Anomalous Transport and Holography



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Outline

- Holography and currents
- Symmetry Breaking I & II
- More Symmetry Breaking (Poincare)
- Quenching
- Questions

Holography: $ds^2 = r^2(-f(r)dt^2 + d\bar{x}^2) + \frac{dr^2}{f(r)r^2}$



[Maldacena] [Witten] [Gubser, Klebanov, Polyakov]

Holography: $ds^2 = r^2(-f(r)dt^2 + d\bar{x}^2) + \frac{dr^2}{f(r)r^2}$



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AdS/CFT

$$\int_{\Phi|_{\partial}=\Phi_{0}} D\Phi e^{iS[\Phi]} = e^{iZ[\Phi_{0}]}$$
$$\frac{\delta^{n}Z[\Phi_{0}]}{\delta\Phi_{0}^{1}(x_{1})\cdots\delta\Phi_{0}^{n}(x_{n})} = \langle O_{1}(x_{1})\dots O_{n}(x_{n}) \rangle$$

Path integral (string theory) on AdS is hard. In practice resort to semi classical limit:

$$S_{grav}[\Phi_0] = Z[\Phi_0]$$

AdS/CFT

Dictionary

AdS	Field Theory
five dimensional	four dimensional
r-direction	RG flow
strongly coupled	weakly coupled
gravity	no gravity
metric	energy momentum tensor
gauge field	current
scalar field	scalar operator

AdS = spherical (hyperbolic) cow of sQGP



CME in Holography

• Anomalies = Chern-Simons terms

$$S = \int d^5x \sqrt{-g} \frac{1}{4} \left(F_R^2 + F_L^2 \right) + \frac{\alpha}{3} \left[A_R \wedge F_R \wedge F_R - A_L \wedge F_L \wedge F_L \right]$$

• Change to V-A basis

$$S = \int \frac{\alpha}{3} \left(A \wedge F_A \wedge F_A + A \wedge F_V \wedge F_V + V \wedge F_A \wedge F_V \right)$$

• Eliminate explicit V-dependence by adding boundary term

$$S_{ct} = c_1 \int_{\partial} V \wedge A \wedge F_V$$

Chose coefficient such that

$$S_{CS} = \int \frac{\alpha}{3} A \wedge [F_A \wedge F_A + 3F_V \wedge F_V]$$

[Rebhan, Stricker, Schmitt] [Gynther, K.L., Pena-Benitez, Rebhan]

 $V-A \neq R-L$

CME in Holography Currents

$$\mathcal{J}_{V}^{\mu} = \sqrt{-g} F_{V}^{\mu r} + 4\alpha \epsilon^{\mu\nu\rho\lambda} A_{\nu} F_{\rho\lambda}^{V}$$
$$\mathcal{J}_{A}^{\mu} = \sqrt{-g} F_{A}^{\mu r} + 4\frac{\alpha}{3} \epsilon^{\mu\nu\rho\lambda} A_{\nu} F_{\rho\lambda}^{A}$$

$$J_V^{\mu} = \sqrt{-g} F_V^{\mu r}$$
$$J_A^{\mu} = \sqrt{-g} F_A^{\mu r}$$

Anomalies

$$\partial_{\mu}\mathcal{J}_{V}^{\mu} = 0$$

$$\partial_{\mu}\mathcal{J}_{A}^{\mu} = -\frac{\alpha}{3}\epsilon^{\mu\nu\rho\lambda} \left(F_{\mu\nu}^{A}F_{\rho\lambda}^{A} + 3F_{\mu\nu}^{V}F_{\rho\lambda}^{V}\right)$$

- Variation of effective action
- Wess-Zumino consistency condition
- Not invariant under axial gauge trafo
- Not unique

$$\partial_{\mu}J_{V}^{\mu} = -2\alpha \,\epsilon^{\mu\nu\rho\lambda}F_{\mu\nu}^{A}F_{\rho\lambda}^{V}$$
$$\partial_{\mu}J_{A}^{\mu} = -\alpha \,\epsilon^{\mu\nu\rho\lambda} \left(F_{\mu\nu}^{A}F_{\rho\lambda}^{A} + F_{\mu\nu}^{V}F_{\rho\lambda}^{V}\right)$$

