

Competition between the rotational effect and the finite-size effect on relativistic fermions

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- HL. Chen, K. Fukushima, XG. Huang, KM, PRD 93, 104052 (2016)
S. Ebihara, K. Fukushima, KM, PLB 764, 94 (2017)
HL. Chen, K. Fukushima, XG. Huang, KM, PRD 96.054032 (2017)

Rotating Relativistic Systems

binary star merger



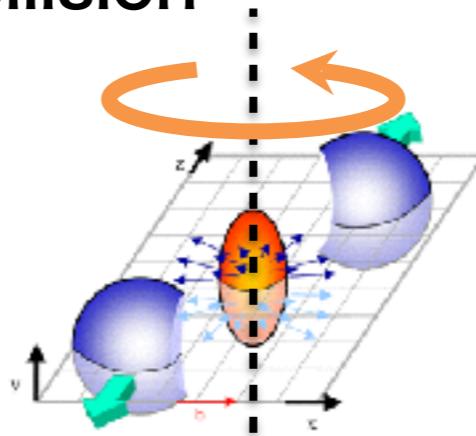
NASA/Tod Strohmayer

black holes

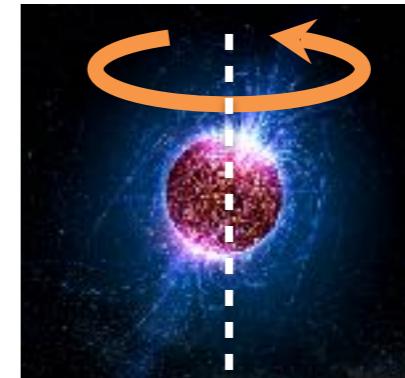


NASA/JPL-Caltech

heavy-ion collision



neutron stars



rotating (and magnetized) QCD systems

Heavy-ion collisions

$$eB \sim 10^{18} \text{ G}$$

Skokov et al. (2009)

$$\Omega \sim 10 \text{ MeV}$$
 (local rotation)

Jiang, Lin, Liao (2016)
Deng, Huang (2016)

Magnetar

$$eB \sim 10^{15} \text{ G}$$
 (surface)

Duncan, Thompson (1992)

$$eB \sim 10^{18-20} \text{ G}$$
 (interior)

Lai, Shapiro (1991) Ferrer et al. (2010)

$$R\Omega \sim 10^{-1}$$

Marshall et al. (2004)

Contents

Part I Finite-size system with Ω

Part II Finite-size system with Ω and eB

Contents (& Goals)

Part I Finite-size system with Ω

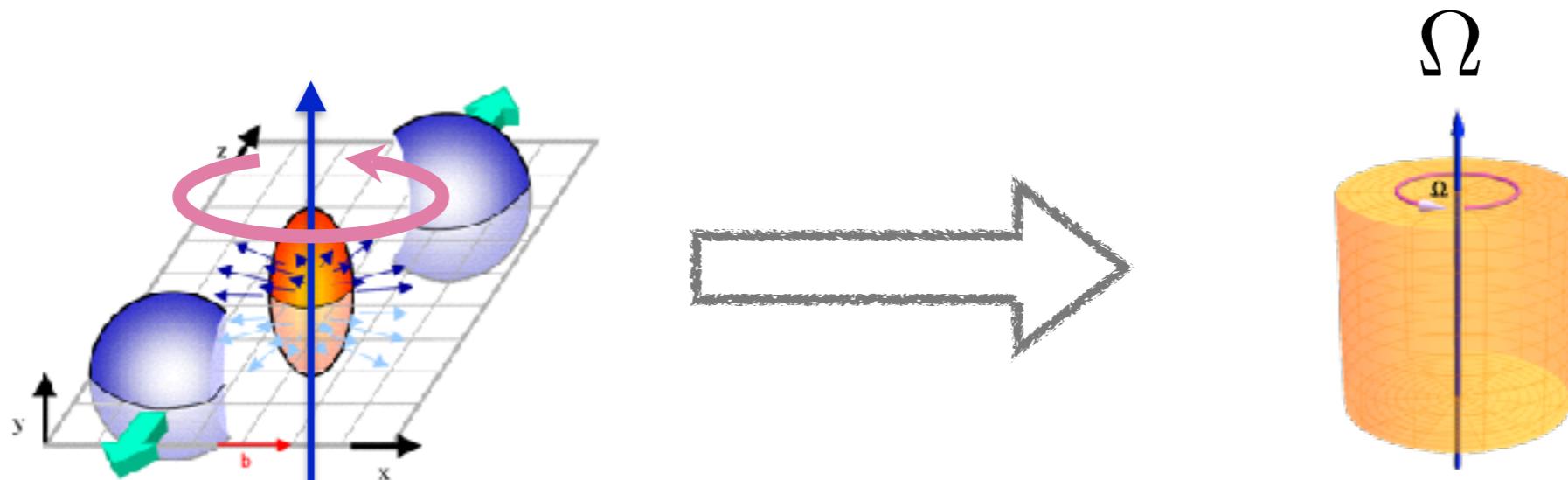
No rotational effect at zero temperature

Part II Finite-size system with Ω and eB

B -field enables rotation to affect thermodynamics

Part I Finite-size system with Ω

Rigidly Rotating System



Chernodub,
Gongyo (2016)

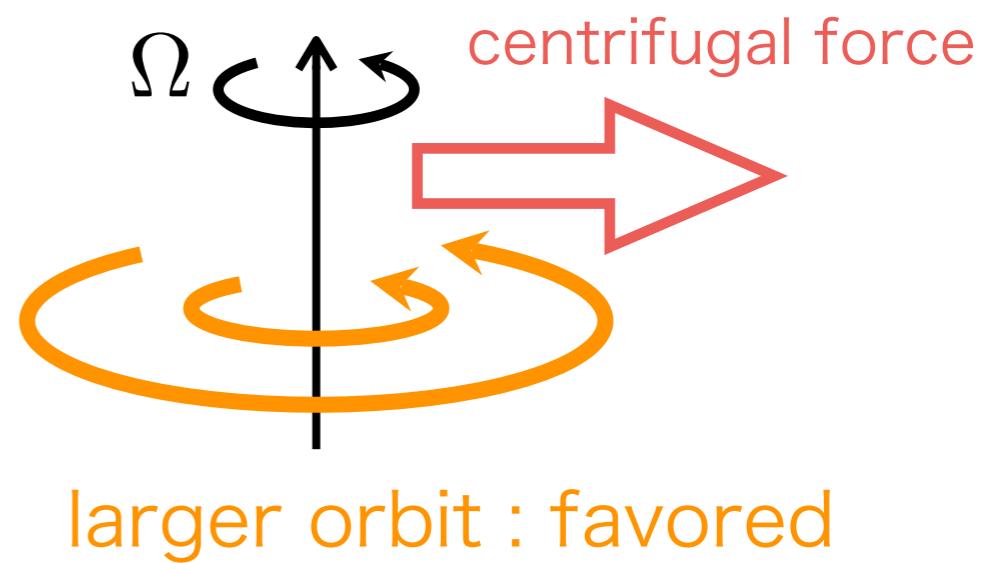
$$H \rightarrow H - \Omega \cdot L$$

μ_{eff}

Ex.) rotating bosons

$$n_{\text{BE}} = \frac{1}{e^{\beta(E-\Omega\ell)} - 1}$$

Landau, Lifshitz (1958)



Finite-size System

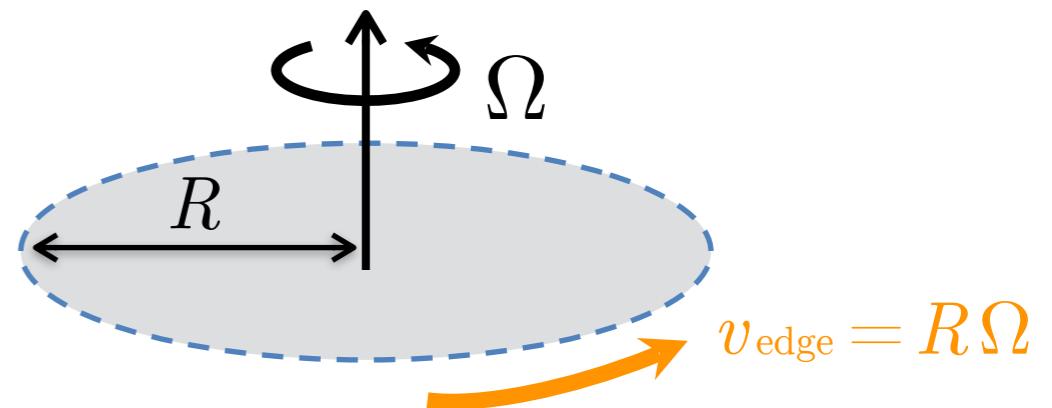
$$n_{\text{BE}} = \frac{1}{e^{\beta(E - \Omega\ell)} - 1}$$

$n_{\text{BE}} < 0$ for large Ω ??

