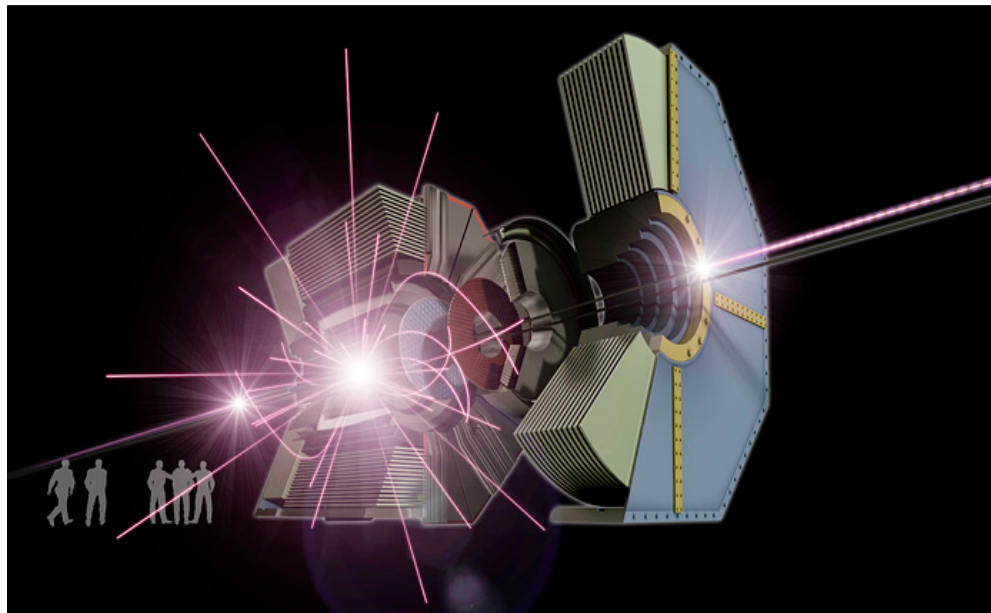
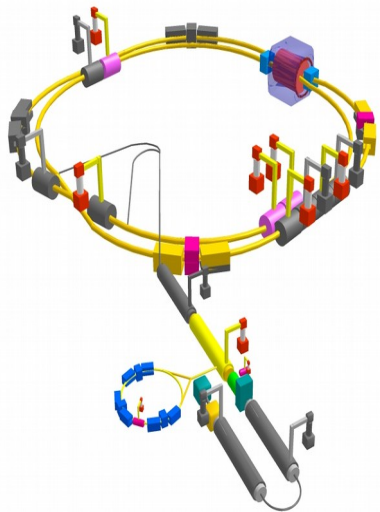


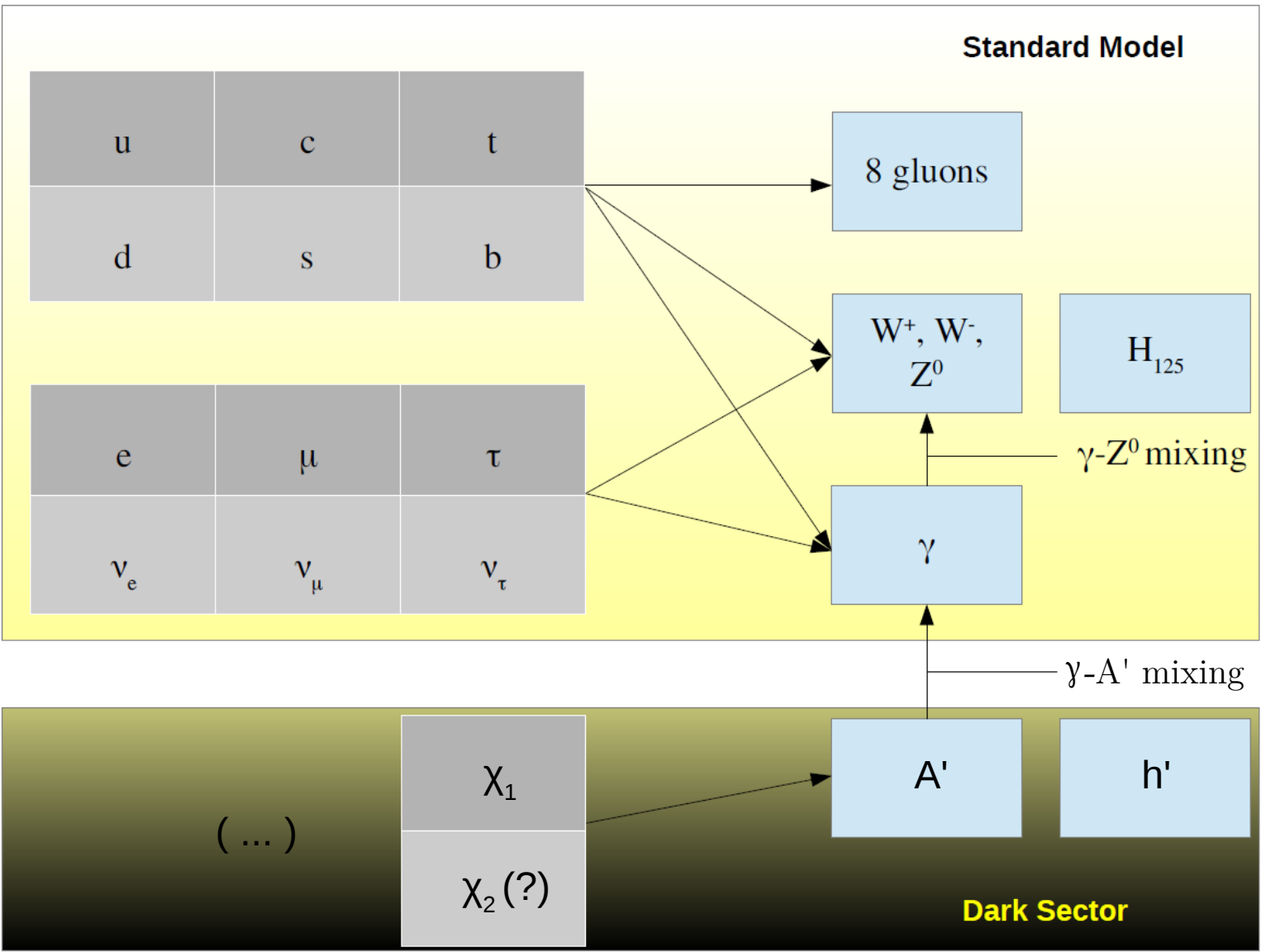
Prospects for low mass dark matter and T/CPT tests

JENNIFER Consortium General Meeting
KEK 06/10/2017

Gianluca Inguglia



Dask sector: how does it look like?



