Global fit to v_{μ} disappearance data with sterile neutrinos

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Standard model of particle physics



Standard model of particle physics



3v in SM

Standard model of particle physics*



*with massive neutrinos



2

$$\nu_{\alpha} = \sum_{k} U_{\alpha k}^{*} \nu_{k}, P(\alpha \to \beta; E, L) = \sum_{k,j} U_{\alpha k}^{*} U_{\beta k} U_{\alpha j} U_{\beta j}^{*} e^{i \frac{\Delta m_{k j}^{2}}{2E}L}$$



Phys.Lett. B782 (2018) 633-640, P.F. de Salas, D.V. Forero, CAT, M. Tórtola, J.W.F. Valle https://globalfit.astroparticles.es/

See also: -Bari-group (Prog.Part.Nucl.Phys. 102 (2018) 48-72) -Nu-fit (JHEP 1901 (2019) 106)

Remaining unknowns are



Phys.Lett. B782 (2018) 633-640, P.F. de Salas, D.V. Forero, CAT, M. Tórtola, J.W.F. Valle



Neutrino mass ordering

Combine with data from decay experiments and cosmological observation using conservative priors to obtain 3.5σ preference for normal ordering

JCAP 1803 (2018) no.03, 011, S. Gariazzo, M. Archidiacono, P.F. de Salas, O. Mena, CAT, M. Tórtola

Front. Astron. Space Sci. 5 (2018) 36, P.F. de Salas, S. Gariazzo, O. Mena, CAT, M. Tórtola



Anomalies in oscillations

LSND

Appearance of electron antineutrinos from a pure source of muon antineutrinos $\overline{\nu}_{\mu} \rightarrow \overline{\nu}_{e}$ for energies 20 MeV < E < 52.8 MeV over 30m



Beam Excess

LSND

PRL 75 (1995) 2650 PRC 54 (1996) 2685 PRL 77 (1996) 3082 PRD 64 (2001) 112007

Karmen PRD 65 (2002) 112001

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LSND saw a 3.8σ excess

No signal seen by **Karmen** at 18m (same technique)

LSND

PRL 75 (1995) 2650 PRC 54 (1996) 2685 PRL 77 (1996) 3082 PRD 64 (2001) 112007

Karmen PRD 65 (2002) 112001



Gallium anomaly

Gallium Radioactive Source Experiments Gallex and Sage



PRC 73 (2006) 045805, **SAGE** PRC 80 (2009) 015807, **SAGE** Nucl.Phys.Proc.Suppl. 168 (2007) 344, Laveder et al PRD 78 (2008) 073009 and PRC 83 (2011) 065504, C. Giunti et al

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Measure electron neutrinos $\nu_e \rightarrow \nu_e$



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Measure electron neutrinos $\nu_e \rightarrow \nu_e$

 $\sim 3\sigma$ deficit of neutrinos



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Reactor Antineutrino Anomaly New reactor antineutrino fluxes, again ~3σ deficit



PRD 83 (2011) 073006, Mention et al PRC 83 (2011) 054615, Mueller et al PRC 84 (2011) 024617, Huber

Beyond three-neutrino oscillations



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We can add a forth neutrino

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This neutrino must be sterile, which means it is a singlet under all standard model gauge groups

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This neutrino must be sterile, which means it is a singlet under all standard model gauge groups

A forth active neutrino is excluded by observations of invisible Z-decays

$$e^+e^- \to Z \to \sum_{j=e,\mu,\tau} \nu_j$$

Phys. Rept. 427 (2006) 257, LEP



3+1 neutrino oscillations We extend the mixing matrix

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \Rightarrow \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$

3+1 neutrino oscillations
We extend the mixing matrix

$$\begin{pmatrix}
U_{e1} & U_{e2} & U_{e3} \\
U_{\mu1} & U_{\mu2} & U_{\mu3} \\
U_{\tau1} & U_{\tau2} & U_{\tau3}
\end{pmatrix} \Rightarrow \begin{pmatrix}
U_{e1} & U_{e2} & U_{e3} \\
U_{\mu1} & U_{\mu2} & U_{\mu3} \\
U_{\tau1} & U_{\tau2} & U_{\tau3} \\
U_{s1} & U_{s2} & U_{s3}
\end{pmatrix} \begin{pmatrix}
U_{e4} \\
U_{\mu4} \\
U_{\mu4} \\
U_{\tau4} \\
U_{s4} \\
U_{s4} \\
\end{pmatrix}$$
DISappearance

$$P_{\alpha\alpha}^{\text{SBL}} \approx 1 - \sin^2(2\theta_{\alpha\alpha}) \sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$$

$$\sin^2(2\theta_{\alpha\alpha}) = 4|U_{\alpha4}|^2(1 - |U_{\alpha4}|^2)$$

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DISappearance

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$$\nu_e \rightarrow \nu_e : |U_{e4}|^2 = \sin^2\theta_{14}$$
@Reactors and Gallium

