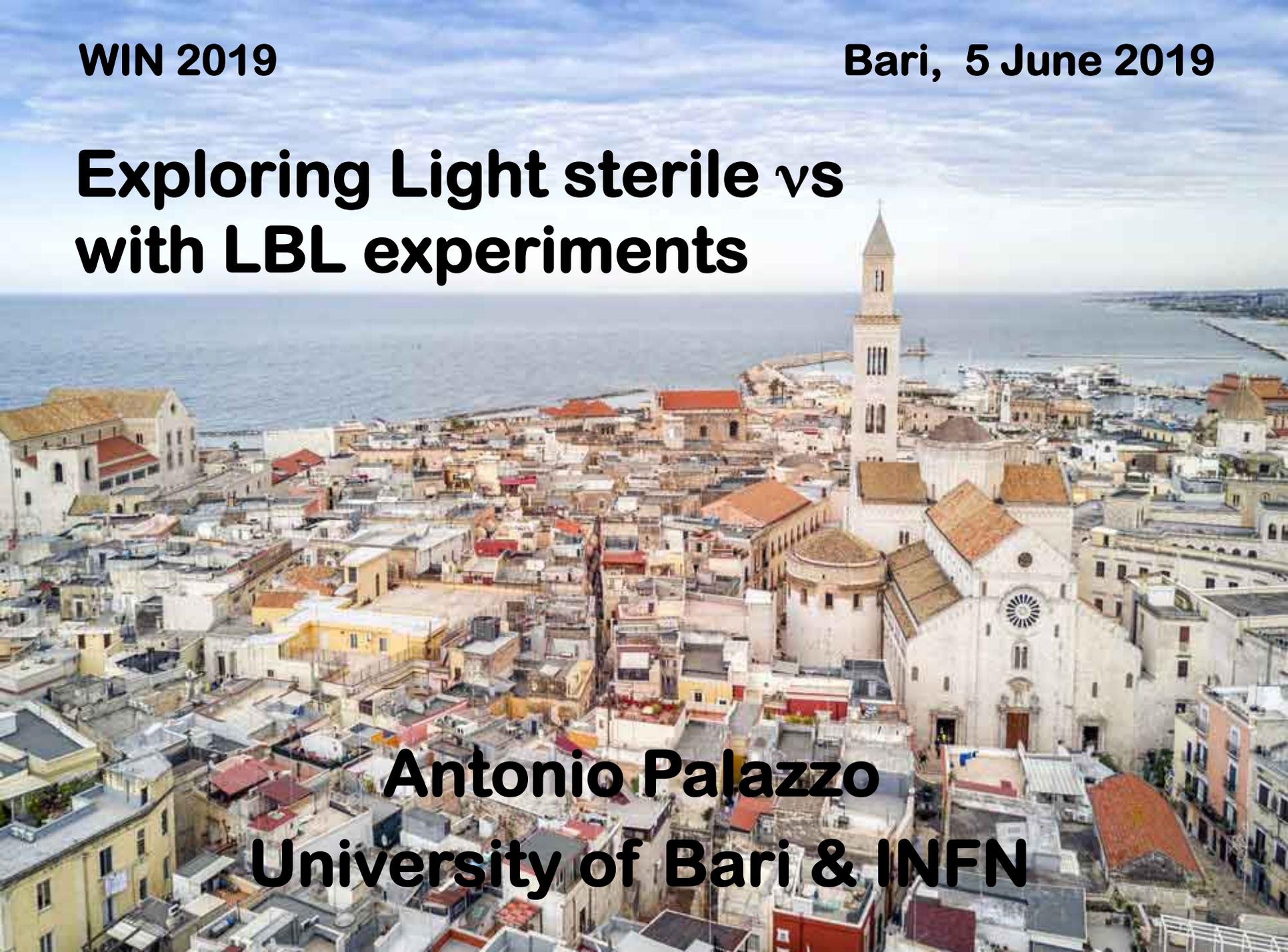


WIN 2019

Bari, 5 June 2019

Exploring Light sterile vs with LBL experiments

**Antonio Palazzo
University of Bari & INFN**



Outline

Introduction

Interference effects mediated by sterile neutrinos

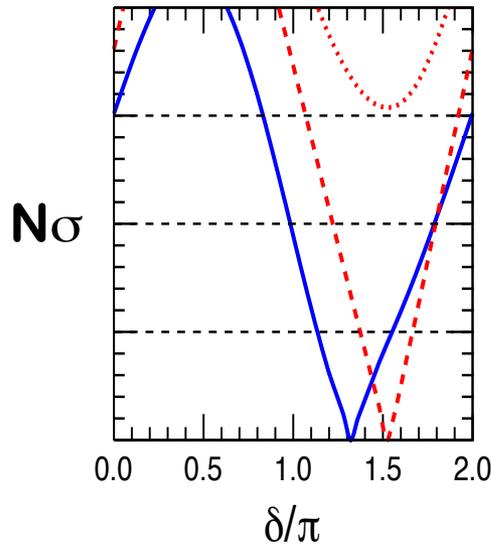
LBL constraints on new CP phases: present

LBL constraints on new CP phases: future

Conclusions

Introduction

It is timely to pose a new question



Capozzi, Lisi, Marrone, A.P,
PPNP 102, 48 (2018)

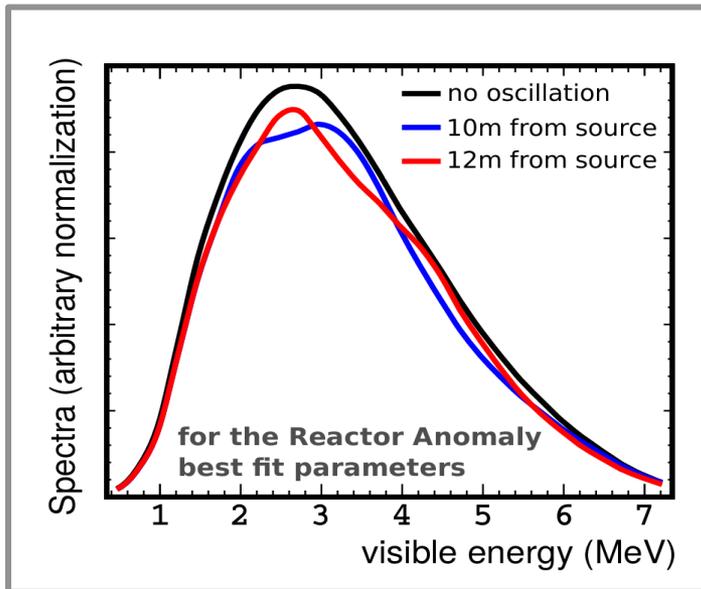
**LBL experiments start
to be sensitive to the
CP violating phase δ**

**Can sterile neutrinos generate observable CP
violating effects at LBL experiments?**

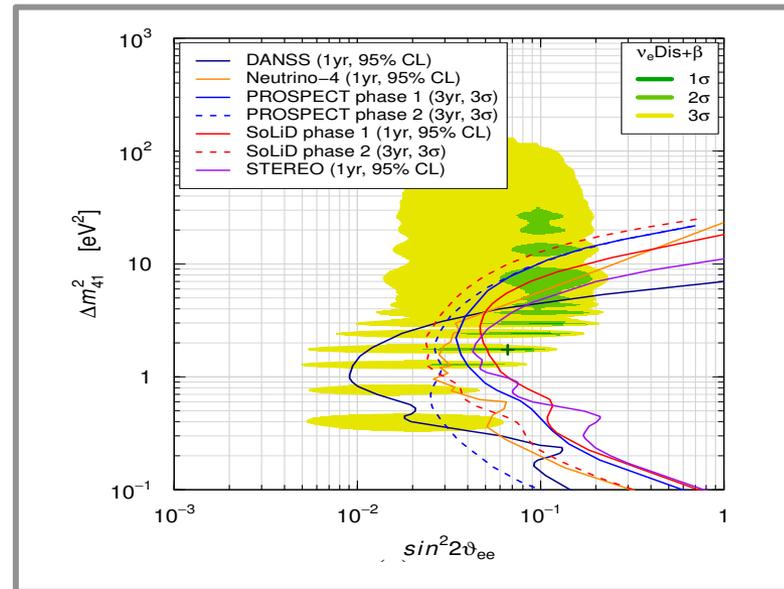
Question basically ignored in the past !

Most probably, the discovery of sterile ν_s can come only from SBL experiments

STEREO



Gariazzo et al., 1703.00860

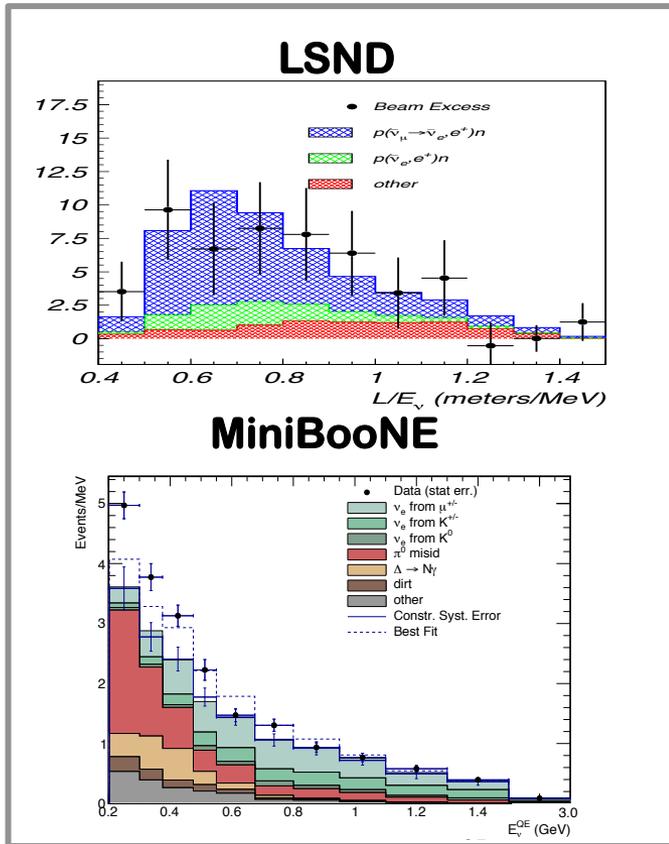


by observing the characteristic oscillation pattern
and we have already some hints...

