

New physics implication from Kaon physics



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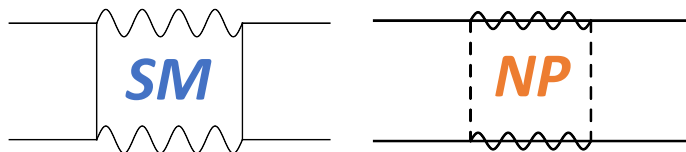


Why Kaon? & What's new?

- Kaon observables are sensitive to NP at a very high scale, which is not accessible at the LHC

- FCNC and CP violation in Kaon system are suppressed in the SM

c.f. meson mixing



$$\mathcal{L}_{eff} = \mathcal{L}^{SM} + \frac{1}{\Lambda_{NP}^2} \sum_i C_i \mathcal{O}_i^{\text{dim6}}$$

If $|C_{NP}| \sim 1$

$$\Lambda_{NP} \sim \begin{cases} \mathcal{O}(10^5 \text{ TeV}) & : K^0 \\ \mathcal{O}(10^4 \text{ TeV}) & : D^0 \\ \mathcal{O}(10^3 \text{ TeV}) & : B_{d,s} \end{cases}$$

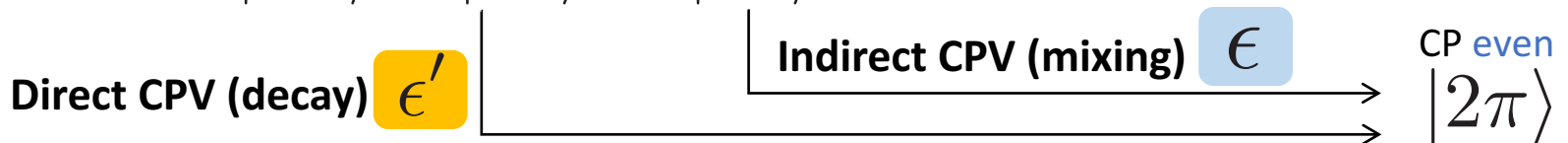
- Several on-going experiments for Kaon observables (KOTO/NA62...)

★ Using recent result of lattice calculation, there is discrepancy in ϵ'/ϵ between SM value and data

ϵ and ϵ'

1964 $K_L \rightarrow 2\pi$ was observed *Discovery of CP violation*

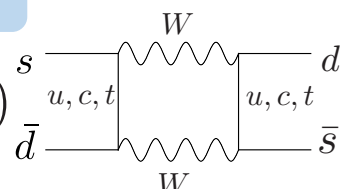
$$|K_L\rangle = \overset{\text{CP odd}}{|K_2\rangle} + \overset{\text{CP even}}{\epsilon}|K_1\rangle$$



$$\eta_{00} \equiv \frac{A(K_L \rightarrow \pi^0 \pi^0)}{A(K_S \rightarrow \pi^0 \pi^0)} \equiv \epsilon - 2\epsilon'$$

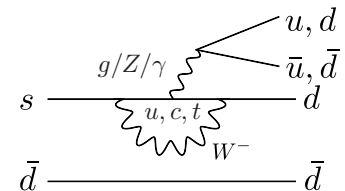
$$\eta_{+-} \equiv \frac{A(K_L \rightarrow \pi^+ \pi^-)}{A(K_S \rightarrow \pi^+ \pi^-)} \equiv \epsilon + \epsilon'$$

Indirect CPV (mixing) ϵ

$$|\epsilon| \simeq \frac{1}{3} (|\eta_{00}| + 2|\eta_{+-}|)$$


Direct CPV (decay) ϵ'

$$\left| \frac{\eta_{00}}{\eta_{+-}} \right|^2 \simeq 1 - 6\text{Re} \left(\frac{\epsilon'}{\epsilon} \right)$$



$$\epsilon = \mathcal{O}(10^{-3}) \quad \text{Re} \left(\frac{\epsilon'}{\epsilon} \right) = \mathcal{O}(10^{-3}) \quad \epsilon' = \mathcal{O}(10^{-6})$$

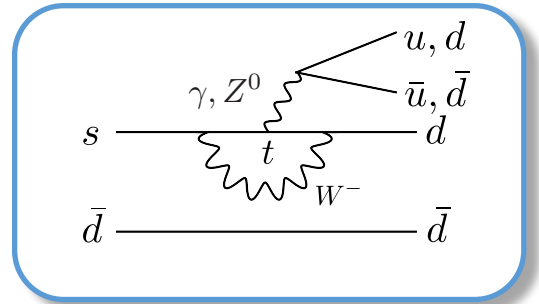
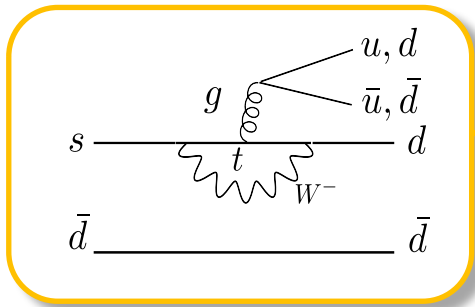
Highly suppressed and sensitive to NP

