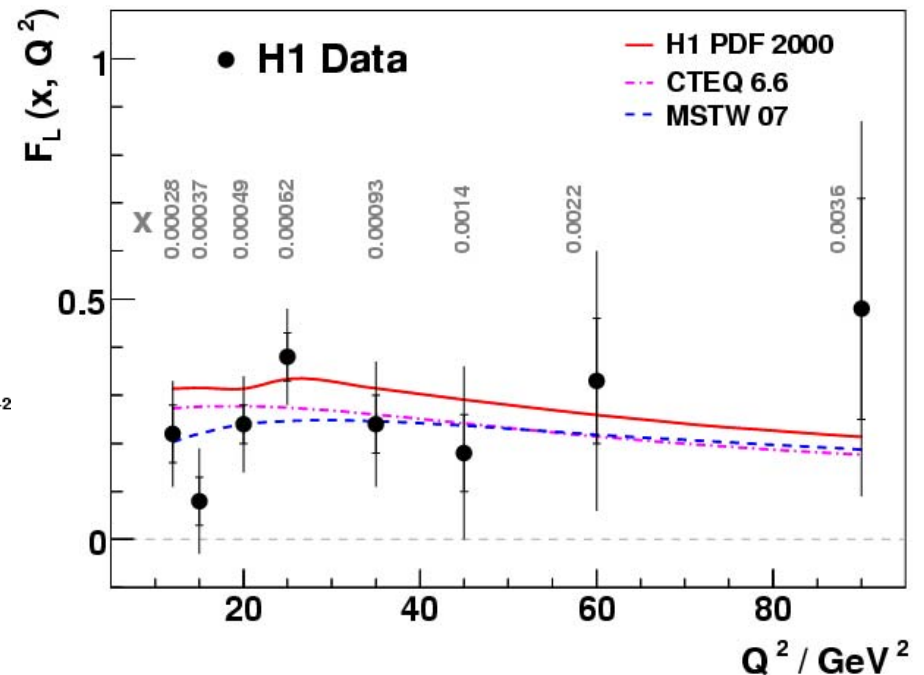
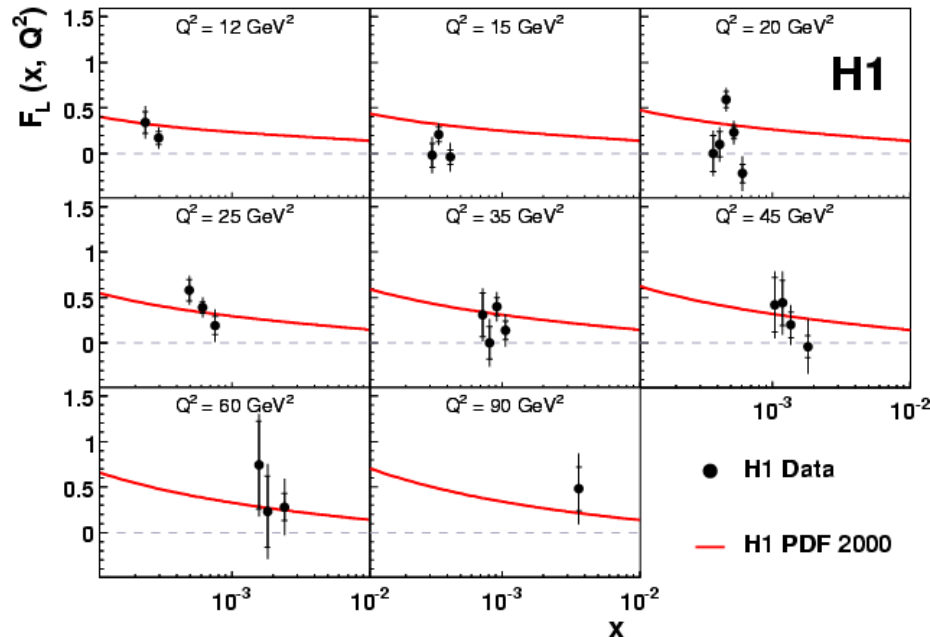
→ determine  $F_L$  and  $F_2$  from linear fits at each  $x$  and  $Q^2$