causality constraint

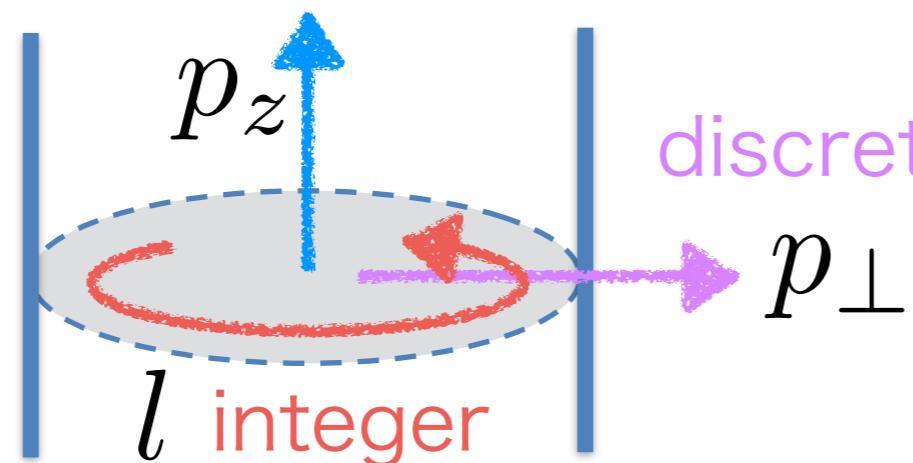
$$v_{\text{edge}} = \Omega R \leq 1$$

→ $R \leq 1/\Omega < \infty$



Rotating systems must be finite-size

continuous



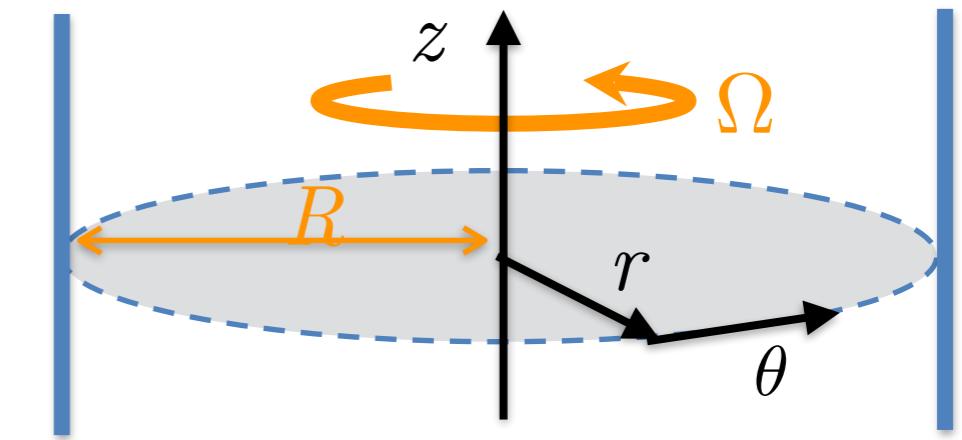
discrete

finite-size effect

Dirac eq. in Cylinders

$$ds^2 = dt^2 - dr^2 - r^2 d\theta^2 - dz^2$$

$$[i\gamma^\mu(\partial_\mu + \Gamma_\mu) - m]\psi = 0$$



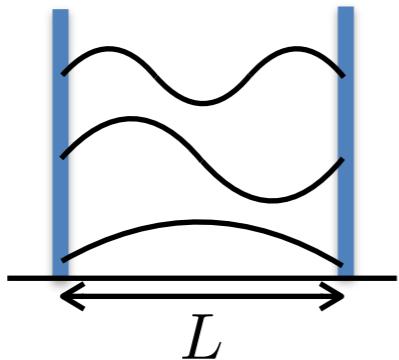
$$\psi = (\psi_1, \psi_2, \psi_3, \psi_4)^T$$

$$\longrightarrow \left[-\partial_0^2 + \partial_z^2 - m^2 + \partial_r^2 + \frac{1}{r}\partial_r + \frac{1}{r^2}\partial_\theta^2 \right] \psi_1 = 0$$

$$\longrightarrow \psi_1 = e^{-ip_0 t + ip_z z + il\theta} J_l(p_\perp r)$$

Momentum Discretization

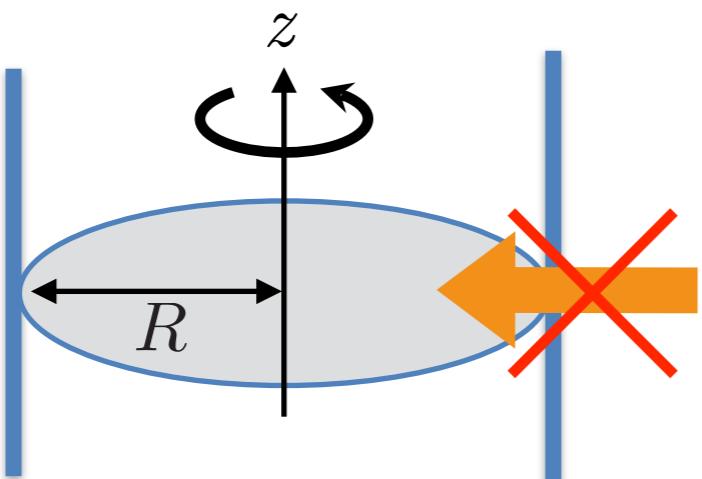
Ex.) Dirichlet b.c.



$$\sin(px)|_{x=L} = 0$$

$$\longrightarrow \quad p = \frac{n\pi}{L} \geq \frac{\pi}{L}$$

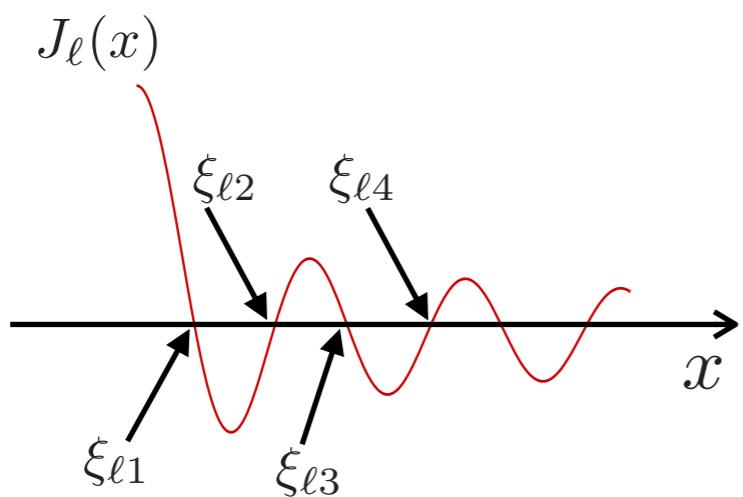
NO incoming current



$$\int dz d\theta \bar{\psi} \gamma^r \psi \Big|_{r=R} = 0$$

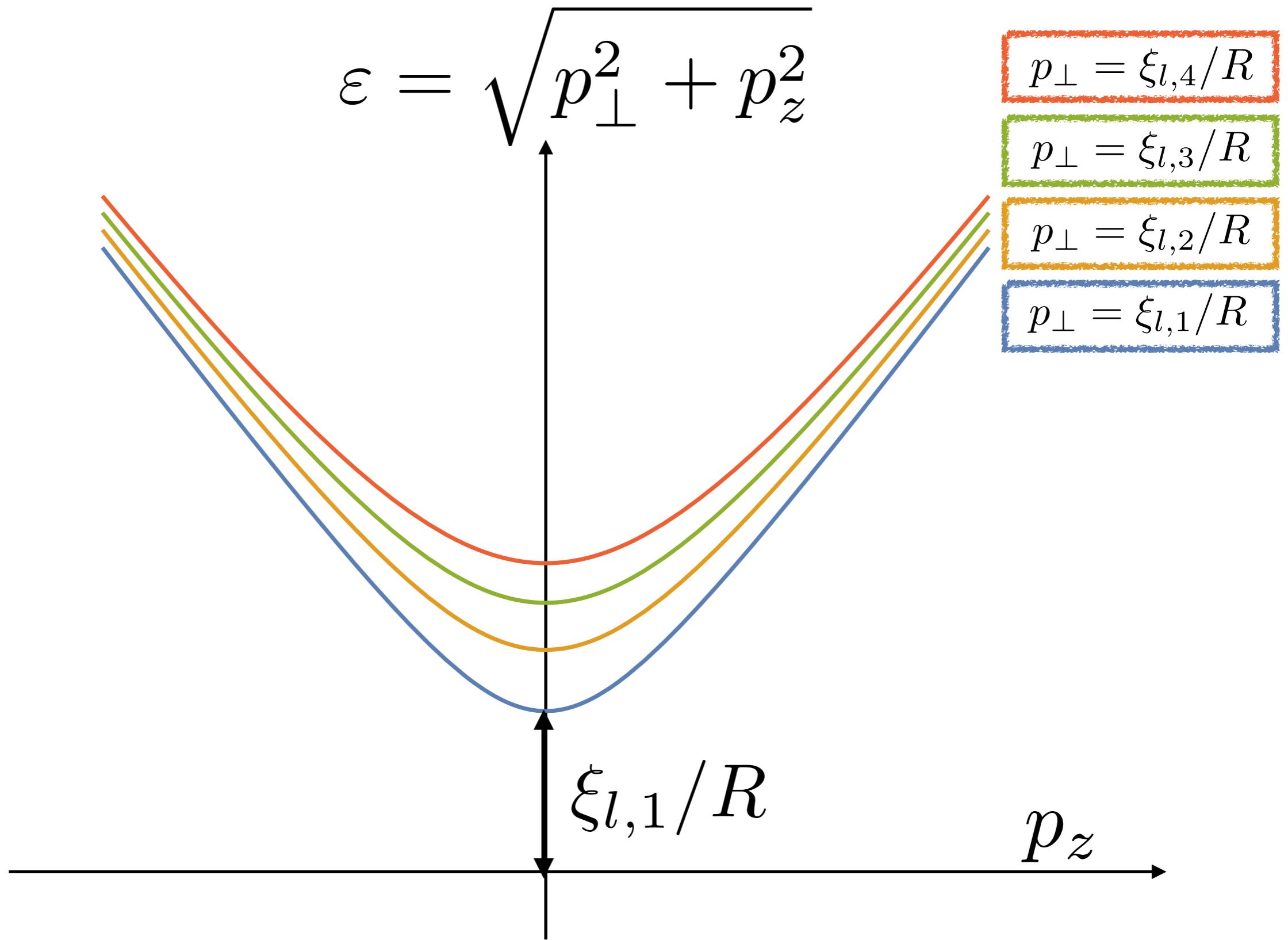
$$\longrightarrow \quad p_\perp \geq \frac{\xi_{l,1}}{R} \simeq \frac{2.4}{R}$$