ϵ'/ϵ

$$A(K^0 \rightarrow (\pi\pi)_{I=0,2}) = A_{0,2} e^{i\delta_{0,2}}$$

$$\frac{\epsilon'}{\epsilon} = - \frac{\omega}{\sqrt{2} |\epsilon|_{\text{exp}} \text{Re}A_0} \left(\text{Im}A_0 - \frac{1}{\omega} \text{Im}A_2 \right)$$

QCD penguin operator **EW penguin operator**



$\Delta I=1/2$ rule $\frac{\text{Re}A_0}{\text{Re}A_2} \equiv \frac{1}{\omega} = 22.46 \quad (\text{exp.})$

In the SM, there is accidental cancellation between $\text{Im}A_0$ and $\text{Im}A_2$ due to the enhancement factor $1/\omega$

EW penguin is comparable to QCD penguin due to the enhancement factor

ϵ'/ϵ discrepancy

$$\langle (\pi\pi)_I | \mathcal{H} | K^0 \rangle = \sum_n \overset{\text{Short distance}}{C_n} \overset{\text{Matrix element}}{\langle (\pi\pi)_I | \mathcal{O}_n | K^0 \rangle}$$

Short distance

- NLO result has been available since early 90's
- NNLO QCD calculation is in progress *Cerda-Sevilla, Gorbahn, Jager, Kokulu 1611.08276*

Long distance (Matrix elements)

- **First lattice result by RBC-UKQCD in 2015** *1502.00263 1505.07863*

From the lattice result, ϵ'/ϵ has been calculated in SM using data for $\text{Re}A_{0,2}$

SM with Lattice $\left(\frac{\epsilon'}{\epsilon}\right)_{\text{SM}} = (1.06 \pm 5.07) \times 10^{-4}$ *Kitahara, Nierste and Tremper, 1607.06727*
c.f. RBC-UKQCD / Buras, Gorbahn, Jager and Jamin 1507.06345

Exp $\left(\frac{\epsilon'}{\epsilon}\right)_{\text{exp}} = (16.6 \pm 2.3) \times 10^{-4}$ *average of NA48 and KTeV*

2.8 σ difference

NP in ϵ'/ϵ ?

ϵ'/ϵ discrepancy

SM with Lattice

$$\left(\frac{\epsilon'}{\epsilon}\right)_{\text{SM}} = (1.06 \pm 5.07) \times 10^{-4}$$

2.8 σ

difference

Exp

$$\left(\frac{\epsilon'}{\epsilon}\right)_{\text{exp}} = (16.6 \pm 2.3) \times 10^{-4}$$

- O_6 & O_8 have dominant effects on ϵ'/ϵ due to chiral enhancement

$$\begin{aligned} \langle (\pi\pi)_0 | \mathcal{O}_6 | K \rangle &\propto B_6^{(1/2)} \\ \langle (\pi\pi)_2 | \mathcal{O}_8 | K \rangle &\propto B_8^{(3/2)} \end{aligned}$$

Non-perturbative parameter

$$\text{QCD penguin } O_6 = (\bar{s}_\alpha d_\beta)_{V-A} \sum (\bar{q}_\beta q_\alpha)_{V+A}$$

$$\text{EW penguin } O_8 = \frac{3}{2} (\bar{s}_\alpha d_\beta)_{V-A} \sum_q e_q (\bar{q}_\beta q_\alpha)_{V+A}$$

- Values extracted from the lattice result

Buras, Buttazzo, Girschbach-Noe and Knegjens 1503.02693

$$B_6^{(1/2)} = 0.57 \pm 0.19 \quad B_8^{(3/2)} = 0.76 \pm 0.05$$

- Error for ϵ'/ϵ is dominated by $B_6^{(1/2)}$

- Two ways of analytic approaches

Large N_c Dual QCD approach

*Buras and Gérard
1507.06326*

$$B_6^{(1/2)} \leq B_8^{(3/2)} < 1 \quad \left(\frac{\epsilon'}{\epsilon}\right)_{\text{SM}} < (6.0 \pm 2.4) \times 10^{-4}$$

1805.11096

ChPT (FSI)

*Gisbert and Pich 1712.06147
hep-ph/0007208*

$$\begin{aligned} B_6^{(1/2)} &\sim 1.5 \\ B_8^{(3/2)} &\sim 0.9 \end{aligned} \quad \left(\frac{\epsilon'}{\epsilon}\right)_{\text{SM}} = (15 \pm 7) \times 10^{-4}$$