Low mass dark matter @ Belle (II): $Y(1S) \rightarrow$ invisible

$Y(nS)$: bound state of a b quark and a b antiquark

$$\frac{BR(Y(1S) \rightarrow \nu \bar{\nu})}{BR(Y(1S) \rightarrow e^+ e^-)} = \frac{27 G^2 M_{Y(1S)}^4}{64 \pi^2 \alpha^2} \left(-1 + \frac{4}{3} \sin^2 \theta_W\right)^2 = 4.14 \times 10^{-4}$$

$$BR(Y(1S) \rightarrow \nu \bar{\nu}) \sim 9.9 \times 10^{-6}$$

→ Low mass dark matter particles however might play a role in the decays of $Y(1S)$, having $Y(1S) \rightarrow \chi\chi$ if kinematic allowed. [Phys. Rev. D **80**, 115019, 2009]

→ Also, new mediators (Z' , A^0 , h^0) or SUSY particles might enhance $Y(1S) \rightarrow \nu\nu(\gamma)$. [Phys. Rev. D **81**, 054025, 2010]

→ In absence of new physics enhancement, Belle2 should be able to observe the SM $Y(1S) \rightarrow \nu\nu$

→ $e^+ e^- \rightarrow Y(3S)$
 \downarrow (4.4%)
 $Y(3S) \rightarrow \pi^+ \pi^- Y(1S)$
 \downarrow
 $Y(1S) \rightarrow$ invisible

→ $e^+ e^- \rightarrow Y(2S)$
 \downarrow (18.1%)
 $Y(2S) \rightarrow \pi^+ \pi^- Y(1S)$
 \downarrow
 $Y(1S) \rightarrow$ invisible

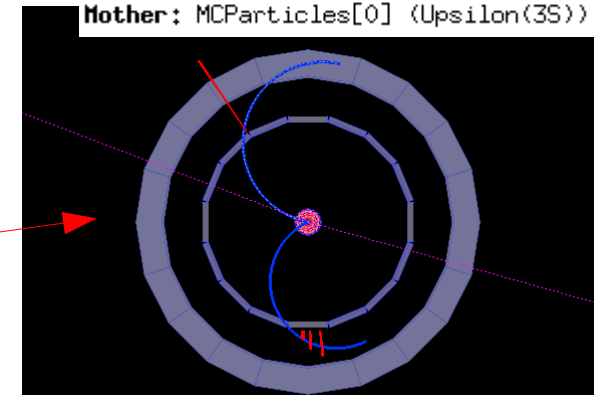
Belle2 Simulation

$Y(3S) \rightarrow \pi^+ \pi^- Y(1S)$,

$Y(1S) \rightarrow \nu\nu$

Charge=1, PDG=211 (pi+)
 pT=0.420365, pZ=0.000692372
 V=(-0.00, -0.00, -0.03)

Mother: MCParticles[0] (Upsilon(3S))



Charge=-1, PDG=-211 (pi-)

pT=0.344016, pZ=0.118851
 V=(-0.00, -0.00, -0.03)

Mother: MCParticles[0] (Upsilon(3S))

$$M_{Y(3S)} = 10.355 \text{ GeV}/c^2, \quad M_{Y(2S)} = 10.023 \text{ GeV}/c^2, \quad M_{Y(1S)} = 9.460 \text{ GeV}/c^2$$

$\sim 900 \text{ MeV}$ available for $P_{\pi\pi}$

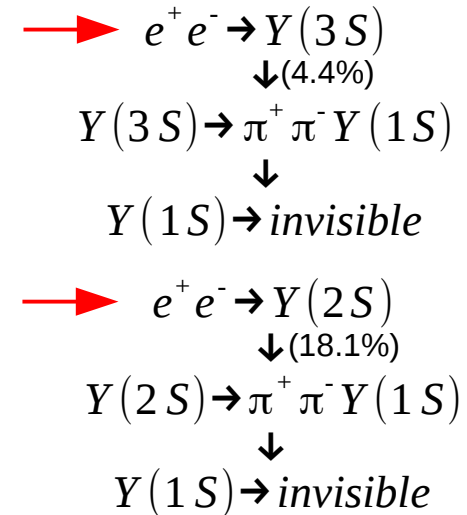
$\sim 540 \text{ MeV}$ available for $P_{\pi\pi}$

Low mass dark matter @ Belle (II): $Y(1S) \rightarrow \text{invisible}$

$$\frac{BR(Y(1S) \rightarrow \nu \bar{\nu})}{BR(Y(1S) \rightarrow e^+ e^-)} = \frac{27 G^2 M_{Y(1S)}^4}{64 \pi^2 \alpha^2} \left(-1 + \frac{4}{3} \sin^2 \theta_w\right)^2 = 4.14 \times 10^{-4}$$

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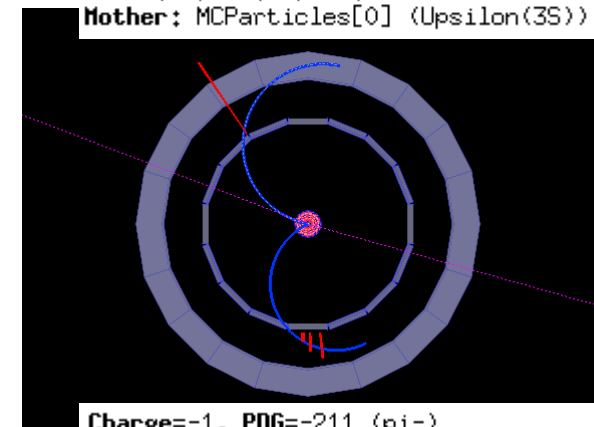


Belle2 Simulation

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$Y(1S) \rightarrow \nu\nu$

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Mother: MCParticles[0] (Upsilon(3S))
```

A signal of $Y(1S) \rightarrow \text{invisible}$ is an excess of events over the background in the M_r distribution at a mass equivalent to that of the $Y(1S)$ ($9.460 \text{ GeV}/c^2$)

$$M_r^2 = s + M_{\pi^+ \pi^-}^2 - 2 \sqrt{s} E_{\pi^+ \pi^-}^{CMS}$$

Low mass dark matter @ Belle II: $Y(1S) \rightarrow$ invisible

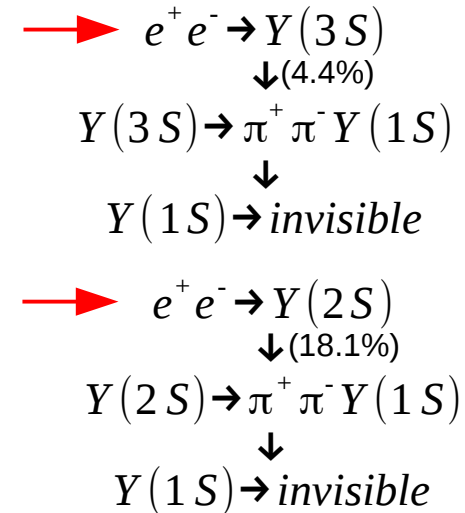
$$\frac{BR(Y(1S) \rightarrow \nu \bar{\nu})}{BR(Y(1S) \rightarrow e^+ e^-)} = \frac{27 G^2 M_{Y(1S)}^4}{64 \pi^2 \alpha^2} \left(-1 + \frac{4}{3} \sin^2 \theta_W\right)^2 = 4.14 \times 10^{-4}$$

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→ Also, new mediators (Z' , A^0 , h^0) or SUSY particles might enhance $Y(1S) \rightarrow \nu\nu(\gamma)$. **[Phys. Rev. D 81, 054025, 2010]**

→ In absence of new physics enhancement, Belle2 should be able to strongly constrain the SM $Y(1S) \rightarrow \nu\nu$

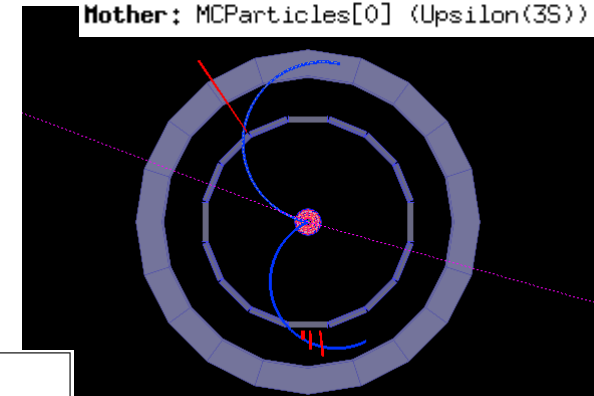


Belle2 Simulation

$Y(3S) \rightarrow \pi^+ \pi^- Y(1S)$,

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Mother: MCParticles[0] (Upsilon(3S))
```