3+1 neutrino oscillations We extend the mixing matrix $\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \Rightarrow \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$ DISappearance $P_{\alpha\alpha}^{\rm SBL} \approx 1 - \sin^2(2\theta_{\alpha\alpha})\sin^2\left(\frac{\Delta m_{41}^2 L}{4E}\right)$ $\sin^2(2\theta_{\alpha\alpha}) = 4|U_{\alpha4}|^2(1-|U_{\alpha4}|^2)$ $\nu_{e} \to \nu_{e} : |U_{e4}|^{2} = \sin^{2} \theta_{14}$ @Reactors and Gallium $\nu_{\mu} \rightarrow \nu_{\mu} : |U_{\mu4}|^2 = \sin^2 \theta_{24} \cos^2 \theta_{14}$ @atmospherics and accelerators

3+1 neutrino oscillations
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APPearance

$$P^{\text{SBL}}_{\alpha\beta} \approx \sin^{2}(2\theta_{\alpha\beta}) \sin^{2}\left(\frac{\Delta m_{41}^{2}L}{4E}\right) \qquad P^{\text{SBL}}_{\alpha\alpha} \approx 1 - \sin^{2}(2\theta_{\alpha\alpha}) \sin^{2}\left(\frac{\Delta m_{41}^{2}L}{4E}\right)$$

$$\sin^{2}(2\theta_{\alpha\beta}) = 4|U_{\alpha4}|^{2}|U_{\beta4}|^{2} \qquad \sin^{2}(2\theta_{\alpha\alpha}) = 4|U_{\alpha4}|^{2}(1 - |U_{\alpha4}|^{2})$$

$$\nu_{e} \rightarrow \nu_{e} : |U_{e4}|^{2} = \sin^{2}\theta_{14}$$
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$$P_{\alpha\beta}^{SBL} \approx \sin^{2}(2\theta_{\alpha\beta}) \sin^{2} \left(\frac{\Delta m_{41}^{2}L}{4E}\right)$$

$$\sin^{2}(2\theta_{\alpha\beta}) = 4|U_{\alpha4}|^{2}|U_{\beta4}|^{2}$$

$$\nu_{\mu} \rightarrow \nu_{e} : \sin^{2}(2\theta_{\mu e}) = 4|U_{e4}|^{2}|U_{\mu4}|^{2}$$

$$\psi_{e} \rightarrow \nu_{e} : |U_{e4}|^{2} = \sin^{2}\theta_{14}$$
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In LBL experiments matter effects are important!

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$$\begin{pmatrix} V_{\rm CC} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \Rightarrow \begin{pmatrix} V_{\rm CC} + V_{\rm NC} & 0 & 0 \\ 0 & V_{\rm NC} & 0 & 0 \\ 0 & 0 & V_{\rm NC} & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

In LBL experiments matter effects are important!



Neutral current potential becomes important

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Neutral current potential becomes important We have to take it into account when diagonalizing the Hamiltonian

v_{μ} disappearance experiments

IceCube and DeepCore



PRL 117 (2016) 071801

IceCube and DeepCore



High-energy regime 0.3 TeV – 20 TeV

PRL 117 (2016) 071801

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High-energy regime 0.3 TeV – 20 TeV Waiting for 7 yr update!

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PRD 95 (2017) 112002

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IceCube and DeepCore



High-energy regime 0.3 TeV – 20 TeV Waiting for 7 yr update!



Low-energy regime 6 GeV – 56 GeV

PRL 117 (2016) 071801

PRD 95 (2017) 112002

IceCube and DeepCore

 $2\Delta \ln \lambda$

0.3



High-energy regime 0.3 TeV – 20 TeV Waiting for 7 yr update!

SK, NO (2015), 90 % C.I SK, NO (2015), 99 % C.L IceCube, NO (2016), 90 % C.I 0.25ceCube, NO (2016), $|\mathrm{U}_{74}|^2 = \sin^2 \theta_{34} \cdot \cos^2 \theta_{24}$ 010 070 070 070 IceCube, IO (2016), 90 % C.L. IceCube, IO (2016), 99 % C.L 0.050.00 10^{-2} 0 2 4 6 8 10^{-3} 10^{-1} $-2\Delta \ln \mathcal{L}$ $|U_{\mu4}|^2 = \sin^2 \theta_{24}$

Low-energy regime 6 GeV – 56 GeV Also constraining θ_{34}

PRD 95 (2017) 112002

PRL 117 (2016) 071801



Two analyses: far-over-near ratio, and two-detector fit

PRL 117 (2016) 151803 PRL 122 (2019) 091803



Two analyses: far-over-near ratio, and two-detector fit For large mass splittings: systematic dominated

PRL 117 (2016) 151803 PRL 122 (2019) 091803



For mass splittings below 20 eV² the bound gets stronger after updating the analysis



For mass splittings below 20 eV² the bound gets stronger after updating the analysis

The effect of the other oscillation parameters is very small in this region

Fit of $\nu_{\mu}/\overline{\nu}_{\mu}$ disappearance data

All data:



MINOS/MINOS+ is the most dominating experiment in the fit

Fit of $\nu_{\mu}/\overline{\nu}_{\mu}$ disappearance data

All data:



DeepCore is the most dominating in constraining $|U_{\tau 4}|^2$



What is the impact on the 3+1 picture?