Result in DQCD approach gives support to lattice result. On the other hand, result in ChPT is consistent with data

- Wait for improved lattice results

ϵ'/ϵ at Lattice study

Amplitude	Exp. data	Lattice QCD
$\text{Re}A_0$ [10^{-7} GeV]	3.322 ± 0.001 [1]	$4.66 \pm 1.00 \pm 1.26$ [2]
$\text{Im}A_0$ [10^{-11} GeV]	—	$-1.90 \pm 1.23 \pm 1.08$ [2]
$\text{Re}A_2$ [10^{-8} GeV]	1.479 ± 0.003 [1]	$1.50 \pm 0.04 \pm 0.14$ [3]
$\text{Im}A_2$ [10^{-13} GeV]	—	$-6.99 \pm 0.20 \pm 0.84$ [3]

[1] Buras et al., 1507.06345
[2] RBC-UKQCD, 1505.07863
[3] RBC-UKQCD, 1502.00263

- $\text{Re}A_0$, $\text{Re}A_2$ are consistent with exp. Data $\rightarrow \Delta I=1/2$ rule is confirmed
- Calculated $I=0$ $\pi\pi$ scattering phase shift of was smaller than measured value

$$\delta_0 = 23.8(4.9)(1.2)^\circ \text{ 2015} \quad (\delta_0)_{\text{exp}} = 38.3(1.3)^\circ$$

\rightarrow New preliminary result *RBC-UKQCD preliminary, 2018*

$$\delta_0 = 30.9(1.5)(3.0)^\circ \quad \text{“Puzzle is resolved”}$$

- Lattice update with higher statistics will appear soon

ϵ'/ϵ beyond the SM

- Motivated by ϵ'/ϵ discrepancy, several new physics models have been studied

<i>Little Higgs Model with T-parity</i>	<i>Blanke, Buras and Recksiegel 1507.06316</i>
<i>Modified Z scenario</i>	<i>Buras, Buttazzo and Kneijens 1507.08672/Buras, 1601.00005</i> <i>Endo, Kitahara, Mishima and KY 1612.08839/Bobeth, Buras, Celis and Jung 1703.04753</i>
<i>Z' models</i>	<i>Buras, Buttazzo, Kneijens 1507.08672 /Buras 1601.00005</i>
<i>331 model</i>	<i>Buras and De Fazio 1512.02869/1604.02344</i>
<i>MSSM Chargino Z penguin</i>	<i>Endo, Mishima, Ueda and KY 1608.01444</i>
<i>Gluino Z penguin</i>	<i>Tanimoto and KY 1603.07960</i> <i>Endo, Goto, Kitahara, Mishima, Ueda and KY 1712.04959</i>
<i>Gluino Box</i>	<i>Kitahara, Nierste and Tremper 1604.07400,1703.05786</i> <i>Crivellin, D'Ambrosio, Kitahara, Nierste 1712.04959</i> <i>Chobanova, D'Ambrosio, Kitahara, Martínez, Santos, Fernández and KY 1711.11030</i>
<i>Vector-like quarks</i>	<i>Bobeth, Buras, Celis and Jung 1609.04783</i>
<i>Right handed current</i>	<i>Cirigliano, Dekens, Vries and Mereghetti 1612.03914</i> <i>Alioli, Cirigliano, Dekens, de Vries and Mereghetti 1703.04751</i>
<i>Leptoquark</i>	<i>Bobeth and Buras 1712.01295</i>
<i>LR symmetric model</i>	<i>Haba, Umeeda and Yamada 1802.09903/1806.0342</i>
<i>Type-III 2HDM</i>	<i>Chen and Nomura 1804.06017/ 1805.07522</i>
<i>Flavorful composite vectors</i>	<i>Matsuzaki, Nishiwaki and KY 1806.02312</i>
<i>Diquark model</i>	<i>Chen and Nomura 1808.04097</i>
<i>3HDM</i>	<i>Marzola and Raidal 1901.08290</i>
<i>General 2HDM</i>	<i>Iguro and Omura, 1905.11778</i>

ϵ'/ϵ beyond the SM

$$\frac{\epsilon'_K}{\epsilon_K} = -\frac{\omega}{\sqrt{2} |\epsilon_K|_{\text{exp}} \text{Re}A_0} \left(\text{Im}A_0 - \frac{1}{\omega} \text{Im}A_2 \right)$$

22.4

- CPV effect
- $\text{Im}A_2$ is enhanced by enhancement factor $1/\omega$
- SM effect is small due to this accidental cancellation

■ NP in $\text{Im}A_0$ or (and) $\text{Im}A_2$

$\text{Im}A_2$... have enhancement factor $1/\omega$

$\text{Im}A_0$... likely to result in huge contribution to ϵ_K

→ NP in $\text{Im}A_2$ is likely

ϵ'/ϵ beyond the SM

■ Master formulae for ϵ'/ϵ *Aebischer, Bobeth, Buras, Gérard and Straub 1807.02520*