→ use relative normalisation (the same for LAr and Spacal) of  $E_p = 460, 575, 920$  GeV from the low  $y$  data for the  $F_L$  measurement

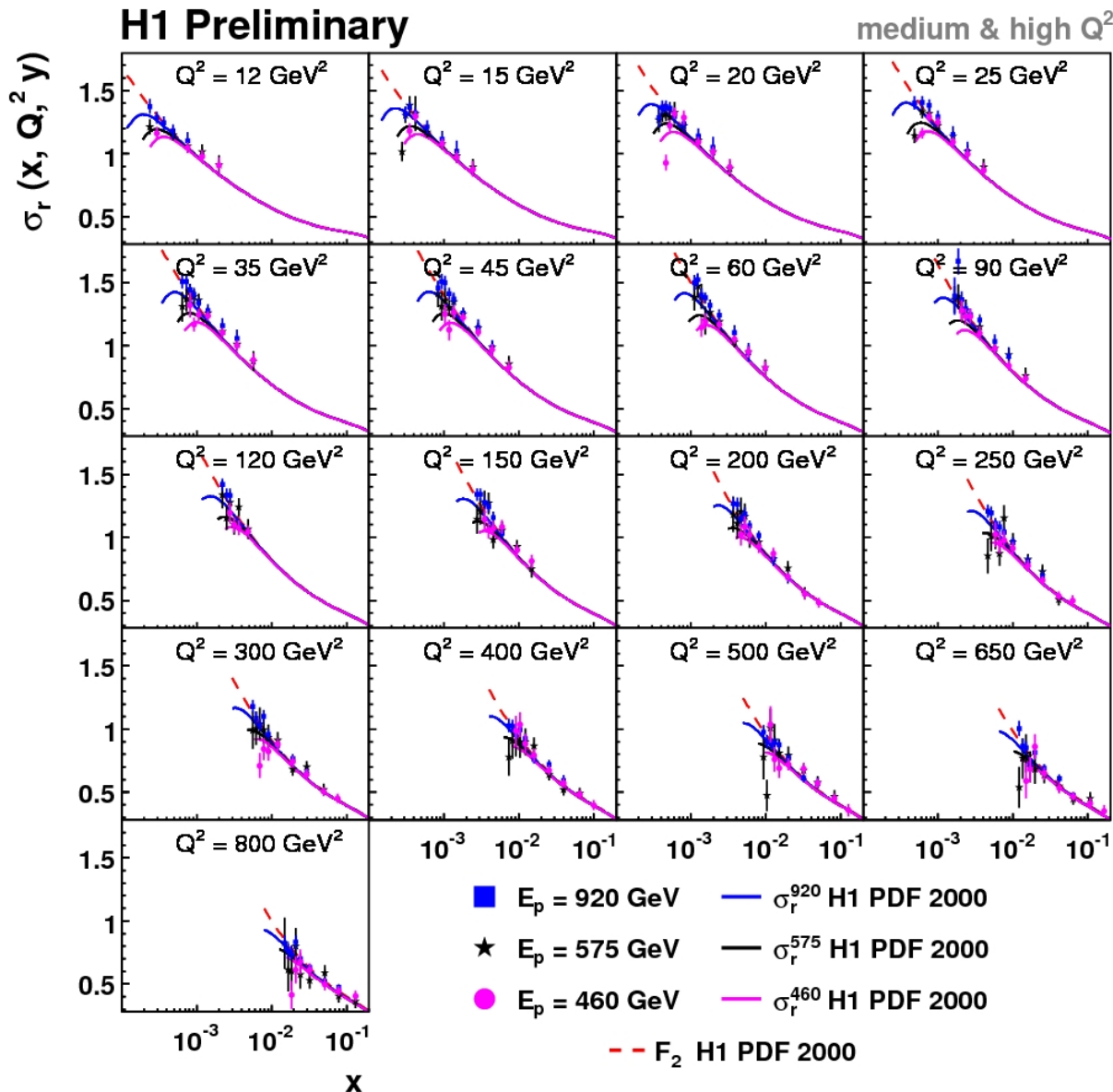
# The published $F_L(x, Q^2)$ and averaged $F_L(Q^2)$ at medium $Q^2$ (H1)