v_e disappearance experiments

NEOS

Single detector, taking ratio to Daya Bay





DANSS

Single movable detector, $\sim 3\sigma$ preference for 3+1





Fit of $\nu_e/\overline{\nu}_e$ **disappearance data DANSS / NEOS**



$$\Delta m_{41}^2 = 1.3 \text{ eV}^2$$
$$\sin^2(2\theta_{ee}) = 0.05$$
$$\sin^2\theta_{14} = 0.01$$

See also: JHEP 1711 (2017) 099, Dentler, M. et al

Fit of $\nu_e/\overline{\nu}_e$ **disappearance data** DANSS / NEOS + Gallium + RAA



Fit of $\nu_e/\overline{\nu}_e$ disappearance data All data:



Fit of $\nu_e/\overline{\nu}_e$ disappearance data New Stereo data:



v_e appearance experiments

MiniBooNE

MiniBooNE was built to check the **LSND** results with a different baseline, but similar L/E



PRL 121 (2018) 221801

MiniBooNE

MiniBooNE was built to check the **LSND** results with a different baseline, but similar L/E

MiniBooNE has no near detector



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MiniBooNE has no near detector



MiniBooNE sees an excess at $\sim 5\sigma$ at low energies



PRL 121 (2018) 221801

Fit of $\nu_e/\overline{\nu}_e$ appearance data All data:



The best fit value of **MiniBooNE** is excluded by **Icarus** and **Opera**

Fit of $\nu_e/\overline{\nu}_e$ appearance data All data:



The best fit value of MiniBooNE is excluded by Icarus and Opera LSND and MiniBooNE only partially agree

Tension in APP vs DIS data

• Data end of 2017



JHEP 1706 (2017) 135, S. Gariazzo, C. Giunti, M. Laveder, Y.F. Li

Tension in APP vs DIS data Data June 2019











Tension in APP vs DIS data Data June 2019



Tension in APP vs DIS data Data June 2019



Tension in APP vs DIS data Only excluding LSND and MB solves the problem



No surprise, because now there is no lower bound

Conclusions



If one insists that all of the LSND and MB excesses comes from sterile mixing, then:

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APP-DIS tension makes the 3+1 global fit unacceptable





If one assumes that (at least part of) the LSND and MB excesses could come from different physics:
Conclusions

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Model-independent indication of light sterile neutrino oscillations in very short-baseline reactor experiments

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Model-independent indication of light sterile neutrino oscillations in very short-baseline reactor experiments

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This raises the question: What are **LSND** and **MiniBoone** observing?







	All	MВ	LSND	MB&LSND
$\chi^2_{\rm min}$	827.4	765.8	802.2	726.9
NDF	760	726	756	722
GoF	4.5%	15%	12%	44%
Δm^2_{41}	1.3	1.3	1.3	1.3
$ U_{e4} ^2$	0.017	0.016	0.012	0.012
$ U_{\mu 4} ^2$	0.019	0.013	0.0037	0
$\sin^2 2\vartheta_{e\mu}$	0.0013	0.00084	0.00017	0
$\sin^2 2\vartheta_{ee}$	0.068	0.064	0.046	0.047
$\sin^2 2\vartheta_{\mu\mu}$	0.073	0.052	0.015	0
$\Delta \chi^2_{\rm NO}$	43.5	31.3	16.2	17.7
NDF_{NO}	4	4	4	4
$n\sigma_{ m NO}$	5.8	4.7	3.0	3.2
$(\chi^2_{\rm min})_{\rm App}$	98.4	49.8	90.3	37.4
NDF_{App}	80	46	76	42
GoF_{App}	7.9%	33%	13%	67%
Δm^2_{41}	0.58	0.93	0.52	0
$\sin^2 2\vartheta_{e\mu}$	0.0068	0.0033	0.0058	0
$(\chi^2_{\rm min})_{\rm Dis}$	690.0	\leftarrow	\leftarrow	\leftarrow
$\mathrm{NDF}_{\mathrm{Dis}}$	678	\leftarrow	\leftarrow	\leftarrow
${\rm GoF}_{\rm Dis}$	37%	\leftarrow	\leftarrow	\leftarrow
Δm_{41}^2	1.3	\leftarrow	\leftarrow	\leftarrow
$ U_{e4} ^2$	0.012	\leftarrow	\leftarrow	\leftarrow
$ U_{\mu 4} ^2$	0	\leftarrow	\leftarrow	\leftarrow
$\sin^2 2\vartheta_{e\mu}$	0	\leftarrow	\leftarrow	\leftarrow
$\sin^2 2\vartheta_{ee}$	0.047	\leftarrow	\leftarrow	\leftarrow
$\sin^2 2\vartheta_{\mu\mu}$	0	\leftarrow	\leftarrow	\leftarrow
$\Delta \chi^2_{\rm PG}$	39.0	26.0	22.4	-0.008
NDF_{PG}	2	2	2	2
GoF_{PG}	3×10^{-9}	2×10^{-6}	1×10^{-5}	100%



1905.11290, S. Gariazzo, P.F. de Salas. S. Pastor