IR gapped mode



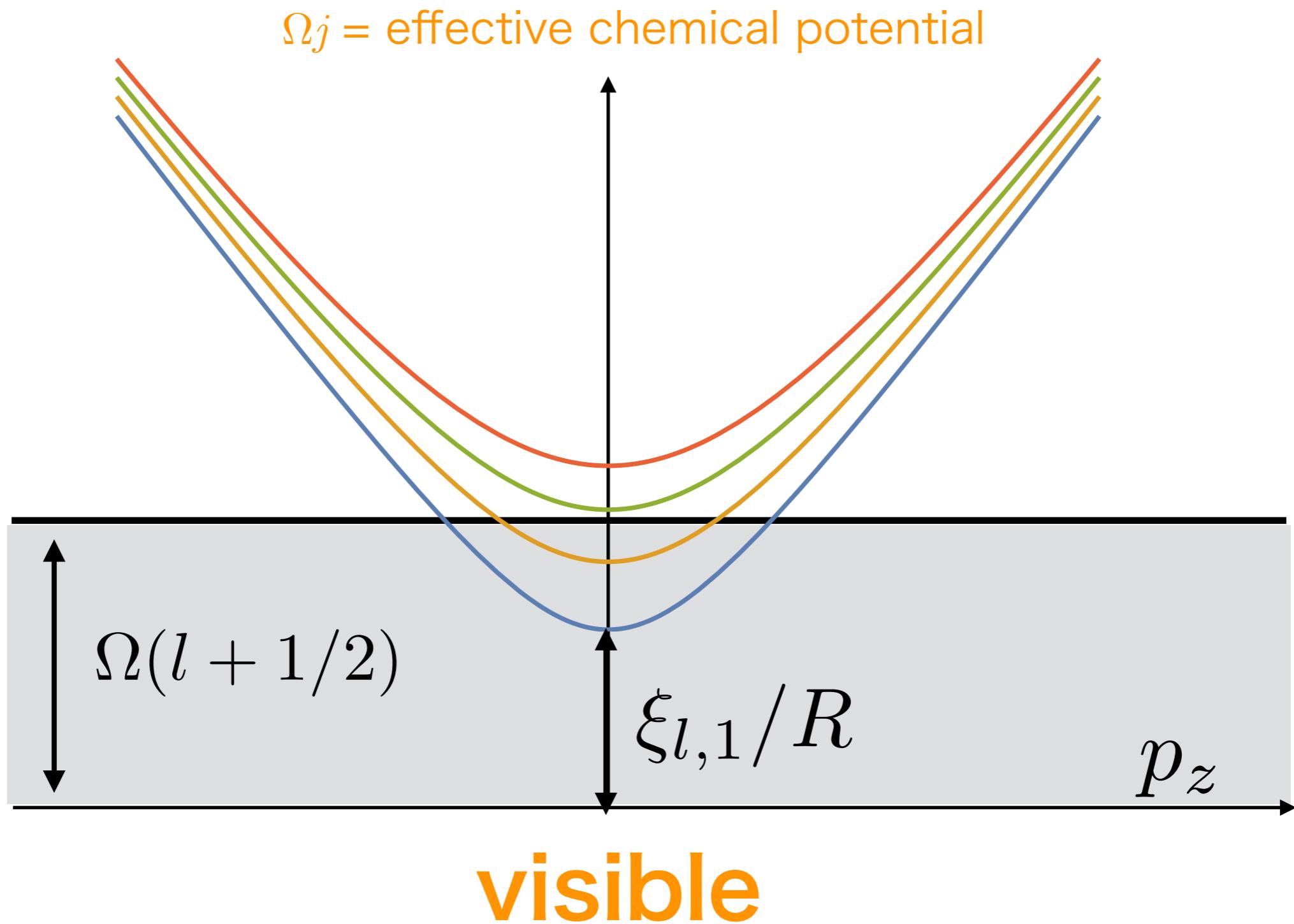
$x = \xi_{l,k}$: the k th root of $J_l(x)$

Rotational Effect at $T = 0$



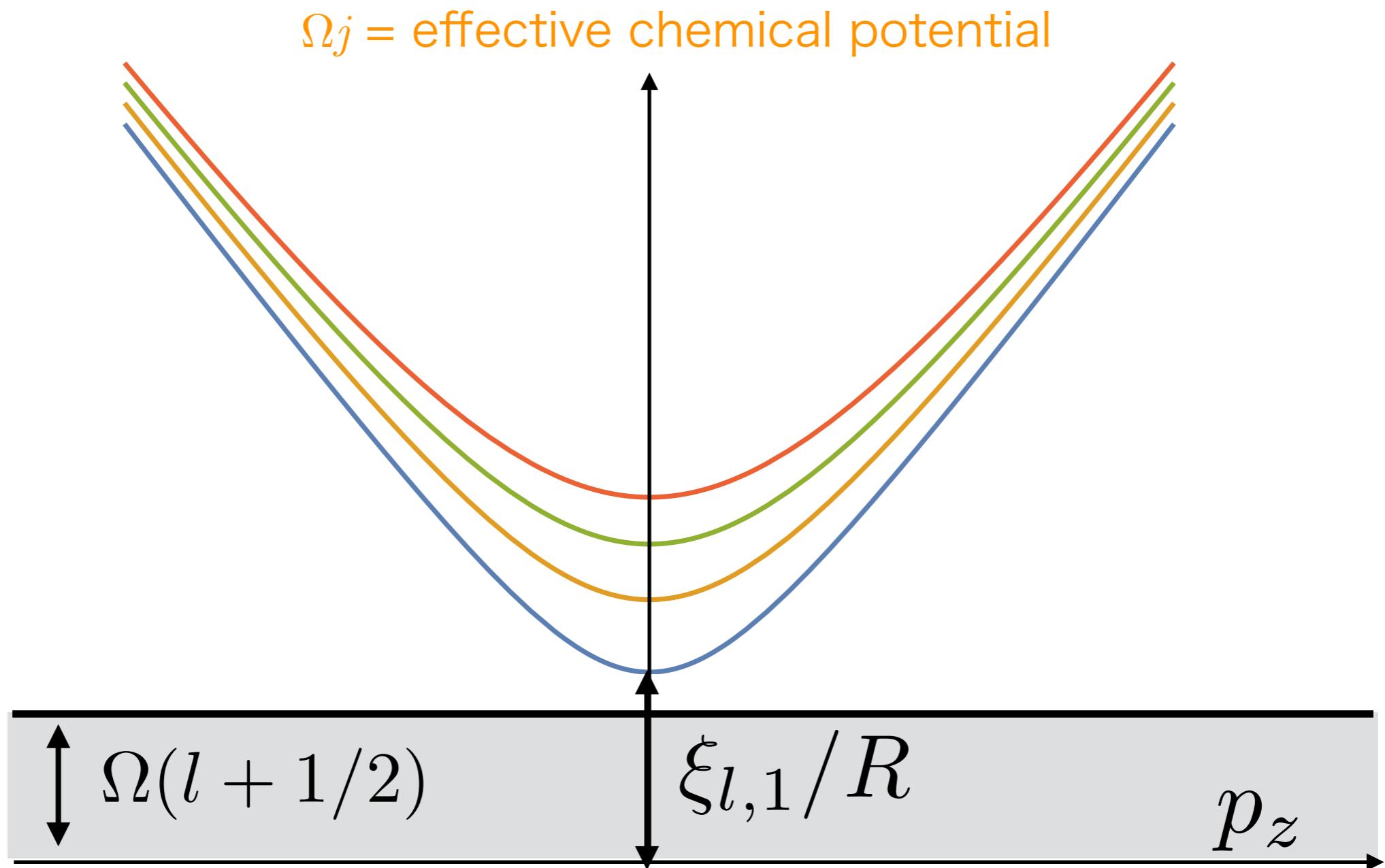
Rotational Effect at $T = 0$

$$f(\varepsilon, j) = \frac{1}{e^{\beta(\varepsilon - \Omega j)} + 1} \rightarrow f(\varepsilon, j) = \theta(\Omega j - \varepsilon)$$



Rotational Effect at $T = 0$

$$f(\varepsilon, j) = \frac{1}{e^{\beta(\varepsilon - \Omega j)} + 1} \rightarrow f(\varepsilon, j) = \theta(\Omega j - \varepsilon)$$

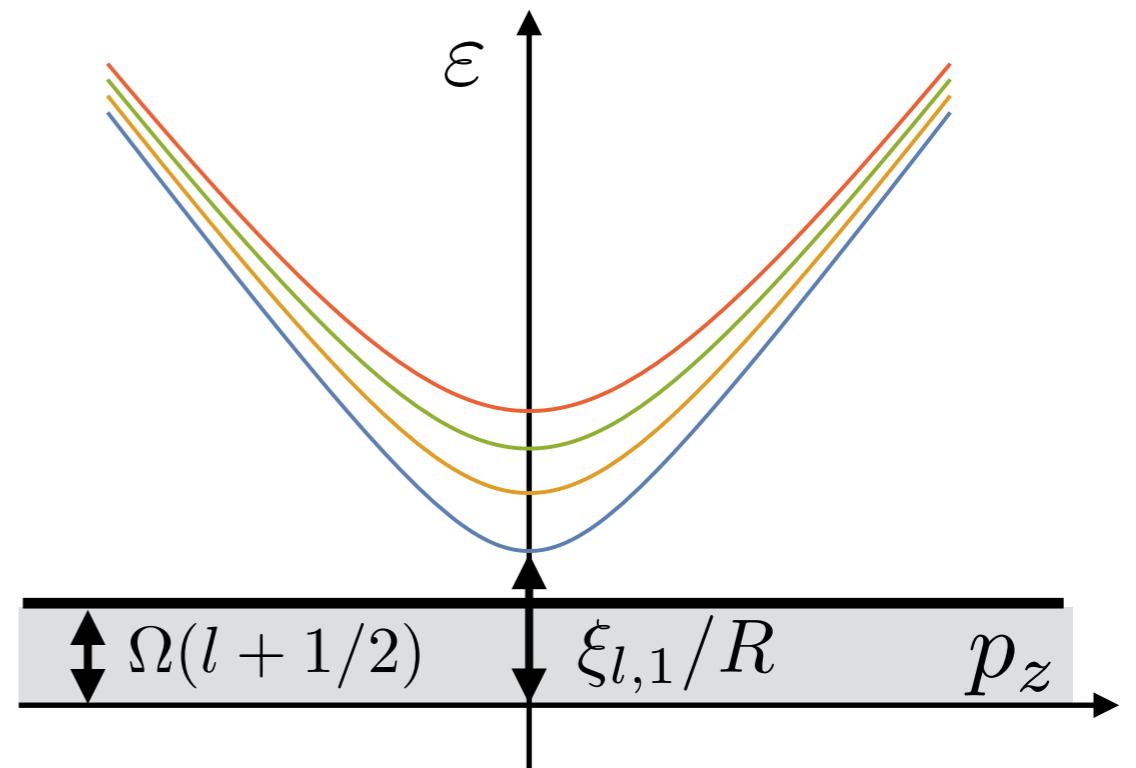


Which is true?

causality $\Omega \leq 1/R$

→ $\boxed{\xi_{l,1}/R > \Omega(l + 1/2)}$

for arbitrary l



finite-size effect



NO rotational effect at $T = 0$

Ebihara, Fukushima, KM (2016)