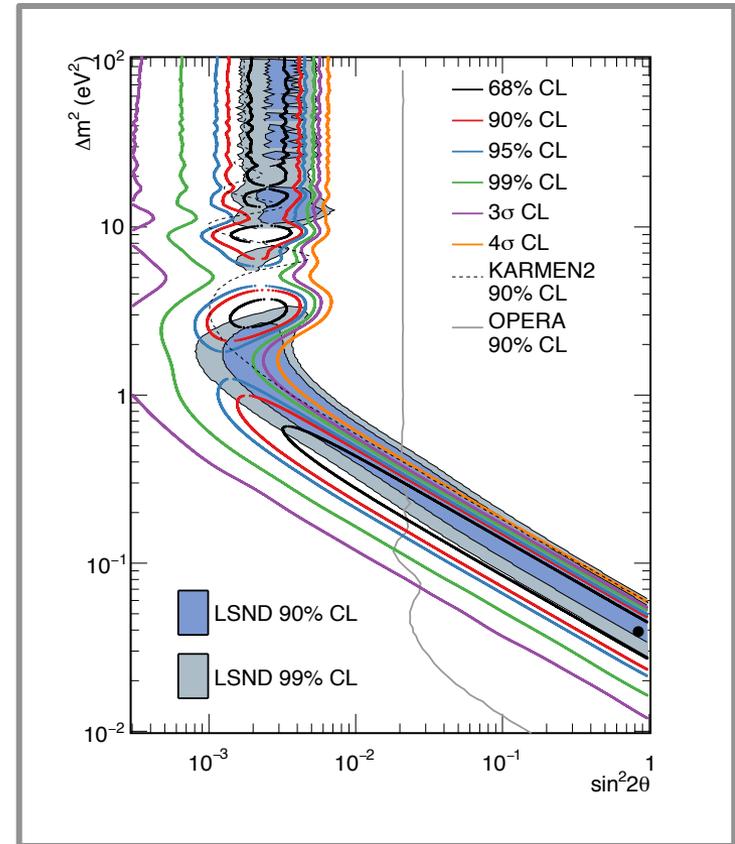
from accelerators

(unexplained ν_e appearance in a ν_μ beam)

3.8σ

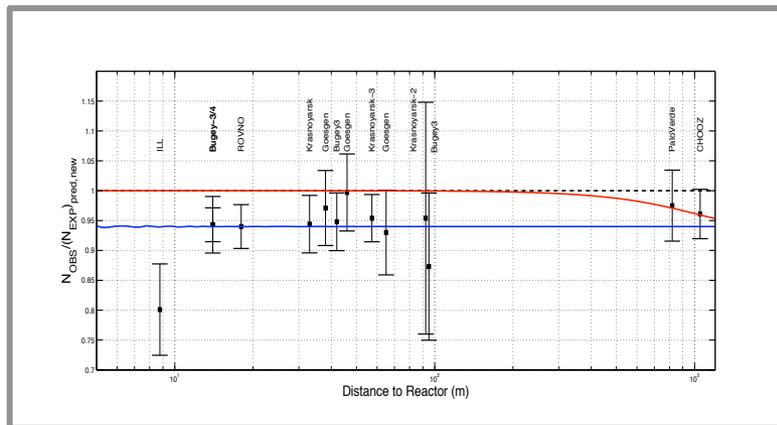


4.8σ

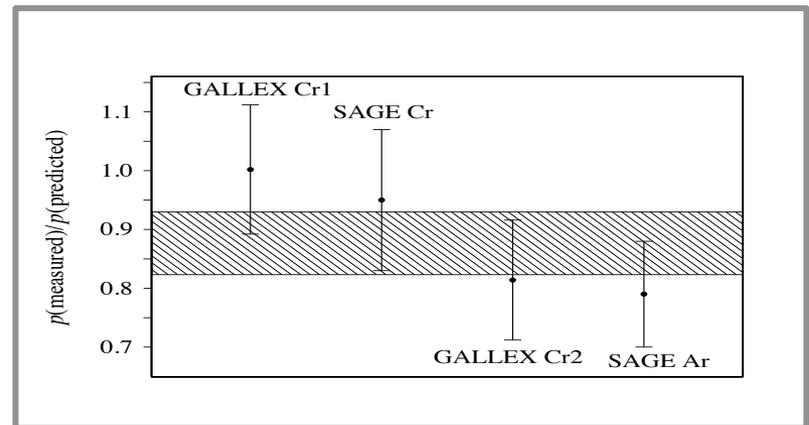


from reactor rates and solar calibration

(unexplained ν_e disappearance)



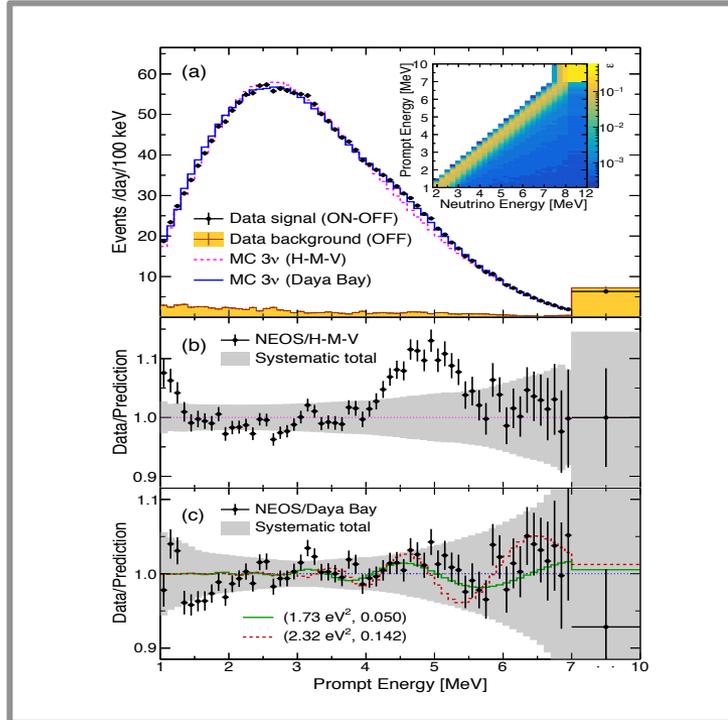
Mention et al. arXiv:1101:2755 [hep-ex]



SAGE coll., PRC 73 (2006) 045805

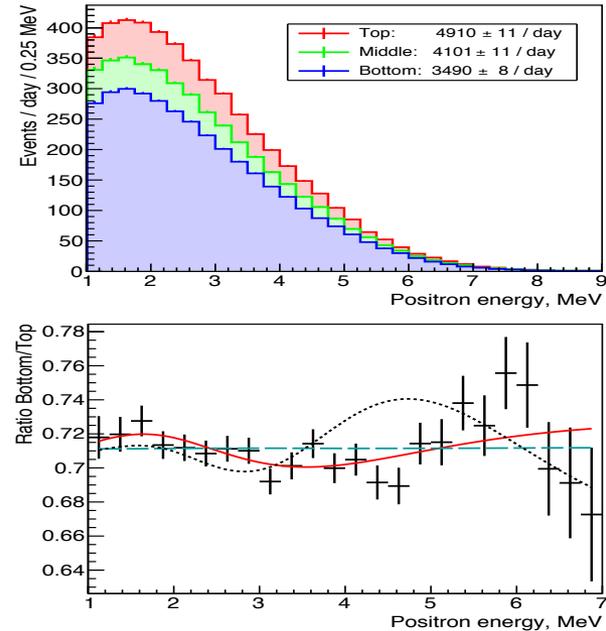
...and recently also from reactor spectra

NEOS arXiv:1610:05134



$$\chi_{4\nu}^2 - \chi_{3\nu}^2 = -6.5$$

DANSS arXiv:1804:04046



$$\chi_{4\nu}^2 - \chi_{3\nu}^2 = -13.1$$

Best fit points very similar: $(\sin^2 2\theta, \Delta m^2) \simeq (0.05, 1.4\text{eV}^2)$