No signal was observed over the expected background and upper limits have been obtained: $BR(Y \rightarrow \nu\nu) < 3 \times 10^{-4}$ (BaBar) and $BR(Y \rightarrow \nu\nu) < 3.0 \times 10^{-3}$ (Belle).

| Process | $L_{int}(ab^{-1})$ | ϵ | $N(Y(1S))$ | $N_{Y(1S) \rightarrow \nu\bar{\nu}}$ | N_{NP} |
|---|----------------------|------------|-------------------|--------------------------------------|------------|
| $\Upsilon(2S) \rightarrow \pi^+ \pi^- \Upsilon(1S)$ | 0.2, $\Upsilon(2S)$ | 0.1-0.2 | 2.3×10^8 | 230-460 | 6900-13800 |
| $\Upsilon(3S) \rightarrow \pi^+ \pi^- \Upsilon(1S)$ | 0.2, $\Upsilon(3S)$ | 0.1-0.2 | 3.2×10^7 | 32-64 | 945-1890 |
| $\Upsilon(4S) \rightarrow \pi^+ \pi^- \Upsilon(1S)$ | 50.0, $\Upsilon(4S)$ | 0.1-0.2 | 5.5×10^6 | 5.5-11 | 165-310 |
| $\Upsilon(5S) \rightarrow \pi^+ \pi^- \Upsilon(1S)$ | 5.0, $\Upsilon(5S)$ | 0.1-0.2 | 7.6×10^6 | 7.6-15.2 | 228-456 |
| $\gamma_{ISR} \Upsilon(2S) \rightarrow (\gamma_{ISR}) \pi^+ \pi^- \Upsilon(1S)$ | 50.0, $\Upsilon(4S)$ | 0.1-0.2 | 1.5×10^8 | 150-300 | 4500-9000 |
| $\gamma_{ISR} \Upsilon(3S) \rightarrow (\gamma_{ISR}) \pi^+ \pi^- \Upsilon(1S)$ | 50.0, $\Upsilon(4S)$ | 0.1-0.2 | 3.5×10^7 | 35-70 | 1050-2100 |

Low mass dark matter @ Belle II: $\Upsilon(1S) \rightarrow$ invisible

| Process | $L_{int}(ab^{-1})$ | ϵ | $N(\Upsilon(1S))$ | $N_{\Upsilon(1S) \rightarrow \nu\bar{\nu}}$ | N_{NP} |
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If we collect $>200\text{fb}^{-1}$ of data @ $\Upsilon(3S)$ [$\Upsilon(2S)$] we should reconstruct between 30 and 300 [~ 200 and ~ 2000] events, assuming 10^{-5} (SM) $< \text{BR}_{\Upsilon \rightarrow \text{invisible}} < 10^{-4}$ (NP) and $\epsilon_{\text{tot}} = 10\%$.

However trigger is not yet taken into considerations and preliminary studies have shown that the trigger (only) efficiency currently is of $\sim 50\%$.

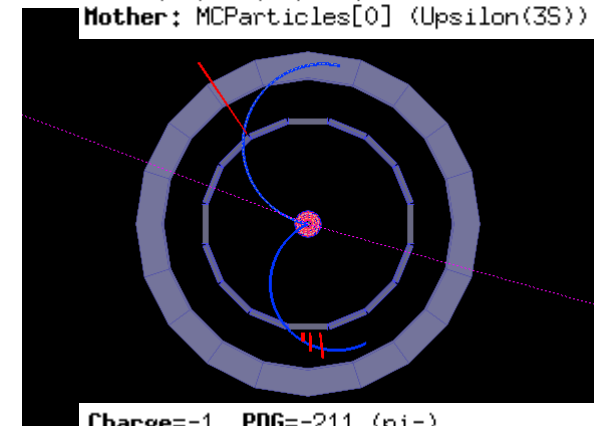
Work also during phase II data taking is needed to perform test on exotic triggers (for the two low momentum π 's).

Another thing that need to be clarified is the data taking at other energies than $\Upsilon(4S)$. If the luminosity is too high the bkgd levels increase as well and triggers have to be scaled accordingly, this is in conflict with the need of low momentum pions trigger: data taking must happen in the early period of phase 3.

Belle2 Simulation

$\Upsilon(3S) \rightarrow \pi^+\pi^-\Upsilon(1S)$,
 $\Upsilon(1S) \rightarrow \nu\bar{\nu}$