Note : visible at high T

CVE $j_5 = \frac{T^2}{12} \Omega$

Vilenkin (1979)

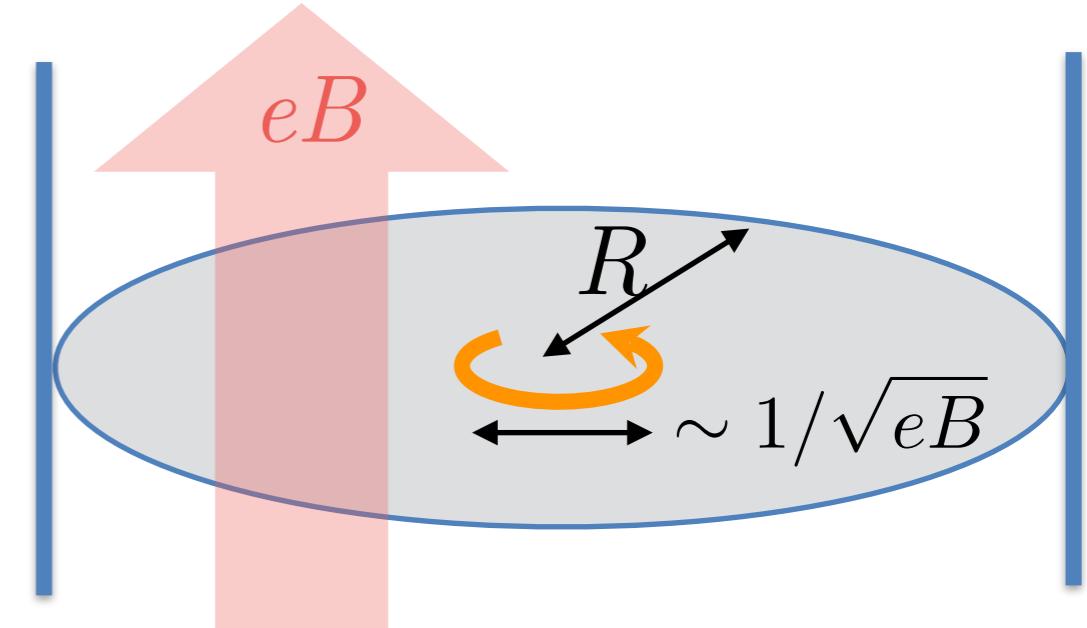
Part II Finite-size system with Ω and eB

Cyclotron Motion

$$(1) \quad 1/\sqrt{eB} \ll R$$

$$\rightarrow p_{\perp} = \sqrt{2neB}$$

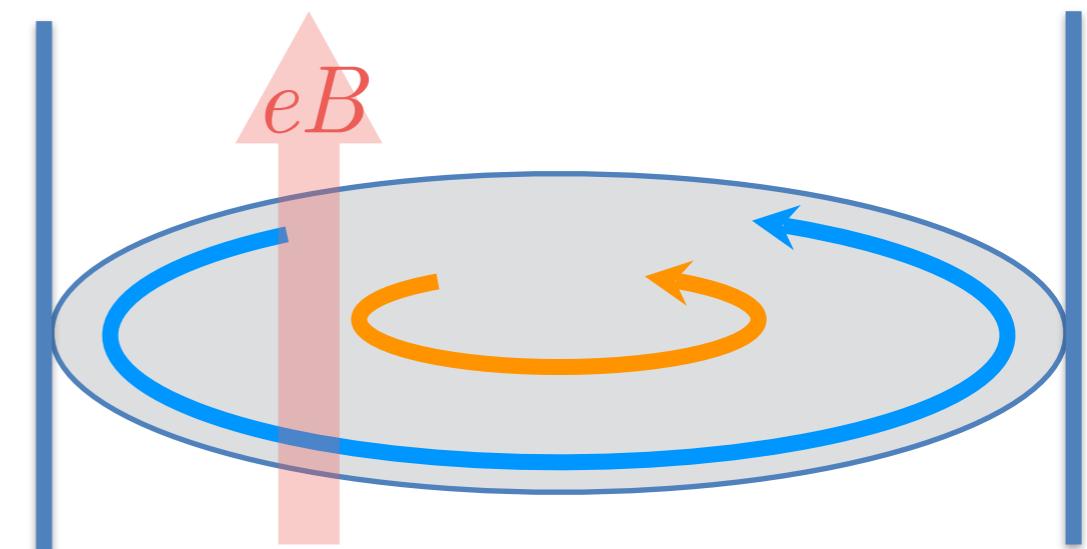
independent of l



$$(2) \quad 1/\sqrt{eB} \gtrsim R$$

small l \rightarrow still $p_{\perp} \simeq \sqrt{2neB}$

large l \rightarrow modified?

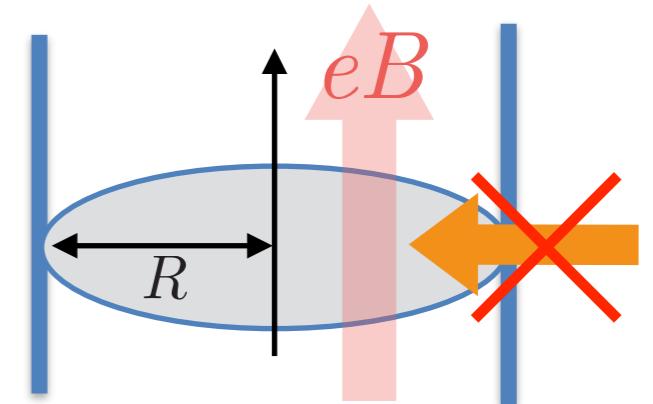


“Incomplete Landau quantization”

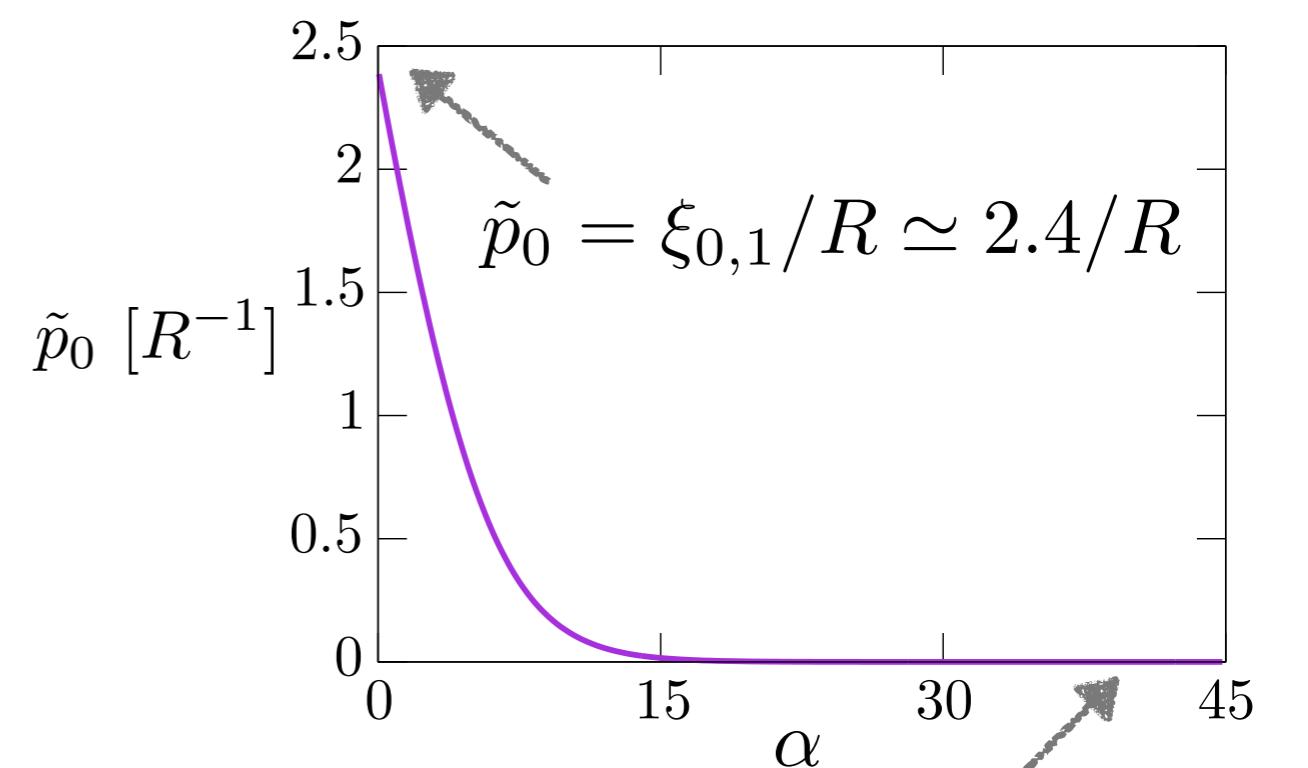
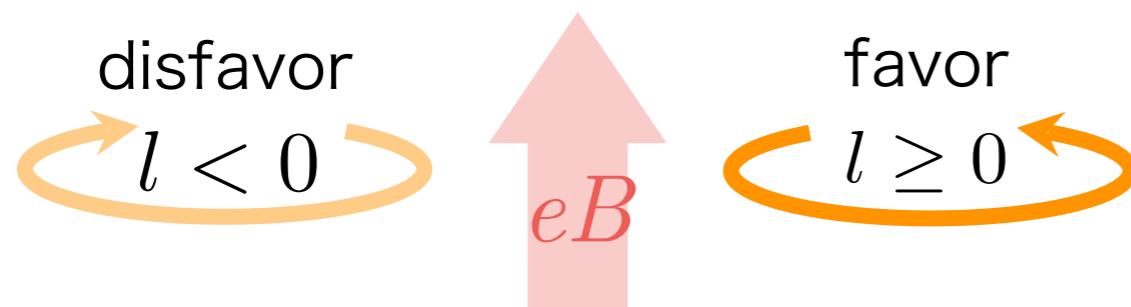
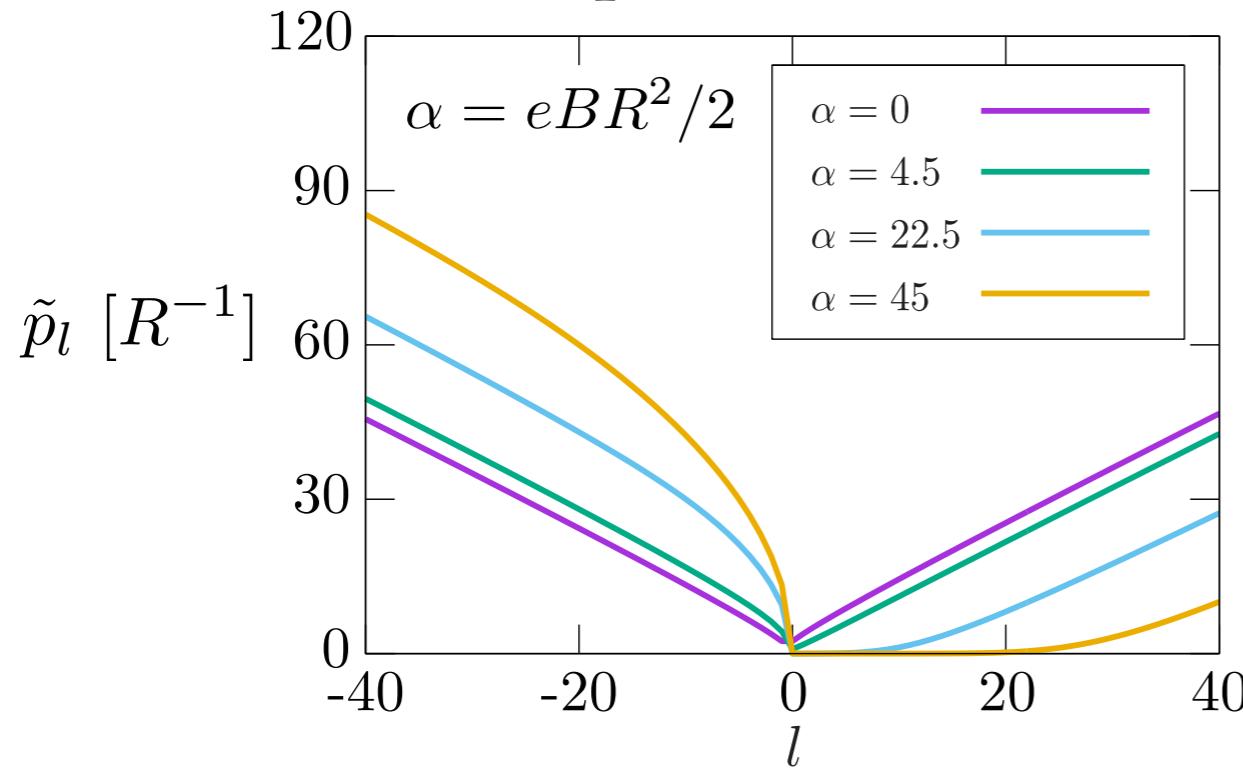
Incomplete Landau Level

