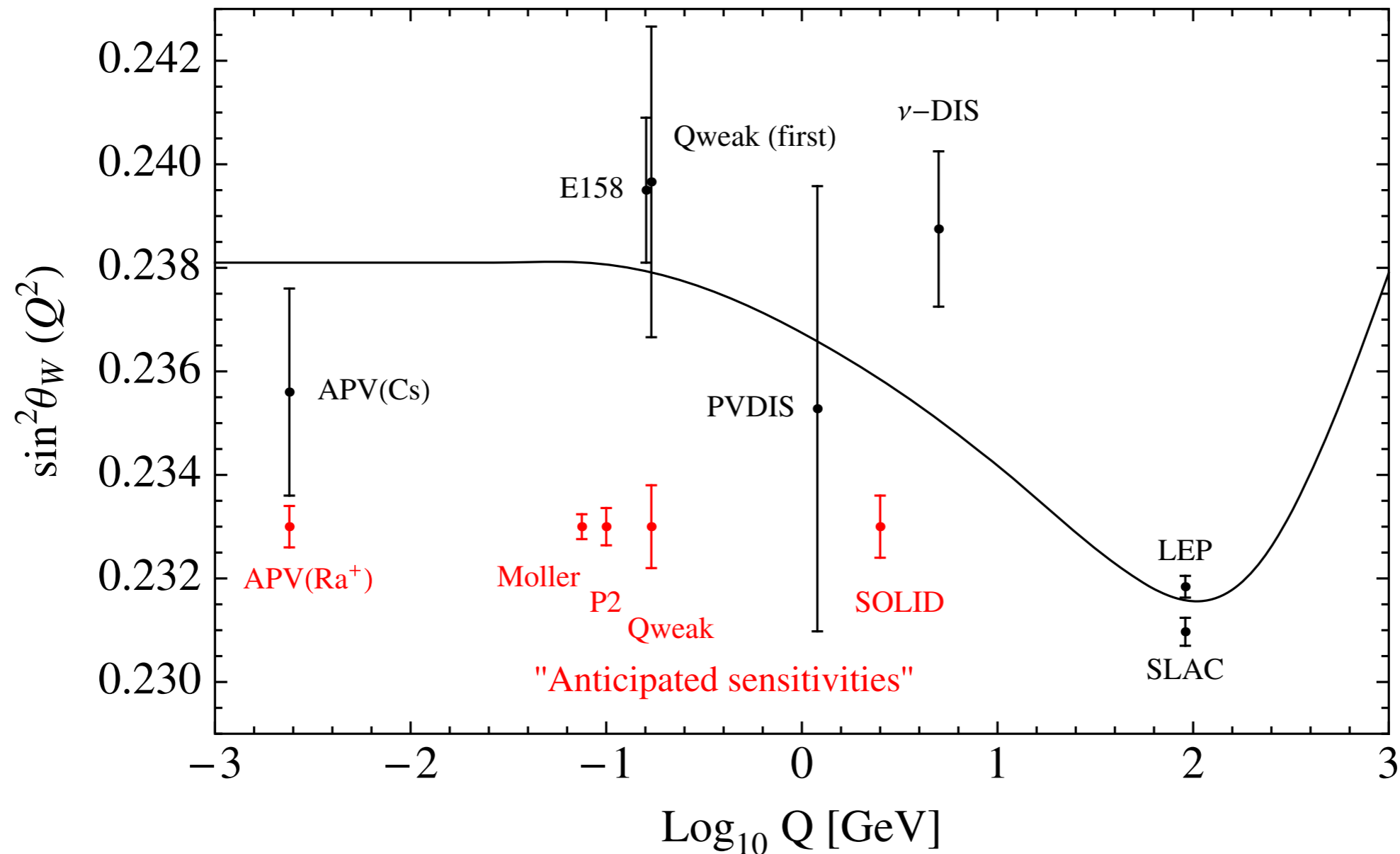

Atomic parity violation with radioactive ions
&
the muX experiment

Andreas Knecht, Paul Scherrer Institut

- ▶ Electric dipole moments: Fr-210, Ra-225, Rn
- ▶ Weak interaction studies: large range of isotopes
- ▶ Atomic parity violation: Fr, Ra

Why atomic parity violation?

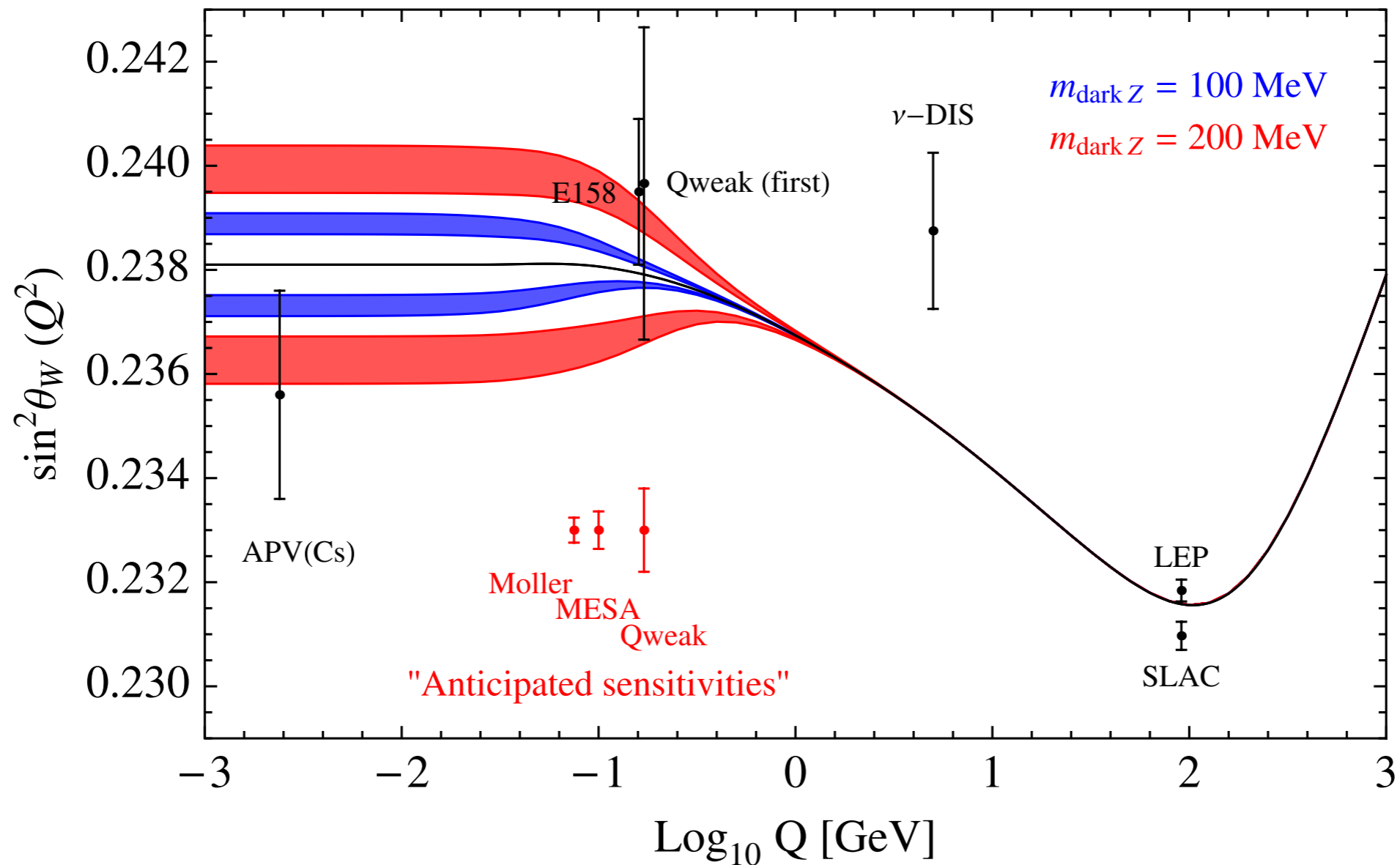
Davoudiasl, Lee, Marciano, Phys. Rev. D **92**, 055005 (2015)



- ▶ Running of the Weinberg angle as a function of momentum transfer
- ▶ APV fixes the low momentum value

Possible New Physics

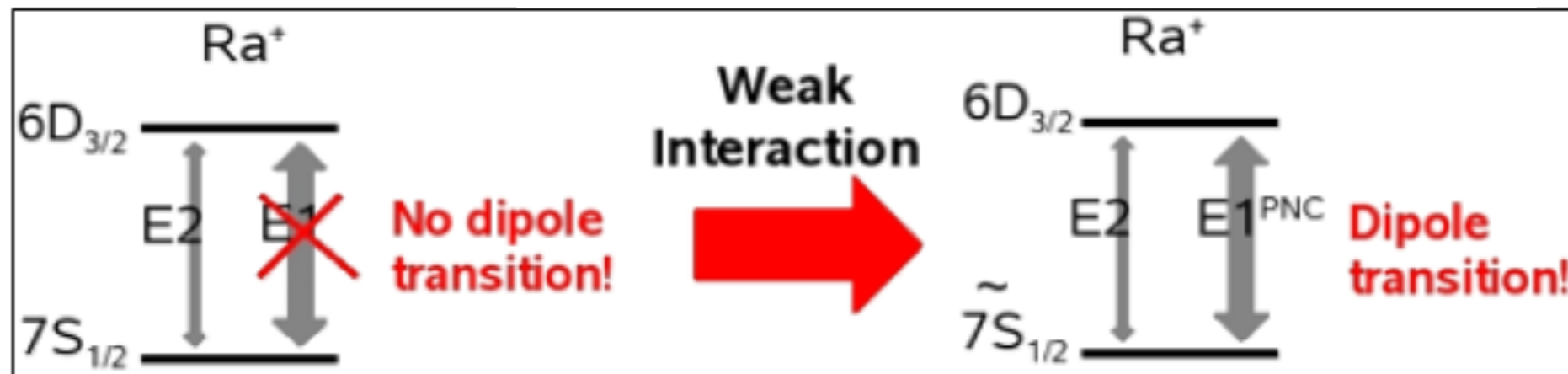
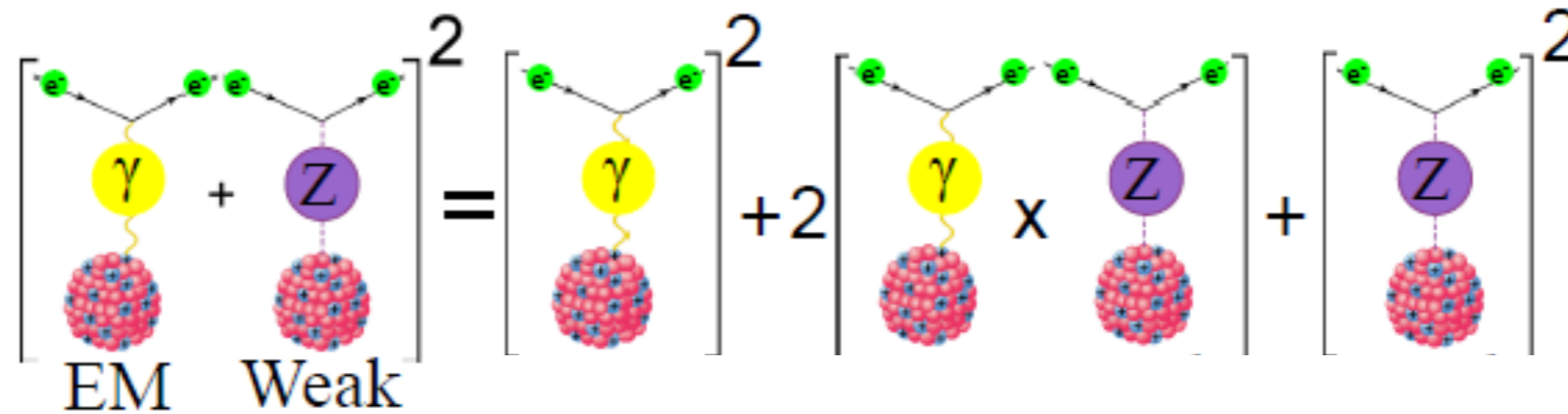
Davoudiasl, Lee, Marciano, Phys. Rev. D **89**, 095006 (2014)



- Possible new physics in the form of a new dark Z boson hides at low momentum!