- Not Variation of effective action
- Not invariant under axial gauge taro
- Unique

AdS has rediscovered theory of *covariant* and *consistent* anomaly [Bardeen, Zumino]

CME in Holography CME and CSE

$$\vec{\mathcal{J}} = \frac{\mu_5 - A_0^5}{2\pi^2} \vec{B}$$
$$\vec{\mathcal{J}}_5 = \frac{\mu}{2\pi^2} \vec{B}$$
$$\vec{\mathcal{J}}_5 = \frac{\mu_5 - \frac{A_0^5}{3}}{2\pi^2} \vec{B}_5$$

• Equilibrium
$$H o H - \mu_5 Q_5 - \mu Q$$

 $A_0^5 = \mu_5$
 $V_0 = \mu$

• Bloch theorem [Yamamoto]

 A_0^5

Symmetry breaking - I CME and CSE

$$\vec{\mathcal{J}} = \frac{\mu_5 - A_0^5}{2\pi^2} \vec{B}$$
$$\vec{\mathcal{J}}_5 = \frac{\mu}{2\pi^2} \vec{B}$$
$$\vec{\mathcal{J}}_5 = \frac{\mu_5 - \frac{A_0^5}{3}}{2\pi^2} \vec{B}_5$$

- Holographic mass term
- Breaks axial symmetry at tree level
- Soft: only in IR
- Counterterms not affected
- Fate of CME et al.?
- In Holography charged scalar

$$S = \int d^5x \sqrt{-g} \left(|D_A \phi|^2 + m^2 |\phi|^2 \right) \quad , \quad \phi \approx \frac{M}{r} + \cdots$$

- Intuition: no chiral transport for large mass (strong breaking)
- BUT: covariant or consistent current ?

Symmetry breaking - I CME and CSE

CME =0 all the way (Bloch!)

[Jimenez-Alba, K.L., Liu, Sun]



Consistent currents vanish !

Same for spontaneous breaking[Amado, Lisker, Yarom]However lattice no change in CSE[Buividovich, Puhr]



• In Holography "linear axion" background (massless scalar)

$$S = \int d^5 x \sqrt{-g} \left(|\partial \phi|^2 \right) \quad , \quad \phi \approx kx + \cdots$$

- Background breaks translations eoms are homogeneous
- Graviton has mass

$$T^{0i} = T^{i0}$$

• Charge (Momentum) = Current (Energy-current)

I. Intuition: Momentum density, broken symmetry

2. Intuition: Energy current is dissipationless

• Include also gravitational anomaly

$$S = \int d^5 x A \wedge \operatorname{tr}(R \wedge R)$$

• Higher derivative term: extrinsic curvature component

$$\delta S = \int_{\partial} d^4 x \theta \operatorname{tr}[R^{(4)} \wedge R^{(4)}) + D(K \wedge DK)]$$

- Asymptotic AdS: extrinsic curvature component vanishes
- On BH horizon gives rise to T^2 terms in CVE etc..
- K is tensor in 4 dim sense

$$\mathcal{I} \to \mathcal{J} - K \wedge DK$$

• Can extrinsic curvature term be seen in UV?

$$ds^{2} = -fr^{2}dt^{2} + \frac{dr^{2}}{r^{2}f} + r^{2}d\vec{x}^{2}$$



Unusual power!