DESY-08-053



→ measured  $F_L$  are above zero and consistent with QCD calculations

# NC cross section in the full $Q^2$ range (H1)

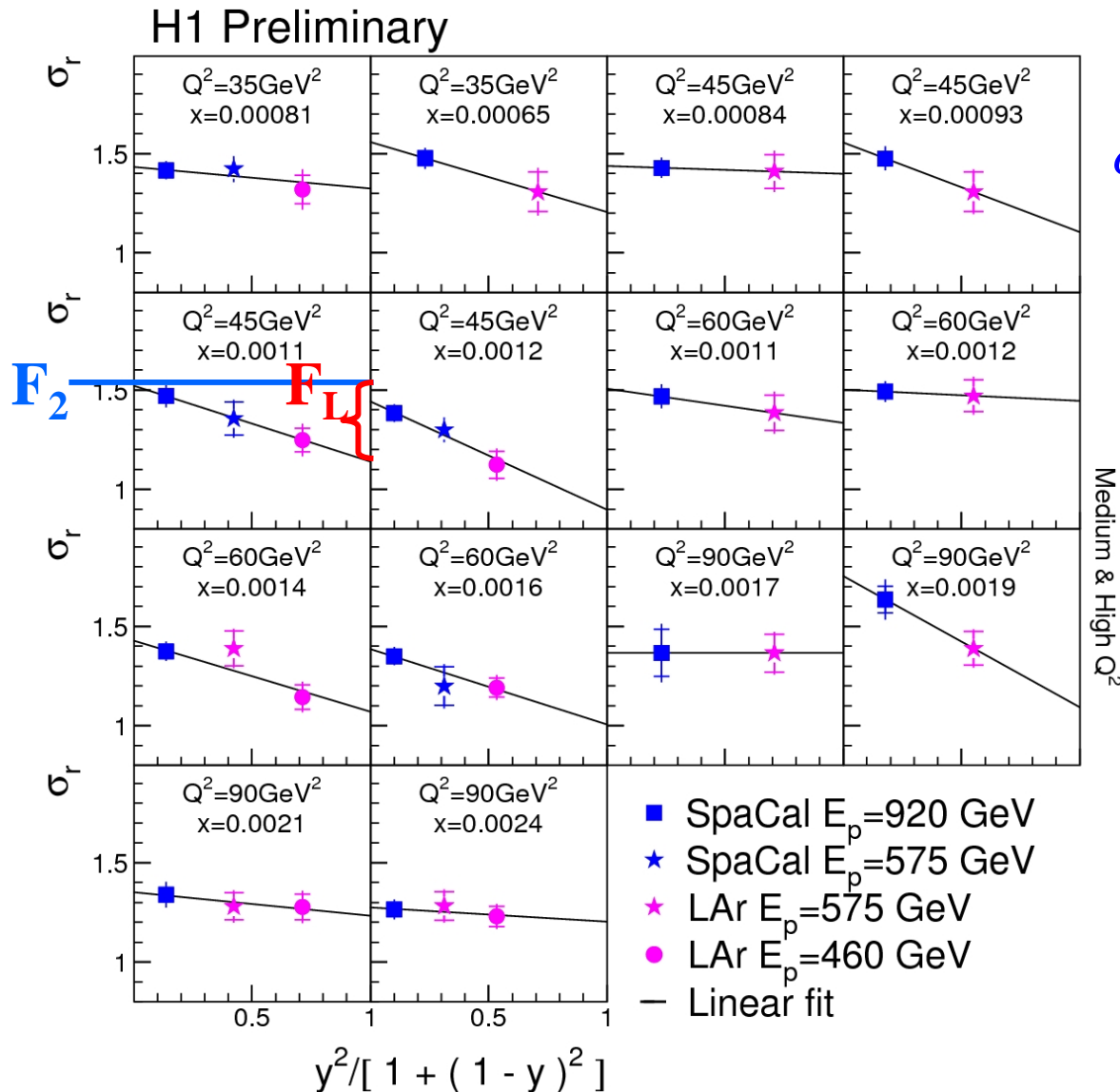


The full range of medium and high  $Q^2$  obtained using Spacal and LAr data

$$E_p = 460, 575, 920 \text{ GeV}$$