Weak Interaction in Atoms

Interference of EM and Weak interactions



$$E1_{\text{PNC}} = K_r Z^3 \quad Q_w = K_r Z^3 (-N + Z(1 - 4\sin^2 \theta_w))$$

Measurement

Atomic Theory

Heavy System

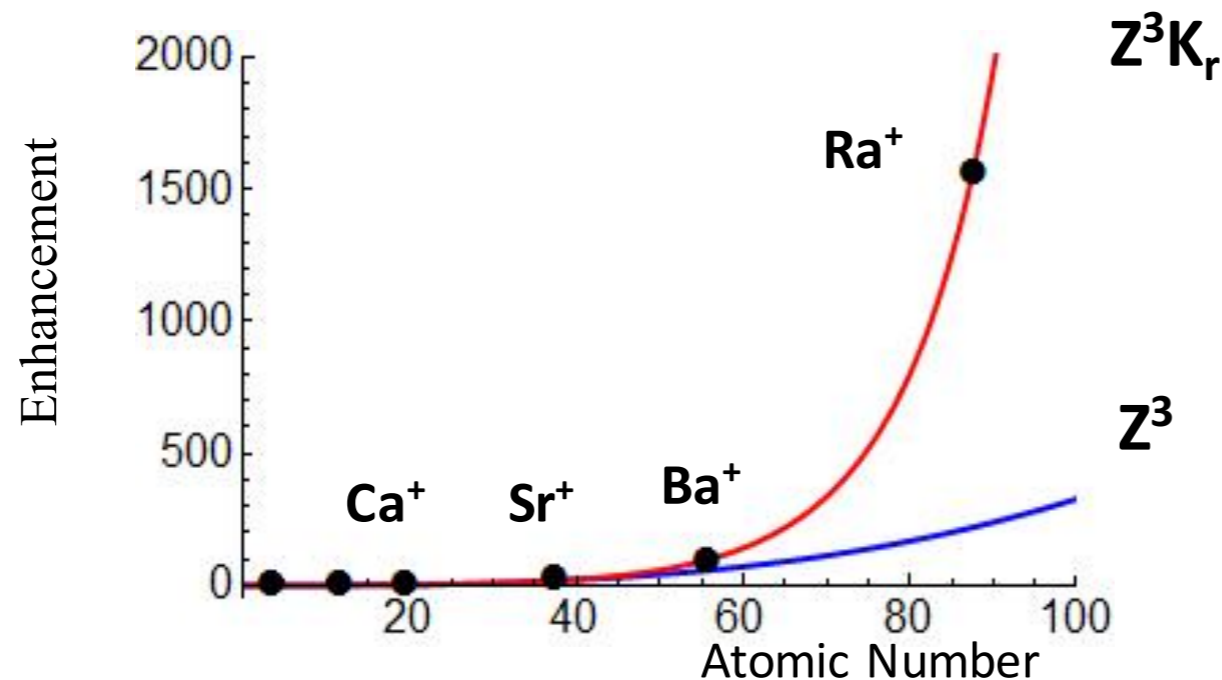
Benefit of Fr, Ra

Scaling of the APV

increase faster than Z^3

(Bouchiat & Bouchiat, 1974)

$$\langle nS_{1/2} | H_W | nP_{1/2} \rangle \propto K_r Z^3 \quad K_r \text{ relativistic enhancement factor}$$



Ra^+ effects

larger by:

20 (Ba^+)

50 (Cs)

L. W. Wansbeek *et al.*,
Phys. Rev. A **78**, 050501
(2008)

→ **5-fold improvement over Cs feasible in 1 day**

Relativistic coupled-cluster (CC) calculation of $E1_{\text{APV}}$ in Ra^+

$$E1_{\text{APV}} = 46.4(1.4) \cdot 10^{-11} \text{iea}_0 (-Q_w/N) \quad (3\% \text{ accuracy})$$

Other results:

$$45.9 \cdot 10^{-11} \text{iea}_0 (-Q_w/N) \quad (\text{R. Pal } et al., \text{ Phys. Rev. A } \mathbf{79}, 062505 (2009), \text{ Dzuba } et al., \text{ Phys. Rev. A } \mathbf{63}, 062101 (2001).)$$

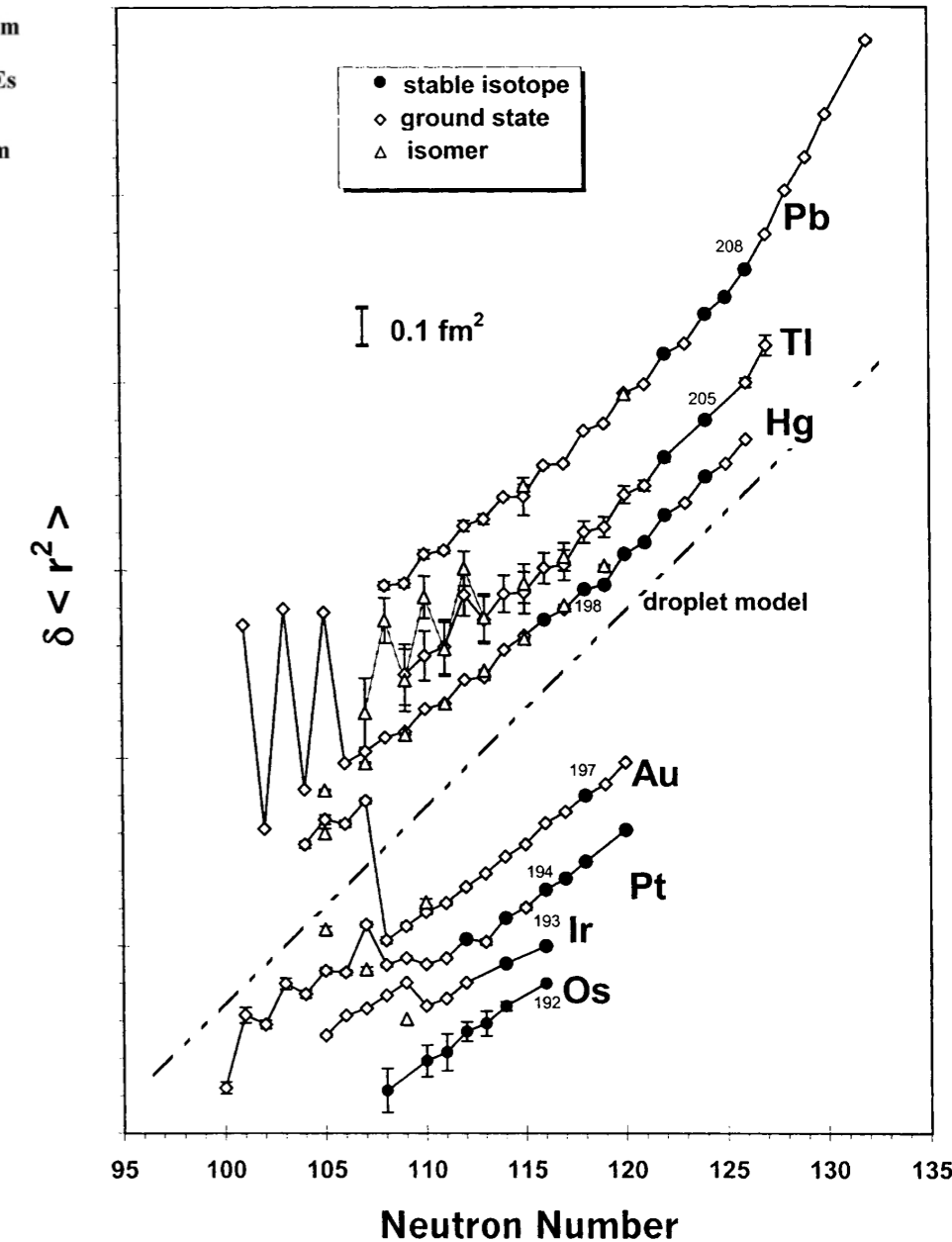
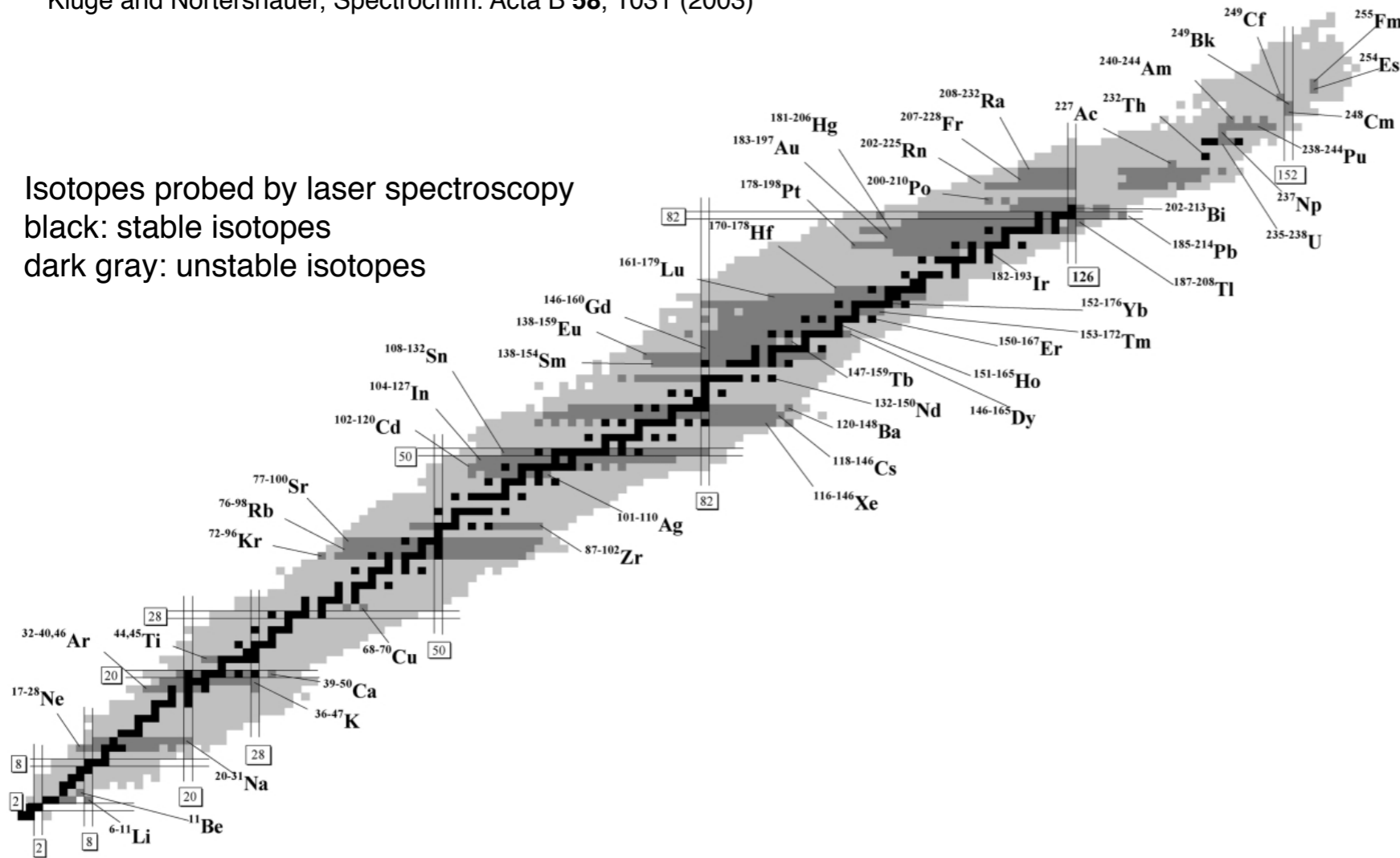
K. Jungmann, L. Willmann, Workshop
on Muonic Atom Spectroscopy (2016)

► Need reliable charge radius at $<0.2\%$ accuracy for atomic theory

Charge Radii from Laser Spectroscopy

Kluge and Nörtershäuer, Spectrochim. Acta B **58**, 1031 (2003)