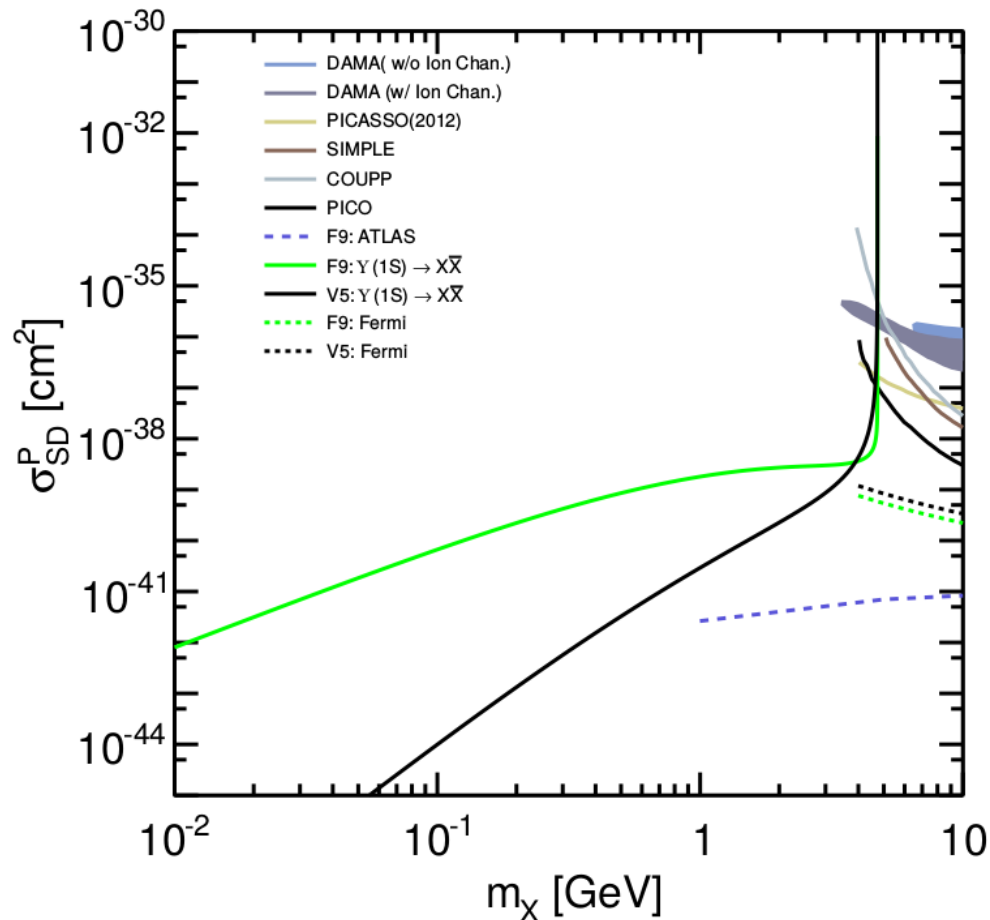
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pT=0.420365, pZ=0.000692372
V=(-0.00, -0.00, -0.03)
Mother: MCParticles[0] (Upsilon(3S))
```



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Charge=-1, PDG=-211 (pi-)
pT=0.344016, pZ=0.118851
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Mother: MCParticles[0] (Upsilon(3S))
```

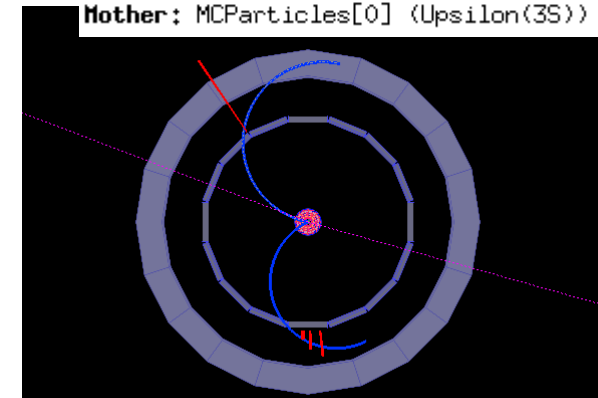
Low mass dark matter @ Belle II: $\Upsilon(1S) \rightarrow \text{invisible}$

| Process | $L_{int}(ab^{-1})$ | ϵ | $N(\Upsilon(1S))$ | $N_{\Upsilon(1S) \rightarrow \nu\bar{\nu}}$ | N_{NP} |
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 Mother: MCParticles[0] (Upsilon(3S))

Low mass dark matter @ Belle II: a dark Z'

Partial width results and BR not published, derived from eqn. 2.12 of Essig et al. JHEP02(2015)157, arXiv:1412.0018 [hep-ph] (contact: Brian Shuve).

The model is a new gauge boson, Z' , which couples to $L_\mu - L_\tau$. The interaction Lagrangian is

$$\mathcal{L} = -g' \bar{\mu} \gamma^\mu Z'_\mu \mu + g' \bar{\tau} \gamma^\mu Z'_\mu \tau - g' \bar{\nu}_{\mu,L} \gamma^\mu Z'_\mu \nu_{\mu,L} + g' \bar{\nu}_{\tau,L} \gamma^\mu Z'_\mu \nu_{\tau,L}.$$

The equations for the partial widths are,

$$\Gamma(Z' \rightarrow \ell^+ \ell^-) = \frac{(g')^2 M_{Z'}}{12\pi} \left(1 + \frac{2M_\ell^2}{M_{Z'}^2}\right) \sqrt{1 - \frac{4M_\ell^2}{M_{Z'}^2}} \theta(M_{Z'} - 2M_\ell),$$

$$\Gamma(Z' \rightarrow \nu_\ell \bar{\nu}_\ell) = \frac{(g')^2 M_{Z'}}{24\pi}.$$

$$Br(Z' \rightarrow invisible) = \frac{2\Gamma(Z' \rightarrow \nu_l \bar{\nu}_l)}{2\Gamma(Z' \rightarrow \nu_l \bar{\nu}_l) + \Gamma(Z' \rightarrow \mu^+ \mu^-) + \Gamma(Z' \rightarrow \tau^+ \tau^-)}$$

1) The branching fraction to one neutrino species is half of the branching fraction to one charged lepton flavour. The reason is, of course, that the Z' only couples to left-handed neutrino chiralities whereas it couples to both left- and right-handed charged leptons.

For $M_{Z'} < 2M_\mu$ $Br(Z' \rightarrow invisible) = 1$.

For $2M_\mu < M_{Z'} < 2M_\tau$ $Br(Z' \rightarrow invisible) \sim 1/2$

For $M_{Z'} > 2M_\tau$ $Br(Z' \rightarrow invisible) \sim 1/3$

Low mass dark matter @ Belle II: a dark Z'