use relative normalisation (the same for LAr and Spacal) of  $E_p = 460, 575, 920$  GeV from the low  $y$  data for the  $F_L$  measurement

# NC cross sections at the same x & Q<sup>2</sup> which involve both the LAr and Spacal data (H1)



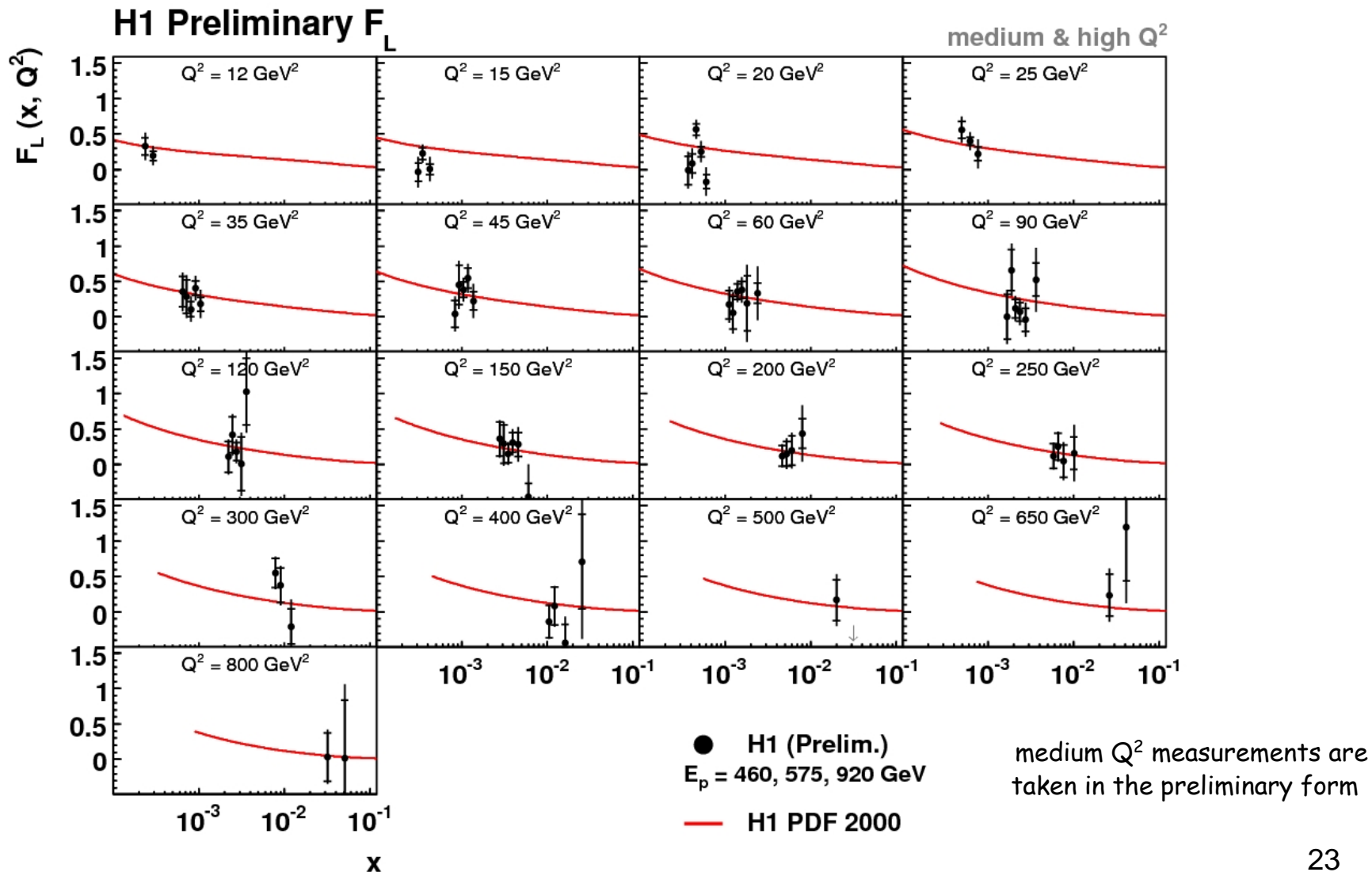
$$\tilde{\sigma}_{NC} = F_2 - \frac{y^2}{1+(1-y)^2} F_L$$