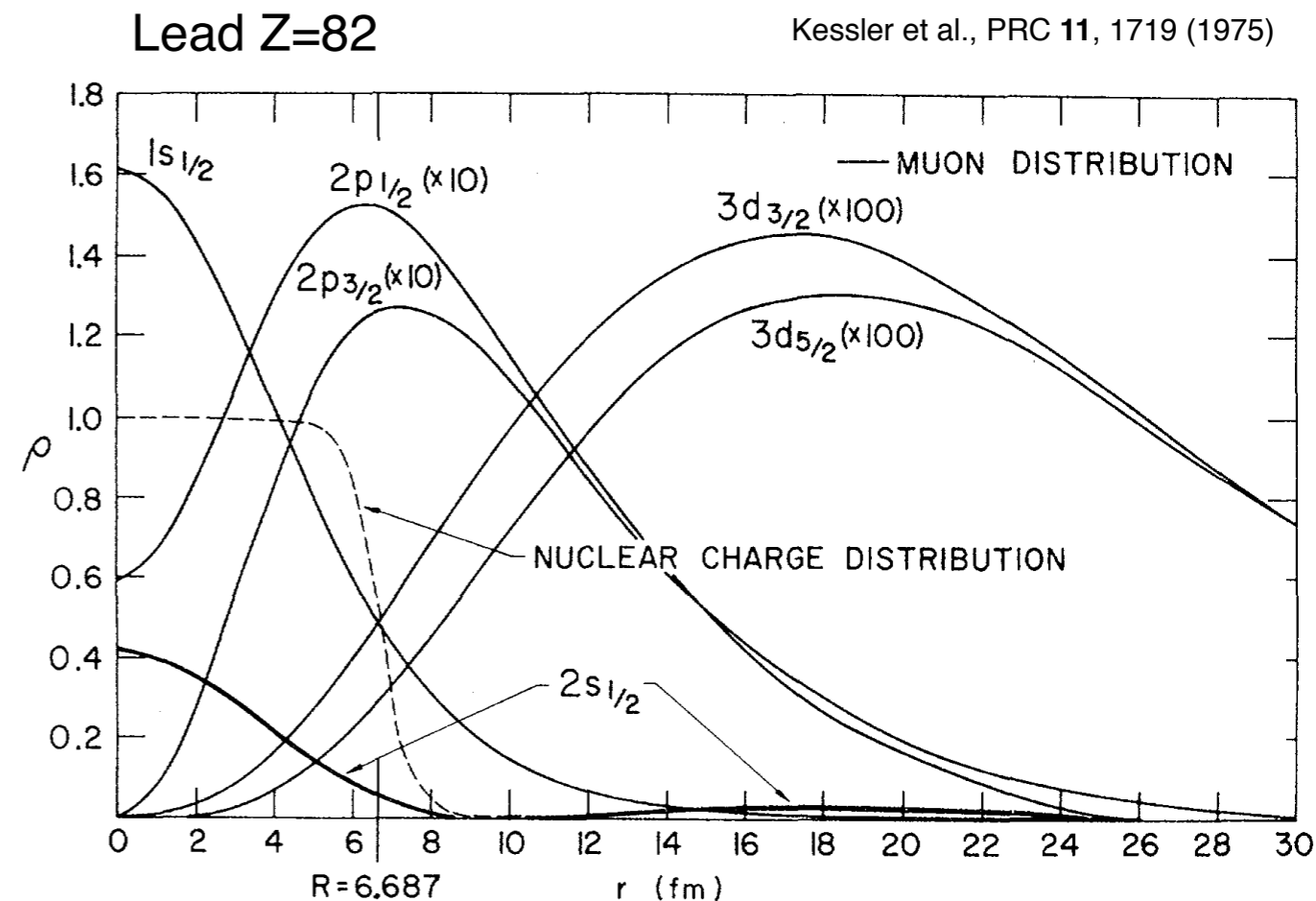
Isotopes probed by laser spectroscopy
 black: stable isotopes
 dark gray: unstable isotopes



- ▶ Wealth of information on nuclear properties from laser spectroscopy
- ▶ Need electron scattering or muonic atom spectroscopy for absolute radii

Muonic Atom Spectroscopy

- ▶ Muonic energy levels highly sensitive to nuclear charge distribution due to large overlap
- ▶ Using QED calculations and model for nuclear charge distribution allows to extract charge radius



Large effect:

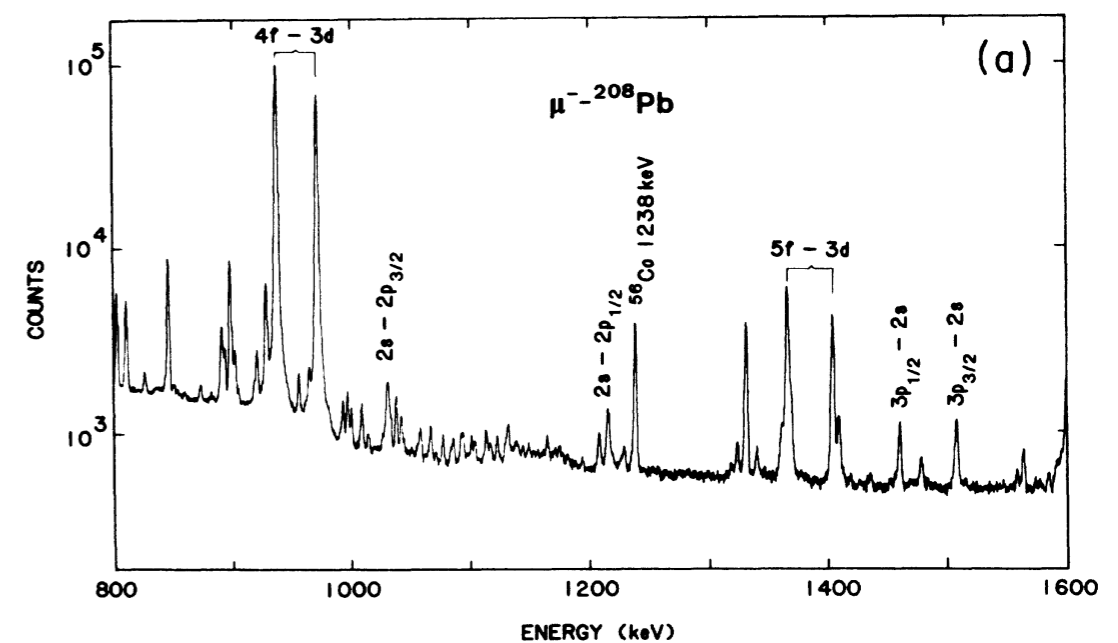
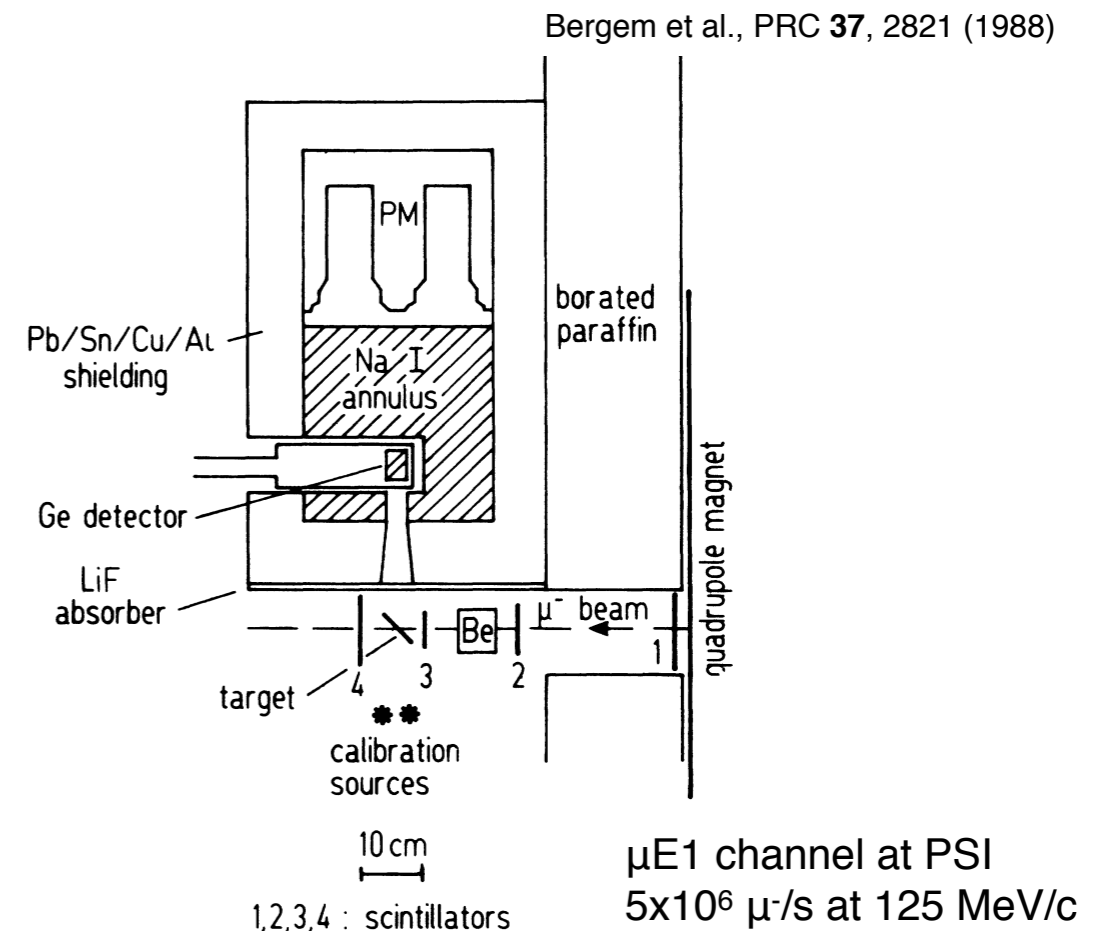
$E_{1s} (Z=82) \sim 19$ MeV (point nucleus)
 $\rightarrow 10.6$ MeV (finite size)

Muonic Atom Spectroscopy

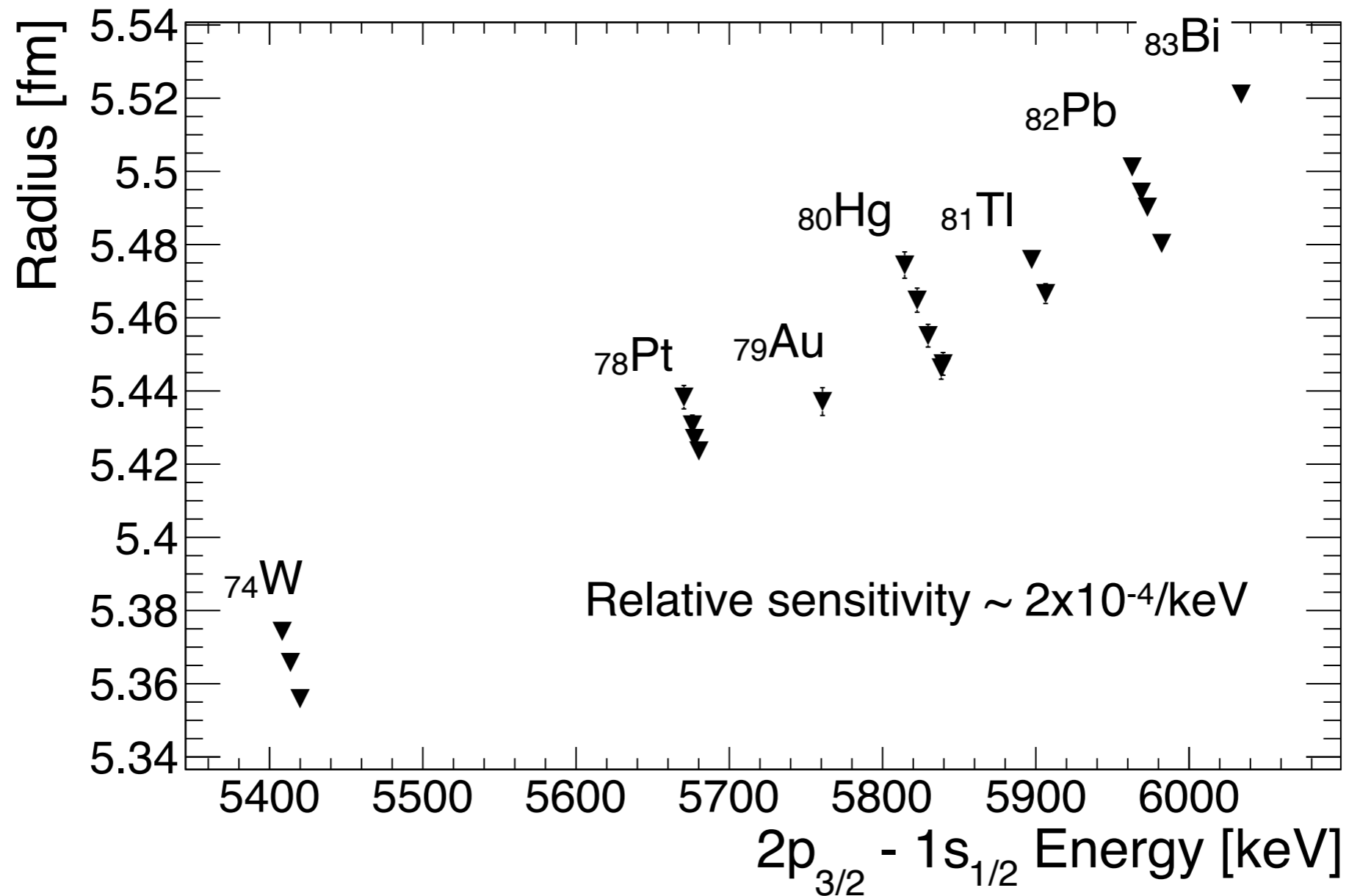
- ▶ Impressive precision in the extracted charge radius can be achieved
- ▶ For ^{208}Pb : $\langle r^2 \rangle^{1/2} = 5.5031(11)$ fm
 2×10^{-4} relative precision

TABLE V. Experimental muonic transition energies (keV) in ^{208}Pb (recoil corrected).