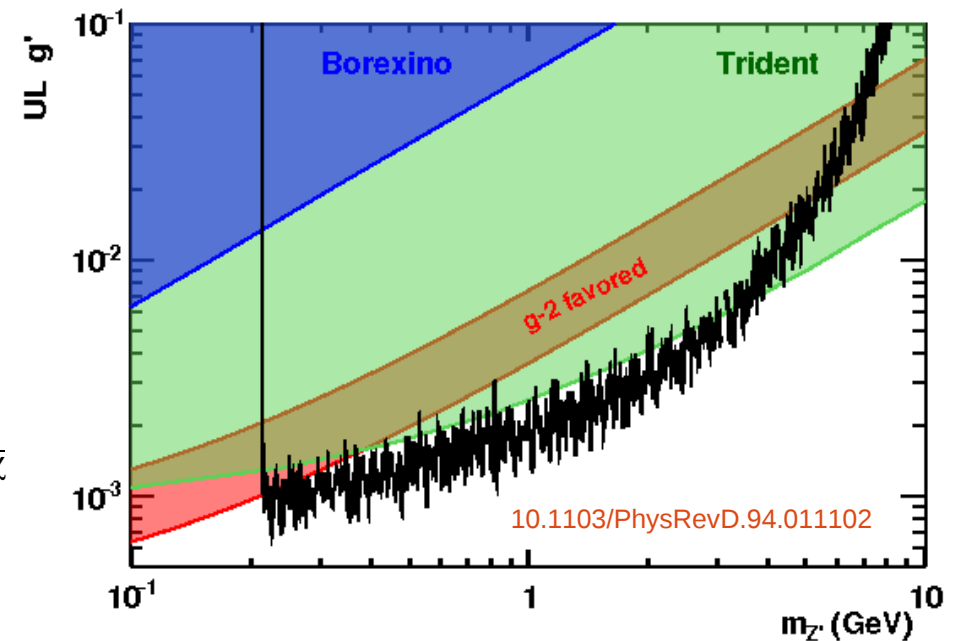
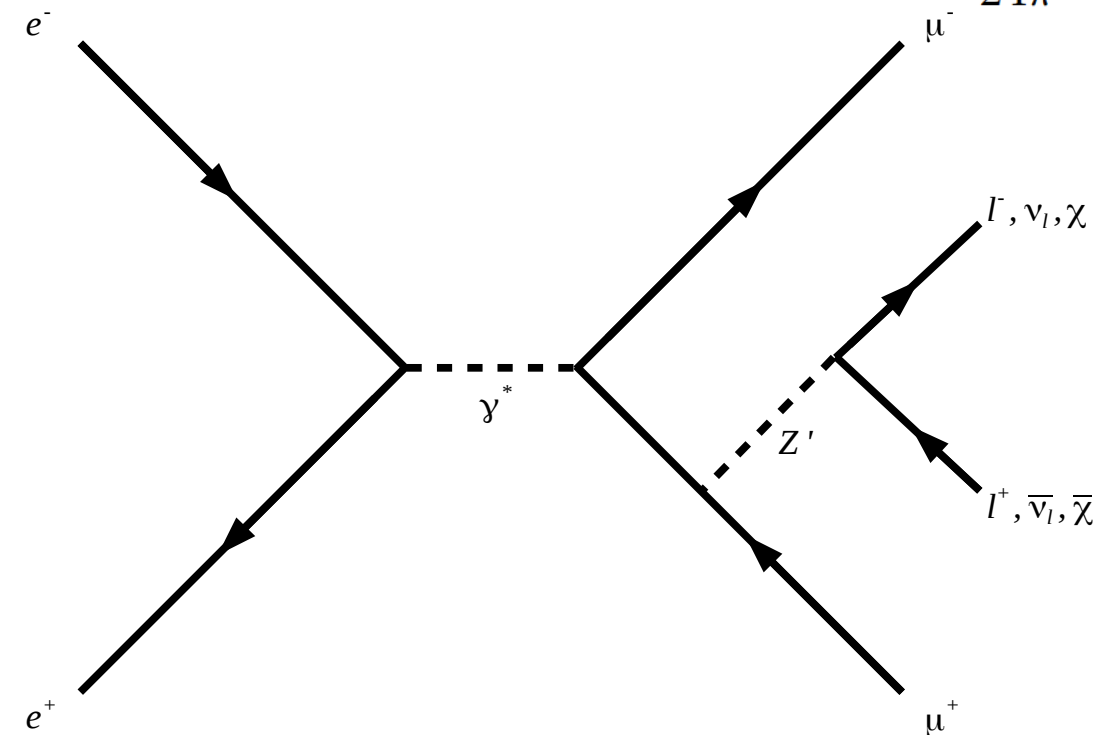
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Beside CP, can Belle II look for T/CPT non invariance?

1203.0930 Belle, CPT

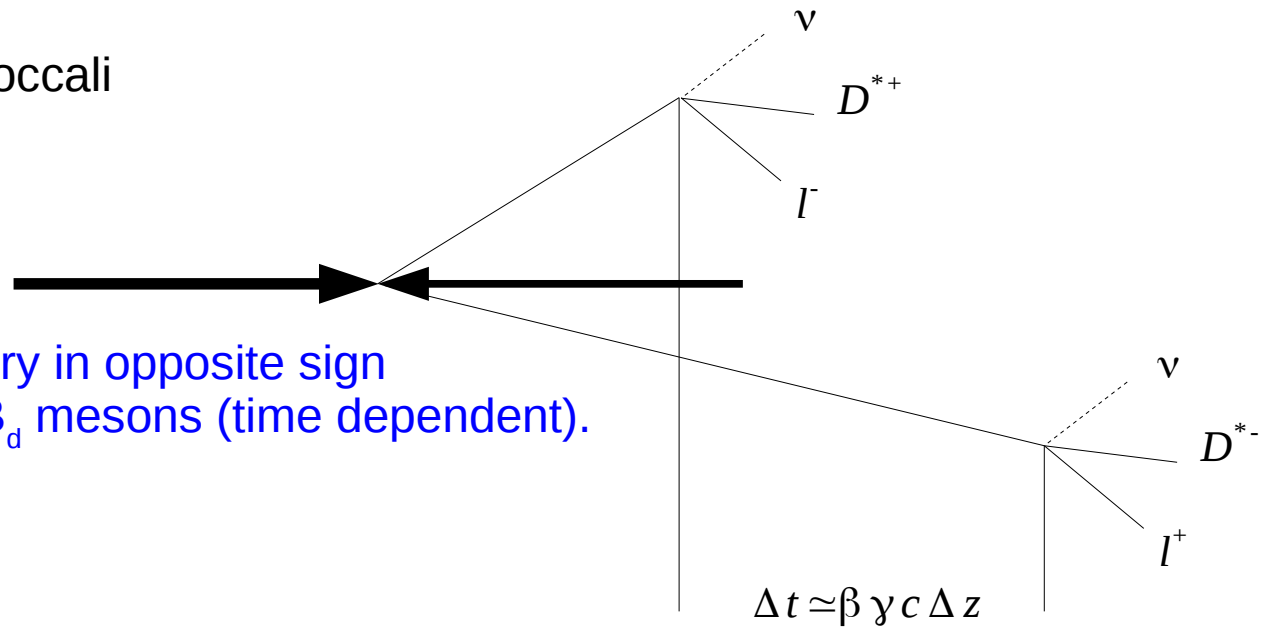
1207.5832 BaBar, T+CPT

1302.4191 A. Bevan, G.I., M. Zoccali

1605.04545 BaBar, CPT

Dual CP/CPT violating asymmetry in opposite sign dilepton events with entangled B_d mesons (time dependent).