Chen, Huang, Fukushima, KM (2017)

$$[i\gamma^\mu(\partial_\mu + ieA_\mu + \Gamma_\mu)]\psi = 0 \quad \text{with}$$

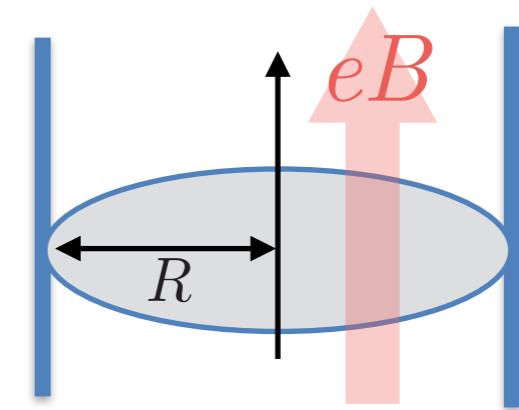
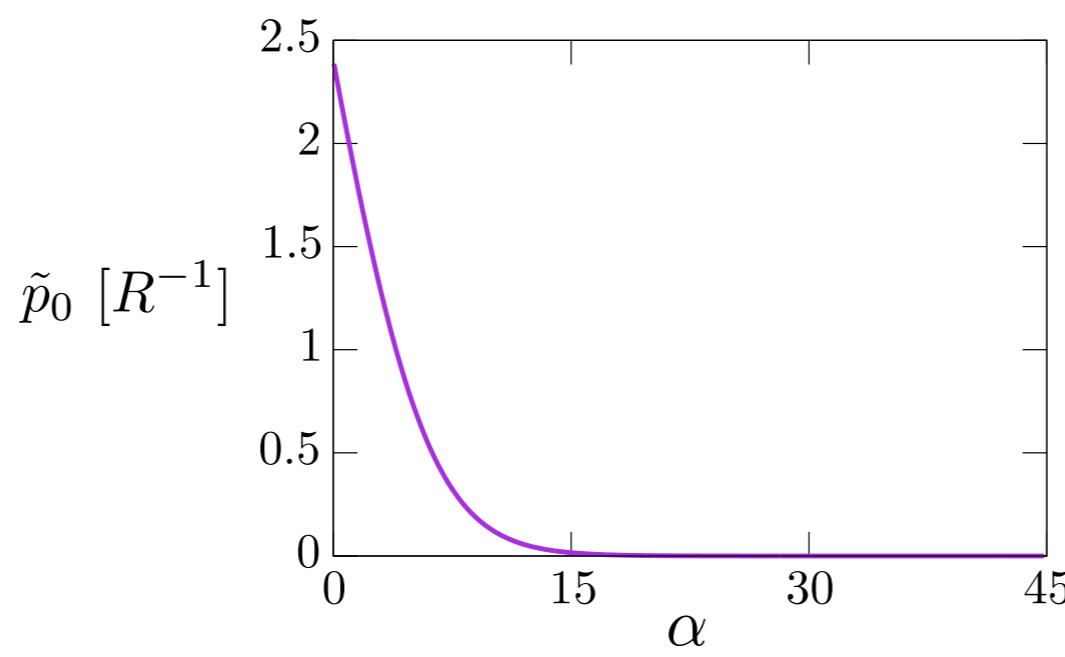


\tilde{p}_l = lowest transverse momentum for l



$\tilde{p}_0 = 0$
Landau zero mode

Gapped to Gapless

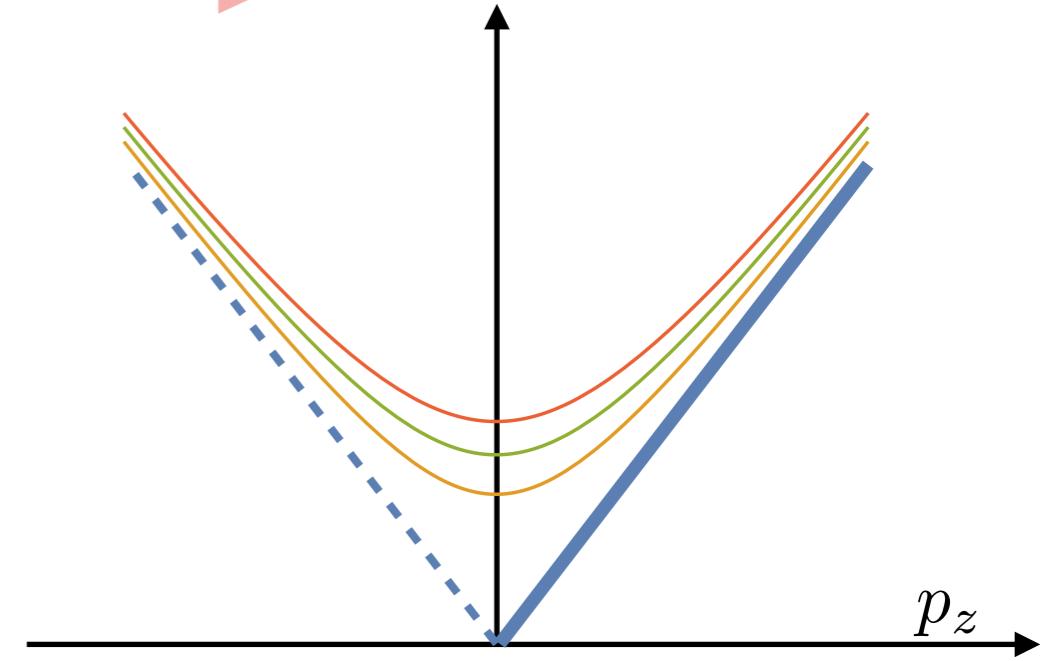
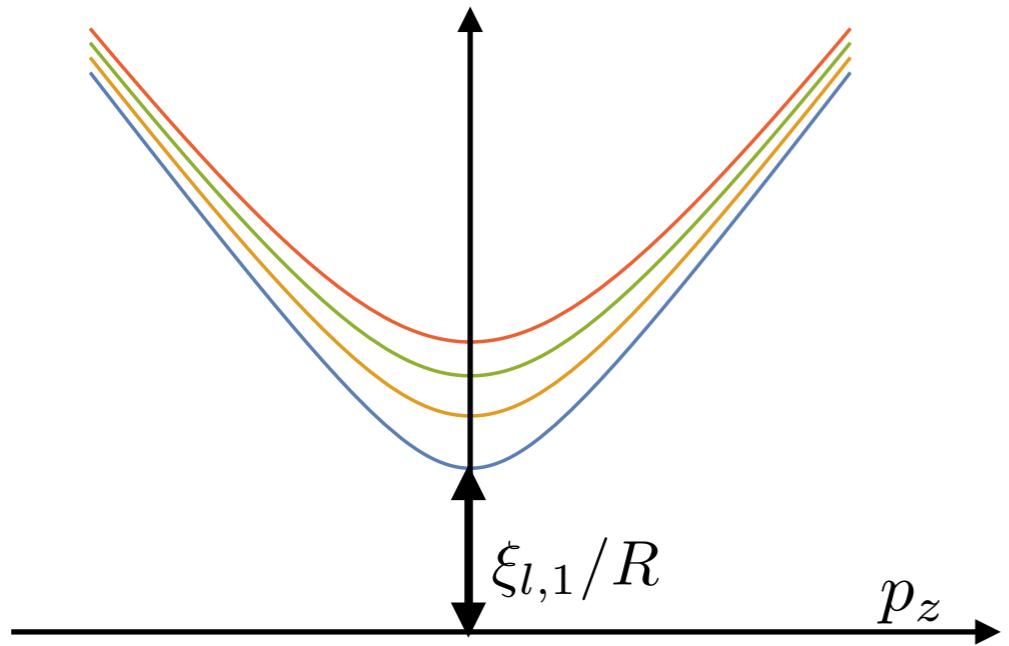