From linear fits  
at each x and Q<sup>2</sup>  
one determines  
F<sub>L</sub> and F<sub>2</sub>

blue points - Spacal  
red points - LAr

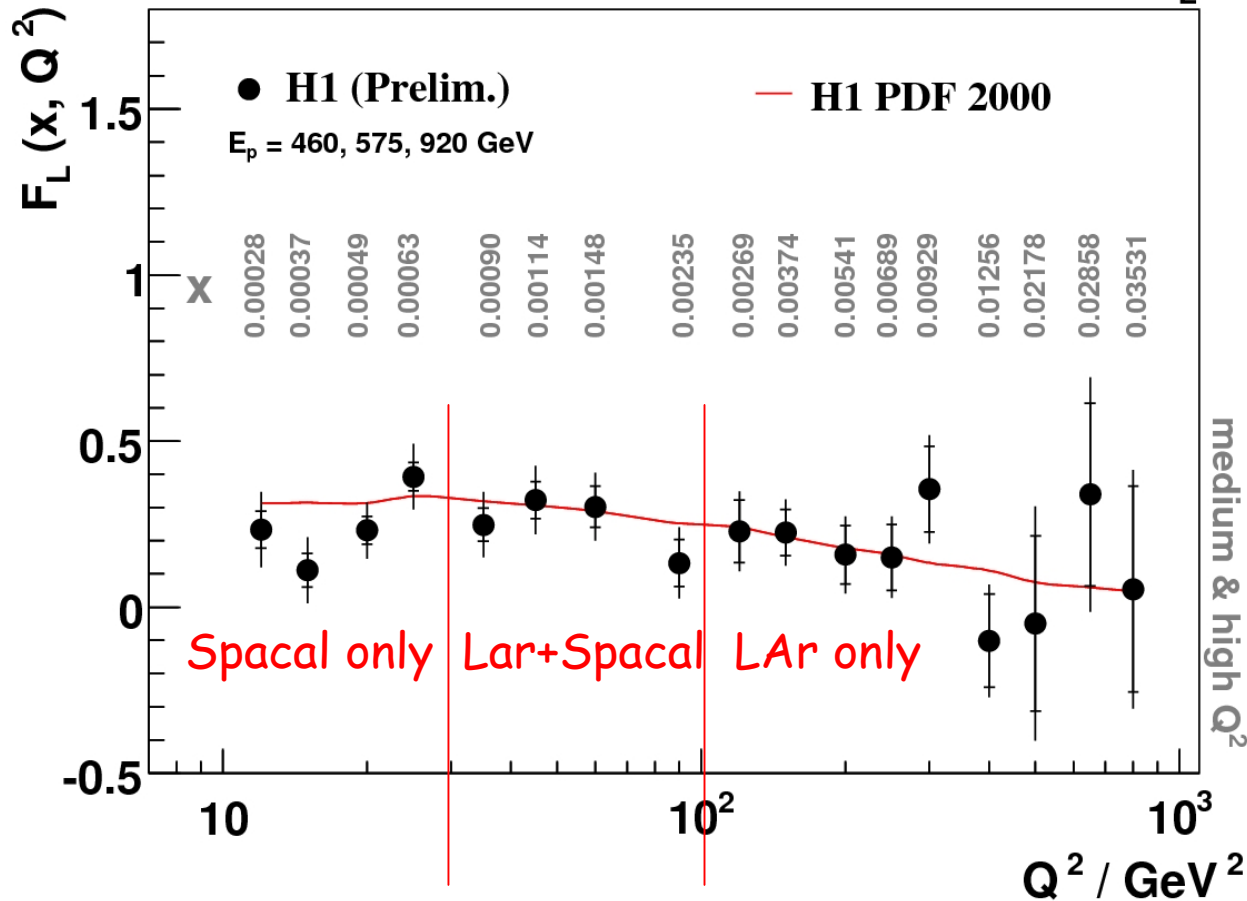
→ nice interplay of the  
two fully independent  
analyses using different  
detectors: Lar and Spacal