$$\Psi = \frac{1}{\sqrt{2}} \left(B_1^0 \bar{B}_2^0 - \bar{B}_1^0 B_2^0 \right)$$



$$i \frac{d}{dt} \begin{pmatrix} |B_{1,H}\rangle \\ |B_{2,L}\rangle \end{pmatrix} = \begin{pmatrix} M_{11} - \frac{i}{2}\Gamma_{11} & M_{12} - \frac{i}{2}\Gamma_{12} \\ M_{12}^* - \frac{i}{2}\Gamma_{12}^* & M_{22}^* - \frac{i}{2}\Gamma_{22}^* \end{pmatrix} \begin{pmatrix} |B^0\rangle \\ |\bar{B}^0\rangle \end{pmatrix}$$

$$|B_{1,H}\rangle = p\sqrt{1-z}|B^0\rangle + q\sqrt{1+z}|\bar{B}^0\rangle$$

$$|B_{2,L}\rangle = p\sqrt{1-z}|B^0\rangle - q\sqrt{1+z}|\bar{B}^0\rangle$$

$$z = \frac{(m_{11} - m_{22}) - \frac{i}{2}(\Gamma_{11} - \Gamma_{22})}{\Delta m - \frac{i}{2}\Delta\Gamma}$$

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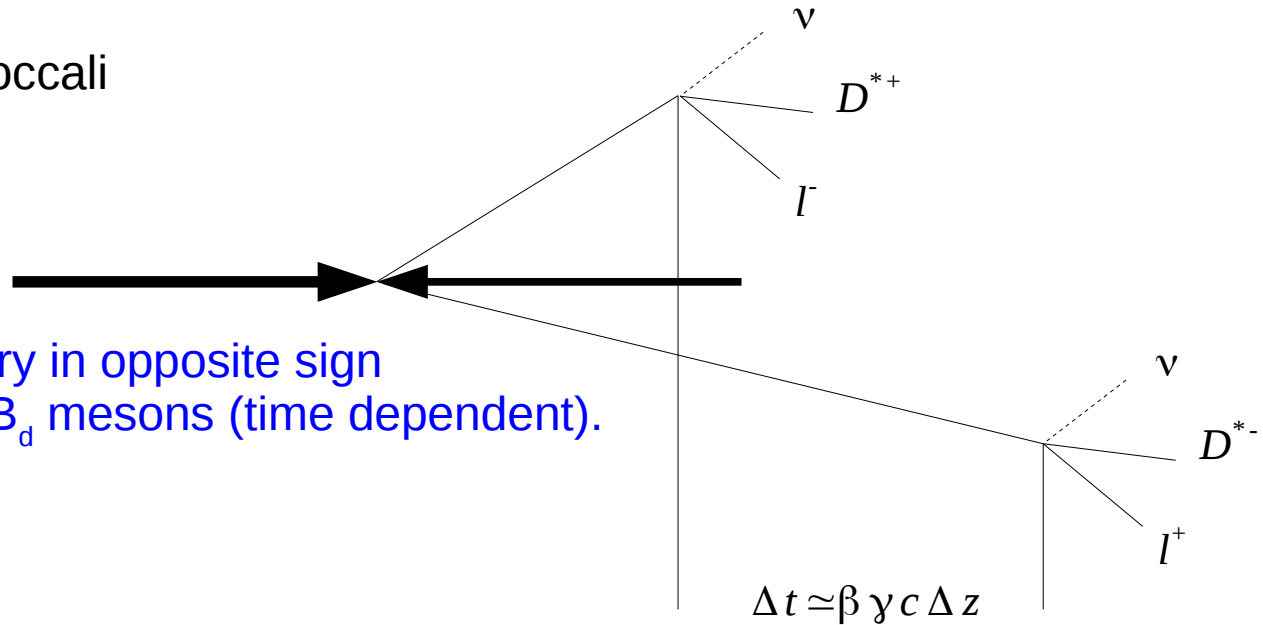
1207.5832 BaBar, T+CPT

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1605.04545 BaBar, CPT

Dual CP/CPT violating asymmetry in opposite sign di-lepton events with entangled B_d mesons (time dependent).

$$\Psi = \frac{1}{\sqrt{2}} \left(B_1^0 \bar{B}_2^0 - \bar{B}_1^0 B_2^0 \right)$$



$$N^{2l} \propto \frac{e^{-\Gamma|\Delta t|}}{2} \left\{ \cosh\left(\frac{\Delta\Gamma\Delta t}{2}\right) - 2\Re z \sinh\left(\frac{\Delta\Gamma\Delta t}{2}\right) + \cos(\Delta m\Delta t) + 2\Im z \sin(\Delta m\Delta t) \right\}$$

$$A_{CPT,CP}(|\Delta t|) = \frac{P(B^0 \rightarrow B^0) - P(\bar{B}^0 \rightarrow \bar{B}^0)}{P(B^0 \rightarrow B^0) + P(\bar{B}^0 \rightarrow \bar{B}^0)} = \frac{N^{2l}(\Delta t > 0) - N^{2l}(\Delta t < 0)}{N^{2l}(\Delta t > 0) + N^{2l}(\Delta t < 0)} \simeq \frac{2\Im z \sin(\Delta m\Delta t) - \Re z \sinh\left(\frac{\Delta\Gamma\Delta t}{2}\right)}{\cosh\left(\frac{\Delta\Gamma\Delta t}{2}\right) + \cos(\Delta m\Delta t)}$$

Expected sensitivity with $5 \text{ ab}^{-1} \sim 3\text{-}4 \times 10^{-3}$

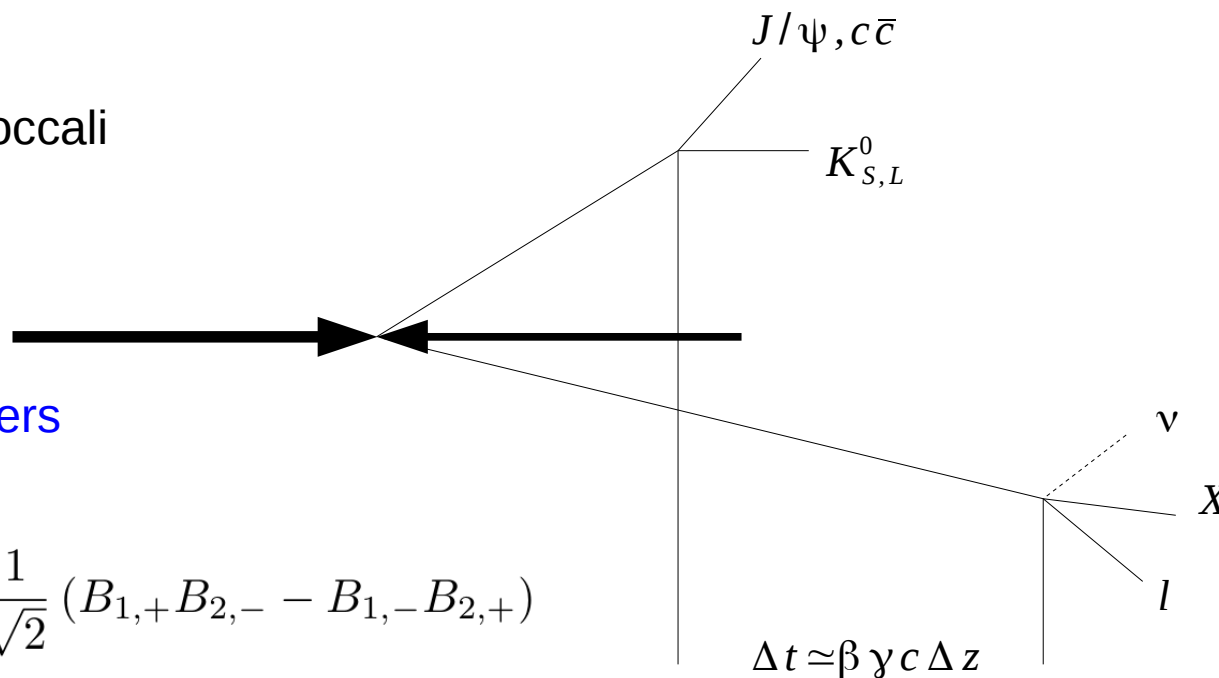
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Test of T invariance using CP filters