$$\alpha = eBR^2/2$$

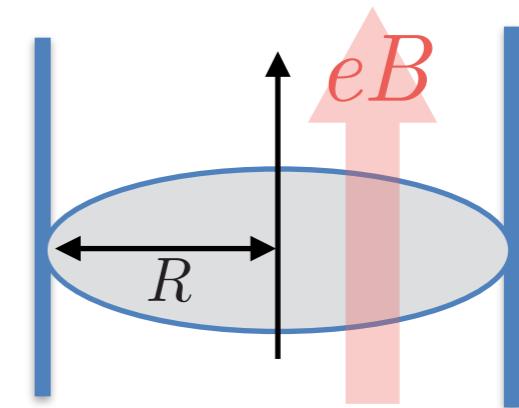
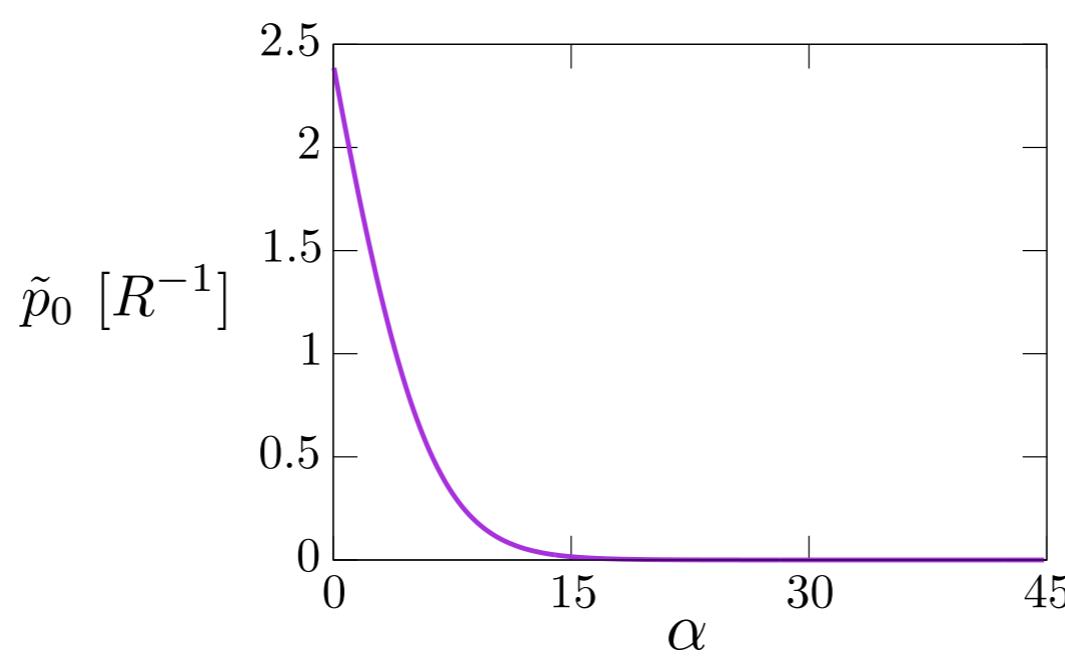
weak

magnetic field

strong



Gapped to Gapless

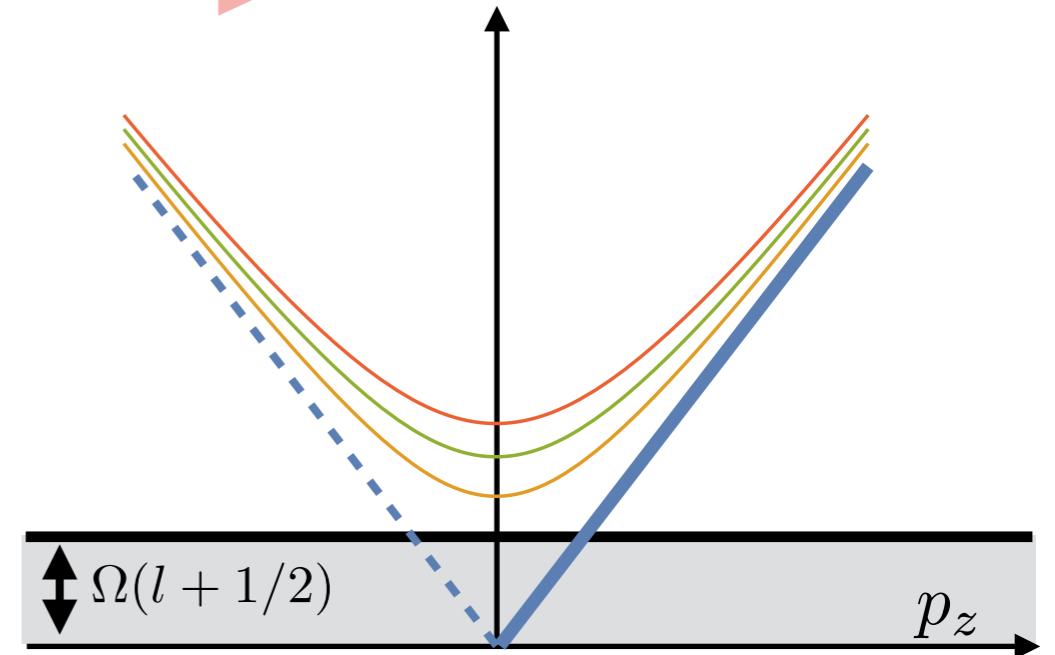
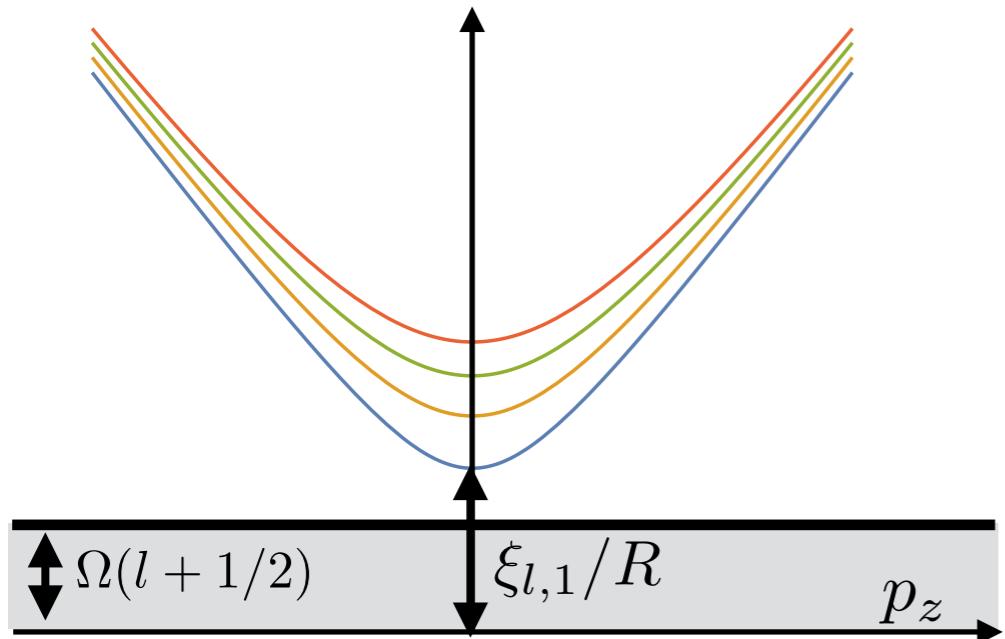


$$\alpha = eBR^2/2$$

weak

magnetic field

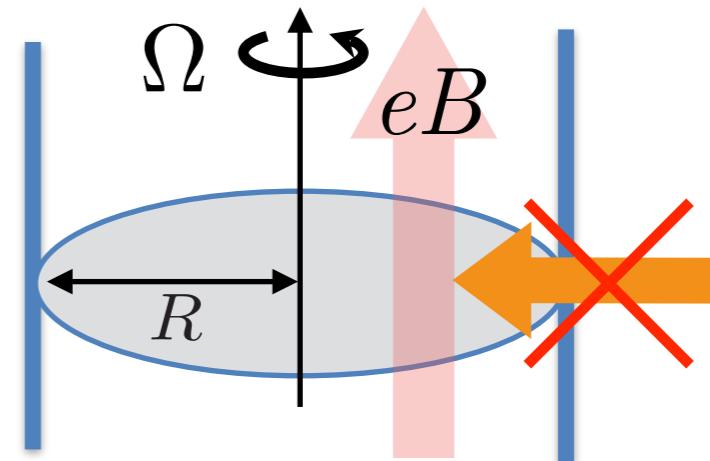
strong



visible rotational effect due to magnetic field

Ex.1) Density Induced by Rotation

$$[i\gamma^\mu(\partial_\mu + ieA_\mu + \Gamma_\mu)]\psi = 0 \quad \text{with}$$



$$\Omega j = \text{effective chemical potential} \quad f_{\pm}(\varepsilon) = \frac{1}{e^{\beta(\varepsilon \mp \Omega j)} + 1}$$

$$\begin{aligned} n(r) &= \langle \psi^\dagger(x)\psi(x) \rangle \\ &= \sum_{p_z, p_\perp} [f_+(\varepsilon) - f_-(\varepsilon)] \times (r\text{-dependence}) \quad \rightarrow \quad n(r=0) \xrightarrow{\sqrt{eB} \gg \Omega} \frac{eB\Omega}{4\pi^2} \\ &\quad \text{temperature independent} \end{aligned}$$