However, SBL have an intrinsic limitation

At SBL atm/sol oscillations are negligible

$$\frac{L}{E} \sim \frac{m}{\text{MeV}}$$

$$\begin{aligned} \Delta_{12} &\simeq 0 \\ \Delta_{13} &\simeq 0 \end{aligned}$$

$$\Delta_{ij} = \frac{\Delta m_{ij}^2 L}{4E}$$

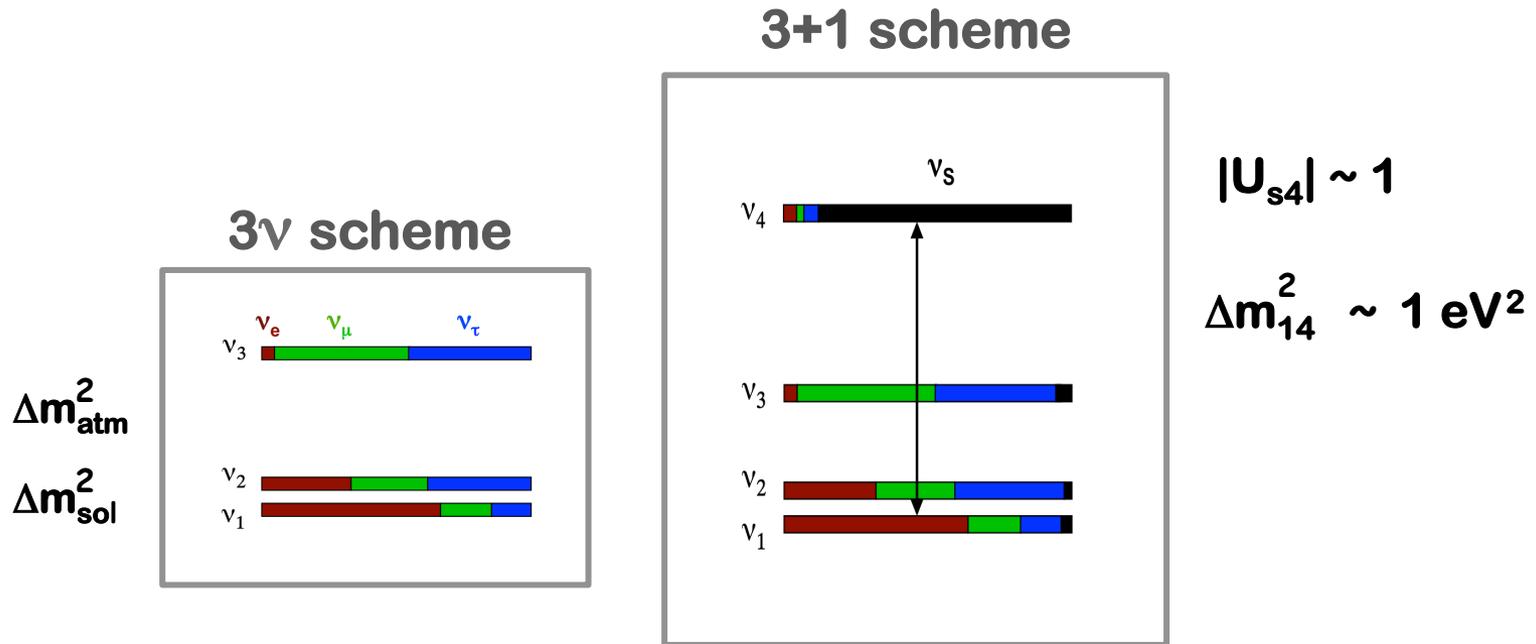
Impossible to observe phenomena of interference between the new frequency ($\Delta_{14} \sim 1$) and atm/sol ones

This limitation can be overcome at LBL's...

Interference effects mediated by sterile ν s

N. Klop & **A.P.**, PRD 91 073017 (2015) arXiv: 1412.7524

How to enlarge the 3-flavor scheme



At LBL the effective 2-flavor SBL description is no more valid and calculations should be done in the 3+1 (or 3+N_s) scheme

Mixing Matrix in the 3+1 scheme

$$U = \tilde{R}_{34} R_{24} \tilde{R}_{14} R_{23} \underbrace{\tilde{R}_{13} R_{12}}_{3\nu}$$

$$R_{ij} = \begin{bmatrix} c_{ij} & s_{ij} \\ -s_{ij} & c_{ij} \end{bmatrix}$$

$$\tilde{R}_{ij} = \begin{bmatrix} c_{ij} & \tilde{s}_{ij} \\ -\tilde{s}_{ij}^* & c_{ij} \end{bmatrix}$$

$$\begin{aligned} s_{ij} &= \sin \theta_{ij} \\ c_{ij} &= \cos \theta_{ij} \\ \tilde{s}_{ij} &= s_{ij} e^{-i\delta_{ij}} \end{aligned}$$

$$3\nu \begin{cases} 3 \text{ mixing angles} \\ 1 \text{ Dirac phase} \\ 2 \text{ Majorana phases} \end{cases}$$

$$3+1 \begin{cases} 6 \\ 3 \\ 3 \end{cases}$$

$$3+N \begin{cases} 3+3N \\ 1+2N \\ 2+N \end{cases}$$

In general, we have additional sources of CPV

LBL transition probability in 3-flavor

$$P_{\nu_\mu \rightarrow \nu_e}^{3\nu} = P^{\text{ATM}} + P^{\text{SOL}} + P^{\text{INT}}$$

In vacuum:

$$P^{\text{ATM}} = 4s_{23}^2 s_{13}^2 \sin^2 \Delta$$

$$P^{\text{SOL}} = 4c_{12}^2 c_{23}^2 s_{12}^2 (\alpha \Delta)^2$$

$$P^{\text{INT}} = 8s_{23}s_{13}c_{12}c_{23}s_{12}(\alpha \Delta) \sin \Delta \cos(\Delta + \delta_{CP})$$

$$\Delta = \frac{\Delta m_{31}^2 L}{4E}, \quad \alpha = \frac{\Delta m_{21}^2}{\Delta m_{31}^2}$$

$$\Delta \sim \pi/2$$

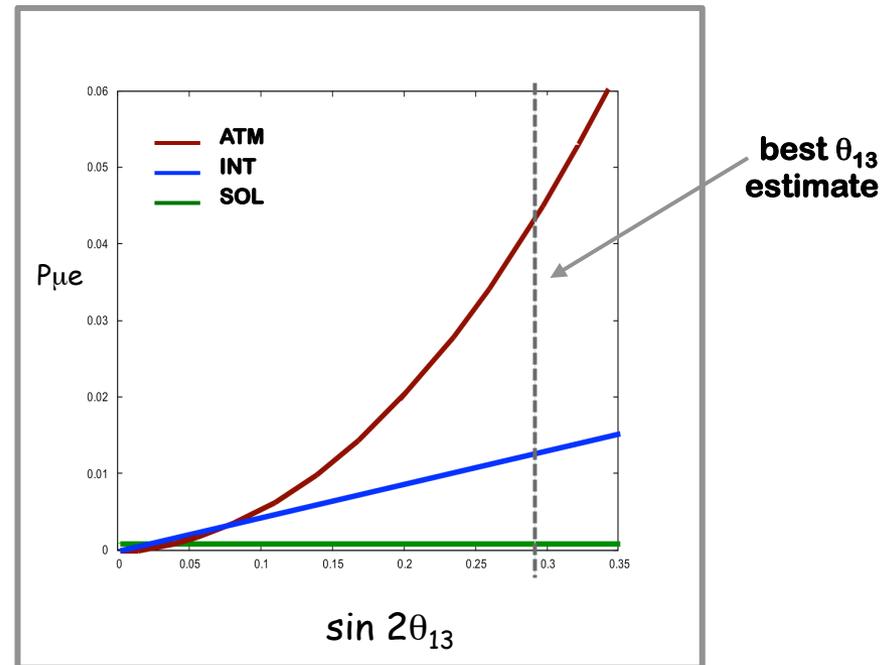
$$\alpha \sim 0.03$$

P^{ATM} leading $\rightarrow \theta_{13} > 0$

P^{INT} subleading \rightarrow dependency on δ

P^{SOL} negligible

T2K osc. maximum E = 0.6 GeV



A new interference term in the 3+1 scheme

N. Klop & A.P., PRD (2015)

- $\Delta_{14} \gg 1$: fast oscillations are averaged out

- But interference of Δ_{14} & Δ_{13} survives and is observable

$$P_{\mu e}^{4\nu} \simeq P^{\text{ATM}} + P_{\text{I}}^{\text{INT}} + P_{\text{II}}^{\text{INT}}$$

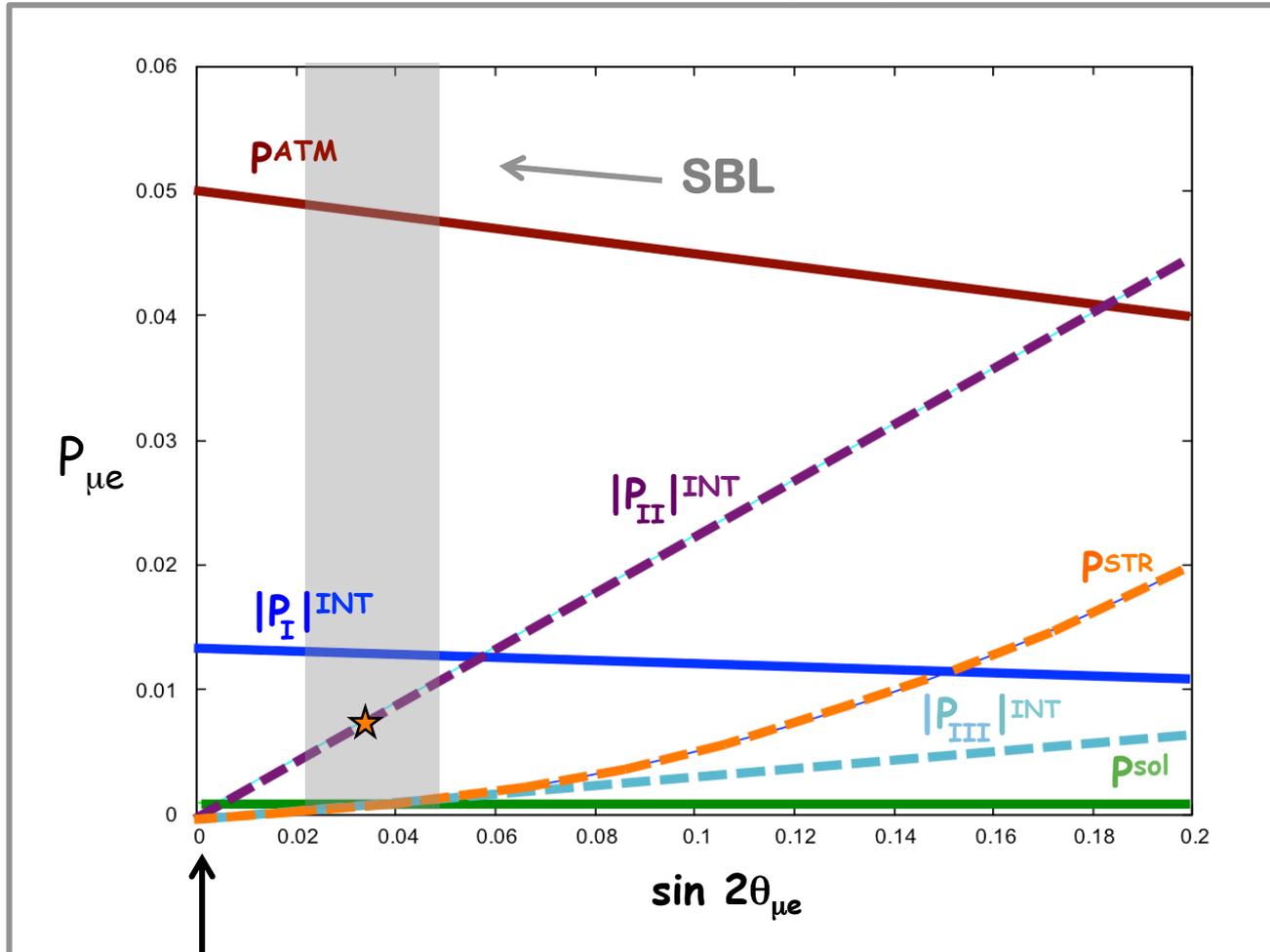
$$\begin{aligned} S_{13} \sim S_{14} \sim S_{24} &\sim 0.15 \sim \epsilon \\ \alpha = \delta m^2 / \Delta m^2 &\sim 0.03 \sim \epsilon^2 \end{aligned}$$