Transition	Kessler (Ref. 9)	Hoehn (Ref. 27)	This experiment
$2p_{3/2}-1s_{1/2}$	5 962.770(420)		5 962.854(90)
$2p_{1/2}-1s_{1/2}$	5 777.910(400)		5 778.058(100)
$3d_{3/2}-2p_{1/2}$	2 642.110(60)	2642.292(23)	2 642.332(30)
$3d_{5/2}-2p_{3/2}$	2 500.330(60)	2500.580(28)	2 500.590(30)
$3d_{3/2}-2p_{3/2}$	2 457.200(200)		2 457.569(70)
$3p_{3/2}-2s_{1/2}$	1 507.480(260)		1 507.754(50)
$3p_{1/2}-2s_{1/2}$			1 460.558(32)
$2s_{1/2}-2p_{1/2}$	1 215.430(260)		1 215.330(30)
$2s_{1/2}-2p_{3/2}$	1 030.440(170)		1 030.543(27)
$5f_{5/2}-3d_{3/2}$	1 404.740(80)		1 404.659(20)
$5f_{7/2}-3d_{5/2}$	1 366.520(80)		1 366.347(19)
$5f_{5/2}-3d_{5/2}$			1 361.748(250)
$4f_{5/2}-3d_{3/2}$	971.850(60)	971.971(16)	971.974(17)
$4f_{7/2}-3d_{5/2}$	937.980(60)	938.113(13)	938.096(18)
$4f_{5/2}-3d_{5/2}$			928.883(14)
$4d_{3/2}-3p_{1/2}$			920.959(28)
$4d_{5/2}-3p_{3/2}$			891.383(22)
$4d_{3/2}-3p_{3/2}$			873.761(63)



Muonic Atom Spectroscopy



- ▶ $2p - 1s$ energy is highly sensitive to charge radius
- ▶ What is the limiting factor?

Muonic Atom Spectroscopy

- ▶ Nuclear polarization is the dominating factor that in the end determines the accuracy of the extracted charge radius
- ▶ Typically assumed uncertainty: 10 - 30%
- ▶ Nuclear excitation spectra important
- ▶ Looking for theorists that want to tackle these calculations with modern methods

TABLE II. Theoretical nuclear polarization corrections in ^{208}Pb .

Energy (MeV)	I^π	$B(E\lambda)\uparrow$ ($e^2b^{2\lambda}$)	$1s_{1/2}$ (eV)	$2s_{1/2}$ (eV)	$2p_{1/2}$ (eV)	$2p_{3/2}$ (eV)	$3p_{1/2}$ (eV)	$3p_{3/2}$ (eV)	$3d_{3/2}$ (eV)	$3d_{5/2}$ (eV)
2.615	3^-	0.612	135	12	90	84	26	26	111	-63
4.085	2^+	0.318	198	20	182	180	76	84	6	4
4.324	4^+	0.155	14	1	8	7	2	2	1	1
4.842	1^-	0.001 56	7	1	-9	-8	0	0	1	1
5.240	3^-	0.130	27	2	16	15	5	5	2	2
5.293	1^-	0.002 04	9	2	-27	-19	0	-1	1	1
5.512	1^-	0.003 80	16	3	-90	-53	-1	-1	1	1
5.946	1^-	0.000 07	0	0	3	-30	0	0	0	0
6.193	2^+	0.050 5	29	3	22	21	7	7	0	0
6.262	1^-	0.000 24	1	0	3	5	0	0	0	0
6.312	1^-	0.000 22	1	0	3	4	0	0	0	0
6.363	1^-	0.000 14	1	0	2	2	0	0	0	0
6.721	1^-	0.000 75	3	1	6	7	0	-1	0	0
7.064	1^-	0.001 56	6	1	9	11	-1	-1	0	0
7.083	1^-	0.000 75	3	1	4	5	-1	-1	0	0
7.332	1^-	0.002 04	8	1	10	11	-2	-2	0	0
Total low-lying states			458	48	233	242	111	117	123	-53
13.5	0^+	0.047 872	906	315	64	38	24	15	1	0
22.8	0^+	0.043 658	546	147	43	26	15	10	0	0
13.7	1^-	0.537 672	1454	221	786	738	255	258	66	54
10.6	2^+	0.761 038	375	37	237	222	67	68	33	30
21.9	2^+	0.566 709	207	21	108	99	29	29	8	7
18.6	3^-	0.497 596	77	7	40	36	11	11	3	2
33.1	3^-	0.429 112	53	5	25	23	7	7	2	1
	$> 3^a$		176	15	80	71	21	21	4	4
Total high-lying states			3794	768	1383	1253	429	419	117	98
Total			4252	816	1616	1495	540	536	240	45

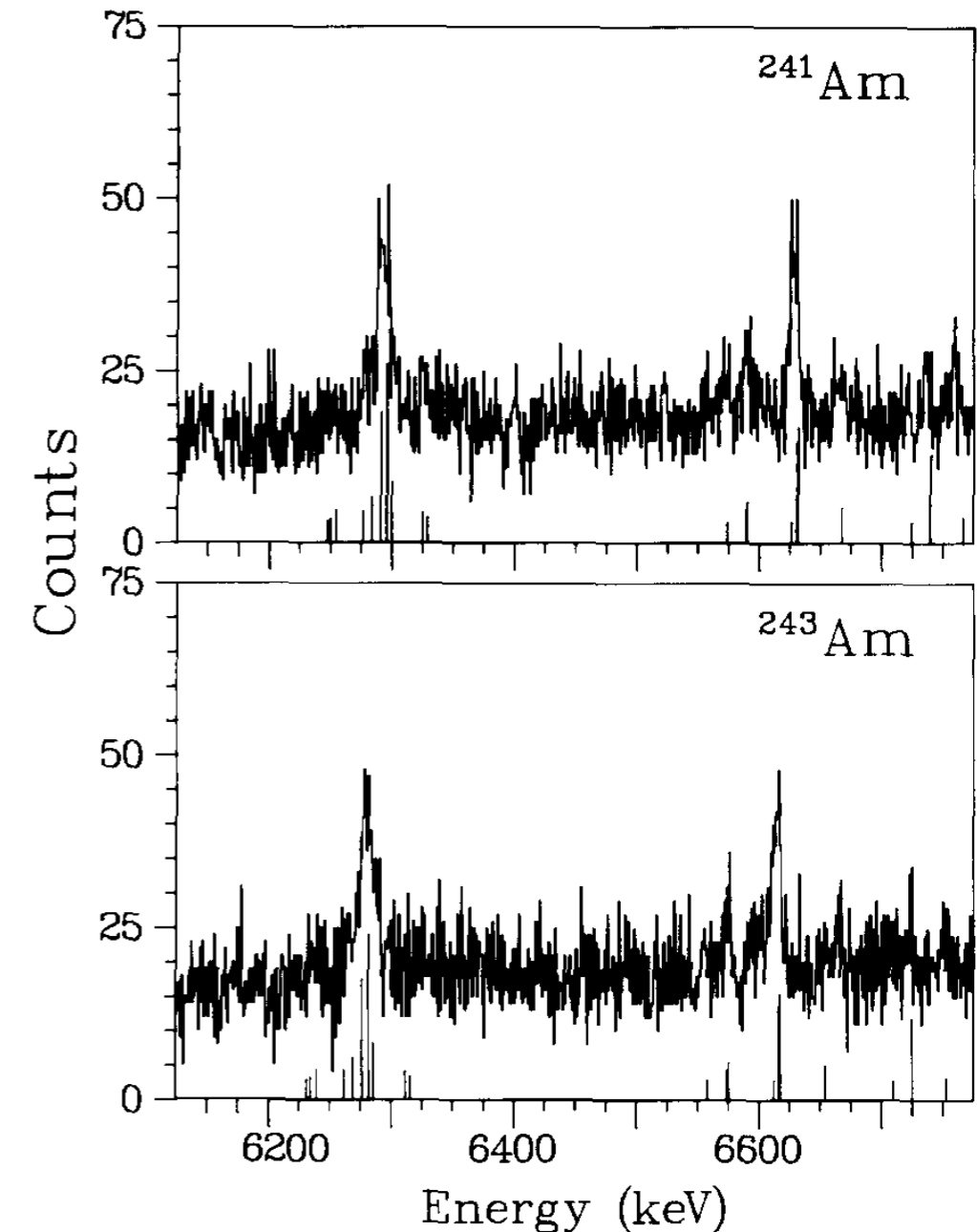
^aValues from Ref. 7. Positive NP values mean that the respective binding energies are increased.

Bergem et al., PRC **37**, 2821 (1988)

What About Radioactive Atoms?

- ▶ Most of the stable isotopes have been measured with muonic atom spectroscopy
- ▶ In a few special cases also radioactive isotopes, e.g. americium
 - ▶ The paper describes the americium target as “modest weight of 1 gram”

Johnson et al., Phys. Lett. **161B**, 75 (1985)

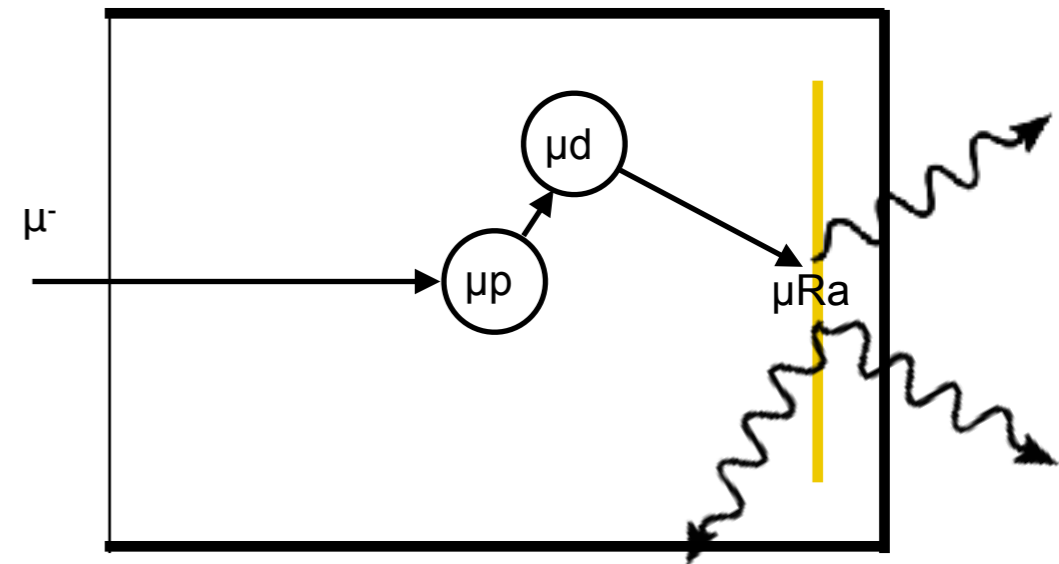


Measurements in ultra-thin targets

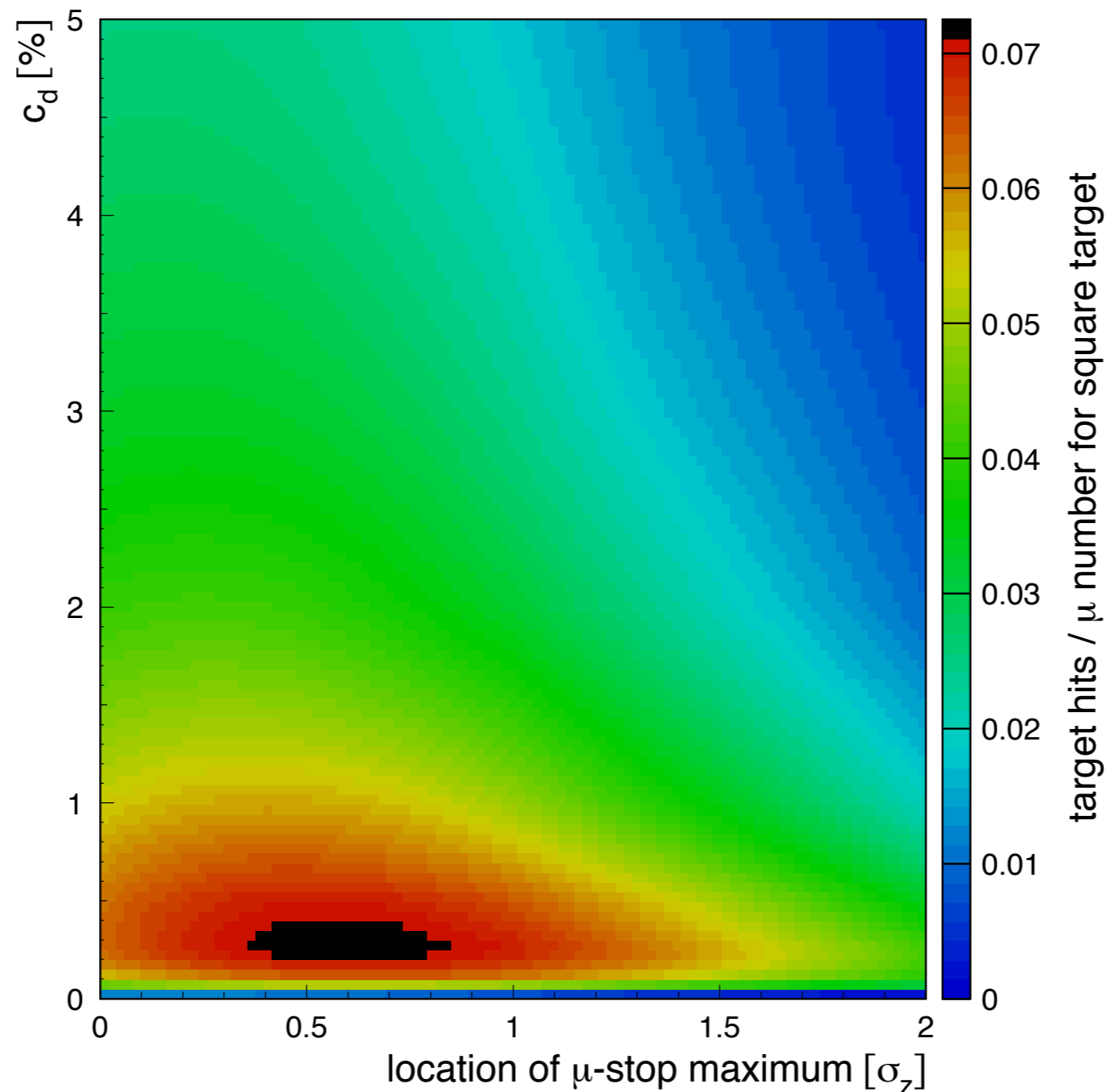
- ▶ Radioprotection laws more strict nowadays
 - ▶ Can only use 0.2 μg of ^{241}Am in PSI experimental hall
 - ▶ Can use 5 μg of ^{226}Ra
- ▶ For “normal” stopping need $O(100 \text{ mg})$
- ▶ Use the “magic” of muonic hydrogen/deuterium atoms and transfer reactions!