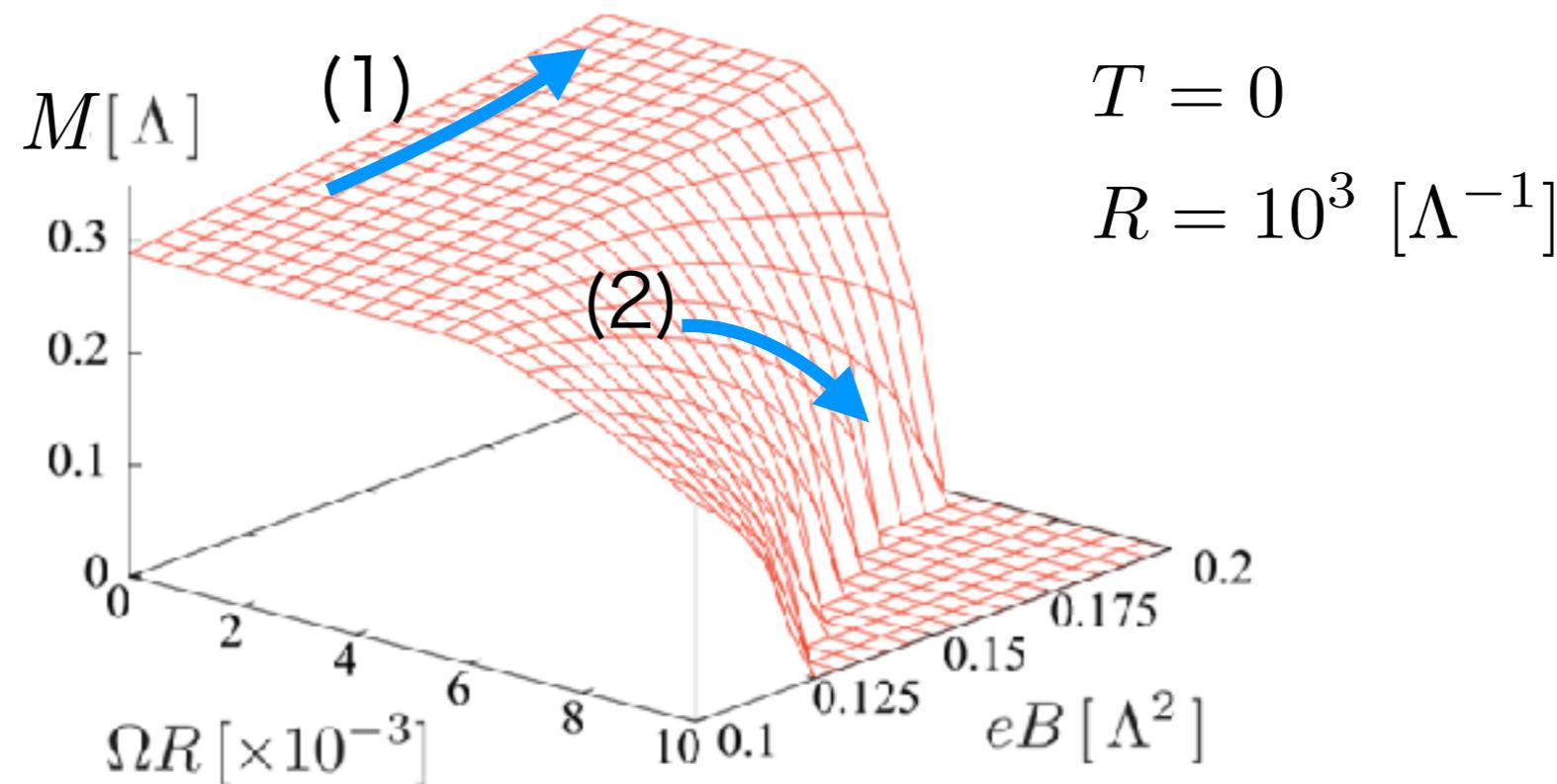
Hattori, Yin (2016)

Ebihara, Fukushima, KM (2017)

Ex.2) Chiral Symmetry Breaking

Chen, Huang, Fukushima, KM (2016)

NJL model (mean field approx.) + homogeneity



(1) eB increases $\longrightarrow M$ increases Magnetic Catalysis

(2) eB increases $\longrightarrow M$ decreases Inverse MC

Cf.) Ebert, Klimenko (1999) Preis, Rebhan, Schmitt (2012)

“rotational magnetic inhibition”

Summary & Outlook

Summary

- No rotational effect at $T = 0$
- B -field leads to a visible rotational effect (even at $T = 0$)
- Rotational magnetic inhibition : inverse phenomenon for the MC

Outlook

- Rotation yields abundant phase structures Chen, Huang, Fukushima, KM (2016)
Jiang, Liao (2016) Chernodub, Gongyo (2016)
Liu, Zahed (2017) Huang, Nishimura, Yamamoto (2017)
- Novel anomalous (magneto-vorticical) correction $T^{ij} = \# eB^i \Omega^j$
Hattori, Huang, KM (in preparation)
- Finite-size system under B -field EdH effect for chiral fermions
Fukushima, Hirono, Huang, Kharzeev, KM (in preparation)

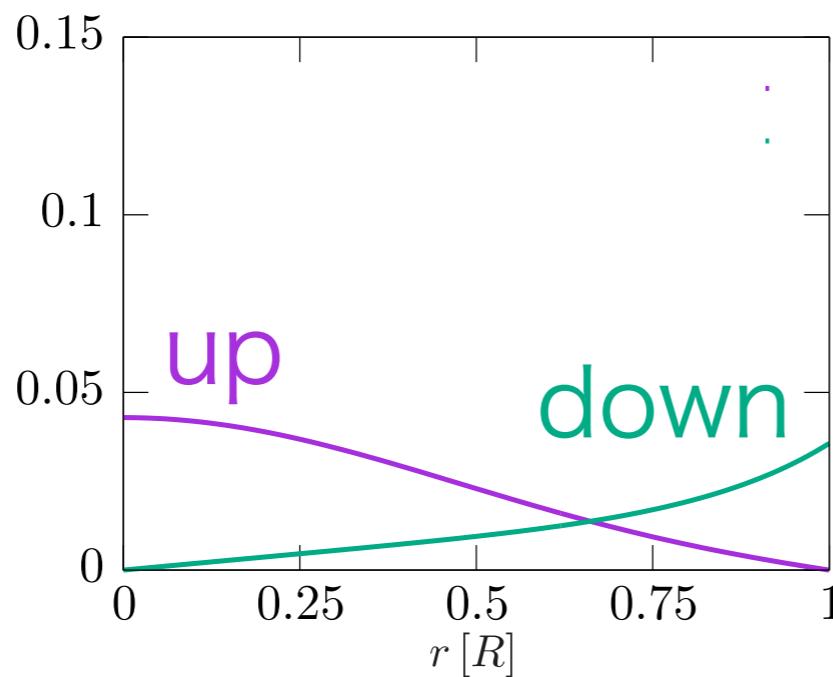
Wave Functions of Lowest Modes

$$\alpha = eBR^2/2$$

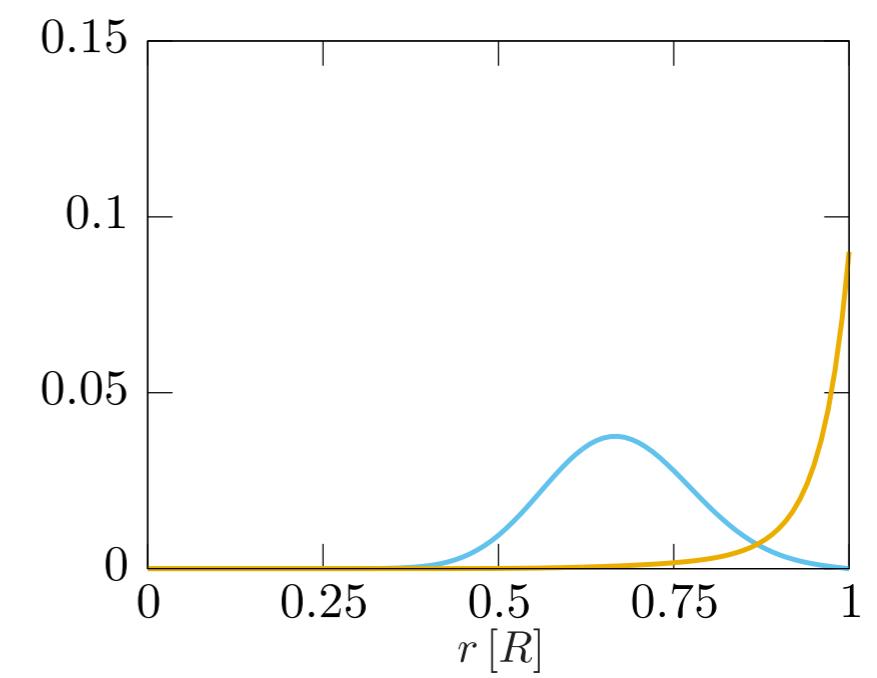
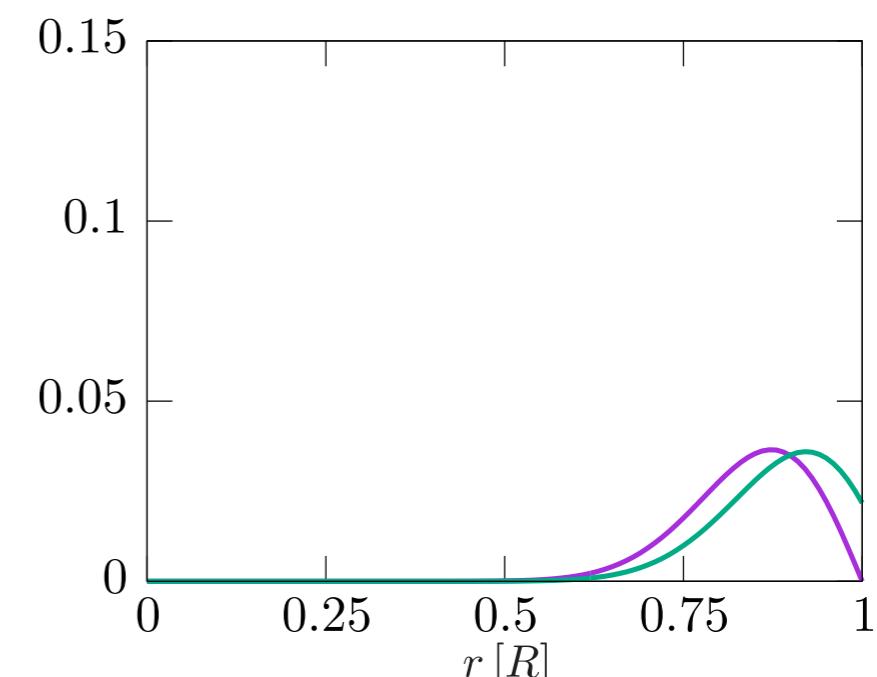
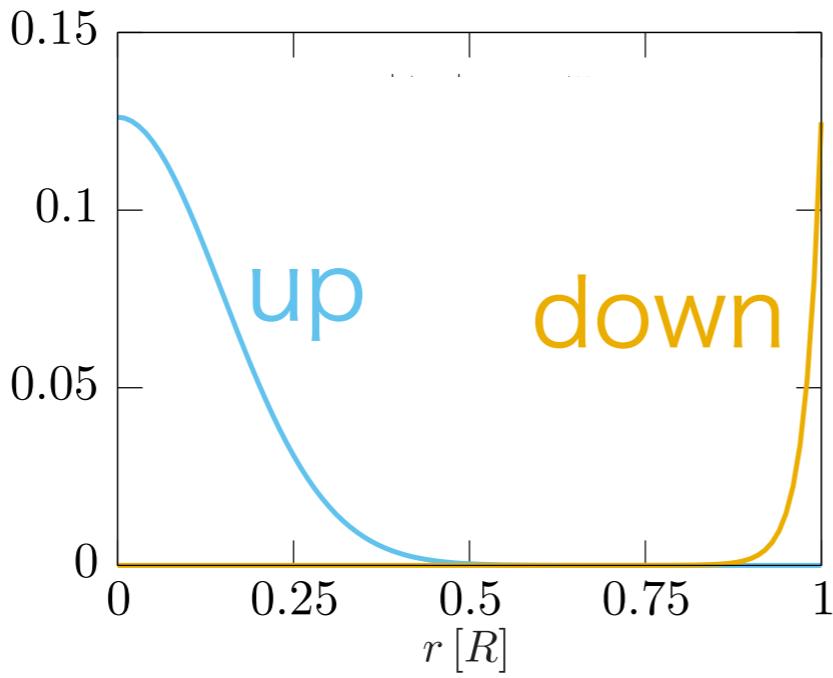
$$l = 0$$