$$\left\{ \begin{aligned} P^{\text{ATM}} &\simeq 4s_{23}^2 s_{13}^2 \sin^2 \Delta && \sim \epsilon^2 \\ P_{\text{I}}^{\text{INT}} &\simeq 8s_{13} s_{23} c_{23} s_{12} c_{12} (\alpha \Delta) \sin \Delta \cos(\Delta + \delta_{13}) && \sim \epsilon^3 \\ P_{\text{II}}^{\text{INT}} &\simeq 4s_{14} s_{24} s_{13} s_{23} \sin \Delta \sin(\Delta + \delta_{13} - \delta_{14}) && \sim \epsilon^3 \end{aligned} \right.$$

Sensitivity to the new CP-phase δ_{14}

Amplitude of the new interference term

N. Klop & A.P., PRD (2015)

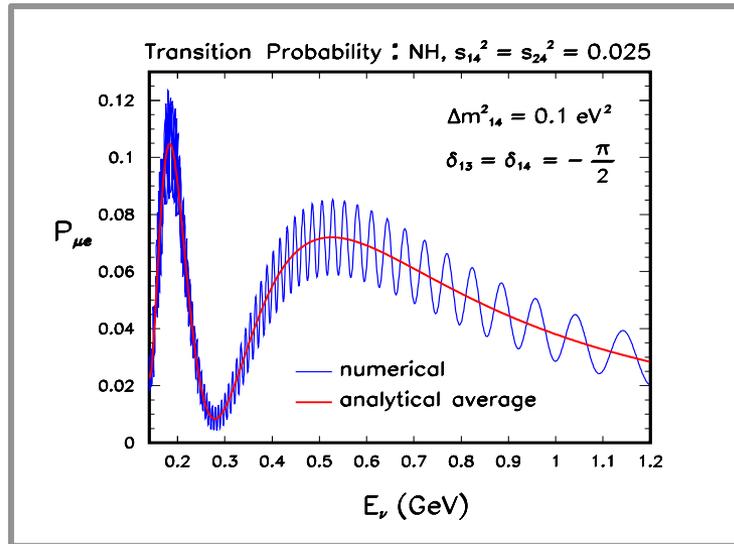


T2K
 $\theta_{13} = 9^\circ$
 $E = 0.6 \text{ GeV}$

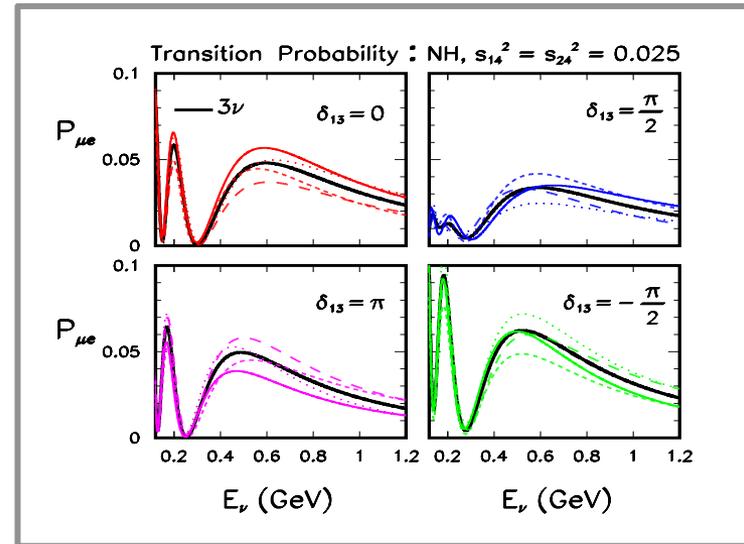
$$\sin^2 2\theta_{\mu e} = 4|U_{e4}|^2|U_{\mu 4}|^2$$

3ν limit

Numerical examples of 4ν probability



The fast oscillations get averaged out due to the finite energy resolution



Different line styles



Different values of δ_{14}

The modifications induced by δ_{14} are almost as large as those induced by the standard CP-phase δ_{13}

Consequences...

LBL constraints on sterile ν s: present

A.P., PRD (RC) 91, 091301 (2015)
arXiv:1503.03966

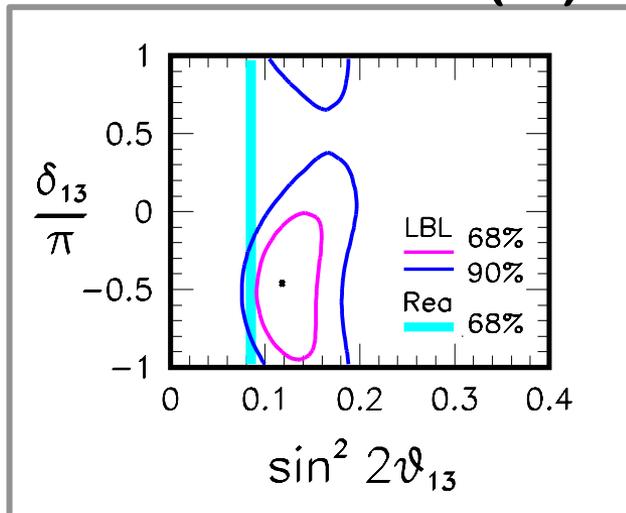
A.P., PLB 757, 142 (2016)
arXiv:1509.03148

Capozzi, Giunti, Laveder & **A.P.**,
PRD 95 (2017)
arXiv:1612.07764

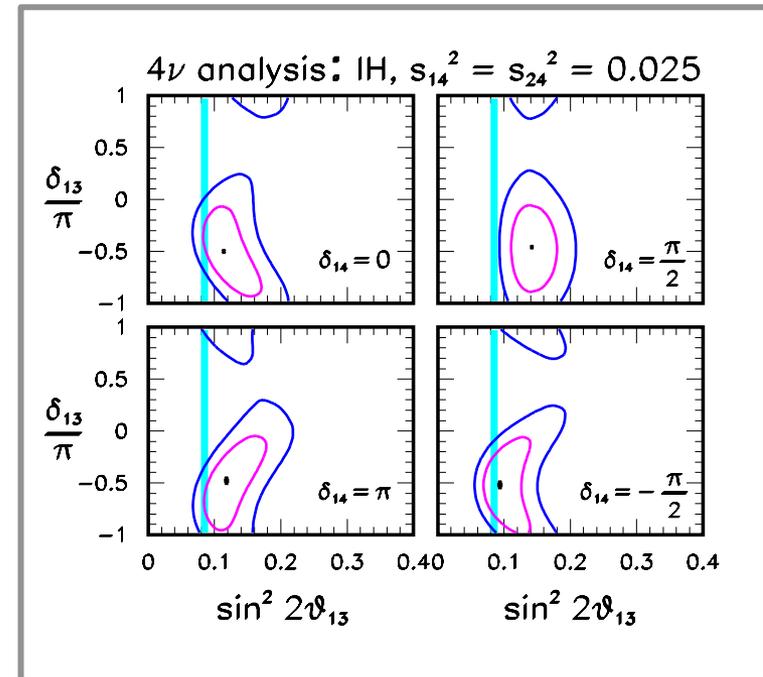
LBL constraints change in the 3+1 scheme

PLB (2016)

3 ν : T2K + NO ν A (IH)