Transfer Reactions

- ▶ Stop in 100 bar hydrogen target with 0.25% deuterium admixture
- ▶ Form muonic hydrogen μp
- ▶ Transfer to deuterium forming μd , gain binding energy of 45 eV
- ▶ Hydrogen gas quasi transparent for μd at ~ 5 eV (Ramsauer-Townsend effect)
- ▶ μd reaches target and transfers to μRa
- ▶ Measure emitted X-rays from cascade



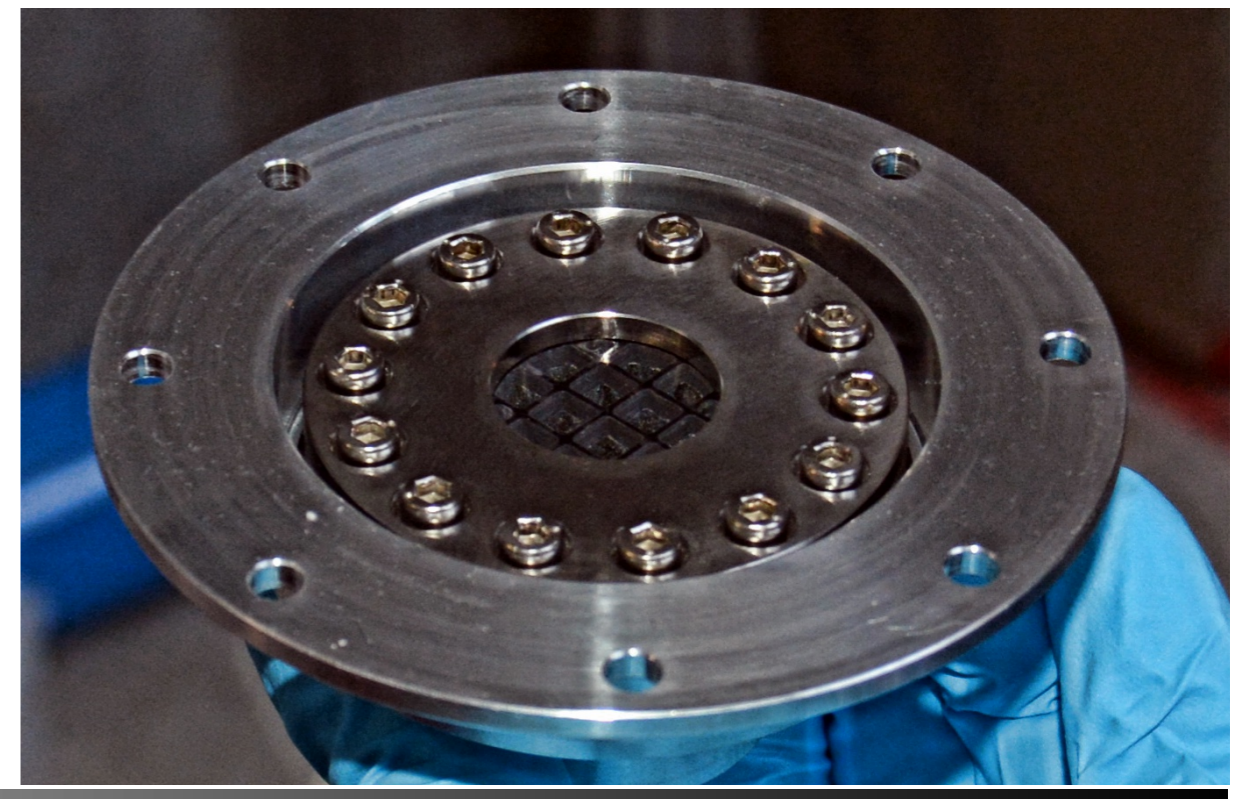
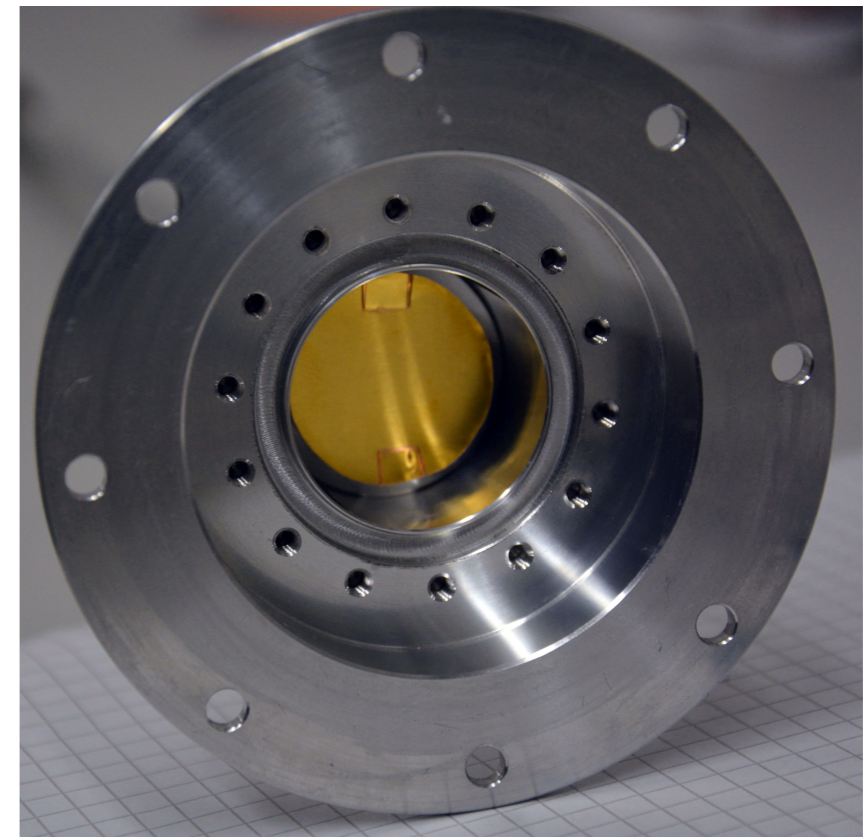
Optimize conditions



- ▶ Maximum efficiency at 0.25% deuterium concentration and a stopping point with part of the stopping distribution inside the target
- ▶ Reach a efficiency of around 7% of incoming muons hitting the target as μd atom

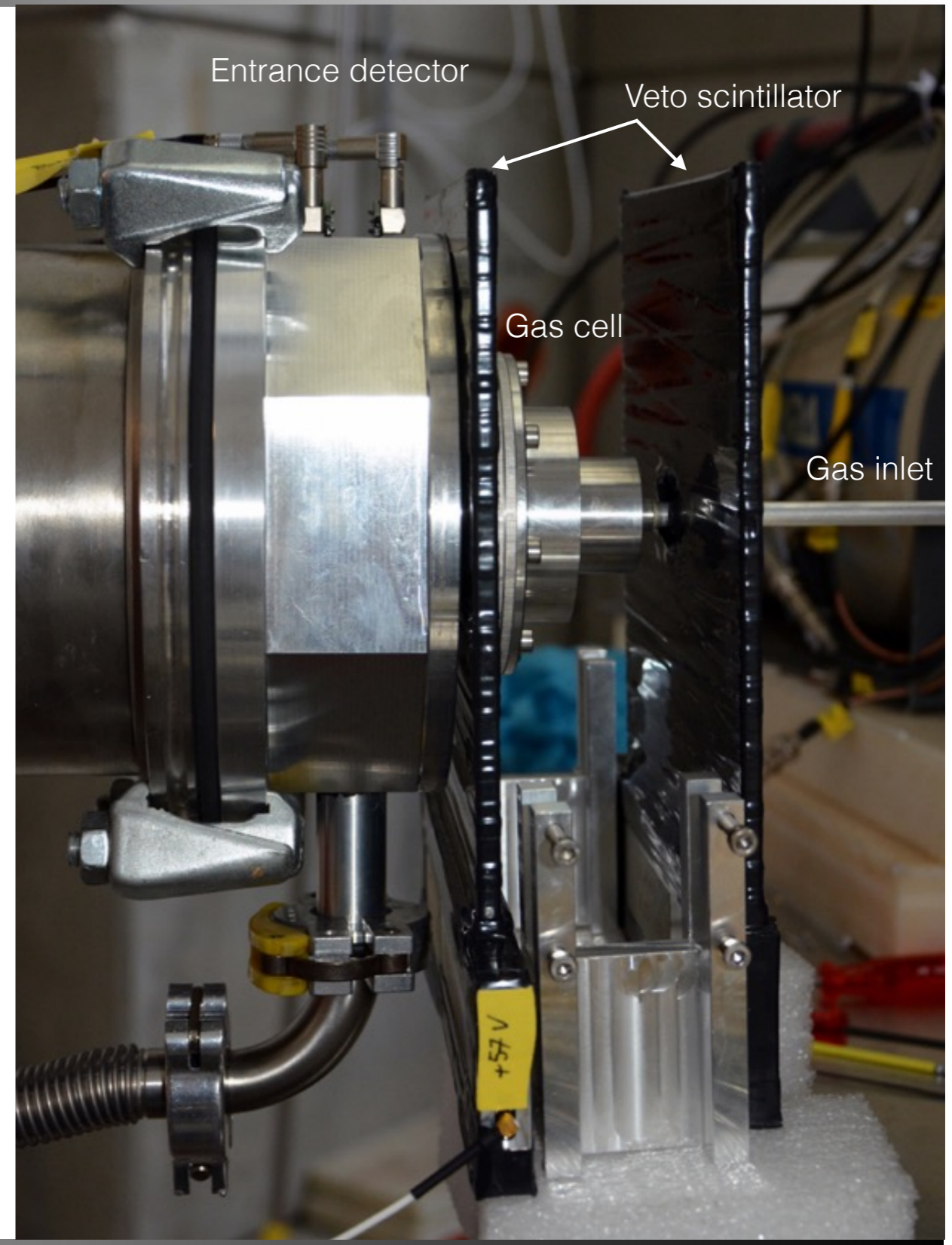
100 bar hydrogen target

- ▶ Target sealed with 0.6 mm carbon fiber window plus carbon fiber/titanium support grid
- ▶ Target holds up to 350 bar
- ▶ 8 mm stopping distribution (FWHM) inside 15 mm gas volume
- ▶ Target disks mounted onto the back of the cell