# $F_L(x, Q^2)$ in the full $Q^2$ range using the LAr and Spacal data (H1)



# Averaged $F_L(Q^2)$ in the full $Q^2$ range (H1)

H1 Preliminary  $F_L$



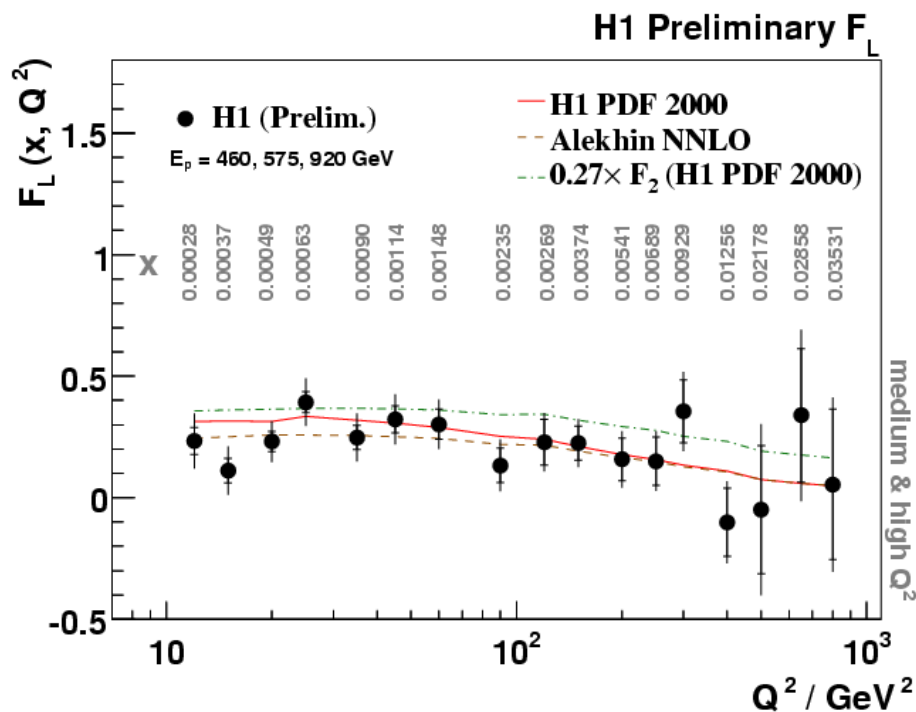
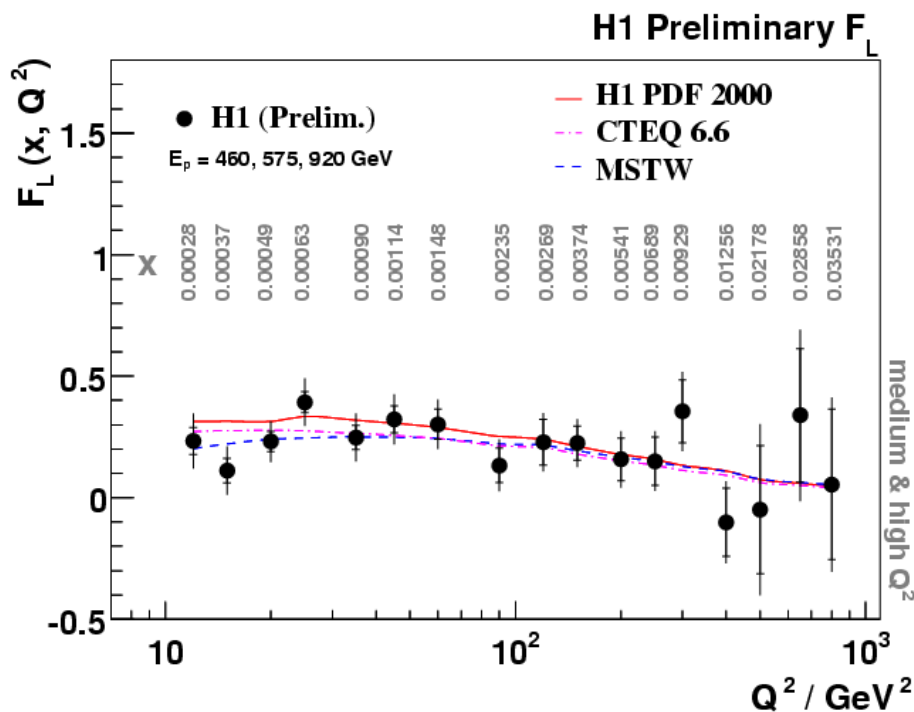
→ Spacal and LAR provide a cross check of the  $F_L$  measurements