$$l = 20$$

$$\alpha = 4.5$$



$$\alpha = 45$$



Magnetic Catalysis

eB increases

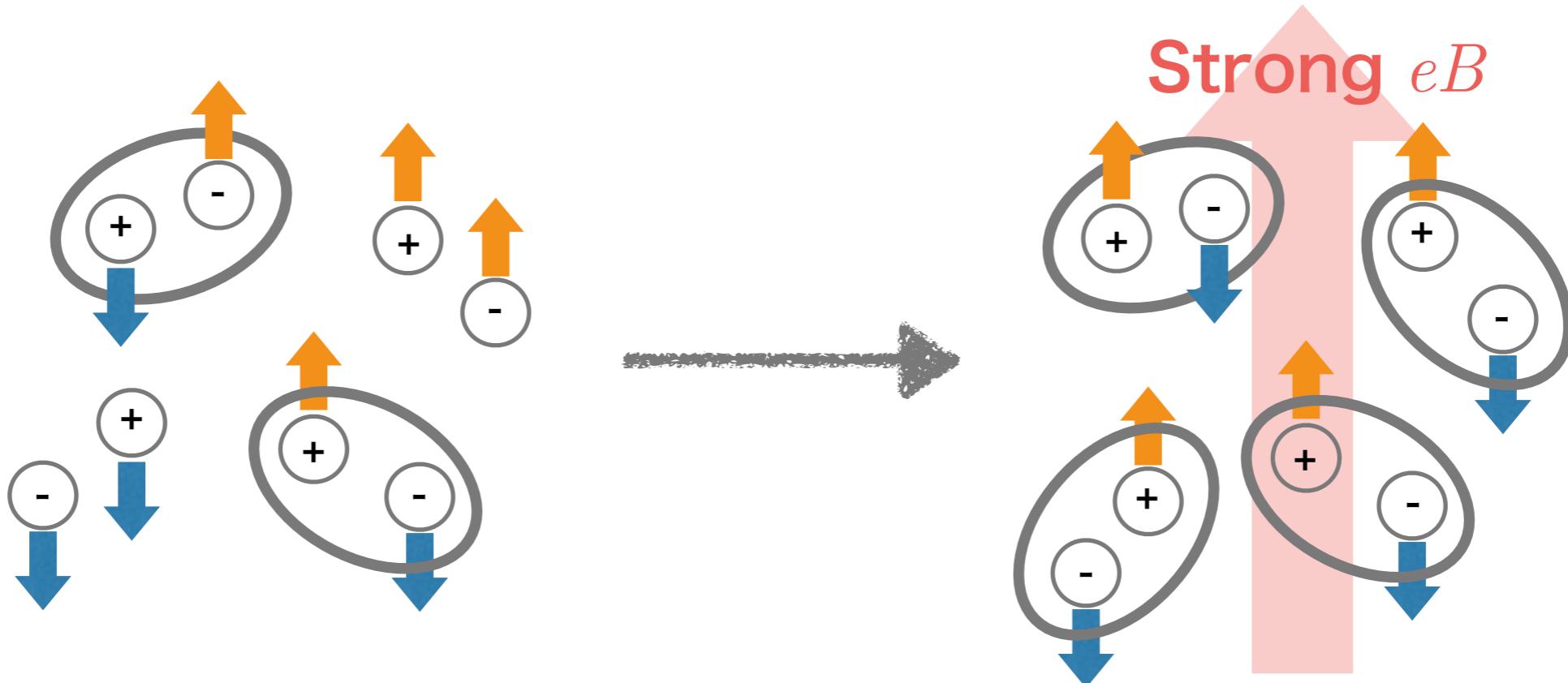
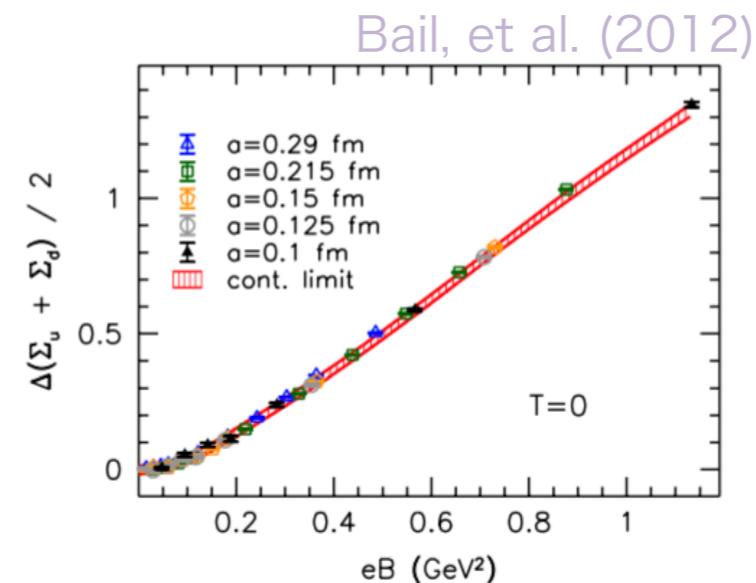


M increases

$$M \sim \Lambda e^{-\frac{2\pi^2}{GeB}}$$

Klimenko (1992)

Gyusynin, Miransky, Shovkovy (1994)



Inverse Magnetic Catalysis

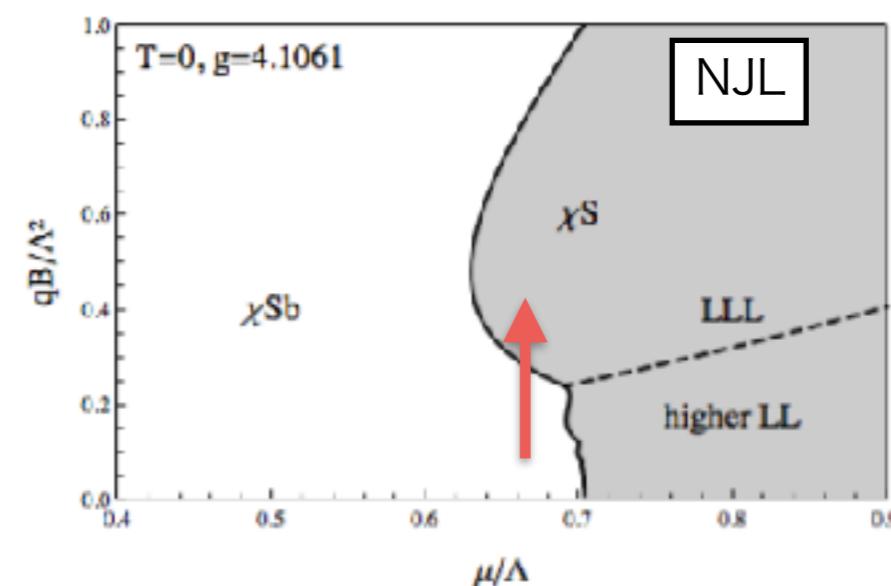
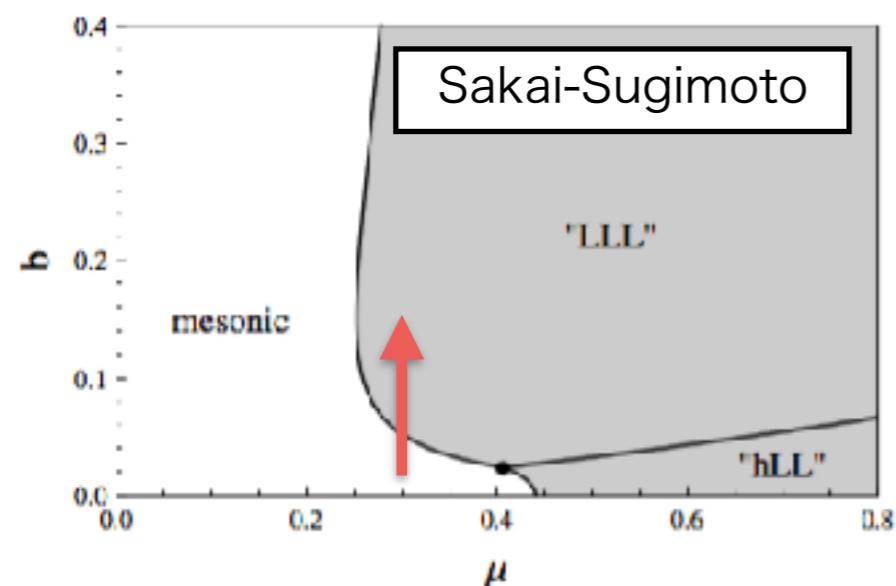
- magnetic field

eB increases $\longrightarrow m$ increases

- magnetic field + density ($\mu \neq 0$)

Ebert, Klimenko (1999)
Preis, Rebhan, Schmitt (2012)

eB increases $\longrightarrow m$ decreases



Preis, Rebhan, Schmitt (2012)