$$\Psi = \frac{1}{\sqrt{2}} \left(B_1^0 \bar{B}_2^0 - \bar{B}_1^0 B_2^0 \right) = \frac{1}{\sqrt{2}} (B_{1,+} B_{2,-} - B_{1,-} B_{2,+})$$

$$A_T = \frac{\Gamma(|1\rangle \rightarrow |2\rangle) - \Gamma(|2\rangle \rightarrow |1\rangle)}{\Gamma(|1\rangle \rightarrow |2\rangle) + \Gamma(|2\rangle \rightarrow |1\rangle)}$$

$$g_{\alpha,\beta}^{\pm} \propto \left[1 + C_{\alpha,\beta}^{\pm} \cos \Delta m \Delta t + S_{\alpha,\beta}^{\pm} \sin \Delta m \Delta t \right]$$

$$S_{\alpha,\beta}^{\pm} = \frac{2 \text{Im} \lambda}{1 + |\lambda|^2} \quad \text{and} \quad C_{\alpha,\beta}^{\pm} = \frac{1 - |\lambda|^2}{1 + |\lambda|^2}$$

$$A_T \simeq \frac{\Delta C_T^{\pm}}{2} \cos \Delta m \Delta t + \frac{\Delta S_T^{\pm}}{2} \sin \Delta m \Delta t$$

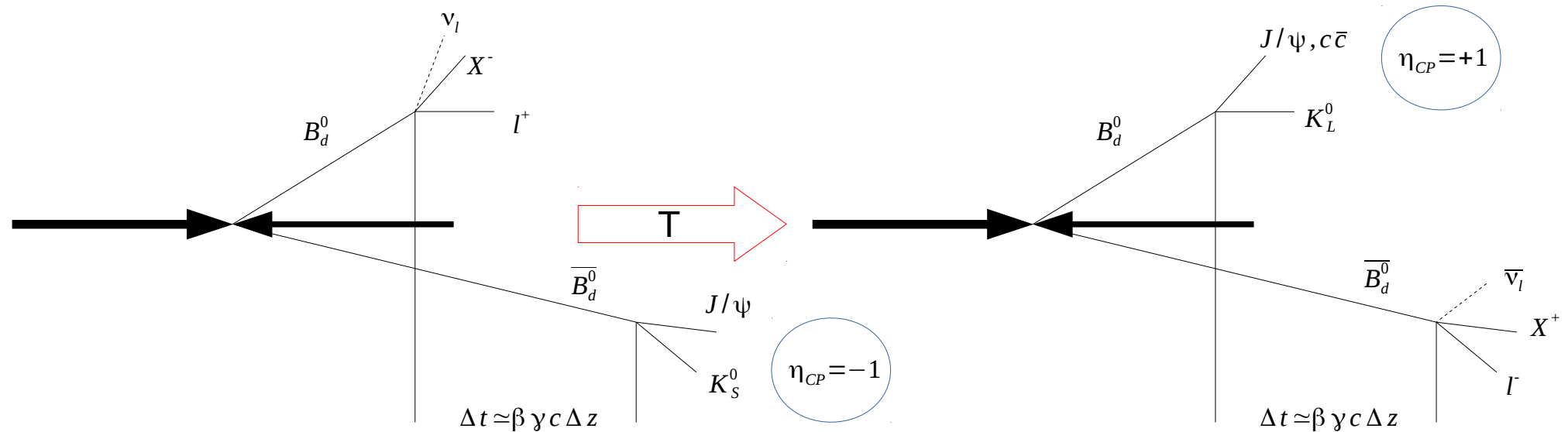
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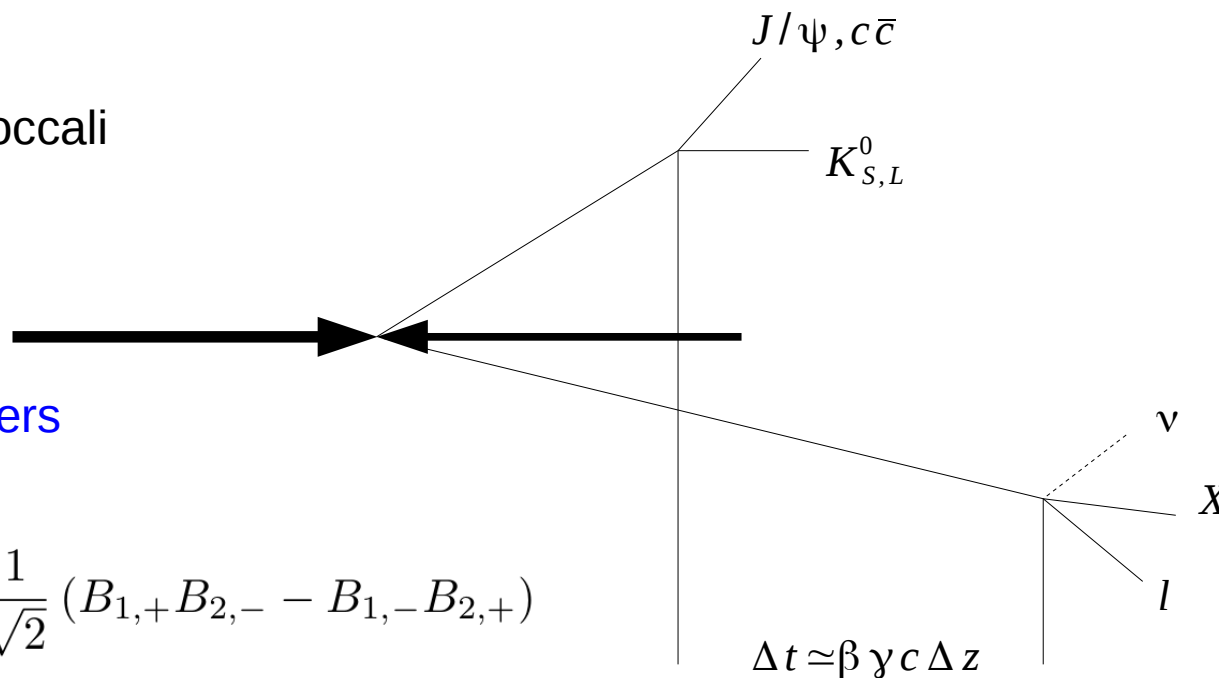
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$$A_T = \frac{\Gamma(|1\rangle \rightarrow |2\rangle) - \Gamma(|2\rangle \rightarrow |1\rangle)}{\Gamma(|1\rangle \rightarrow |2\rangle) + \Gamma(|2\rangle \rightarrow |1\rangle)}$$