Entrance & veto detectors

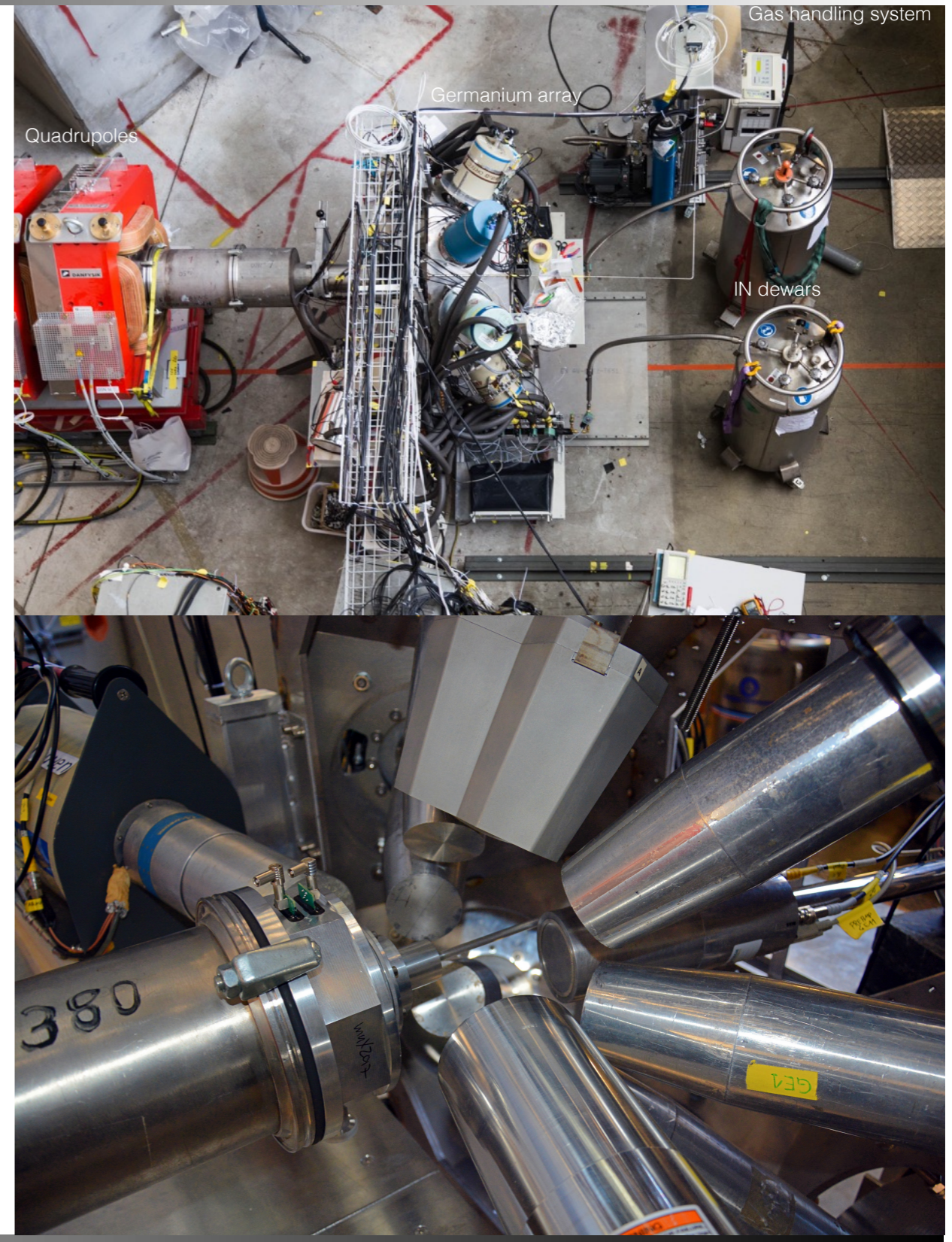
- ▶ Entrance detector to see incoming muon
- ▶ Veto scintillators to form anti-coincidence with decay electron



Optimize detection efficiency

- ▶ piE1 beam line at PSI
- ▶ ~ 10 kHz μ^- at 28 MeV/c

- ▶ 11 germanium detectors in an array from French/UK loan pool, Leuven, PSI



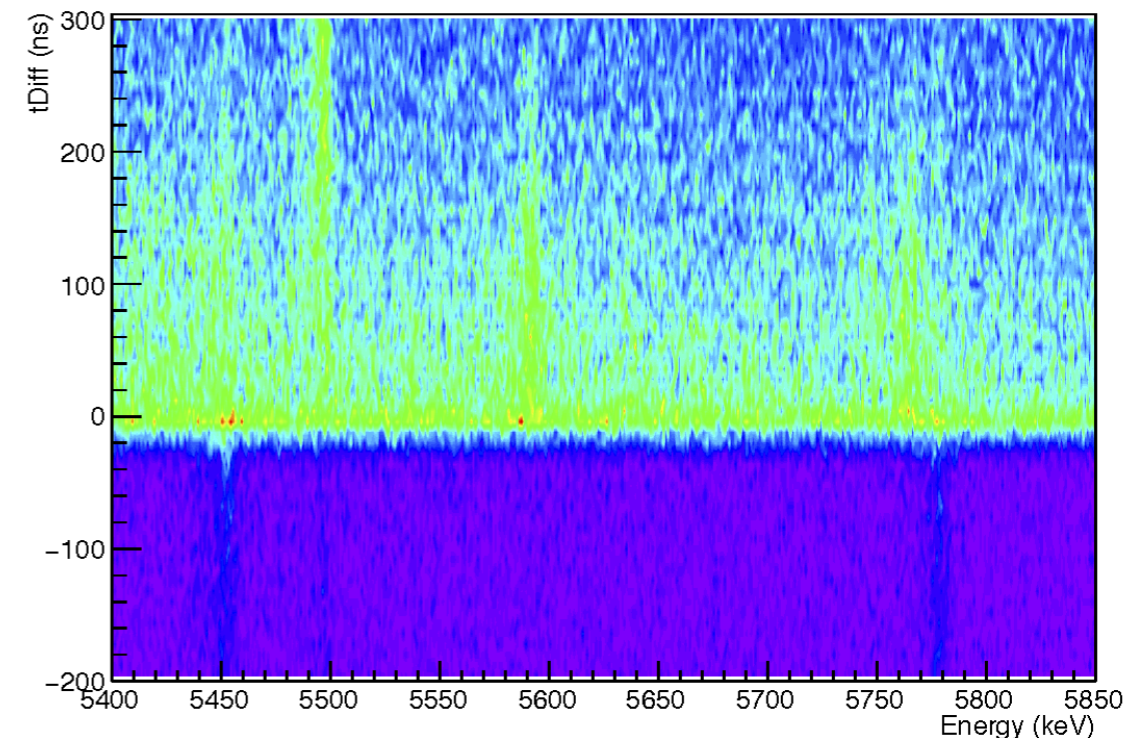
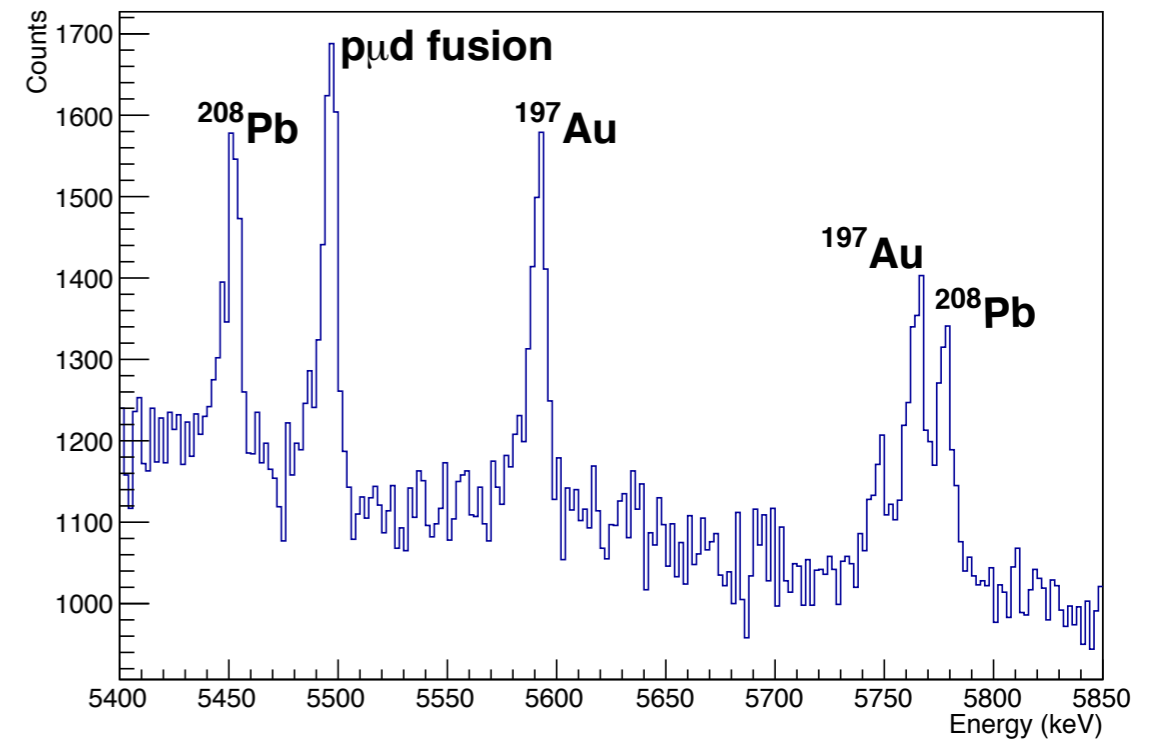
Optimize target conditions

Target	Size	Backing	N_γ / N_μ	ϵ
50 nm Au	4.9 cm ²	Cu	$(10.9 \pm 0.3) \times 10^{-5}$	10.0%
10 nm Au	4.9 cm ²	Cu	$(6.9 \pm 0.2) \times 10^{-5}$	6.3%
3 nm Au	4.9 cm ²	Cu	$(3.6 \pm 0.1) \times 10^{-5}$	3.3%
3 nm Au	4.9 cm ²	kapton	$(3.2 \pm 0.1) \times 10^{-5}$	2.9%
3 nm Au	1 cm ²	Cu	$(1.3 \pm 0.1) \times 10^{-5}$	1.2%

- ▶ Detected 2p-1s gammas per incoming muon for various targets
- ▶ Not all μ d converted in thin targets
- ▶ Can still reliably see gammas from 5 μ g gold target (1 cm², 3 nm)

Measurement

- ▶ Measurement with 5 μg gold target as proof-of-principle
- ▶ Data taken during 18.5 h
- ▶ Ready for radioactive radium target this year



Other radioactive isotopes

Isotope	Half-life	Max. Activity	Max. Mass
^{226}Ra	1600 y	200 kBq	5 μg
^{248}Cm	350'000 y	5 kBq	32 μg
^{209}Po	102 y	200 kBq	0.3 μg

- ▶ Isotopes without measured charge radius
- ▶ Maximum activity based on current regulations and without major modifications to experimental area infrastructure (100 x approval limit)

Benefit from more absolute measurements

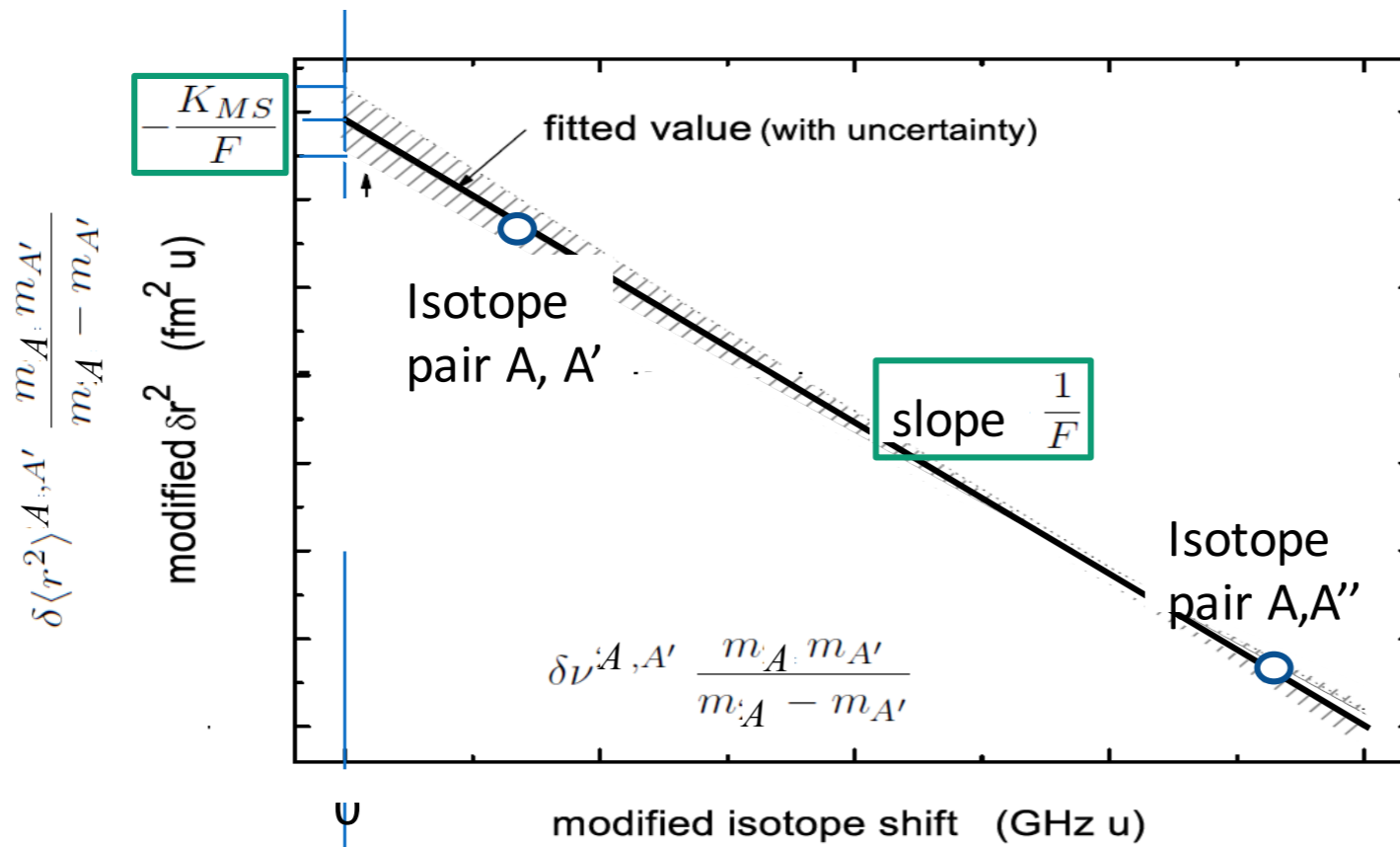
=> Modified King plot

M. Kowalska, Workshop on Muonic Atom Spectroscopy (2016)

- When data for at least 3 isotopes exists (i.e. stable isotopes):
 - Combine absolute radii (transitions in muonic atoms and/or electron scattering) and isotope shifts in optical transitions to derive more precise F and K_{MS} values

$$\delta\nu^{A,A'} \frac{m_A m_{A'}}{m_A - m_{A'}} = K_{MS} + F \delta\langle r^2 \rangle^{A,A'} \frac{m_A m_{A'}}{m_A - m_{A'}}$$

$$\delta\langle r^2 \rangle^{A,A'} \frac{m_A m_{A'}}{m_A - m_{A'}} = -\frac{K_{MS}}{F} + \frac{1}{F} \delta\nu^{A,A'} \frac{m_A m_{A'}}{m_A - m_{A'}}$$



But if there are fewer stable isotopes ...
See Na, Mn, Cu, Ga ...