4 ν →

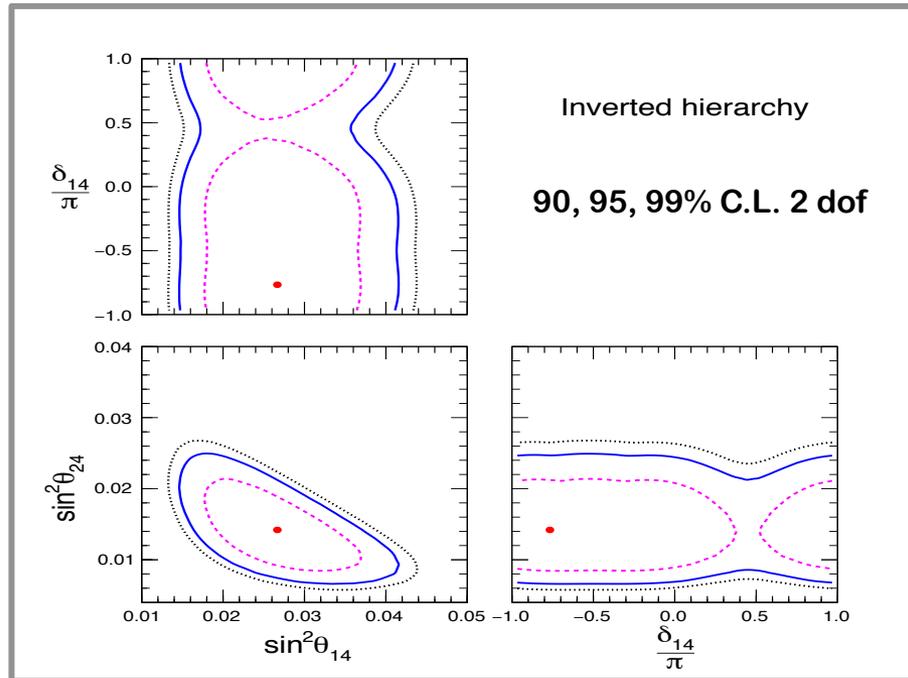


- The level of (dis-)agreement of LBL & Rea. depends on δ_{14}
- In this analysis θ_{14} and θ_{24} are fixed at the SBL best fit values
- These results call for a more refined analysis ...

Joint SBL and LBL constraints on $[\theta_{14}, \theta_{24}, \delta_{14}]$

PRD (2017)

SBL + LBL

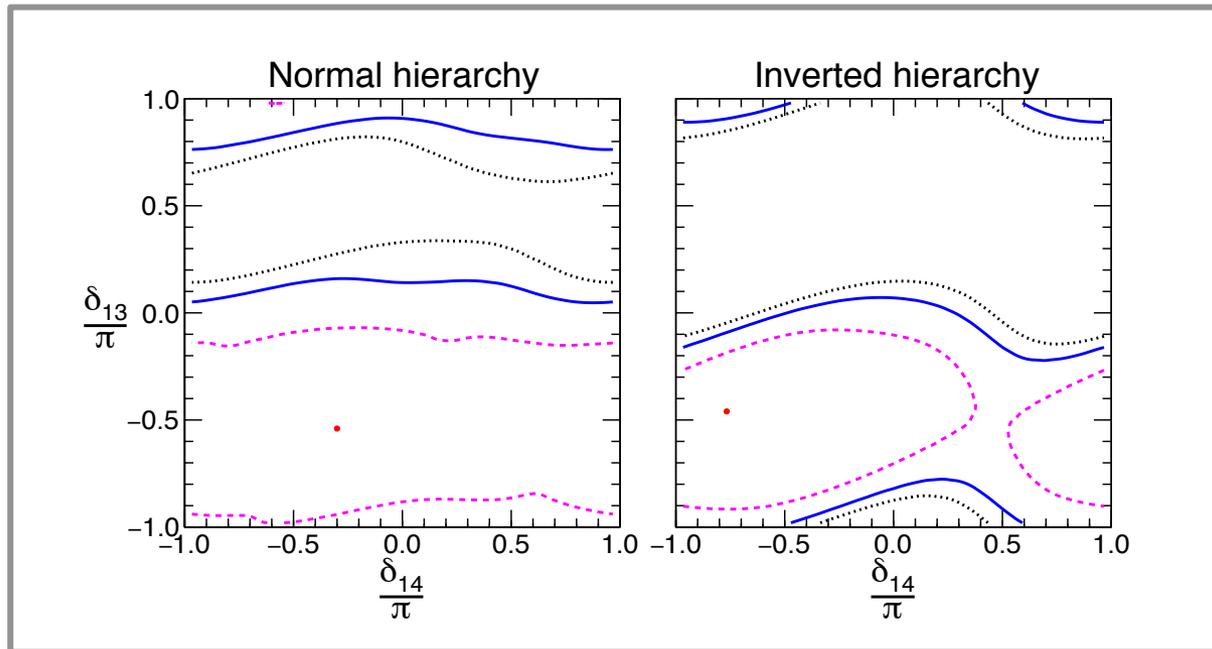


- $[\theta_{14}, \theta_{24}]$ determined by SBL experiments
- δ_{14} constrained by LBL experiments

Constraints on the two CP-phases

PRD (2017)

SBL + LBL

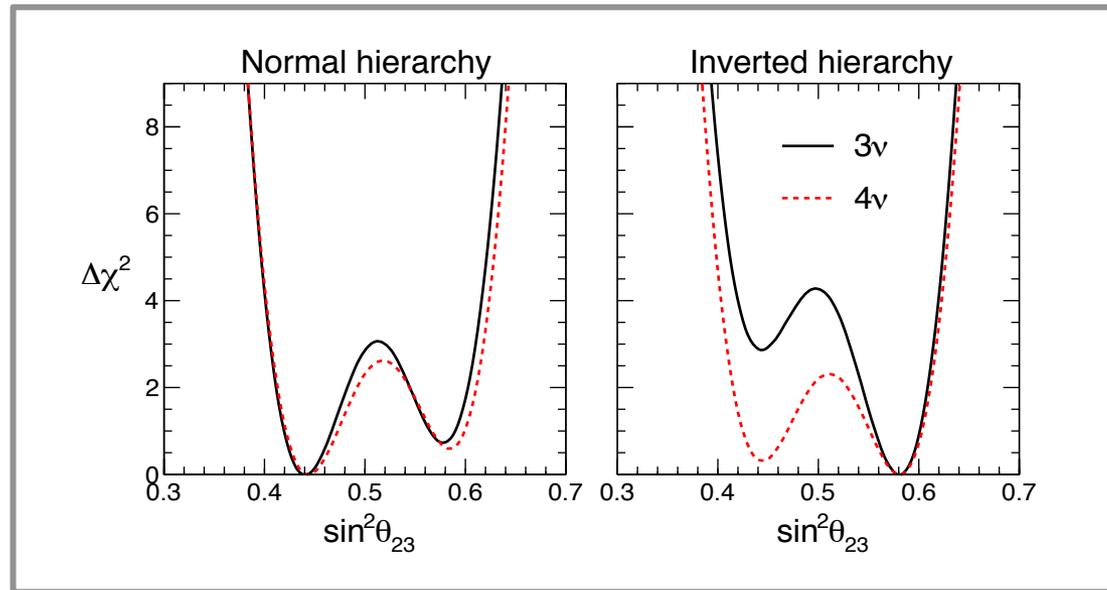


- δ_{13} is more constrained than δ_{14}
- Best fit values: $\delta_{13} \sim \delta_{14} \sim -\pi/2$
- This information cannot be extracted from SBL alone !

Impact of sterile neutrinos on θ_{23}

PRD (2017)

SBL + LBL



Indication for non-maximal θ_{23} persists in 3+1 scheme

Preference for θ_{23} octant disappears in 3+1 scheme

Octant fragility seems to be a general feature (see later)

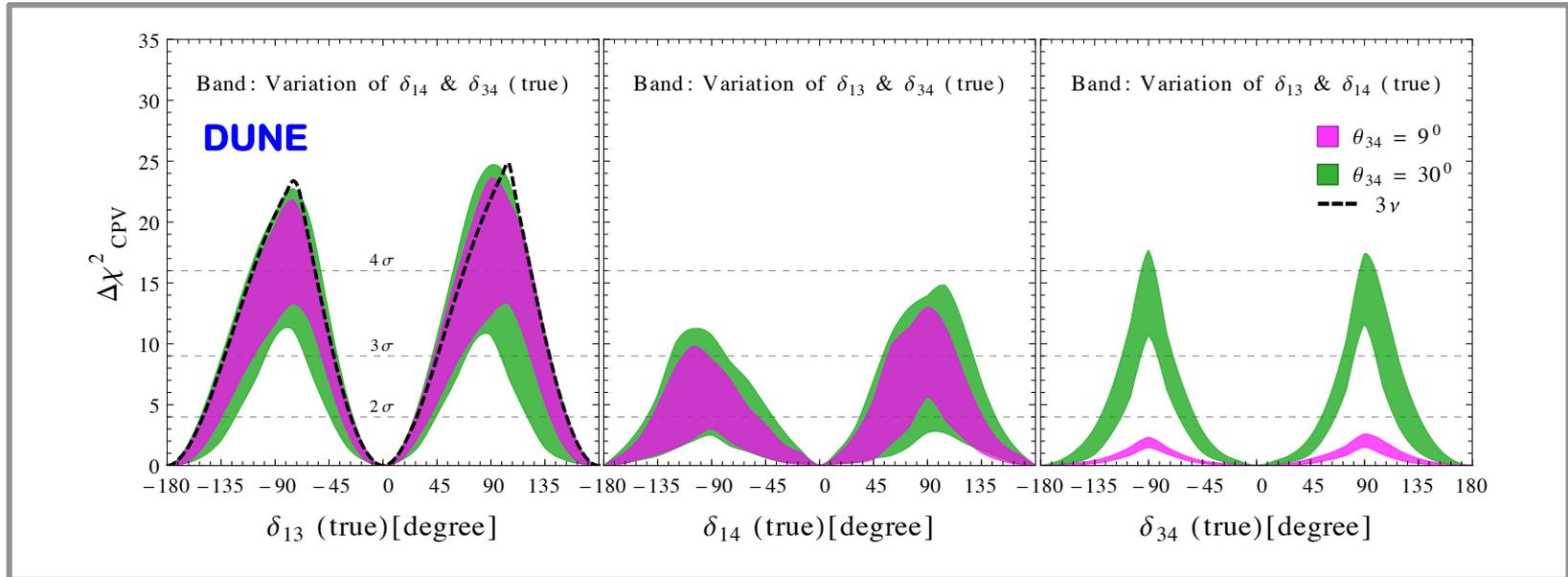
Looking to the future

Agarwalla, Chatterjee, Dasgupta, **A.P.**,
arXiv: 1601.05995 (JHEP 2016)