→ overall correlated systematics between  $F_L$  points is  $\delta F_L \approx 0.05-0.10$

medium  $Q^2$  measurements are taken in the preliminary form



# Comparison of $F_L$ from H1 with recent theory predictions



$F_L = 0.27 \times F_2$  is motivated by  
 Schildknecht et al. arXiv:0806.0202

- $F_L$  measurements are in a good agreement with the NLO/NNLO QCD calculations
- extension to  $Q^2 < 10 \text{ GeV}^2$  will provide an important constraint

# Summary

*The longitudinal structure function  $F_L(x, Q^2)$  is measured at HERA in a model independent way using low  $E_p$  data*

H1:

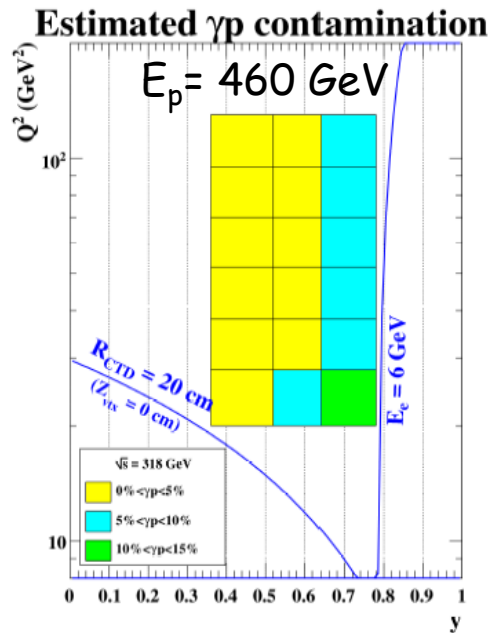
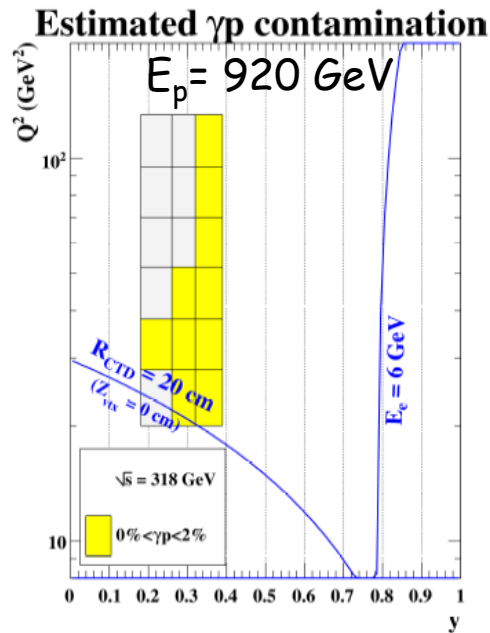
- measured at medium and high  $Q^2$  :  $12 \leq Q^2 \leq 800 \text{ GeV}^2$   
using the  $e^+p$  2007 data collected with  $E_p = 460, 575$  and  $920 \text{ GeV}$
- nice interplay of the two fully independent analyses which use two different detectors: LAr and Spacal
- measured  $F_L(x, Q^2)$  is in agreement with the recent theoretical calculations in the QCD framework