Isotopes at SPES

- ▶ Isotopes without muonic data with sufficient yield at SPES to reach μg ion numbers in a few days
- ▶ Sufficiently long lived for transport
- ▶ Need to carefully study activity and regulations for lower half-life isotopes
- ▶ We'll continue to develop the method towards lower target masses
 - more radioactive elements can be measured

		Yield (SPES, 200 μA)	Half-life [s]
Sn	123	1.28E+10	1.12E+07
Sn	125	3.5E+10	8.33E+05
Sn	126	4.21E+10	3.16E+12
Te	132	4.21E+10	2.77E+05
I	131	5.47E+10	6.93E+05
I	133	1.89E+11	7.49E+04
Cs	134	4.08E+09	6.52E+07
Cs	135	1.68E+10	7.26E+13
Cs	136	4.96E+10	1.14E+06
Cs	137	1.07E+11	9.49E+08
Ba	140	6.05E+10	1.10E+06

Conclusions

- ▶ Particle physics interest in radioactive ions in connection with the study of EDM, weak interaction and atomic parity violation
- ▶ Muonic atom spectroscopy able to measure absolute charge radii needed for atomic theory in APV
- ▶ muX collaboration developed method of using microgram targets for muonic atom spectroscopy enabling the measurement of high-activity targets
- ▶ Several isotopes with sufficient yield at SPES identified that could be measured with our method

A. Adamczak¹, A. Antognini^{2,3}, N. Berger⁴, T. Cocolios⁵, R. Dressler²,
R. Eichler², P. Indelicato⁶, K. Jungmann⁷, K. Kirch^{2,3}, A. Knecht²,
J. Krauth⁴, J. Nuber², A. Papa², R. Pohl⁴, M. Pospelov^{8,9},
E. Rapisarda², P. Reiter¹⁰, N. Ritjoho^{2,3}, S. Roccia¹¹, N. Severijns⁵,
A. Skawran^{2,3}, F. Wauters⁴, and L. Willmann⁷

¹Institute of Nuclear Physics, Polish Academy of Sciences, Krakow,
Poland

²Paul Scherrer Institut, Villigen, Switzerland

³ETH Zürich, Switzerland

⁴University of Mainz, Germany

⁵KU Leuven, Belgium

⁶LKB Paris, France

⁷University of Groningen, The Netherlands

⁸University of Victoria, Canada

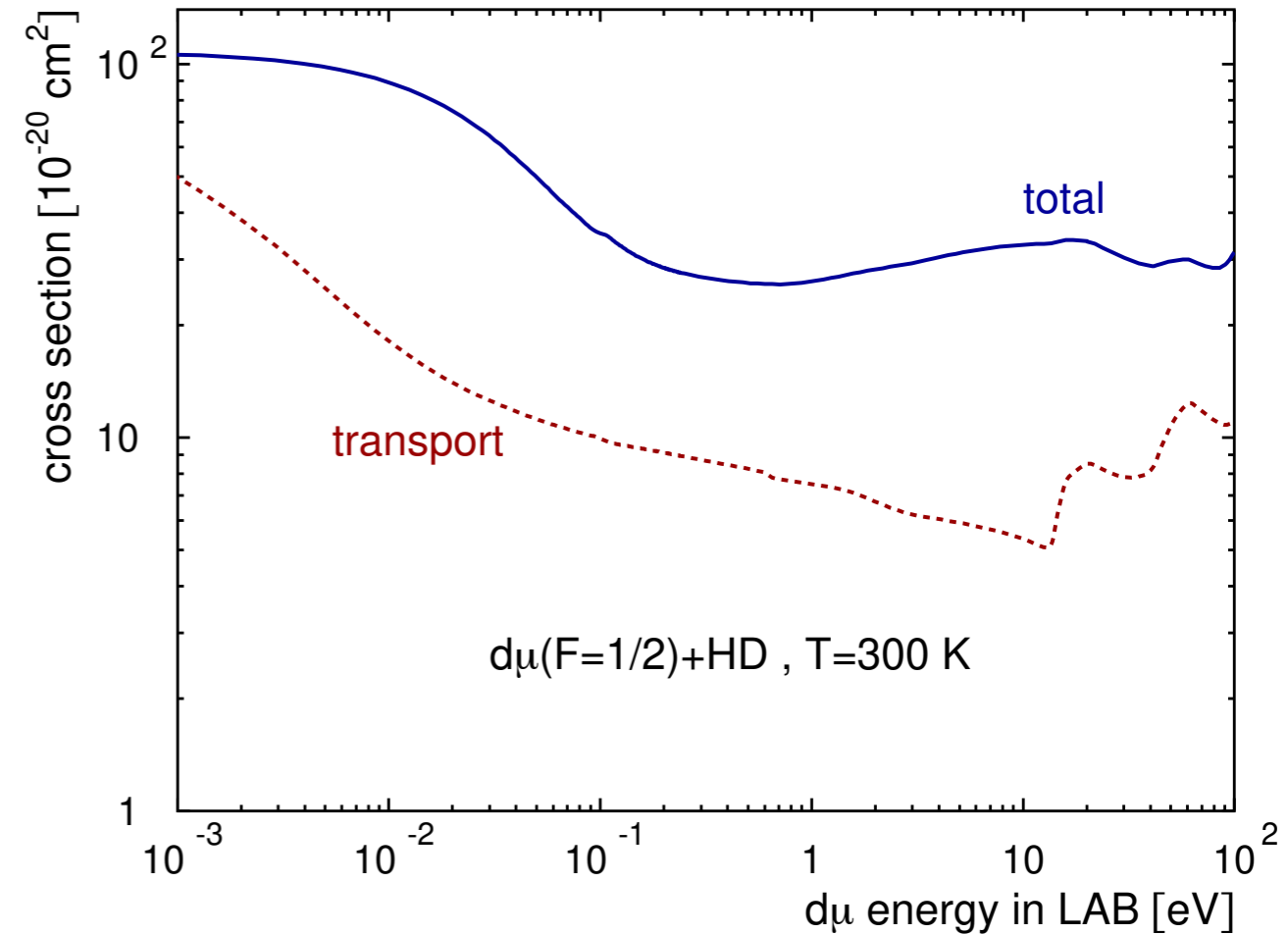
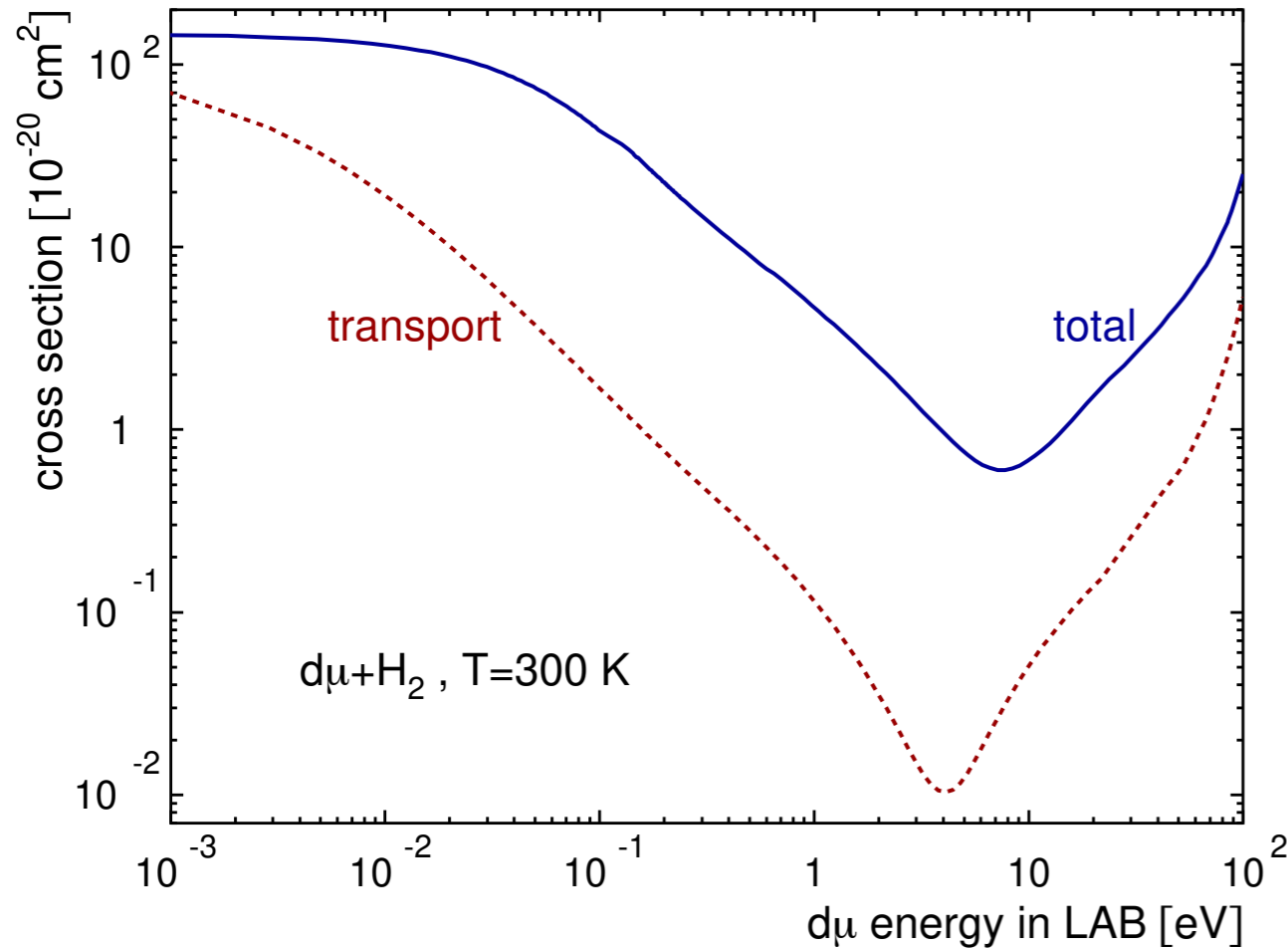
⁹Perimeter Institute, Waterloo, Canada