Master formula including BSM operators is derived with DQCD

$$\left(\frac{\epsilon'}{\epsilon}\right)_{\text{BSM}} = \sum_i P_i(\mu_W) \text{Im} [C_i(\mu_W) - C'_i(\mu_W)]$$

P_i : Hadronic matrix elements + RG effects

Most efficient operators explaining ϵ'/ϵ anomaly

$O_{VLR}^u = (\bar{s}^\alpha \gamma_\mu P_L d^\alpha)(\bar{u}^\beta \gamma^\mu P_R d^\beta)$	$\left. \begin{array}{l} \text{SM type operators} \\ \text{HME calculated by Lattice \& DQCD} \\ \text{Generate O6(ImA0) \& O8(ImA2)} \end{array} \right\}$	<p>NP scenario</p> <p>New heavy vectors Modified Z penguin (MSSM, VLQ, LHT, ...)</p>
$\tilde{O}_{VLR}^u = (\bar{s}^\alpha \gamma_\mu P_L d^\beta)(\bar{u}^\beta \gamma^\mu P_R d^\alpha)$		
$O_{VLR}^d = (\bar{s}^\alpha \gamma_\mu P_L d^\alpha)(\bar{d}^\beta \gamma^\mu P_R d^\beta)$		
$\tilde{O}_{VLR}^d = (\bar{s}^\alpha \gamma_\mu P_L d^\beta)(\bar{d}^\beta \gamma^\mu P_R d^\alpha)$		
$O_{TLL}^u = (\bar{s}^\alpha \sigma_{\mu\nu} P_L d^\alpha)(\bar{u}^\beta \sigma^{\mu\nu} P_L d^\beta)$	$\left. \begin{array}{l} \text{New scalar \& tensor Operators} \\ \text{HME calculated by only DQCD} \end{array} \right\}$	<p>Heavy scalars</p>
$\tilde{O}_{TLL}^u = (\bar{s}^\alpha \sigma_{\mu\nu} P_L d^\beta)(\bar{u}^\beta \sigma^{\mu\nu} P_L d^\beta)$		
$O_{TLL}^d = (\bar{s}^\alpha \sigma_{\mu\nu} P_L d^\alpha)(\bar{d}^\beta \sigma^{\mu\nu} P_L d^\beta)$		
$O_{SLR}^u = (\bar{s}^\alpha P_L d^\alpha)(\bar{u}^\beta P_R u^\beta)$		

✂️ Chrome magnetic operator $\langle O_{8g} \rangle$ (calculated by Lattice & DQCD) would be suppressed

*ETM collaboration, 1712.09824
Buras and Gérard 1803.08052*

ε'/ε beyond the SM

■ SMEFT study : (*SM effective field theory*) [$SU(2) \times U(1)$ inv.] ($\mu_{EW} < \mu < \mu_{NP}$)

-Model independent approach *Aebischer, Bobeth, Buras and Straub 1808.00466*

The constraints from K^0 and D^0 mixing as well as EDM are potentially important

- Z penguin scenario *Bobeth, Buras, Celis and Jung 1703.04753*

Endo, Kitahara, Mishima and KY 1612.08839/ Endo, Goto, Kitahara, Mishima, Ueda and KY 1712.04959

$\Delta S=1$ operators generate $\Delta S=2$ contributions, through top-Yukawa enhanced RG evolution

$$\left(H^\dagger i \overleftrightarrow{D}_\mu H \right) (\bar{s}_R \gamma^\mu d_R) \xrightarrow{\text{RG evolution}} \begin{aligned} & (\bar{s}_R \gamma^\mu d_R) Z_\mu \\ & \Delta F=2 \text{ operator} \\ & (\bar{s}_L \gamma_\mu d_L) (\bar{s}_R \gamma^\mu d_R) \end{aligned}$$

ε'/ε beyond the SM

■ NP in ε'/ε also affect other observables

- $\Delta S=2$ process ϵ_K and ΔM_K give severe constraint

Some model need tuning to avoid this constraint

- Kaon rare decay $K_L \rightarrow \pi^0 \nu \bar{\nu}$ and $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ could be good probe

← Pure imaginary. Strong correlation with ε'/ε

- B observables have correlation (and become constraint) in some models

- Other observables (EDM)

■ Different implications (correlations & predictions) for other observables appear depending on models \Rightarrow Possibility of model discriminations