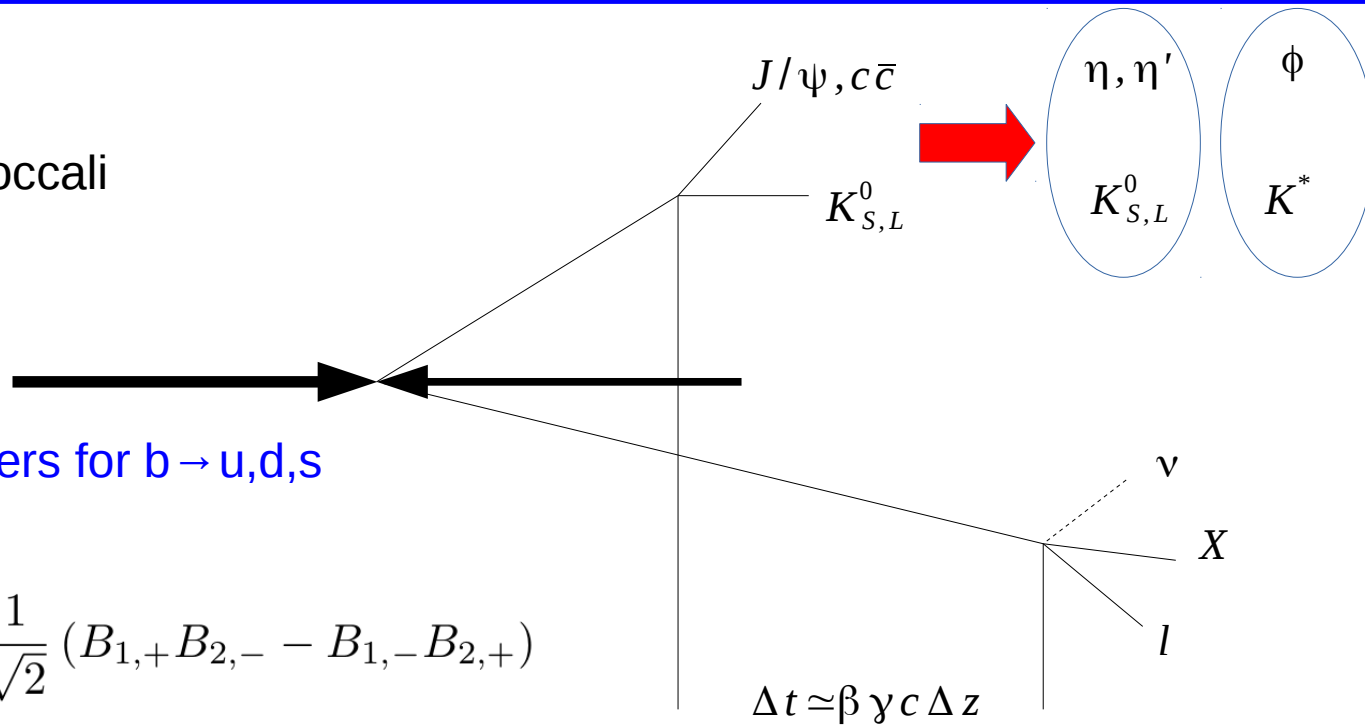
$$g_{\alpha,\beta}^{\pm} \propto \left[1 + C_{\alpha,\beta}^{\pm} \cos \Delta m \Delta t + S_{\alpha,\beta}^{\pm} \sin \Delta m \Delta t \right]$$

$$S_{\alpha,\beta}^{\pm} = \frac{2 \text{Im} \lambda}{1 + |\lambda|^2} \quad \text{and} \quad C_{\alpha,\beta}^{\pm} = \frac{1 - |\lambda|^2}{1 + |\lambda|^2}$$

$$A_T \simeq \frac{\Delta C_T^{\pm}}{2} \cos \Delta m \Delta t + \frac{\Delta S_T^{\pm}}{2} \sin \Delta m \Delta t$$

Beside CP, can Belle II look for T/CPT non invariance?

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 1207.5832 BaBar, T+CPT
 1302.4191 A. Bevan, G.I., M. Zoccali
 1605.04545 BaBar, CPT



Test of T invariance using CP filters for $b \rightarrow u, d, s$

$$\Psi = \frac{1}{\sqrt{2}} (B_1^0 \bar{B}_2^0 - \bar{B}_1^0 B_2^0) = \frac{1}{\sqrt{2}} (B_{1,+} B_{2,-} - B_{1,-} B_{2,+})$$

$$A_T = \frac{\Gamma(|1\rangle \rightarrow |2\rangle) - \Gamma(|2\rangle \rightarrow |1\rangle)}{\Gamma(|1\rangle \rightarrow |2\rangle) + \Gamma(|2\rangle \rightarrow |1\rangle)}$$

$$g_{\alpha,\beta}^{\pm} \propto \left[1 + C_{\alpha,\beta}^{\pm} \cos \Delta m \Delta t + S_{\alpha,\beta}^{\pm} \sin \Delta m \Delta t \right]$$

$$S_{\alpha,\beta}^{\pm} = \frac{2 \text{Im} \lambda}{1 + |\lambda|^2} \quad \text{and} \quad C_{\alpha,\beta}^{\pm} = \frac{1 - |\lambda|^2}{1 + |\lambda|^2}$$

$$A_T \simeq \frac{\Delta C_T^{\pm}}{2} \cos \Delta m \Delta t + \frac{\Delta S_T^{\pm}}{2} \sin \Delta m \Delta t$$

Channel suppressed by a factor 10,
 but considering $\sim 5 \text{ab}^{-1}$ of data one
 can expect similar precision on $b \rightarrow uds$
 as obtained by BaBar for $b \rightarrow c$
 transitions

$$\Delta S_T^+ = -1.37 \pm 0.14 \text{ (stat.)} \pm 0.06 \text{ (syst.)},$$

$$\Delta S_T^- = 1.17 \pm 0.18 \text{ (stat.)} \pm 0.11 \text{ (syst.)}$$

Conclusions

- Important opportunities for new physics searches ahead.
- For invisible $Y(1S)$ decays we need to have precise plan for data taking at the $Y(3S)$, and this is still missing. Triggers will depend also on when data are taken, ideally the sooner the better..In general tests of exotics triggers will be performed during data taking from phase 2.
- T/CPT tests could and should be performed during the operation of Belle II, and a few channels have been identified to complement and improve existing measurements.
- JENNIFER contribution will have a high impact on these projects:
 - 4-5 people (including myself, excluding undergraduates) are expected to work on these studies starting in mid 2018 to ~2022
 - For $Y(3S)$ data taking 1-2 people should be based at KEK.

Thank you for your attention!