ZEUS:

- measured in the range  $24 \leq Q^2 \leq 110 \text{ GeV}^2$   
using the  $e^+p$  2007 data collected with  $E_p = 460$  and  $920 \text{ GeV}$
- measured  $F_L(x, Q^2)$  consistent within errors with QCD calculations but also with  $F_L=0$

→ *more to come:  $F_L$  at  $Q^2 < 10 \text{ GeV}^2$  (H1), analysis of  $E_p=575 \text{ GeV}$  data (ZEUS),  $F_L^D$ , ...*

# Experimental challenge: $\gamma p$ bkg at high $y$



ZEUS:

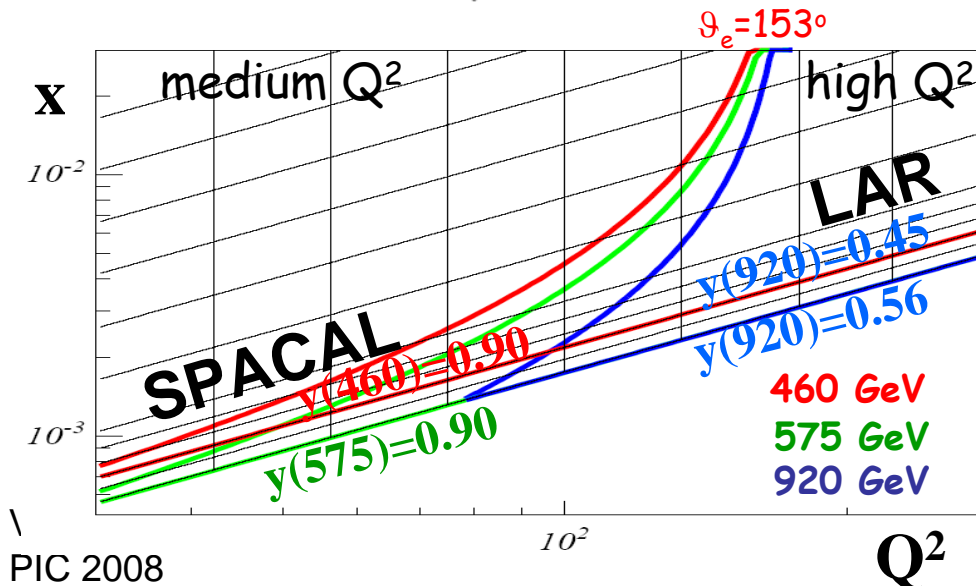
$\gamma p$  background contribution  
in the  $Q^2$ - $y$  bins used for FL

< 2% for  $E_p = 920 \text{ GeV}$  ( $y < 0.40$ )  
10-15% for  $E_p = 460 \text{ GeV}$  ( $y \approx 0.76$ )

H1:

the same binning in  $x$  and  $Q^2$   
for all  $E_p$  and LAr/Spacal

- measurements up to  $y = 0.90$   
where  $\gamma p$  bkg is up to 50% and more



ERA

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# Electron identification & background suppression at high $y$

Electron is identified by compactness of the cluster in calorimeter and track pointing to the cluster.

further reduction of  $\gamma p$  background keeping high eff. for electron:

## *Spacal sample*

- distance between extrapolated track and the electron cluster  $D < 6$  cm
- energy fraction behind the electron cluster  $E_h/E_e < 0.15$

## *LAR sample at $E_e < 6$ GeV*

- small transverse size of the electron cluster in LAr:  $E_{cra} < 4$  cm
- matching between track momentum and cluster energy:  $0.7 < E_t^{cluster}/P_t^{track} < 1.5$