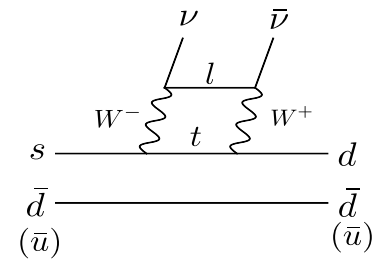
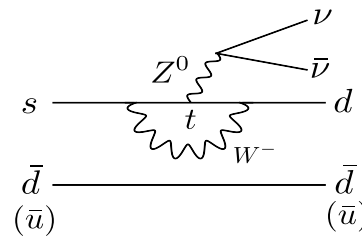
$K \rightarrow \pi \nu \bar{\nu}$



Correlation with B anomalies

$K_L \rightarrow \pi^0 \nu \bar{\nu}$ and $K^+ \rightarrow \pi^+ \nu \bar{\nu}$

- Highly suppressed in SM : $BR_{SM} \sim 10^{-11}$
- Theoretically clean (Hadronic matrix element can be estimated using isospin sym.)
- $K_L \rightarrow \pi^0 \nu \bar{\nu}$ is purely CP violating mode



$K^+ \rightarrow \pi^+ \nu \bar{\nu}$



NA62@CERN

- NA62 at CERN observed one event in 2016 data

$$BR(K^+ \rightarrow \pi^+ \nu \bar{\nu})_{SM} = (9.11 \pm 0.72) \times 10^{-11}$$

$$BR(K^+ \rightarrow \pi^+ \nu \bar{\nu})_{exp} = (1.73_{-1.05}^{+1.15}) \times 10^{-10}$$

BNL 949/E787

$$< 14 \times 10^{-10} (95\% C.L.) \quad \text{New 2018}$$

- Expected about 20 SM events from the 2017-2018 data sample

$K_L \rightarrow \pi^0 \nu \bar{\nu}$



KOTO@J-PARC

- KOTO at J-PARC reported result from the 2015 data last summer

$$BR(K_L \rightarrow \pi^0 \nu \bar{\nu})_{SM} = (3.00 \pm 0.30) \times 10^{-11}$$

$$BR(K_L \rightarrow \pi^0 \nu \bar{\nu})_{exp} < 2.6 \times 10^{-8} (90\% C.L.)$$

E391a

$$< 3.0 \times 10^{-9} (90\% C.L.) \quad \text{New 2018}$$

- KOTO-phase2 aims to measure at 10% accuracy

$\epsilon'/\epsilon \Leftrightarrow K \rightarrow \pi \nu \nu$ - Examples -

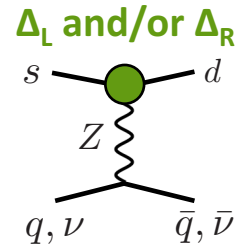
Z scenario

Buras, Buttazzo and Kneijens 1507.08672 / Buras 1601.00005 / Bobeth, Buras, Celis and Jung 1703.04753

There are interesting correlations between Kaon observables depending on the chiral structure of coupling (LH and/or RH)

$$\text{Re}\left(\frac{\epsilon'}{\epsilon}\right) \propto -\text{Im} \Delta_L^{sd} - 3 \text{Im} \Delta_R^{sd} + \dots$$

CPV \Rightarrow Strong correlation



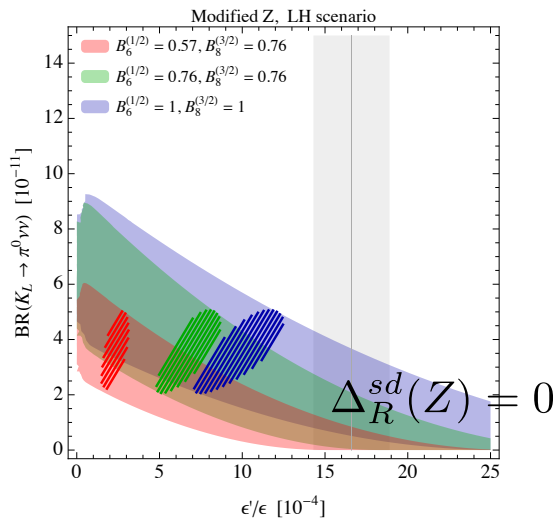
$$\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \propto |X + \dots|^2$$

$$\text{BR}(K_L \rightarrow \pi^0 \nu \bar{\nu}) \propto (\text{Im} X)^2$$

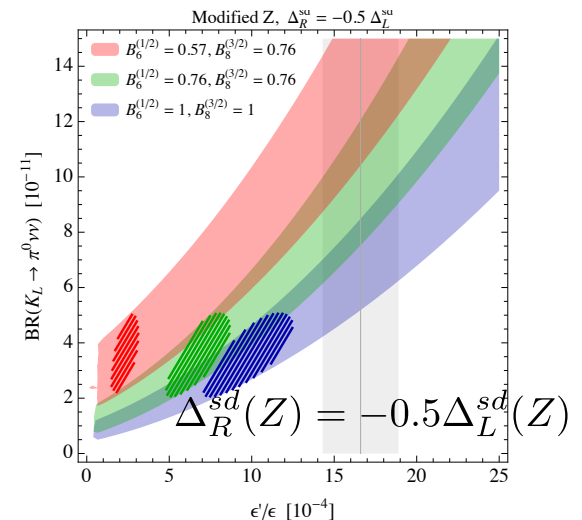
$$X = X(x_t)_{\text{SM}} + \frac{\pi^2}{2M_W^2 G_F^2} \frac{\Delta_L^{\nu\nu}}{V_{ts}^* V_{td} M_Z^2} (\Delta_L^{sd} + \Delta_R^{sd})$$