• Extrinsic curvature as additional variable

$$\delta S_{on-shell} = \int_{\partial} \sqrt{-g} (t^{\mu\nu} \delta g_{\mu\nu} + u^{\mu\nu} \delta K_{\mu\nu})$$

• Energy momentum tensor (Ward identity)

$$\Theta^{\mu\nu} = t^{\mu\nu} + u^{\mu\lambda} K^{\nu}_{\lambda}$$

• New term is due to gravitational Chern-Simons term

• CME and CVE without new term

$$T^{0i} = \left(8\alpha\mu^2 + 32\lambda T^2 - 4\lambda k^2\right)B^i$$
$$J^i = \left(8\alpha\mu^2 + 32\lambda T^2\right)2\omega^i$$

- Impossible in unitary theory
- CVE = 2 CME for energy current by Kubo formulas
- Including the new term

$$T^{0i} = \left(8\alpha\mu^2 + 32\lambda T^2\right)B^i$$

- All is well!
- Energy current being dissipation wins!

[Copetti, Fernandez-Pendas, K.L., Megias]

[Copetti, Fernandez-Pendas]

Quenching the CME

• Natural question: anomaly induced transport **far** from equilibrium physics

• Possible importance for Heavy Ion Collisions (magnetic field has already decayed in hydrodynamic regime)

 Holography allows both: study fast time evolution (quenches and anomalous transport

• Study CME via gravitational Chern-Simons term

What to look for

- "Minimal" setup: inject energy
- Equilibrium: energy temperature $T_0 \rightarrow T$
- CME in energy-momentum tensor
- First near equilibrium = hydro

$$T_{\mu\nu} = (\epsilon + p)u_{\mu}u_{\nu} + p\eta_{\mu\nu} + \hat{\sigma}_B(u_{\mu}B_{\nu} + u_{\nu}B_{\mu}),$$

$$J_{\mu} = \rho u_{\mu} + \sigma_B B_{\mu},$$

$$J_{\mu}^X = \rho_X u_{\mu} + \sigma_{B,X} B_{\mu}$$

• Landau frame

 $\hat{\sigma}_B = 0$ $\sigma_B = 24\alpha\mu - \frac{\rho}{\epsilon + p} \left(12\alpha\mu^2 + 32\lambda\pi^2 T^2\right)$ $\sigma_{B,X} = -\frac{\rho_X}{\epsilon + p} \left(12\alpha\mu^2 + 32\lambda\pi^2 T^2\right)$

What to look for

• Energy current = Momentum density

 $T_{0i} = T_{i0}$

• Momentum density = conserved charge

 $32\lambda\pi^2 T_0^2 \vec{B} = (\epsilon + p)\vec{v}$

• Monitor response in tracer U(I) current

$$\vec{J}_X = 32 \frac{\rho_X}{\epsilon + p} (T_0^2 - T^2) \pi^2 \lambda \vec{B}$$

• Removing constants: benchmark near equilibrium curve

$$j_X = \frac{T^2/T_0^2 - 1}{T^4/T_0^4}$$



Holographic quench

$$S = \frac{1}{2\kappa^2} \int d^5x \sqrt{-g} \left(\mathcal{R} - 2\Lambda - \frac{1}{4}F^2 - \frac{1}{4q^2}F_X^2 - \frac{1}{2}(\partial\phi)^2 + \lambda\epsilon^{MNOPQ}A_M R_{BNO}^A R_{APQ}^B \right)$$

- Tracer U(I) in decoupling limit
- Holographic quench $\phi_0(t, \vec{x}) = \frac{1}{2}\eta \left(1 + \tanh \frac{t}{\tau}\right)$

$$ds^{2} = \frac{1}{z^{2}} \left(-f(t,z)e^{-2\delta(t,z)}dt^{2} + \frac{dz^{2}}{f(t,z)} + d\vec{x}^{2} \right)$$

$$F_{xy} = B \qquad \qquad F_{X,0z} = \rho_X z \, e^{-\delta(t,z)}$$

• To linear order in B

$$J_X^\mu = \lim_{z \to 0} \sqrt{-g} \, F_X^{\mu z}$$

Holographic quench

• Fast quenches $2\tau T < 1$



Holographic quench

Very slow quenches



Summary

- Holography is efficient discovery tool for transport
- Fate of anomalous transport under symmetry breaking
- New questions:
 - Why do consistent currents vanish?
 - What is the extrinsic curvature in field theory?
- Anomalous transport far from equilibrium (QGP)