Agarwalla, Chatterjee, **A.P.**,
arXiv: 1603.03759 (JHEP 2016)
arXiv: 1607.01745 (PLB 2016)
arXiv: 1605.04299 (PRL 2017)
arXiv: 1801.04855 (JHEP 2018)
arXiv: 1906.XXXX (in preparation)

CPV discovery potential

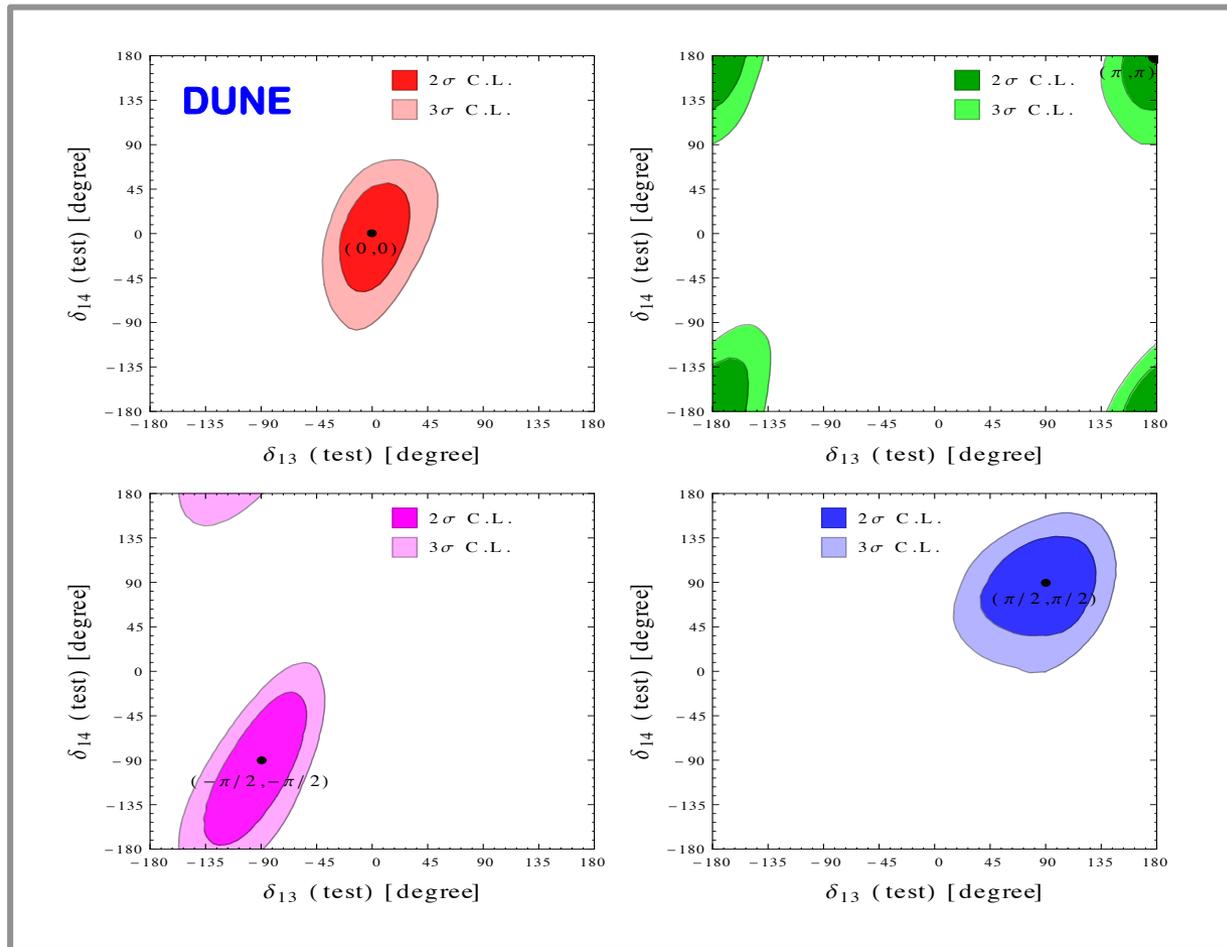
JHEP 2016



- Sensitivity to CPV induced by δ_{13} reduced in 3+1 scheme
- Potential sensitivity also to the new CP-phases δ_{14} e δ_{34}
- Clear hierarchy in the sensitivity: $\delta_{13} > \delta_{14} > \delta_{34}$ for $\theta_{14} = \theta_{24} = \theta_{34} = 9^\circ$