$$|\epsilon_K| \propto \text{Im} \left[(\Delta_L^{sd})^2 + (\Delta_R^{sd})^2 - 240 \Delta_L^{sd} \Delta_R^{sd} \right]$$

LH Scenario
 \Rightarrow negative correlation



LH+RH Scenario
 \Rightarrow positive correlation



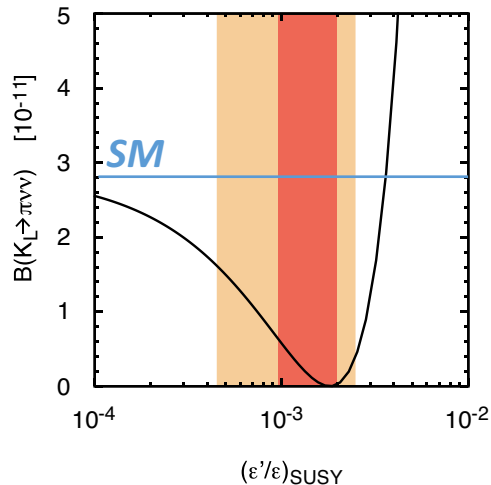
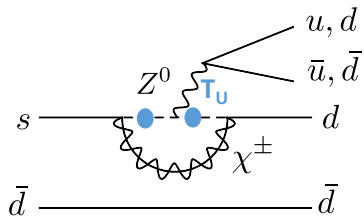
$\epsilon'/\epsilon \Leftrightarrow K \rightarrow \pi \nu \bar{\nu}$ - Examples -

Chargino Z penguin

Endo, Mishima, Ueda and KY
1608.01444

Large trilinear couplings bring enhancement of ϵ'/ϵ

LH Z scenario \Rightarrow negative correlation btwn ϵ'/ϵ and $K_L \rightarrow \pi^0 \nu \bar{\nu}$



$\epsilon'/\epsilon \Leftrightarrow$ SUSY scale $<$ 4-6 TeV

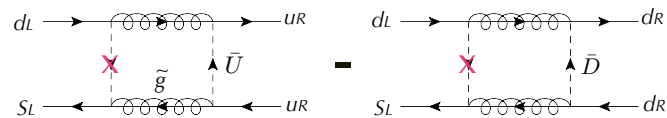
$$BR(K_L \rightarrow \pi^0 \nu \bar{\nu}) < 0.6 \text{ SM}$$

$$BR(K^+ \rightarrow \pi^+ \nu \bar{\nu}) \text{ O}(10 \sim 100\%) \text{ effect}$$

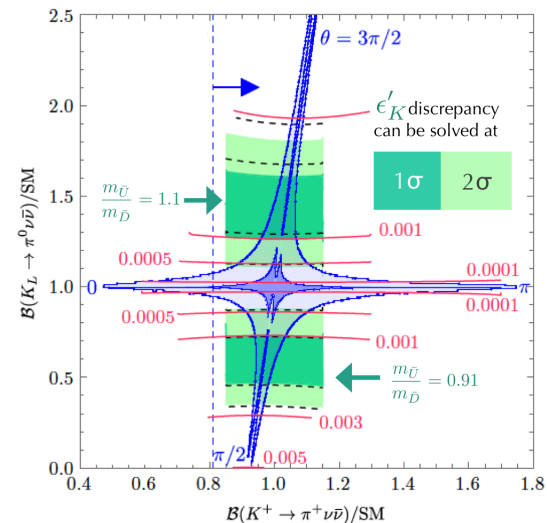
Glauino box

Crivellin, D'Ambrosio, Kitahara and Nierste
1703.05786

Large isospin breaking ($m_{\tilde{U}} \neq m_{\tilde{D}}$) gives effect on $\text{Im}A_2$



$m_{\tilde{q}_1} = 1.5 \text{ TeV}, m_L = 300 \text{ GeV}$



\Rightarrow **KOTO**

\Rightarrow **NA62**

Different correlations between ϵ'/ϵ and $K \rightarrow \pi \nu \bar{\nu}$ may allow to distinguish among models