¹⁰Institut für Kernphysik, Universität zu Köln, Germany

¹¹CSNSM, Université Paris Sud, CNRS/IN2P3, Orsay Campus, France

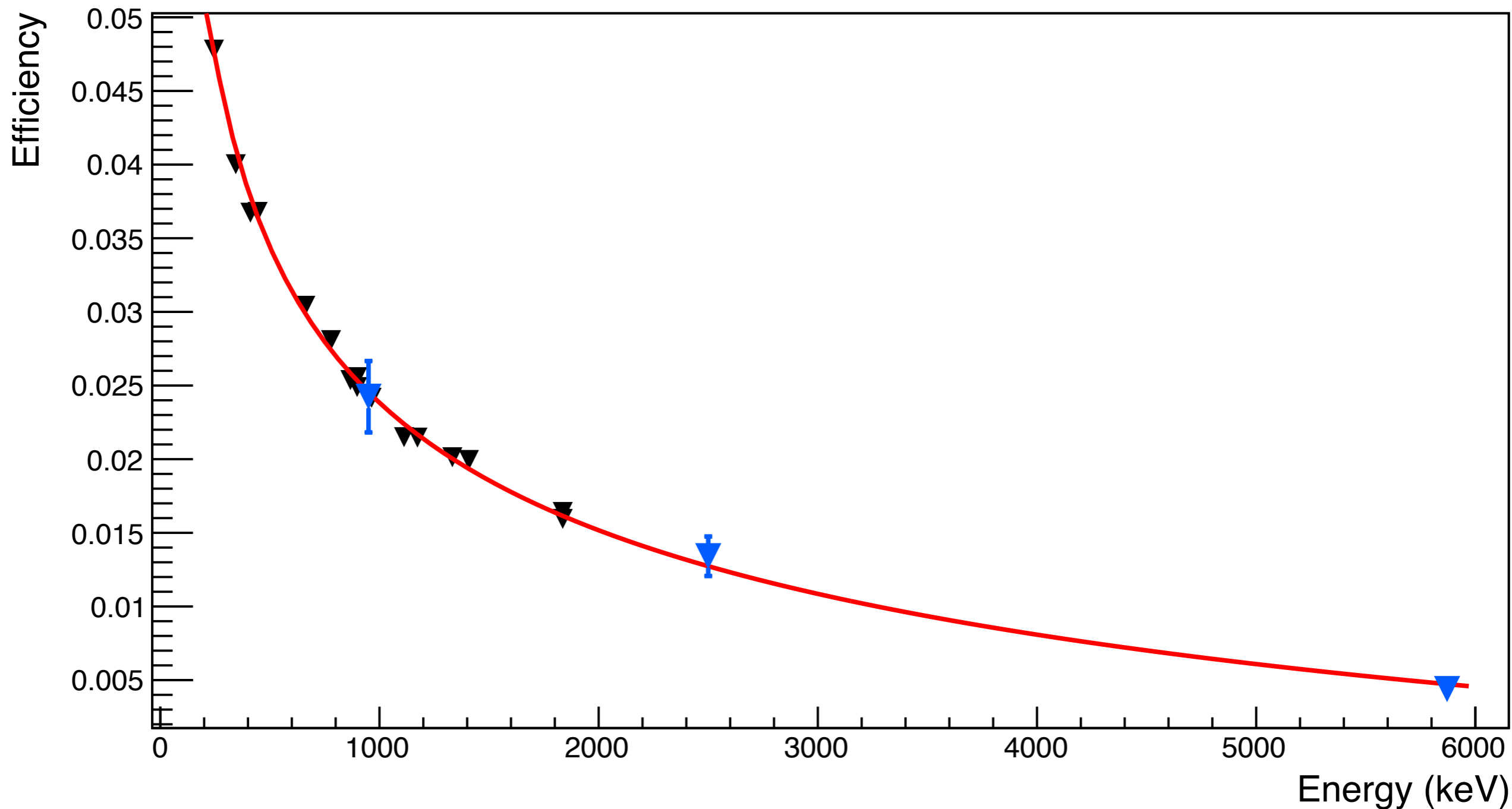
Backup

Scattering cross sections



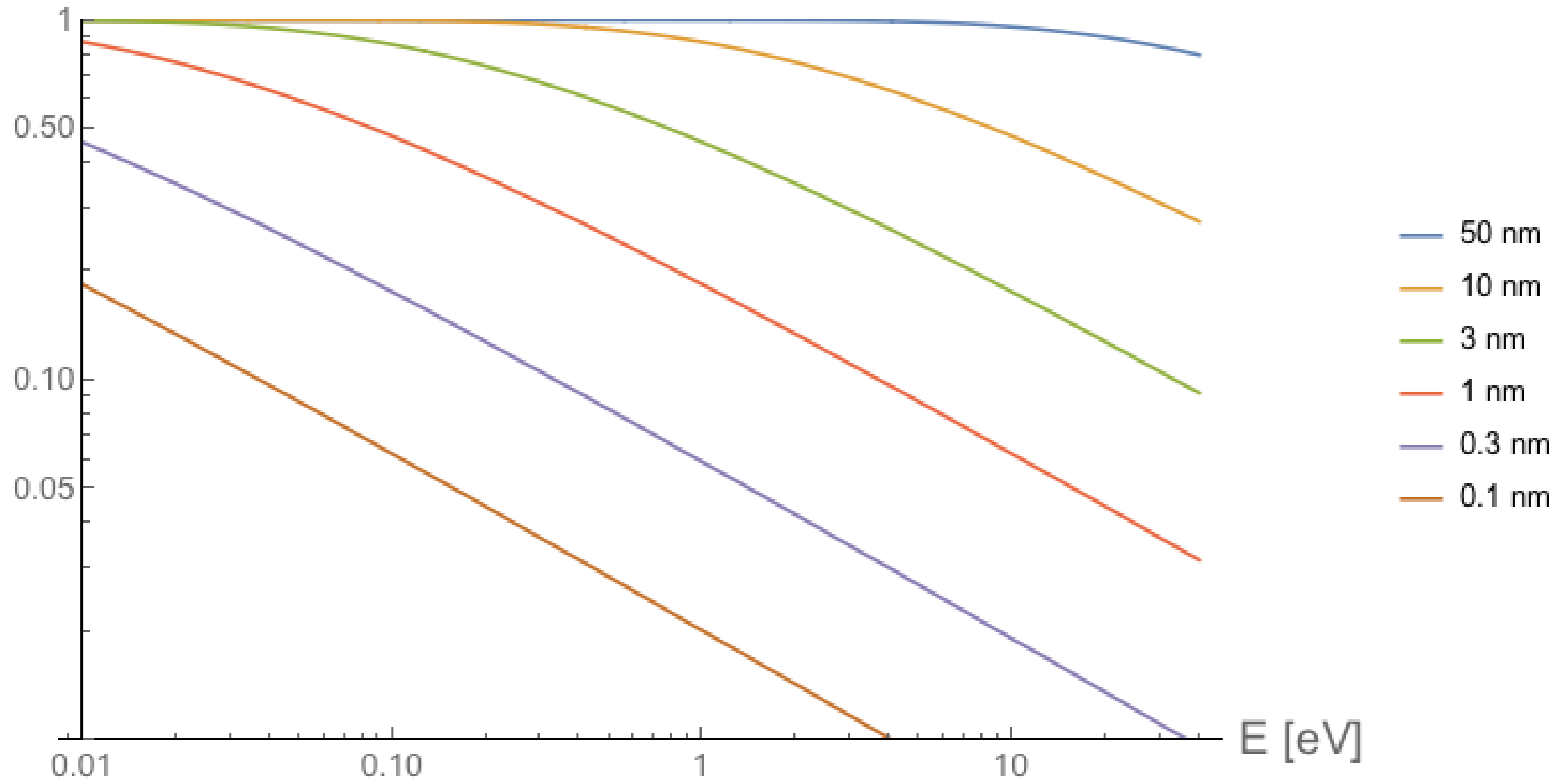
Array Detection Efficiency

Detector Efficiency Ge1-10



Transfer Probability in Gold

Transfer Prob.



Weinberg angle

The **Weinberg angle** or **weak mixing angle** is a parameter in the **Weinberg–Salam** theory of the **electroweak interaction**, part of the **Standard Model** of particle physics, and is usually denoted as θ_W . It is the angle by which **spontaneous symmetry breaking** rotates the original W^0 and B^0 **vector boson** plane, producing as a result the Z^0 boson, and the **photon**.

$$\begin{pmatrix} \gamma \\ Z^0 \end{pmatrix} = \begin{pmatrix} \cos \theta_W & \sin \theta_W \\ -\sin \theta_W & \cos \theta_W \end{pmatrix} \begin{pmatrix} B^0 \\ W^0 \end{pmatrix}$$

It also gives the relationship between the masses of the **W and Z bosons** (denoted as m_W and m_Z).

$$m_Z = \frac{m_W}{\cos \theta_W}.$$

The angle can be expressed in terms of the $SU(2)_L$ and $U(1)_Y$ **couplings** (**weak isospin** g and **weak hypercharge** g' , respectively),

$$\cos \theta_W = \frac{g}{\sqrt{g^2 + g'^2}} \quad \text{and} \quad \sin \theta_W = \frac{g'}{\sqrt{g^2 + g'^2}}.$$

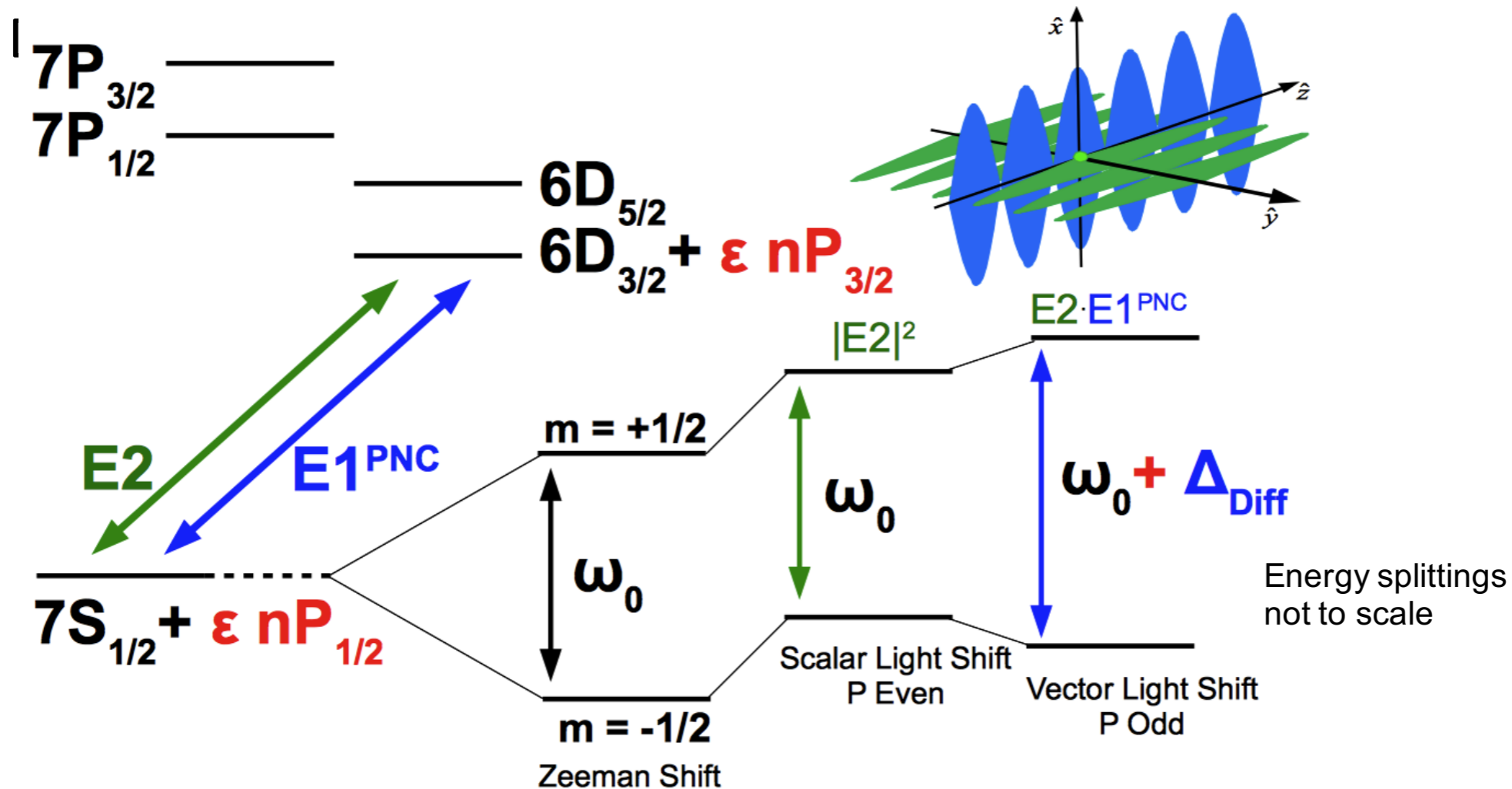
The electric charge is then expressible in terms of it, $e = g \sin \theta_W = g' \cos \theta_W$; see the Figure.

As the value of the mixing angle is currently determined empirically, it has been mathematically defined as^[1]

$$\cos \theta_W = \frac{m_W}{m_Z}.$$

https://en.wikipedia.org/wiki/Weinberg_angle

Experimental Method



$$|\Omega|^2 = |\Omega_{m'm}^{E2} + \Omega_{m'm}^{PNC}|^2 \sim |\Omega_{m'm}^{E2}|^2 + 2\text{Re}|\Omega_{m'm}^{PNC}\Omega_{m'm}^{E2}|$$

N. Fortson, Phys. Rev. Lett. 70, 2383-2386 (1993)

Radium Activity

- ▶ 5 μg corresponds to 200 kBq of ^{226}Ra and all daughters
- ▶ Highest gamma emitters: ^{214}Pb , ^{214}Bi
- ▶ Gamma rate: ~ 400 kHz