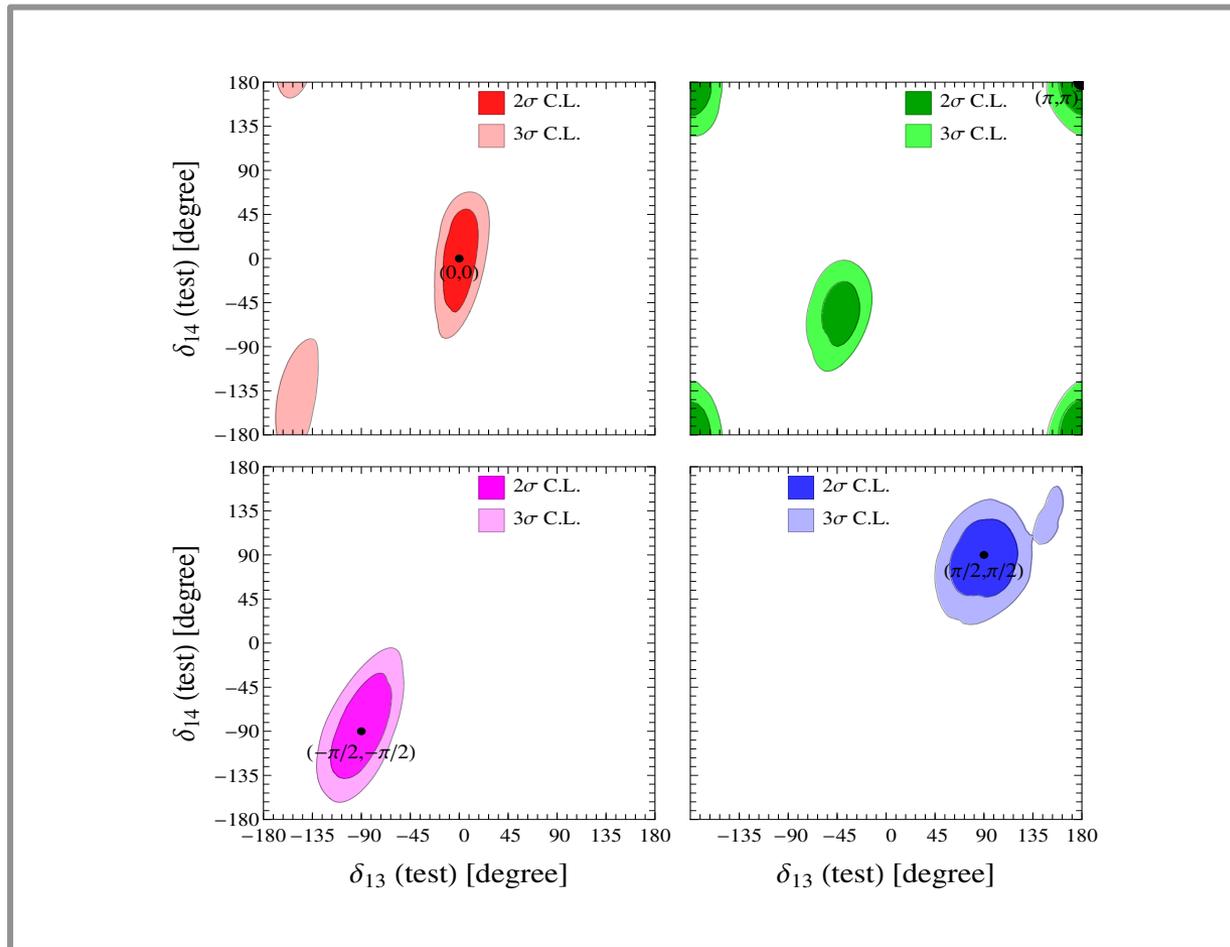
Reconstruction of the CP phases in DUNE

JHEP 2016



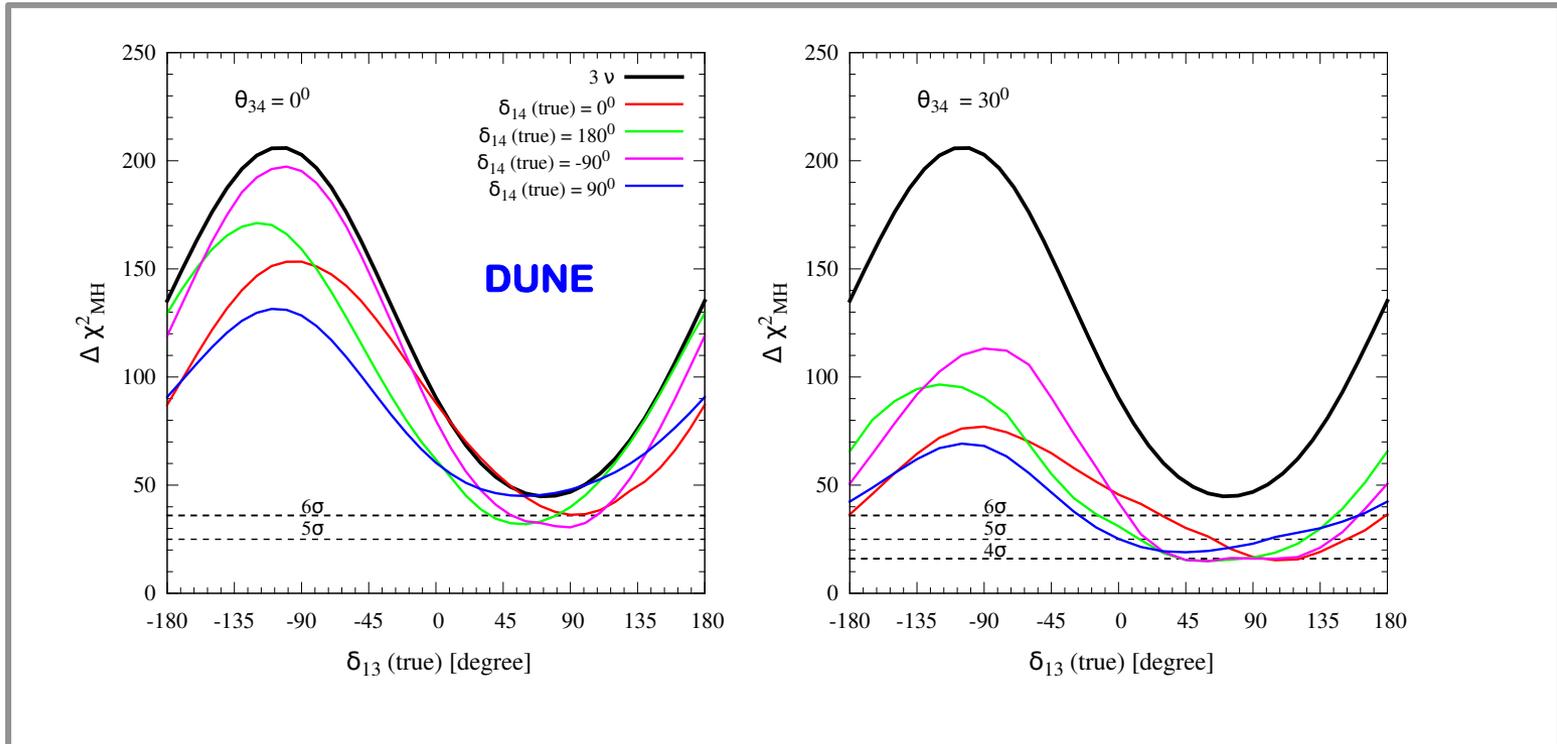
Reconstruction of the CP phases in T2HK

JHEP 2018



Discovery potential of mass hierarchy

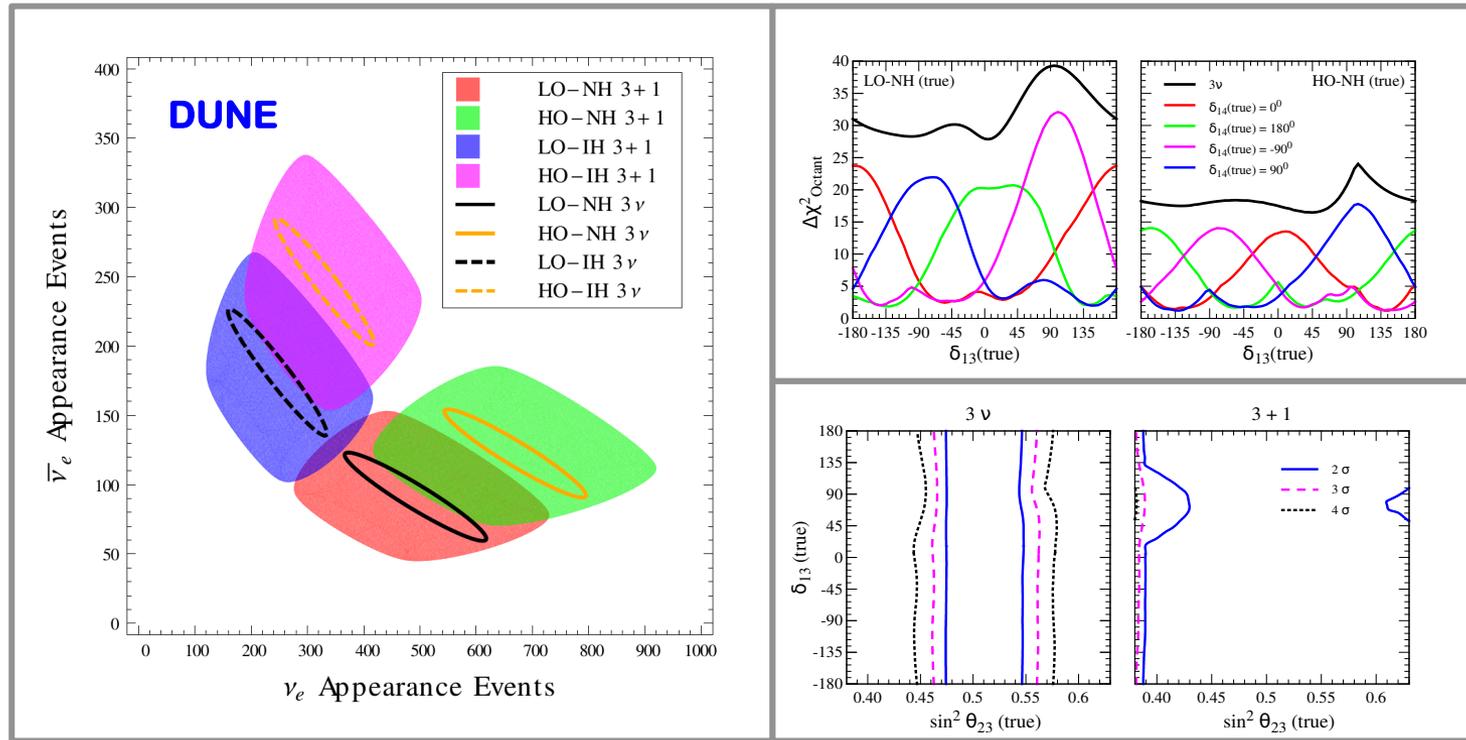
JHEP 2016



Degradation of sensitivity but 4 σ level preserved

Octant of θ_{23} in danger with a sterile neutrino

PRL 2017



Distinct ellipses (3ν) become overlapping blobs (3+1)

For unfavorable combinations of δ_{13} & δ_{14} sensitivity is lost

Conclusions

- **Sterile neutrinos are sources of additional CPV**
- **Consequences for the LBL estimates of the standard parameters (MH, CP-phase δ , octant of θ_{23})**
- **Full exploration of new CP-phase (δ_{14}, δ_{34}) possible only with LBL's**
- **LBL experiments complementary to the SBL ones**

**Thank you
for your attention!**

Back up slides

CPV and averaged oscillations

$$A_{\alpha\beta}^{\text{CP}} \equiv P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$$

$$A_{\alpha\beta}^{\text{CP}} = -16 J_{\alpha\beta}^{12} \sin \Delta_{21} \underbrace{\sin \Delta_{13} \sin \Delta_{32}}$$

if

$$\Delta \equiv \Delta_{13} \simeq \Delta_{23} \gg 1$$

osc. averaged out by finite E resol.

→

$$\langle \sin^2 \Delta \rangle = 1/2$$

It can be:

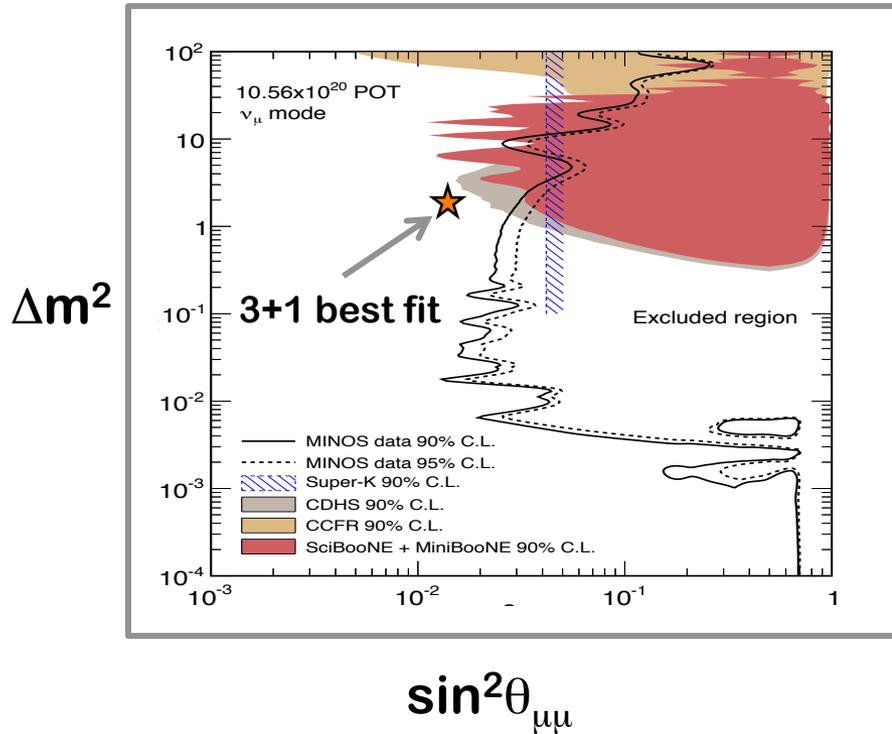
$$A_{\alpha\beta}^{\text{CP}} \neq 0$$

(if $\sin \delta = \emptyset$)