B anomalies

Lepton flavor universality Violation (LFUV) in semi-leptonic **B** decays

$$b \rightarrow c\tau\nu$$

$$R_{D^{(*)}} = \frac{\mathcal{B}(B \rightarrow D^{(*)}\tau\nu)}{\mathcal{B}(B \rightarrow D^{(*)}\ell\nu)}$$

~ 3 σ excess over the SM

$$b \rightarrow s\ell\ell$$

$$R_{K^{(*)}} = \frac{\mathcal{B}(B \rightarrow K^{(*)}\mu^+\mu^-)}{\mathcal{B}(B \rightarrow K^{(*)}e^+e^-)}$$

~ 2.5 σ less over the SM

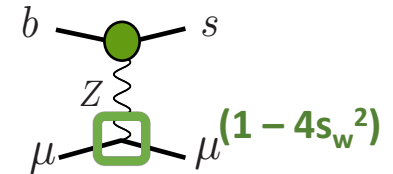
Correlation with ε'/ε ?

$\varepsilon'/\varepsilon \Leftrightarrow B$ anomalies

Possibility of simultaneous explanation of them are discussed in several models

- Z model is not favored by anomalies in $b \rightarrow s$ transitions, which suggest large C_9^{NP}

$$O_9 = (\bar{s}_L \gamma_\mu b_L)(\bar{\mu} \gamma^\mu \mu)$$



In Z model, it is hard to produce large C_9^{NP} due to smallness of the vector coupling to charged lepton

- In Leptoquark model, which is one of strong candidate of NP model realizing B anomalies, it is difficult to explain ε'/ε because of bounds from rare Kaon decays *Christoph and Buras 1712.01295*
- 2HDM + νR can address $R_{K^{(*)}}$ and ε'/ε *Iguro and Omura, 1905.11778*
- 3HDM + νR can access $RD^{(*)}, RK^{(*)}$ and ε'/ε *Marzo, Marzola and Raidal 1901.08290*
- Composite model can access $R_{K^{(*)}}$ and ε'/ε *Matsuzaki, Nishiwaki and KY 1806.02312* ★



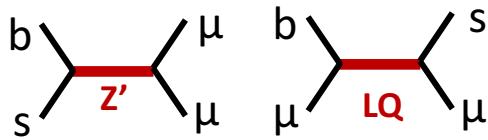
$\epsilon'/\epsilon \Leftrightarrow B$ anomalies - Example -

Flavorful composite vectors

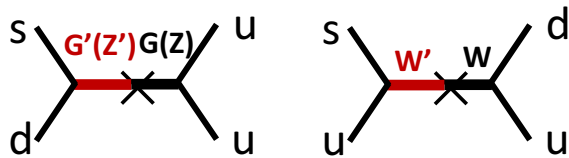
Matsuzaki, Nishiwaki and KY 1806.02312

New vector particles : $G', Z', W',$ Leptoquark(LQ) are included as **composite vectors**

$b \rightarrow s \mu \mu$



$\epsilon'/\epsilon (K \rightarrow \pi \pi)$

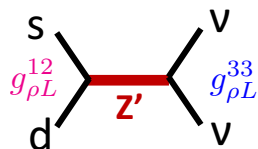


Flavor texture: **Assume pure imaginary** (to avoid ϵ_K constraint)

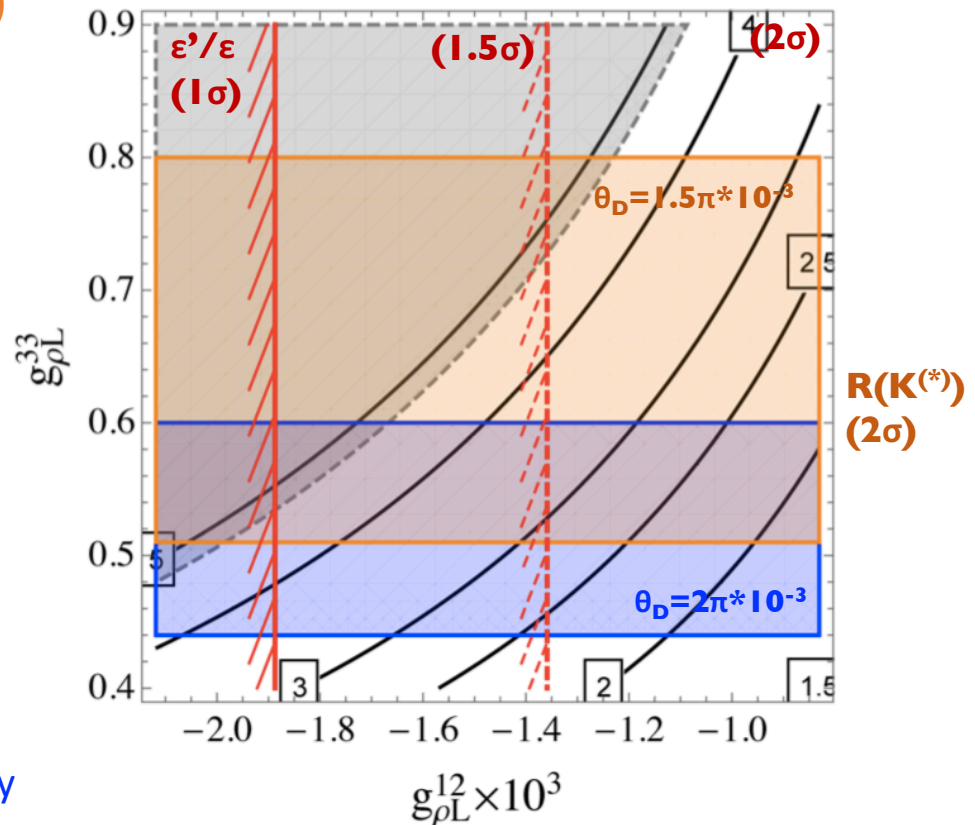
$$g_{\rho L}^{ij} = \begin{pmatrix} 0 & ig_{\rho L}^{12} & 0 \\ i(g_{\rho L}^{12})^* & 0 & 0 \\ 0 & 0 & g_{\rho L}^{33} \end{pmatrix}^{ij} \quad \epsilon'/\epsilon \quad B \text{ anomaly}$$

The correlation b/w ϵ'/ϵ and B obs. appear in $K \rightarrow \pi \nu \nu$

$K \rightarrow \pi \nu \nu$



$Br[K^+ \rightarrow \pi^+ \nu \nu]/SM$



$R(K^*)$ & ϵ'/ϵ (2σ)

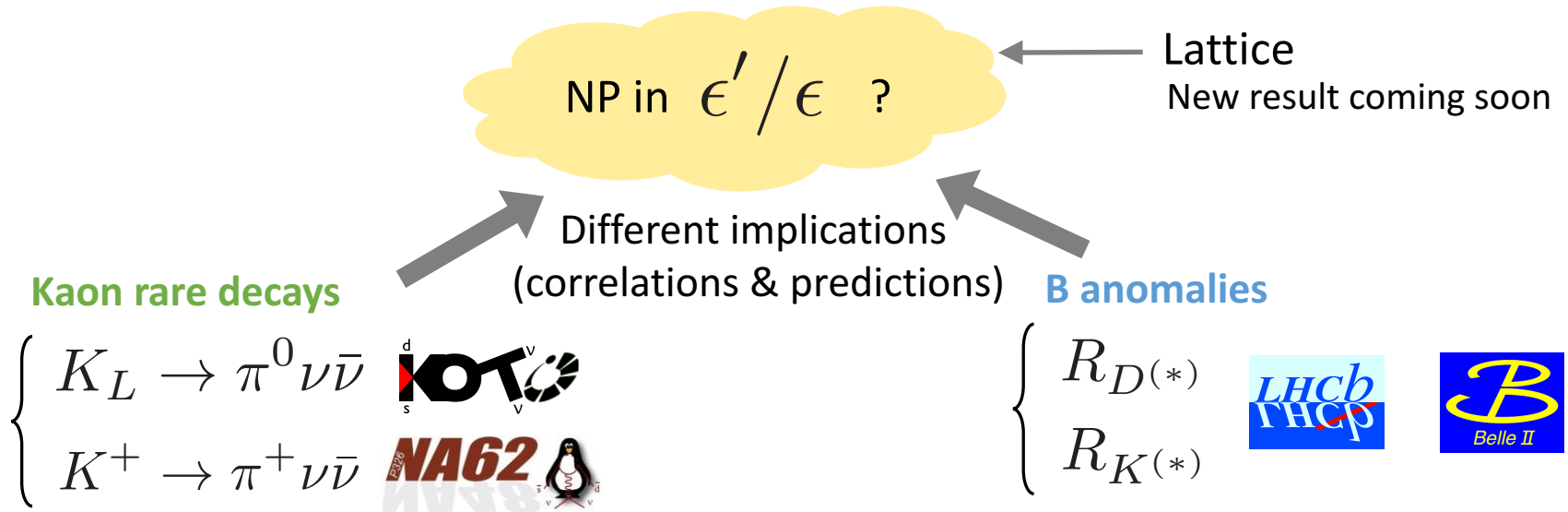
$\Leftrightarrow 1.5 < BR(K^+ \rightarrow \pi^+ \nu \nu)/SM$

$\Leftrightarrow 10 < BR(K_L \rightarrow \pi^0 \nu \nu)/SM$

Summary

Kaon physics is highly suppressed and sensitive to NP

2.8 σ discrepancy in direct CPV of Kaon ϵ'/ϵ



Many experiments are on-going, and could allow us to discriminate NP models

Kaon physics will continue to offer a powerful probe for NP!