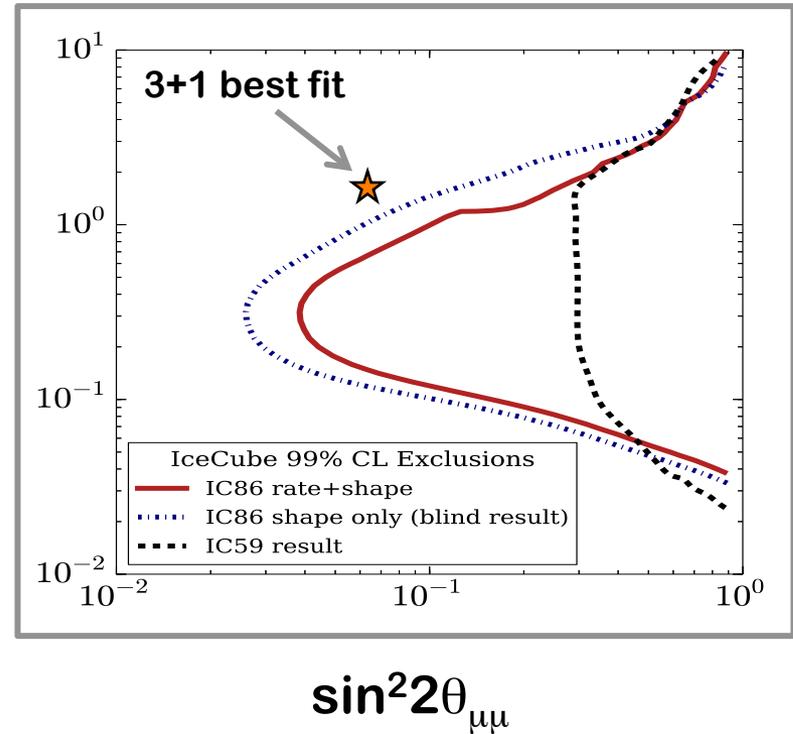
The bottom line is that if one of the three ν_i is ∞ far from the other two ones this does not erase CPV
(relevant for the 4 ν case)

No anomaly in ν_μ disappearance

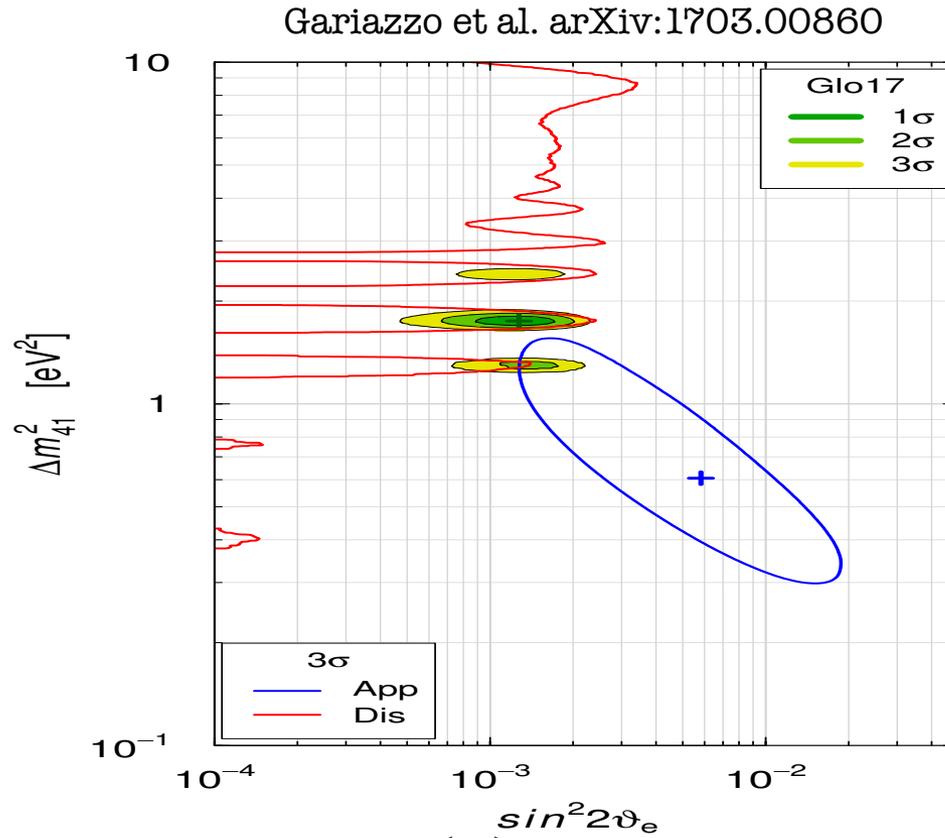
SBL & MINOS (NC)



IceCube

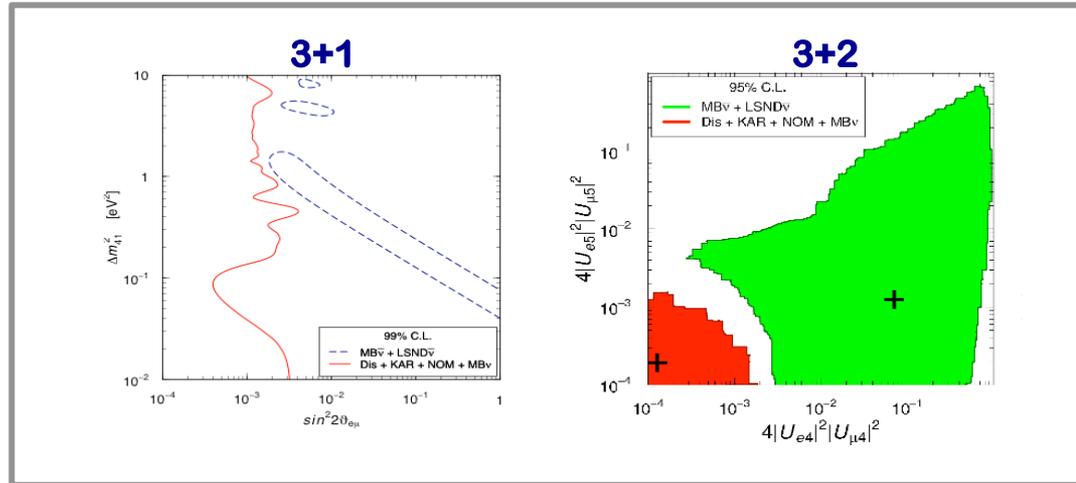


Global SBL data fit in the 3+1 scheme



There is strong internal tension

Tension in all ν_s models



Giunti
&
Laveder

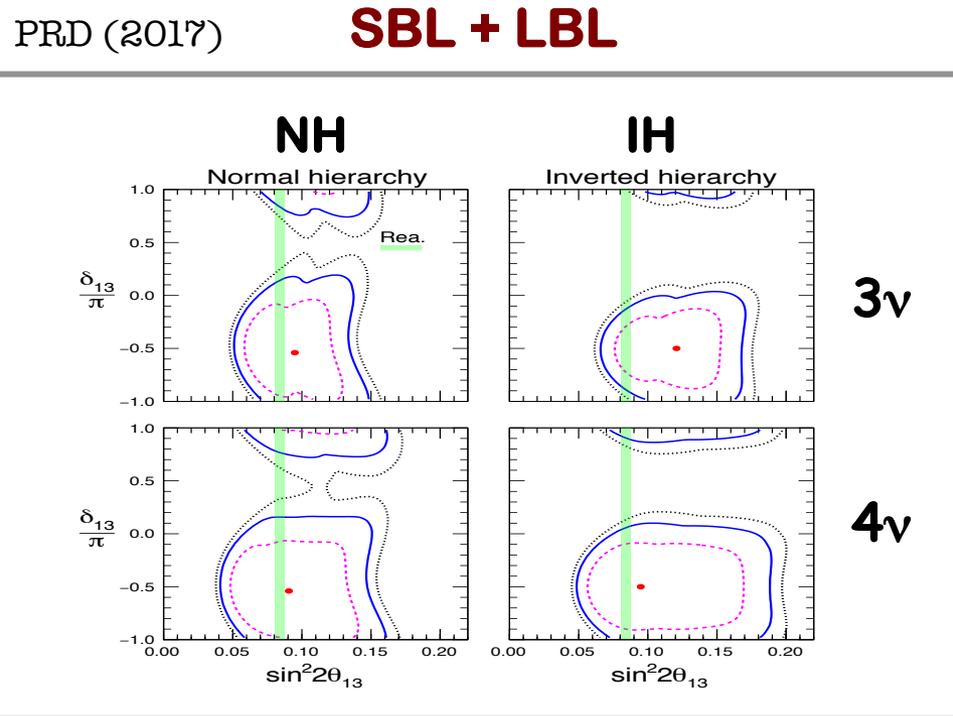
arXiv:1107.1452

$\nu_\mu \rightarrow \nu_e$ **positive**
 $\nu_e \rightarrow \nu_e$ **positive**
 $\nu_\mu \rightarrow \nu_\mu$ **negative**

$|U_{e4}||U_{\mu4}| > 0$
 $|U_{e4}| > 0$
 $|U_{\mu4}| \sim 0$

$$\sin^2 2\theta_{e\mu} \simeq \frac{1}{4} \sin^2 2\theta_{ee} \sin^2 2\theta_{\mu\mu} \simeq 4|U_{e4}|^2|U_{\mu4}|^2$$

Impact on the standard parameters $[\theta_{13}, \delta_{13}]$



- Allowed range for θ_{13} from LBL alone gets enlarged
- Values preferred for $\delta_{13} \equiv \delta$ basically unaltered
- Mismatch (in IH) of LBL and Reactors decreases in 3+1